

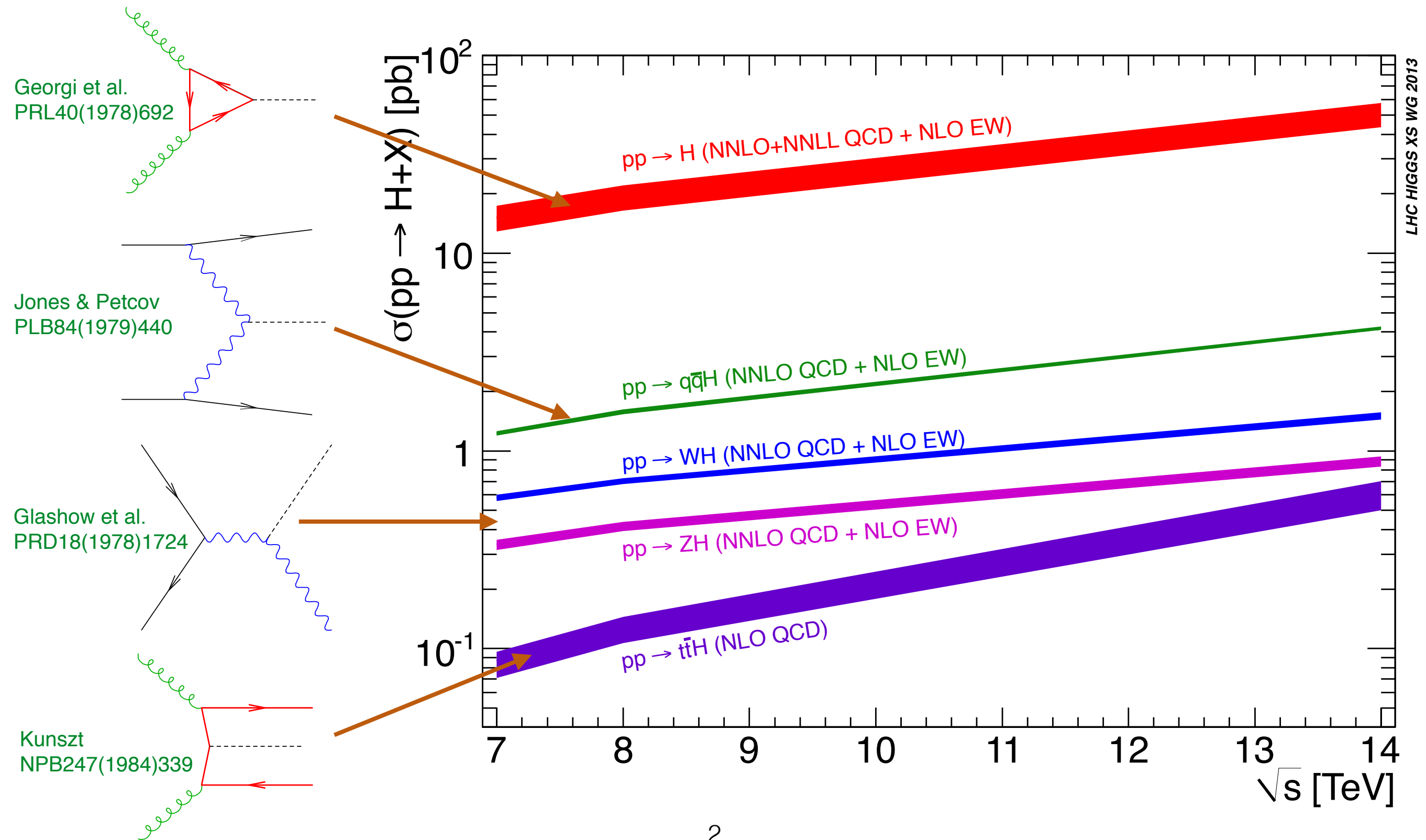
# Higher-Order QCD Effects in Higgs Boson Production

Bryan Webber



UNIVERSITY OF  
CAMBRIDGE

# Higgs production cross sections



The challenges:

(taken from [R. Tanaka, talk at Aspen Higgs WS 03/13])

## ggF, VBF, WH/ZH, ttH, BSM Higgs

**Higgs Cross Sections**  
(inclusive/exclusive)

**Differential K-factors**  
(effect of jet-veto etc.)

**QCD correction  $N^k\text{LO} + N^m\text{LL}$**   
**EW correction, Mixed QCD-EW**

**Heavy Higgs Line Shape**  
**SM Backgrounds & Interferences**

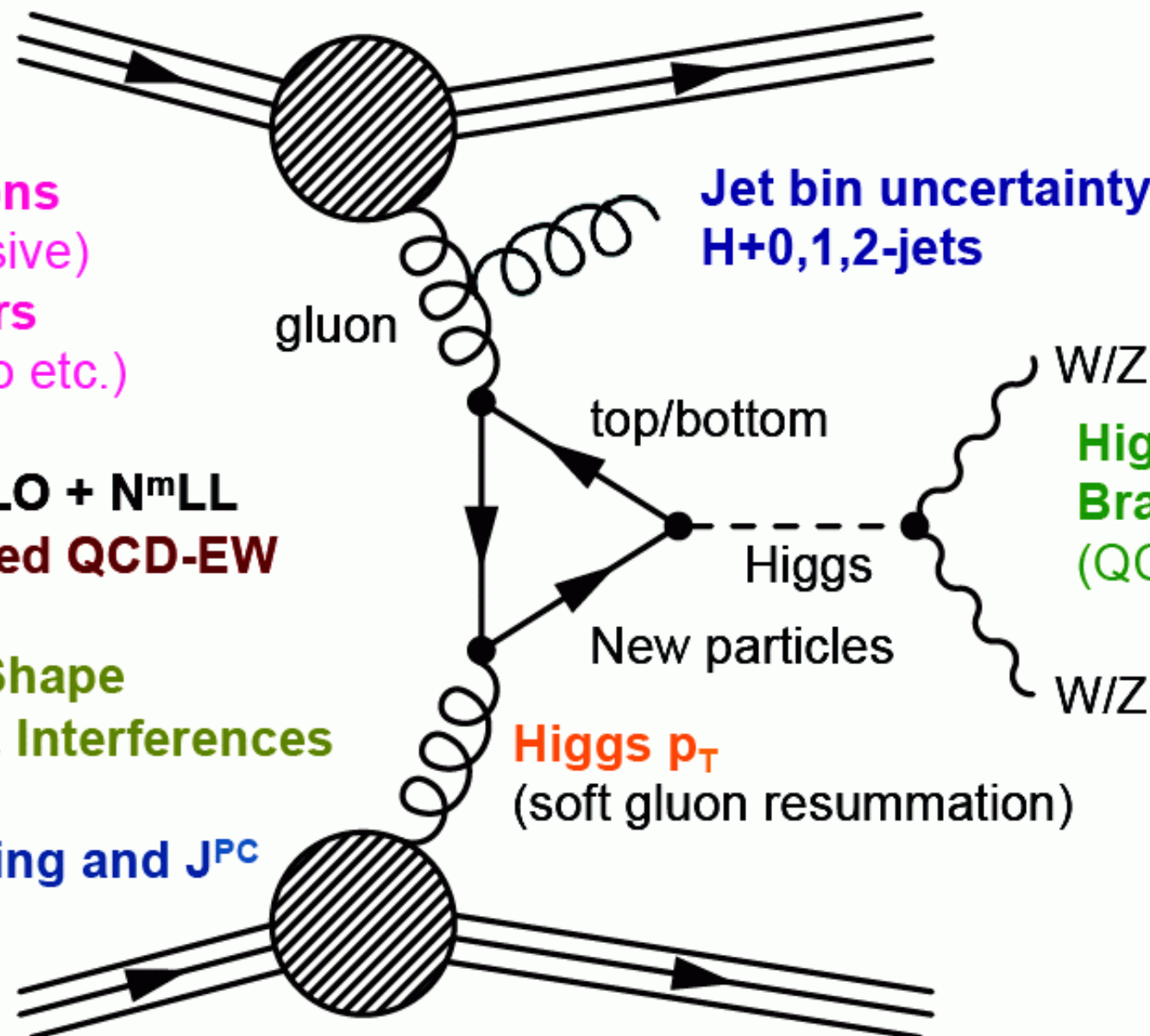
**Higgs Mass, Coupling and  $J^{PC}$**

**PDF+ $\alpha_s$  uncertainties**

**Renormalization/Factorization scale dependence**

**Jet bin uncertainty**  
**H+0,1,2-jets**

**Higgs decay**  
**Branching ratios**  
(QCD/EW corr.)



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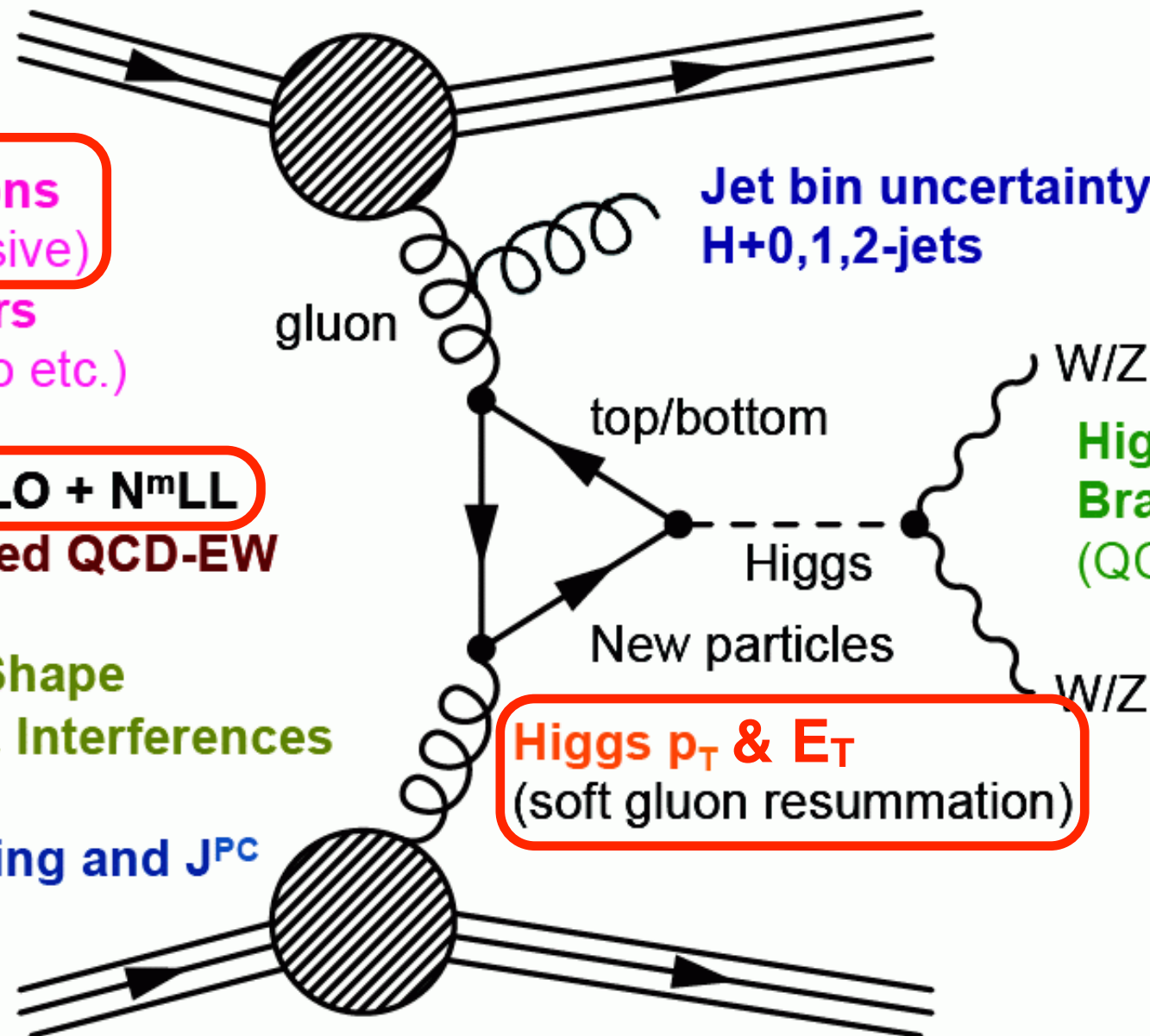
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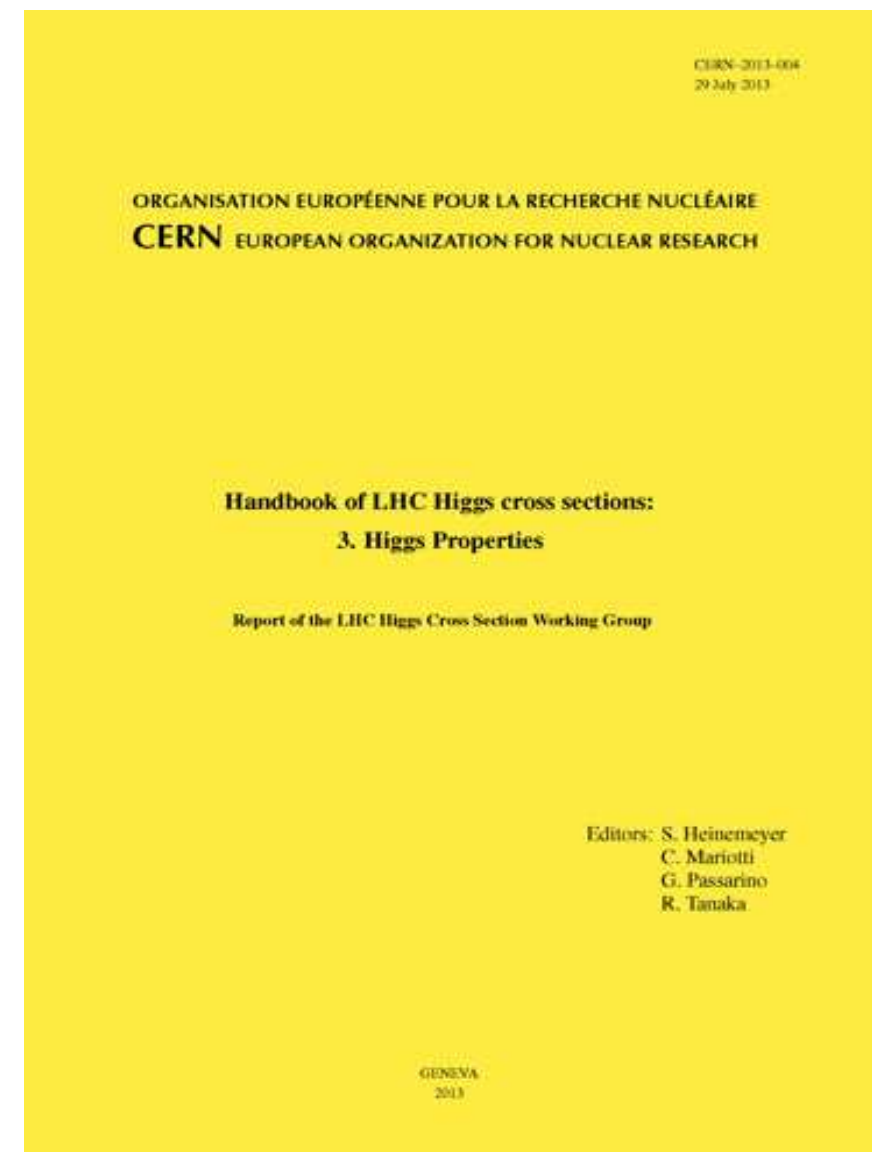
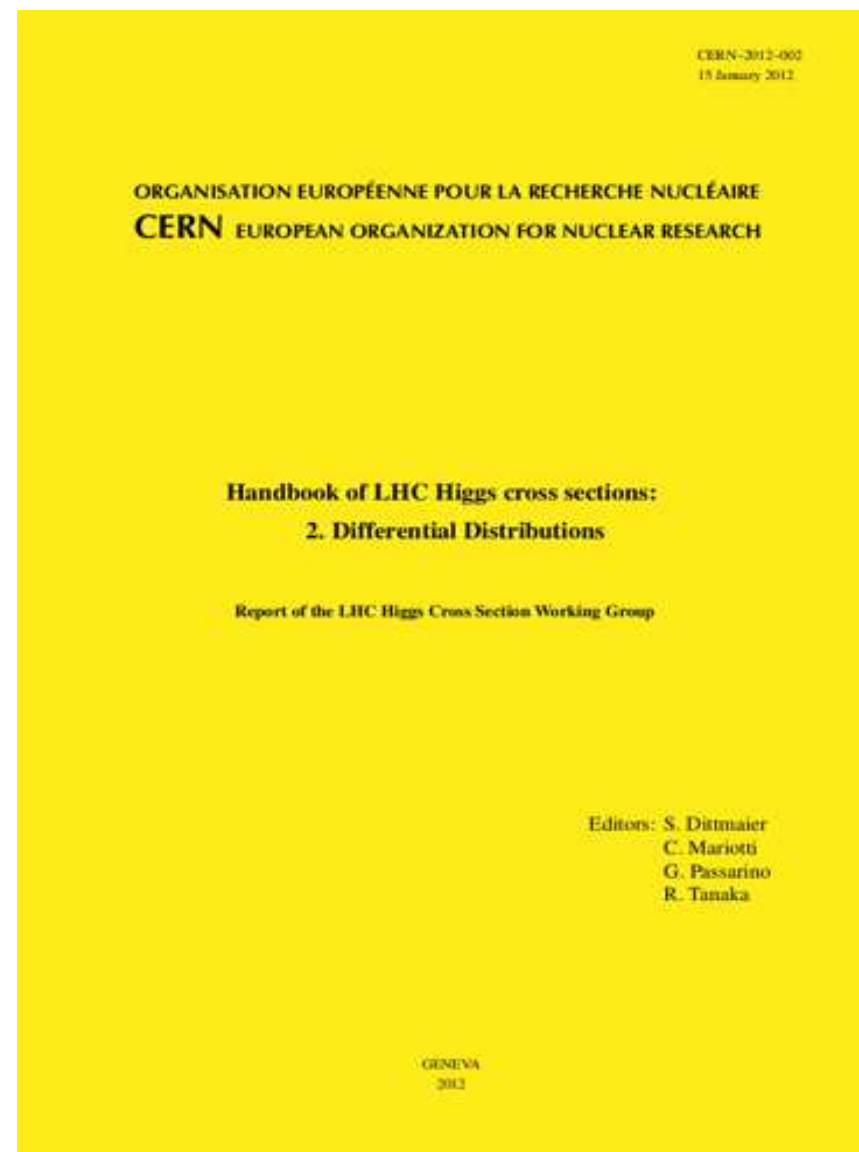
**Higgs decay**  
**Branching ratios**  
(QCD/EW corr.)

**Higgs  $p_T$  &  $E_T$**   
(soft gluon resummation)

**PDF+ $\alpha_s$  uncertainties**  
**Renormalization/Factorization scale dependence**



# LHC Higgs Cross Section Working Group



arXiv:1101.0593, 1201.3084, 1307.1347

# Higgs Production Cross Section

# Inclusive (QCD) Cross Section

$$\sigma_H(s) = \sum_{i,j} \int dx_1 dx_2 f_i(x_1, \mu_F^2) f_j(x_2, \mu_F^2) \hat{\sigma}_{ij}(x_1 x_2 s, m_H^2, \alpha_S(\mu_R^2), \mu_F^2, \mu_R^2)$$

**Table B.10:** ggF cross sections at the LHC at 8 TeV and corresponding scale and PDF+ $\alpha_s$  uncertainties computed according to the PDF4LHC recommendation.

HXSWG vol.3

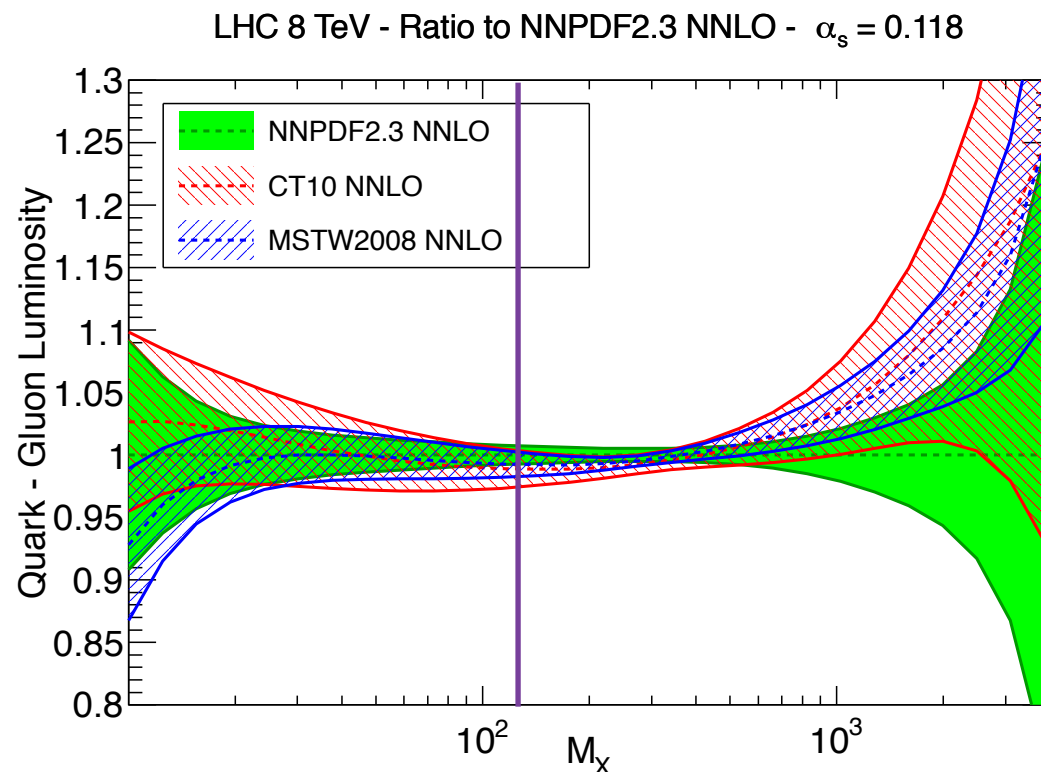
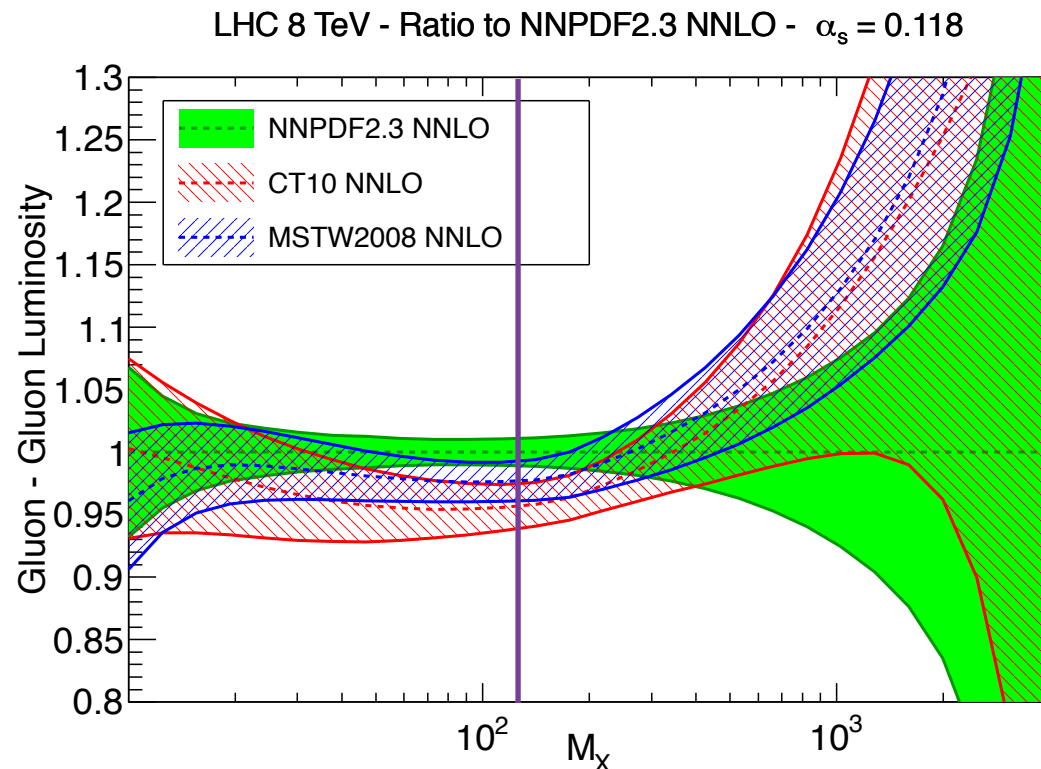
$M_H$ [GeV]	$\sigma$ [pb]	QCD Scale [%]		PDF+ $\alpha_s$ [%]	
124.4	19.45	+7.2	-7.9	+7.5	-6.9
124.5	19.42	+7.2	-7.9	+7.5	-6.9
124.6	19.39	+7.2	-7.9	+7.5	-6.9
124.7	19.36	+7.2	-7.9	+7.5	-6.9
124.8	19.33	+7.2	-7.8	+7.5	-6.9
124.9	19.30	+7.2	-7.8	+7.5	-6.9
125.0	19.27	+7.2	-7.8	+7.5	-6.9
125.1	19.24	+7.2	-7.8	+7.5	-6.9
125.2	19.21	+7.2	-7.8	+7.5	-6.9
125.3	19.18	+7.2	-7.8	+7.5	-6.9
125.4	19.15	+7.2	-7.8	+7.5	-6.9
125.5	19.12	+7.2	-7.8	+7.5	-6.9
125.6	19.09	+7.2	-7.8	+7.5	-6.9
125.7	19.06	+7.2	-7.8	+7.5	-6.9
125.8	19.03	+7.2	-7.8	+7.5	-6.9
125.9	19.00	+7.2	-7.8	+7.5	-6.9
126.0	18.97	+7.2	-7.8	+7.5	-6.9
126.1	18.94	+7.2	-7.8	+7.5	-6.9
126.2	18.91	+7.2	-7.8	+7.5	-6.9
126.3	18.88	+7.2	-7.8	+7.5	-6.9
126.4	18.85	+7.2	-7.8	+7.5	-6.9
126.5	18.82	+7.2	-7.8	+7.5	-6.9

NNLO   $\sigma_{\text{ggF}}(8 \text{ TeV}) = 19.1 \pm 2.0 \text{ pb}$



# PDF Uncertainties

Ball et al., 1211.5142



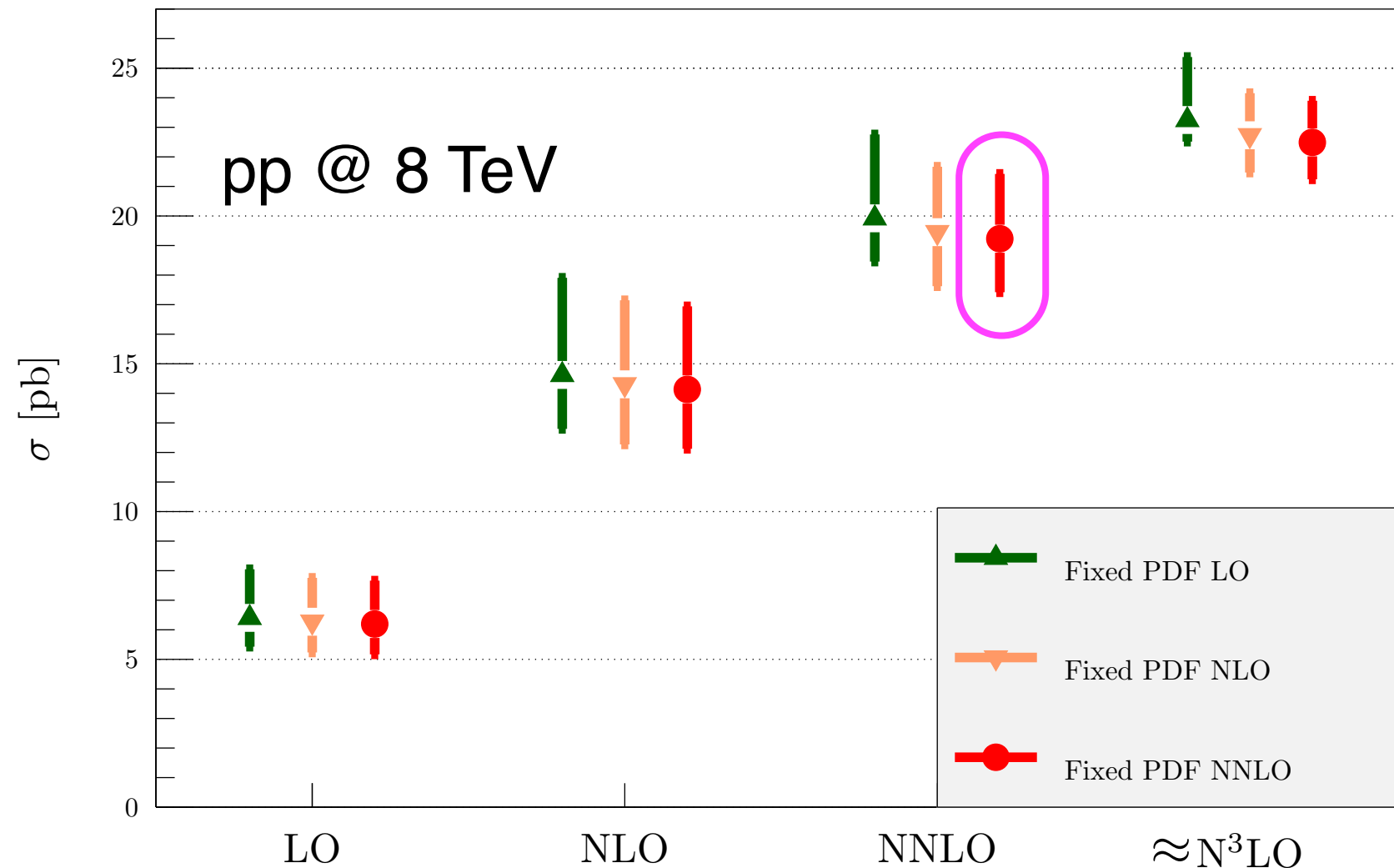
- Parton luminosity  $\mathcal{L}_{ij}(M_X^2, s) = \int dx_1 dx_2 f_i(x_1, M_X^2) f_j(x_2, M_X^2) \delta(x_1 x_2 s - M_X^2)$
- Relevant PDFs (relatively) well known at  $x \sim M_H/\sqrt{s}$
- Some disagreement with CT10  $\mathcal{L}_{gg}$
- Remains true at 13 TeV
- Can be improved (in principle)



# Inclusive (QCD) Cross Section

Ball et al., 1303.3590

Forte, Isigro, Vita, 1312.6688



- Higher-order effects are larger than x2 scale variation estimates

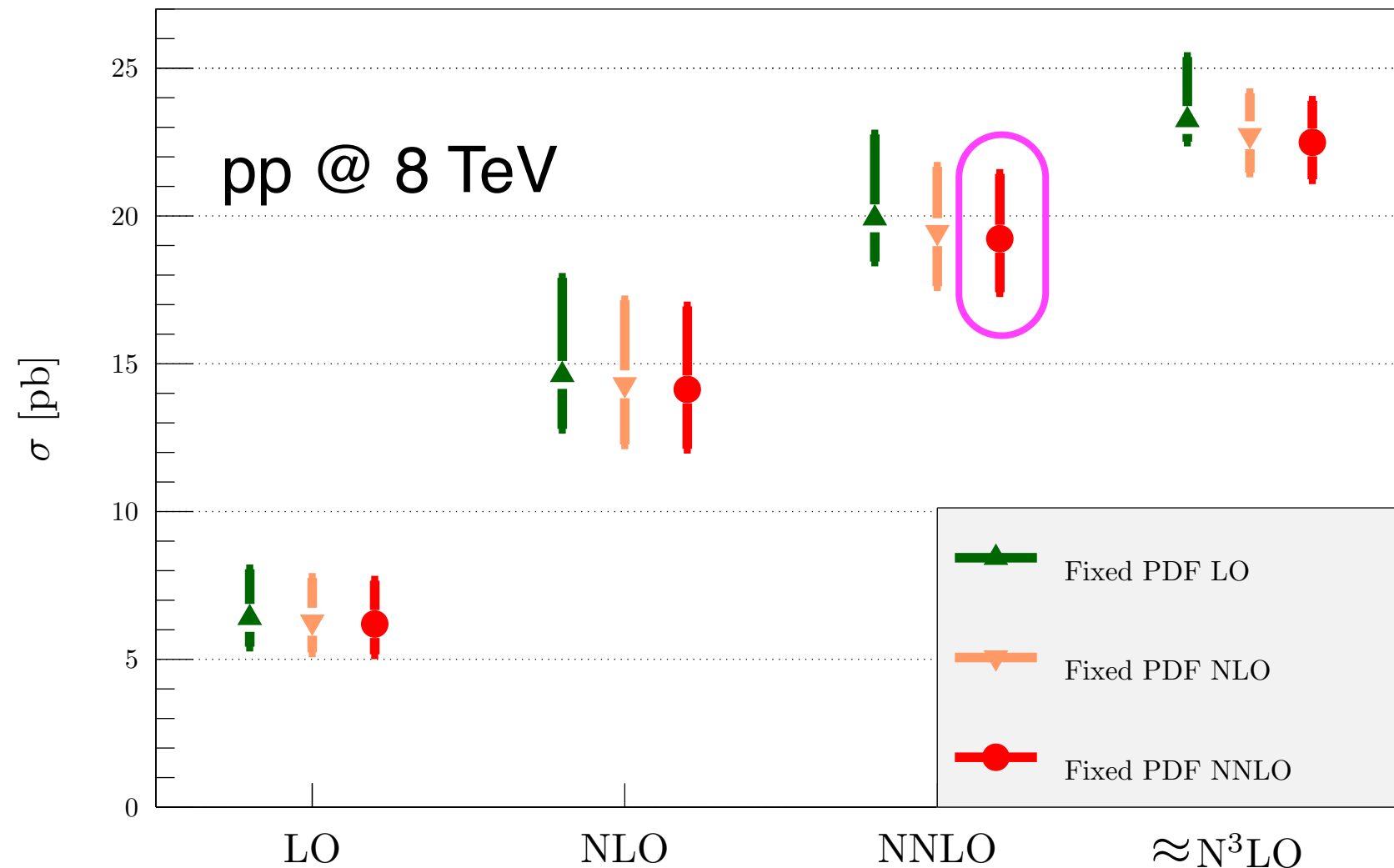
- $$\sigma_{ggF} \approx \sigma_{LO} \left[ 1 + \kappa \{ \lambda \alpha_S + (\lambda \alpha_S)^2 + (\lambda \alpha_S)^3 + \dots \} \right]$$

$\lambda \approx 5.6$

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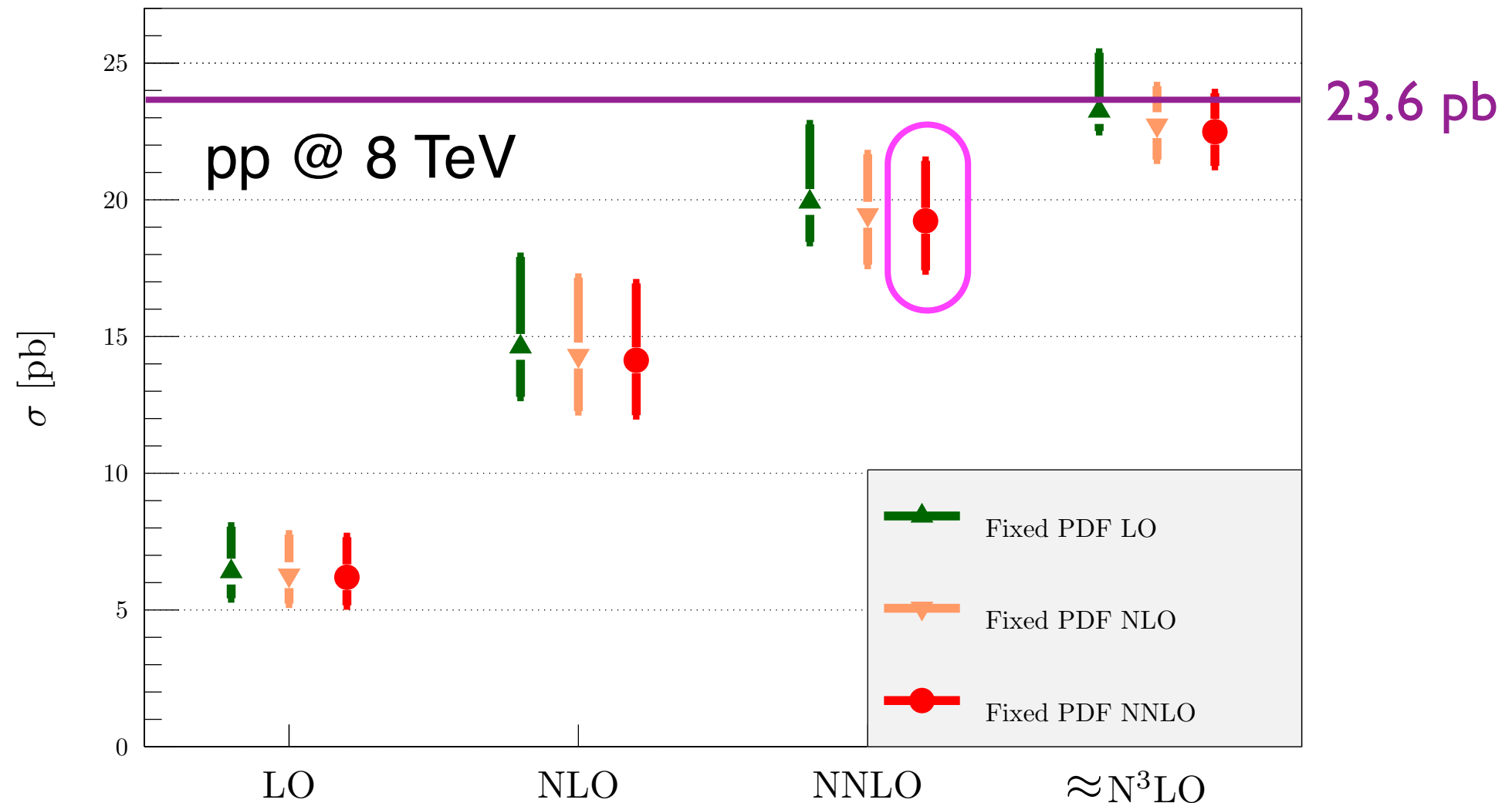
- $$\sigma_{ggF} \approx \sigma_{LO} \left[ 1 + \kappa \{ \lambda \alpha_S + (\lambda \alpha_S)^2 + (\lambda \alpha_S)^3 + \dots \} \right] \approx \sigma_{LO} \left[ 1 - \kappa + \frac{\kappa}{1 - \lambda \alpha_S} \right]$$

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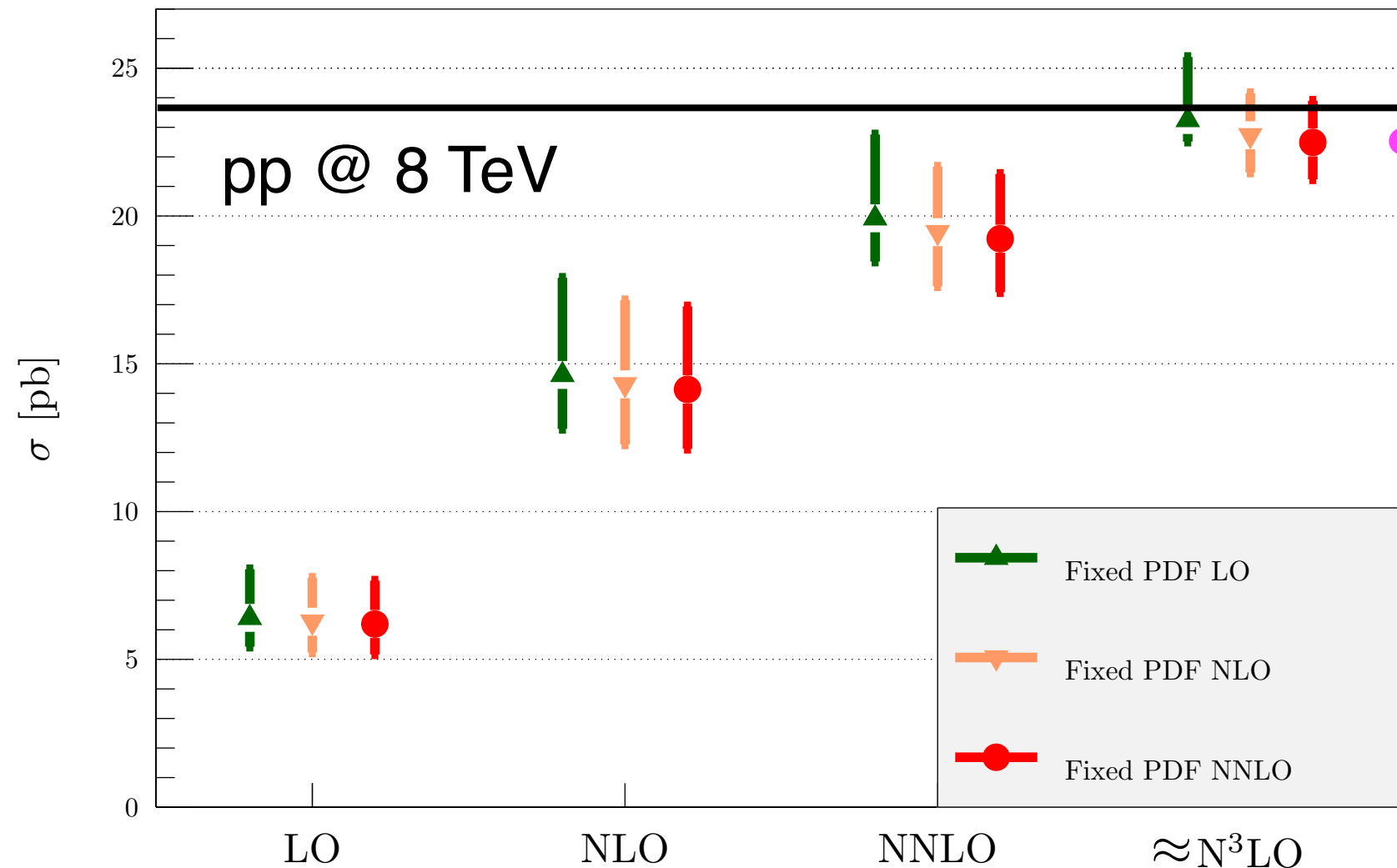


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- $$\sigma_{ggF} \approx \sigma_{LO} \left[ 1 + \kappa \{ \lambda \alpha_S + (\lambda \alpha_S)^2 + (\lambda \alpha_S)^3 + \dots \} \right] \approx \sigma_{LO} \left[ 1 - \kappa + \frac{\kappa}{1 - \lambda \alpha_S} \right]$$

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# Inclusive (QCD) Cross Section



NNLO  $\rightarrow \sigma_{ggF}(8 \text{ TeV}) = 19.1 \pm 2.0 \text{ pb}$

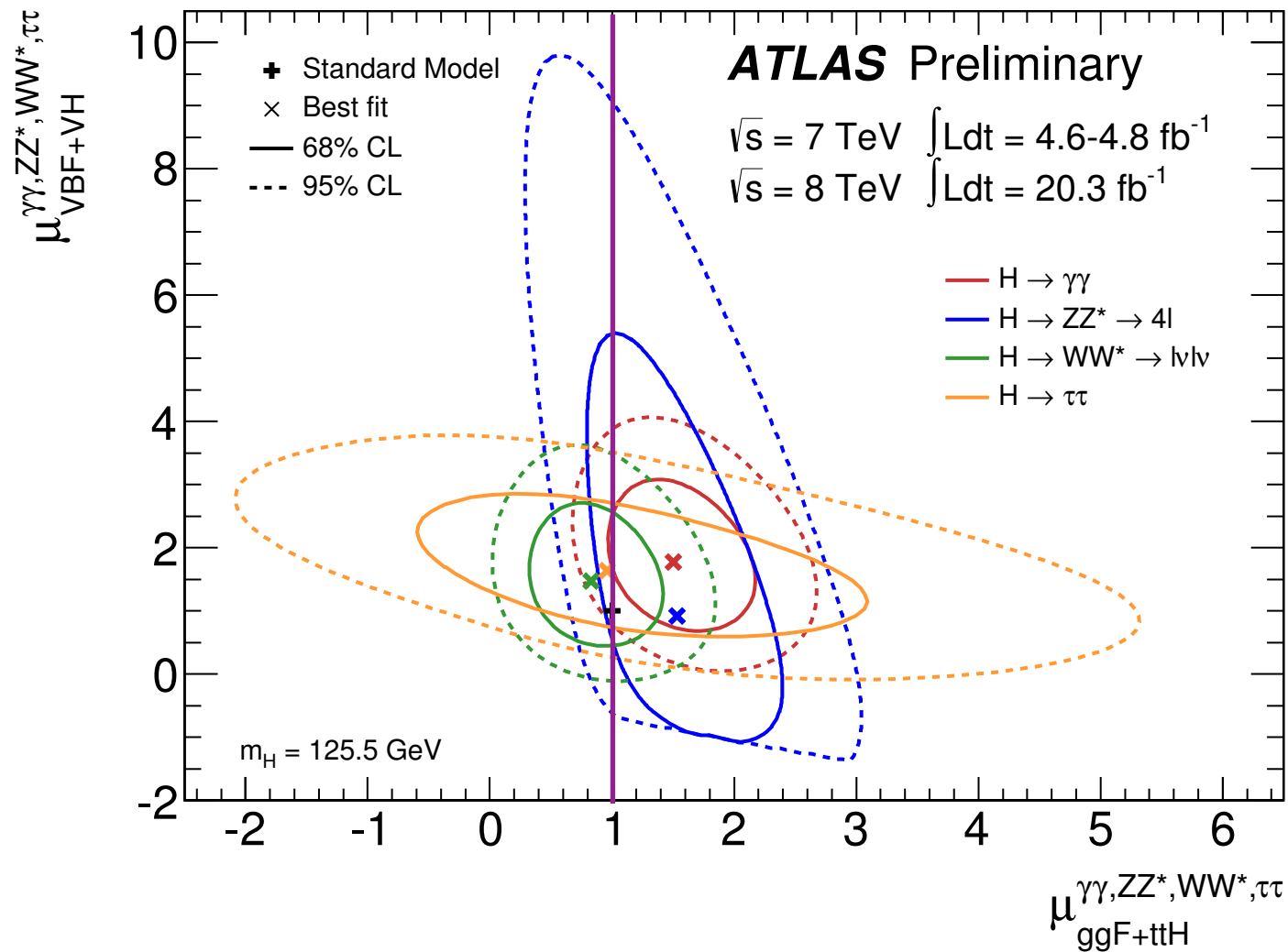
Series extrapolation:  $23.6 \pm ?? \text{ pb}$

David & Passarino, 1307.1843:  $22.5 \pm 2.6 \text{ pb}$

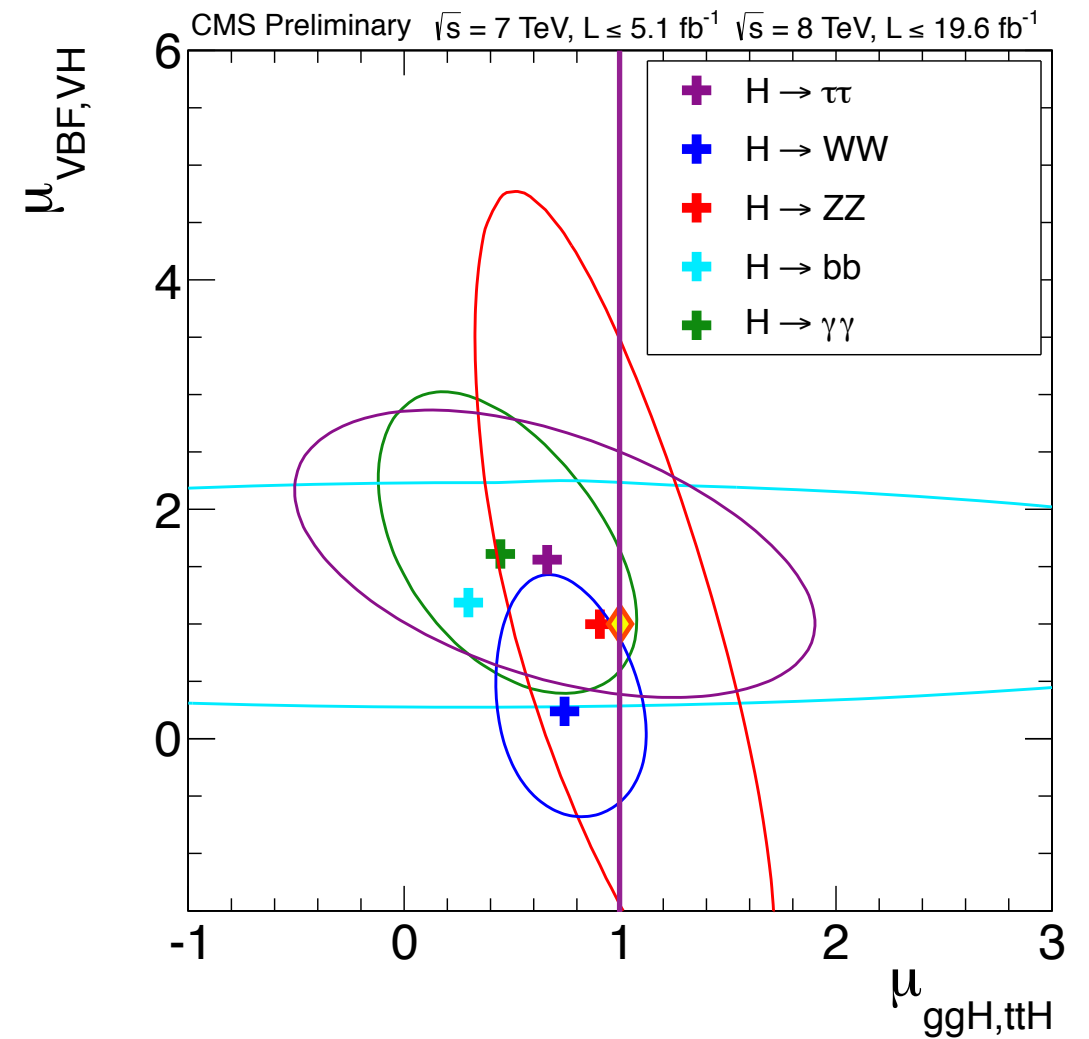
Full  $N^3LO$  coming soon (Anastasiou et al.)

# Inclusive (QCD) Cross Section

- ATLAS, but not CMS, find ggF excess in  $\gamma\gamma$  and  $ZZ^*$  channels



ATLAS-CONF-2014-009 (Moriond EW)



CMS PAS HIG-13-005

# Cross Sections at 13 TeV

HXSWG 05/04/2014

$\sqrt{s} = 13.0 \text{ TeV}$

## gluon-gluon Fusion Process

- All cross sections are in complex-pole-scheme from the dFG program. They are computed at NNLL QCD and NLO EW.

$m_H$ (GeV)	Cross Section (pb)	+QCD Scale %	-QCD Scale %	+(PDF+ $\alpha_s$ ) %	-(PDF+ $\alpha_s$ ) %
125.0	43.92	+7.4	-7.9	+7.1	-6.0
125.5	43.62	+7.4	-7.9	+7.1	-6.0
126.0	43.31	+7.4	-7.9	+7.1	-6.0

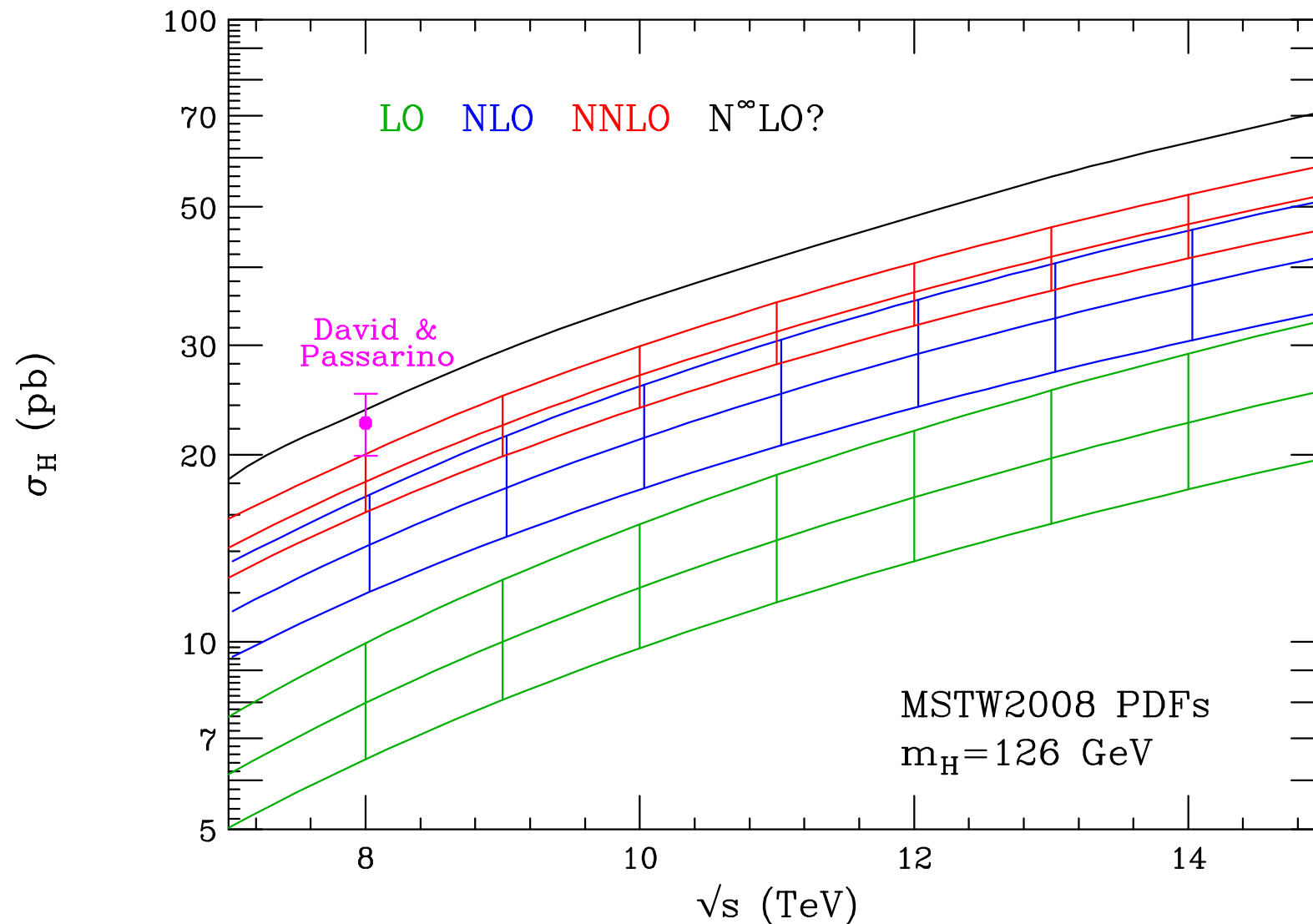
## VBF Process

- At NNLO QCD and NLO EW. All cross sections are in complex-pole-scheme.

$m_H$ (GeV)	Cross Section (pb)	+QCD Scale %	-QCD Scale %	+(PDF+ $\alpha_s$ ) %	-(PDF+ $\alpha_s$ ) %
125.0	3.748	+0.7	-0.7	+3.2	-3.2
125.5	3.727	+1.0	-0.7	+3.4	-3.4
126.0	3.703	+1.3	-0.6	+3.1	-3.1

# Cross Section vs Energy

<http://theory.fi.infn.it/grazzini/hcalculators.html>



NNLO  $\rightarrow \sigma_{ggF}(13 \text{ TeV}) = 43.3 \pm 4.3 \text{ pb}$

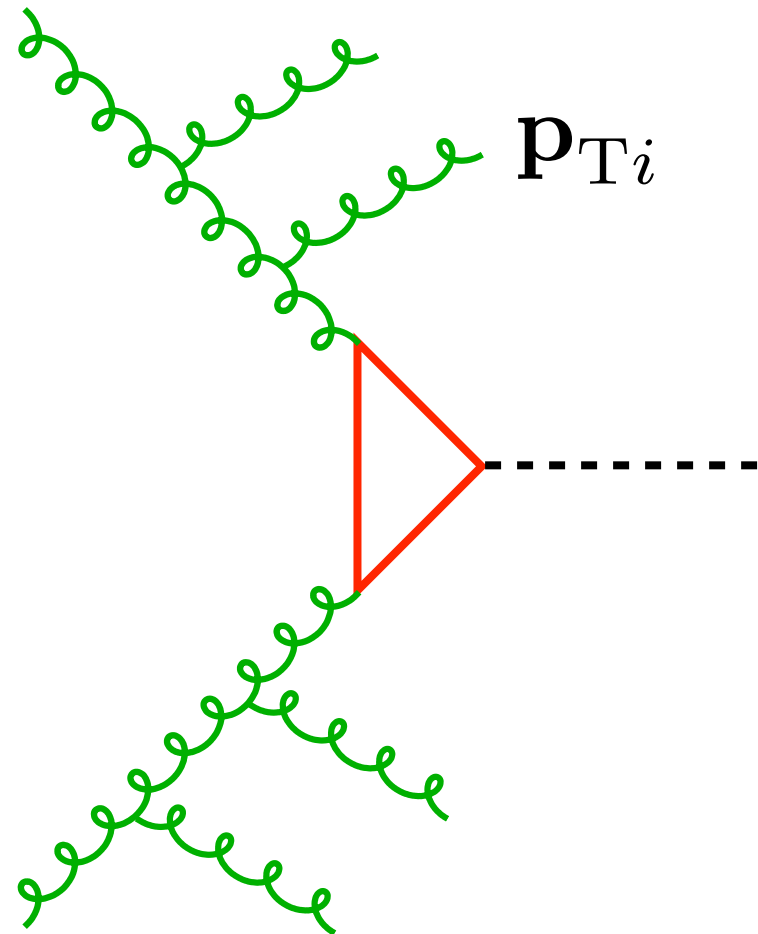
Series extrapolation:  $\sigma_{ggF}(13 \text{ TeV}) = 55.9 \pm ?? \text{ pb}$

My best guess:  $\sigma_{ggF}(13 \text{ TeV}) = 53 \pm 8 \text{ pb}$



# Higgs Differential Cross Sections

# Higgs $q_T$ & $E_T$



- Higgs transverse momentum

$$\mathbf{q}_T = - \sum \mathbf{p}_{Ti}$$

Bozzi et al. 0705.3887

Mantry & Petriello, 0911.4135

Catani & Grazzini, 1011.3918

de Florian et al. 1109.2109

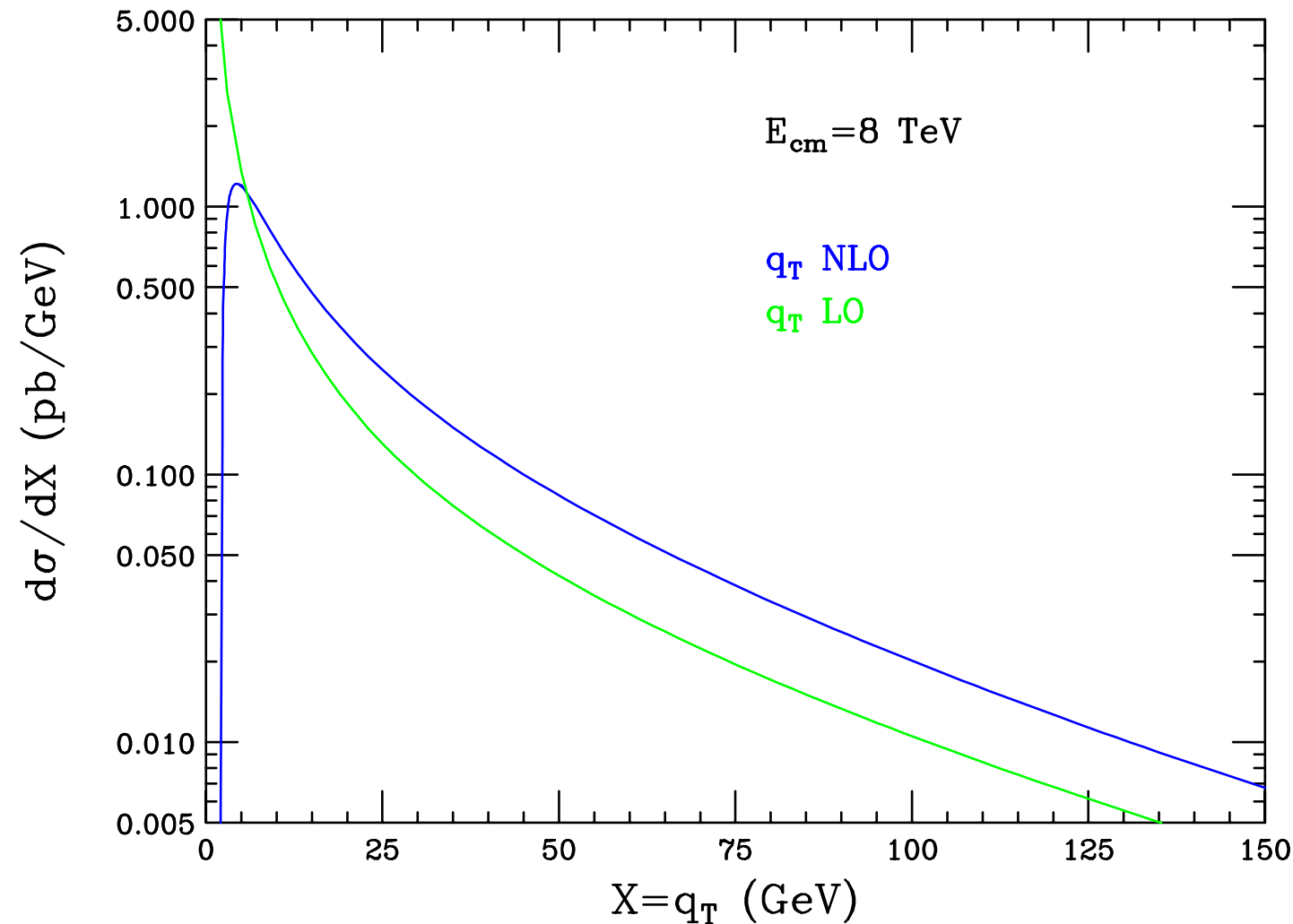
- Radiated transverse energy

$$E_T = \sum |\mathbf{p}_{Ti}|$$

Papaefstathiou, Smillie, BW, 1002.4375

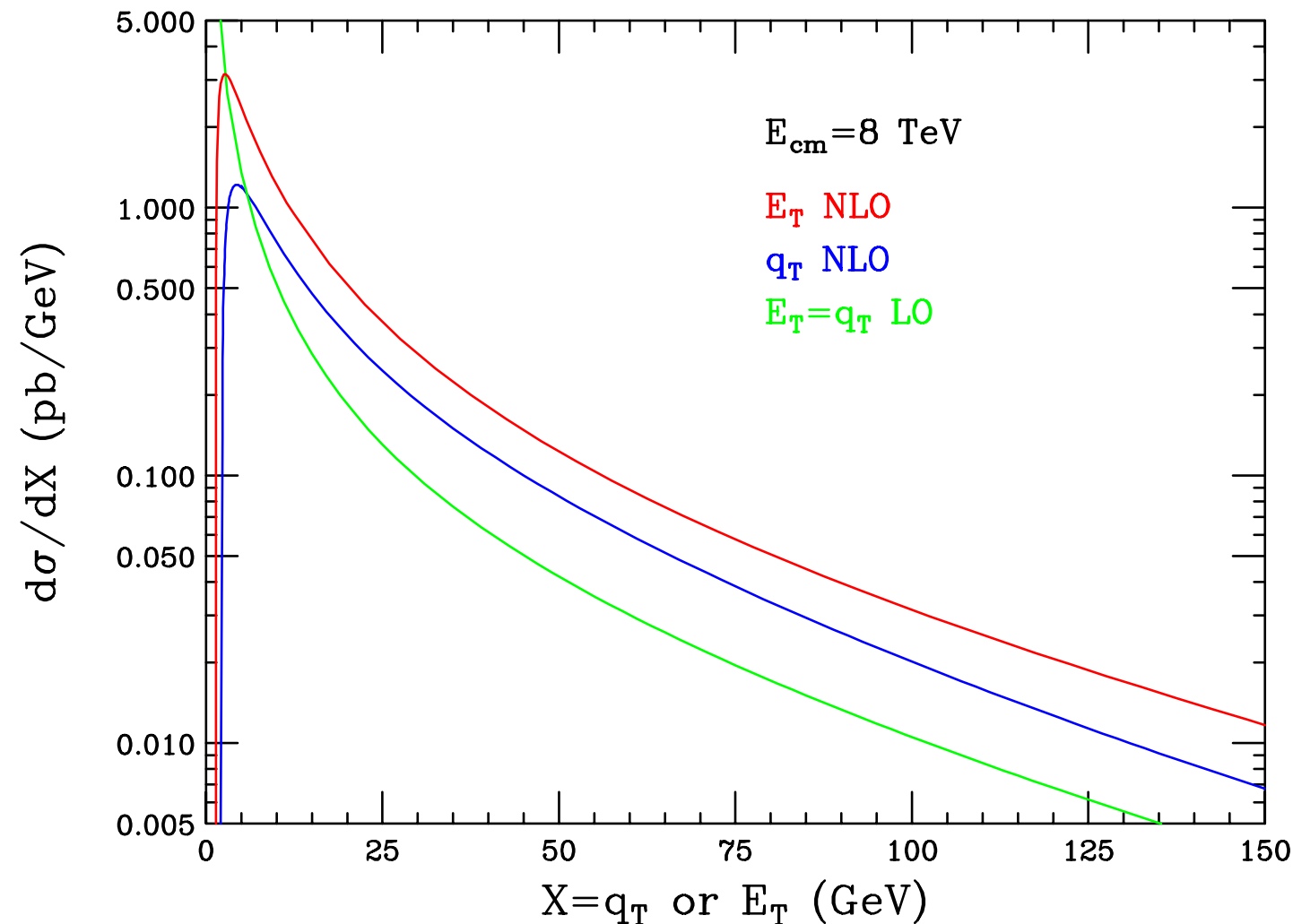
+Grazzini, 1403.3394

# Higgs $q_T$ (fixed order)



- $(\text{N})\text{LO} \xrightarrow{q_T \rightarrow 0} (-)\infty$
- Large logs of  $m_H^2/q_T^2$  need resummation

# Higgs $q_T$ & $E_T$ (fixed order)



- $(N)LO \xrightarrow{E_T \rightarrow 0} (-)\infty$
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# Resummation of Higgs $q_T$

$$d\sigma = \int dx_1 dx_2 f_a(x_1, \mu) f_b(x_2, \mu) d\hat{\sigma}_{ab}(x_1 x_2 s, \mu, \dots)$$

$$\frac{1}{\hat{\sigma}_{gg}} \frac{d^2 \hat{\sigma}_{gg}}{d\mathbf{q}_T^2} \sim \delta^2(\mathbf{q}_T) + \alpha_S \int d^2 \mathbf{p}_T \left[ \frac{A_g}{\mathbf{p}_T^2} \ln \frac{m_H^2}{\mathbf{p}_T^2} + \frac{B_g}{\mathbf{p}_T^2} \right]_+ \delta^2(\mathbf{q}_T + \mathbf{p}_T) + \dots$$

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# Resummation of Higgs $q_T$

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$$\sim \int \frac{d^2 \mathbf{b}}{(2\pi)^2} e^{i\mathbf{b} \cdot \mathbf{q}_T} \exp \left\{ \alpha_S \int d^2 \mathbf{p}_T \left[ \frac{A_g}{\mathbf{p}_T^2} \ln \frac{m_H^2}{\mathbf{p}_T^2} + \frac{B_g}{\mathbf{p}_T^2} \right] (e^{i\mathbf{b} \cdot \mathbf{p}_T} - 1) \right\}$$



# Resummation & matching of Higgs $q_T$

$$d\sigma = \int dx_1 dx_2 f_a(x_1, \mu) f_b(x_2, \mu) d\hat{\sigma}_{ab}(x_1 x_2 s, \mu, \dots)$$

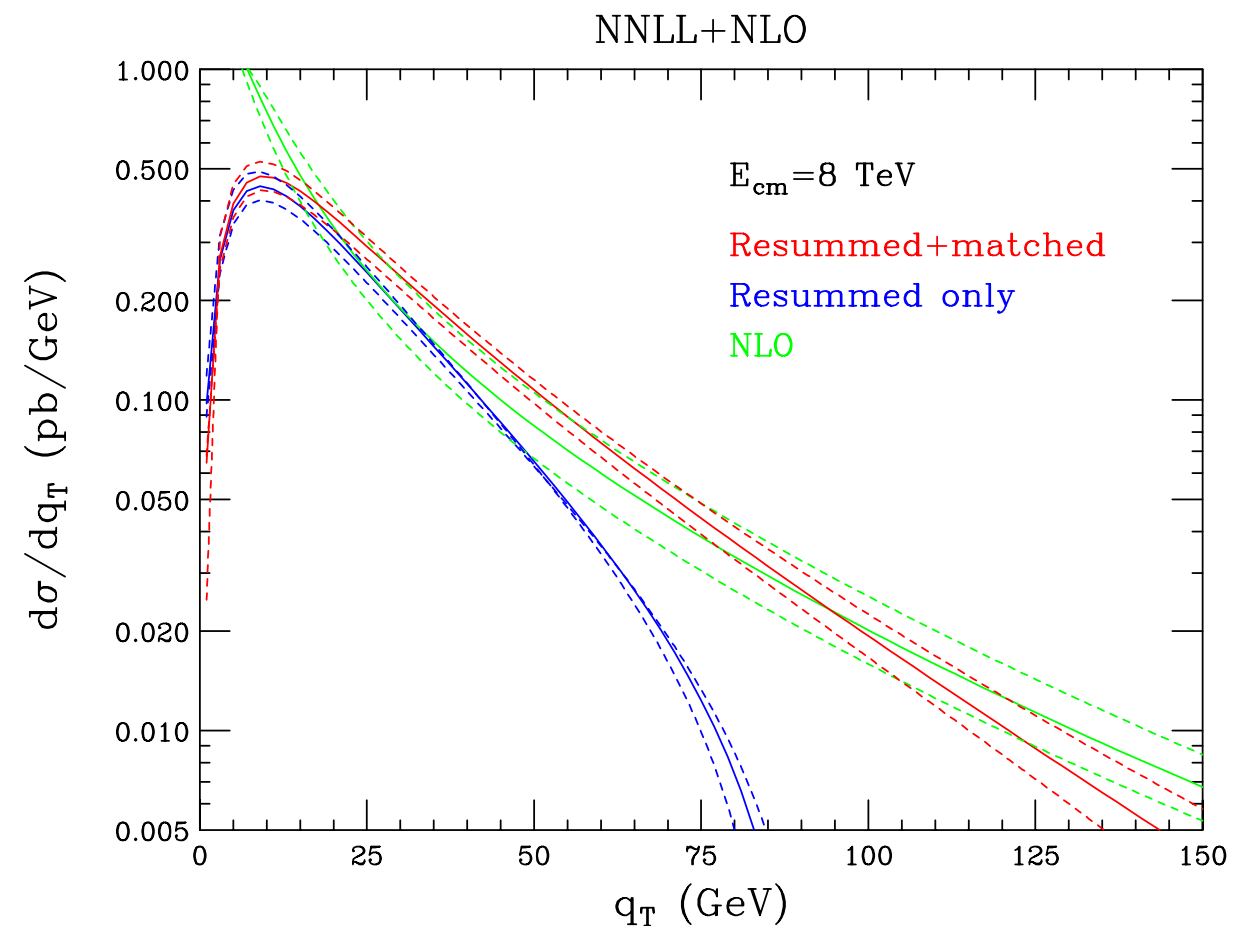
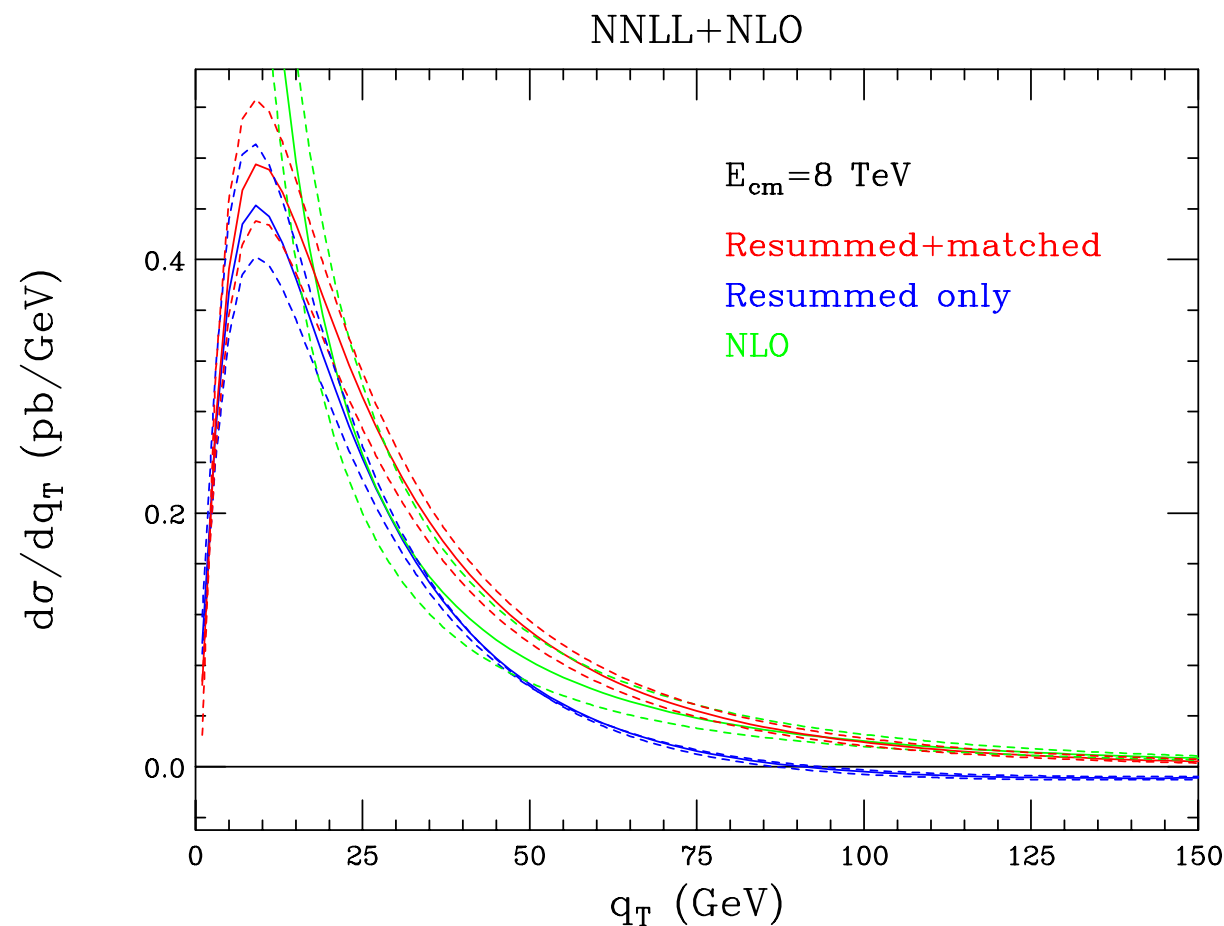
$$\frac{1}{\hat{\sigma}_{gg}} \frac{d^2 \hat{\sigma}_{gg}}{d\mathbf{q}_T^2} \sim \delta^2(\mathbf{q}_T) + \alpha_S \int d^2 \mathbf{p}_T \left[ \frac{A_g}{\mathbf{p}_T^2} \ln \frac{m_H^2}{\mathbf{p}_T^2} + \frac{B_g}{\mathbf{p}_T^2} \right]_+ \delta^2(\mathbf{q}_T + \mathbf{p}_T) + \dots$$

$$\sim \int \frac{d^2 \mathbf{b}}{(2\pi)^2} e^{i\mathbf{b} \cdot \mathbf{q}_T} \left\{ 1 + \alpha_S \int d^2 \mathbf{p}_T \left[ \frac{A_g}{\mathbf{p}_T^2} \ln \frac{m_H^2}{\mathbf{p}_T^2} + \frac{B_g}{\mathbf{p}_T^2} \right] (e^{i\mathbf{b} \cdot \mathbf{p}_T} - 1) + \dots \right\}$$

$$\sim \int \frac{d^2 \mathbf{b}}{(2\pi)^2} e^{i\mathbf{b} \cdot \mathbf{q}_T} \exp \left\{ \alpha_S \int d^2 \mathbf{p}_T \left[ \frac{A_g}{\mathbf{p}_T^2} \ln \frac{m_H^2}{\mathbf{p}_T^2} + \frac{B_g}{\mathbf{p}_T^2} \right] (e^{i\mathbf{b} \cdot \mathbf{p}_T} - 1) \right\}$$

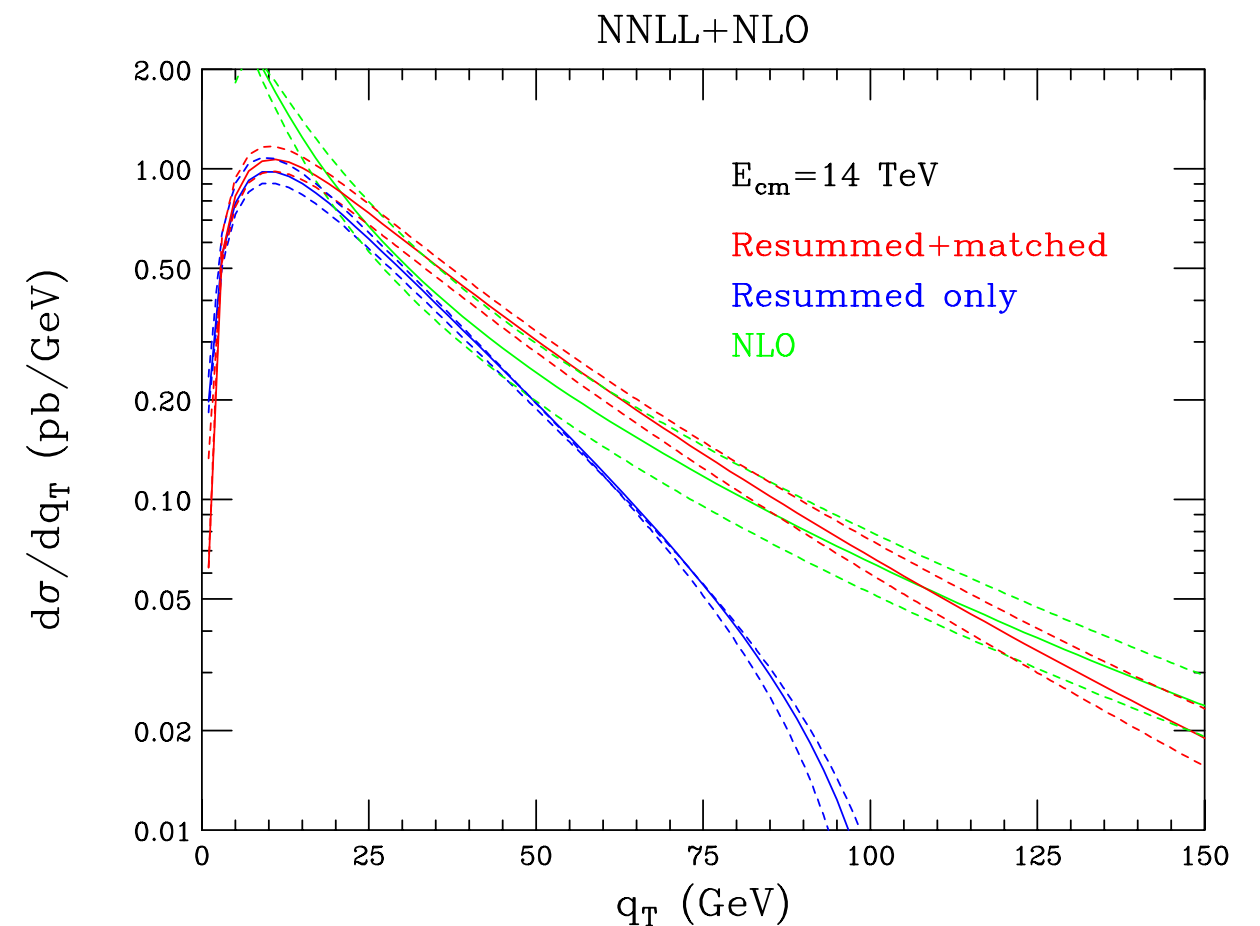
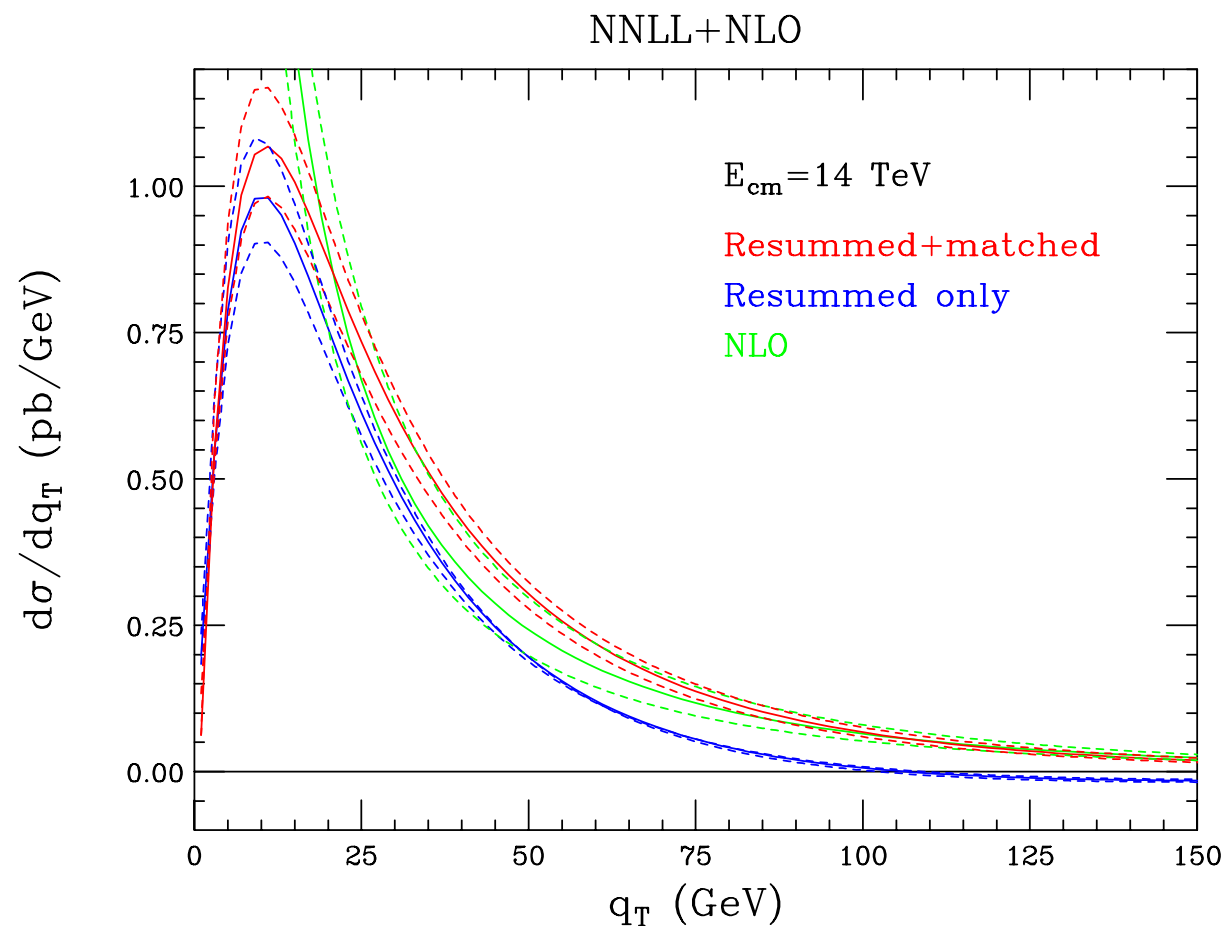
$$\frac{d\sigma}{dq_T} = \left[ \frac{d\sigma}{dq_T} \right]_{\text{resum}} - \left[ \frac{d\sigma}{dq_T} \right]_{\text{resum,NLO}} + \left[ \frac{d\sigma}{dq_T} \right]_{\text{NLO}}$$

# Higgs transverse momentum: 8 TeV



- Peak at  $\sim 10 \text{ GeV}$ :  $\log(m_H^2/q_T^2) \sim 5.1$
- Resummation affects spectrum out to larger  $q_T$

# Higgs transverse momentum: 14 TeV



- Peak at  $\sim 10 \text{ GeV}$ :  $\log(m_H^2/q_T^2) \sim 5.1$
- Resummation affects spectrum out to larger  $q_T$

# Resummation of Higgs $q_T$

$$d\sigma = \int dx_1 dx_2 f_a(x_1, \mu) f_b(x_2, \mu) d\hat{\sigma}_{ab}(x_1 x_2 s, \mu, \dots)$$

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# Resummation of Higgs $E_T$

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$$\frac{1}{\hat{\sigma}_{gg}} \frac{d\hat{\sigma}_{gg}}{dE_T} \sim \delta(E_T) + \alpha_S \int d^2 \mathbf{p}_T \left[ \frac{A_g}{\mathbf{p}_T^2} \ln \frac{m_H^2}{\mathbf{p}_T^2} + \frac{B_g}{\mathbf{p}_T^2} \right]_+ \delta(E_T - |\mathbf{p}_T|) + \dots$$

$$\sim \int \frac{d\tau}{2\pi} e^{i\tau E_T} \exp \left\{ \alpha_S \int d^2 \mathbf{p}_T \left[ \frac{A_g}{\mathbf{p}_T^2} \ln \frac{m_H^2}{\mathbf{p}_T^2} + \frac{B_g}{\mathbf{p}_T^2} \right] (e^{-i\tau |\mathbf{p}_T|} - 1) \right\}$$

# Resummation of Higgs $E_T$

$$\frac{1}{\hat{\sigma}_{gg}} \frac{d\hat{\sigma}_{gg}}{dE_T} \sim \int_{-\infty}^{+\infty} \frac{d\tau}{2\pi} e^{i\tau E_T} \exp \left\{ \alpha_S \int d^2\mathbf{p}_T \left[ \frac{A_g}{\mathbf{p}_T^2} \ln \frac{m_H^2}{\mathbf{p}_T^2} + \frac{B_g}{\mathbf{p}_T^2} \right] \left( e^{-i\tau|\mathbf{p}_T|} - 1 \right) \right\}$$

- Defined for  $E_T \lesssim 0$

# Resummation of Higgs $E_T$

$$\frac{1}{\hat{\sigma}_{gg}} \frac{d\hat{\sigma}_{gg}}{dE_T} \sim \int_{-\infty}^{+\infty} \frac{d\tau}{2\pi} e^{i\tau E_T} \exp \left\{ \alpha_S \int d^2\mathbf{p}_T \left[ \frac{A_g}{\mathbf{p}_T^2} \ln \frac{m_H^2}{\mathbf{p}_T^2} + \frac{B_g}{\mathbf{p}_T^2} \right] \left( e^{-i\tau|\mathbf{p}_T|} - 1 \right) \right\}$$

- Defined for  $E_T \gtrless 0$
- For  $E_T < 0$ , can close  $\tau$ -contour in lower half-plane



# Resummation of Higgs $E_T$

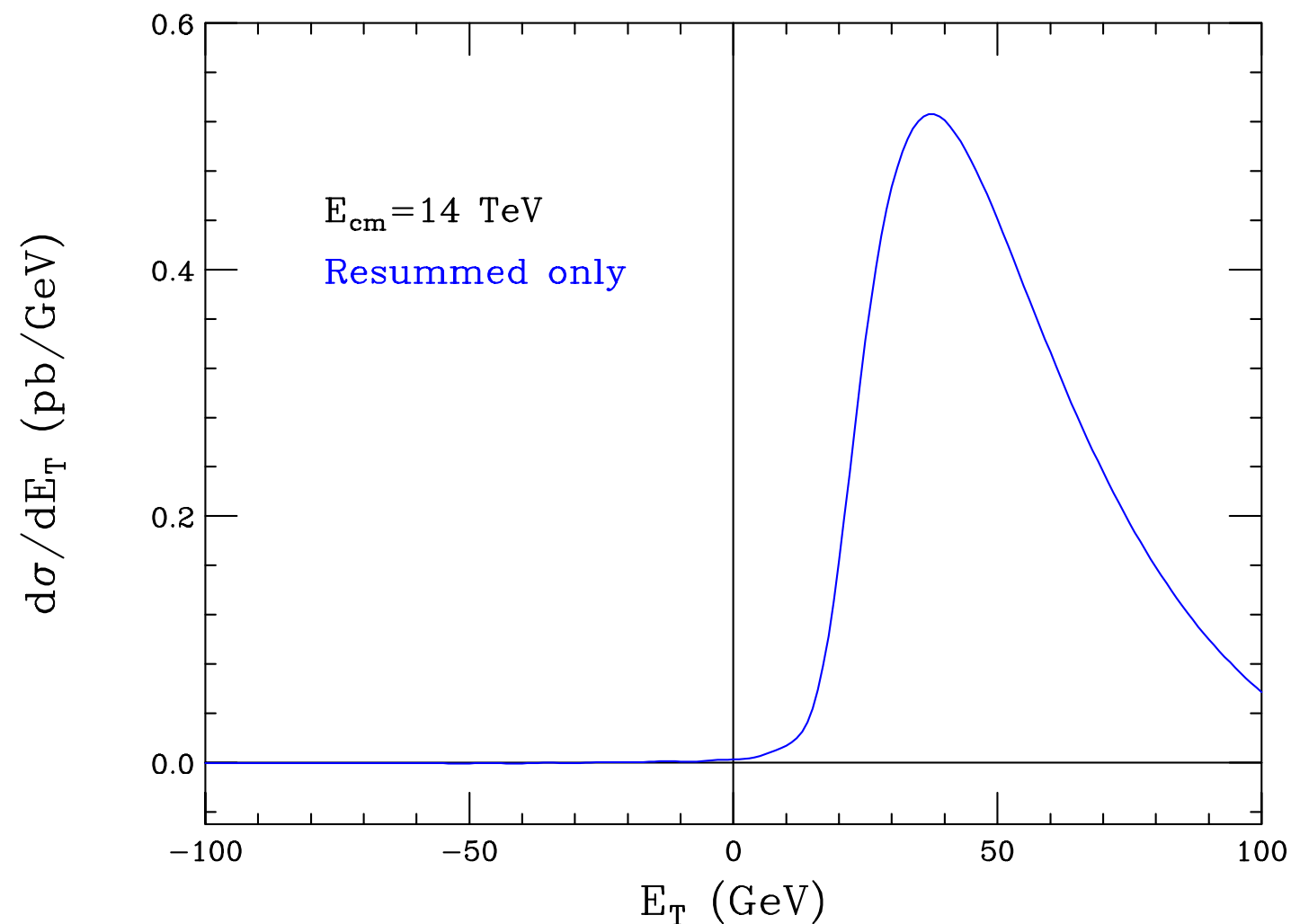
$$\frac{1}{\hat{\sigma}_{gg}} \frac{d\hat{\sigma}_{gg}}{dE_T} \sim \int_{-\infty}^{+\infty} \frac{d\tau}{2\pi} e^{i\tau E_T} \exp \left\{ \alpha_S \int d^2\mathbf{p}_T \left[ \frac{A_g}{\mathbf{p}_T^2} \ln \frac{m_H^2}{\mathbf{p}_T^2} + \frac{B_g}{\mathbf{p}_T^2} \right] \left( e^{-i\tau|\mathbf{p}_T|} - 1 \right) \right\}$$

- Defined for  $E_T \lesssim 0$
- For  $E_T < 0$ , can close  $\tau$ -contour in lower half-plane
- No singularities in lower half-plane

# Resummation of Higgs $E_T$

$$\frac{1}{\hat{\sigma}_{gg}} \frac{d\hat{\sigma}_{gg}}{dE_T} \sim \int_{-\infty}^{+\infty} \frac{d\tau}{2\pi} e^{i\tau E_T} \exp \left\{ \alpha_S \int d^2\mathbf{p}_T \left[ \frac{A_g}{\mathbf{p}_T^2} \ln \frac{m_H^2}{\mathbf{p}_T^2} + \frac{B_g}{\mathbf{p}_T^2} \right] \left( e^{-i\tau|\mathbf{p}_T|} - 1 \right) \right\}$$

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- For  $E_T < 0$ , can close  $\tau$ -contour in lower half-plane
- No singularities in lower half-plane



# Resummation & matching of Higgs $E_T$

$$\left[ \frac{d\sigma_H}{dQ^2 dE_T} \right]_{\text{res.}} = \frac{1}{2\pi} \sum_{a,b} \int_0^1 dx_1 \int_0^1 dx_2 \int_{-\infty}^{+\infty} d\tau e^{-i\tau E_T} f_{a/h_1}(x_1, \mu) f_{b/h_2}(x_2, \mu) W_{ab}^H(x_1 x_2 s; Q, \tau, \mu)$$

$$W_{ab}^H(s; Q, \tau, \mu) = \int_0^1 dz_1 \int_0^1 dz_2 C_{ga}(\alpha_S(\mu), z_1; \tau, \mu) C_{gb}(\alpha_S(\mu), z_2; \tau, \mu) \delta(Q^2 - z_1 z_2 s) \sigma_{gg}^H(Q, \alpha_S(Q)) S_g(Q, \tau)$$

$$S_g(Q, \tau) = \exp \left\{ -2 \int_0^Q \frac{dq}{q} \left[ 2A_g(\alpha_S(q)) \ln \frac{Q}{q} + B_g(\alpha_S(q)) \right] (1 - e^{iq\tau}) \right\}$$

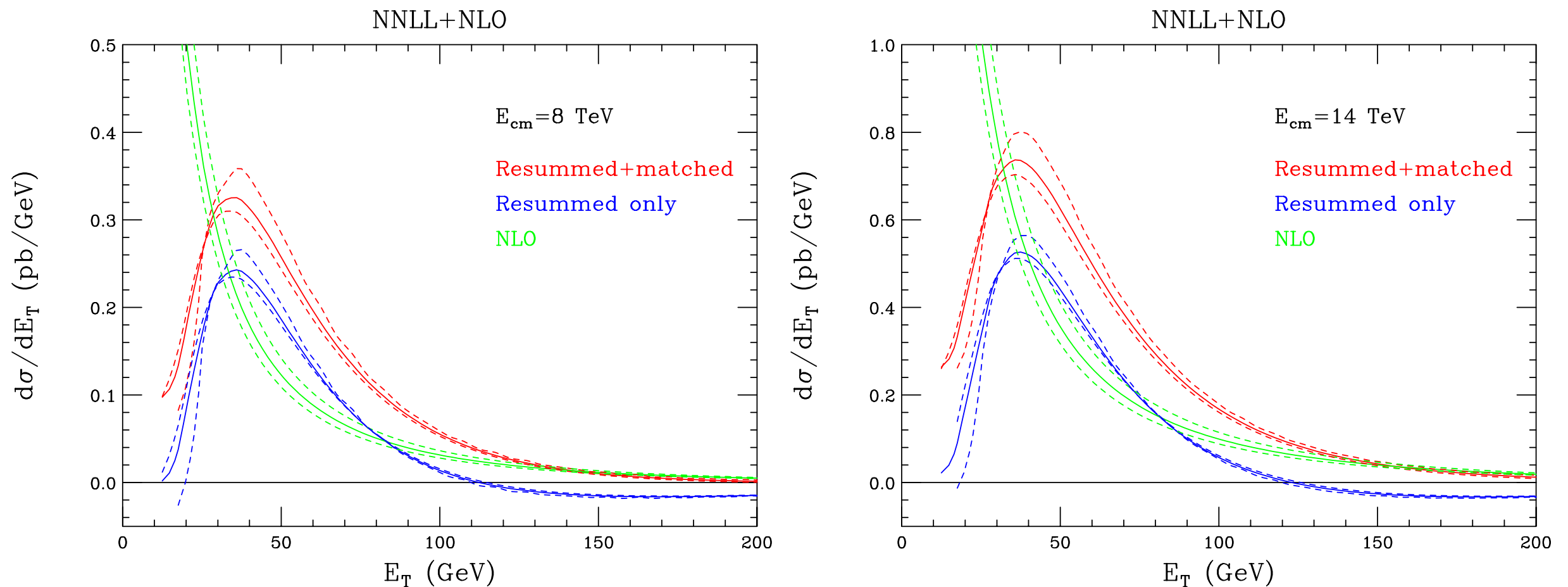
$$A_g(\alpha_S) = \sum_{n=1}^{\infty} \left( \frac{\alpha_S}{\pi} \right)^n A_g^{(n)} ,$$

$$B_g(\alpha_S) = \sum_{n=1}^{\infty} \left( \frac{\alpha_S}{\pi} \right)^n B_g^{(n)} ,$$

$$C_{ga}(\alpha_S, z) = \delta_{ga} \delta(1 - z) + \sum_{n=1}^{\infty} \left( \frac{\alpha_S}{\pi} \right)^n C_{ga}^{(n)}(z)$$

$$\frac{d\sigma_H}{dE_T} = \left[ \frac{d\sigma_H}{dE_T} \right]_{\text{resum}} - \left[ \frac{d\sigma_H}{dE_T} \right]_{\text{resum,NLO}} + \left[ \frac{d\sigma_H}{dE_T} \right]_{\text{NLO}}$$

# Transverse energy distribution

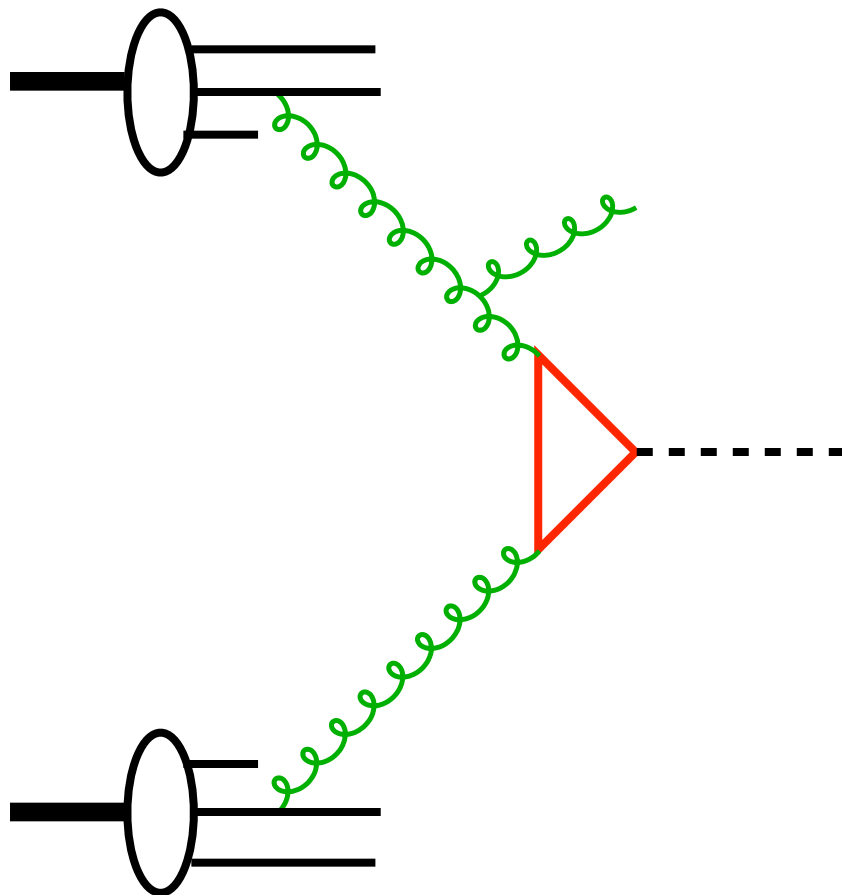


- Peak at  $\sim 35$  GeV:  $\log(m_H^2/E_T^2) \sim 2.6$
- Resummation affects spectrum out to much larger  $E_T$
- Unlike  $q_T$ , the **Underlying Event** also contributes...

# Monte Carlo Simulation

# Monte Carlo simulation

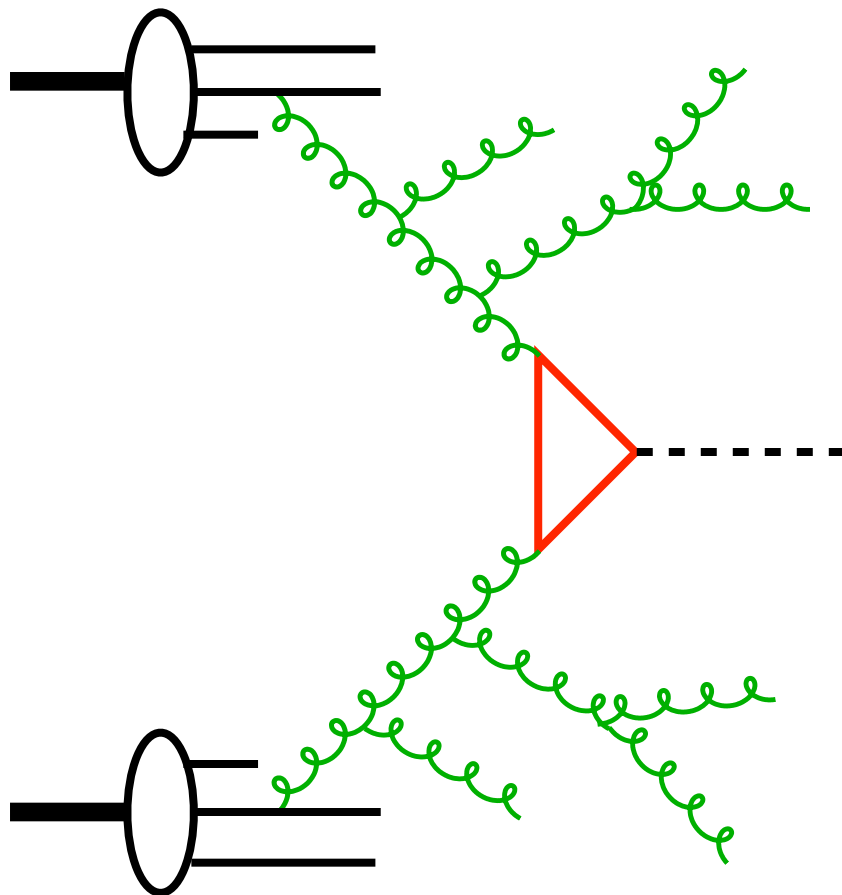
- aMC@NLO: NLO with matching



Frixione, BW, hep-ph/0204244

Frederix et al., 1110.4738

# Monte Carlo simulation



- aMC@NLO: NLO with matching to ...
- Parton showers ( $\approx$  higher orders)

Frixione, BW, hep-ph/0204244

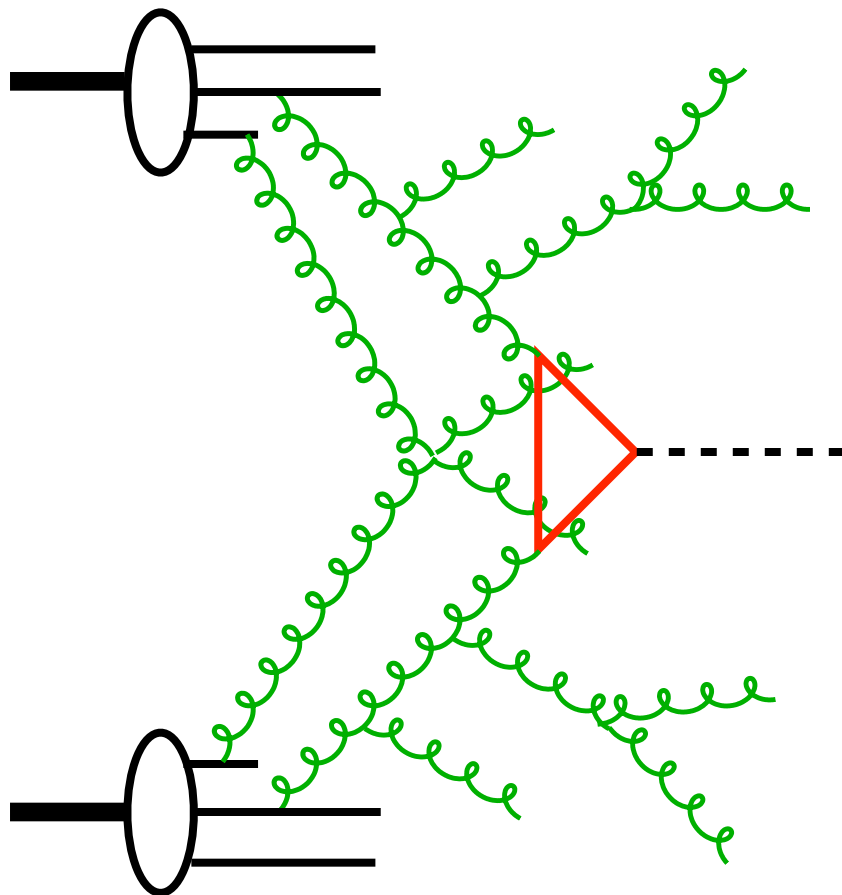
Frederix et al., 1110.4738

Bahr et al., 0812.0529,

<http://projects.hepforge.org/herwig/>



# Monte Carlo simulation



- aMC@NLO: NLO with matching to ...
- Parton showers ( $\approx$  higher orders)
- Underlying event model

Frixione, BW, hep-ph/0204244

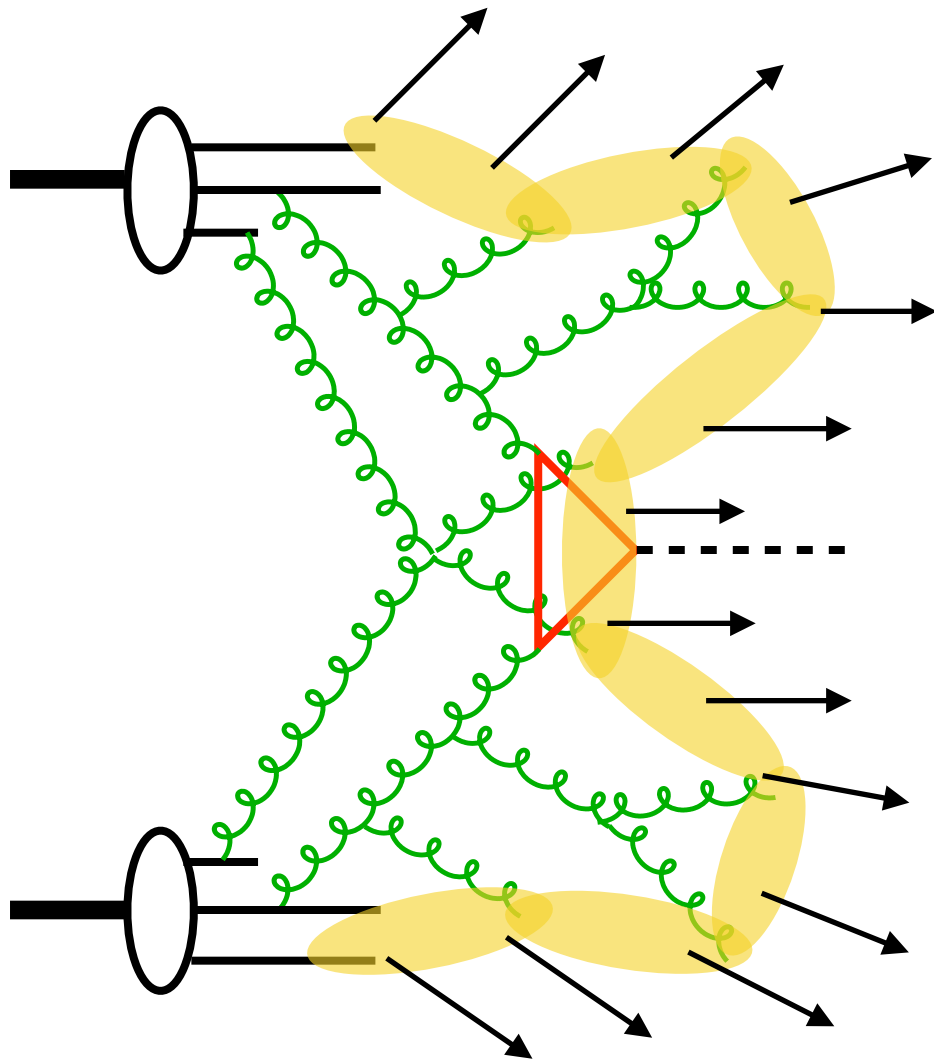
Frederix et al., 1110.4738

Bahr et al., 0812.0529,

<http://projects.hepforge.org/herwig/>

Borozan, Seymour, hep-ph/0207283

# Monte Carlo simulation



- aMC@NLO: NLO with matching to ...
- Parton showers ( $\approx$  higher orders)
- Underlying event model
- Cluster hadronization model

Frixione, BW, hep-ph/0204244

Frederix et al., 1110.4738

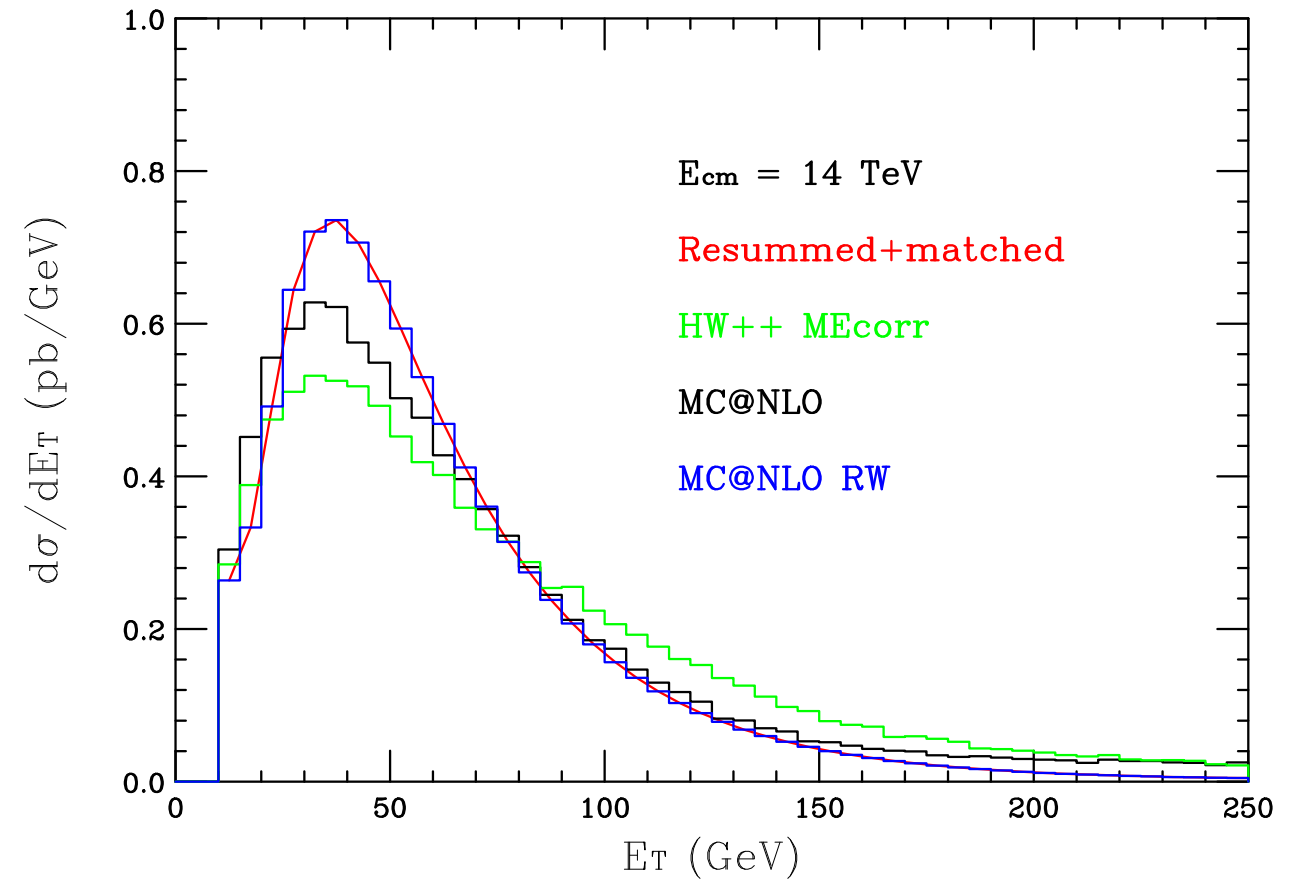
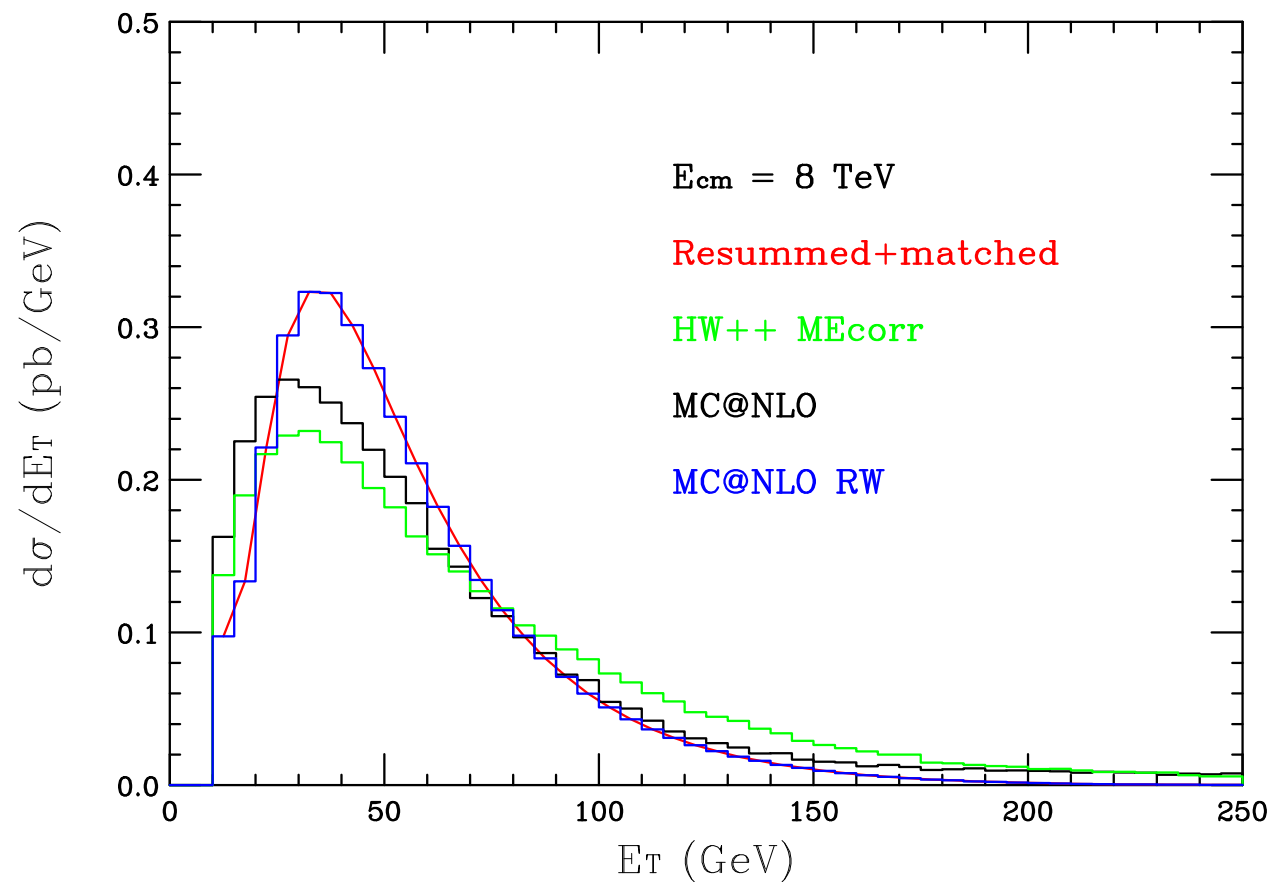
Bahr et al., 0812.0529,

<http://projects.hepforge.org/herwig/>

Borožan, Seymour, hep-ph/0207283

BW, Nucl.Phys. B238(1984)492

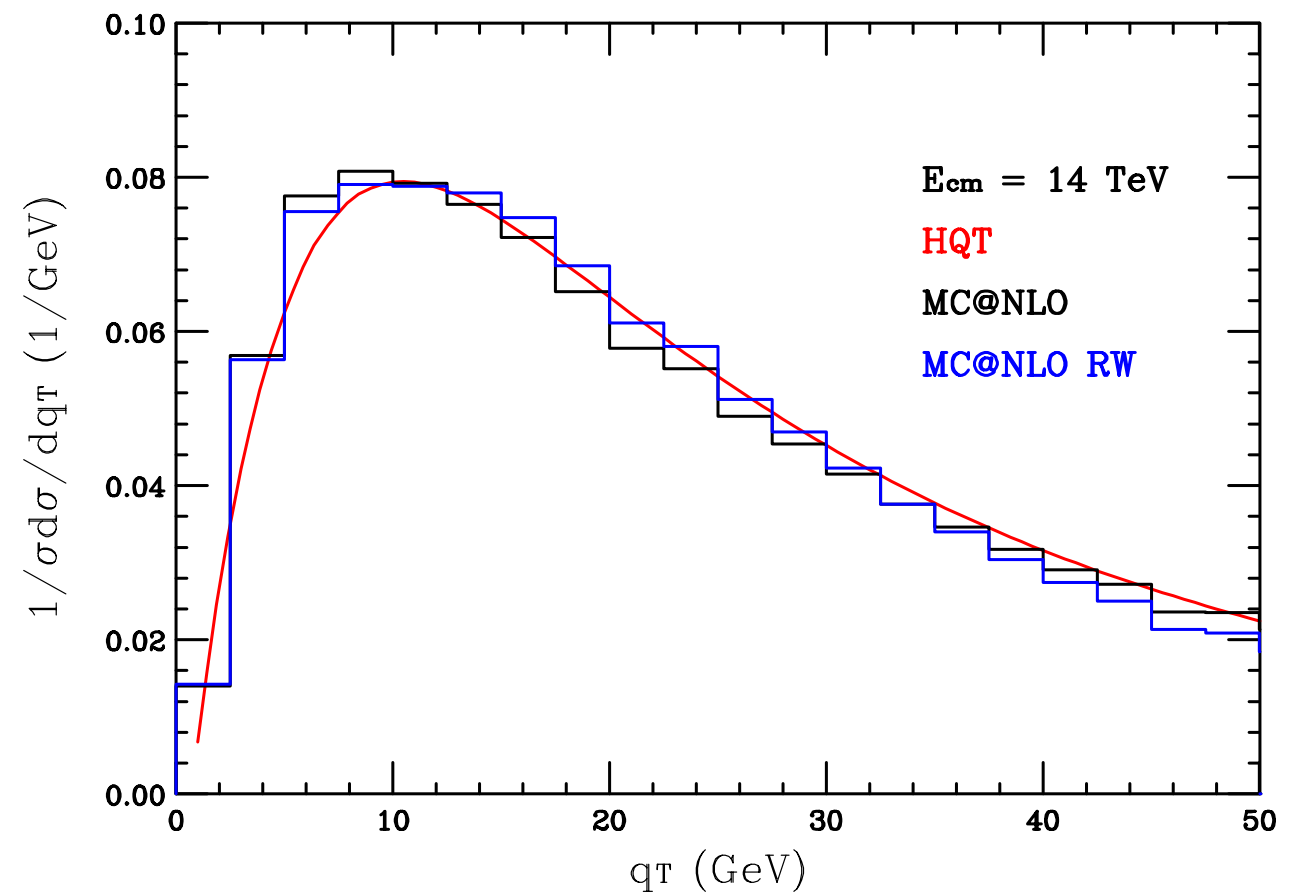
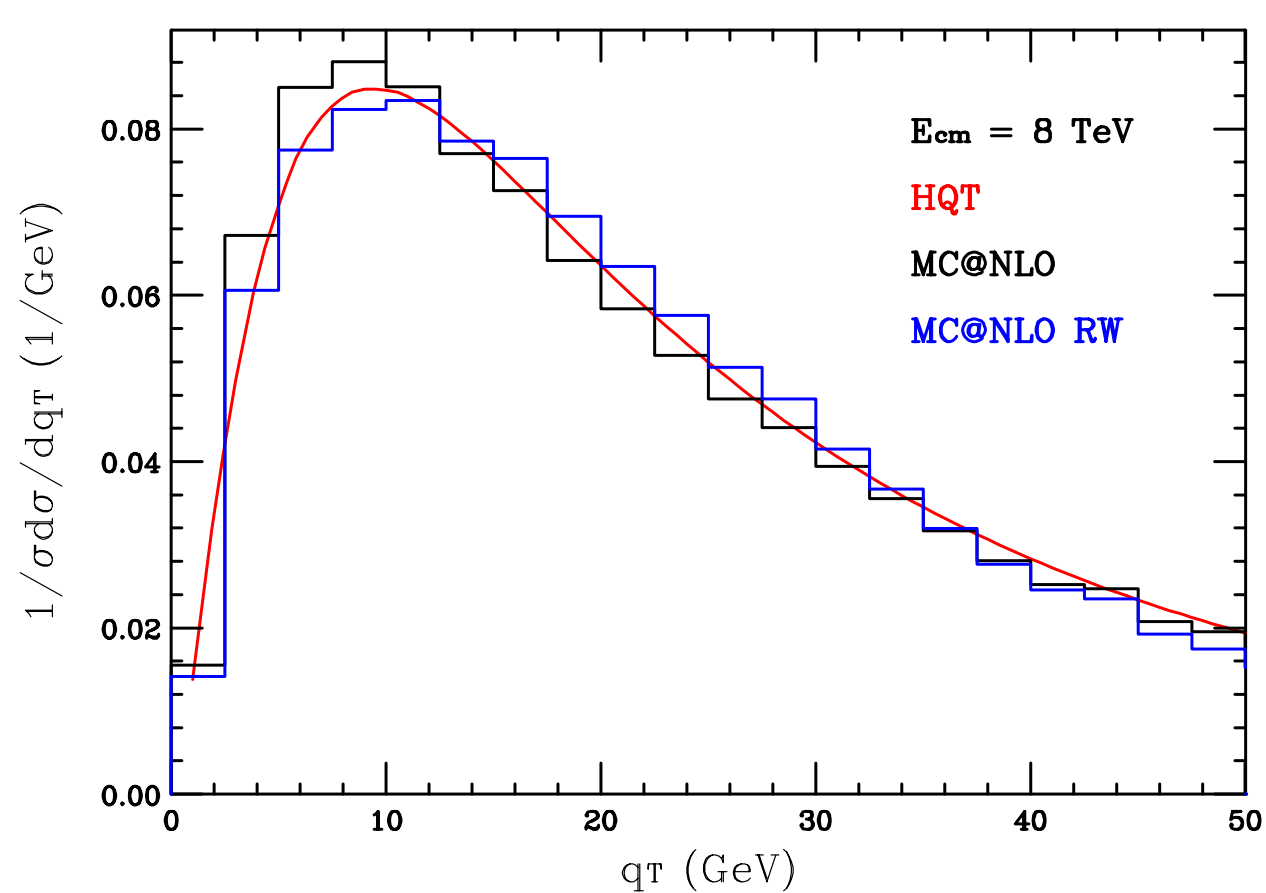
# Monte Carlo Higgs $E_T$



- RW = reweighted to agree with resummed+matched  $E_T$
- Underlying event and hadronization NOT included

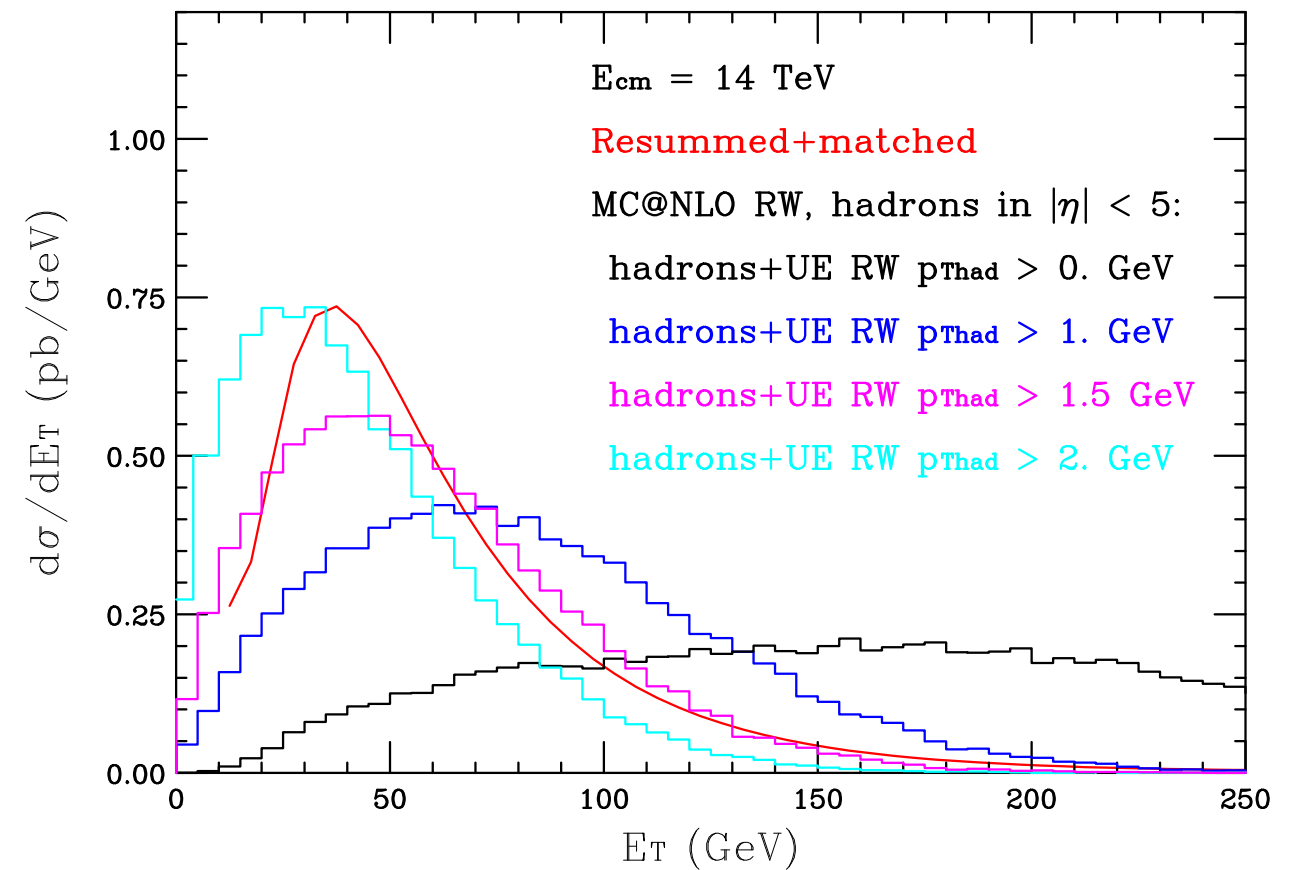
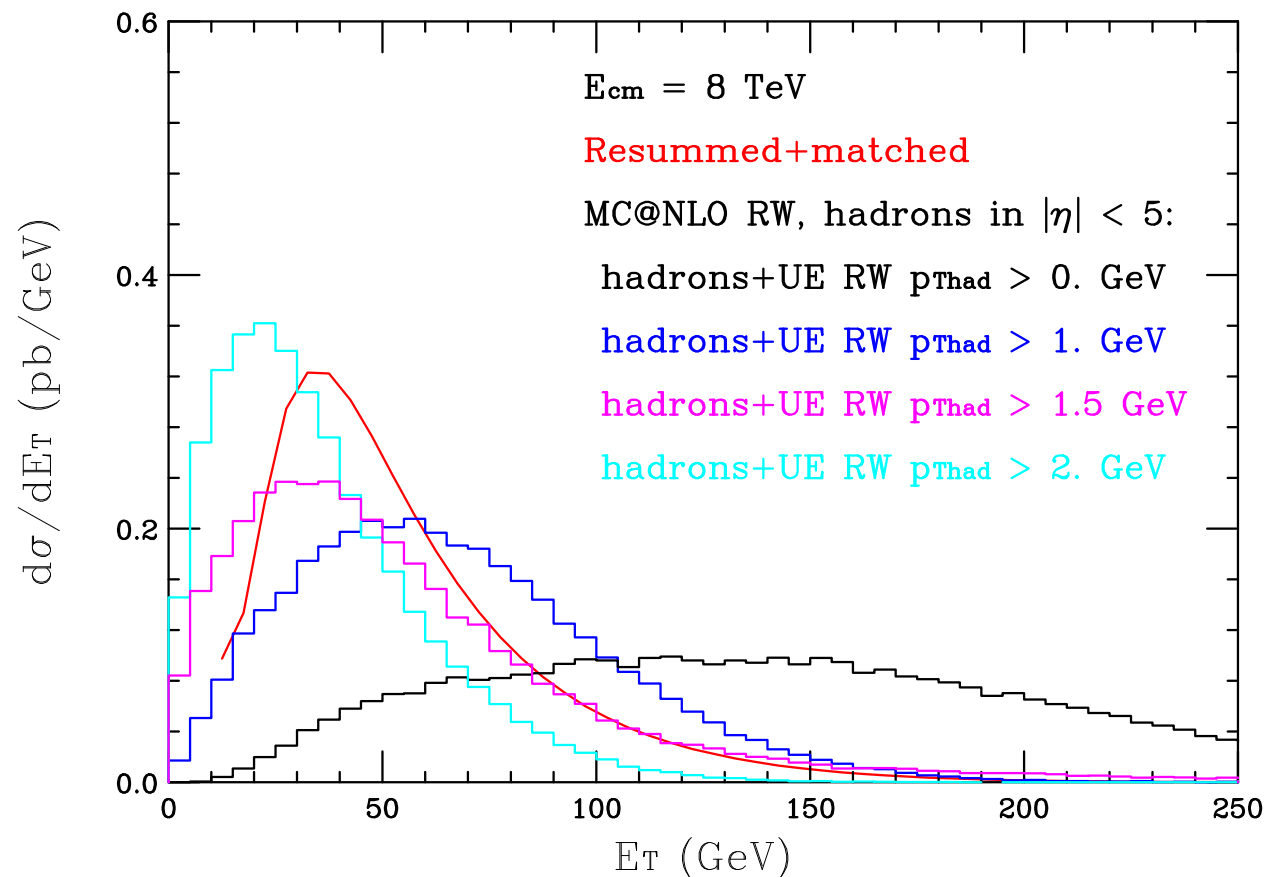
Grazzini, Papaefstathiou, Smillie, BW, 1403.3394

# Monte Carlo Higgs $q_T$



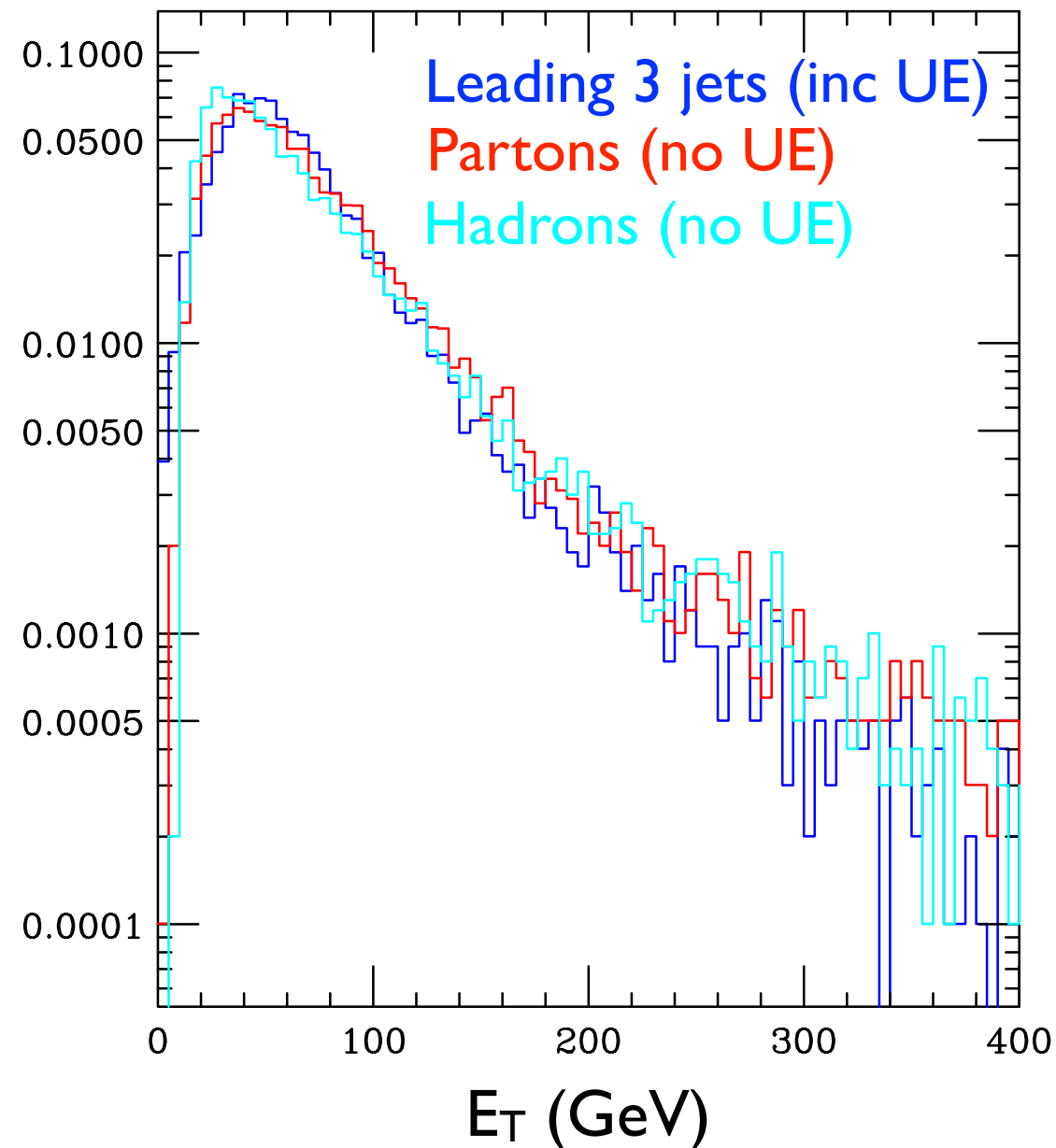
- RW = reweighted to agree with resummed+matched  $E_T$
- Reweighting improves agreement with HQT (= resummed+matched  $q_T$ )

# Monte Carlo Higgs $E_T$



- RW = reweighted to agree with resummed+matched  $E_T$
- Underlying event and hadronization INCLUDED
- Strong dependence on minimum hadron  $p_T$

# Higgs $E_T$ from jets?

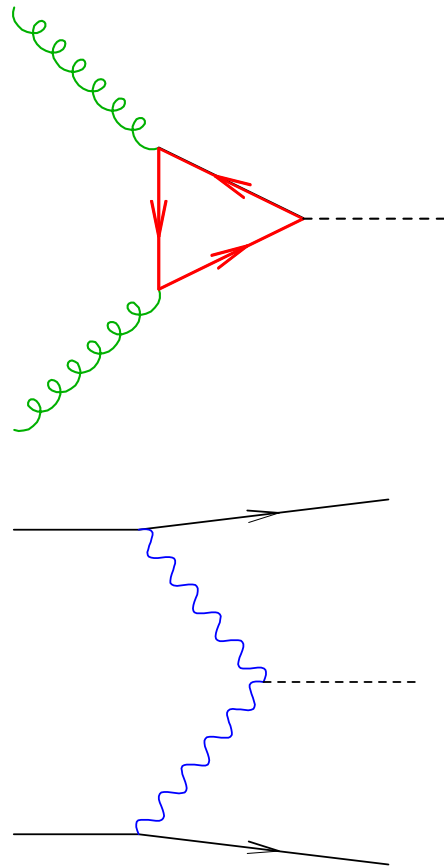


- Suggested by G Salam:
- Parton level  $\approx E_T$  of leading  $n$  jets (anti-kt,  $R=0.7$ )
- Less sensitive to underlying event and hadronization

Papaefstathiou, BW, prelim.

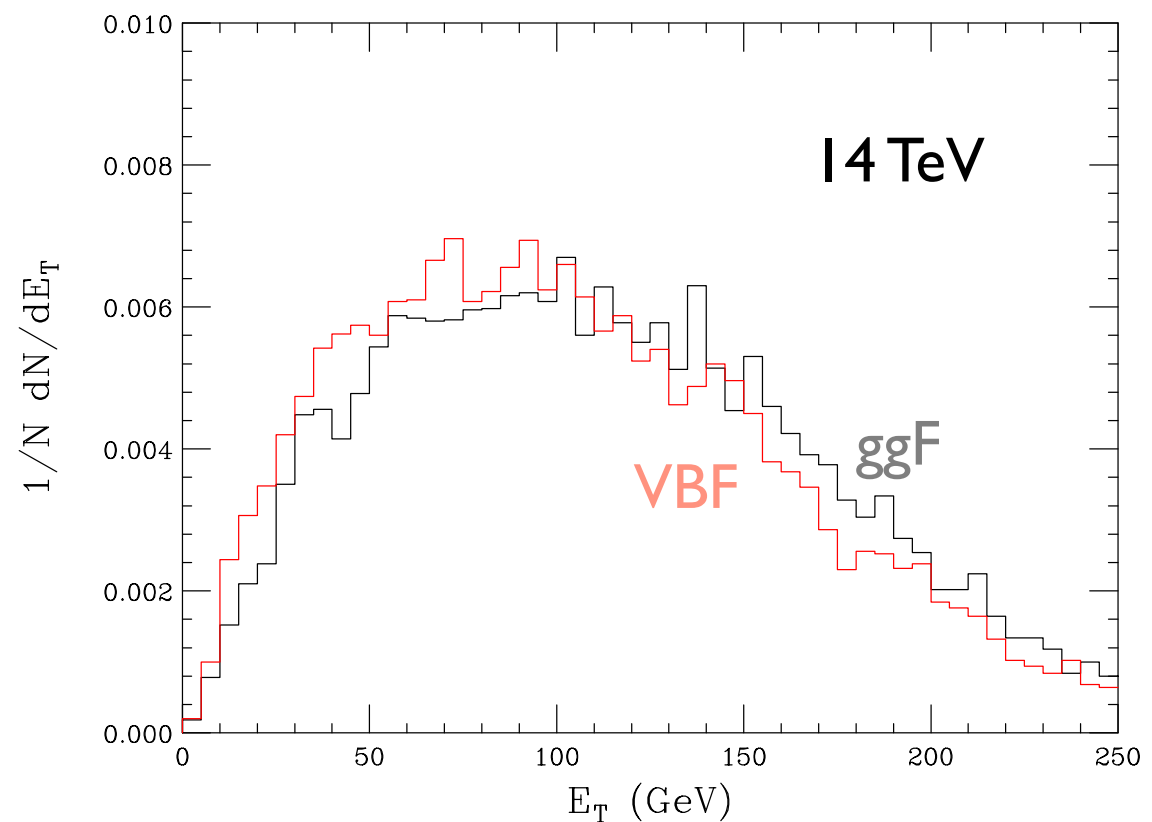
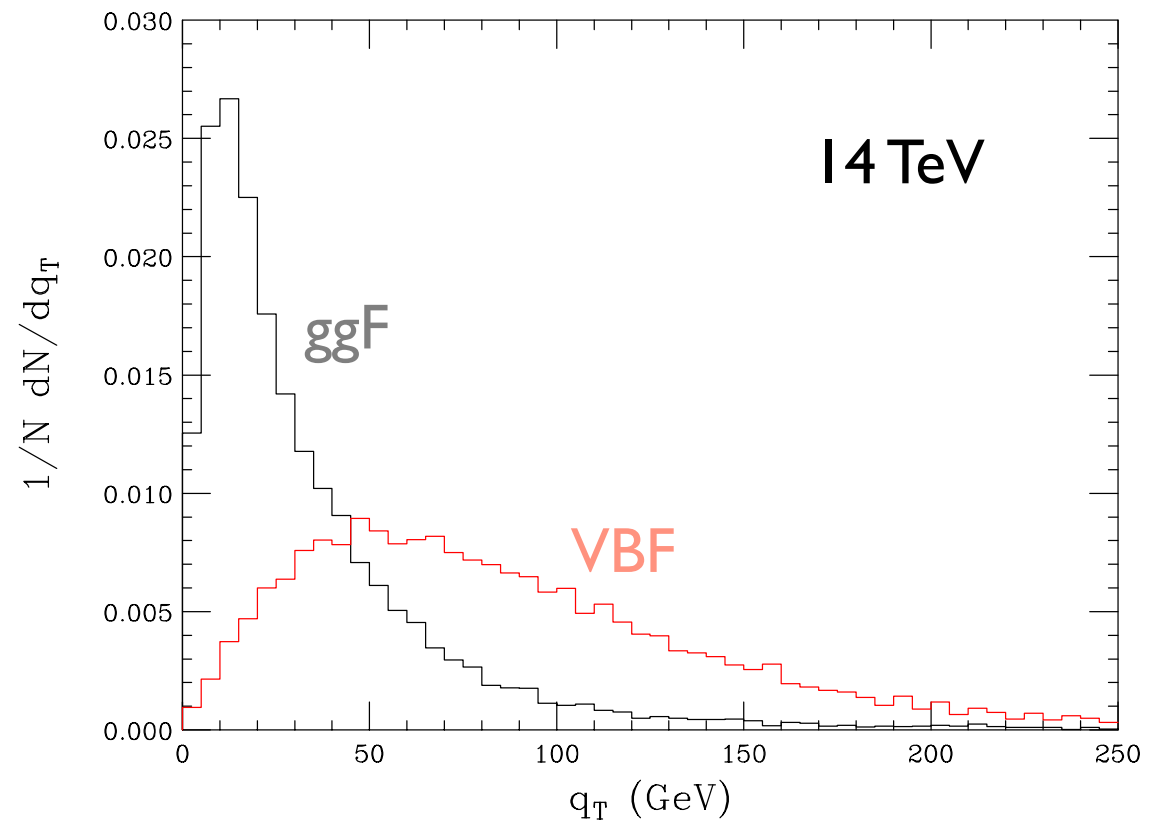
# Selecting VBF

# Monte Carlo Higgs $q_T$ & $E_T$



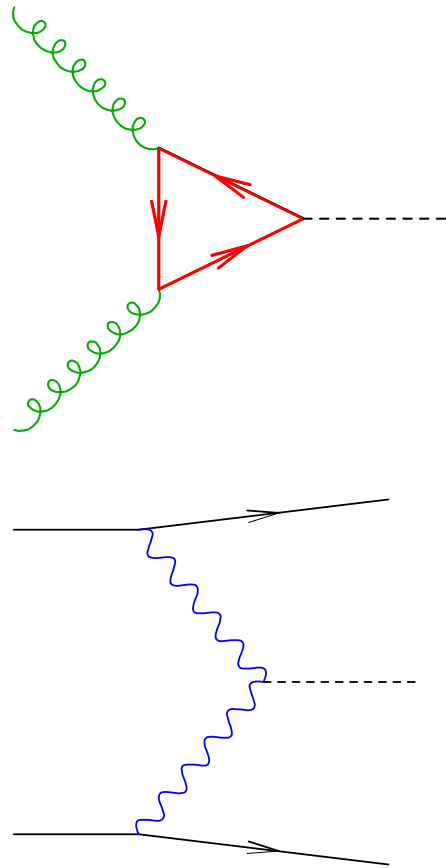
- Massless coloured vs massive colourless exchange
  - Big difference in  $q_T$
  - Small difference in  $E_T$

A Papaefstathiou, BW

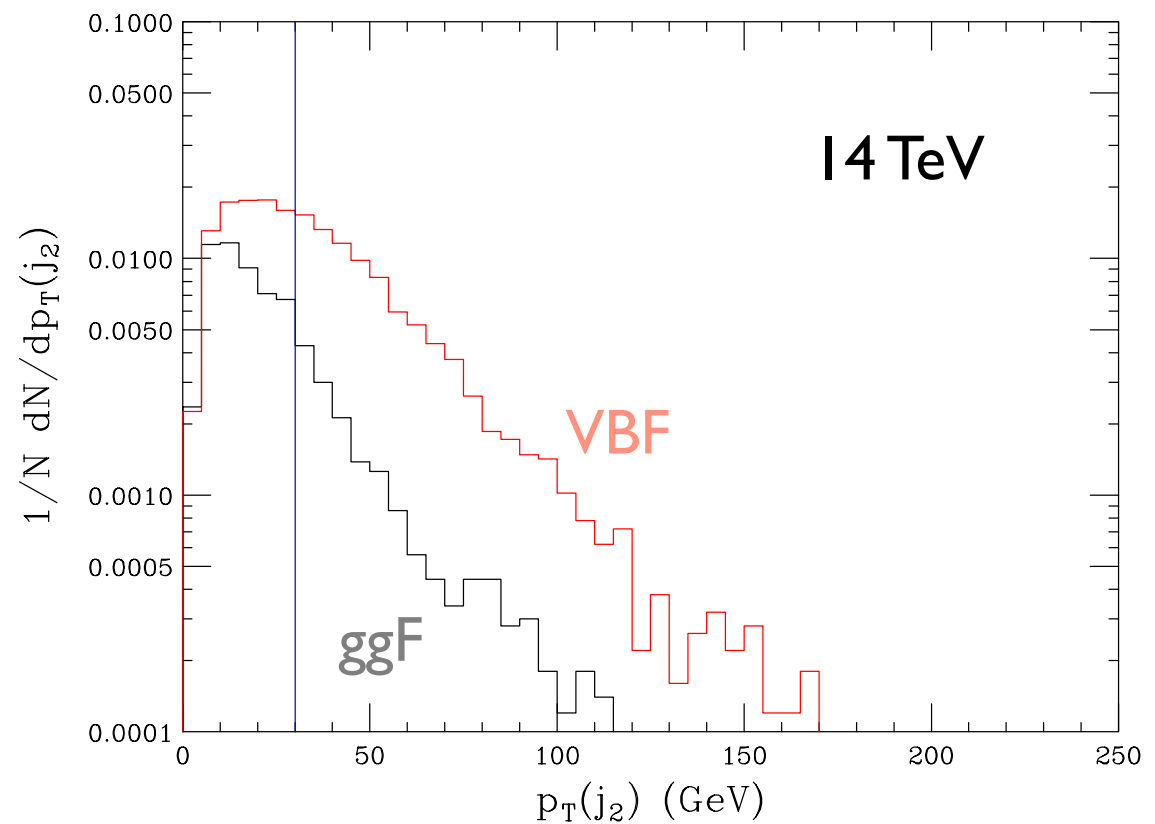
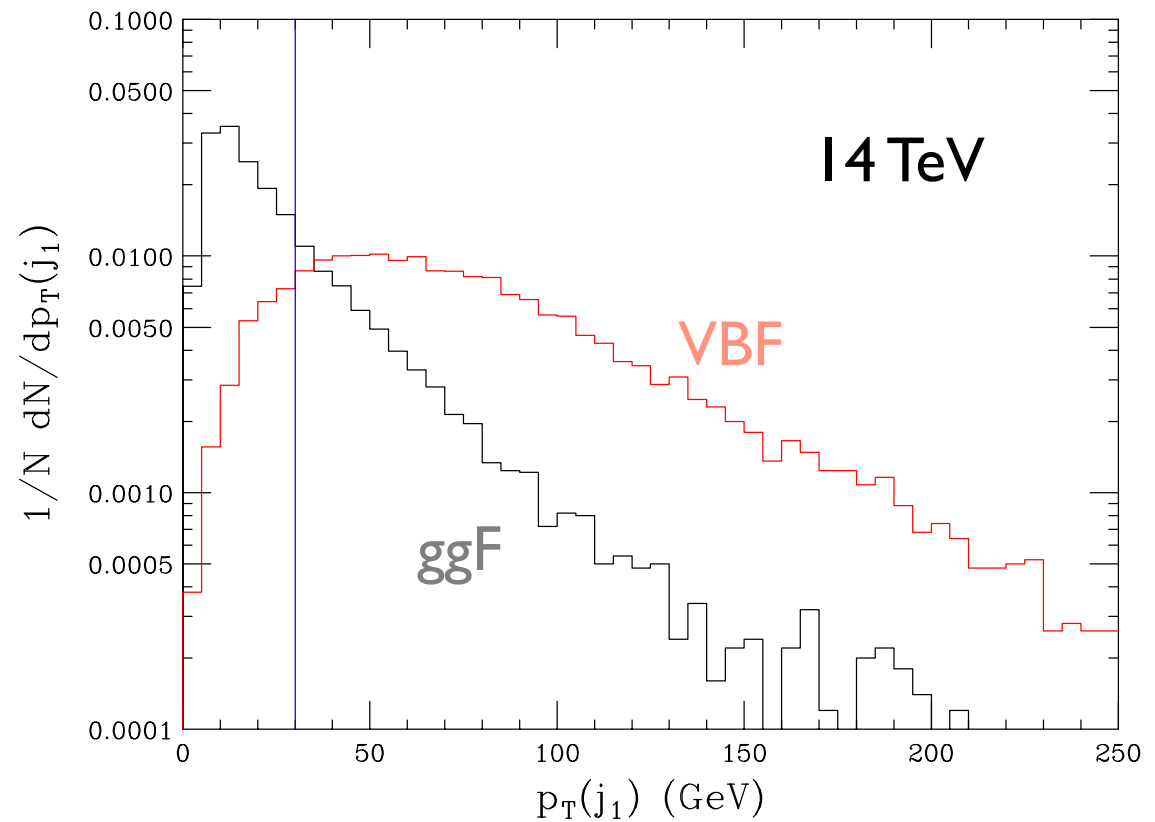




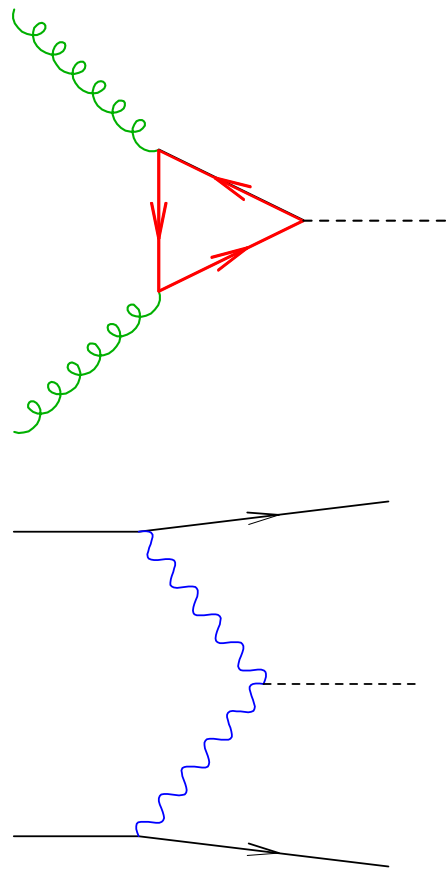
# Leading jets $p_T$



- Two hard leading jets in VBF

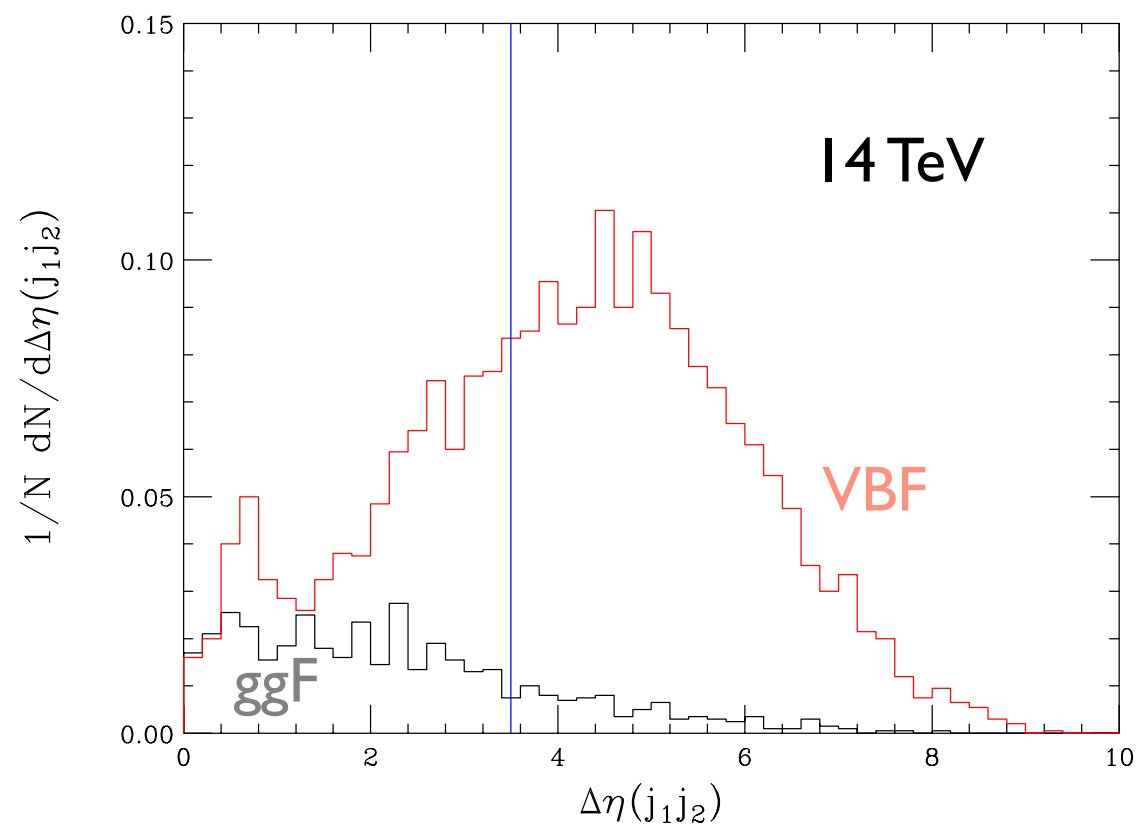
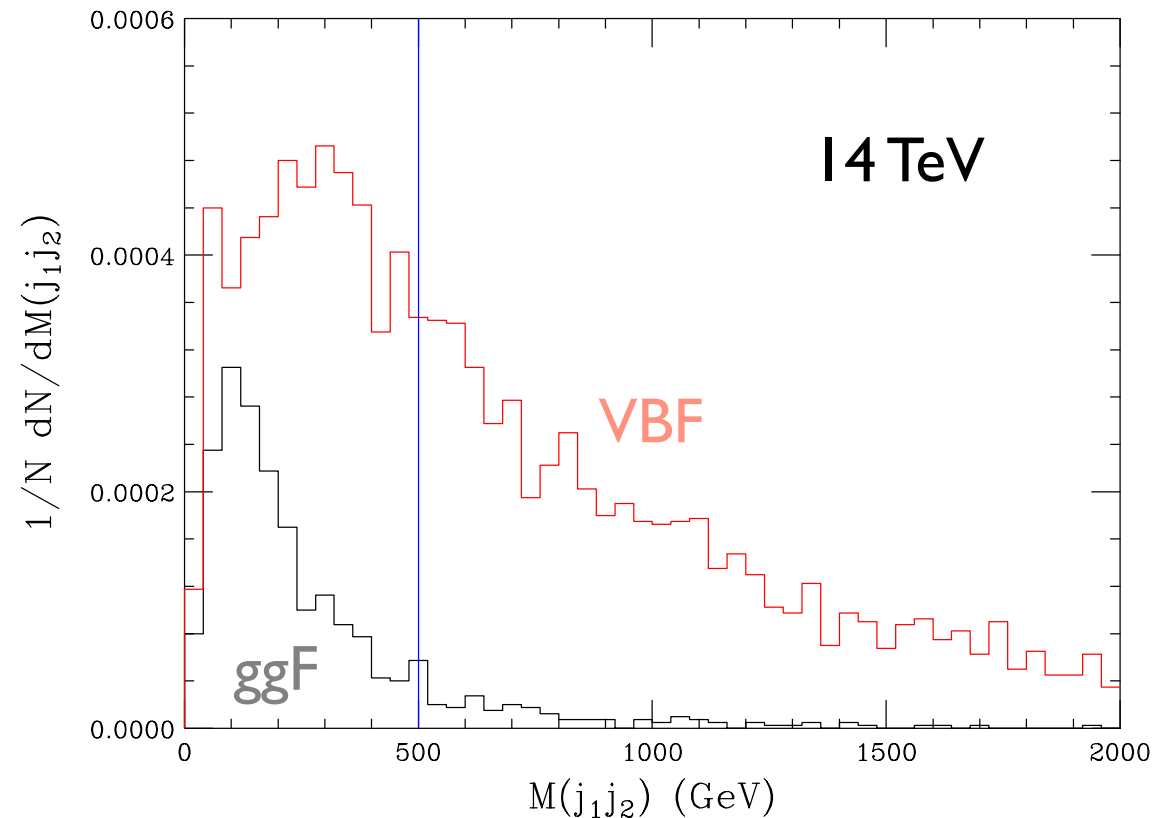


# Leading jets $M_{jj}$ and $\Delta\eta_{jj}$



- Cuts ( $\sim$ CMS):

- ▶  $p_T(j_1), p_T(j_2) > 30$  GeV
- ▶  $M(j_1 j_2) > 500$  GeV
- ▶  $\Delta\eta(j_1 j_2) > 3.5$



# Selecting VBF

- Cuts enhance VBF/ggF by  $\sim 25$
- For  $300 \text{ fb}^{-1}$  at 14 TeV:

Mode	BR%	ggF(raw)	VBF(raw)	ggF(cut)	VBF(cut)
<b>bb</b>	56.10	10,000,000	720,000	96,000	170,000
<b>WW</b>	23.13	4,200,000	300,000	40,000	71,000
<b>gg</b>	8.49	1,500,000	110,000	15,000	26,000
<b>tt</b>	6.16	1,100,000	79,000	11,000	19,000
<b>ZZ</b>	2.90	520,000	37,000	5,000	8,900
<b>cc</b>	2.83	510,000	37,000	4,800	8,700
$\gamma\gamma$	0.23	41,000	2,900	390	710
<b>Z<math>\gamma</math></b>	0.16	29,000	2,100	270	490

# Conclusions

# Conclusions

- Higgs ggF cross section at 13 TeV is still very uncertain
  - ▶ My estimate:  $\sigma_{\text{ggF}}(13 \text{ TeV}) = 53 \pm 11 \text{ pb}$
- Higgs transverse momentum resummed to NNLL+NLO
  - ▶ Peak  $q_T \sim 10 \text{ GeV}$ , independent of energy
- Radiated transverse energy resummed to NNLL+NLO
  - ▶ Peak  $E_T \sim 35 \text{ GeV}$  in associated transverse energy
  - ▶ Contribution from Underlying Event reduced for jets
- VBF has broader  $q_T$  but similar  $E_T$  distribution
  - ▶ Jet cuts can enhance VBF/ggF by  $\sim 25$

**Thanks for  
listening!**