# **Moonshine and String Theory**

#### **Timm Wrase**



#### IPMU

#### August 7, 2014

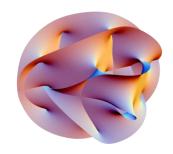
Based on: N. Paquette, TW to appear M. Cheng, X. Dong, J. Duncan, S. Harrison, S. Kachru, TW 1406.5502 TW 1402.2973 M. Cheng, X. Dong, J. Duncan, J. Harvey, S. Kachru, TW 1306.4981

# Outline

Introduction to moonshine



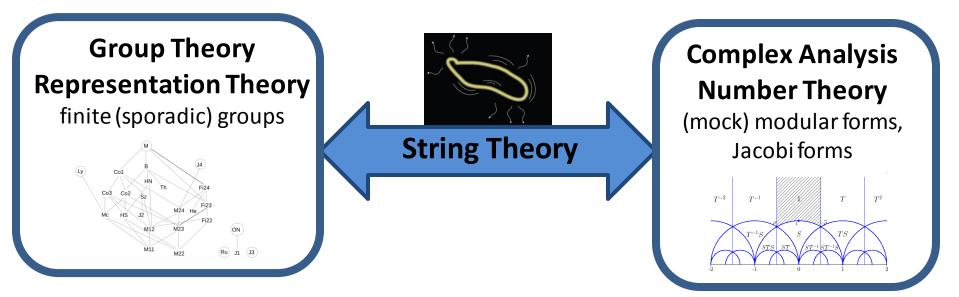
 Mathieu Moonshine and string compactifications



• New moonshine phenomena



# Moonshine



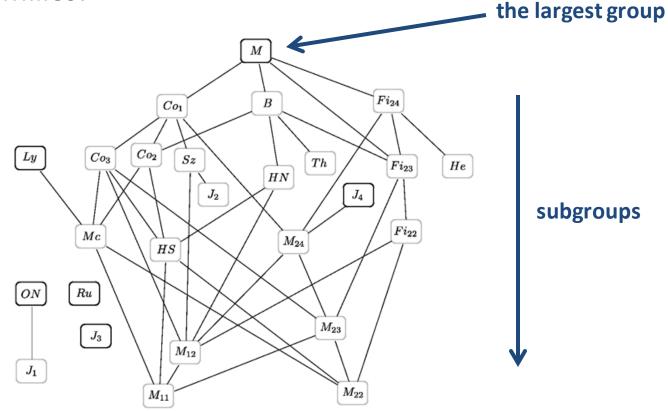
There are 18 infinite families, e.g.

• Alternating group of n elements  $A_n$ e.g.  $A_3$ : (123)  $\leftrightarrow$  (231)  $\leftrightarrow$  (312)

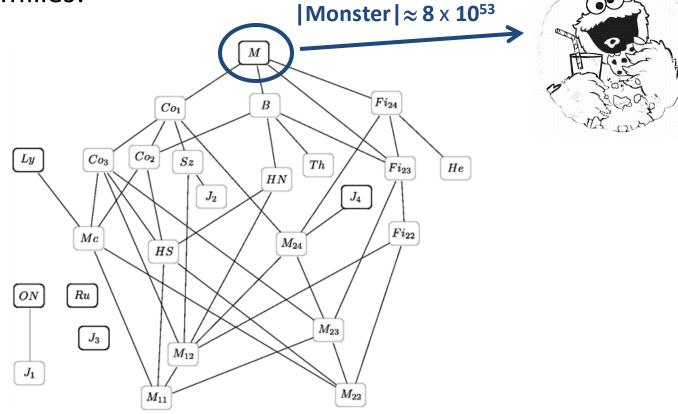
- Cyclic groups of prime order  $\rm C_p$ 

e.g.: 
$$C_p = \mathbb{Z}_p = \langle e^{2\pi i/p} \rangle$$

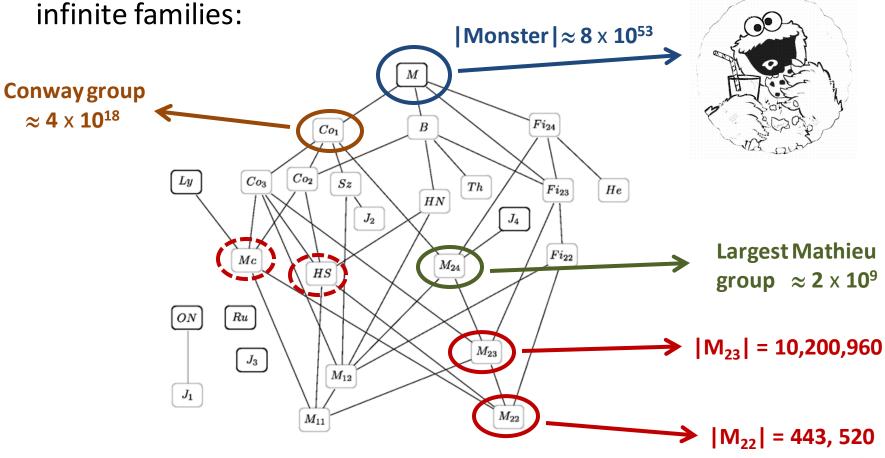
There are also 26 so called **sporadic groups** that do not come in infinite families:



There are also 26 so called **sporadic groups** that do not come in infinite families:



There are also 26 so called sporadic groups that do not come in



## Modular Forms

Modular function of weight k

$$f\left(\frac{a\tau+b}{c\tau+d}\right) = (c\tau+d)^k f(\tau), \qquad \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in SL(2,\mathbb{Z})$$

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Jacobi form of weight k and index m

$$f\left(\frac{a\tau+b}{c\tau+d},\frac{z}{c\tau+d}\right) = (c\tau+d)^{k} e^{\frac{2\pi i m c z^{2}}{c\tau+d}} f(\tau,z)$$
$$f(\tau,z+\lambda\tau+\mu) = e^{-2\pi i m (\lambda^{2}\tau+\lambda z)} f(\tau,z), \qquad \lambda,\mu \in \mathbb{Z}$$

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$$f(\tau,z+\lambda\tau+\mu) = e^{-2\pi i m (\lambda^{2}\tau+\lambda z)} f(\tau,z), \qquad \lambda,\mu \in \mathbb{Z}$$

Can Fourier expand

$$f(\tau, z) = f(q = \exp[2\pi i \tau], y = \exp[2\pi i z]) = \sum_{n \ge 0} \sum_{r^2 \le 4mn} c(n, r) q^n y^r$$

- The irreducible representations of the Monster group have dimensions 1, 196 883, 21 296 876, ...
- The J-function, that appears in many places in string theory, enjoys the expansion

$$J(q) = \frac{1}{q} + 196884q + 21493760q^2 + \dots$$

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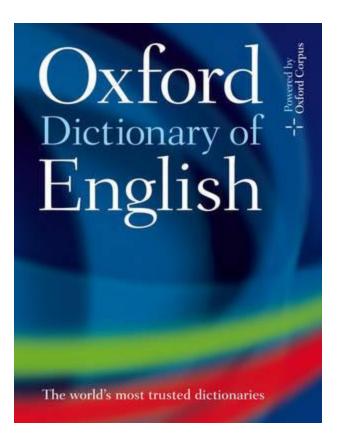
$$\boxed{1 + 196883} \qquad \boxed{1 + 196883} + 21296876$$

• as observed by John McKay in the late 70's

# moon SHīn/ 10

*noun informal* noun: **moonshine** 

1. foolish talk or ideas.



# moon SHīn/ 10

*noun informal* noun: **moonshine** 

1. foolish talk or ideas.

2. NORTH AMERICAN illicitly distilled or smuggled liquor.



This surprising connection can be explained by string theory: Frenkel, Lepowsky, Meurman 1988

 The (left-moving) bosonic string compactified on a Z<sub>2</sub> orbifold of ℝ<sup>24</sup>/Λ with Λ the Leech lattice has as its 1-loop partition function the J(q)-function

$$Z(q) = \operatorname{Tr}_{\mathrm{H}} q^{L_0 - \frac{c}{24}} = J(q) = \frac{1}{q} + 196884q + 21493760q^2 + \dots$$

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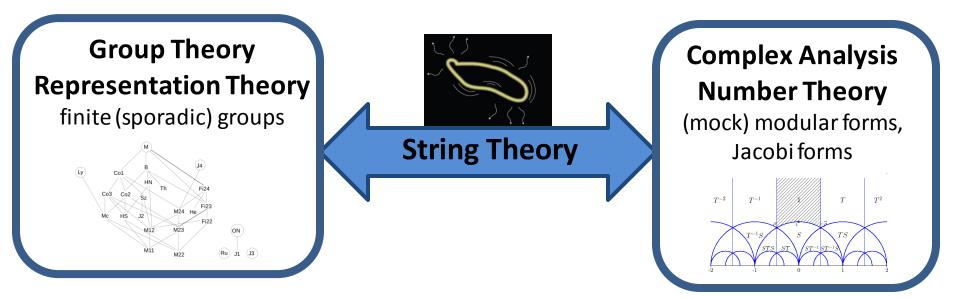
$$Z(q) = \operatorname{Tr}_{H} q^{L_{0} - \frac{c}{24}} = J(q) = \frac{1}{q} + 196884q + 21493760q^{2} + \dots$$

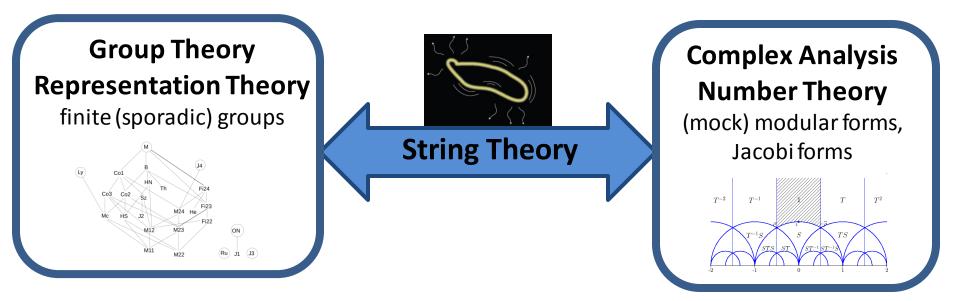
• The symmetry group of the compactification space  $\mathbb{R}^{24}/\Lambda/\mathbb{Z}_2$  is the Monster group.

Since we have a Virasoro algebra we can expand the J(q)-function in terms of Virasoro characters (traces of Verma modules)

$$\operatorname{ch}_{h=0}(q) = \frac{q^{-c/24}}{\prod_{n=2}^{\infty}(1-q^n)}, \qquad \operatorname{ch}_h(q) = \frac{q^{h-c/24}}{\prod_{n=1}^{\infty}(1-q^n)}$$

$$J(q) = \frac{1}{q} + 196884 q + 21493760 q^{2} + \dots$$
  
= 1 ch<sub>0</sub>(q) + 196883 ch<sub>2</sub>(q) + 21296876 ch<sub>3</sub>(q) + \dots





#### Very interesting for mathematicians!

Compactification of the bosonic string:

- $\implies$  we have a tachyon (instability)
- $\implies$  spacetime theory has no fermions

Additionally

• Only two spacetime dimensions are non-compact

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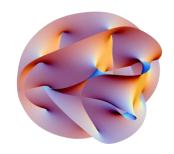
Not so interesting for physicists!

# Outline

Introduction to moonshine



 Mathieu Moonshine and string compactifications



• New moonshine phenomena



 In 2010 Eguchi, Ooguri and Tachikawa discovered a new moonshine phenomenon that connects K3 to the largest Mathieu group M<sub>24</sub>

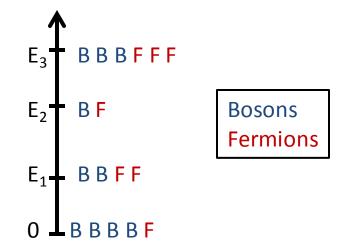
Eguchi, Ooguri, Tachikawa 1004.0956

• They considered a N=(4,4) SCFT with K3 target and calculate an index that is called elliptic genus

• The Witten index

$$Z_{\text{Witten}} = \text{Tr}(-1)^F$$

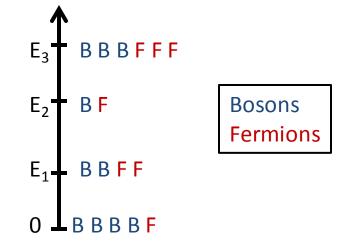
 $= n_B - n_F$ 

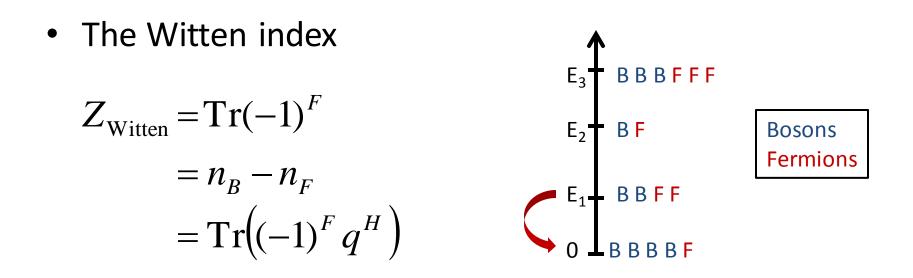


• The Witten index

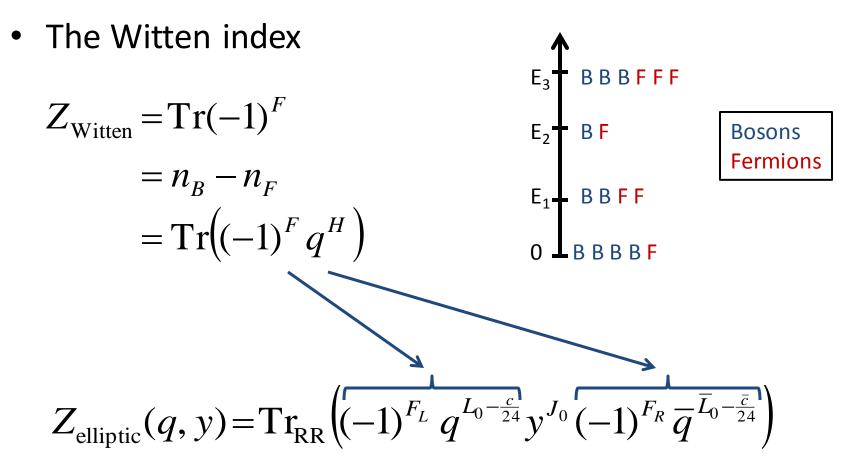
$$Z_{\text{Witten}} = \text{Tr}(-1)^F$$

$$= n_B - n_F$$
$$= \operatorname{Tr}\left((-1)^F q^H\right)$$





• An index is invariant under deformations of the theory, e.g. masses go to zero



No dependence on  $\overline{q}$  !

Chemical potential for U(1) in left-moving N=2 theory

$$Z_{\text{elliptic}}(q, y) = \operatorname{Tr}_{RR}\left((-1)^{F_L} q^{L_0 - \frac{c}{24}} y^{J_0} (-1)^{F_R} \overline{q}^{\overline{L}_0 - \frac{\overline{c}}{24}}\right)$$

Witten index: No dependence on  $\overline{q}$ 

$$Z_{\text{elliptic}}^{\text{K3}}(q, y) = 8 \left( \frac{\theta_2(q, y)^2}{\theta_2(q, 1)^2} + \frac{\theta_3(q, y)^2}{\theta_3(q, 1)^2} + \frac{\theta_4(q, y)^2}{\theta_4(q, 1)^2} \right)$$

T. Eguchi, H. Ooguri, A. Taormina, S.-K. Yang Nucl. Phys. B 315, 193 (1989)

#### We have N=(4,4) world sheet supersymmetry

 $\implies$  expand in N=4 Virasoro characters

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N=4 Virasoro characters are defined by the trace over the highest weight state and all its descendants

$$\operatorname{ch}_{h,l}(q, y) = \operatorname{Tr}\left((-1)^{F_L} q^{L_0 - \frac{c}{24}} y^{J_0}\right)$$

For the case h=c/24 there are short BPS multiplets

$$Z_{\text{elliptic}}^{\text{K3}}(q, y) = 8 \left( \frac{\theta_2(q, y)^2}{\theta_2(q, 1)^2} + \frac{\theta_3(q, y)^2}{\theta_3(q, 1)^2} + \frac{\theta_4(q, y)^2}{\theta_4(q, 1)^2} \right)$$

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T. Eguchi, K. Hikami 0904.0911

$$Z_{\text{elliptic}}^{\text{K3}} = 24 \operatorname{ch}_{h=\frac{1}{4},l=0}^{\text{short}} - 2 \operatorname{ch}_{h=\frac{1}{4},l=\frac{1}{2}}^{\log} + \sum_{n=1}^{\infty} A_n \operatorname{ch}_{h=\frac{1}{4}+n,l=\frac{1}{2}}^{\log}$$

 $A_n = \{90, 462, 1440, \dots\}$ 

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$$23 + 1$$
T. Eguchi, H. Ooguri, Y. Tachikawa 1004.0956

$$A_n = \{45 + \overline{45}, 231 + \overline{231}, 770 + \overline{770}, \dots\}$$

Irreps of M<sub>24</sub>

**Dimensions of** 

Does this imply a connection between  $M_{24}$  and K3?

- The geometric symmetries of K3 are contained in  $M_{23}{\subset}\,M_{24}$ 

Mukai, Kondo 1988, 1998

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Gaberdiel, Hohenegger, Volpato 1106.4315

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Gaberdiel, Hohenegger, Volpato 1106.4315

• However, all the  $A_n$  are sums of dimensions of irreps of  $M_{24}$  with positive coefficients

Gannon 1211.5531

## Mathieu Moonshine

- K3 has played a central role in string compactifications and string dualities
- What are implications we can derive from Mathieu moonshine for string compactifications?
- Has the elliptic genus of K3 already appeared in the string theory literature?

## Mathieu Moonshine

- K3 has played a central role in string compactifications and string dualities
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- Has the elliptic genus of K3 already appeared in the string theory literature?

- Consider the heterotic E<sub>8</sub> x E<sub>8</sub> string theory compactified on K3 x T<sup>2</sup>
- We need to embed 24 instantons into E<sub>8</sub> x E<sub>8</sub> → (12+n,12-n) n = 0,1,...,12 to satisfy the Bianchi identity for H<sub>3</sub>

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- The resulting four dimensional theories preserves N=2 spacetime supersymmetry
- The 1-loop corrections to the prepotential are related to the new supersymmetric index Z<sub>new</sub>

Dixon, Kaplunovsky, Louis, de Wit, Lüst, Stieberger, Antoniadis, Narain, Taylor, Gava, Kiritsis, Kounnas, Harvey, Moore, ....

$$h(S,T,U) = h^{\text{tree}} + h^{1-loop} + O(e^{-2\pi i S})$$

For K3 x T<sup>2</sup> compactifications we have for the standard embedding that preserves N=(4,4)Affine E<sub>8</sub> SO(12) characters  $Z_{\text{new}}(q;T,U) = \frac{i}{2} \frac{\Theta_{\Gamma_{2,2}}(T,U)}{n(q)^4} \frac{E_4(q)}{n(q)^8} \left| \left( \frac{\theta_2(q)}{n(q)} \right)^6 Z_{\text{elliptic}}^{\text{K3}}(q,-1) \right|$  $+\left(\frac{\theta_3(q)}{n(q)}\right)^6 q^{\frac{1}{4}} Z_{\text{elliptic}}^{\text{K3}}(q,-q^{\frac{1}{2}}) + \left(\frac{\theta_4(q)}{n(q)}\right)^6 q^{\frac{1}{4}} Z_{\text{elliptic}}^{\text{K3}}(q,q^{\frac{1}{2}})$ SO(12) characters Harvey, Moore hep-th/9510182

*T* is the complexified Kähler modulus, *U* the complex structure modulus of the T<sup>2</sup>

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So in particular the [...] part has an "SO(12)xM<sub>24</sub>"-expansion: exactly the same M<sub>24</sub> as in Mathieu Moonshine due to N=(4,4)

For K3 x T<sup>2</sup> compactifications we have for the standard embedding that preserves N=(4,4) Affine E<sub>8</sub>  $\bigvee$  $Z_{\text{new}}(q;T,U) = \frac{i}{2} \frac{\Theta_{\Gamma_{2,2}}(T,U)}{\eta(q)^4} \frac{E_4(q)}{\eta(q)^8} \frac{E_6(q)}{\eta(q)^{12}} = \frac{i}{2} \frac{\Theta_{\Gamma_{2,2}}(T,U)E_4(q)E_6(q)}{\eta(q)^{24}}$ 

Harvey, Moore hep-th/9510182

So in particular the  $E_6(q)$  has an "SO(12)xM<sub>24</sub>"-expansion

For K3 x T<sup>2</sup> compactifications we have for the standard embedding that preserves N=(4,4)

Take away message:

 $Z_{\text{new}}$  depends on *T* and *U* and is connected to  $Z_{\text{elliptic}}$  and therefore to  $M_{24}$ 

*T* is the complexified Kähler modulus, *U* the complex structure modulus of the T<sup>2</sup>

The 1-loop correction to the prepotential is roughly determined by

$$\Delta(T,U) = \int \frac{d^2 \tau}{\tau_2} Z_{\text{new}}(q = e^{2\pi i \tau}; T, U) (Q^2 - \frac{1}{8\pi \tau_2})$$

and knows about  $M_{24}$  since  $Z_{new}$  does

M. Cheng, X. Dong, J. Duncan, J. Harvey, S. Kachru, TW 1306.4981

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The modular invariance of  $\tau_2 Z_{\text{new}}(Q^2 - \frac{1}{8\pi_2})$  actually tells us that there is a unique solution. So for all instanton embeddings (12+n,12-n) the answer is the same.

Kiritsis, Kounnas, Petropoulos, Rizos hep-th/9608034 Henningson, Moore hep-th/9608145

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We have to solve the following second order differential equation Harvey, Moore hep-th/9510182

$$-\operatorname{Re}\left(\partial_{T}\partial_{U}h^{1-loop} + \frac{1}{T_{1}U_{1}}(1 - T_{1}\partial_{T} - U_{1}\partial_{U})h^{1-loop}\right) - \frac{1}{\pi}\operatorname{Re}(\log[J(iT) - J(iU)])$$

$$= \frac{1}{2\pi}\int \frac{d^{2}\tau}{\tau_{2}} \left(-iZ_{new}(q;T,U) \cdot (Q_{E_{8}}^{2} - \frac{1}{8\pi_{2}}) - b(E_{8})\right) + \frac{b(E_{8})}{2\pi}(\log[2T_{1}U_{1}] + 4\operatorname{Re}(\log[\eta(iT)\eta(iU)]))$$

The solution is given by

$$h^{1-loop} = -\frac{1}{3}U^3 + C + \sum_{k,l} c(kl) Li_3(e^{2\pi i(kT+lU)})$$

where the polylogarithm is given by  $Li_3(x) = \sum_{n=1}^{\infty} \frac{x^n}{n^3}$ 

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and the expansion coefficients are the same as in our index (they go along for the ride when integrating)

$$Z_{\text{new}}(q;T,U) = \frac{i}{2} \frac{\Theta_{\Gamma_{2,2}}(T,U)E_4(q)E_6(q)}{\eta(q)^{24}} = \frac{i}{2} \Theta_{\Gamma_{2,2}}(T,U) \left(\sum_{m \ge -1} c(m)q^m\right)$$
  
Dimensions of M<sub>24</sub> 49

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Dimensions of M<sub>24</sub> (appearing in a spacetime quantity)

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Dimensions of M<sub>24</sub> 50

#### **String duality**

Heterotic string on K3 x T<sup>2</sup> with instanton embedding (12+n,12-n)

Type IIA string theory on elliptic fibrations over F<sub>n</sub> (Hirzebruch surface)

#### **String duality**

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Fn

dilaton S  $\leftarrow$  Size of base S<sup>2</sup>

Type IIA string theory on elliptic fibrations over  $F_{n:}$ 

- Prepotential receives instanton corrections
- These are determined by the Gromow-Witten invariants  $\approx$  curve counting (S<sup>2</sup>, T<sup>2</sup>, ...)

Type IIA string theory on elliptic fibrations over F<sub>n</sub>:

- Prepotential receives instanton corrections
- These are determined by the Gromow-Witten invariants  $\approx$  curve counting (S<sup>2</sup>, T<sup>2</sup>, ...) M. Cheng, X. Dong, J. Duncan, J. Harvey, S. Kachru, TW 1306.4981  $h(S,T,U) = -STU - \frac{1}{3}U^3 + C + \sum_{k,l} c(kl)Li_3(e^{2\pi i(kT+lU)}) + O(e^{-2\pi iS})$ M. Alim. F. Scheidegger 1205 1784 M. Alim, E. Scheidegger 1205.1784 A. Klemm, J. Manschot, T. Wotschke 1205.1795 Dimensions of  $M_{24}$ (appearing in a spacetime quantity) 54

#### **String duality**

Type IIA string theory on elliptic fibrations over F<sub>n</sub> (Hirzebruch surface) Type IIB string theory

on mirror manifold

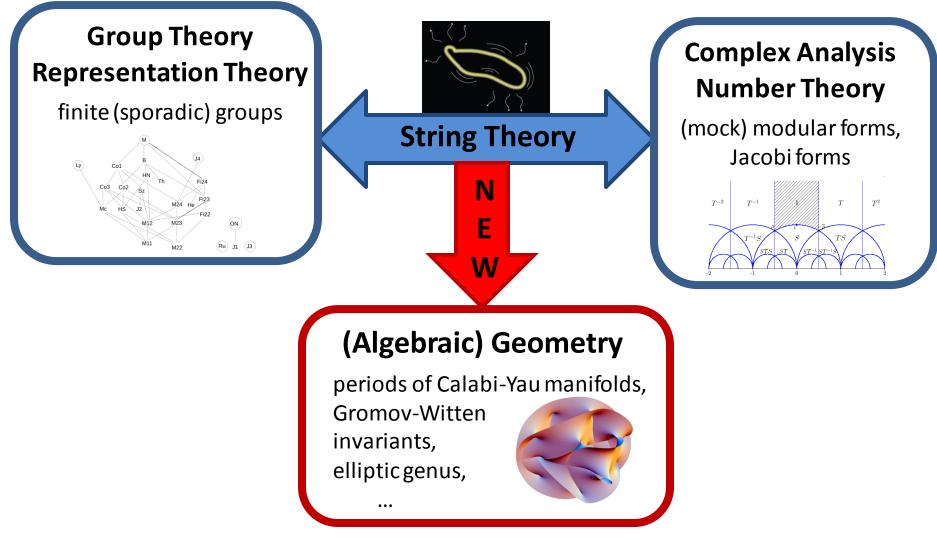
#### **String duality**

Type IIA string theory on elliptic fibrations over F<sub>n</sub> (Hirzebruch surface) Type IIB string theory on mirror manifold

 $CY_3$  manifold  $X_n \leftarrow P CY_3$  manifold  $Y_n$ 

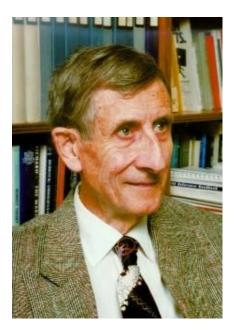
Gromov-Witten invariants Periods of the holomorphic 3-form  $\Omega$ 

### New math connections



"I have a sneaking hope, a hope unsupported by any facts or any evidence, that sometime in the twenty-first century physicists will stumble upon the Monster group, built in some unsuspected way into the structure of the Universe."

– Freeman Dyson (1983)



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The (bulk) moduli dependent 1-loop correction to the gauge kinetic function arises only from N=2 subsectors! Dixon, Louis, Kaplunovsky Nuclear Physics B 355 (1991)

Example  $T^{6}/\mathbb{Z}_{6-11} = T^{2} \times T^{2} \times T^{2}/\mathbb{Z}_{6-11}$ :

$$\mathbb{Z}_{6-II} = \langle g \rangle, \quad g: (z_1, z_2, z_3) \to (e^{\pi i/3} z_1, e^{2\pi i/3} z_2, -z_3)$$

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has two N=2 subsector

$$\mathbb{Z}_{3} = \{1, g^{2}, g^{4}\}, \qquad g^{2} : (z_{1}, z_{2}, z_{3}) \to (e^{2\pi i/3} z_{1}, e^{4\pi i/3} z_{2}, z_{3})$$
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For which the internal space is  $T^4/\mathbb{Z}_3 \ge T^2$  or  $T^4/\mathbb{Z}_2 \ge T^2$ respectively and therefore an orbifold limit of  $T^2 \ge K3$ .

N=2 sectors lead to 1-loop corrections  

$$f_{\alpha}^{1-\text{loop}}(T,U) = \sum_{i=1,2,3} \frac{|G_i'|}{|G|} \left[ -\frac{1}{2} \partial_{T_i} \partial_{U_i} h_i^{1-\text{loop}}(T_i, U_i) \right]$$

$$h_i^{\text{N=1 gauge}} = -\frac{1}{8\pi^2} \log[J(iT_i) - J(iU_i) - \frac{b_{\alpha,i}^{N=2}}{4\pi^2} (\log[\eta(iT_i)\eta(iU_i)]) \right]$$

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N=1 gauge  
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where the prepotential was calculated above

$$h^{1-loop}(T,U) = -\frac{1}{3}U^3 + C + \sum_{k,l} c(kl)Li_3(e^{2\pi i(kT+lU)})$$
  
Dimensions of M<sub>24</sub>

Group $\mathbb{Z}_N$	Generator $\frac{1}{N}(\varphi_1,\varphi_2,\varphi_3)$	$\mathcal{N} = 2 \mod \mathbf{u}$	
$\mathbb{Z}_3$	$\frac{1}{3}(1,1,1)$	-	
$\mathbb{Z}_4$	$\frac{1}{4}(1,1,2)$	$T_{3}, U_{3}$	
$\mathbb{Z}_{6-I}$	$\frac{1}{6}(1,1,4)$	$T_3$	No N=2 sectors
$\mathbb{Z}_{6-II}$	$rac{1}{6}(1,2,3)$	$T_2, T_3, U_3$	
$\mathbb{Z}_7$	$\frac{1}{7}(1,2,4)$	-	
$\mathbb{Z}_{8-I}$	$\frac{1}{8}(1,2,5)$	$T_2$	
$\mathbb{Z}_{8-II}$	$\frac{1}{8}(1,3,4)$	$T_3, U_3$	
$\mathbb{Z}_{12-I}$	$\frac{1}{12}(1,4,7)$	$T_2$	
$\mathbb{Z}_{12-II}$	$\frac{1}{12}(1,5,6)$	$T_3, U_3$	

$\mathbb{Z}_N \times \mathbb{Z}_M$	1 <sup>st</sup> generator $\frac{1}{N}(\varphi_1,\varphi_2,\varphi_3)$	$2^{\mathrm{nd}}$ generator $\frac{1}{M}(\hat{\varphi}_1, \hat{\varphi}_2, \hat{\varphi}_3)$	$\mathcal{N} = 2 \mod \mathbf{u}$
$\mathbb{Z}_2 \times \mathbb{Z}_2$	$rac{1}{2}(1,0,1)$	$\frac{1}{2}(0,1,1)$	$T_1, U_1, T_2, U_2, T_3, U_3$
$\mathbb{Z}_2 \times \mathbb{Z}_4$	$rac{1}{2}(1,0,1)$	$rac{1}{4}(0,1,3)$	$T_1, U_1, T_2, T_3$
$\mathbb{Z}_2 \times \mathbb{Z}_6$	$rac{1}{2}(1,0,1)$	$rac{1}{6}(0,1,5)$	$T_1, U_1, T_2, T_3$
$\mathbb{Z}_2 \times \mathbb{Z}_6'$	$rac{1}{2}(1,0,1)$	$rac{1}{6}(1,1,4)$	$T_1, T_2, T_3$
$\mathbb{Z}_3 \times \mathbb{Z}_3$	$\frac{1}{3}(1,0,2)$	$\frac{1}{3}(0,1,2)$	$T_1, T_2, T_3$
$\mathbb{Z}_3 \times \mathbb{Z}_6$	$\frac{1}{3}(1,0,2)$	$rac{1}{6}(0,1,5)$	$T_1, T_2, T_3$
$\mathbb{Z}_4 \times \mathbb{Z}_4$	$rac{1}{4}(1,0,3)$	$rac{1}{4}(0,1,3)$	$T_1, T_2, T_3$
$\mathbb{Z}_6 \times \mathbb{Z}_6$	$rac{1}{6}(1,0,5)$	$rac{1}{6}(0,1,5)$	$T_1, T_2, T_3$

Four dimensional N=1 models obtained from orbifold compactifications of the heterotic  $E_8 \times E_8$  string theory receive universal 1-loop corrections to their gauge kinetic functions that are related to M<sub>24</sub>

TW 1402.2973

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For all T<sup>6</sup>/ $\mathbb{Z}_N$ , N $\neq$ 3,7, and all T<sup>6</sup>/ $\mathbb{Z}_N$ x $\mathbb{Z}_M$ 

$$f(S,T,U) \approx S + \partial_T \partial_U \sum_{k,l} c(kl) Li_3(e^{2\pi i(kT+lU)}) + \dots + O(e^{-2\pi iS})$$
  
Dimensions of M<sub>24</sub>

N. Paquette, TW work in progress:

• The holomorphic 3-form  $\Omega$  plays a role in flux compactifications

Gukov, Vafa, Witten hep-th/9906070 Giddings, Kachru, Polchinski hep-th/0105097

$$W = \int H \wedge \Omega + \dots$$

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• The Yukawa couplings in heterotic models are given by the third derivative of  $\Omega$  with respect to the moduli Hosono, Klemm, Theisen, Yau hep-th/9308122

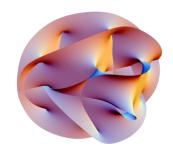
$$Y_{IJK} \approx \partial_I \partial_J \partial_K h(S, T, U)$$

# Outline

Introduction to moonshine



 Mathieu Moonshine and string compactifications





Consider eight (left-moving) bosons and fermions compactified on the orbifold  $T^8/\mathbb{Z}_2 = \mathbb{R}^8/\Lambda_{E8}/\mathbb{Z}_2$ 

$$\mathbb{Z}_2: (X^I, \psi^I) \to -(X^I, \psi^I)$$

The partition function in the NS sector is

Frenkel, Lepowsky, Meurman 1985

$$Z(q) = \frac{1}{\sqrt{q}} + 276\sqrt{q} + 2048q + 11202q^{\frac{3}{2}} + \dots$$
  
Dimensions of representations  
of Conway group



- The Conway symmetry is not manifest in this description
- The theory is equivalent to a theory of 24 chiral fermions orbifolded by a  $\mathbb{Z}_2$  symmetry:  $\psi^I \rightarrow -\psi^I$

J. Duncan math/0502267

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- The theory is equivalent to a theory of 24 chiral fermions orbifolded by a  $\mathbb{Z}_2$  symmetry:  $\psi^I \rightarrow -\psi^I$ J. Duncan math/0502267
- One can construct an N=1 superalgebra that breaks the Spin(24) symmetry to Co<sub>0</sub>
- This explains this Conway moonshine

- The toroidal orbifold  $T^8/\mathbb{Z}_2 = \mathbb{R}^8/\Lambda_{E8}/\mathbb{Z}_2$  preserves N=4 worldsheet supersymmetry
- Calculate the partition function and expand in *N*=4 characters:

$$Z = \operatorname{Tr}\left((-1)^{F} q^{L_{0}-c/24} y^{J_{0}}\right) = 21 \operatorname{ch}_{h=\frac{1}{2},l=0}^{\operatorname{short}} + \operatorname{ch}_{h=\frac{1}{2},l=1}^{\operatorname{short}}$$

$$+560 \operatorname{ch}_{h=\frac{3}{2},l=\frac{1}{2}}^{\log} + 8470 \operatorname{ch}_{h=\frac{5}{2},l=\frac{1}{2}}^{\log} + \dots$$

$$+ 210 \operatorname{ch}_{h=\frac{3}{2},l=1}^{\log} + 4444 \operatorname{ch}_{h=\frac{5}{2},l=1}^{\log} + \dots$$

• Two infinite series, coefficients unrelated to Conway

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#### All coefficients are dimensions of the Mathieu group M<sub>22</sub>!

- The toroidal orbifold  $T^8/\mathbb{Z}_2 = \mathbb{R}^8/\Lambda_{E8}/\mathbb{Z}_2$  preserves *N=4* worldsheet supersymmetry
- Let us generalize the above idea and choose chiral super-Virasoro algebras with N > 1 for 24 fermions

M. Cheng, X. Dong, J. Duncan, S. Harrison, S. Kachru, TW 1406.5502

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- Let us generalize the above idea and choose chiral super-Virasoro algebras with N > 1 for 24 fermions M. Cheng, X. Dong, J. Duncan, S. Harrison, S. Kachru, TW 1406.5502
- We find that N=2 (N=4) super-Virasoro algebras break the symmetry group to subgroups of Co<sub>0</sub> that fix a 2-plane (3-plane) in the 24 dimensional representation of Co<sub>0</sub>

- There are a variety of groups that do not act on a 2- (or 3-) plane in the 24 dimensional representation of Co<sub>0</sub>
- For example: Stabilizers of 2-plane Stabilizers of 3-plane  $M_{23}$   $M_{22}$  Higman-Sims  $U_4(3)$  McLaughlin  $U_6(2)$

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- $\implies$  Coefficient in expansion in N=2 (N=4) characters can be decomposed in irreps of these groups

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+ 770 ch<sup>long</sup><sub>$$h=\frac{3}{2},l=1$$</sub> + 13915 ch<sup>long</sup> <sub>$h=\frac{5}{2},l=1$</sub>  + ...

$$+231 \mathrm{ch}_{h=\frac{3}{2},l=2}^{\log}+5796 \mathrm{ch}_{h=\frac{5}{2},l=2}^{\log}+\dots$$

All coefficients are dimensions of  $M_{23}$ , HS, McL and  $U_6(2)$ !

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All coefficients are dimensions of  $M_{22}$  and  $U_4(3)$ !

• The twined partition functions

$$Z_{g} = \operatorname{Tr}\left(g\left(-1\right)^{F} q^{L_{0}-c/24} y^{J_{0}}\right)$$

have special properties only for  $M_{22}$  and  $M_{23}$ :

The Mathieu groups  $M_{22}$  and  $M_{23}$  satisfy the extra moonshine property that the mock modular forms in all twined partition functions have only poles at the cusp at  $\tau = i \infty$ 

(They can most likely be constructed via Rademacher sums.)

Theory of 24 chiral fermions orbifolded by a Z<sub>2</sub> symmetry:

$$\mathbb{Z}_2: \psi^I \to -\psi^I$$

gives rise to variety of moonshine phenomena:

Symmetry group: Spin(24)  $\supset$  Co<sub>0</sub>  $\supset$  M<sub>23</sub>  $\supset$  M<sub>22</sub>

Expansion in: N=0  $\longrightarrow$  N=1  $\longrightarrow$  N=2  $\longrightarrow$  N=4

# Conclusion

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#### **THANK YOU!**