

# Cosmological Reionization Simulations for LOFAR

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# People involved: The LOFAR-EoR Core

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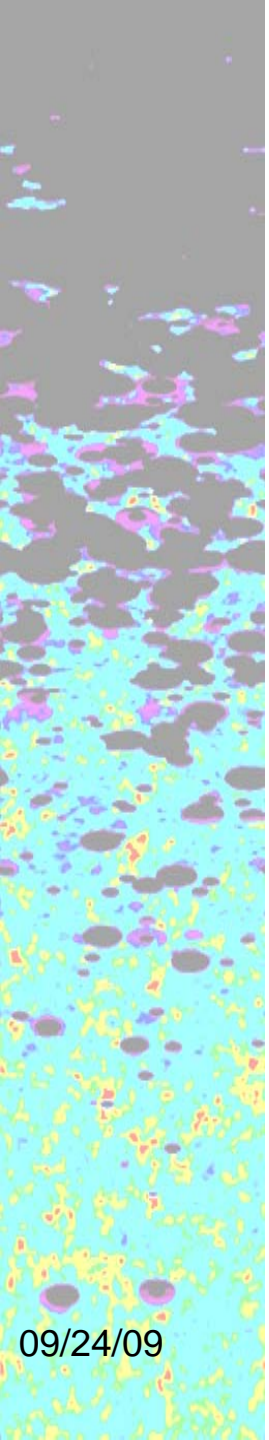
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# Outline

- Introduction
  - Glossary of the relevant physics and the astrophysics
- “The goal”
- Implementation & Results
- What comes next?

# Glossary

- Brief history of the Universe
- Observable: 21-cm spin-flip transition
- Brightness temperature
- Radiative transfer basics



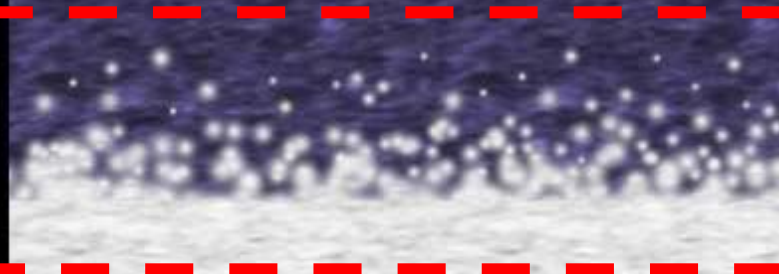
0.4 Myr  
( $z \sim 1100$ )



**COSMIC MICROWAVE  
BACKGROUND**

**DARK AGES**

400 Myr  
( $z \sim 10$ )



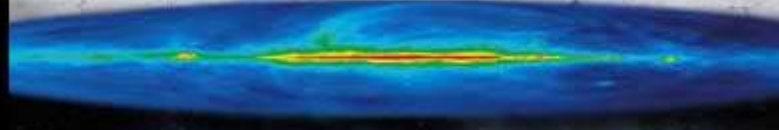
**EPOCH OF  
REIONIZATION**

2 Gyr  
( $z \sim 3$ )



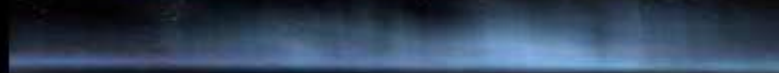
**EXTRAGALACTIC  
FOREGROUNDS**

13.7 Gyr  
( $z \sim 0$ )



**GALACTIC  
FOREGROUNDS**

time since  
Big Bang



**IONOSPHERE**



**LOFAR TELESCOPE**



**BLUEGENE STELLA**

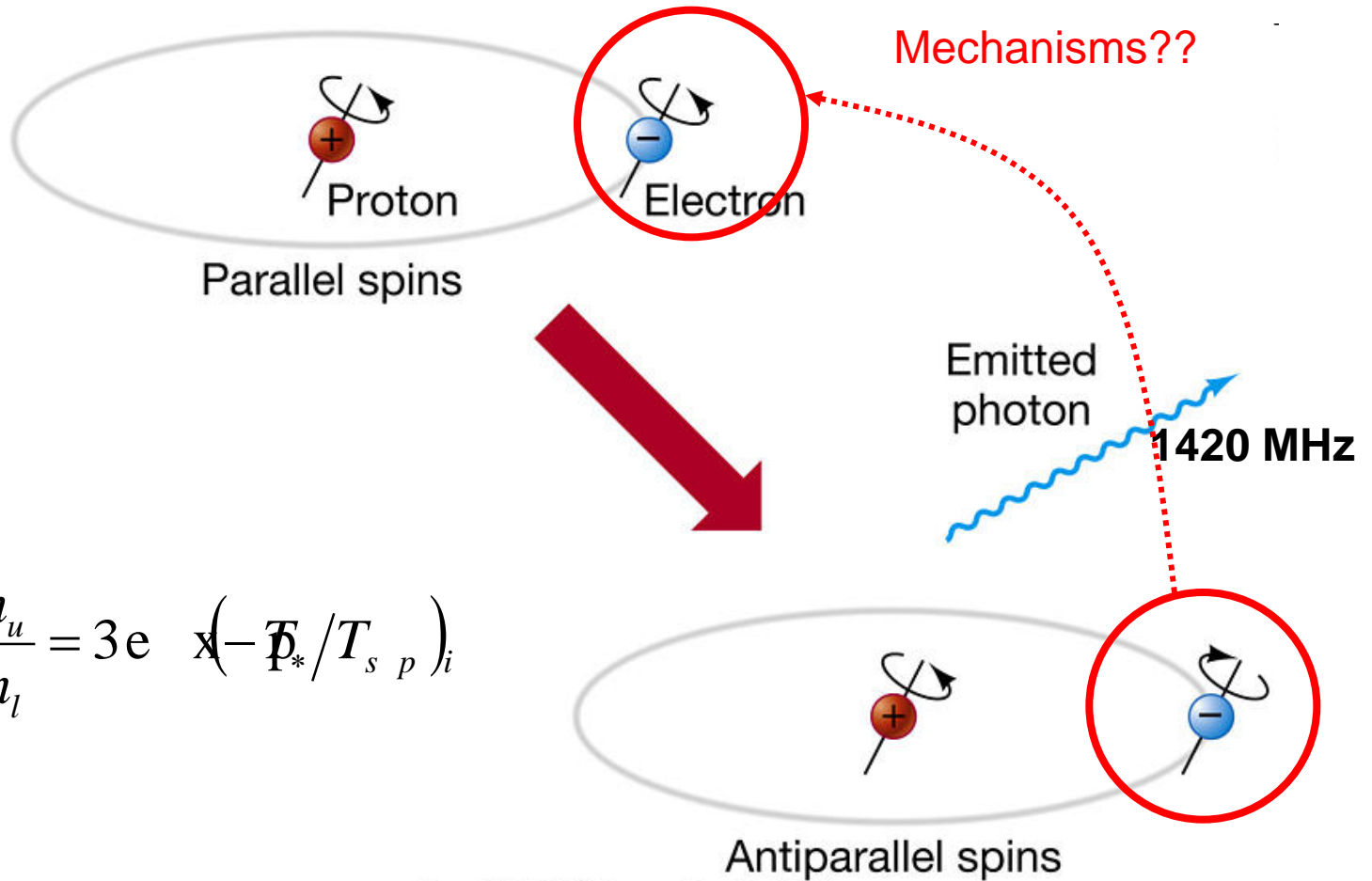
*Courtesy:*

**Vibor Jelic**

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100001101101111001100101  
110010110011010100101000
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# 21-cm Physics



$$\frac{n_u}{n_l} = 3e^{-\frac{h\nu}{kT_{sp}}} \left(1 + \frac{A_{ul}}{T_{sp} \lambda^3} \right)^{-1}$$

# 21cm Physics: Spin Temp

$$T_s = \frac{T_{\text{CMB}} + y_k T_k + y_\alpha T_\alpha}{1 + y_k + y_\alpha}$$

Field: 1958

$K \rightarrow$  kinetic and  $\alpha \rightarrow$  light

We assume  $T_k = T_\alpha$

If  $y_k$  or  $y_\alpha \gg 1$  then  $T_s = T_k$

If  $y_k$  and  $y_\alpha \ll 1$  then  $T_s = T_{\text{CMB}}$

**Remember!!**

**It is very well possible;**

**$T_k < T_{\text{CMB}}$**

# $\delta T_b$ : Brightness temperature

$$\delta T_b = (20 \text{ mK}) (1 + \delta) \left( \frac{X_{HI}}{h} \right) \left( 1 - \frac{T_{CMB}}{T_{spin}} \right) \left( \frac{\Omega_b h^2}{0.0223} \right) \left[ \left( \frac{1+z}{10} \right) \left( \frac{0.24}{\Omega_m} \right) \right]^{1/2}$$

**Cosmology**      **Astrophysics**

**Radiative Transfer**



# The Goal

Develop an efficient *radiative transfer* scheme to explore the various plausible reionization scenarios and predict the nature of the redshifted **21-cm signal** from the **Epoch of Reionization** .....specially for the **LOFAR case**.



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Where do we start?

Study the effect of various types  
of sources on ionization and  
temperature of their surrounding  
IGM

# Problems studied

- **(Mini)qsos:**

**Mass range of BH:**  $1000 - 10^9 M_{\odot}$

**SED:**  $E^{-\alpha}$

$$E_{\text{low}} < E < 10^4 \text{ eV}$$

$$E_{\text{low}} = 10.4 \text{ or } 200 \text{ eV}$$

**Luminosity :** Eddington X Efficiency

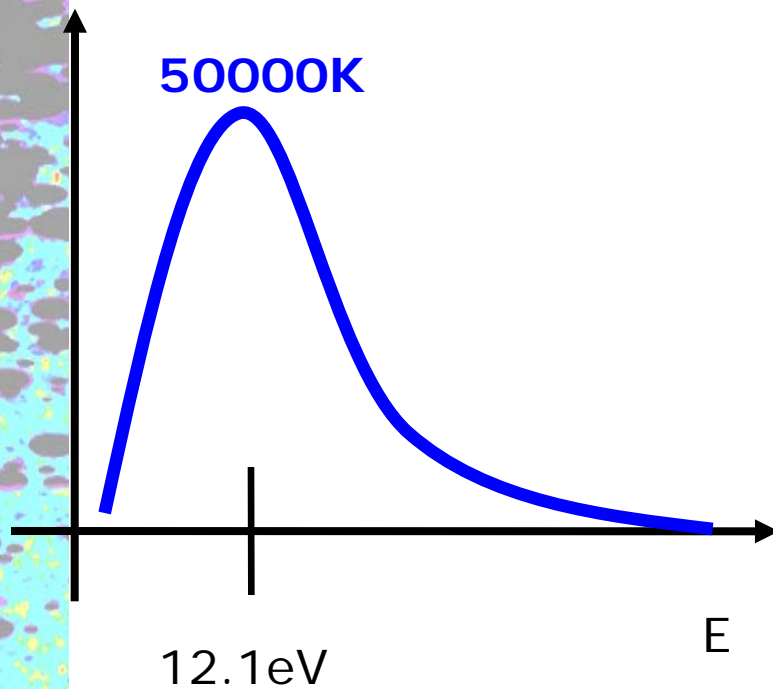
- **Stars:**

**Mass range:**  $10 - 10^4 M_{\odot}$

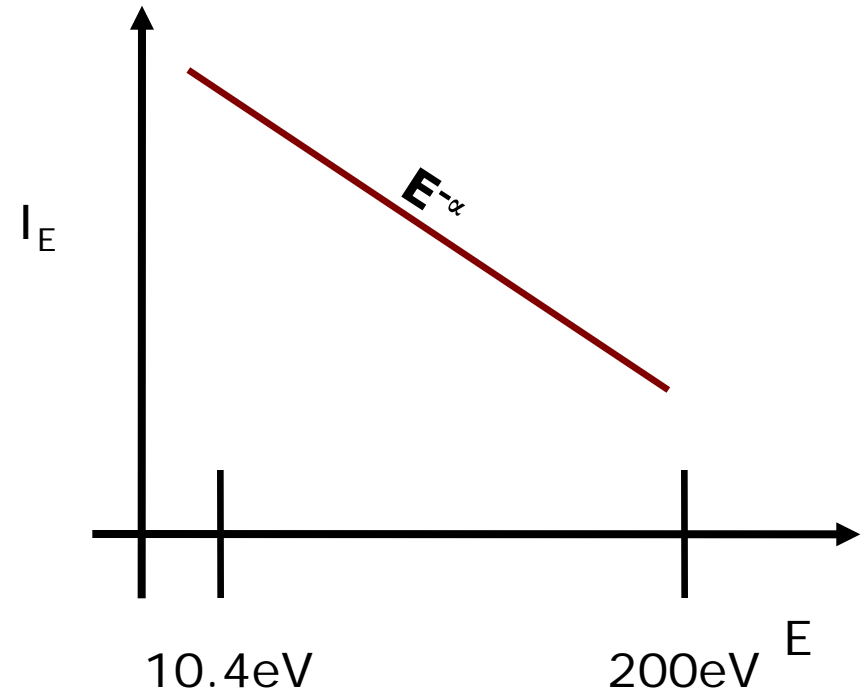
**SED:** Blackbody at 50000 K.

**Luminosity :** Schaerer (2004)

# Types of “first sources”



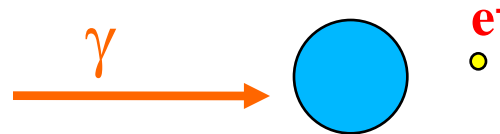
Stars



Power-law source

# Why X-ray heats

**UV photons**



**Large cross section but ejected electron has low energy**

**X-ray photons**



**Low cross section but ejected electron has high energy**

# RT-Equations

Thomas & Zaroubi, 2008, MNRAS

- Rate equations:

$$\frac{d n_i}{d t} = \Gamma(T) - \alpha(T)$$

$$\left\{ \begin{array}{l} n_i \rightarrow \text{HII, HeII, HeIII} \\ \Gamma \rightarrow \text{Ionizations} \\ \alpha \rightarrow \text{Recombinations} \end{array} \right.$$

- Temperature evolution:

$$\frac{d T}{d t} = H_P T + H_{Ec} - C_{cm} - C_{Hm} - C_{upe} - C_{bl} - C_{fb}$$

Details : Fukugita & Kawasaki (1994)



## Ionization

$$\frac{d}{dt} \left[ \frac{n(\text{H II})}{n_{\text{H}}} \right] = \frac{R_{1c} n(\text{H I})}{n_{\text{H}}} - \frac{\alpha_{2, \text{H II}} n_e^2}{n_{\text{H}}}$$

$$R_{2c} = \gamma_{2c} + \beta_{2, \text{H I}} n_e + \int_{\varepsilon_{2, \text{H I}}} c \sigma_{2f, \text{H I}} n_{\gamma}(\varepsilon_{\gamma}) d\varepsilon_{\gamma}$$

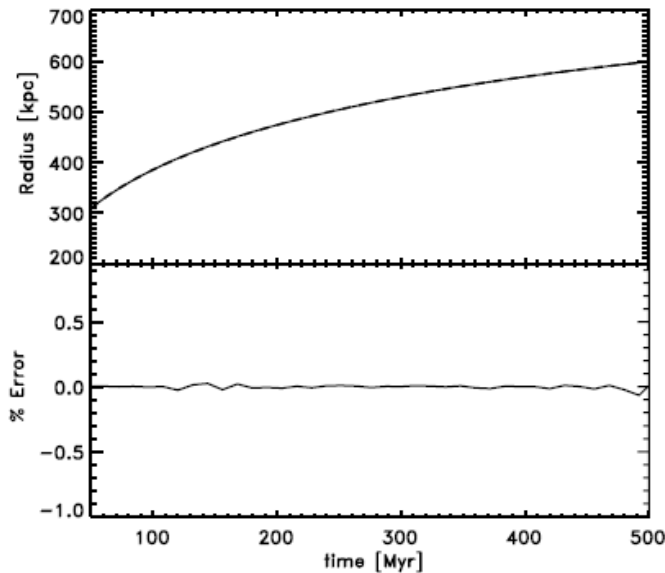
$$\begin{aligned} \frac{d}{dt} \left[ \frac{n(\text{He II})}{n_{\text{He}}} \right] &= \frac{n(\text{He I})}{n_{\text{He}}} \int_{\varepsilon_{\text{He I}}} d\varepsilon_{\gamma} \sigma_{\text{bf, He I}} c n_{\gamma}(\varepsilon_{\gamma}) \\ &+ \beta_{\text{He I}} n_e \frac{n(\text{He I})}{n_{\text{He}}} - \beta_{\text{He II}} n_e \frac{n(\text{He II})}{n_{\text{He}}} \\ &- \alpha_{\text{He II}} n_e \frac{n(\text{He II})}{n_{\text{He}}} + \alpha_{\text{He III}} n_e \frac{n(\text{He III})}{n_{\text{He}}} \\ &- \xi_{\text{He II}} n_e \frac{n(\text{He II})}{n_{\text{He}}}, \end{aligned}$$

$$\begin{aligned} \frac{d}{dt} \left[ \frac{n(\text{He III})}{n_{\text{He}}} \right] &= \frac{n(\text{He II})}{n_{\text{He}}} \int_{\varepsilon_{\text{He II}}} d\varepsilon_{\gamma} \sigma_{\text{bf, He II}} c n_{\gamma}(\varepsilon_{\gamma}) \\ &+ \beta_{\text{He II}} n_e \frac{n(\text{He II})}{n_{\text{He}}} - \alpha_{\text{He III}} n_e \frac{n(\text{He III})}{n_{\text{He}}}, \end{aligned}$$

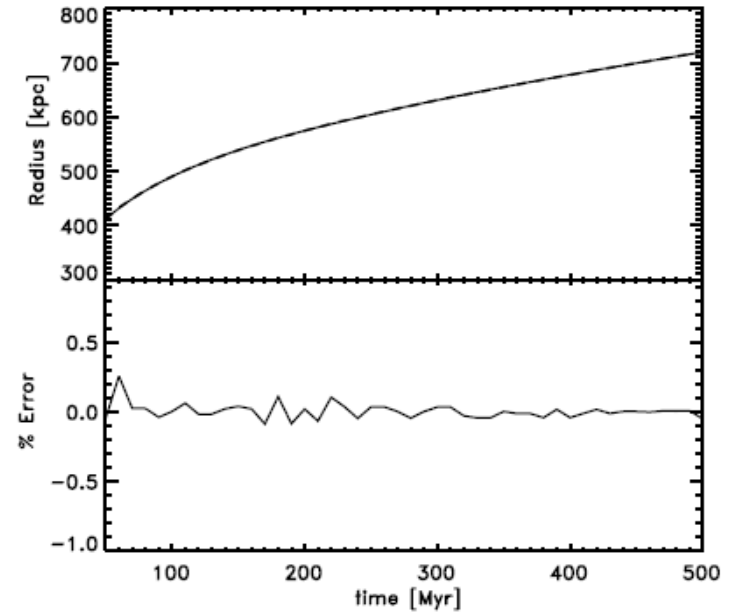
## Temperature

$$\begin{aligned} \frac{3}{2} \frac{d}{dt} \left( \frac{kT_e n_B}{\mu} \right) &= \sum_{i=\text{H I}, \text{He I}, \text{He II}} n(i) c \int (\varepsilon_{\gamma} - \varepsilon_i) n_{\gamma} \sigma_{\text{bf}, i} d\varepsilon_{\gamma} \\ &- \sum_{i=\text{H I}, \text{He I}, \text{He II}} \xi_i n_e n(i) \\ &- \sum_{i=\text{H II}, \text{He II}, \text{He III}} \eta_i n_e n(i) \\ &- \omega_{\text{He II}} n_e n(\text{He III}) \\ &- \sum_{i=\text{H I}, \text{He I}, \text{He II}} \psi_i n_e n(i) \\ &- \lambda_c \\ &- \theta_{\text{ff}} [n(\text{H II}) + n(\text{He II}) + 4n(\text{He III})] n_e \\ &- \frac{15}{2} \frac{a}{a} \left( \frac{kT_e n_B}{\mu} \right), \end{aligned}$$

# Benchmarking RT CODE



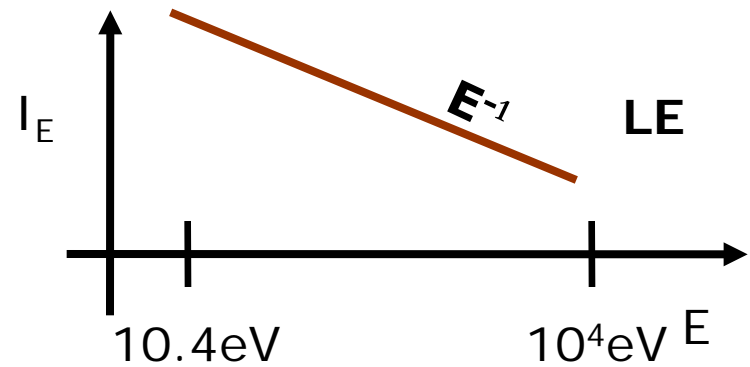
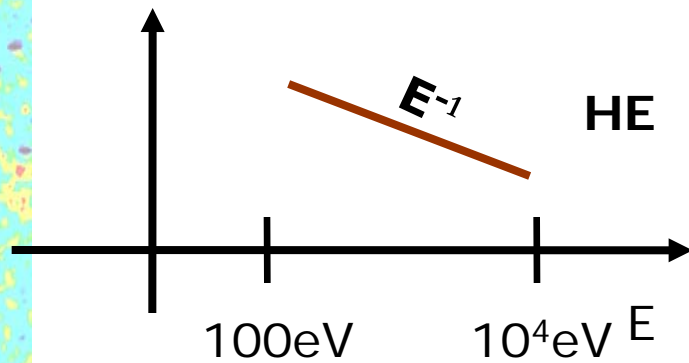
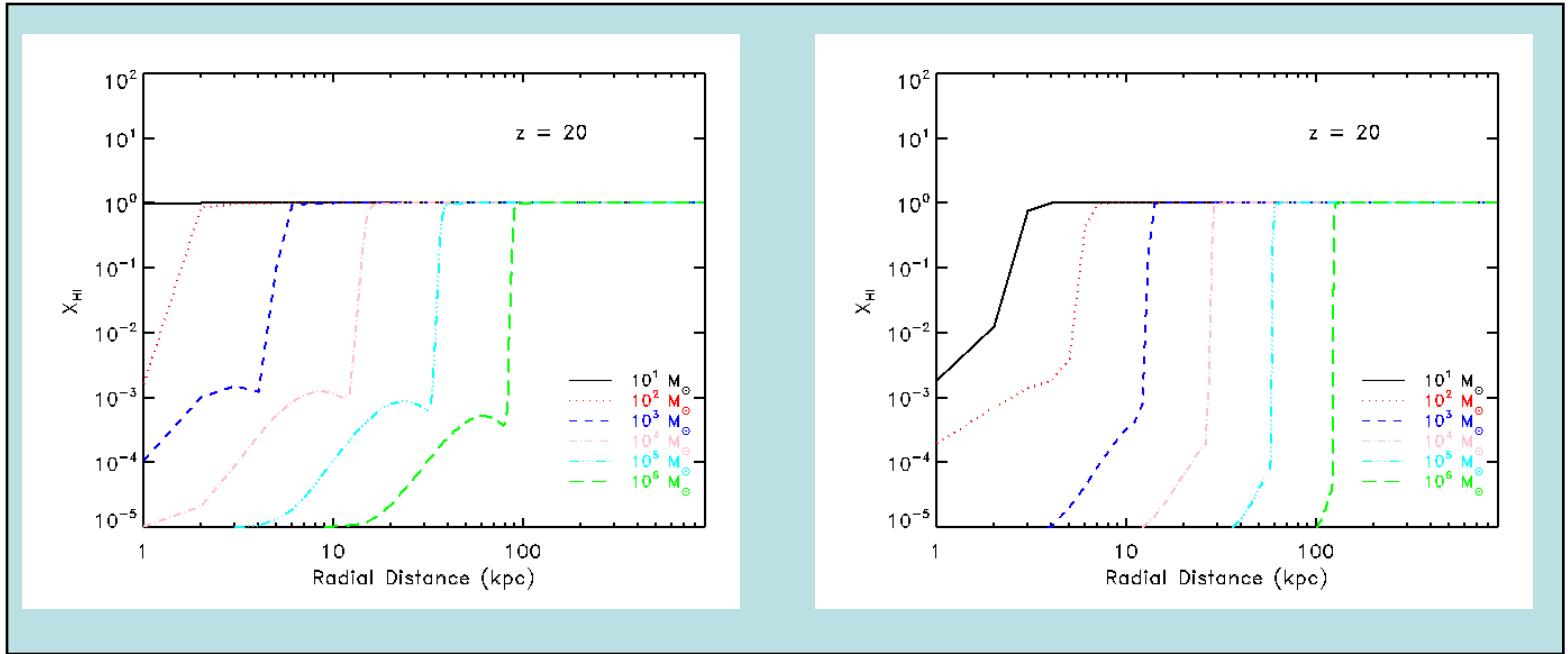
Stromgen sphere



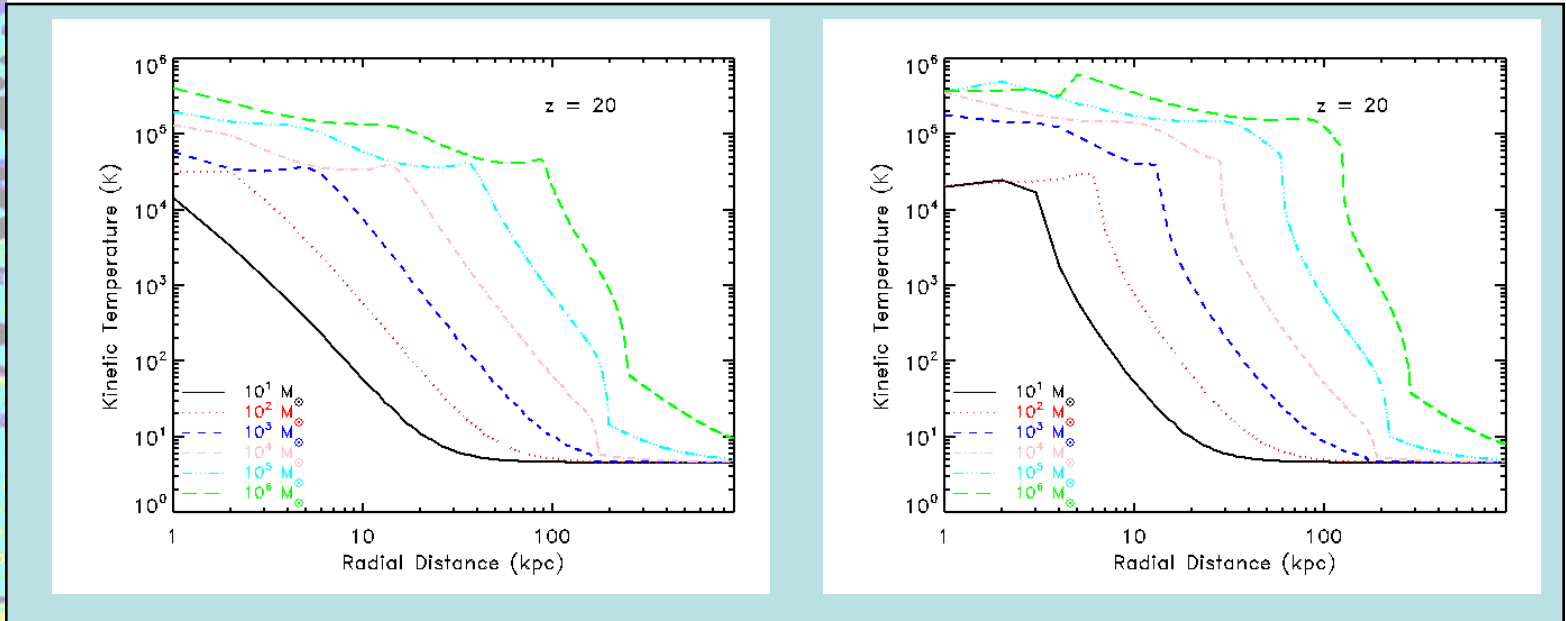
Shapiro & Giroux: expanding Universe



# Miniqso: Ionization



# Miniqso: Heating



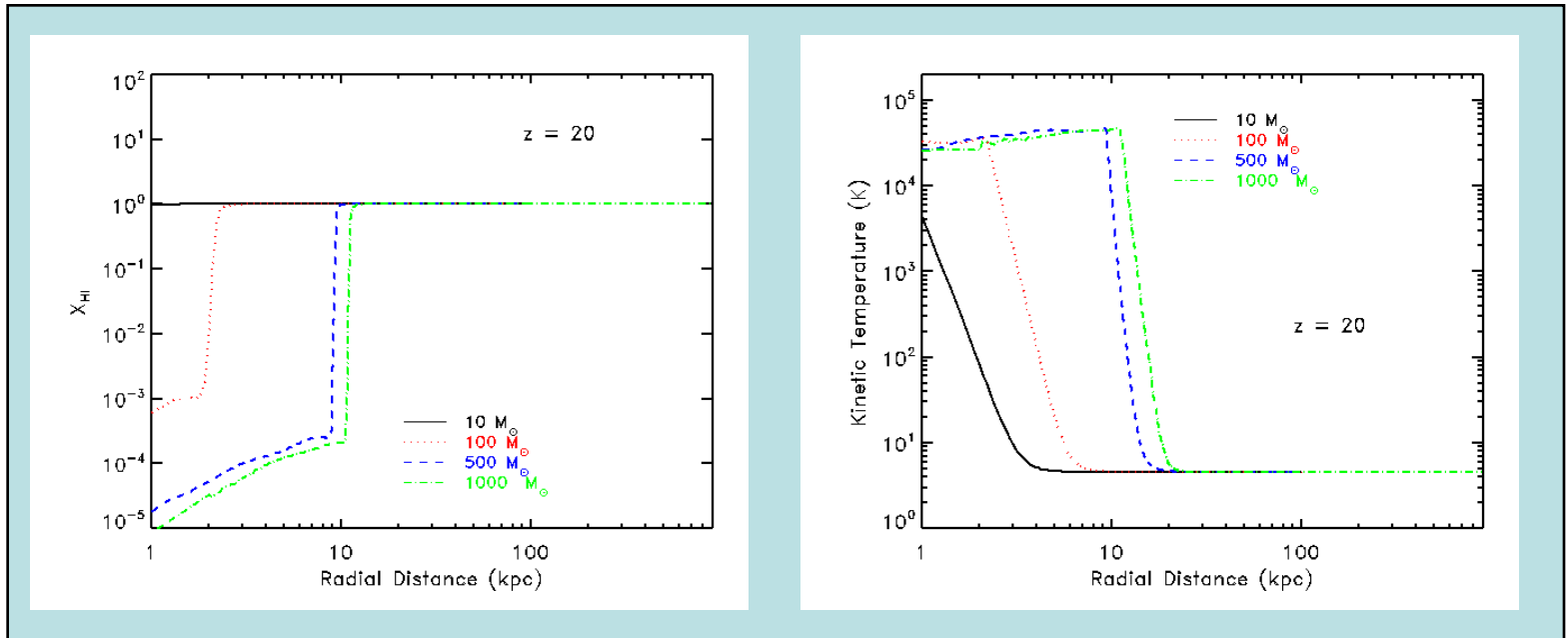
$$E_{\text{low}} = 200\text{eV}$$

$$E_{\text{low}} = 10.4\text{eV}$$

# Stars

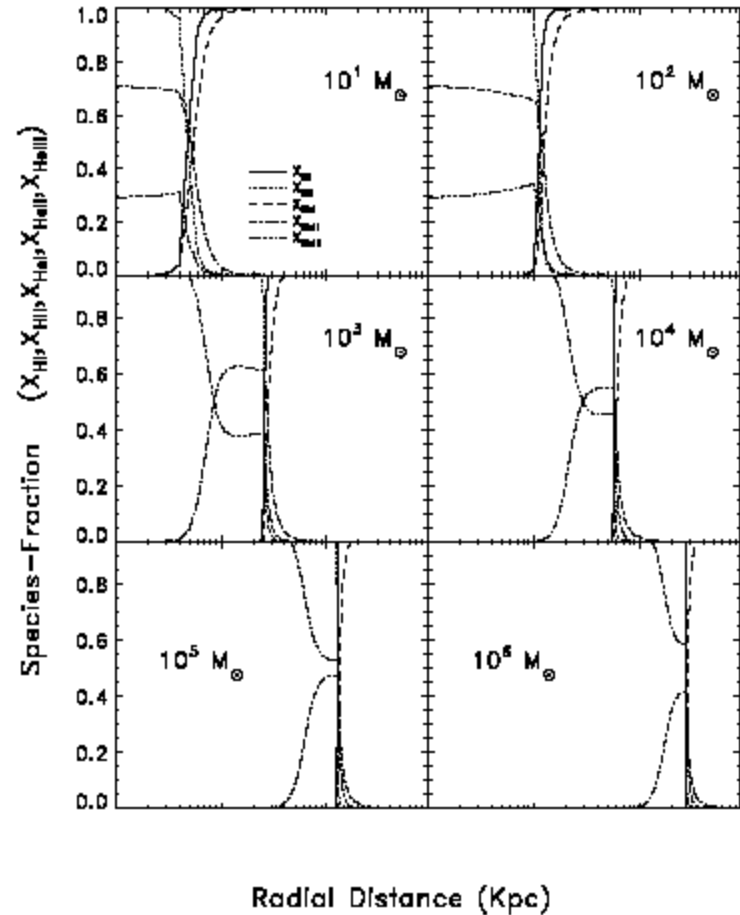
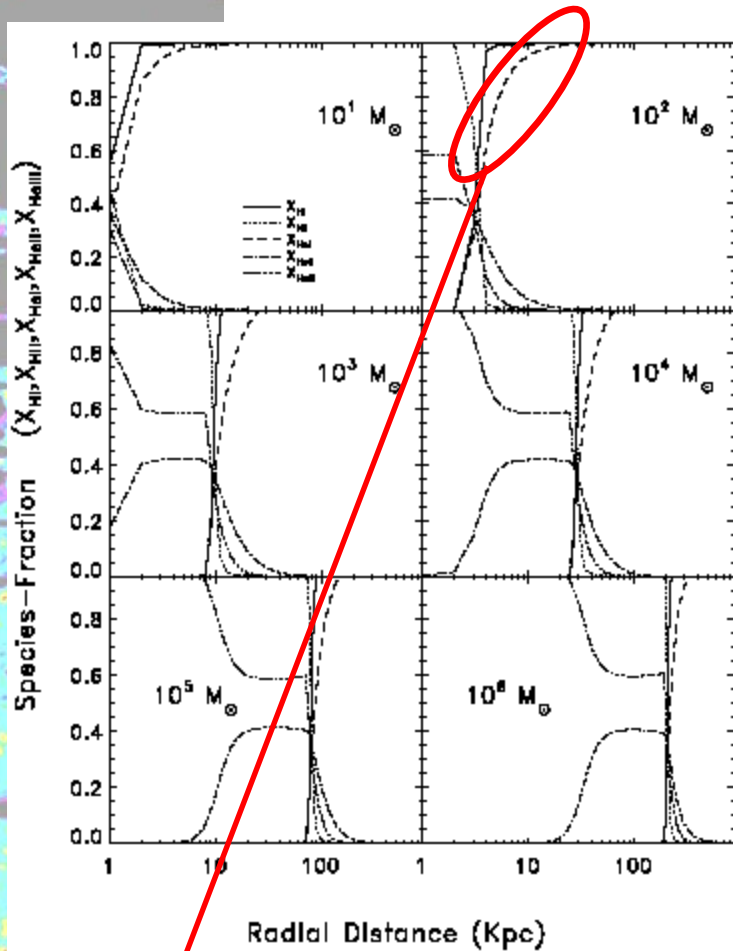
## Ionization

## Heating



Interesting....

## Quasars: ionization



He front

HE



Thomas & Zaroubi, in prep

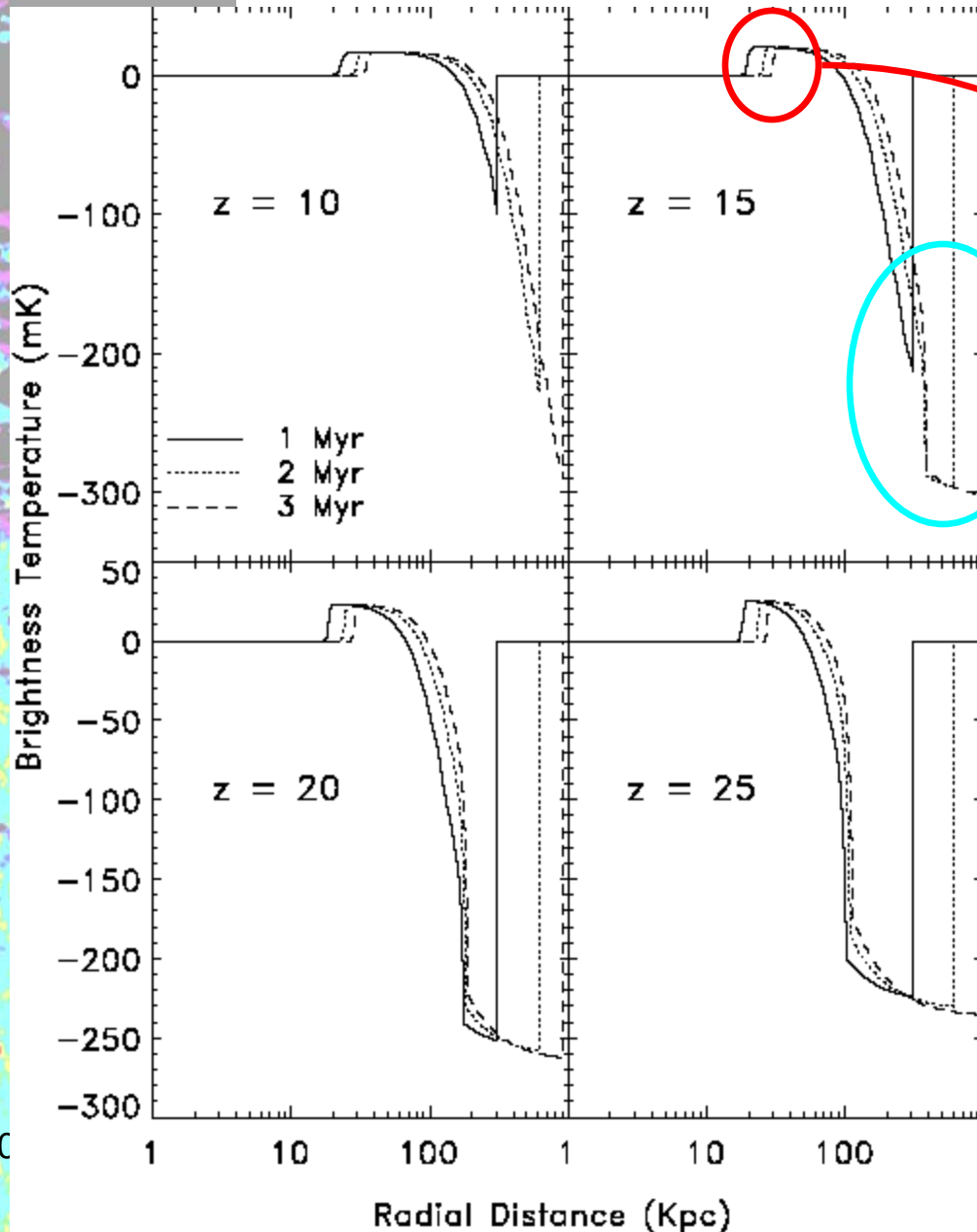
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# First sources & their signatures : *Brightness Temperature*



~20 to 30 mK  
(emission)

~-250 to -350 mK  
(absorption)

# Uncertainties encountered

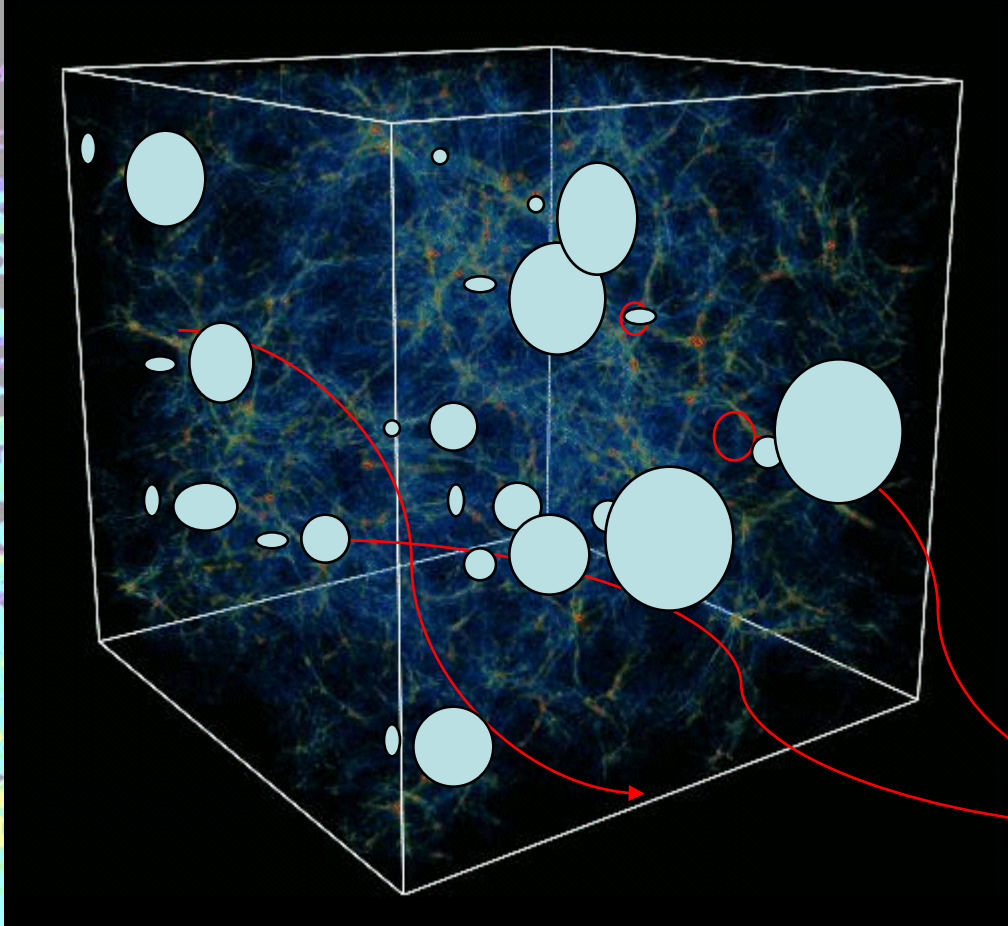
- Energy spectrum
- Span of energies
- Ly $\alpha$  coupling issues
- Ly $\alpha$  forest properties

**Simulate as many scenarios as possible**

- Redshift of turn-ON
- Spectral indices...

# Pseudo 3D -RT

Thomas et al., 2008, MNRAS



**BEARS -**

**Bubble**

**Expansion**

**Around**

**Radiative**

**Sources**

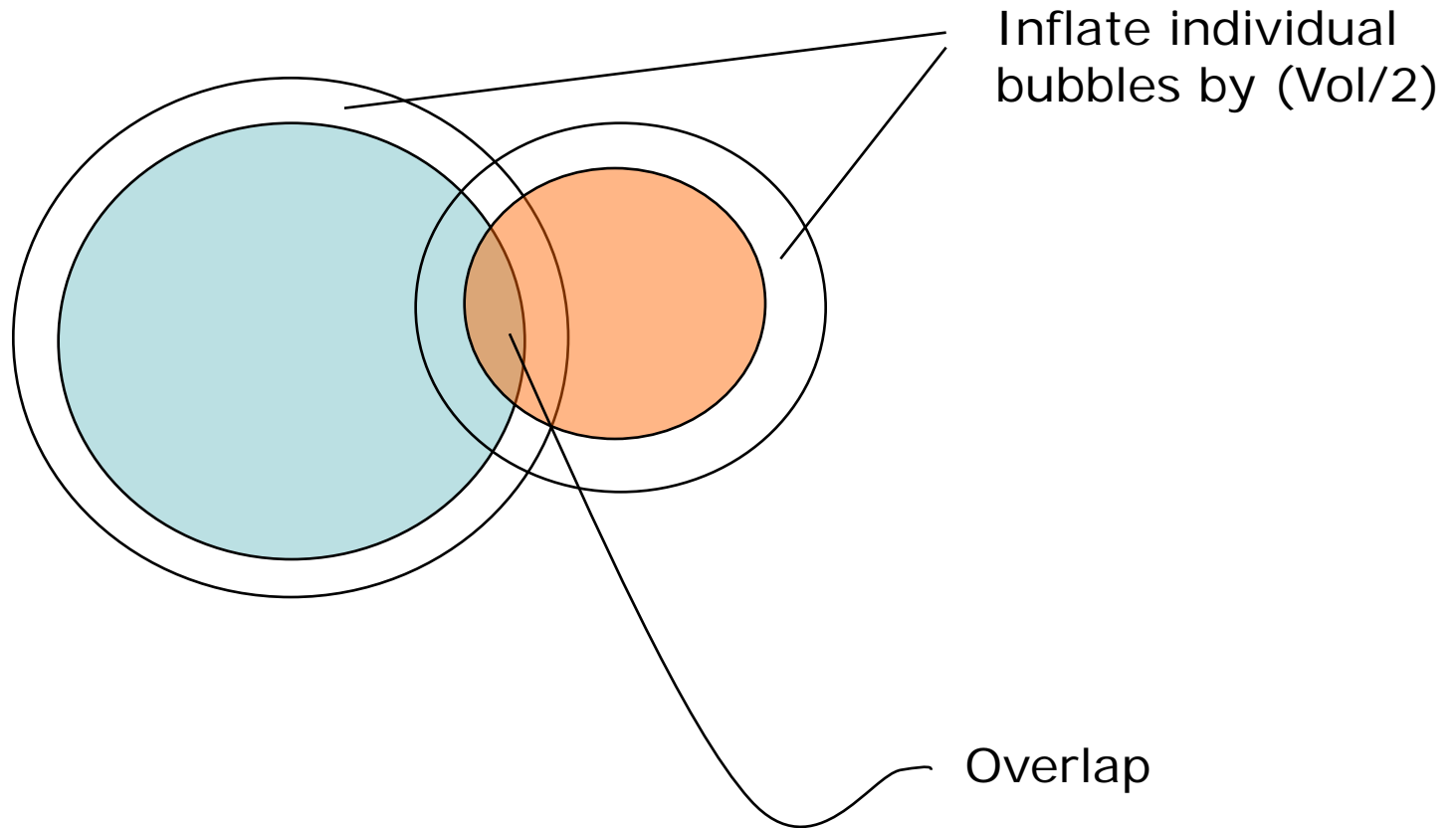
**Identify the  
haloes**

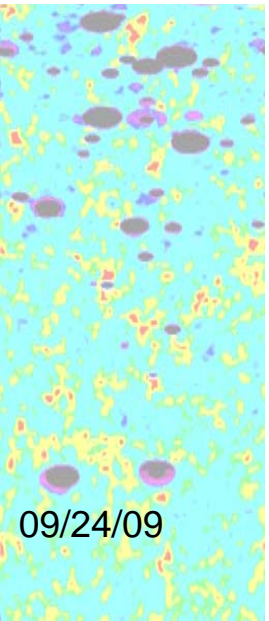
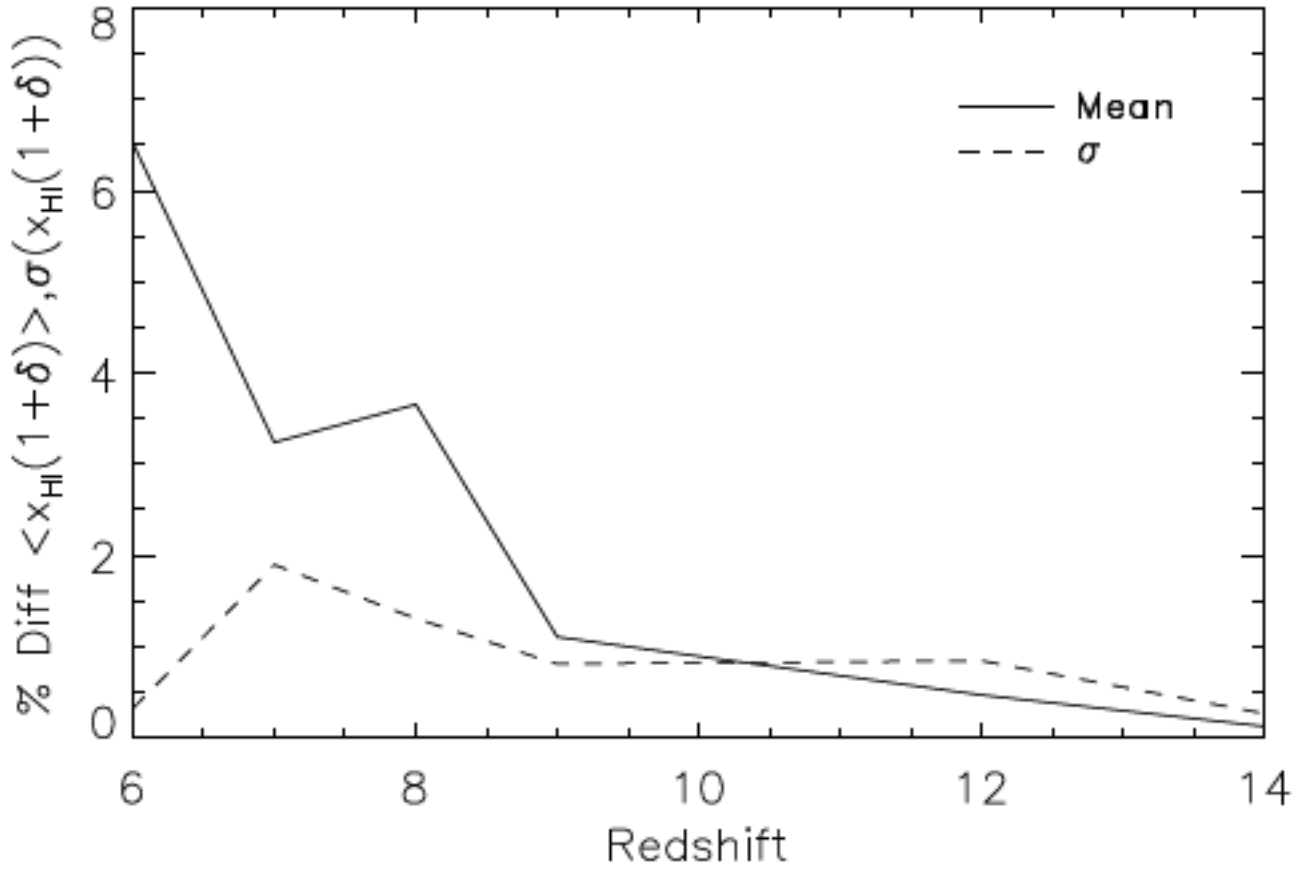
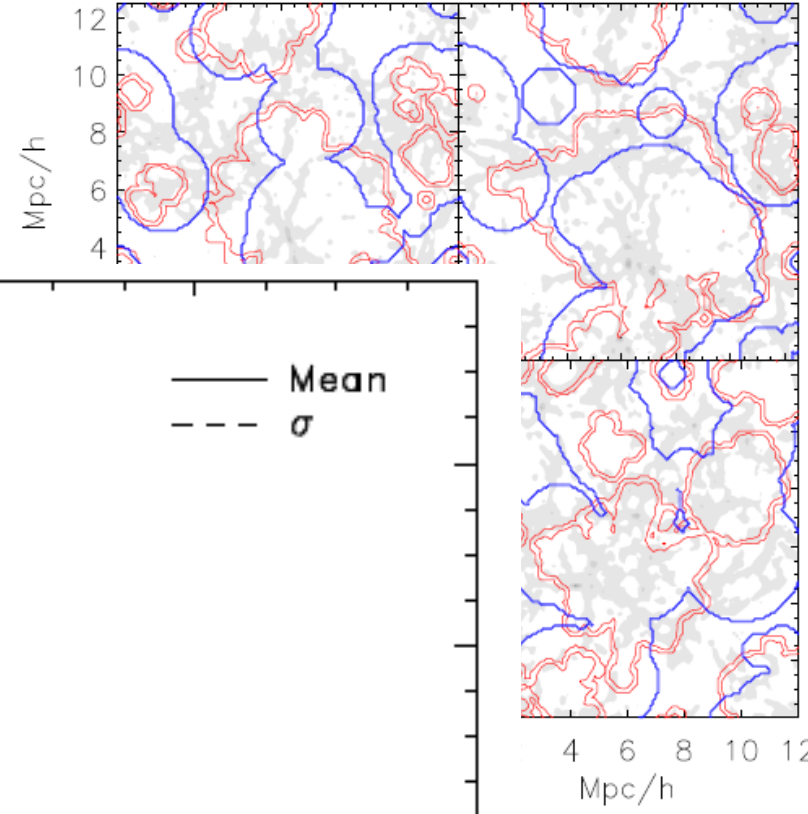
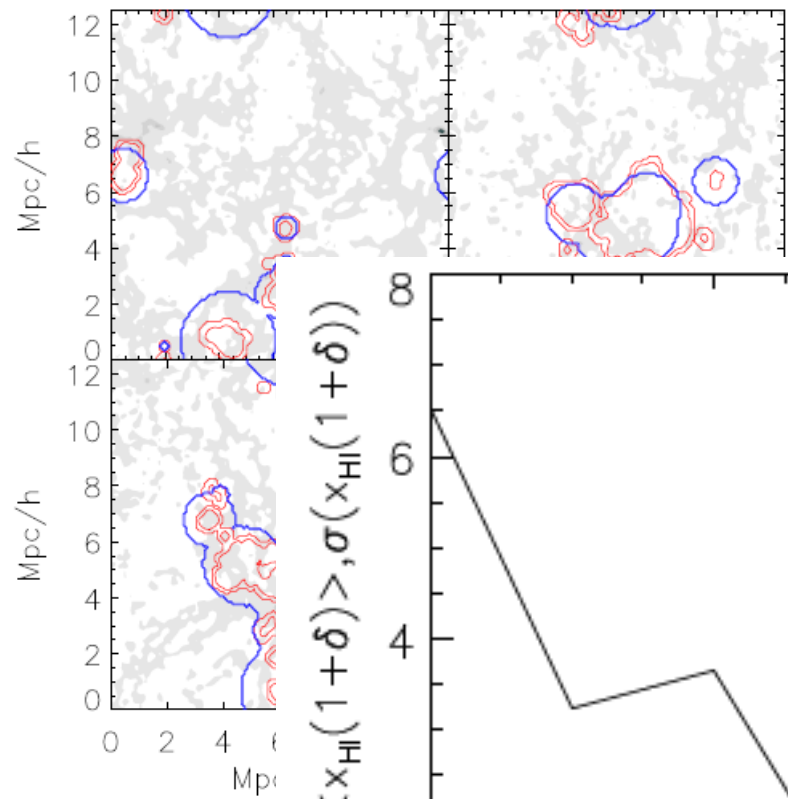
# Why should this approach work?

- $6 < z < 12$  : The Universe is still considerably homogeneous and isotropic even at small scales.
- Overdensity around the source is corrected for.



# Correction for overlap: Photon conservation





5

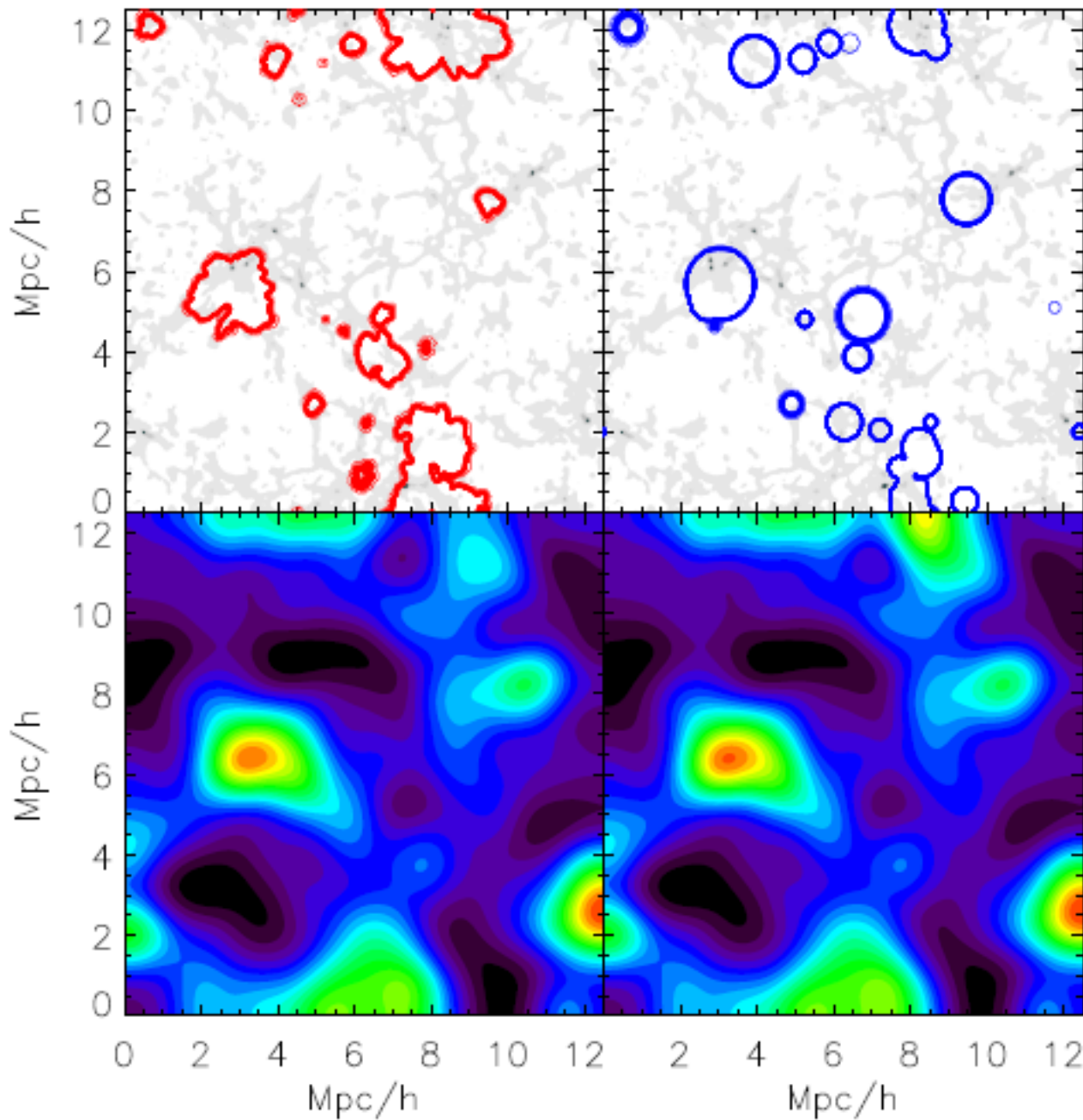
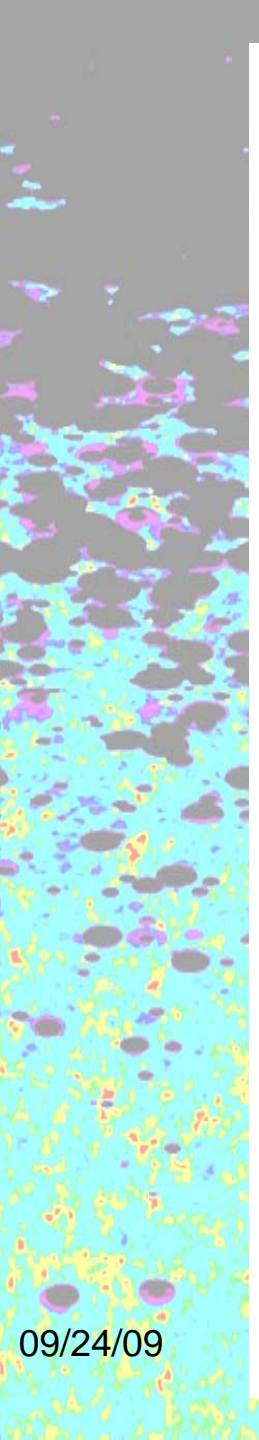
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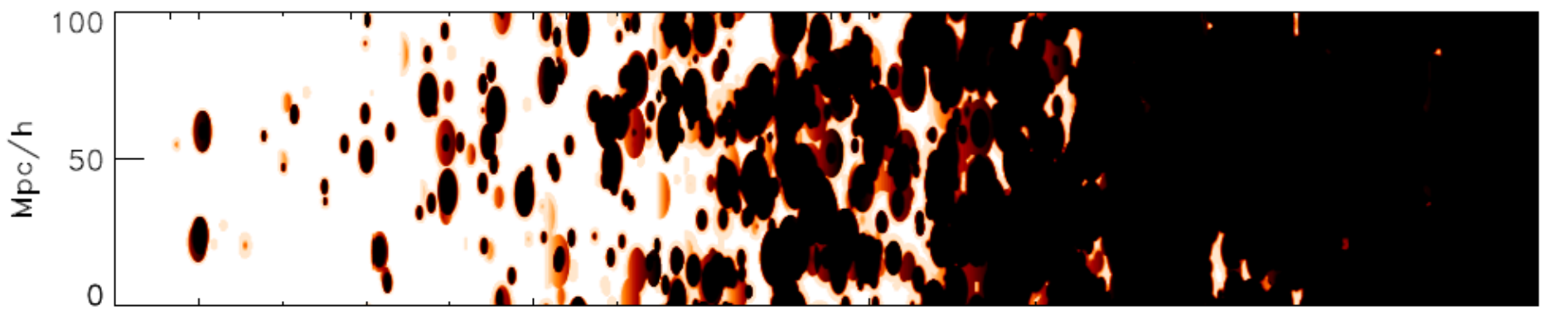


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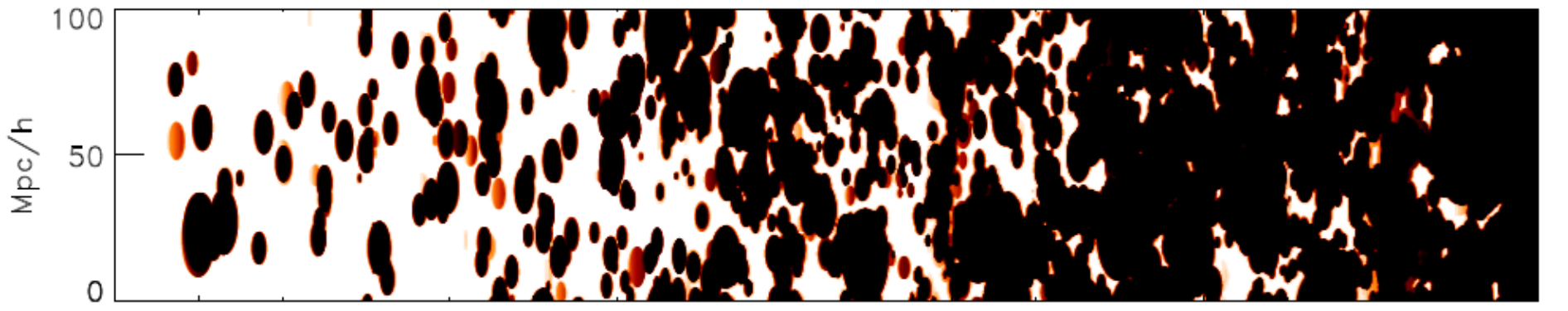


Redshift

**QUASARS**



**STARS**



Frequency [MHz]



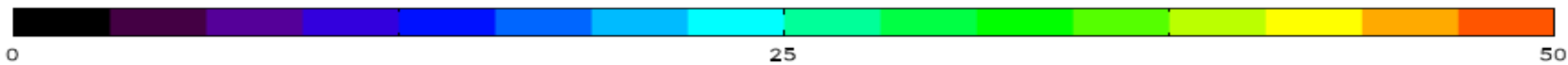
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Redshift

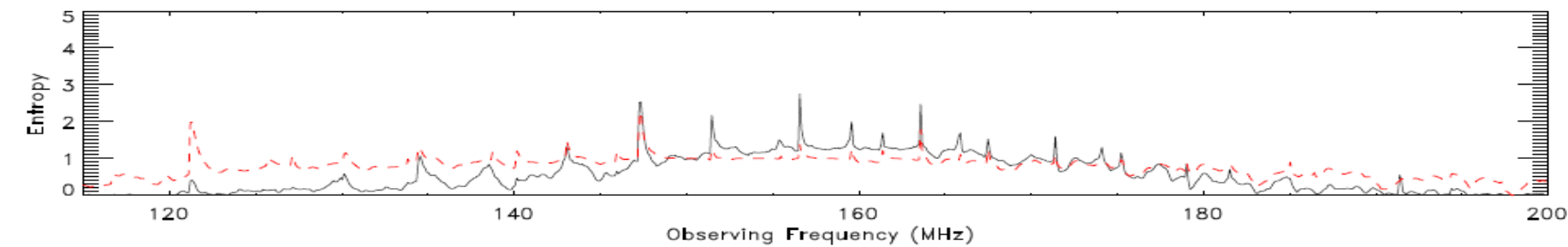
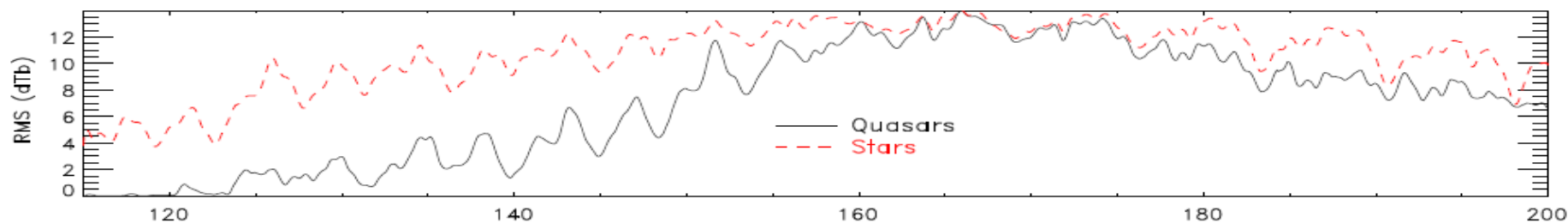
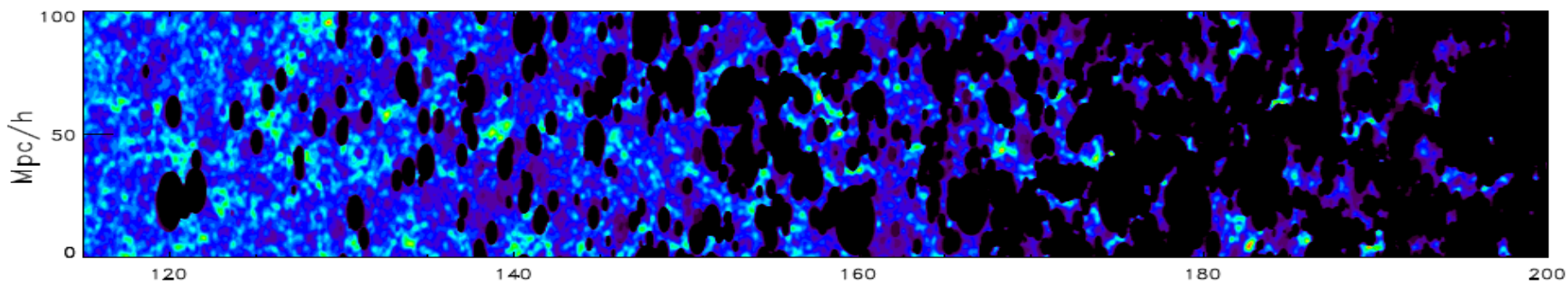
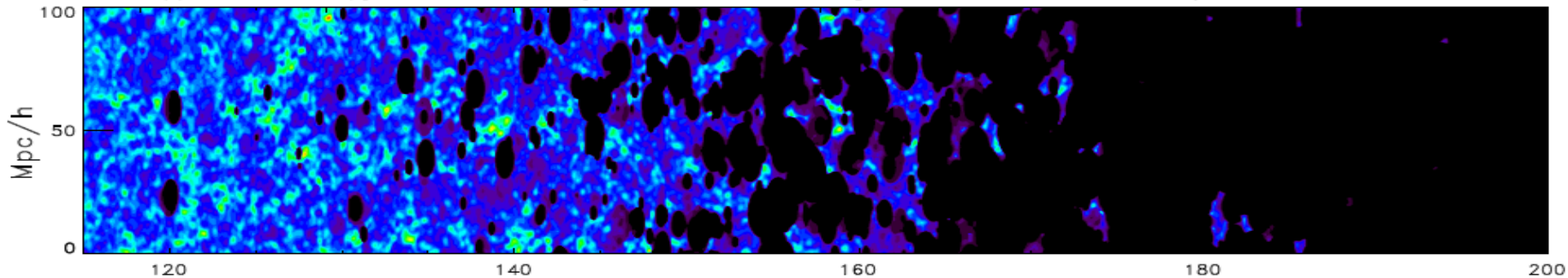
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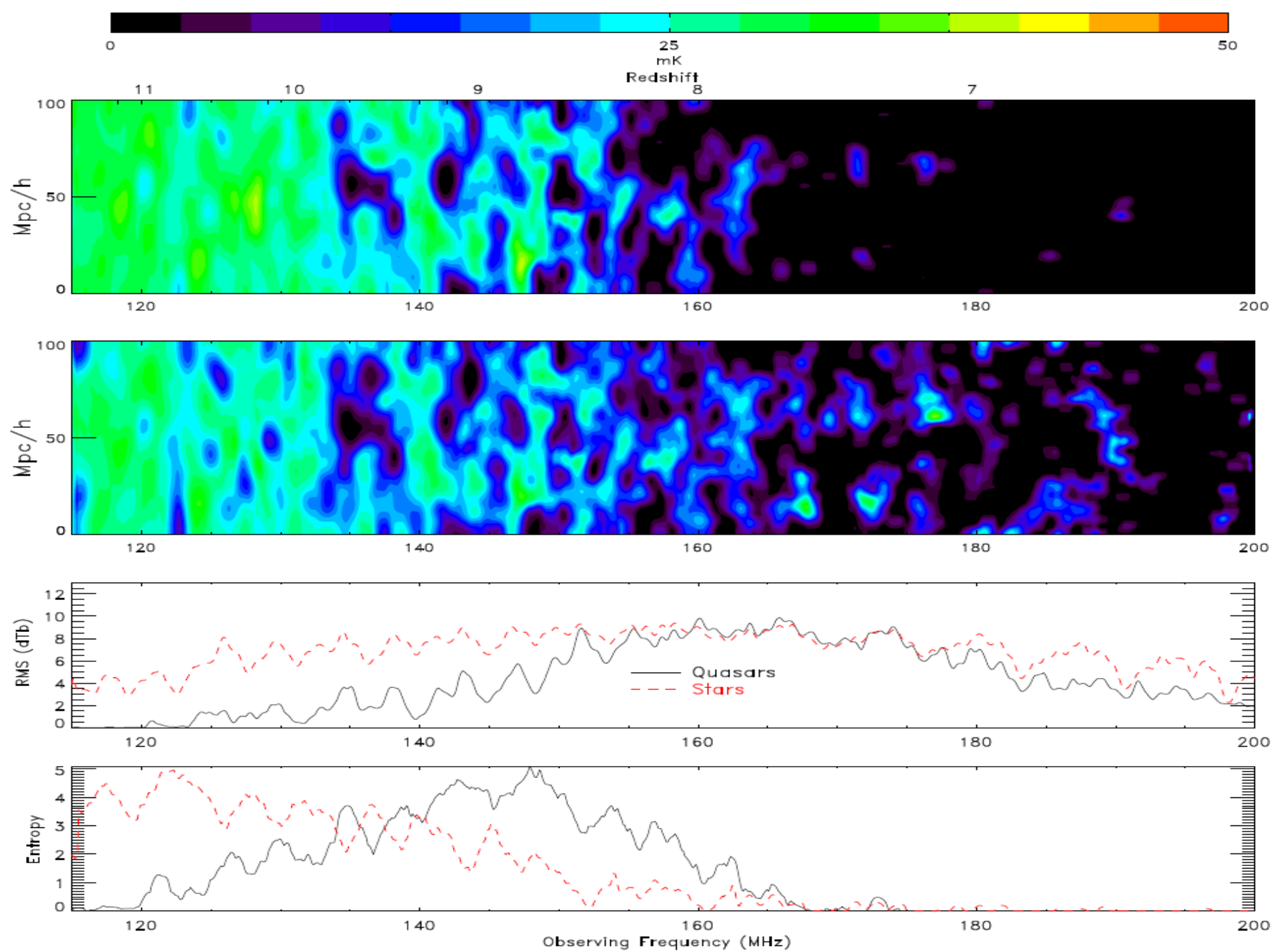
10

9

8

7





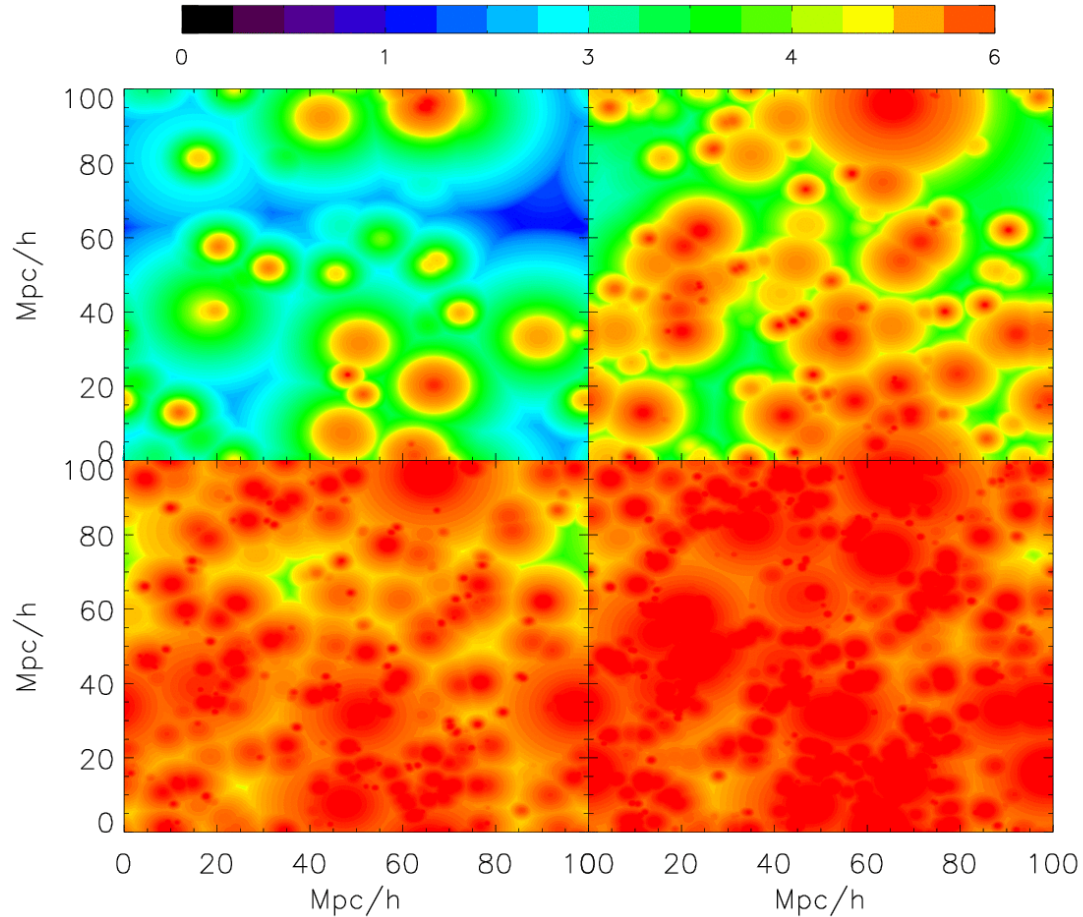
# Inclusion of temperature

- Previously the spin temperature was assumed to decoupled from the CMB.

$$\delta T_b = (20 \text{ mK}) (1 + \delta) \left( \frac{X_{HI}}{h} \right) \left( 1 - \frac{T_{CMB}}{T_{spin}} \right) \left( \frac{\Omega_b h^2}{0.0223} \right) \left[ \left( \frac{1+z}{10} \right) \left( \frac{0.24}{\Omega_m} \right) \right]^{1/2}$$

- The IGM in the Universe can be heated very differently depending on the sources.
- Therefore we need a self-consistent manner of including the temperature.

# Quasars: Temperature



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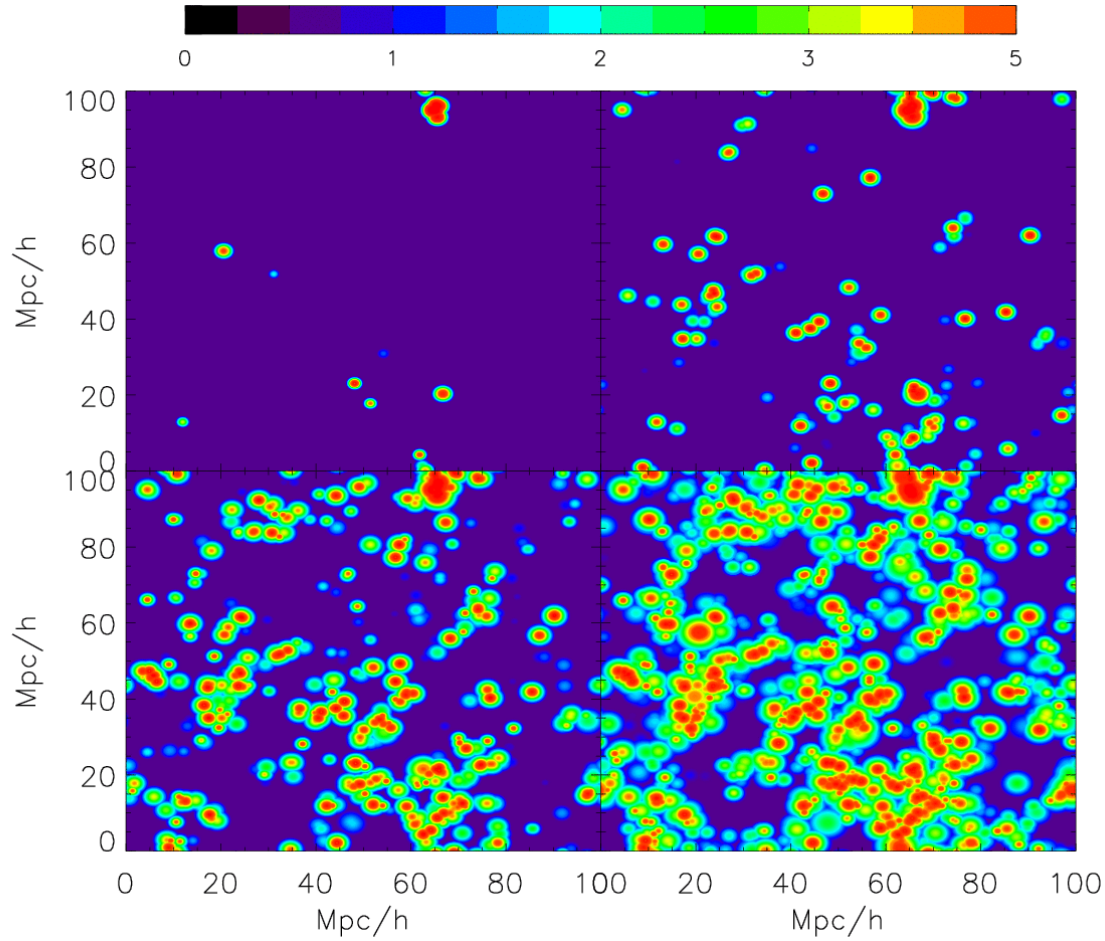
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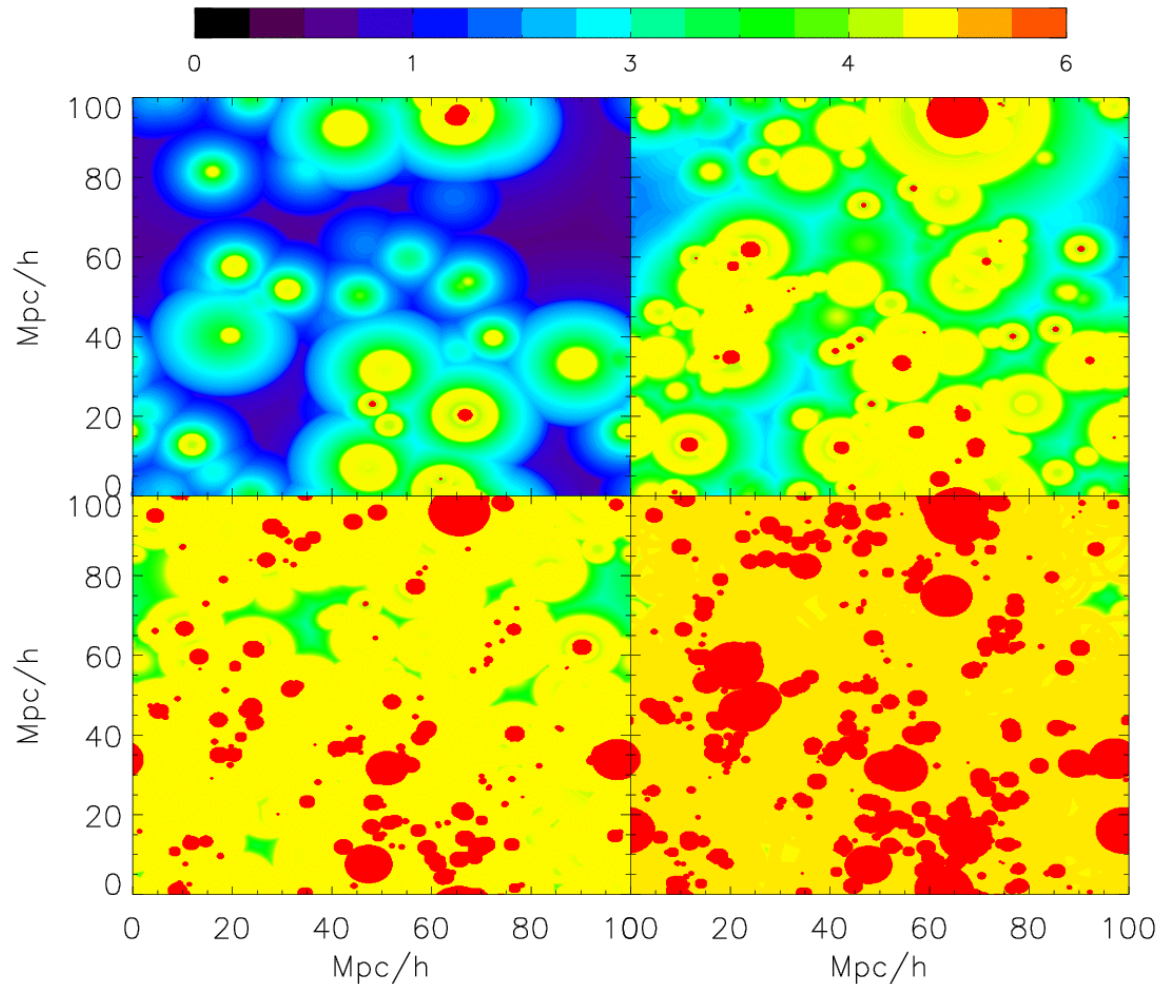
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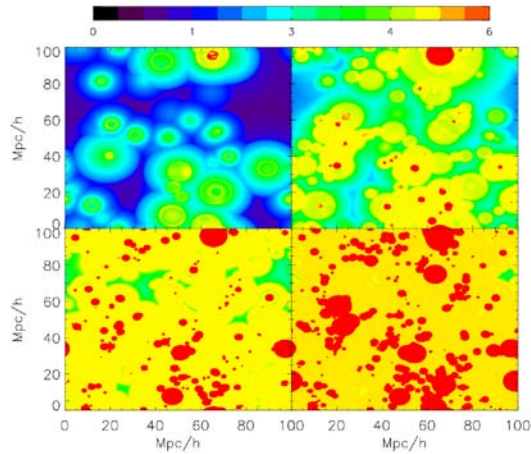
# Stars: Temperature



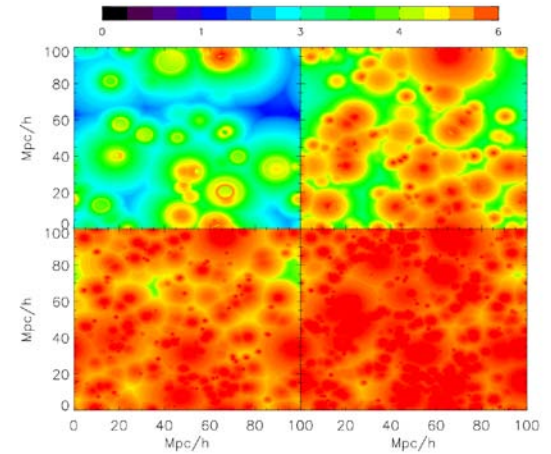
# Combi : Temperature



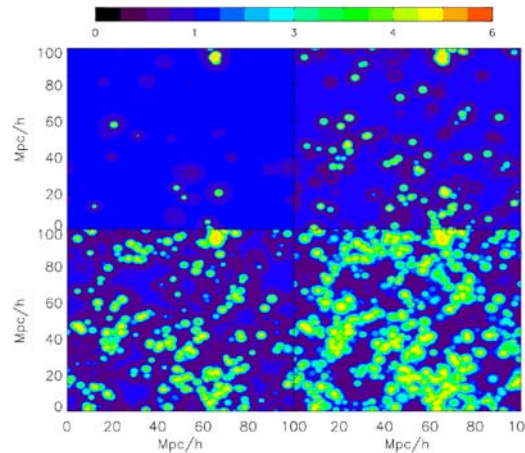
# Spin Temperatures



Combi

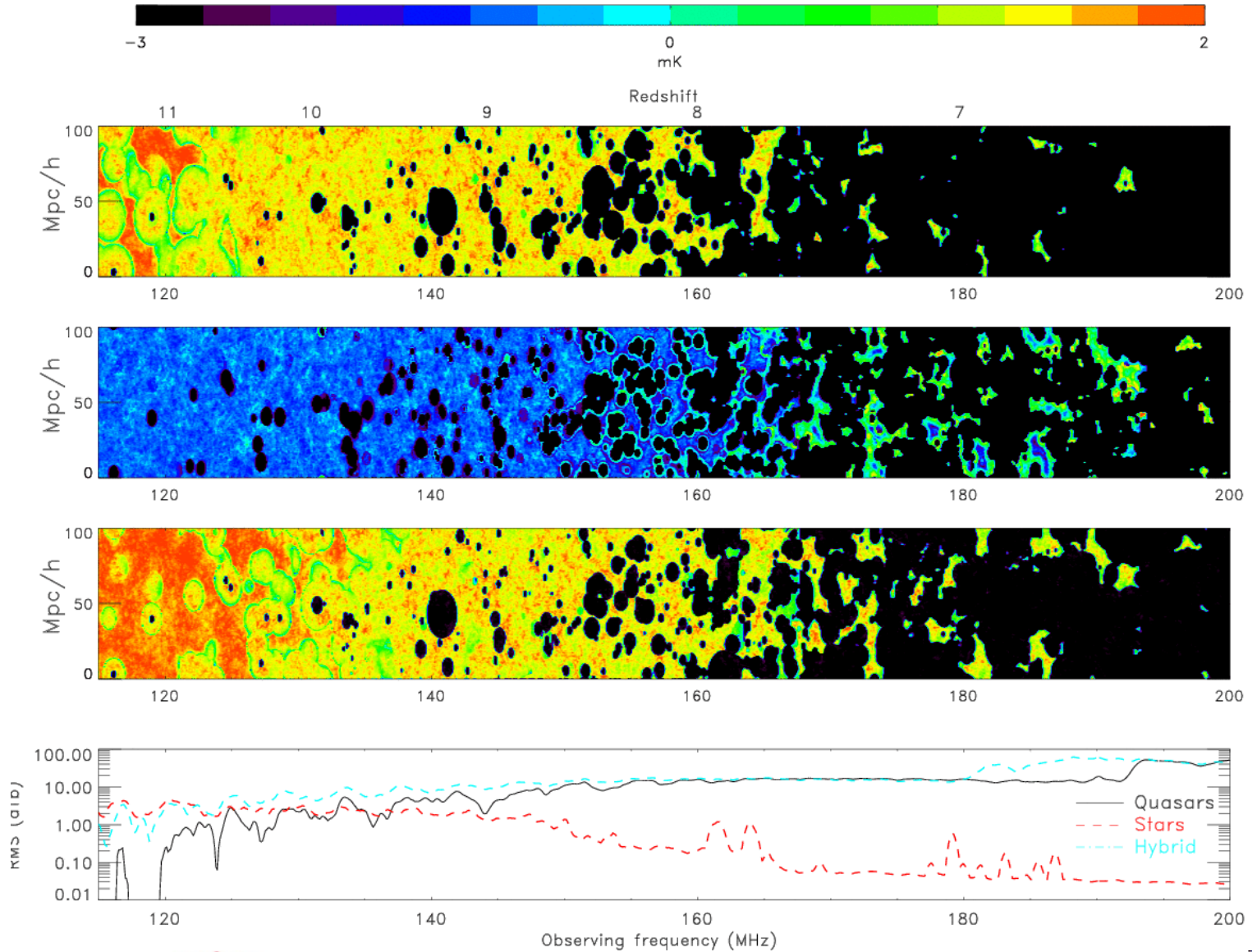


Qsos



Stars

# EoR Histories - I



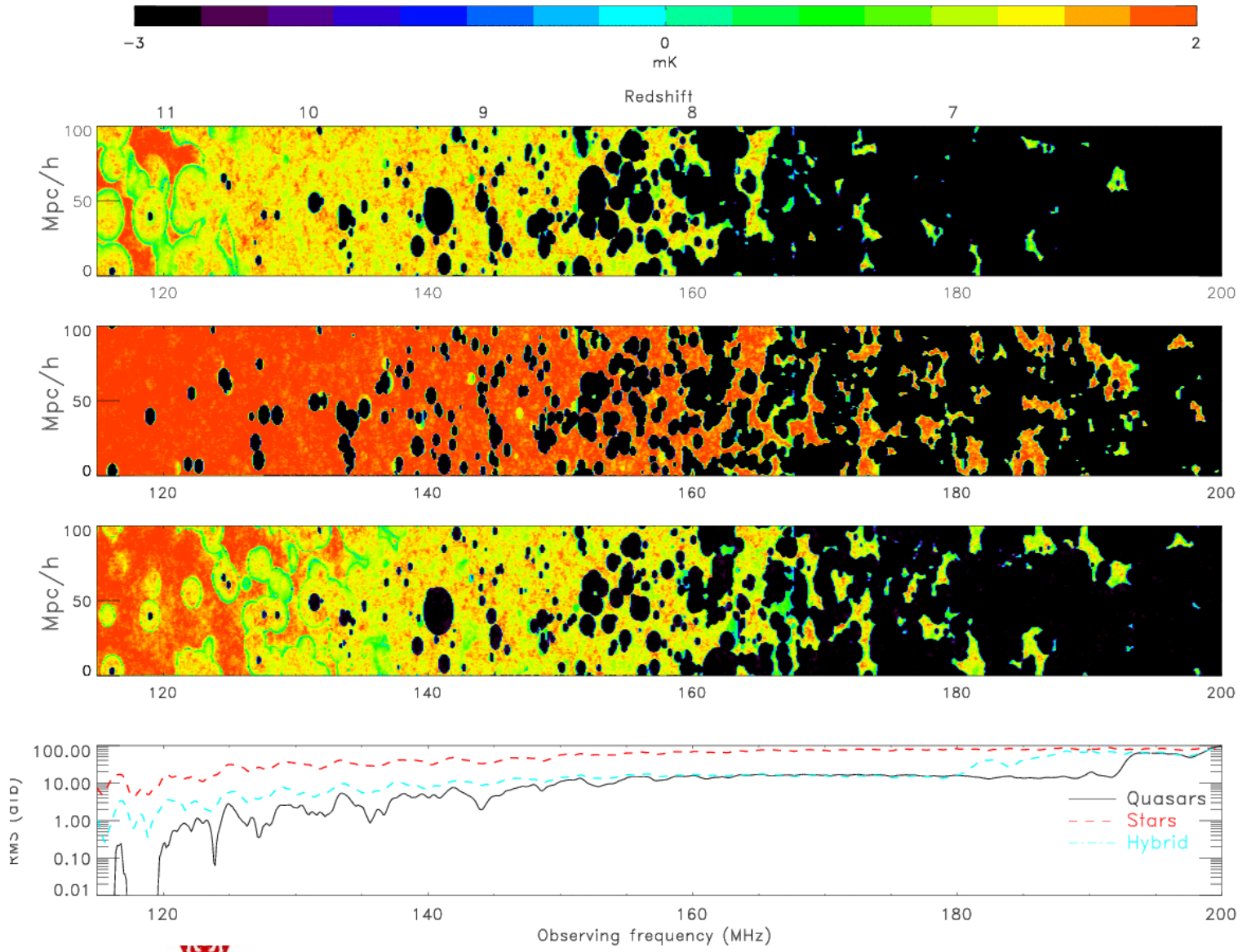
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# EoR Histories - II



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# To summarize

- Soft X-ray background is not a very strict constrain.
- Influence of Qsos on heating the IGM is significant.
- Various SEDs manifests themselves differently in the ionization and heating profiles around them. (a 1-D CODE will suffice to study these properties)
- We need to span a very large parameter space to get begin to get an idea of the numerous plausible reionization signals.
- BEARS is a good approximation to perform radiative transfer at high redshifts
- Tracking the temperature (spin) evolution self-consistently does make a lot of difference.

# Looking forward...

- LOFAR field:  $5^\circ \times 5^\circ$  ... translates to about 1 Gpc-comoving at EoR redshifts which will allow the measurement of:
  - Matter power spectrum at large  $z$
  - Sources of ionization and heating
  - The speed at which ionization occurs
  - Influence of reionization on galaxy formation at later stages.
- Cross-correlations studies.... not good to patch-up smaller boxes.
- Thinking of newer ways to accomplish larger boxes and better resolutions.

Thank you!

&

Stay tuned .....  
First fringes from LOFAR  
CORE in 2009

