Searching for 5-dimensional Nontrivial UV Fixed Point

Jin-Beom Bae

Ewha Woman's University & Institute for Basic Sciences

Based on ArXiv:1412.6549 work with Soo-Jong Rey

October 28, 2015

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

イロン 不同 とくほう イロン

-

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

◆□ > ◆□ > ◆ □ > ◆ □ > ● ● ● ● ●





æ



3D : Free Theory

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

æ



・ 同 ト ・ ヨ ト ・ ヨ ト

э



3-dimensional Ising theory Second order Phase Transition

同 ト イ ヨ ト イ ヨ ト



- ₹ 🖬 🕨



**□ > < = > <** 

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

<ロ> <部> < 部> < き> < き> < き</p>

 $\star$  Deform 3-dimensional free theory by quartic term,

 $\star$  Deform 3-dimensional free theory by quartic term,

$$S = \int d^3x \underbrace{\frac{1}{2} (\partial_\mu \phi^i(x))^2}_{\text{Free theory}}$$

イロン 不同 とくほう イロン

 $\star$  Deform 3-dimensional free theory by quartic term,

$$S = \int d^3x \underbrace{\frac{1}{2} (\partial_\mu \phi^i(x))^2}_{\text{Free theory}} + \underbrace{\frac{\lambda}{4!} (\phi^i(x)\phi^i(x))^2}_{\text{Quartic deformation}}$$

- 4 同 6 4 日 6 4 日 6

 $\star$  Deform 3-dimensional free theory by quartic term,

$$\begin{split} S &= \int d^3 x \underbrace{\frac{1}{2} (\partial_\mu \phi^i(x))^2}_{\text{Free theory}} + \underbrace{\frac{\lambda}{4!} (\phi^i(x) \phi^i(x))^2}_{\text{Quartic deformation}} \\ \left\{ \begin{array}{c} \dim((\phi^i(x) \phi^i(x))^2) = 2\\ \dim(\lambda) = 1 \end{array} \right\} \xrightarrow{} & \text{Non-perturbative}\\ \text{(with respect to } \lambda) @ IR \end{split}$$

-

 $\star$  Deform 3-dimensional free theory by quartic term,

$$\begin{split} S &= \int d^3x \underbrace{\frac{1}{2} (\partial_{\mu} \phi^i(x))^2}_{\text{Free theory}} + \underbrace{\frac{\lambda}{4!} (\phi^i(x) \phi^i(x))^2}_{\text{Quartic deformation}} \\ \left\{ \begin{array}{c} \dim((\phi^i(x) \phi^i(x))^2) = 2\\ \dim(\lambda) = 1 \end{array} \right\} \xrightarrow{} & \text{Non-perturbative}\\ (\text{with respect to } \lambda) @ IR \end{split}$$

\* Alternative :  $D = 4 - \varepsilon$  expansion.

 $\star$  Deform 3-dimensional free theory by quartic term,

$$\begin{split} S &= \int d^3x \underbrace{\frac{1}{2} (\partial_{\mu} \phi^i(x))^2}_{\text{Free theory}} + \underbrace{\frac{\lambda}{4!} (\phi^i(x) \phi^i(x))^2}_{\text{Quartic deformation}} \\ \left\{ \begin{array}{c} \dim((\phi^i(x) \phi^i(x))^2) = 2\\ \dim(\lambda) = 1 \end{array} \right\} \xrightarrow{} & \text{Non-perturbative}\\ (\text{with respect to } \lambda) @ IR \end{split}$$

\* Alternative :  $D = 4 - \varepsilon$  expansion.

 $\star$  The beta function of  $\mathit{O}(\mathit{N})$   $\phi^4\text{-theory}$  : ( $\varepsilon=4-\mathit{D})$ 

$$\beta(\lambda) = -\varepsilon\lambda + (N+8)\frac{\lambda^2}{8\pi^2} + \mathcal{O}(\lambda^3)$$

Fixed point at  $\lambda^* = 0$ (Gaussian) and  $\lambda^* = \frac{8\pi^2}{N+8}\varepsilon$  (Wilson-Fischer).

## Critical Exponents from $\epsilon$ -expansion

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

- \* @ \* \* 注 \* \* 注 \*

э

### Critical Exponents from $\epsilon$ -expansion

 $\star$  At Wilson-Fischer fixed point, the anomalous dimension of  $\phi$  and  $\phi^2$  are(by  $\varepsilon\text{-expansion})$  :

$$\begin{split} \Delta_{\phi} &= \frac{D}{2} - 1 + \gamma_{\phi} = 1 - \frac{\varepsilon}{2} + \frac{N+2}{4(N+8)^2} \varepsilon^2 + \mathcal{O}(\varepsilon^3) \\ \Delta_{\phi^2} &= D - 2 + \gamma_{\phi^2} = 2 - \frac{6}{N+8} \varepsilon + \mathcal{O}(\varepsilon^2) \end{split}$$

▲ □ ▶ ▲ □ ▶ ▲ □ ▶

 $\star$  At Wilson-Fischer fixed point, the anomalous dimension of  $\phi$  and  $\phi^2$  are(by  $\varepsilon\text{-expansion})$  :

$$\begin{split} \Delta_{\phi} &= \frac{D}{2} - 1 + \gamma_{\phi} = 1 - \frac{\varepsilon}{2} + \frac{N+2}{4(N+8)^2} \varepsilon^2 + \mathcal{O}(\varepsilon^3) \\ \Delta_{\phi^2} &= D - 2 + \gamma_{\phi^2} = 2 - \frac{6}{N+8} \varepsilon + \mathcal{O}(\varepsilon^2) \end{split}$$

These quantities are related to the *critical exponents*  $\eta$  and  $\nu$ , by  $\Delta_{\phi} = \frac{D}{2} - 1 + \frac{\eta}{2}$  and  $\Delta_{\phi^2} = D - \frac{1}{\nu}$ .

◆□ ▶ ◆□ ▶ ◆ □ ▶ ◆ □ ● ● ● ● ● ●

 $\star$  At Wilson-Fischer fixed point, the anomalous dimension of  $\phi$  and  $\phi^2$  are(by  $\varepsilon\text{-expansion})$  :

$$\begin{split} \Delta_{\phi} &= \frac{D}{2} - 1 + \gamma_{\phi} = 1 - \frac{\varepsilon}{2} + \frac{N+2}{4(N+8)^2} \varepsilon^2 + \mathcal{O}(\varepsilon^3) \\ \Delta_{\phi^2} &= D - 2 + \gamma_{\phi^2} = 2 - \frac{6}{N+8} \varepsilon + \mathcal{O}(\varepsilon^2) \end{split}$$

These quantities are related to the *critical exponents*  $\eta$  and  $\nu$ , by  $\Delta_{\phi} = \frac{D}{2} - 1 + \frac{\eta}{2}$  and  $\Delta_{\phi^2} = D - \frac{1}{\nu}$ .

\* This  $\varepsilon$ -perturbation carried out up to 7-loop. [Guida, Zinn-Justin 98] For  $\epsilon = 1$ , Borel resummation technique available.

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

・ロト ・回ト ・モト ・モト

æ

\* Helium-4 Specific heat measurement

[Lipa, Swanson, Nissen, Chui, Israelsson 96]



\* Helium-4 Specific heat measurement

[Lipa, Swanson, Nissen, Chui, Israelsson 96]



\* Measured value of the exponent  $\alpha$  is -0.0127(3).

\* Helium-4 Specific heat measurement

[Lipa, Swanson, Nissen, Chui, Israelsson 96]

・ 同 ト ・ ヨ ト ・ ヨ ト



\* Measured value of the exponent  $\alpha$  is -0.0127(3).

\* By scaling relation  $2 - \alpha = \nu D$ , measured exponent  $\alpha$  is related to the  $\nu$ (0.6709). This measured value agrees to the  $\varepsilon$ -expansion result of O(2) model in 3-dimension(0.671  $\pm$  0.005).

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

・ロト ・四ト ・ヨト ・ヨト

Critical exponent		
$\eta$ $\nu$		

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

<ロ> <同> <同> < 同> < 同>



Critical exponent	$\varepsilon$ expansion(1998)	
η	0.03650(500)	
ν	0.63050(250)	

э

□ ▶ ▲ 臣 ▶ ▲ 臣



Critical exponent	$\varepsilon$ expansion(1998)	Monte-Carlo(2010)	
$\eta$	0.03650(500)	0.03627(10)	
$\nu$	0.63050(250)	0.63002(10)	

- 4 回 > - 4 回 > - 4 回 >

э



[El-Showk, Paulos, Poland, Rychkov, Simmons-Duffin, Vichi 14]



イロト イポト イヨト イヨト

э



[El-Showk, Paulos, Poland, Rychkov, Simmons-Duffin, Vichi 14]



イロト イポト イヨト イヨト

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

<ロ> <同> <同> < 同> < 同>

э

Basic Strategy

\* Consider four point correlation function of same scalar operators  $\langle O(x_1)O(x_2)O(x_3)O(x_4) \rangle$  in Conformal Field Theory

< □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □

#### Basic Strategy

\* Consider four point correlation function of same scalar operators  $\langle O(x_1)O(x_2)O(x_3)O(x_4) \rangle$  in Conformal Field Theory

\* Assumes Unitarity and Crossing symmetry

イロト 不得 とうせい かほとう ほ

Basic Strategy

\* Consider four point correlation function of same scalar operators  $\langle O(x_1)O(x_2)O(x_3)O(x_4) \rangle$  in Conformal Field Theory

\* Assumes Unitarity and Crossing symmetry

Input

< ロ > < 同 > < 回 > < 回 > < □ > <

Basic Strategy

\* Consider four point correlation function of same scalar operators  $\langle O(x_1)O(x_2)O(x_3)O(x_4) \rangle$  in Conformal Field Theory

\* Assumes Unitarity and Crossing symmetry



直 ト イヨ ト イヨ ト

Basic Strategy

\* Consider four point correlation function of same scalar operators  $\langle O(x_1)O(x_2)O(x_3)O(x_4) \rangle$  in Conformal Field Theory

\* Assumes Unitarity and Crossing symmetry



▶ < ∃ > <</p>
# Conformal Bootstrap Program

Basic Strategy

\* Consider four point correlation function of same scalar operators  $\langle O(x_1)O(x_2)O(x_3)O(x_4) \rangle$  in Conformal Field Theory

\* Assumes Unitarity and Crossing symmetry



Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

<ロ> <同> <同> < 同> < 同>

Bootstrapping various dimensions

(人間) (人) (人) (人) (人) (人)

Bootstrapping various dimensions

 $\star$  D=2 : Ising theory, Minimal model.  $_{[\rm Rychkov, \ 11]}$ 

・ 同 ト ・ ヨ ト ・ ヨ ト

Bootstrapping various dimensions

- \* D = 2: Ising theory, Minimal model. [Rychkov, 11]
- $\star D = 3$ : Ising theory, O(N) class,  $O(N) \times O(N)$  class.

[El-Showk, Paulos, Poland, Rychkov, Simmons-Duffin, Vichi, 12], [Kos, Poland, Simmons-Duffin, 14] [Kos, Poland, Simmons-Duffin, 13], [Kos, Poland, Simmons-Duffin, Vichi, 15] [Nakayama, Ohtsuki, 14]

 $\mathcal{N} = 8$  Superconformal Bootstrap, Gross-Neveu.

Bootstrapping various dimensions

- \* D = 2 : Ising theory, Minimal model. [Rychkov, 11]
- $\star$  D = 3 : Ising theory, O(N) class,  $O(N) \times O(N)$  class.

[El-Showk, Paulos, Poland, Rychkov, Simmons-Duffin, Vichi, 12], [Kos, Poland, Simmons-Duffin, 14] [Kos, Poland, Simmons-Duffin, 13], [Kos, Poland, Simmons-Duffin, Vichi, 15] [Nakayama, Ohtsuki, 14]

 $\mathcal{N} = 8$  Superconformal Bootstrap, Gross-Neveu. [Chester, Lee, Pufu, Yacobi, 14] [Iliesiu, Kos, Poland, Pufu, Simmons-Duffin, Yacoby, 15]

 $\begin{array}{l} \star \ D=4: \ \mathcal{N}=4 \ \text{Superconformal Bootstrap,} \\ \text{[Beem, Rastelli, van Rees, 13]} \\ \mathcal{N}=2 \ \text{Superconformal Bootstrap,} \\ \text{[Beem, Lemos, Liendo, Rastelli, van Rees, 14], [Madalena Lemos, Pedro Liendo, 15]} \\ \mathcal{N}=1 \ \text{Superconformal Bootstrap.} \\ \text{[Poland, Stergiou, 15], [Poland, Simmons-Duffin, 11]} \end{array}$ 

Bootstrapping various dimensions

- \* D = 2 : Ising theory, Minimal model. [Rychkov, 11]
- $\star$  D = 3 : Ising theory, O(N) class,  $O(N) \times O(N)$  class.

[El-Showk, Paulos, Poland, Rychkov, Simmons-Duffin, Vichi, 12], [Kos, Poland, Simmons-Duffin, 14] [Kos, Poland, Simmons-Duffin, 13], [Kos, Poland, Simmons-Duffin, Vichi, 15] [Nakayama, Ohtsuki, 14]

 $\mathcal{N} = 8$  Superconformal Bootstrap, Gross-Neveu. [Chester, Lee, Pufu, Yacobi, 14] [Iliesiu, Kos, Poland, Pufu, Simmons-Duffin, Yacoby, 15]

- $\begin{array}{l} \star \ D=4: \mathcal{N}=4 \ \text{Superconformal Bootstrap,} \\ \text{[Beem, Rastelli, van Rees, 13]} \\ \mathcal{N}=2 \ \text{Superconformal Bootstrap,} \\ \text{[Beem, Lemos, Liendo, Rastelli, van Rees, 14], [Madalena Lemos, Pedro Liendo, 15]} \\ \mathcal{N}=1 \ \text{Superconformal Bootstrap.} \\ \text{[Poland, Stergiou, 15], [Poland, Simmons-Duffin, 11]} \end{array}$
- $\star D = \frac{p}{q} : \text{Extended } 4 \varepsilon \text{ expansion.}$ [El-Showk, Paulos, Poland, Rychkov, Simmons-Duffin, Vichi, 13]

Bootstrapping various dimensions

- \* D = 2 : Ising theory, Minimal model. [Rychkov, 11]
- \* D = 3: Ising theory, O(N) class,  $O(N) \times O(N)$  class. [EI-Showk, Paulos, Poland, Rychkov, Simmons-Duffin, Vichi, 12], [Kos, Poland, Simmons-Duffin, 14]

[Kos, Poland, Simmons-Duffin, 13], [Kos, Poland, Simmons-Duffin, Vichi, 12], [Kos, Poland, Simmons-Duffin, 13], [Nakayama, Ohtsuki, 14]

 $\mathcal{N} = 8$  Superconformal Bootstrap, Gross-Neveu. [Chester, Lee, Pufu, Yacobi, 14] [Iliesiu, Kos, Poland, Pufu, Simmons-Duffin, Yacoby, 15]

- $\begin{array}{l} \star \ D=4: \mathcal{N}=4 \ \text{Superconformal Bootstrap,} \\ \substack{[\text{Beem, Rastelli, van Rees, 13]}\\ \mathcal{N}=2 \ \text{Superconformal Bootstrap,} \\ \substack{[\text{Beem, Lemos, Liendo, Rastelli, van Rees, 14], [Madalena Lemos, Pedro Liendo, 15]}\\ \mathcal{N}=1 \ \text{Superconformal Bootstrap.} \\ \substack{[\text{Poland, Stergiou, 15], [Poland, Simmons-Duffin, 11]} \end{array}$
- $\star D = \frac{p}{q} : \text{Extended } 4 \varepsilon \text{ expansion.}$ [El-Showk, Paulos, Poland, Rychkov, Simmons-Duffin, Vichi, 13]
- \* D = 6 :  $\mathcal{N} = (2, 0)$  Superconformal bootstrap. [Beem, Lemos, Rastelli, van Rees, 15]

Bootstrapping various dimensions

- \* D = 2 : Ising theory, Minimal model. [Rychkov, 11]
- ★ D = 3 : Ising theory, O(N) class, O(N) × O(N) class. [EI-Showk, Paulos, Poland, Rychkov, Simmons-Duffin, Vichi, 12], [Kos, Poland, Simmons-Duffin, 14] [Kos, Poland, Simmons-Duffin, 13], [Kos, Poland, Simmons-Duffin, Vichi, 15]

[Nakayama, Ohtsuki, 14]

 $\mathcal{N} = 8$  Superconformal Bootstrap, Gross-Neveu. [Chester, Lee, Pufu, Yacobi, 14] [Iliesiu, Kos, Poland, Pufu, Simmons-Duffin, Yacoby, 15]

- $\begin{array}{l} \star \ D=4: \mathcal{N}=4 \ \text{Superconformal Bootstrap,} \\ \text{[Beem, Rastelli, van Rees, 13]} \\ \mathcal{N}=2 \ \text{Superconformal Bootstrap,} \\ \text{[Beem, Lemos, Liendo, Rastelli, van Rees, 14], [Madalena Lemos, Pedro Liendo, 15]} \\ \mathcal{N}=1 \ \text{Superconformal Bootstrap.} \\ \text{[Poland, Stergiou, 15], [Poland, Simmons-Duffin, 11]} \end{array}$
- $\star D = \frac{p}{q} : \text{Extended } 4 \varepsilon \text{ expansion.}$ [El-Showk, Paulos, Poland, Rychkov, Simmons-Duffin, Vichi, 13]
- \* D = 6 :  $\mathcal{N} = (2,0)$  Superconformal bootstrap. [Beem, Lemos, Rastelli, van Rees, 15]
- \* (Today) UV fixed point of D = 5 O(N) class.

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

<ロ> <部> < 部> < き> < き> < き</p>



<ロ> <同> <同> < 同> < 同>

D-dimensional Flat space







Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

< ロ > < 同 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ >



- 4 同 6 4 日 6 4 日 6



 $\star$  Time translation along  $\tau$  direction replaced by dilatation along r direction in flat space.

- 4 同 6 4 日 6 4 日 6



 $\star$  Time translation along  $\tau$  direction replaced by dilatation along r direction in flat space.

 $\rightarrow$  Every quantum states in conformal field theory are labeled by dilatation eigenvalue  $\Delta$  and spin /.



 $\star$  Time translation along  $\tau$  direction replaced by dilatation along r direction in flat space.

 $\rightarrow$  Every quantum states in conformal field theory are labeled by dilatation eigenvalue  $\Delta$  and spin /.

\*  $P^{\mu}$  acts as raising operator while  $K^{\mu}$  acts as lowering operator.

$$D(P^{\mu}|\Omega\rangle) = (\Delta + 1)|\Omega\rangle, \quad D(K^{\mu}(|\Omega\rangle) = (\Delta - 1)|\Omega\rangle$$

・ロト ・ 同 ト ・ ヨ ト ・ ヨ ト ・



 $\star$  Time translation along  $\tau$  direction replaced by dilatation along r direction in flat space.

 $\rightarrow$  Every quantum states in conformal field theory are labeled by dilatation eigenvalue  $\Delta$  and spin /.

 $\star$   $P^{\mu}$  acts as raising operator while  $K^{\mu}$  acts as lowering operator.

$$D(P^{\mu}|\Omega\rangle) = (\Delta + 1)|\Omega\rangle, \quad D(K^{\mu}(|\Omega\rangle) = (\Delta - 1)|\Omega\rangle$$

\* The primary state  $|\Omega^{P}\rangle$  is defined by  $K^{\mu}(0)|\Omega^{P}\rangle = 0$ .

 $\star$  Conformal symmetry fixes structure of the correlation functions.

◆ロ > ◆母 > ◆臣 > ◆臣 > ○日 ○ ○ ○ ○

 $\star$  Conformal symmetry fixes structure of the correlation functions.

$$\left\langle \phi_i(x_1)\phi_j(x_2)\right\rangle = \frac{\delta_{ij}}{|x_{12}|^{-2\Delta_{\phi}}}$$

◆ロ > ◆母 > ◆臣 > ◆臣 > ○日 ○ ○ ○ ○

 $\star$  Conformal symmetry fixes structure of the correlation functions.

$$\begin{split} \left\langle \phi_i(\mathbf{x}_1)\phi_j(\mathbf{x}_2)\right\rangle &= \frac{\delta_{ij}}{|\mathbf{x}_{12}|^{-2\Delta_{\phi}}}\\ \left\langle \phi_i(\mathbf{x}_1)\phi_j(\mathbf{x}_2)\phi_k(\mathbf{x}_3)\right\rangle &= \frac{f_{ijk}}{|\mathbf{x}_{12}|^{\Delta_{\phi_i}+\Delta_{\phi_j}-\Delta_{\phi_k}}|\mathbf{x}_{23}|^{\Delta_{\phi_k}+\Delta_{\phi_i}-\Delta_{\phi_j}}} \end{split}$$

◆ロ → ◆母 → ◆ 臣 → ◆ 臣 → 今 Q @

 $\star$  Conformal symmetry fixes structure of the correlation functions.

$$\begin{split} \left\langle \phi_{i}(x_{1})\phi_{j}(x_{2})\right\rangle &= \frac{\delta_{ij}}{|x_{12}|^{-2\Delta_{\phi}}} \\ \left\langle \phi_{i}(x_{1})\phi_{j}(x_{2})\phi_{k}(x_{3})\right\rangle &= \frac{f_{ijk}}{|x_{12}|^{\Delta_{\phi_{i}}+\Delta_{\phi_{j}}-\Delta_{\phi_{k}}}|x_{23}|^{\Delta_{\phi_{j}}+\Delta_{\phi_{k}}-\Delta_{\phi_{i}}}|x_{31}|^{\Delta_{\phi_{k}}+\Delta_{\phi_{i}}-\Delta_{\phi_{j}}} \\ \left\langle \phi_{i}(x_{1})\phi_{j}(x_{2})\phi_{k}(x_{3})\phi_{l}(x_{4})\right\rangle &= \frac{1}{|x_{12}|^{\Delta_{i}+\Delta_{j}}}\frac{|x_{23}|^{\Delta_{k}+\Delta_{j}}}{|x_{34}|^{\Delta_{k}+\Delta_{l}}}\left(\frac{|x_{24}|}{|x_{14}|}\right)^{\Delta_{12}}\left(\frac{|x_{14}|}{|x_{13}|}\right)^{\Delta_{13}} \end{split}$$

 $\star$  Conformal symmetry fixes structure of the correlation functions.

$$\begin{split} \left\langle \phi_{i}(x_{1})\phi_{j}(x_{2})\right\rangle &= \frac{\delta_{ij}}{|x_{12}|^{-2\Delta}\phi} \\ \left\langle \phi_{i}(x_{1})\phi_{j}(x_{2})\phi_{k}(x_{3})\right\rangle &= \frac{f_{ijk}}{|x_{12}|^{\Delta}\phi_{i}^{+\Delta}\phi_{j}^{-\Delta}\phi_{k}|_{x_{23}}|^{\Delta}\phi_{j}^{+\Delta}\phi_{k}^{-\Delta}\phi_{i}|_{x_{31}}|^{\Delta}\phi_{k}^{+\Delta}\phi_{i}^{-\Delta}\phi_{j}} \\ \left\langle \phi_{i}(x_{1})\phi_{j}(x_{2})\phi_{k}(x_{3})\phi_{l}(x_{4})\right\rangle &= \frac{1}{|x_{12}|^{\Delta_{i}^{+}\Delta_{j}}|_{x_{34}}|^{\Delta_{k}^{+}\Delta_{l}}} \left(\frac{|x_{24}|}{|x_{14}|}\right)^{\Delta_{12}} \left(\frac{|x_{14}|}{|x_{13}|}\right)^{\Delta_{13}} G(u \equiv \frac{x_{12}^{2}x_{24}^{2}}{x_{13}^{2}x_{24}^{2}}, v \equiv \frac{x_{14}^{2}x_{23}^{2}}{x_{13}^{2}x_{24}^{2}}) \end{split}$$

◆ロ → ◆母 → ◆ 臣 → ◆ 臣 → 今 Q @

 $\star$  Conformal symmetry fixes structure of the correlation functions.

$$\begin{split} \left\langle \phi_{i}(x_{1})\phi_{j}(x_{2})\right\rangle &= \frac{\delta_{ij}}{|x_{12}|^{-2\Delta_{\phi}}} \\ \left\langle \phi_{i}(x_{1})\phi_{j}(x_{2})\phi_{k}(x_{3})\right\rangle &= \frac{f_{ijk}}{|x_{12}|^{\Delta_{\phi_{i}}+\Delta_{\phi_{j}}-\Delta_{\phi_{k}}} |x_{23}|^{\Delta_{\phi_{j}}+\Delta_{\phi_{k}}-\Delta_{\phi_{i}}} |x_{31}|^{\Delta_{\phi_{k}}+\Delta_{\phi_{i}}-\Delta_{\phi_{j}}} \\ \left\langle \phi_{i}(x_{1})\phi_{j}(x_{2})\phi_{k}(x_{3})\phi_{l}(x_{4})\right\rangle &= \frac{1}{|x_{12}|^{\Delta_{i}+\Delta_{j}}} \left(\frac{|x_{24}|}{|x_{14}|}\right)^{\Delta_{12}} \left(\frac{|x_{14}|}{|x_{13}|}\right)^{\Delta_{13}} G(u \equiv \frac{x_{12}^{2}x_{24}^{2}}{x_{13}^{2}x_{24}^{2}}, v \equiv \frac{x_{14}^{2}x_{23}^{2}}{x_{13}^{2}x_{24}^{2}} \end{split}$$

\* Operator Product Expansion(OPE) in CFT is :

$$\phi(\mathbf{x}) \times \phi(\mathbf{y}) = \sum_{\mathcal{O}} f_{\phi\phi\mathcal{O}} P(\mathbf{x} - \mathbf{y}, \partial_{\mathbf{y}}) \mathcal{O}(\mathbf{y})$$

▲ロ▶ ▲冊▶ ▲ヨ▶ ▲ヨ▶ ヨ のの⊙

\* Conformal symmetry fixes structure of the correlation functions.

$$\begin{split} \left\langle \phi_{i}(x_{1})\phi_{j}(x_{2})\right\rangle &= \frac{\delta_{ij}}{|x_{12}|^{-2\Delta_{\phi}}} \\ \left\langle \phi_{i}(x_{1})\phi_{j}(x_{2})\phi_{k}(x_{3})\right\rangle &= \frac{f_{ijk}}{|x_{12}|^{\Delta_{\phi_{i}}+\Delta_{\phi_{j}}-\Delta_{\phi_{k}}} |x_{23}|^{\Delta_{\phi_{j}}+\Delta_{\phi_{k}}-\Delta_{\phi_{i}}} |x_{31}|^{\Delta_{\phi_{k}}+\Delta_{\phi_{i}}-\Delta_{\phi_{j}}} \\ \left\langle \phi_{i}(x_{1})\phi_{j}(x_{2})\phi_{k}(x_{3})\phi_{l}(x_{4})\right\rangle &= \frac{1}{|x_{12}|^{\Delta_{i}+\Delta_{j}}|x_{34}|^{\Delta_{k}+\Delta_{l}}} \left(\frac{|x_{24}|}{|x_{14}|}\right)^{\Delta_{12}} \left(\frac{|x_{14}|}{|x_{13}|}\right)^{\Delta_{13}} G(u \equiv \frac{x_{12}^{2}x_{24}^{2}}{x_{13}^{2}x_{24}^{2}}, v \equiv \frac{x_{14}^{2}x_{23}^{2}}{x_{13}^{2}x_{24}^{2}} \end{split}$$

\* Operator Product Expansion(OPE) in CFT is :

$$\phi(x) \times \phi(y) = \sum_{\mathcal{O}} f_{\phi\phi\mathcal{O}} P(x - y, \partial_y) \mathcal{O}(y)$$

\* Applying 12/34 channel OPE twice,

$$\left\langle \overleftarrow{\phi(x_1)\phi(x_2)} \overleftarrow{\phi(x_3)\phi(x_4)} \right\rangle = \sum_{\mathcal{O}} (f_{\phi\phi\mathcal{O}})^2 \frac{g_{\Delta,l}(u,v)}{\sum_{x_{12}}^{\Delta_{\phi}} \sum_{x_{34}}^{\Delta_{\phi}}} = \frac{\mathcal{G}(u,v)}{\sum_{x_{12}}^{\Delta_{\phi}} \sum_{x_{34}}^{\Delta_{\phi}}}$$

Function  $g_{\Delta,l}(u, v)$  is called by *conformal block*.

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

ヘロト ヘヨト ヘヨト ヘヨト

 $\star$  Closed form of the even dimensional conformal block :  $$_{[Dolan,\ Osborn\ 00]}$$ 

$$g_{\Delta,l}(z,\bar{z}) = \begin{cases} f_{\Delta+l}(z)f_{\Delta-l}(\bar{z}) + (z \leftrightarrow \bar{z}) & D=2\\ \frac{z\bar{z}}{z-\bar{z}}[k_{\Delta+l}(z)k_{\Delta-k-2}(\bar{z}) - (z \leftrightarrow \bar{z})] & D=4\\ k_{\beta}(x) \equiv x_{2}^{\frac{\beta}{2}}F_{1}(\frac{\beta}{2},\frac{\beta}{2};\beta;x) \end{cases}$$

 $\star$  Closed form of the even dimensional conformal block :  $$_{[Dolan, Osborn 00]}$$ 

$$g_{\Delta,l}(z,\bar{z}) = \begin{cases} f_{\Delta+l}(z)f_{\Delta-l}(\bar{z}) + (z\leftrightarrow\bar{z}) & D=2\\ \frac{z\bar{z}}{z-\bar{z}}[k_{\Delta+l}(z)k_{\Delta-k-2}(\bar{z}) - (z\leftrightarrow\bar{z})] & D=4\\ k_{\beta}(x) \equiv x_{2}^{-2}F_{1}(\frac{\beta}{2},\frac{\beta}{2};\beta;x) \end{cases}$$

\* This closed form is solution of conformal Casimir equation. [Dolan, Osborn 11]

$$\frac{1}{2}M^{AB}M_{BA}\mathcal{O}_{\Delta}^{(l)} = \frac{1}{2}c_{\Delta,l}\mathcal{O}_{\Delta}^{(l)}, \quad c_{\Delta,l} = l(l+D-2) + \Delta(\Delta-D)$$
$$\longrightarrow \mathcal{D}_{z,\bar{z}}G_{\Delta,l} = \frac{1}{2}c_{\Delta,l}G_{\Delta,l}, \quad G_{\Delta,l} \sim u^{\frac{1}{2}(\Delta-l)}(1+\mathcal{O}(u,1-v))$$

◆ロ > ◆母 > ◆臣 > ◆臣 > ○ 臣 ○ のへで

\* Closed form of the even dimensional conformal block : [Dolan, Osborn 00]

$$g_{\Delta,l}(z,\bar{z}) = \begin{cases} f_{\Delta+l}(z)f_{\Delta-l}(\bar{z}) + (z\leftrightarrow\bar{z}) & D=2\\ \frac{z\bar{z}}{z-\bar{z}}[k_{\Delta+l}(z)k_{\Delta-k-2}(\bar{z}) - (z\leftrightarrow\bar{z})] & D=4\\ k_{\beta}(x) \equiv x_{2}^{\frac{\beta}{2}}F_{1}(\frac{\beta}{2},\frac{\beta}{2};\beta;x) \end{cases}$$

\* This closed form is solution of conformal Casimir equation. [Dolan, Osborn 11]

$$\frac{1}{2}M^{AB}M_{BA}\mathcal{O}_{\Delta}^{(l)} = \frac{1}{2}c_{\Delta,l}\mathcal{O}_{\Delta}^{(l)}, \quad c_{\Delta,l} = l(l+D-2) + \Delta(\Delta-D)$$
$$\longrightarrow \mathcal{D}_{z,\bar{z}}G_{\Delta,l} = \frac{1}{2}c_{\Delta,l}G_{\Delta,l}, \quad G_{\Delta,l} \sim u^{\frac{1}{2}(\Delta-l)}(1+\mathcal{O}(u,1-v))$$

 $\star$  For odd dimension, no closed form expression is know. Alternatively, numerical value of conformal block is available via recursion relation.

[El-Showk, Paulos, Poland, Rychkov, Simmons-Duffin, Vichi 12], [Kos, Poland, Simmons-Duffin 13]

◆□ ▶ ◆□ ▶ ◆ □ ▶ ◆ □ ● ● ● ● ● ●

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

<ロ> <同> <同> <同> < 同> < 同> < 同> <

æ

\* Conformal symmetry fix the insertion points  $x_1, x_3, x_4$ .



<ロ> <同> <同> < 同> < 同>

\* Conformal symmetry fix the insertion points  $x_1, x_3, x_4$ .



\* Under this setting, two variables  $z, \bar{z}$  are related to the conformal cross ratios by  $u = z\bar{z}, v = (1 - z)(1 - \bar{z})$ .

\* Conformal symmetry fix the insertion points  $x_1, x_3, x_4$ .



\* Under this setting, two variables  $z, \bar{z}$  are related to the conformal cross ratios by  $u = z\bar{z}, v = (1 - z)(1 - \bar{z})$ .

\* Conformal block is two-variable function for *arbitrary spacetime dimension except* D = 1. Therefore, higher dimensional extension is very straightforward.

#### Bootstrap Constraint

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

・ロト ・回ト ・モト ・モト

æ

#### Bootstrap Constraint

\* Unitarity : 
$$\Delta \ge \frac{D}{2} - 1$$
 for spin 0 field,  
 $\Delta \ge D + l - 2$  for spin I field.  
 $f_{\phi\phi\mathcal{O}}^2 > 0.$ 

[Minwalla 97]

(日)

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

#### Bootstrap Constraint

\* Unitarity : 
$$\Delta \geq \frac{D}{2} - 1$$
 for spin 0 field,  
 $\Delta \geq D + I - 2$  for spin I field. [Minwalla 97]  
 $f_{\phi\phi\mathcal{O}}^2 > 0.$ 

\* Crossing symmetry of the 4-point correlation function :



- 4 同 6 4 日 6 4 日 6
#### Bootstrap Constraint

\* Unitarity : 
$$\Delta \geq \frac{D}{2} - 1$$
 for spin 0 field,  
 $\Delta \geq D + I - 2$  for spin I field. [Minwalla 97]  
 $f_{\phi\phi\mathcal{O}}^2 > 0.$ 

\* Crossing symmetry of the 4-point correlation function :



\* The Sum Rule from Unitarity and Crossing symmetry constraint : [Rattazzi, Rychkov, Tonni, Vichi 08]

$$u^{\Delta} - v^{\Delta} = \sum_{\substack{\Delta, l \\ \Delta \ge \Delta_{unit}}} f_{\phi\phi\mathcal{O}}^{2} \left( v^{\Delta} g_{\Delta,l}(u, v) - u^{\Delta} g_{\Delta,l}(v, u) \right)$$

### Bootstrap Constraint II

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

æ

### Bootstrap Constraint II

#### $\star$ Intermediate states(in OPE) are consist of



(日)

### Bootstrap Constraint II

#### $\star$ Intermediate states(in OPE) are consist of



\* The final Sum Rule in bootstrap program :

$$\begin{bmatrix} u^{\Delta} - v^{\Delta} = \sum_{\substack{\Delta, l \\ \Delta \ge \Delta_{\epsilon}(l=0), \\ \Delta \ge \Delta_{unit}(l \neq 0)}} f_{\phi\phi\mathcal{O}}^{2} (v^{\Delta}g_{\Delta,l}(u,v) - u^{\Delta}g_{\Delta,l}(v,u)) \end{bmatrix}$$

Searching for 5-dimensional Nontrivial UV Fixed Point

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

・ロト ・回ト ・モト ・モト

æ

#### \* The Full Bootstrap Constraint : Sum Rule

$$\underbrace{u^{\Delta} - v^{\Delta}}_{\mathcal{F}_{0}(u,v)} = \sum_{\overline{\partial}(\Delta,l)} \underbrace{f^{2}_{\phi\phi\mathcal{O}}}_{>0} \underbrace{(v^{\Delta}g_{\Delta,l}(u,v) - u^{\Delta}g_{\Delta,l}(v,u))}_{\mathcal{F}(u,v)}$$

・ロト ・回ト ・ヨト ・ヨト

 $\star$  The Full Bootstrap Constraint : Sum Rule  $\bigoplus$  it's Descendants

$$\underbrace{\underline{u^{\Delta} - v^{\Delta}}_{\mathcal{F}_{0}(u,v)} = \sum_{\eth(\Delta,l)} \underbrace{f^{2}_{\phi\phi\mathcal{O}}}_{>0} \underbrace{(v^{\Delta}g_{\Delta,l}(u,v) - u^{\Delta}g_{\Delta,l}(v,u))}_{\mathcal{F}(u,v)}}_{\mathcal{F}(u,v)}$$
$$\underbrace{\underbrace{\partial^{m}_{z}\partial^{n}_{\bar{z}}\mathcal{F}_{0}(z,\bar{z})}_{\mathcal{F}^{n,m}_{0}(z,\bar{z})} = \sum_{\eth(\Delta,l)} \underbrace{f^{2}_{\phi\phi\mathcal{O}}}_{>0} \underbrace{\underbrace{\partial^{m}_{z}\partial^{n}_{\bar{z}}\mathcal{F}(z,\bar{z})}_{\mathcal{F}^{n,m}_{0}(z,\bar{z})}}_{\mathcal{F}^{n,m}(z,\bar{z})}$$

イロン イロン イヨン イヨン

 $\star$  The Full Bootstrap Constraint : Sum Rule  $\bigoplus$  it's Descendants

$$\underbrace{\underbrace{u^{\Delta} - v^{\Delta}}_{\mathcal{F}_{0}(u,v)} = \sum_{\eth(\Delta,l)} \underbrace{f^{2}_{\phi\phi\mathcal{O}}}_{>0} \underbrace{(v^{\Delta}g_{\Delta,l}(u,v) - u^{\Delta}g_{\Delta,l}(v,u))}_{\mathcal{F}(u,v)}}_{\mathcal{F}(u,v)}$$
$$\underbrace{\underbrace{\partial_{z}^{m}\partial_{\bar{z}}^{n}\mathcal{F}_{0}(z,\bar{z})}_{\mathcal{F}_{0}^{n,m}(z,\bar{z})} = \sum_{\eth(\Delta,l)} \underbrace{f^{2}_{\phi\phi\mathcal{O}}}_{>0} \underbrace{\partial_{z}^{m}\partial_{\bar{z}}^{n}\mathcal{F}(z,\bar{z})}_{\mathcal{F}^{n,m}(z,\bar{z})}}$$

\* Define linear functional  $\Lambda : \Lambda(\mathcal{F}^{m,n}) \equiv \sum_{\substack{m,n \\ m,n}}^{m+n=k} c_{m,n} \mathcal{F}^{m,n},$  $\Lambda(\mathcal{F}_0^{m,n}) \equiv \sum_{\substack{m,n \\ m,n}}^{m+n=k} c_{m,n} \mathcal{F}_0^{m,n}$ 

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

イロト 不得 とうせい かほとう ほ

★ The Full Bootstrap Constraint : Sum Rule ⊕ it's Descendants

$$\underbrace{\underbrace{u^{\Delta} - v^{\Delta}}_{\mathcal{F}_{0}(u,v)} = \sum_{\eth(\Delta,l)} \underbrace{f^{2}_{\phi\phi\mathcal{O}}}_{>0} \underbrace{\underbrace{(v^{\Delta}g_{\Delta,l}(u,v) - u^{\Delta}g_{\Delta,l}(v,u))}_{\mathcal{F}(u,v)}}_{\mathcal{F}(u,v)}}_{\mathcal{F}(u,v)}$$
$$\underbrace{\underbrace{\partial^{m}_{z}\partial^{n}_{\bar{z}}\mathcal{F}_{0}(z,\bar{z})}_{\mathcal{F}^{n,m}(z,\bar{z})} = \sum_{\eth(\Delta,l)} \underbrace{f^{2}_{\phi\phi\mathcal{O}}}_{>0} \underbrace{\underbrace{\partial^{m}_{z}\partial^{n}_{\bar{z}}\mathcal{F}(z,\bar{z})}_{\mathcal{F}^{n,m}(z,\bar{z})}}_{\mathcal{F}^{n,m}(z,\bar{z})}}$$

\* Define linear functional  $\Lambda : \Lambda(\mathcal{F}^{m,n}) \equiv \sum_{\substack{m,n \\ m,n}}^{m+n=k} c_{m,n} \mathcal{F}^{m,n},$  $\Lambda(\mathcal{F}_0^{m,n}) \equiv \sum_{\substack{m,n \\ m,n}}^{m+n=k} c_{m,n} \mathcal{F}_0^{m,n}$ 

$$\Lambda\big(\mathcal{F}_0^{n,m}(z,\bar{z})\big) = \sum_{\eth(\Delta,l)} \underbrace{f_{\phi\phi\mathcal{O}}^2}_{>0} \Lambda\big(\mathcal{F}^{n,m}(z,\bar{z})\big)$$

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

◆□ ▶ ◆□ ▶ ◆ □ ▶ ◆ □ ● ● ● ● ● ●

★ The Full Bootstrap Constraint : Sum Rule ⊕ it's Descendants

$$\underbrace{\underbrace{u^{\Delta} - v^{\Delta}}_{\mathcal{F}_{0}(u,v)} = \sum_{\eth(\Delta,l)} \underbrace{f^{2}_{\phi\phi\mathcal{O}}}_{>0} \underbrace{\underbrace{(v^{\Delta}g_{\Delta,l}(u,v) - u^{\Delta}g_{\Delta,l}(v,u))}_{\mathcal{F}(u,v)}}_{\mathcal{F}(u,v)}}_{\mathcal{F}(u,v)}$$
$$\underbrace{\underbrace{\partial^{m}_{z}\partial^{n}_{\bar{z}}\mathcal{F}_{0}(z,\bar{z})}_{\mathcal{F}^{n,m}_{0}(z,\bar{z})} = \sum_{\eth(\Delta,l)} \underbrace{f^{2}_{\phi\phi\mathcal{O}}}_{>0} \underbrace{\underbrace{\partial^{m}_{z}\partial^{n}_{\bar{z}}\mathcal{F}(z,\bar{z})}_{\mathcal{F}^{n,m}_{0}(z,\bar{z})}}_{\mathcal{F}^{n,m}_{0}(z,\bar{z})}$$

\* Define linear functional  $\Lambda : \Lambda(\mathcal{F}^{m,n}) \equiv \sum_{\substack{m,n \\ m,n}}^{m+n=k} c_{m,n} \mathcal{F}^{m,n},$  $\Lambda(\mathcal{F}_0^{m,n}) \equiv \sum_{\substack{m,n \\ m,n}}^{m+n=k} c_{m,n} \mathcal{F}_0^{m,n}$ 

$$\overline{\Lambda(\mathcal{F}_0^{n,m}(z,\bar{z}))} = \sum_{\eth(\Delta,l)} \frac{f_{\phi\phi\mathcal{O}}^2}{>0} \overline{\Lambda(\mathcal{F}^{n,m}(z,\bar{z}))}$$

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

イロン 不同 とくほう イロン

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

・ロト ・回ト ・モト ・モト

æ



Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

<ロ> <同> <同> < 同> < 同>

æ



\* RHS of sum rule : Linear combination (with positive coefficient  $f^2_{\phi\phi\mathcal{O}}$ ) of  $\Lambda(\mathcal{F}^{n,m}(z,\bar{z}))$  forms cone.

 $\star$  LHS of sum rule : Red vector.

▲ □ ▶ ▲ □ ▶ ▲ □ ▶



\* RHS of sum rule : Linear combination (with positive coefficient  $f^2_{\phi\phi\mathcal{O}}$ ) of  $\Lambda(\mathcal{F}^{n,m}(z,\bar{z}))$  forms cone.

\* LHS of sum rule : Red vector.

[Rattazzi, Rychkov, Tonni, Vichi 08]

Linear Programming Find a vector  $\vec{c} = (c_1, c_2, \dots, c_n)$  which minimize  $\vec{c}^T \cdot \vec{b} = c_1 b_1 + c_2 b_2 + \dots + c_n b_n$  subject to constraints  $\vec{b}^T \cdot A > \vec{a}$ .

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

・ロト ・回ト ・モト ・モト

æ



<ロ> <同> <同> < 同> < 同>

æ



• If the red vector always inside the cone regardless of  $c_{m,n}$ , that conformal field theory is *consistent* with bootstrap constraints.

/□ ▶ < 글 ▶ < 글



• If the red vector always inside the cone regardless of  $c_{m,n}$ , that conformal field theory is *consistent* with bootstrap constraints.

• If there is a set  $c_{m,n}$  which makes red vector outside the cone, that conformal field theory is *inconsistent* with bootstrap constraints.

- 4 周 ト 4 戸 ト 4 戸 ト

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

Basic Strategy

イロン イロン イヨン イヨン

Basic Strategy

Inputs :  $\Delta_{\phi}$ ,  $\Delta_{\epsilon}$ 

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

#### Basic Strategy

Is it consistent with Unitarity and Crossing symmetry?

Inputs :  $\Delta_{\phi}$ ,  $\Delta_{\epsilon}$ 

(By Linear Programming or Semi-definite Programming)

- ∢ ≣ ▶

#### Basic Strategy

Is it consistent with Unitarity and Crossing symmetry?

Inputs :  $\Delta_{\phi}, \Delta_{\epsilon}$ 

(By Linear Programming or Semi-definite Programming)



< ∃ →

#### Basic Strategy



Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

-

## Bootstrapping 3D $\mathbb{Z}_2$ theory

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

◆ロ > ◆母 > ◆臣 > ◆臣 > ● ● ● ● ●

## Bootstrapping 3D $\mathbb{Z}_2$ theory

 $\star$  Bootstrapping 3D  $\mathbb{Z}_2$  theory shows

[El-Showk, Paulos, Poland, Rychkov, Simmons-Duffin, Vichi 12]



## Bootstrapping 3D $\mathbb{Z}_2$ theory

 $\star$  Bootstrapping 3D  $\mathbb{Z}_2$  theory shows

[El-Showk, Paulos, Poland, Rychkov, Simmons-Duffin, Vichi 12]



\* The boundary has kink, whose location is  $(\Delta_{\phi}, \Delta_{\epsilon}) = (0.518, 1.413)$ . This numerical value indeed agrees to scaling dimension of relevant primary operators  $(\phi, \epsilon)$ .

/⊒ ► < ∃ ►

## Sum Rule in O(N) global symmetry

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

(日)

## Sum Rule in O(N) global symmetry

\* Operator product of two primary scalar fields decomposed by

$$\phi_i \times \phi_j \sim \sum_{S} \delta_{ij} \mathcal{O} + \sum_{T} \mathcal{O}_{(ij)} + \sum_{A} \mathcal{O}_{[ij]}$$

イロト 不得 トイヨト イヨト 二日

## Sum Rule in O(N) global symmetry

\* Operator product of two primary scalar fields decomposed by

$$\phi_i \times \phi_j \sim \sum_{S} \delta_{ij} \mathcal{O} + \sum_{T} \mathcal{O}_{(ij)} + \sum_{A} \mathcal{O}_{[ij]}$$

\* Reflecting this structure, sum rule is promoted :

[Kos, Poland, Simmons-Duffin, 13]

(人間) (人) (人) (人) (人) (人)

$$\sum_{S} c_{\Delta,l} V_{S,\Delta,l} + \sum_{T} c_{\Delta,l} V_{T,\Delta,l} + \sum_{A} c_{\Delta,l} V_{A,\Delta,l} = 0$$

where

$$\begin{split} \mathcal{V}_{\mathcal{S},\Delta,I} &= \begin{pmatrix} 0 \\ \mathcal{F}_{\Delta,I}^{-}(u,v) \\ \mathcal{F}_{\Delta,I}^{+}(u,v) \end{pmatrix}, \quad \mathcal{V}_{\mathcal{T},\Delta,I} \begin{pmatrix} \mathcal{F}_{\Delta,I}^{-}(u,v) \\ (1-\frac{2}{N})\mathcal{F}_{\Delta,I}^{-}(u,v) \\ -(1+\frac{2}{N})\mathcal{F}_{\Delta,I}^{+}(u,v) \end{pmatrix}, \quad \mathcal{V}_{A,\Delta,I} &= \begin{pmatrix} -\mathcal{F}_{\Delta,I}^{-}(u,v) \\ \mathcal{F}_{\Delta,I}^{-}(u,v) \\ -\mathcal{F}_{\Delta,I}^{+}(u,v) \end{pmatrix} \\ \mathcal{F}^{\pm} &\equiv v^{\Delta} \mathcal{G}_{\Delta,I}(u,v) \pm u^{\Delta} \mathcal{G}_{\Delta,I}(v,u) \end{split}$$

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

# Bootstrapping 3D O(N) theory

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

<ロ> <部> < 部> < き> < き> < き</p>

# Bootstrapping 3D O(N) theory

\* Bootstrapping 3D O(N) theory shows

[Kos, Poland, Simmons-Duffin 13]



Image: A image: A

-

э

# Bootstrapping 3D O(N) theory

\* Bootstrapping 3D O(N) theory shows

[Kos, Poland, Simmons-Duffin 13]



 $\star$  The boundary shows kink again, whose location agrees to the critical exponents  $\eta,\nu.$ 

- 4 同 ト 4 ヨ ト 4 ヨ

## Bootstrapping 3D O(N) theory, Mixed correlator

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

◆□ > ◆□ > ◆豆 > ◆豆 > ̄豆 = つへで

## Bootstrapping 3D O(N) theory, Mixed correlator



Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

### 5-dimensional Renormalization Group Flow

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point
\* At near free theory, quartic interaction is *irrelevant deformation*. Therefore, IR theory corresponds to the free theory. The situation is differ from 3-dimensional renormalization group flow.

-

\* At near free theory, quartic interaction is *irrelevant deformation*. Therefore, IR theory corresponds to the free theory. The situation is differ from 3-dimensional renormalization group flow.



直 ト イヨ ト イヨ ト

\* At near free theory, quartic interaction is *irrelevant deformation*. Therefore, IR theory corresponds to the free theory. The situation is differ from 3-dimensional renormalization group flow.



→ < Ξ → <</p>

\* At near free theory, quartic interaction is *irrelevant deformation*. Therefore, IR theory corresponds to the free theory. The situation is differ from 3-dimensional renormalization group flow.



\* Extension of  $\varepsilon$  expansion into D = 5 corresponds to  $\varepsilon = -1$ . However, this causes negative coupling at fixed point, which means unstable fixed point.

▲ □ ▶ ▲ □ ▶ ▲ □ ▶

\* At near free theory, quartic interaction is *irrelevant deformation*. Therefore, IR theory corresponds to the free theory. The situation is differ from 3-dimensional renormalization group flow.



\* Extension of  $\varepsilon$  expansion into D = 5 corresponds to  $\varepsilon = -1$ . However, this causes negative coupling at fixed point, which means unstable fixed point.

 $\star$  Mean Field Theory expect only free theory at above upper critical dimension( $D_c = 4$ ).

▲ 同 ▶ ▲ 国 ▶ ▲ 国 ▶ …

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

æ

\* The bulk theory of O(N) symmetric theory conjectured to be higherspin theory in  $AdS_{D+1}$ . [Klebanov, Polyakov 02]

$$\left\langle exp\left(\int d^{D}x \ h_{0}^{(\mu_{1}\mu_{2}\cdots\mu_{s})} \qquad \underbrace{J_{(\mu_{1}\mu_{2}\cdots\mu_{s})}}_{(\mu_{1}\mu_{2}\cdots\mu_{s})}\right)\right\rangle = e^{W[h_{0}]}$$

O(N) singlet conserved current

◆□ ▶ ◆□ ▶ ◆ □ ▶ ◆ □ ● ● ● ● ● ●

\* The bulk theory of O(N) symmetric theory conjectured to be higherspin theory in  $AdS_{D+1}$ . [Klebanov, Polyakov 02]

$$\left\langle exp\left(\int d^{D}x \ h_{0}^{(\mu_{1}\mu_{2}\cdots\mu_{s})} \underbrace{J_{(\mu_{1}\mu_{2}\cdots\mu_{s})}}_{O(N) \text{ singlet conserved current}}\right)\right\rangle = e^{W[h_{0}]}$$

 $\star$  The mass dimension of spin-0 field is given by

$$\Delta_{\pm}=rac{D}{2}\pm\sqrt{rac{D^2}{4}+m^2L^2}$$

\* The bulk theory of O(N) symmetric theory conjectured to be higherspin theory in  $AdS_{D+1}$ . [Klebanov, Polyakov 02]

$$\left\langle exp\left(\int d^{D}x \ h_{0}^{(\mu_{1}\mu_{2}\cdots\mu_{s})} \underbrace{J_{(\mu_{1}\mu_{2}\cdots\mu_{s})}}_{O(N) \text{ singlet conserved current}}\right)\right\rangle = e^{W[h_{0}]}$$

 $\star$  The mass dimension of spin-0 field is given by

$$\Delta_{\pm} = \frac{D}{2} \pm \sqrt{\frac{D^2}{4} + m^2 L^2}$$

\* At free theory, mass dimension of scalar current is D-2. Consistent AdS radius is  $L^2 = \frac{4-2D}{m^2}$ .

◆□ ▶ ◆□ ▶ ◆ □ ▶ ◆ □ ● ● ● ● ● ●

\* The bulk theory of O(N) symmetric theory conjectured to be higherspin theory in  $AdS_{D+1}$ . [Klebanov, Polyakov 02]

 $\left\langle exp\left(\int d^{D}x \ h_{0}^{(\mu_{1}\mu_{2}\cdots\mu_{s})} \underbrace{J_{(\mu_{1}\mu_{2}\cdots\mu_{s})}}_{O(N) \text{ singlet conserved current}}\right)\right\rangle = e^{W[h_{0}]}$ 

 $\star$  The mass dimension of spin-0 field is given by

$$\Delta_{\pm}=rac{D}{2}\pm\sqrt{rac{D^2}{4}+m^2L^2}$$

\* At free theory, mass dimension of scalar current is D - 2. Consistent AdS radius is  $L^2 = \frac{4-2D}{m^2}$ .

\* We have two solutions  $\Delta_+ = D - 2$  and  $\Delta_- = 2$ , which indicates UV/IR fixed point of *D*-dimensional O(N) theory. This is consistent with unitary bound as far as 2 < D < 6.

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

イロン イロン イヨン イヨン

3

 $\star$  Hubbard-Stratonovich transformation is applied for large-N limit.

$$S = \int d^{D}x \Big(\frac{1}{2}(\partial \phi^{i})^{2} + \frac{1}{2}\sigma \phi^{i}\phi^{i} - \frac{\sigma^{2}}{4\lambda}\Big)$$

イロト 不得 とうせい かほとう ほ

 $\star$  Hubbard-Stratonovich transformation is applied for large-N limit.

$$S = \int d^{D}x \left(\frac{1}{2}(\partial \phi^{i})^{2} + \frac{1}{2}\sigma \phi^{i} \phi^{i} - \frac{\sigma^{2}}{4\lambda}\right)$$

\* Leading  $\frac{1}{N}$  1-loop computation result is

[Fei, Giombi, Klebanov, 14]

$$\begin{split} \Delta_{\phi} &= \frac{D}{2} - 1 + \frac{1}{N}\eta_1, \quad \Delta_{\sigma} = 2 + \frac{1}{N}\frac{4(D-1)(D-2)}{D-4}\eta_1 \\ \eta_1 &= \frac{2^{D-3}(D-4)\Gamma(\frac{D-1}{2})\text{sin}(\frac{\pi D}{2})}{\pi^{3/2}\Gamma(\frac{D}{2}+1)} \end{split}$$

ヘロン 人間 とくほと 人ほとう

3

 $\star$  Hubbard-Stratonovich transformation is applied for large-N limit.

$$S = \int d^{D}x \left(\frac{1}{2}(\partial \phi^{i})^{2} + \frac{1}{2}\sigma \phi^{i} \phi^{i} - \frac{\sigma^{2}}{4\lambda}\right)$$

\* Leading  $\frac{1}{N}$  1-loop computation result is

[Fei, Giombi, Klebanov, 14]

$$\begin{split} \Delta_{\phi} &= \frac{D}{2} - 1 + \frac{1}{N} \eta_1, \quad \Delta_{\sigma} = 2 + \frac{1}{N} \frac{4(D-1)(D-2)}{D-4} \eta_1 \\ \eta_1 &= \frac{2^{D-3}(D-4)\Gamma(\frac{D-1}{2})\text{sin}(\frac{\pi D}{2})}{\pi^{3/2}\Gamma(\frac{D}{2}+1)} \end{split}$$

This computation is valid for arbitrary D. Among them, not only 2 < D < 4, but also 4 < D < 6, anomalous dimension is positive. (Consistence with unitary bound.)

◆□ ▶ ◆□ ▶ ◆ □ ▶ ◆ □ ● ● ● ● ● ●

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

イロン イロン イヨン イヨン

3

 $\star$  At 5-dimension and large N limit, anomalous dimension of  $\phi$  and  $\sigma\equiv\lambda\phi^i\phi^i$  are

$$\Delta_{\phi} = \frac{3}{2} + \underbrace{\frac{0.216152}{N} - \frac{4.342}{N^2} - \frac{121.673}{N^3} + \cdots}_{\begin{array}{c} \frac{\eta}{2} \end{array}}_{\Delta_{\sigma}} \\ \Delta_{\sigma} = 2 + \underbrace{\frac{10.3753}{N} + \frac{206.542}{N^2} + \cdots}_{\begin{array}{c} 3 - \frac{1}{1} \end{array}}_{\begin{array}{c} \end{array}}_{\begin{array}{c} \end{array}}$$

▲ロ▶ ▲冊▶ ▲ヨ▶ ▲ヨ▶ ヨ のの⊙

 $\star$  At 5-dimension and large N limit, anomalous dimension of  $\phi$  and  $\sigma\equiv\lambda\phi^i\phi^i$  are

$$\begin{split} \Delta_{\phi} &= \frac{3}{2} + \underbrace{\frac{0.216152}{N} - \frac{4.342}{N^2} - \frac{121.673}{N^3} + \cdots}_{\begin{array}{c} \frac{\eta}{2} \end{array}} \\ \Delta_{\sigma} &= 2 + \underbrace{\frac{10.3753}{N} + \frac{206.542}{N^2} + \cdots}_{\begin{array}{c} 3 - \frac{1}{\nu} \end{array}} \end{split}$$

 $\star$  When N below 35, anomalous dimension of  $\phi$  became negative. Hence, unitarity violated.

▲ロ▶ ▲冊▶ ▲ヨ▶ ▲ヨ▶ ヨ のの⊙

 $\star$  At 5-dimension and large N limit, anomalous dimension of  $\phi$  and  $\sigma\equiv\lambda\phi^i\phi^i$  are

$$\begin{split} \Delta_{\phi} &= \frac{3}{2} + \underbrace{\frac{0.216152}{N} - \frac{4.342}{N^2} - \frac{121.673}{N^3} + \cdots}_{\begin{array}{c} \frac{\eta}{2} \end{array}} \\ \Delta_{\sigma} &= 2 + \underbrace{\frac{10.3753}{N} + \frac{206.542}{N^2} + \cdots}_{\begin{array}{c} 3 - \frac{1}{\nu} \end{array}} \end{split}$$

 $\star$  When N below 35, anomalous dimension of  $\phi$  became negative. Hence, unitarity violated.

Question) If we knows exact  $\Delta_{\phi}$ , then unitarity would be violated by specific value of *N*? or not?

◆□ ▶ ◆□ ▶ ◆ □ ▶ ◆ □ ● ● ● ● ● ●

 $\star$  At 5-dimension and large N limit, anomalous dimension of  $\phi$  and  $\sigma\equiv\lambda\phi^i\phi^i$  are

$$\begin{split} \Delta_{\phi} &= \frac{3}{2} + \underbrace{\frac{0.216152}{N} - \frac{4.342}{N^2} - \frac{121.673}{N^3} + \cdots}_{\begin{array}{c} \frac{\eta}{2} \end{array}} \\ \Delta_{\sigma} &= 2 + \underbrace{\frac{10.3753}{N} + \frac{206.542}{N^2} + \cdots}_{\begin{array}{c} 3 - \frac{1}{\nu} \end{array}} \end{split}$$

 $\star$  When N below 35, anomalous dimension of  $\phi$  became negative. Hence, unitarity violated.

Question) If we knows exact  $\Delta_{\phi}$ , then unitarity would be violated by specific value of *N*? or not?

 $\star$  Likewise 3-dimensional IR fixed point, can we see appearance of UV fixed point via conformal bootstrap?

イロト 不得 トイヨト イヨト 二日

# Bootstrapping 5-dimensional O(N) CFT

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

<ロ> <同> <同> < 同> < 同>

3

# Bootstrapping 5-dimensional O(N) CFT

\* The straightforward application of numerical conformal bootstrap into 5-dimensional O(N = 500) theory gives



Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

# Bootstrapping 5-dimensional O(N) CFT

 $\star$  The straightforward application of numerical conformal bootstrap into 5-dimensional O(N = 500) theory gives



 $\star$  This does not specify UV fixed point on its boundary as 3-dimensional bootstrap does.

< ∃ >

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

э



Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

<ロ> <部> < 部> < き> < き> < き</p>

★ Black line → Real Spectrum of 3D IR conformal field theory. ( $\epsilon = 1.41275(25), \ \epsilon' = 3.84(4), \ \epsilon'' = 4.67(11) \cdots$ ) [Gliozzi, Rago 14]



▲ロ▶ ▲冊▶ ▲ヨ▶ ▲ヨ▶ ヨ のの⊙

★ Black line → Real Spectrum of 3D IR conformal field theory. ( $\epsilon = 1.41275(25), \epsilon' = 3.84(4), \epsilon'' = 4.67(11)\cdots$ ) [Gliozzi, Rago 14] Orange Area → Operators in the Sum Rule(over the  $\Delta_{\epsilon}$ ).



◆□ ▶ ◆□ ▶ ◆ □ ▶ ◆ □ ● ● ● ● ● ●

\* Black line  $\rightarrow$  Real Spectrum of 3D IR conformal field theory. ( $\epsilon = 1.41275(25)$ ,  $\epsilon' = 3.84(4)$ ,  $\epsilon'' = 4.67(11) \cdots$ ) [Gliozzi, Rago 14] Orange Area  $\rightarrow$  Operators in the Sum Rule(over the  $\Delta_{\epsilon}$ ). If  $\Delta_{\epsilon}$  exceeds  $\epsilon = 1.41275(25)$ , 3D Ising theory outted.



◆□ > ◆□ > ◆臣 > ◆臣 > 善臣 - のへで

\* Black line  $\rightarrow$  Real Spectrum of 3D IR conformal field theory. ( $\epsilon = 1.41275(25)$ ,  $\epsilon' = 3.84(4)$ ,  $\epsilon'' = 4.67(11) \cdots$ ) [Gliozzi, Rago 14] Orange Area  $\rightarrow$  Operators in the Sum Rule(over the  $\Delta_{\epsilon}$ ). If  $\Delta_{\epsilon}$  exceeds  $\epsilon = 1.41275(25)$ , 3D Ising theory outted.





Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

\* Black line  $\rightarrow$  Real Spectrum of 3D IR conformal field theory. ( $\epsilon = 1.41275(25), \epsilon' = 3.84(4), \epsilon'' = 4.67(11) \cdots$ ) [Gliozzi, Rago 14] Orange Area  $\rightarrow$  Operators in the Sum Rule(over the  $\Delta_{\epsilon}$ ). If  $\Delta_{\epsilon}$  exceeds  $\epsilon = 1.41275(25), 3D$  Ising theory outted.



 $\star$  Modify the spectrum assumptions in the Sum Rule.



\* Black line  $\rightarrow$  Real Spectrum of 3D IR conformal field theory. ( $\epsilon = 1.41275(25)$ ,  $\epsilon' = 3.84(4)$ ,  $\epsilon'' = 4.67(11) \cdots$ ) [Gliozzi, Rago 14] Orange Area  $\rightarrow$  Operators in the Sum Rule(over the  $\Delta_{\epsilon}$ ). If  $\Delta_{\epsilon}$  exceeds  $\epsilon = 1.41275(25)$ , 3D Ising theory outted.



 $\star$  Modify the spectrum assumptions in the Sum Rule. Additionally, we introduced the new scale  $\Delta_{\rm gap}.$ 



\* Black line  $\rightarrow$  Real Spectrum of 3D IR conformal field theory. ( $\epsilon = 1.41275(25)$ ,  $\epsilon' = 3.84(4)$ ,  $\epsilon'' = 4.67(11) \cdots$ ) [Gliozzi, Rago 14] Orange Area  $\rightarrow$  Operators in the Sum Rule(over the  $\Delta_{\epsilon}$ ). If  $\Delta_{\epsilon}$  exceeds  $\epsilon = 1.41275(25)$ , 3D Ising theory outted.



\* Modify the spectrum assumptions in the Sum Rule. Additionally, we introduced the new scale  $\Delta_{gap}$ . This setup is the stronger spectrum assupptions.



Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

ヘロト ヘ部ト ヘヨト ヘヨト

3

Basic Strategy

<ロ> <同> <同> < 同> < 同>

э

Basic Strategy

Three inputs

 $\Delta_{\phi}$ ,  $\Delta_{\min}$ ,  $\Delta_{gap}$ 

Fixed  $\Delta_{\rm gap}$ 

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

・ロト ・回ト ・ヨト ・ヨト

3

#### Basic Strategy

Three inputs

 $\Delta_{\phi}$ ,  $\Delta_{\min}$ ,  $\Delta_{gap}$ 

Fixed  $\Delta_{\rm gap}$ 

Is it consistent with Unitarity and Crossing symmetry ?

(By Linear Programming or Semi-definite Programming)

→ < = →

#### Basic Strategy

Three inputs

 $\Delta_{\phi}$ ,  $\Delta_{\min}$ ,  $\Delta_{gap}$ 

Fixed  $\Delta_{\rm gap}$ 

Is it consistent with Unitarity and Crossing symmetry ?

(By Linear Programming or Semi-definite Programming)



- ∢ ≣ ▶
## Conformal Bootstrap in nutshell

#### Basic Strategy

Three inputs

 $\Delta_{\phi}$ ,  $\Delta_{\min}$ ,  $\Delta_{gap}$ 

Fixed  $\Delta_{\rm gap}$ 

Is it consistent with Unitarity and Crossing symmetry ?

(By Linear Programming or Semi-definite Programming)



- ∢ ≣ ▶

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

<ロ> <部> < 部> < き> < き> < き</p>

 $\star$  The 2-parameter method carved out more region indeed. For N=500 and  $\Delta_{\rm gap}=7.0,$  our result shows

- ◆母 ▶ ◆臣 ▶ ◆臣 ▶ ○臣 ○ のへ(?)

 $\star$  The 2-parameter method carved out more region indeed. For N = 500 and  $\Delta_{\rm gap}$  = 7.0, our result shows



Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

 $\star$  The 2-parameter method carved out more region indeed. For N=500 and  $\Delta_{\rm gap}=7.0,$  our result shows



\* The location of low tip agrees to the result of  $\frac{1}{N}$  expansion via Hubbard-Stratonovich transformation.

- ● ● ●

## Bootstrapping 5-dimensional CFT for various N

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

## Bootstrapping 5-dimensional CFT for various N

\* We played similar game for various high value of N = 500, 1000, 2000. The location of low tip corresponds to  $\frac{1}{N}$  expansion results.

・ 同 ト ・ ヨ ト ・ ヨ ト

-

### Bootstrapping 5-dimensional CFT for various N

\* We played similar game for various high value of N = 500, 1000, 2000. The location of low tip corresponds to  $\frac{1}{N}$  expansion results.



Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

▶ ∢ ≣ ▶

### Conclusion and Outlook

Jin-Beom Bae Searching for 5-dimensional Nontrivial UV Fixed Point

\* UV fixed point of 5-dimensional O(N) theory is predictable from Holographic principle. The critical exponents of the UV theory is computable at large-N limit.

▲ロ▶ ▲冊▶ ▲ヨ▶ ▲ヨ▶ ヨ のの⊙

\* UV fixed point of 5-dimensional O(N) theory is predictable from Holographic principle. The critical exponents of the UV theory is computable at large-N limit.

 $\star$  With more restrictive assumption on spectrum, we show that conformal bootstrap also utilized to figure out IR /UV fixed point for 5-dimensional theory.

◆□ ▶ ◆□ ▶ ◆ □ ▶ ◆ □ ● ● ● ● ● ●

\* UV fixed point of 5-dimensional O(N) theory is predictable from Holographic principle. The critical exponents of the UV theory is computable at large-N limit.

 $\star$  With more restrictive assumption on spectrum, we show that conformal bootstrap also utilized to figure out IR /UV fixed point for 5-dimensional theory.

\* Mixed Bootstrap Program on 5-dimensional O(N)?

◆□ ▶ ◆□ ▶ ◆ □ ▶ ◆ □ ● ● ● ● ● ●