

Where are we heading?

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Two half-talks

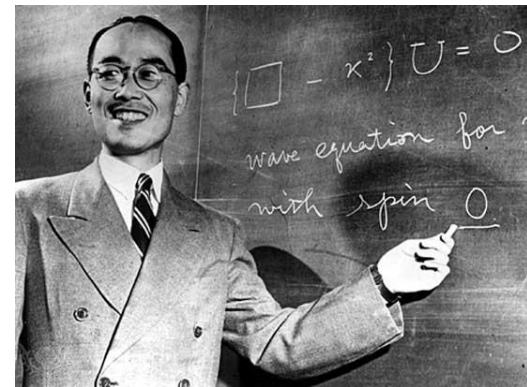
- A brief, broad brush status report of particle physics and what the future could be like
- The role of symmetries in physics and how it is changing

The Standard Model is extremely successful

- Many experimental tests of the model
- No known discrepancy between theory and experiment
- Unprecedented accuracy

Open problems with the SM

- Where did the spectrum of particles come from?
 - Gauge group
 - Quarks and leptons quantum numbers
 - Generations
- What determines the electroweak scale (Higgs, W , Z masses)?
- Where did the Yukawa couplings come from?
 - Lead to fermion masses
 - Quarks mixing angles
 - CP violation
 - ...



Open problems with the SM

- Hierarchies (more below)
 - Hierarchy of quark and lepton masses (they span 5 orders of magnitudes)
 - Pattern of CKM angles (why are they small?)
 - Strong CP problem ($\theta_{\text{QCD}} < 10^{-11}$)
 - Electroweak scale and Higgs mass
- Dark matter
- Neutrino masses and mixing angles (not small). Beyond the SM, but reflects physics at much higher energies. Will not discuss today.



More data soon

- The LHC is operating at higher energy – 6.5 TeV per beam and higher luminosity
- Data from other experiments (precision measurements, dark matter searches, cosmology, astrophysics)

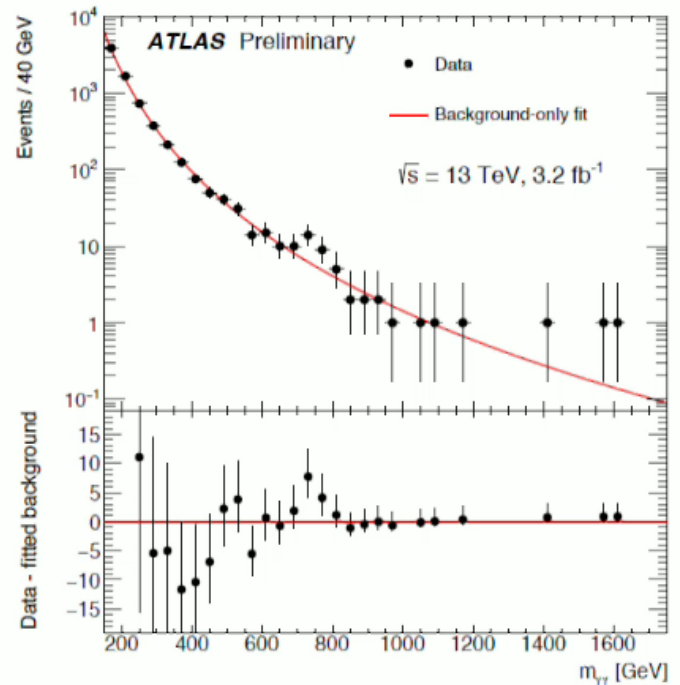
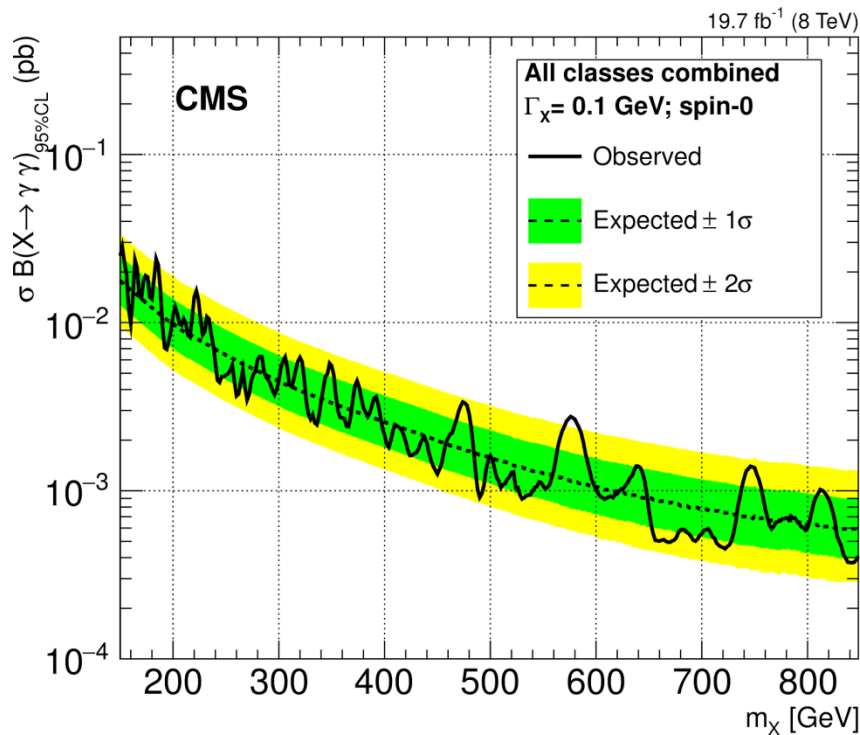


Options for the near future

- Nothing beyond the Standard Model with its single Higgs
- Going beyond the Standard Model
 - Discrepancies in the Higgs production rate and/or the various decay modes branching ratios
 - Small discrepancies in other processes
 - Additional particles

First sign of physics beyond the SM?

Di-photon resonance at 750GeV (reassurance in Moriond)?



Possible conceptual extensions of the SM

- Supersymmetry – it is weakly coupled
- Strong coupling dynamics for electroweak breaking – Technicolor, warped extra dimensions (i.e. strongly coupled field theory that is dual to a weakly coupled gravitational theory)
- Something else we have not yet thought about

One line status report (with many caveats)

The measured Higgs mass $\sim 125\text{GeV}$ is uncomfortably high for (minimal) supersymmetry and uncomfortably low for strong dynamics.

More details below

Hierarchy problem/Naturalness

- Dimensional analysis usually works – observables are given typically by the scale of the problem times a number of order one.
- Dirac's large numbers problem: Why is the proton so much lighter than the Planck scale?

$$M_p \ll M_{Planck}$$



Hierarchy problem/Naturalness

This particular problem is now understood as following from asymptotic freedom

$$\frac{M_p}{M_{Planck}} \sim e^{-\frac{a}{g^2}} \ll 1$$

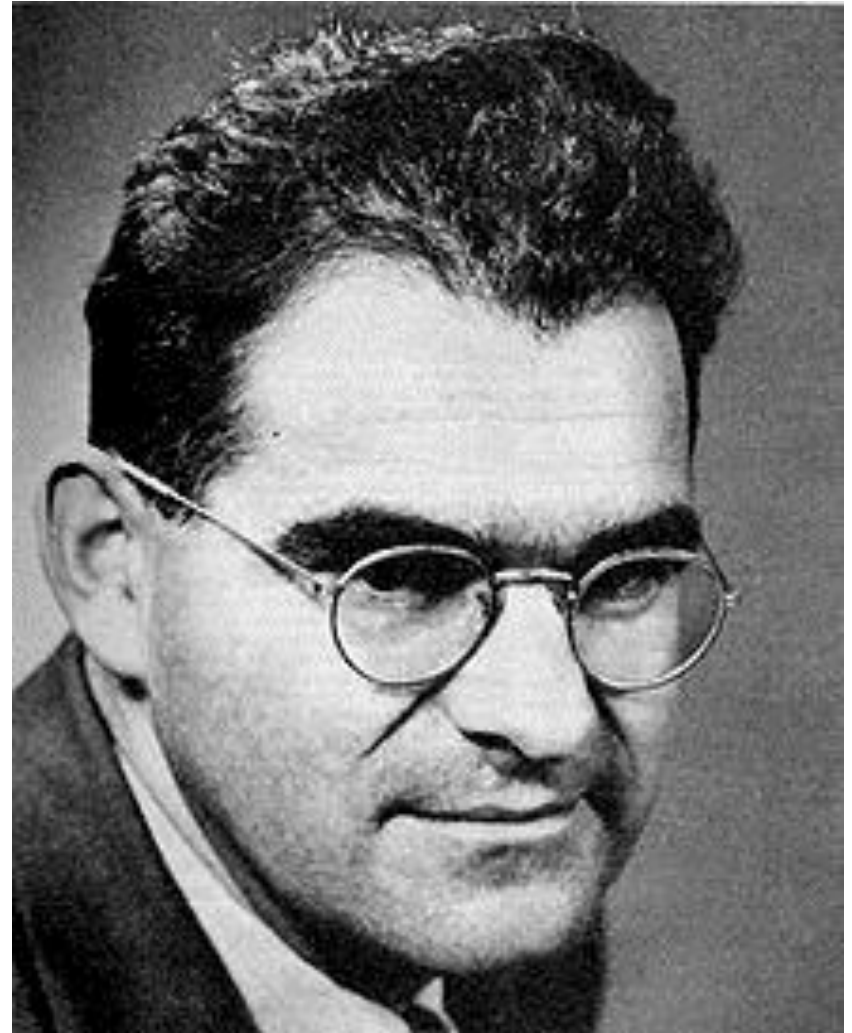
Its newer version involves the electroweak scale

$$M_W, M_Z, M_H \ll M_{Planck}$$

More generally, the intuitive hierarchy problem: where did very small dimensionless numbers come from?

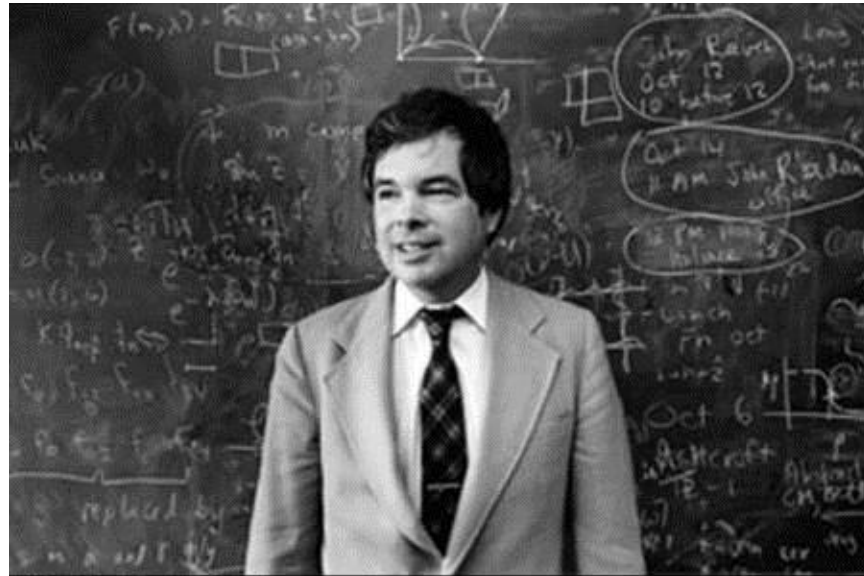
Hierarchy problem/Naturalness

We should avoid
quantum field theories
with quadratic
divergences.
Logarithmic divergences
are OK.
(Weisskopf)



Hierarchy problem/Naturalness

- Small scalar masses are unnatural (Wilson)
 - It is like being very close to a phase transition
 - Scalar mass terms suffer from large quadratic divergences



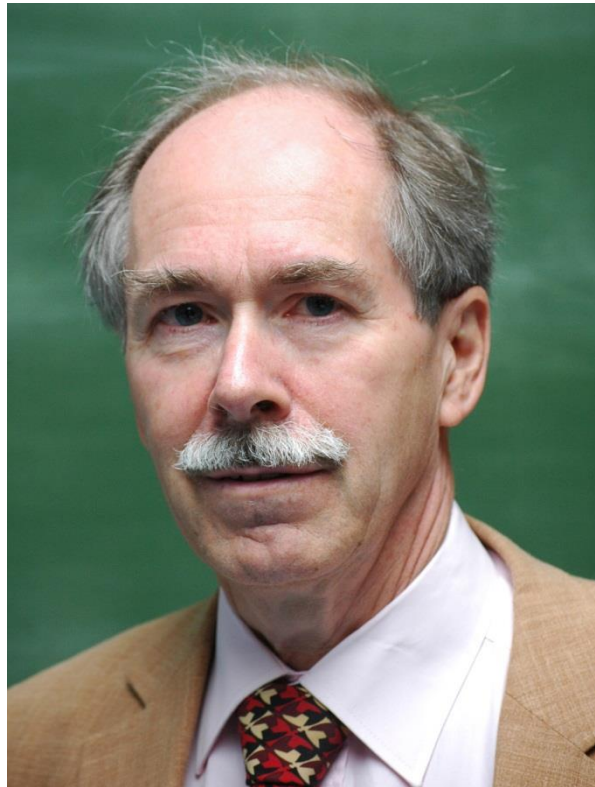
Hierarchy problem/Naturalness

- Alternatively, they are extremely sensitive to small changes of the parameters of the theory at high energy – delicate unnatural cancellations between high energy parameters (Weinberg)



Hierarchy problem/Naturalness

- A dimensionless parameter is naturally small only if the theory is more symmetric when it is exactly zero ('t Hooft) – technical naturalness.



Hierarchy problem/Naturalness

- The intuitive problem
 - Where did small numbers come from?
 - Why doesn't dimensional analysis work? All dimensionless numbers should be of order one.
 - Can postpone the solution to higher energies
- The technical problem
 - Even if in some approximation we find a hierarchy, higher order corrections can destabilize it.
 - Quantum fluctuations tend to restore dimensional analysis.
 - Must solve at the same scale

Hierarchy problem/Naturalness

- Hierarchy in fermion masses and mixing angles
 - Only the intuitive problem – enhanced symmetry when they vanish.
 - The origin (explanation) can arise from extremely high energy physics .
- Strong CP problem
 - Both the intuitive and the technical issue – no enhanced symmetry when $\theta_{\text{QCD}} = 0$
 - Only logarithmic divergence (with small coefficient)
 - The explanation must involve low energy physics.
Axions? $m_{up} = 0$? Something else?

Hierarchy problem/Naturalness

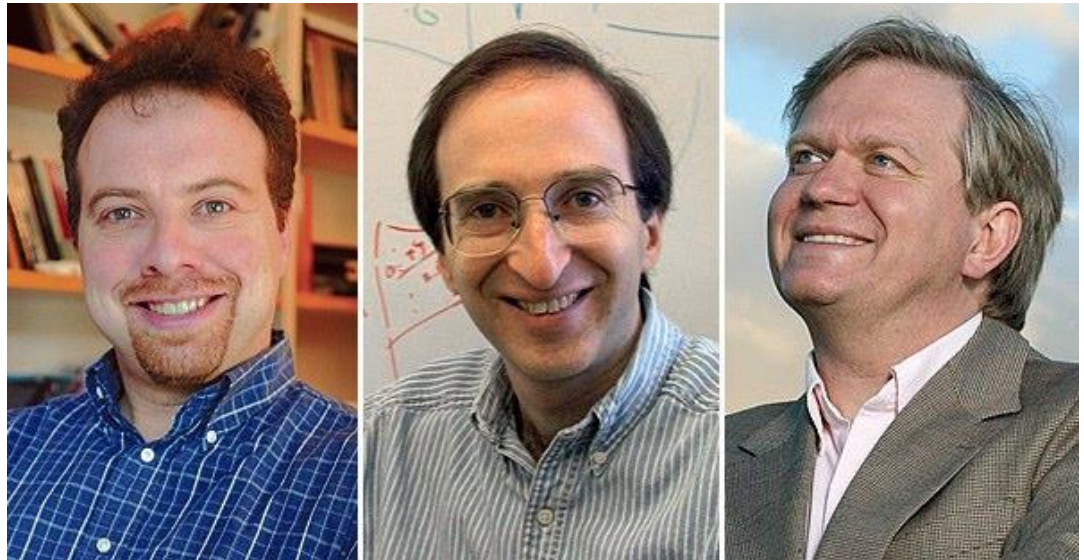
- Higgs mass and the electroweak scale
 - Quadratic divergences – sensitivity to high energy physics
 - No symmetry is restored when they vanish. (The $SU(2) \times U(1)$ symmetry is always present but might be spontaneously broken.)
 - Both the intuitive and the technical problems
 - Hence, expect to solve it at low energies

The biggest hierarchy problem



The biggest hierarchy problem

- The cosmological constant is quartically divergent – it is fine tuned to 120 decimal points.
- We used to think that it vanishes. We did not have a mechanism explaining why it is zero, but we hoped that one day we would find a principle setting it to zero.
- Now that we know that it is nonzero, our naturalness prejudice is being shaken.



Natural solutions to the Higgs hierarchy problem: Technicolor

- Technicolor is basically dead
 - Precision measurements (the S and T parameters) and the measured m_H disfavor it.
 - More intuitively, the measured mass of the Higgs tells us that it is weakly coupled. Strong coupling solutions like Technicolor tend to lead to a strongly coupled Higgs.
 - More sophisticated composite Higgs models could work, but they are somewhat complicated and contrived.

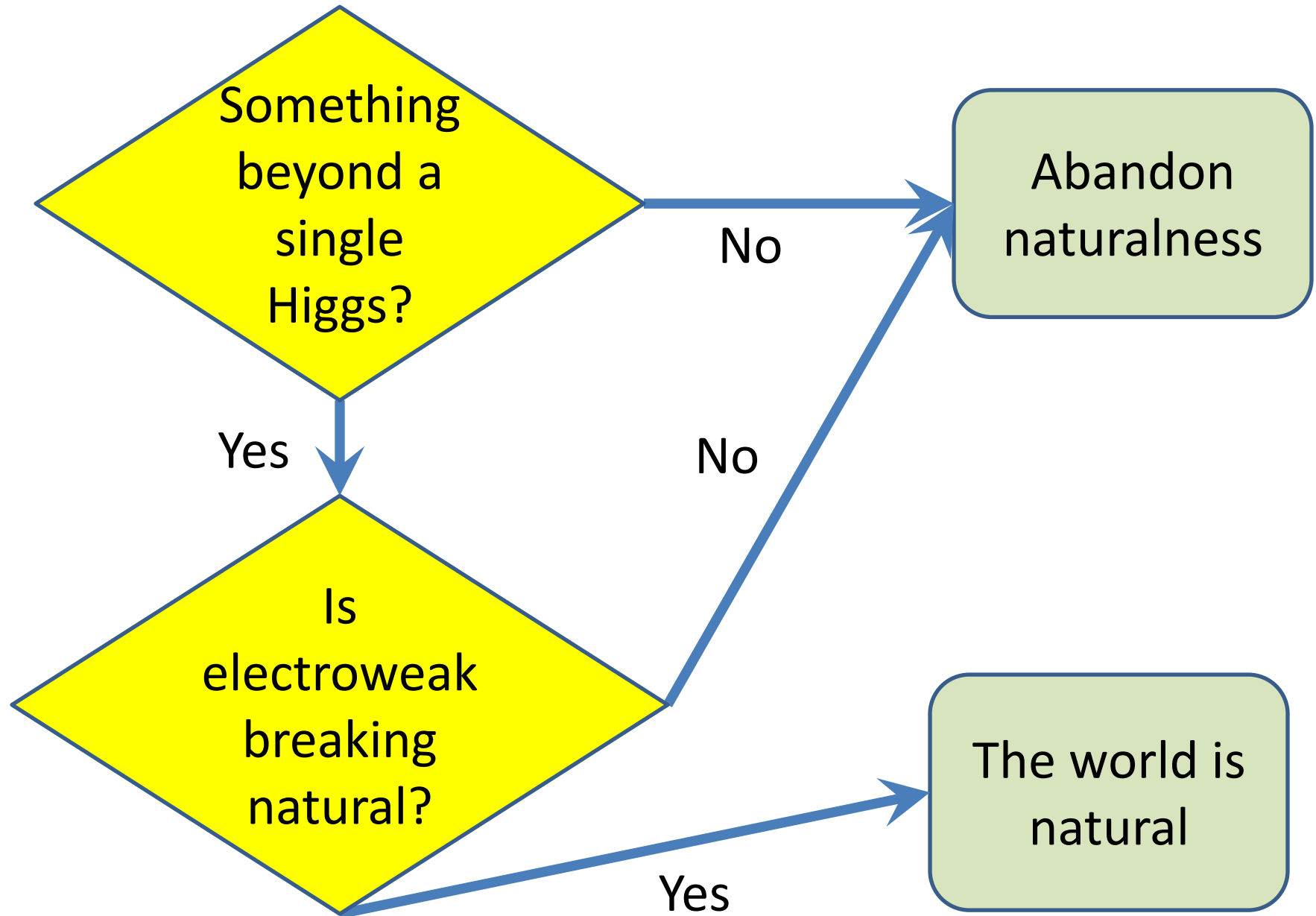
Natural solutions to the Higgs hierarchy problem: Supersymmetry

It is hard to make SUSY fully natural.

In the Minimal SUSY Standard Model the Higgs self-coupling is related to the gauge coupling:

- Classically $m_{Higgs} \leq m_Z$
- Quantum corrections can lift the Higgs mass, but for reaching 125GeV we need
 - heavy stop
 - large A-terms
 - going beyond the minimal model
- Everyone of these is possible, but problematic.

Options for the near future



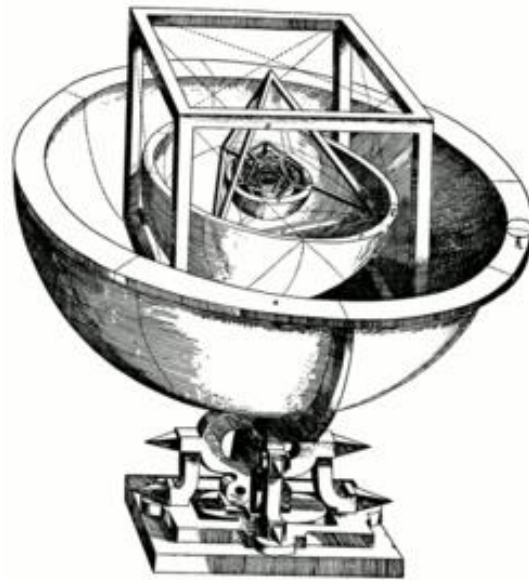
If TeV Physics is unnatural

Leading option: landscape of vacua (and perhaps the A-word)

- The world is much bigger than we think (a multiverse)
- The laws of physics are different in different places – the laws of physics are environmental
- Predicting or explaining the parameters of the SM (e.g. the electron mass) is like predicting the sizes of the orbits of the planets.

A historical reminder

Kepler had a beautiful mathematical explanation of the sizes of the orbits of the planets in terms of the 5 Platonic solids.



This turned out to be the wrong question.

If TeV Physics is unnatural

- Should we attempt to solve other naturalness questions (strong CP, ratios of fermion masses and mixing angles)?
- What will be the right questions to ask and to explore?
- Some might say that we should stop looking for deeper truth at shorter distances. Instead, some or all the parameters are environmental and should not be explained.
- End of reductionism?

Conclusions

The LHC can find:

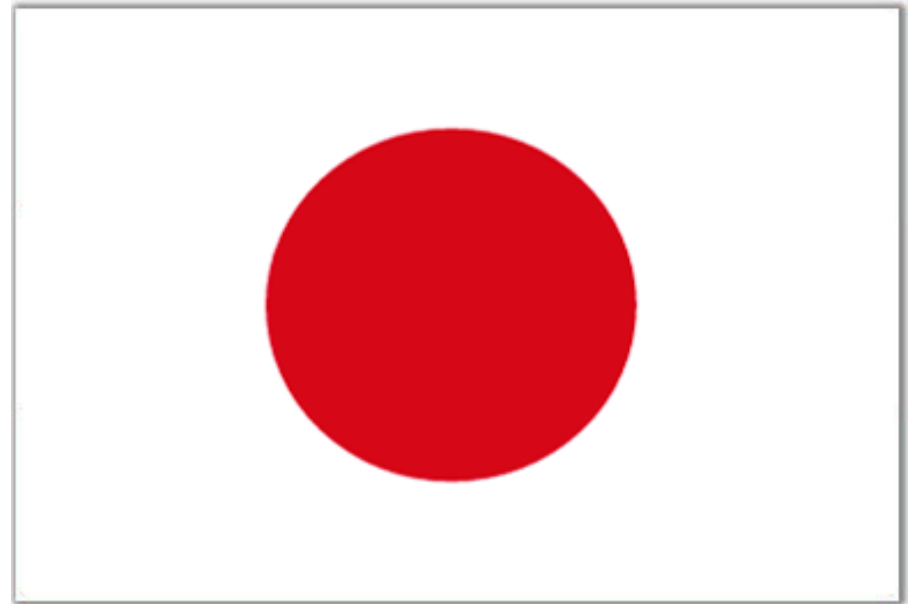
- No discrepancy with the minimal Standard Model
- New physics beyond the minimal Standard Model that does not address the stability of the weak scale
- A natural explanation of the weak scale
 - Supersymmetry
 - Strong dynamics
 - Something we have not yet thought about

Conclusions

All these options are interesting

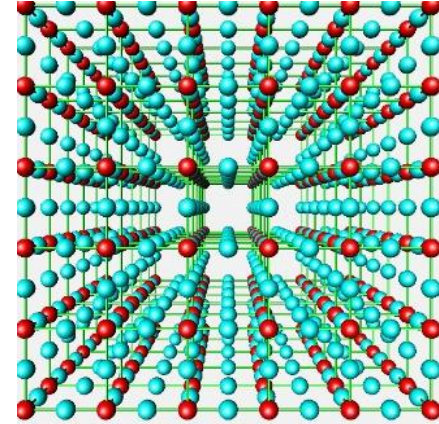
- They give us correct reliable information about Nature.
- They point to a deep physical principle with far reaching philosophical consequences about the Universe. Is our world natural? Is it special? End of reductionism?
- We are in a win-win situation. Every outcome is interesting.

Symmetries



Physicists love symmetries

- Crystallography
- Lorentz
- Flavor $SU(2)$, $SU(3)$
 - Consequence of light quarks. Quarks are deep.
- Color
 - Gauge symmetry is deep



Gauge symmetry is deep

- Largest symmetry (a group for each point in spacetime)
- Useful in making the theory manifestly Lorentz invariant, unitary and local (and hence causal)
- Appears in
 - Maxwell theory, the Standard Model
 - General Relativity
 - Many condensed matter systems
 - Deep mathematics (fiber bundles)

But

- Because of Gauss law the Hilbert space is gauge invariant. (More precisely, it is invariant under small gauge transformation; large gauge transformations are central.)
- Hence: **gauge symmetry is not a symmetry.**
 - It does not act on anything.
- A better phrase is **gauge redundancy.**

Gauge symmetries cannot break

- Not a symmetry and hence cannot break
- For spontaneous symmetry breaking we need an infinite number of degrees of freedom transforming under the symmetry. Not here.



- This is the deep reason there is no massless Nambu-Goldstone boson when gauge symmetries are “broken.”

Gauge symmetries cannot break



For weakly coupled systems (e.g. Landau-Ginsburg theory of superconductivity, or the weak interactions) the language of spontaneous gauge symmetry breaking is appropriate and extremely useful [Stueckelberg, Anderson, Brout, Englert, Higgs].

Global symmetries can emerge as accidental symmetries at long distance.

Then they are approximate.

Exact gauge symmetries can be emergent.

Examples of emergent gauge symmetry

- Simple dualities in free theories
 - In $3d$ a compact scalar is dual to Maxwell theory, whose gauge symmetry is **emergent**.
 - In $4d$ Maxwell theory is dual to a magnetic Maxwell theory. Its gauge symmetry is **emergent**.
- Common in condensed matter physics
 - Fractional Hall effect
 - Particle-vortex duality in $2+1$ dimensions
 - ...

Duality in interacting field theories

$N = 4$ supersymmetry

- This is a scale invariant theory characterized by a gauge group G and a complex coupling constant $\tau = \frac{\theta}{2\pi} + \frac{4\pi}{g^2} i$ for each factor in G .
- For simply laced G the theory with τ is the same as with $\tau + 1$ (shift θ by 2π) and the same as with $-1/\tau$ (generating $SL(2, \mathbf{Z})$). The latter maps strong to weak coupling.
- The duality is an exact equivalence of theories.
 - Same spectrum of states
 - Same spectrum of operators
 - Same correlation functions

Duality in interacting field theories

$N = 4$ supersymmetry

- The gauge symmetry of the dual description is emergent!
- Which of the two gauge symmetries is fundamental?
- Which set of gluons is elementary?
- Perhaps neither gauge symmetry is fundamental.
- Notion of “elementary particle” is ill-defined.

Interacting gauge theories

Start at short distance with a gauge group G . Depending on the details we end up at long distance with:

- IR freedom – a free theory based on G (same theory)
- A nontrivial fixed point. Interacting conformal field theory – no notion of particles.
- An approximately free (IR free) theory of bound states
- An empty theory – gap (possibly topological order)

All these options are realized in QCD for various numbers of flavors. (The approximately free theory is a theory of pions.)

Duality in interacting field theories

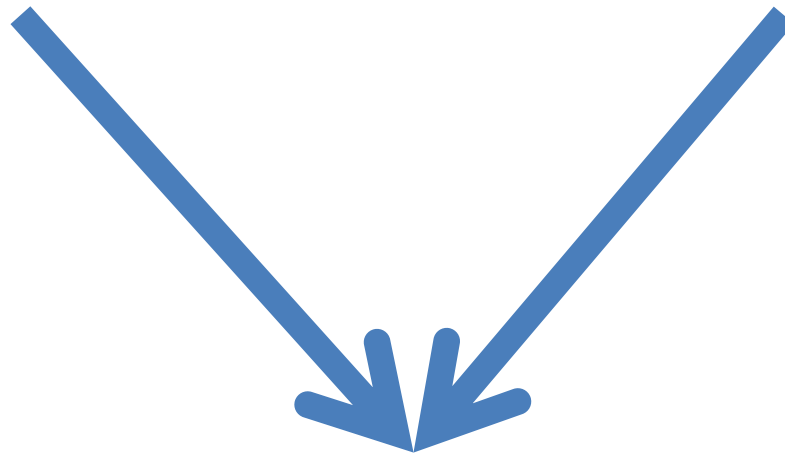
$N = 1$ supersymmetry

Electric theory

G

Magnetic theory

\tilde{G}



Non-trivial IR fixed point

Duality in interacting field theories

$N = 1$ supersymmetry

Another option:

Electric theory

Based on G



Approximately free theory (IR free)

Based on \tilde{G}

Duality in interacting field theories

$N = 1$ supersymmetry

In the **UV** an asymptotically free theory based on G

In the **IR** an IR free theory based on \tilde{G}

At low energies QCD has pions. This theory has a non-Abelian gauge theory.

- The gauge fields of \tilde{G} are composite.
- Their gauge symmetry is emergent.
- There is no ambiguity in the IR gauge symmetry – approximately free massless gauge fields.

Many more examples. Duality and emergent gauge symmetries are ubiquitous.

Emergent general covariance and emergent spacetime

- So far we discussed duality between two field theories
- String-string duality
 - T-duality
 - S-duality
 - U-duality
- String-fields duality
 - Matrix models for low dimensional string theories
 - BFSS M(atrrix) model
 - AdS/CFT
 - More generally gauge-gravity duality

Conclusions

- Gauge symmetries can come and go.
- They can emerge.
- It is often convenient to use them to make the description manifestly Lorentz invariant, unitary and local.
 - But there can be different such descriptions.
- Gauge symmetry is not fundamental.
- Look for a formulation of field theory that makes the duality manifest.
 - We should not be surprised by duality!