Tufts Mathematics—String Theory



Increased complexity by surgery

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Increasing the complexity of free particle motion by breaking its mechanical nature

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School of Arts and Sciences CONTINUERS TY CONTINUE AS ITY CONTINUES AND CONTI

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Hadamard and negatively curved surfaces (1898) The set *E* of vectors at a base point whose geodesics are bounded

Les surfaces à courbures opposées et leurs lignes géodésique. J. Math. Pures Appl. 4 (1898): 27–73.

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Picture from Hadamard

The manner in which these sets arise clearly recalls that of the sets encountered by Mr. Poincaré, introduced more explicitly to the subject by Mr. Bendixson, then studied by Mr. Cantor and which, while being perfect, are not *dense* in any interval. Here the angles λ play the role of the intervals called (a_v, b_v) by Mr. Cantor."

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The Maxwell–Boltzmann Ergodic Hypothesis

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James Clerk Maxwell (1831–1879) and Ludwig Boltzmann (1844–1906) aimed to give a rigorous formulation of the kinetic theory of gases and statistical mechanics.





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Boltzmann's H-theorem says that the time and space (ensemble) averages of an observable (a function on the phase space) agree.

Vorlesungen über Gastheorie I & II. Leipzig: Ambrosius Barth, **1896–1898**, or Gesamtausgabe, Graz: Akademische Druck und Verlagsanstalt, 1981, 1.

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Boltzmann's H-theorem assumed the Ergodic Hypothesis:

The trajectory of the point representing the state of the system in phase space passes through every point on the constant-energy hypersurface of the phase space.



Tufts The Quasi-Ergodic Hypothesis

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Poincaré and many physicists doubted the Ergodic Hypothesis since no example satisfying it had been exhibited.

Poincaré: Sur la théorie cinétique des gas. Revue Générale des Sciences pures et appliqueés 5 (1894), 513-521

Paul and Tatiana Ehrenfest proposed the alternative **Quasi-Ergodic Hypothesis**: *The trajectory of the point representing the state of the system in phase space is dense on the constant energy hypersurface of the phase space.* Begriffliche Grundlagen der statistischen Auffassung in der Mechanik.

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In Encyklopaedie der Mathematischen Wissenschaften, Leipzig 1912: Teubner. 4: Art. 32, 1–90.

The quest was on...

Ein mechanisches System mit quasiergodischen Bahnen

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Emil Artin:

Ein mechanisches System mit quasiergodischen Bahnen Abhandlungen aus dem Mathematischen Seminar der Hamburgischen Universität, vol. 3 (1924), pp. 170-175.

"May it be permitted to point to a simple mechanical system with 2 degrees of freedom and quasiergodic orbits upon which the author came in the course of a correspondence with Mr. G. Herglotz...

... From this, one already obtains that the "quasiergodic chains" have the cardinality of the continuum. More! According to results of Mr. Celestyn Burstin almost all numbers ξ have a "quasiergodic" continued-fraction expansion. Therefore, almost all of the geodesic lines through a point of the surface are quasiergodic."

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"Let us remark on the physical realizability. One obtains... the surface of rotation of a tractrix (curve of pursuit) of a string of length 1. It is known to have curvature K = -1, so our half-plane can be partially developed onto this surface...But with that we have our physical realization...Our mechanical system can be interpreted... as the force-free motion of a point particle (the point being constrained to remain on the surface)."

Tufts Free particle motion on...



Tufts Breaking the configuration space...





Foulon–Ding–Geiges–Weinstein– Handel–Thurston surgery



Tufts Agenda

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- Foulon–Ding–Geiges–Weinstein–Handel–Thurston surgery is a contact surgery on the unit tangent bundle of a surface.
- This forces orbit complexity.
- It produces several contact 3-flows.
- From the fiber flow a flow with quadratic or exponential complexity.
- From the geodesic flow a contact structure whose every Reeb flow has positive entropy.
- If the surgered geodesic flow is hyperbolic, then it has increased orbit growth.

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There are manifolds with 2 contact flows, one exponentially complex, one polynomially.

The entire talk is in dimension 3.

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Definition (Contact, Anosov, algebraic)

A 1-form *A* is contact : $\Leftrightarrow A \land dA$ is a volume. Its Reeb field R_A is defined by $R_A \in \ker dA \& A(R_A) = 1$. Its flow $\varphi : \mathbb{R} \times M \to M$ preserves *A* (contact flow)

 $\Phi \text{ is an Anosov flow if } TM = \mathbb{R}X \oplus E^+ \oplus E^-,$ flow unstable stable $\exists C > 0, \ \lambda \in (0,1) \quad \forall t > 0:$

 $||D\varphi^{-t}| \in E^+|| \le C\lambda^t$ and $||D\varphi^t| \in E^-|| \le C\lambda^t$.

Algebraic : \Leftrightarrow finitely covered by the geodesic flow of a surface [or the suspension of a diffeomorphism of the 2-torus]¹.

¹not contact—hence out of scope!



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Anosov flows are abundant (surgeries by Franks–Williams, Christy, Fried, Handel–Thurston, Goodman, Bonatti–Langevin, Barbot, Fenley, Béguin–Bonatti–Yu).

But, experts thought surgeries could not produce contact flows

The method for showing that this surgery yields contact flows was new: It involves a suitable deformation of the contact form and a time-change adapted to this deformation.

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Theorem (Exotic contact Anosov flows)

Geodesic flows of negatively curved surfaces admit smooth Dehn surgeries that produce new (analytic) contact Anosov flows. These surgeries include the Handel–Thurston surgery, and the resulting flow always has the following properties:

I *It acts on a manifold that is not a unit tangent bundle.*

It is not topologically orbit equivalent to an algebraic flow.

It has "exponentially more" closed orbits than its progenitor...

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4 ... even in any one homology class.

"Exotic" properties

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Theorem (Contact Anosov flows on hyperbolic manifolds)

 \exists contact Anosov flows on hyperbolic 3-manifolds². For them³ each closed orbit is isotopic to infinitely many others⁴ and any Reeb flow⁵ with the same contact structure has exponential orbit growth, infinite free homotopy classes, and exponential growth in each homology class.

²If $M \sim K$ is a hyperbolic manifold, *e.g.*, *K* projects to a filling geodesic, then all but finitely many of our Dehn surgeries produce a hyperbolic manifold. ³and for any topologically orbit equivalent contact Anosov flow ⁴Barthelmé–Fenley: Freely homotopic closed orbits are isotopic —so there is no knot theory of closed orbits in a free homotopy class. ⁵ $A(R_A) = 1 \& dA(R_A, \cdot) = 0$

Mechanisms due to Fenley and Barbot

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These flows are \mathbb{R} -covered, *i.e.*, on the universal cover the stable and unstable leaf spaces are \mathbb{R} .

That each orbit lies in exactly one (center-)un/stable leaf defines an embedding of \mathbb{O}^{Φ} in $\mathscr{L}^{u} \times \mathscr{L}^{s}$. Its image \mathbb{O} is either

- $\mathscr{L}^u \times \mathscr{L}^s$ —"product flow" (orbit-equivalent to suspension of toral automorphism), or
- the open set between graphs of homeomorphisms $\alpha, \beta: \mathscr{L}^u \to \mathscr{L}^s$ —"skewed".



 $\mathbb{R}\text{-}covered\ cases\ [Picture\ by\ Fenley]}$

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Mechanisms due to Fenley and Barbot



Periodic orbit growth

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The orbit space of a skewed R-covered flow [Pictures by Barthelmé]

Fenley: If Φ is skewed and *M* is a hyperbolic manifold, then every closed orbit is freely homotopic to infinitely many others. Barthelmé: Actually, isotopic.

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Tufts Handel–Thurston surgery



Tufts Surgery in phase space



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Annulus $\Sigma = S^1 \times (\pi/2 - \epsilon, \pi/2 + \epsilon) \subset S^1 \times S^1$ before surgery

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Annulus $\Sigma = S^1 \times (\pi/2 - \epsilon, \pi/2 + \epsilon) \subset S^1 \times S^1$ after surgery

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 γ : Normal vectors to closed geodesic = knot tangent to the contact structure (Legendrian, $A(\gamma') \equiv 0$) and transverse to both E^- and E^+ .

Smooth annulus Σ : $\gamma \subset \Sigma \pitchfork X$, $\Sigma \pitchfork E^- \oplus E^+$ away from γ

 \exists coordinates (*s*, *w*) on Σ such that $S^1 \ni s =$ parameter for γ , and

on $\Lambda := \bigcup_{t \in (-\eta, \eta)} \varphi^t(\Sigma)$ [*t* = transverse parameter = flow time] $A = dt + w ds, \quad dA = dw \wedge ds, \quad A \wedge dA = dt \wedge dw \wedge ds.$

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Define (1, q)-Dehn surgery $(q \in \mathbb{N})$ by the transition map

$$F \colon S^1 \times (-\epsilon, \epsilon) \to S^1 \times (-\epsilon, \epsilon), \quad \boxed{(s, w) \mapsto (s + f(w), w)}$$

with

$$f: [-\epsilon, \epsilon] \to S^1, \quad w \mapsto \exp(-iqg(w/\epsilon)).$$



Then $F_*A \wedge dA = A \wedge dA$: the new flow X_F preserves volume.

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Fixing the contact form "Splitting the difference"

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But $F_*A = A + wf'(w) dw$. To fix, "split the difference":

 $h(t, w) := \frac{1}{2} \lambda(t) \int_{-\epsilon}^{w} x f'(x) \, dx \text{ on } (-\eta, \eta) \times (-\epsilon, \epsilon) \text{ and } h = 0 \text{ outside.}$ $\lambda: \mathbb{R} \to [0,1] \text{ is a smooth bump function}$

Then $dh_{|_{t=0}} = \frac{1}{2}wf'(w) dw$; $A_h := A \pm dh$ for $\pm t \ge 0$ gives

 $F_*(A_h) = F_*(A - dh) = F_*A - F_*dh = (A + 2dh) - dh = A + dh = A_h$

 \Rightarrow *A*_{*h*} defines a contact form on the surgered manifold.

Its Reeb field is a time-change $X_h :=$

$$\frac{X_F}{1\pm \mathrm{d}h(X_F)}.$$

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Did I mention hyperbolicity?



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The flow is hyperbolic (only) if $q \in \mathbb{N}$ by the cone criterion—the surgery transition kicks cones into cones...

... and indeed in such a way that:

Proposition (Lyapunov exponents)

The positive Lyapunov exponent of X_F is no less than the positive Lyapunov exponent of X, so

 $h_{\text{Liouville}}(X_F) \ge h_{\text{Liouville}}(X)$

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by the Pesin entropy formula.

Tufts Periodic orbit growth

Increased complexity by Suppose the original geodesic flow X_0 comes from a surface of constant curvature. Foulon Entropy Rigidity constant curvature $h_{\text{top}}(X_0) = h_{\text{Liouville}}(X_0) \le h_{\text{Liouville}}(X_h) < h_{\text{top}}(X_h)$ "no less Lyapunov exponent" Periodic orbit \Rightarrow More periodic orbits: $h_{top}(X_h) > h_{top}(X_0)$ $\lim_{t \to \infty} th_{\rm top}(\varphi) P_t(\varphi) e^{-th_{\rm top}(\varphi)} = 1, \quad i.e., \quad P_t(\varphi) \sim \frac{e^{th_{\rm top}(\varphi)}}{th_{\rm ton}(\varphi)}.$ NB: := #closed orbits of length up to t

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This contact *structure* forces orbit growth Other Reeb fields

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Given a contact structure $\xi = \ker \alpha$, consider another contact form $\alpha' = f \alpha$ where $f \in \mathscr{C}^{\infty}(M, \mathbb{R} \setminus \{0\})$. Then $\ker \alpha' = \ker \alpha = \xi$.

But $d\alpha' = df \wedge \alpha + f d\alpha$ —so ker $d\alpha'$ is quite different. The condition $\iota_{R_{\alpha'}} d\alpha' = 0$ implies that R_{α} and $R_{\alpha'}$ are not collinear unless *f* is constant.

A Reeb field on a contact manifold (M,ξ) is the Reeb field of any contact form α with $\xi = \ker \alpha$. "A contact form on (M,ξ) "

By Lieberman's Theorem, these are exactly the nowhere-vanishing vector fields transverse to ξ whose flows preserve ξ .

This contact *structure* forces orbit growth (Not new: if a Reeb flow is Anosov then others have positive entropy)

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Theorem (Periodic orbit growth in homotopy classes)

 $(M_F, A_h) = contact manifold from surgery (not assumed hyperbolic), X_h Anosov, <math>\rho$ a primitive free homotopy class containing an X_h -periodic orbit, λ hypertight⁶ contact form. $\exists a > 0, b \in \mathbb{R}$ with $N_T^{\rho}(\lambda) \ge a \ln(T) + b$ for all T > 0? [We use estimates by Barthelmé–Fenley.]

Theorem (Entropy without hyperbolicity—after Alves) ("separating")

 $(M_F, A_h) = contact manifold from surgery along a simple geodesic (so <math>M_F$ is not hyperbolic), A_h not assumed Anosov. If λ is a contact form on $(M_F, \ker(A_h))$, then $h_{top}(R_{\lambda}) > 0$.

If we drop "hypertight" and assume M_F hyperbolic we get "ln(ln T)."

 ${}^{7}N_{T}^{\rho}(\lambda) =$ number of closed orbits in ρ of length up to T.

⁶: \Leftrightarrow "no contractible periodic Reeb orbit."

More contact flows

An S^1 -family of contact Anosov flows on $S\mathbb{H}$

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 $\begin{pmatrix} 0 & \frac{1}{2} \\ \frac{1}{2} & 0 \end{pmatrix} \begin{pmatrix} 0 & -\frac{1}{2} \\ \frac{1}{2} & 0 \end{pmatrix} \begin{pmatrix} \frac{1}{2} & 0 \\ 0 & -\frac{1}{2} \end{pmatrix}$ Structure equations: $[V, X] =: H, \quad [H, X] = V, \quad [H, V] = X$ horizontal geodesic vertical Contact flows: $B := dA(V, \cdot)$ is a contact form with $R_B = H$; so is $E := E_{\theta} := \cos \theta A + \sin \theta B_{\theta}^{8}$ and $P := R_E = \cos\theta X + \sin\theta H$ is Anosov: $Q := \cos\theta H - \sin\theta X$ and $\zeta^{\pm} := Q \pm V$ give so $0 = [P, f\zeta^{\pm}] = (\dot{f} \mp f)\zeta^{\pm} \Rightarrow \boxed{f = e^{\pm t}}.$ $[P, \zeta^{\pm}] = \mp \zeta^{\pm}$ $f\zeta^{\pm}$ invarian

 $(S^1$ -family of Anosov flows) $V = R_C$ with $C := dA(H, \cdot)$ gives a periodic flow, the fiber rotation.

 ${}^8E \wedge dE(P,V,Q) \equiv 1$

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More contact flows

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 $V \xrightarrow{\text{surgery}}$

• parabolic flow (quadratic growth) using a simple geodesic



• (nonuniformly) hyperbolic flow otherwise (Heberle).



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Foulon rigidity: $h_{\text{Liouville}}(\varphi^t) < h_{\text{top}}(\varphi^t)$ —but no control of the gap. Before:



 \Rightarrow More periodic orbits:

$$h_{top}(\varphi^t) > h_{top}(g^t).$$

B:
$$\lim_{t \to \infty} th_{top}(\varphi) P_t(\varphi) e^{-th_{top}(\varphi)} = 1$$
, *i.e.*, $P_t(\varphi) \sim \frac{e^{th_{top}(\varphi)}}{th_{top}(\varphi)}$

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The entropy gap quantified—an open problem (Idea: Bishop–Hughes–Vinhage–Yang @MRC2017)



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A double torus

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A double torus

The entropy gap quantified—an open problem (Idea: Bishop–Hughes–Vinhage–Yang @MRC2017)

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Use the configuration space with cutting sequences à la Series:



The Vinhage picture

Thank you for your attention! Feedback invited on the book draft at https://tinyurl.com/HypFlows

