

Massive Galileons and Vainshtein screening: a numerical perspective

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arXiv:2008.01456

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with J. Braden, C. Burrage, B. Coltman, B. Elder, T. Padilla, T. Wilson



The problem

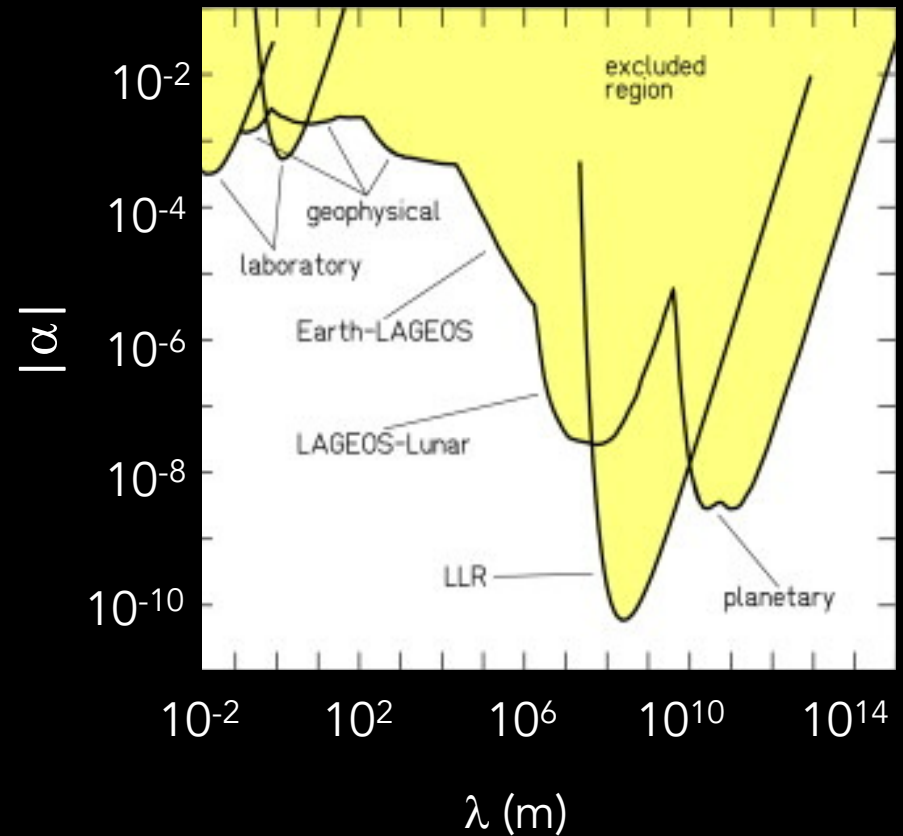
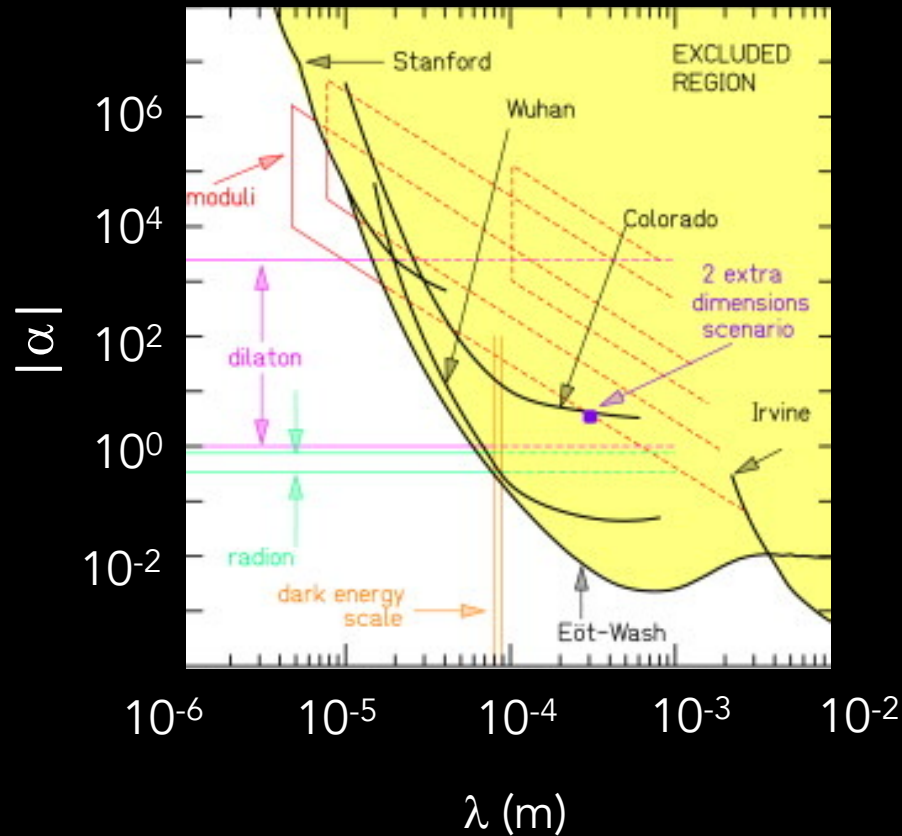
Vainstein screening and
UV completion

Scalar extensions and 5th forces

- Theories attempting to solve the **cosmological constant problem** come with extra scalar fields
- Scalar extensions come with extra mediated **5th forces**
- **Do we see extra forces** experimentally?

Scalar extensions and 5th forces

$$V(r) = -\alpha G \frac{m_1 m_2}{r} e^{-\frac{r}{\lambda}}$$



E.G.Adelberger J.H.Gundlach B.R.Heckel S.Hoedl, S.Schlamming, Progr. in Part. and Nucl. Phys., 62, 1, 2009, 102-134; C. Will, Living Reviews in Relativity 2006.

Scalar extensions and 5th forces

→ *Canonical* scalar fields *with small mass* disfavoured

$$\mathcal{L} = -\frac{1}{2}(\partial\phi)^2 - \frac{1}{2}m^2\phi^2 - \frac{\phi}{M}T^\mu{}_\mu$$

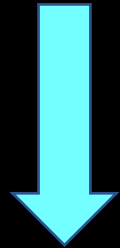
Scalar extensions and 5th forces

- If scalar field has non-trivial **self-interactions**, 5th force may be **screened** (= suppressed) where it is known to be small...
- ... while still being able to mediate a **long-range force** where it can be relevant at current **cosmological scales**

$$\mathcal{L} = -\frac{1}{2}Z^{\mu\nu}(\phi, \partial\phi, \dots)\partial_\mu\partial_\nu\phi - V(\phi) + g(\phi)T^\mu{}_\mu$$

Scalar extensions and 5th forces

$$V(r) = -\alpha G \frac{M_s m}{r} e^{-\frac{r}{\lambda}}$$



$$V(r) = -\alpha(\phi_{bg}) G \frac{f(M_s) m}{r} e^{-\frac{r}{\lambda(\phi_{bg})}}$$

Debye screening of electromagnetism

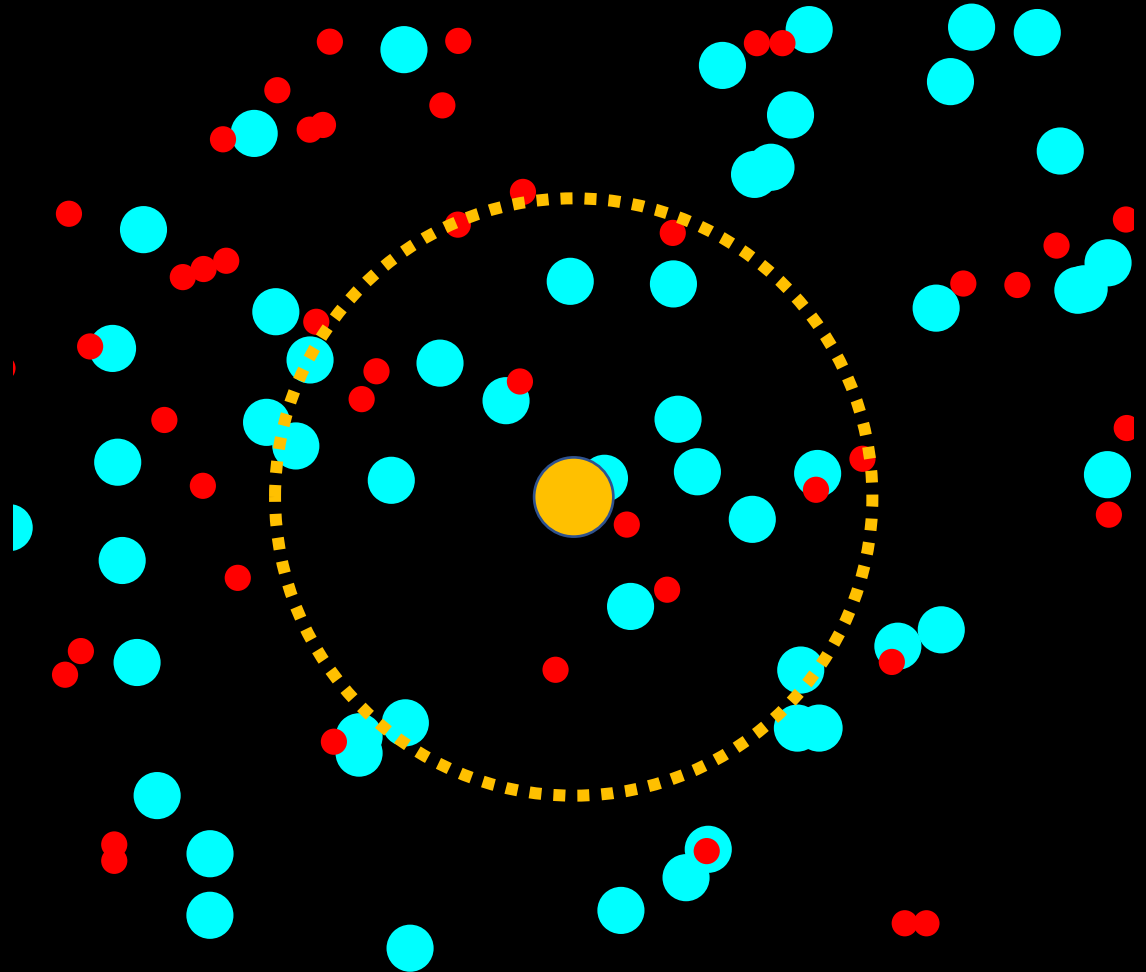
The influence of a charged particle in a plasma is only felt up to the Debye length $\sim 1/m_{\text{Debye}}$

$$m_{\text{Debye}}^2 = \frac{8\pi n e^2}{T}$$

n electron density

e electron charge

T plasma temperature



Chameleon screening



Inside galaxy: density high,
Chameleons short-ranged

Outside galaxy: density low,
Chameleons long-ranged



Symmetron screening

$$V_{\text{eff}}(\phi) = \frac{1}{2} \left(\frac{\rho}{M^2} - \mu^2 \right) \phi^2 + \frac{\lambda}{4!} \phi^4$$

Inside galaxy: density high,
Symmetron decoupled

$$\phi \rightarrow 0$$

Outside galaxy: density low,
Symmetron coupled

$$\phi \rightarrow \mu/\sqrt{\lambda}$$

Local matter coupling: $\sim \phi/M^2$

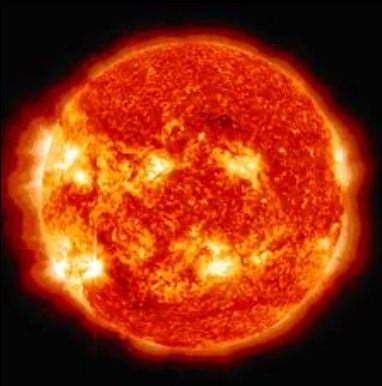
Vainshtein screening

Standard gravity $\square\phi = \frac{\rho}{M}$

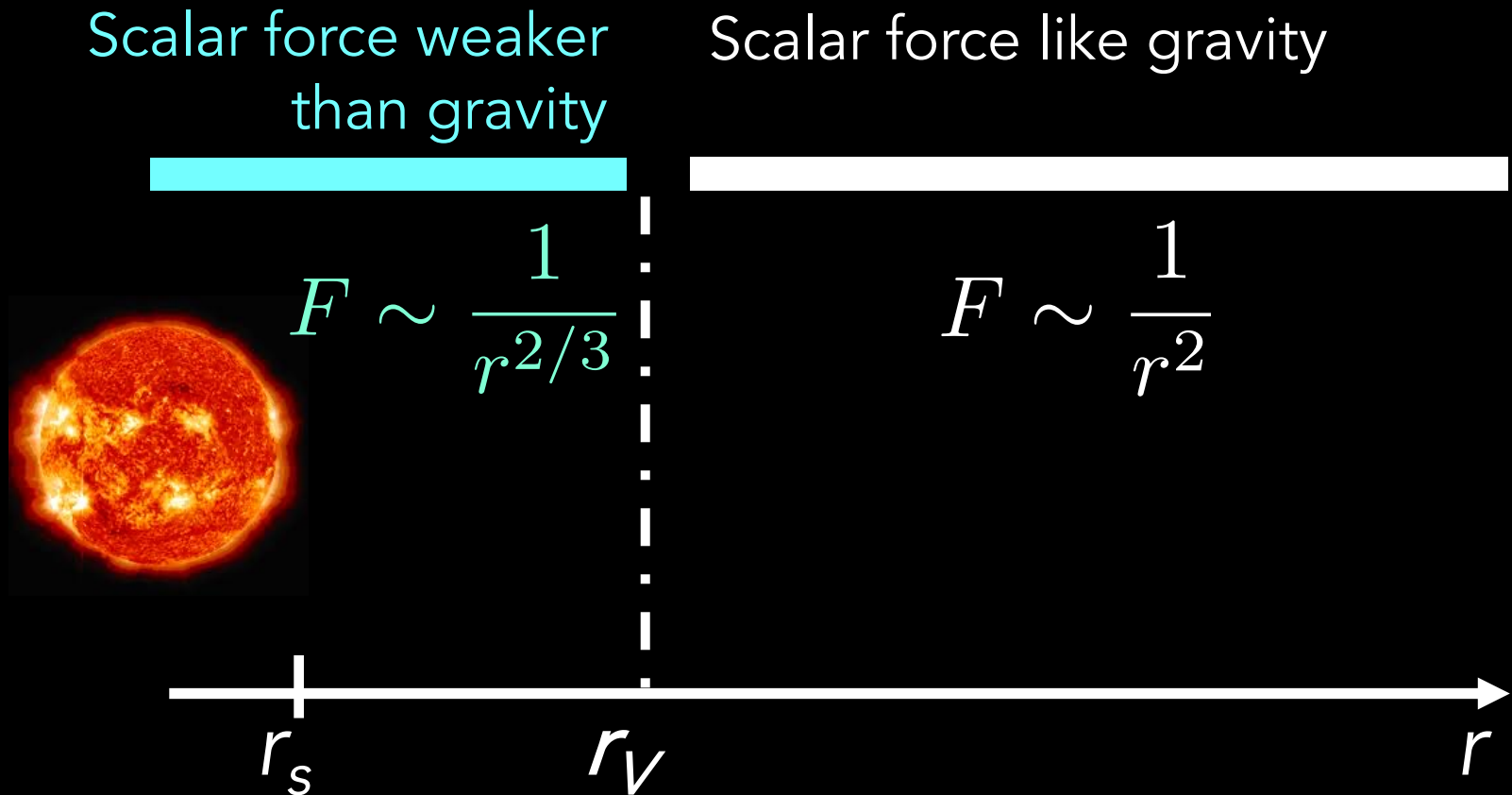
Cubic Galileons (with Vainshtein) $\square\phi + (\square\phi)^2 - (\partial_\mu\partial_\nu\phi)^2 = \frac{\rho}{M}$

High derivative, non-linear terms dominate

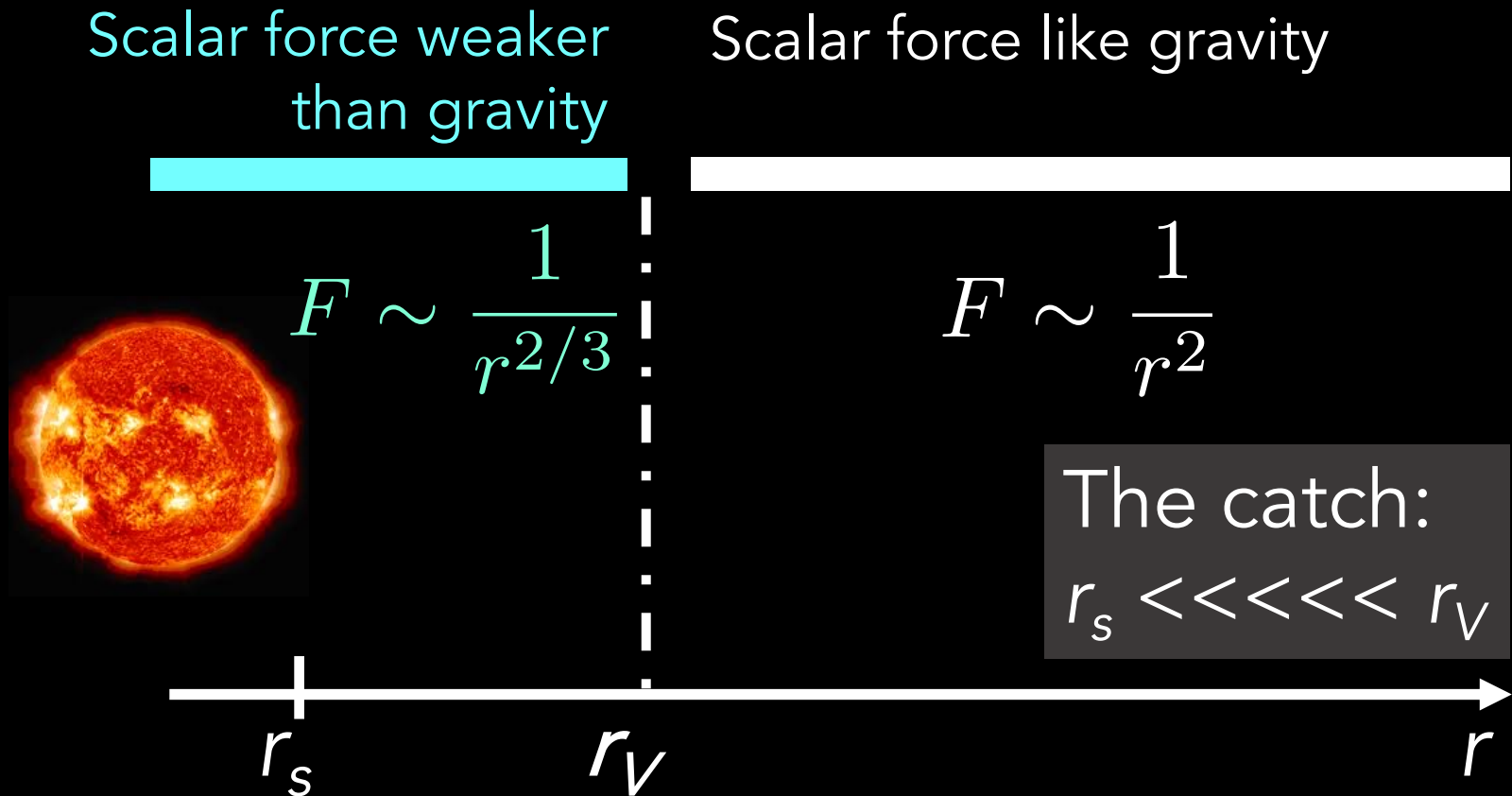
Standard terms dominate



Vainshtein screening



Vainshtein screening



Vainshtein screening

$$\mathcal{L} = -\frac{1}{2}(\partial\phi)^2 - \frac{1}{2\Lambda^3}(\partial\phi)^2\Box\phi$$



Typically, Vainshtein screening is taken to be an effective field theory...

...but if specific higher-order kinetic terms are large, then we are out of a perturbative regime...

Vainshtein screening

$$\mathcal{L} = -\frac{1}{2}(\partial\phi)^2 - \frac{1}{2\Lambda^3}(\partial\phi)^2\Box\phi$$



why are these higher-order kinetic term macroscopically important but others are not?

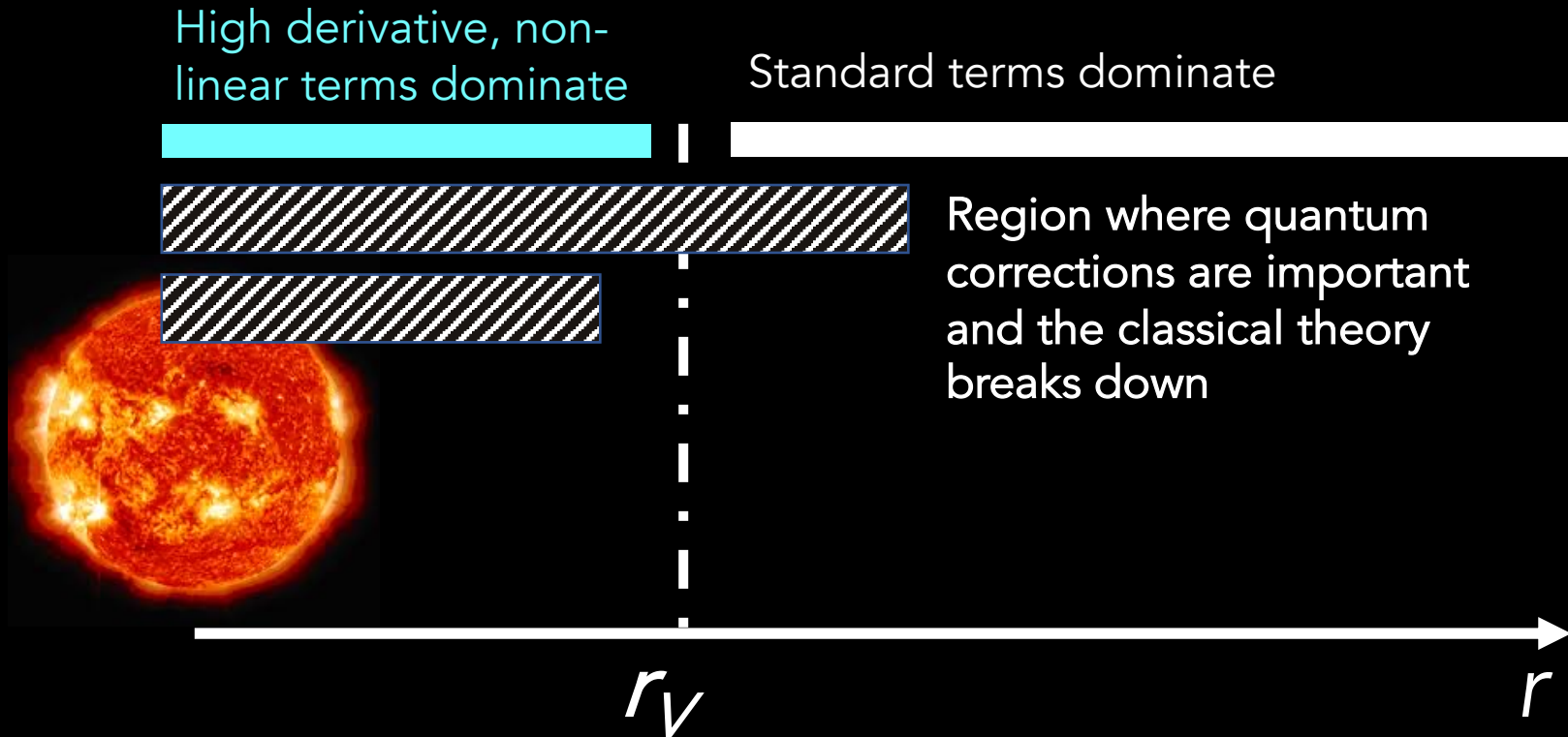
Vainshtein screening

$$\mathcal{L} = -\frac{1}{2}(\partial\phi)^2 - \frac{1}{2\Lambda^3}(\partial\phi)^2\Box\phi$$



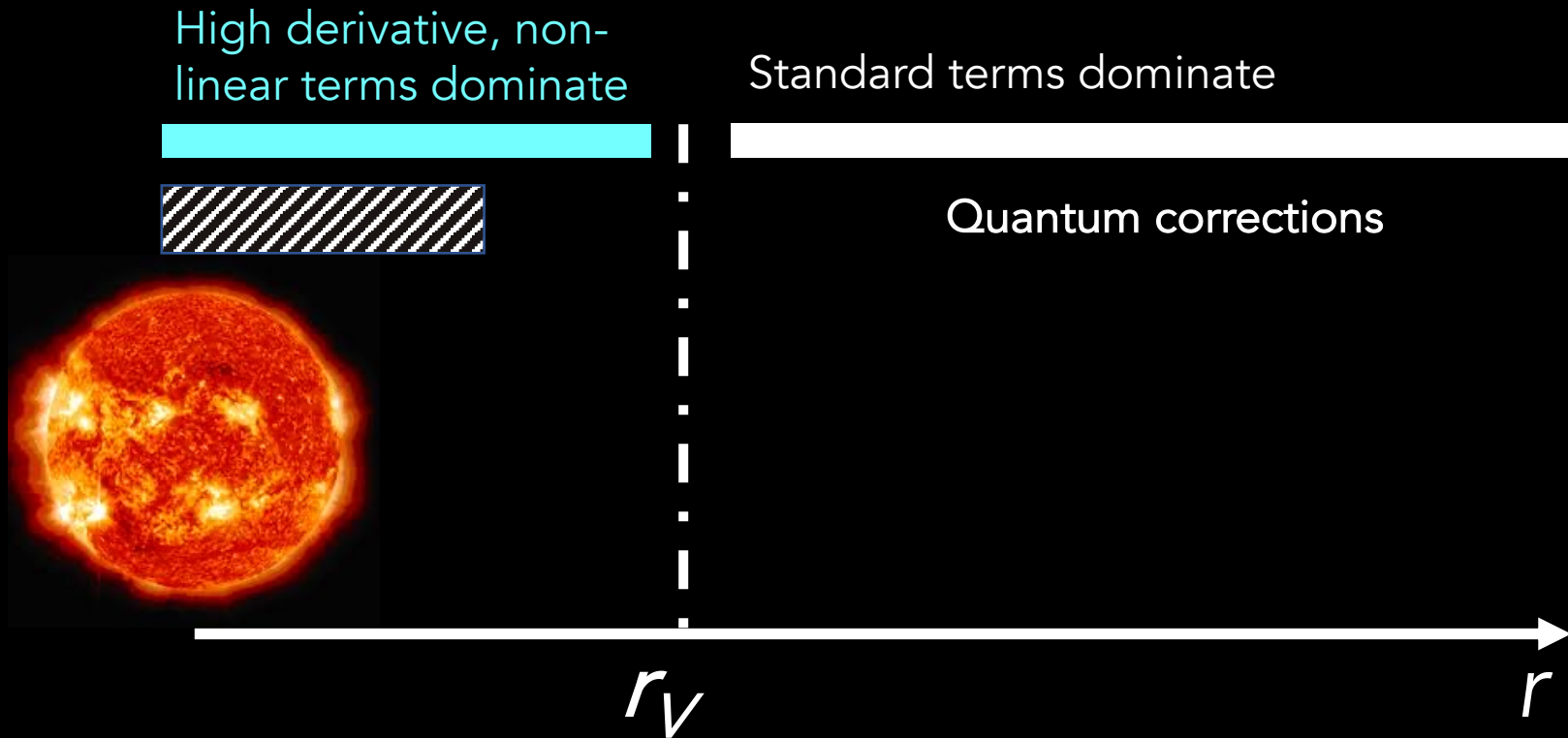
You can only really tell if you
know the full UV sector

Vainshtein screening – more accurate picture



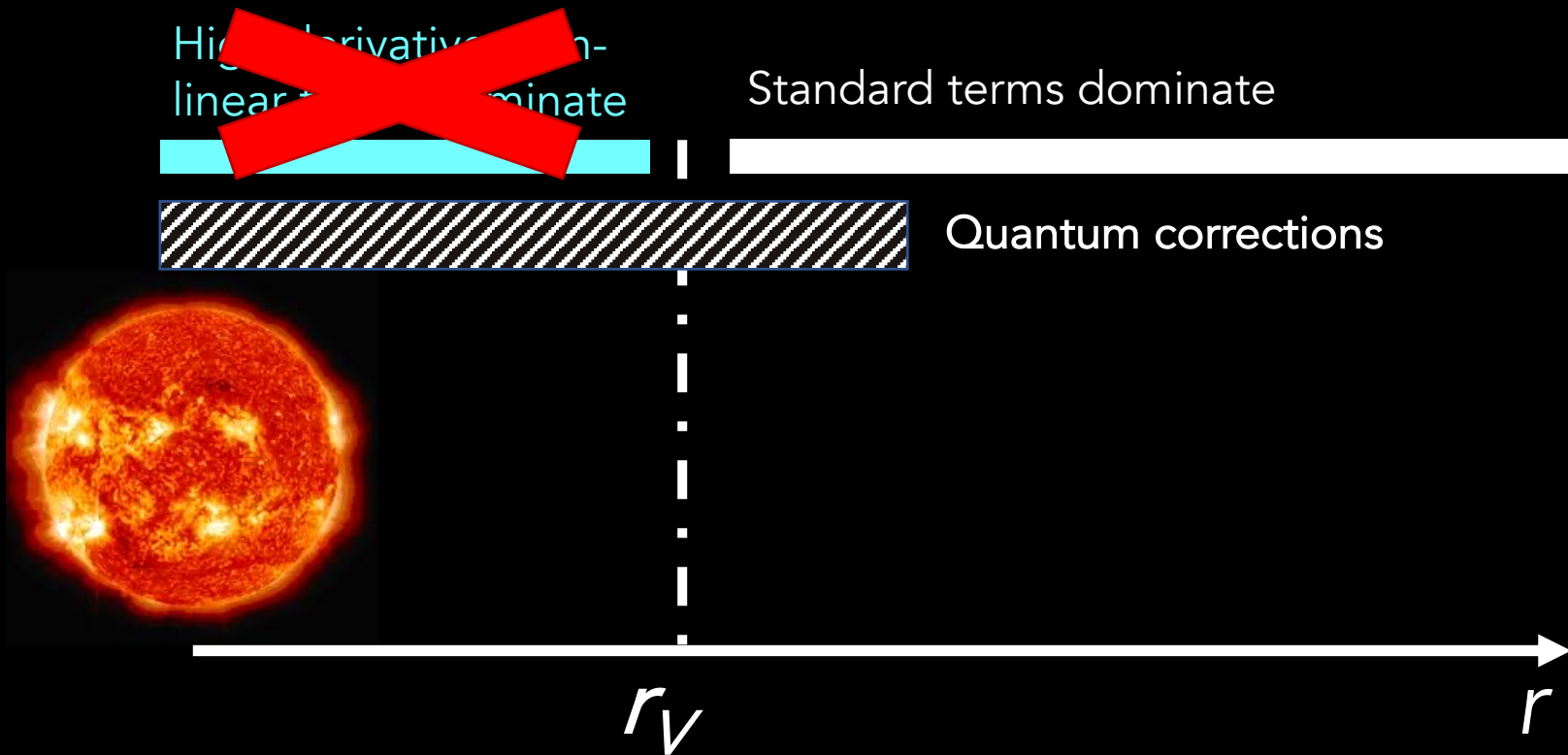
Vainshtein screening – more accurate picture

Λ_3 theories: classical nonlinearities kick in before quantum corrections;
however, standard UV completion not possible



Vainshtein screening – more accurate picture

Λ_5 theories: standard UV completion possible;
however, quantum corrections kick in before
classical nonlinearities



The goal

- Compare a low-energy theory with Vainshtein screening against its UV completion
- Does the screening survive the extension?

A UV-complete theory of massive Galileons

$$\mathcal{L} = -\frac{1}{2}(\partial\phi)^2 - \frac{1}{2}m^2\phi^2 - \frac{1}{2}(\partial H)^2 - \frac{1}{2}M^2 H^2 \\ -\alpha H\Box\phi - \frac{\lambda}{4!}H^4 - \frac{\rho}{M_{pl}}\phi$$

Theory of two coupled massive scalar fields, one self-coupled
Important higher-order kinetic terms

Does it screen?

A UV-complete theory of massive Galileons

$$\begin{cases} \square\phi - m^2\phi - \alpha\square H = \frac{\rho}{M_{pl}} \\ \square H - M^2 H - \alpha\square\phi - \frac{\lambda}{3!}H^3 = 0 \end{cases}$$

Theory of two coupled massive scalar fields, one self-coupled
Important higher-order kinetic terms

Does it screen?

Low-energy theory

From this high-energy theory, it is possible to obtain the IR theory

$$\square\pi - m^2\pi - \frac{\lambda\alpha^4}{3!} \frac{\square(\square\pi^3)}{M^8} + \dots = \frac{\rho}{M_{pl}}$$

which relies on the series of operators

$$O_n \propto (-1)^{n+1} \lambda^n \alpha^{2n+1} \frac{\square(\square\pi^{2n+1})}{M^{6n+2}}$$

to converge.

Does it look anything like the
high-energy theory?

The goal

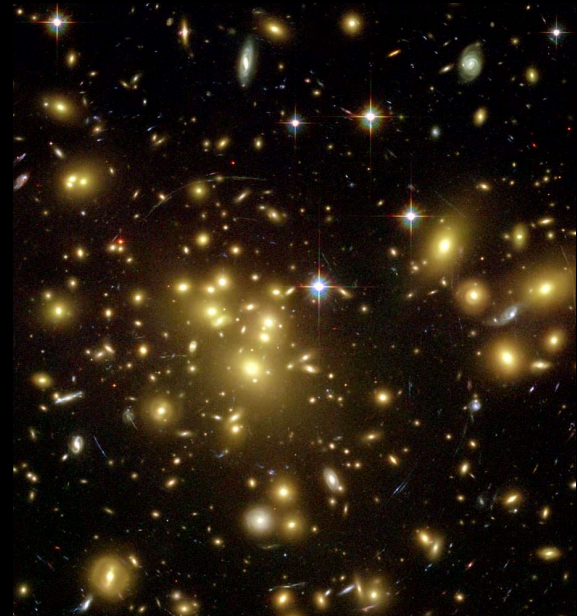
- Let's compute the scalar force for the high-energy and low-energy theories, to see which one, if any, displays Vainshtein screening
- Let's compute the O_n operators to see if the series converge

The problem

- We need to solve complex equations characterised by large non-linearities, sharp transitions, complex boundary conditions

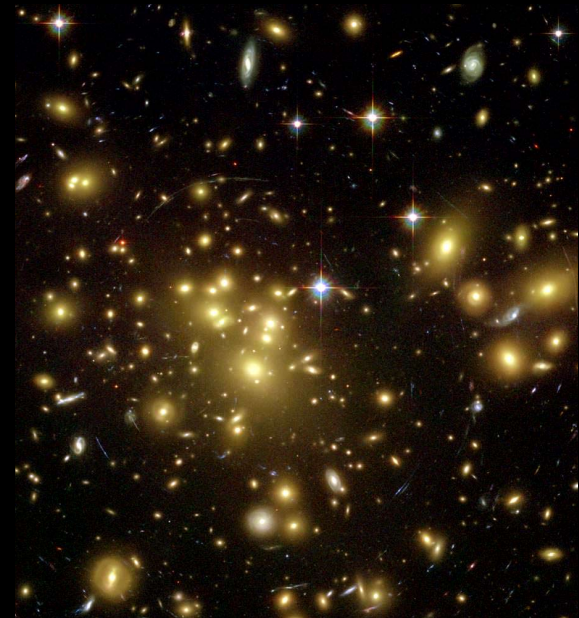


VS



The problem

- We need to solve complex equations characterised by large non-linearities, sharp transitions, complex boundary conditions





Method

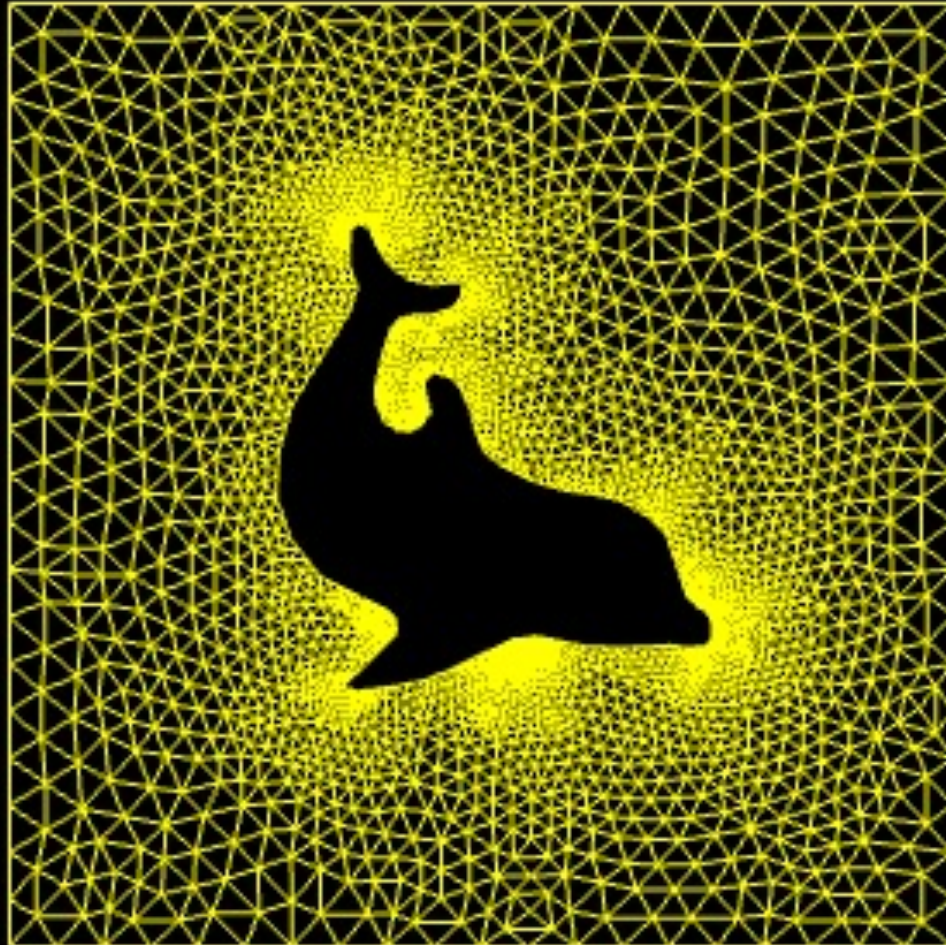
φ enics: screening with the
finite element method

ϕ enics

- Finite-element code for the solution of the full equation of motion of models of screening
- Builds on FEniCS library
- Computes field profiles, fifth force, (horrible) high-order operators accurately

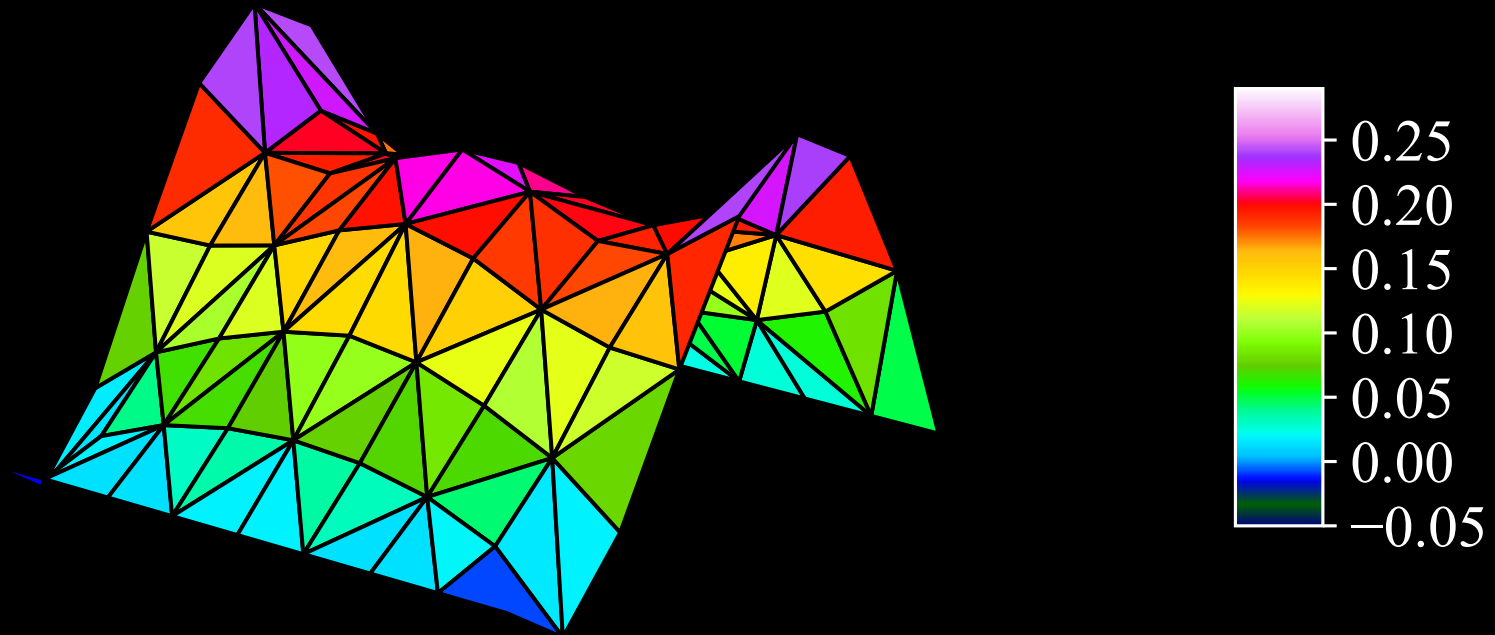
The finite element method

Step (1): discretise domain



The finite element method

Functions can be approximated piecewise by polynomials:



$$f = \sum_i^N f(P_i) e_i$$

P_i vertices
 e_i basis polynomials

The finite element method

Step (2): turn original (*strong*) equation into its **weak** form

- Example: Poisson equation

Strong problem: find u such that

$$-\nabla^2 u = f$$

$$u = u_D$$

$$-\frac{\partial u}{\partial n} = g$$

in Ω (spatial domain)

in $\partial\Omega_D$ (Dirichlet boundary)

in $\partial\Omega_N$ (Neumann boundary),

n outward normal direction to the boundary

Weak formulation: find u such that

$$-\langle \nabla^2 u, v \rangle = \langle f, v \rangle \quad \forall v \in \hat{V},$$

$\hat{V}$ space of 'test' functions,
 $\langle \cdot, \cdot \rangle$ inner product

With the inner product $\langle f, g \rangle = \int_{\Omega} f g \, dx$:

$$-\int_{\Omega} \nabla^2 u \, v = \int_{\Omega} f v \quad \forall v \in \hat{V}$$

using integration by parts:

$$\int_{\Omega} \nabla u \nabla v - \int_{\partial\Omega} \frac{\partial u}{\partial n} v \, ds = \int_{\Omega} f v$$

The finite element method

Integrating by parts and applying Neumann boundary conditions...

$$\int_{\Omega} \nabla u \nabla v = - \int_{\partial\Omega_N} g v \, ds + \int_{\Omega} f v$$

$$a(u, v) := \int_{\Omega} \nabla u \nabla v$$

bilinear form

$$L(v) := - \int_{\partial\Omega_N} g v \, ds + \int_{\Omega} f v$$

linear form

The finite element method

Step (3): turn complicated equation into algebraic system

Because a and L are bilinear and linear, respectively, they only need to be defined on basis vectors:

find u such that:

$$a(u, v) = L(v) \quad \forall v \in \hat{V}$$

Expanding in the discretised basis:

$$a(u, e_i) = L(e_i) \quad \forall i = 1, \dots, N$$

$$\sum_{j=1}^N u_j a(e_j, \hat{e}_i) = L(\hat{e}_i) \quad \forall i = 1, \dots, N$$



$$AU = b$$

Coefficients of u in
discretised basis

The finite element method

Step (4): if equation is nonlinear, solve iteratively using e.g. Newton's method

$$F(u) = 0$$



$$F(u_k) + J(u_k)(u - u_k) = 0$$

Starting from some initial guess u_k ,
Then substituting $u \rightarrow u_k$ until convergence



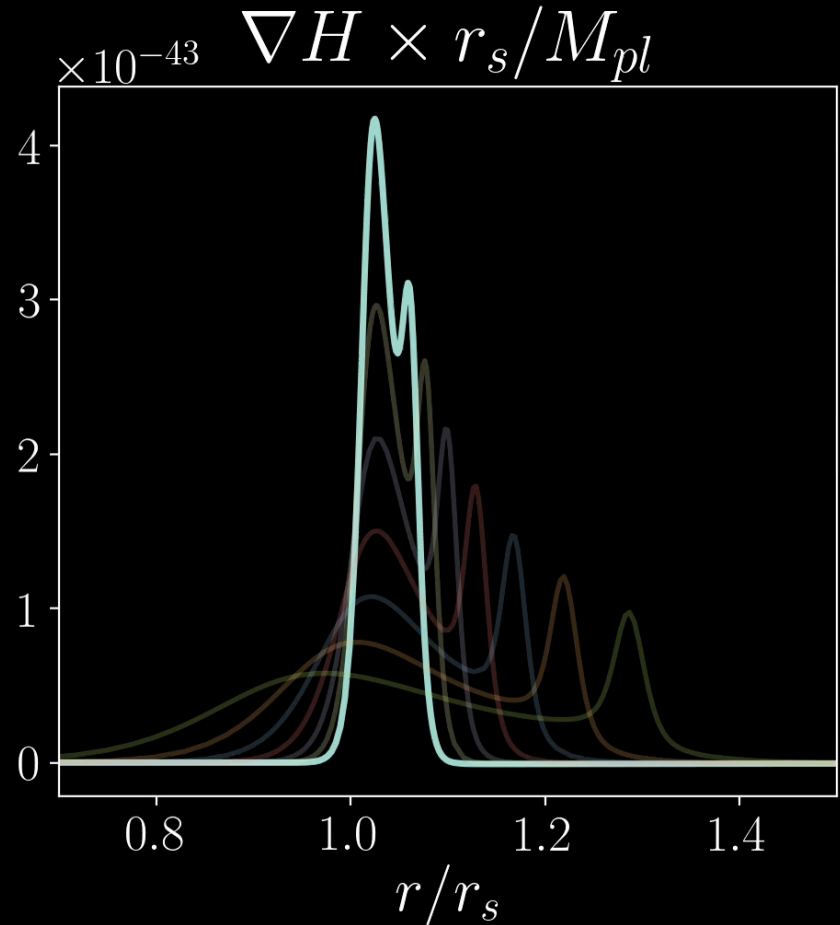
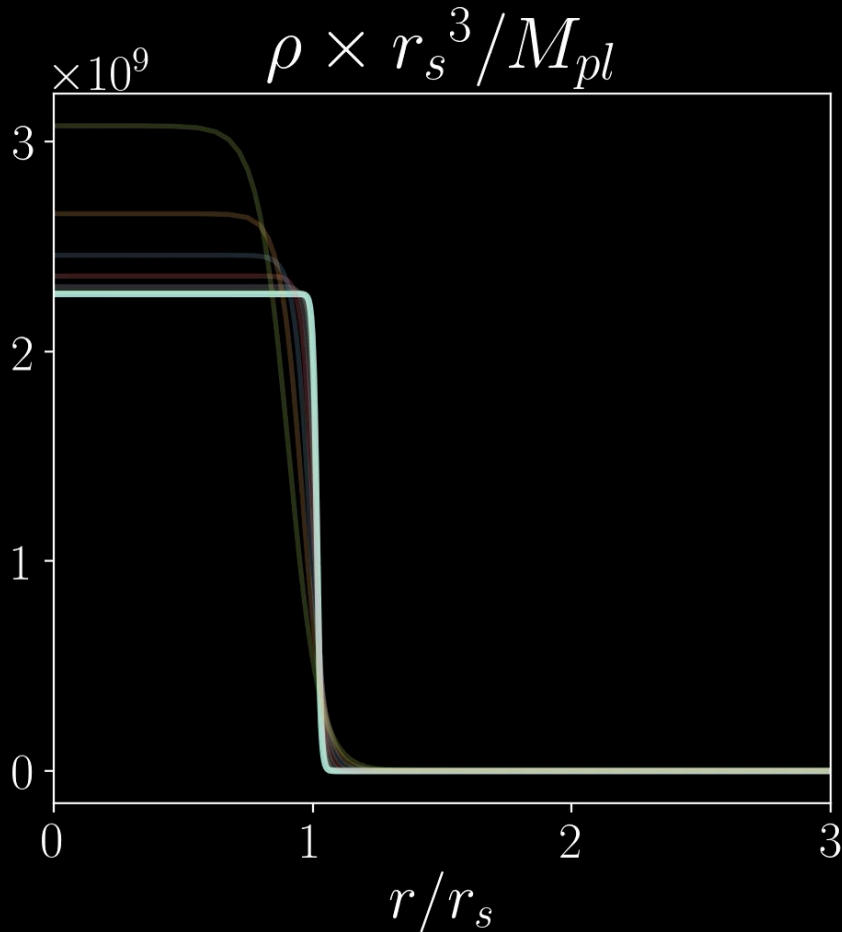
Prelude to results

φ enics: screening with the
finite element method

Features at the source-vacuum transition

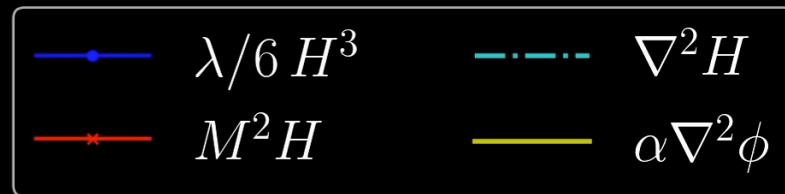
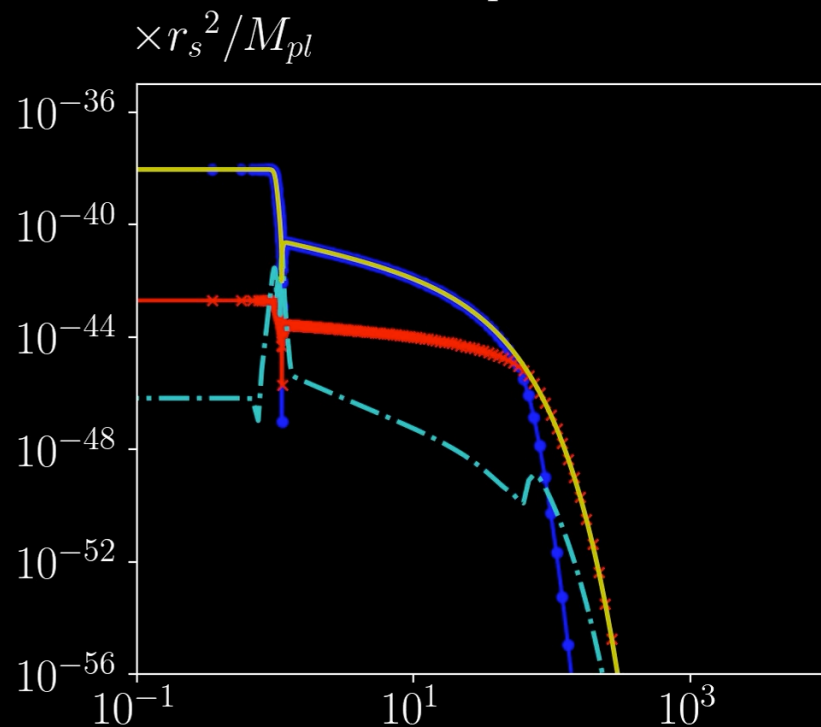
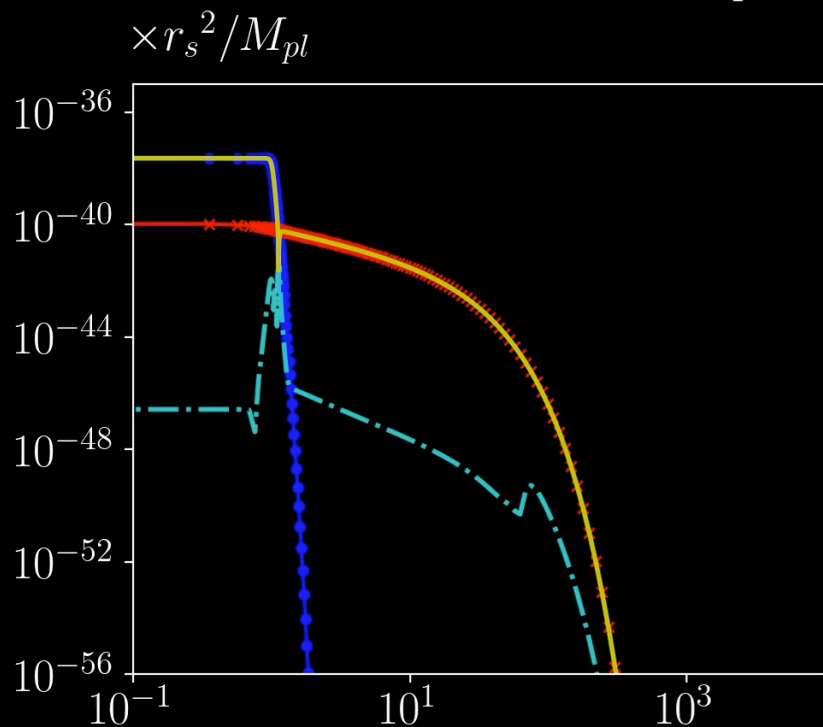
$$mr_s = 1/10$$

$$Mr_s = 10$$

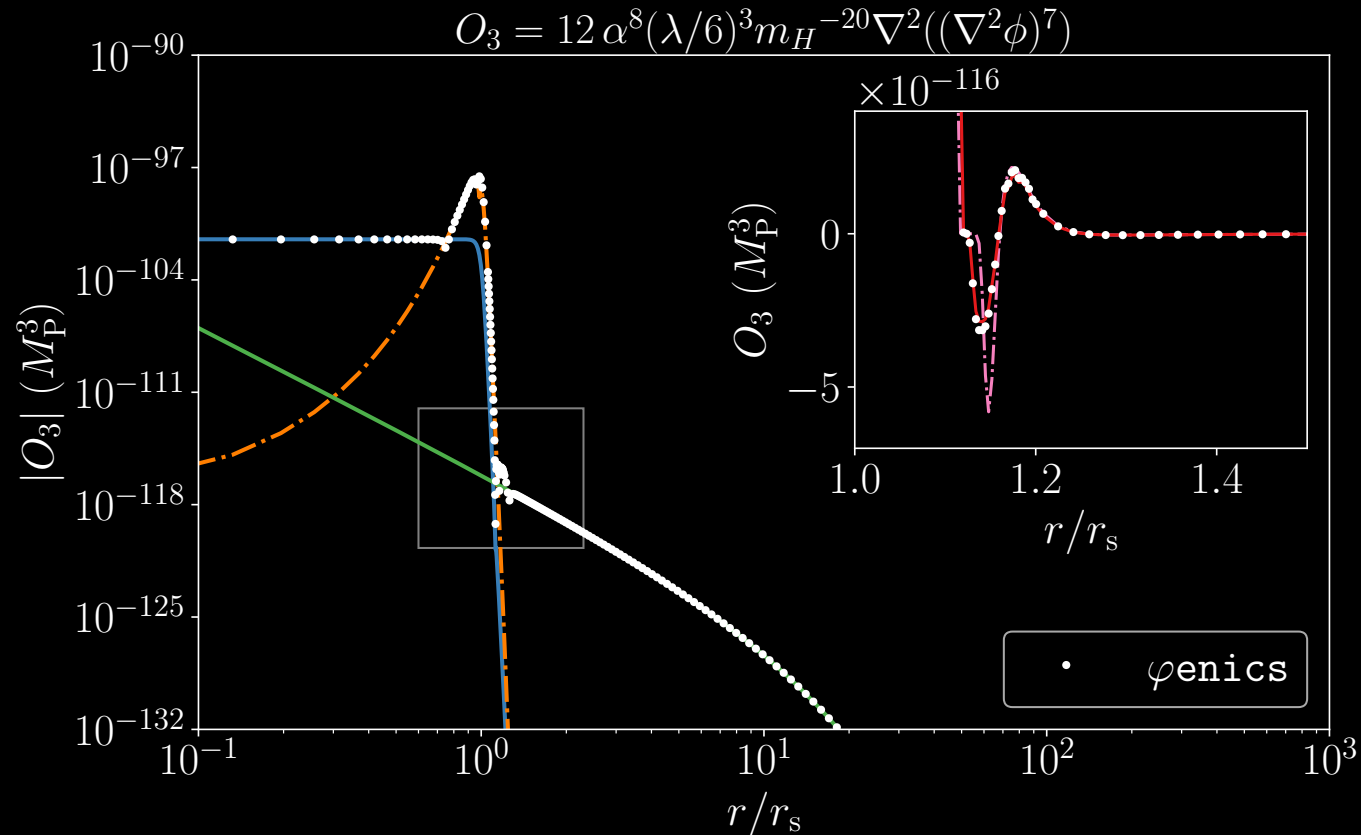


Terms in the equations across all regimes

$$m = 10^{-48} M_{pl} \quad M = 3.2 \times 10^{-47} M_{pl}$$



High order operators

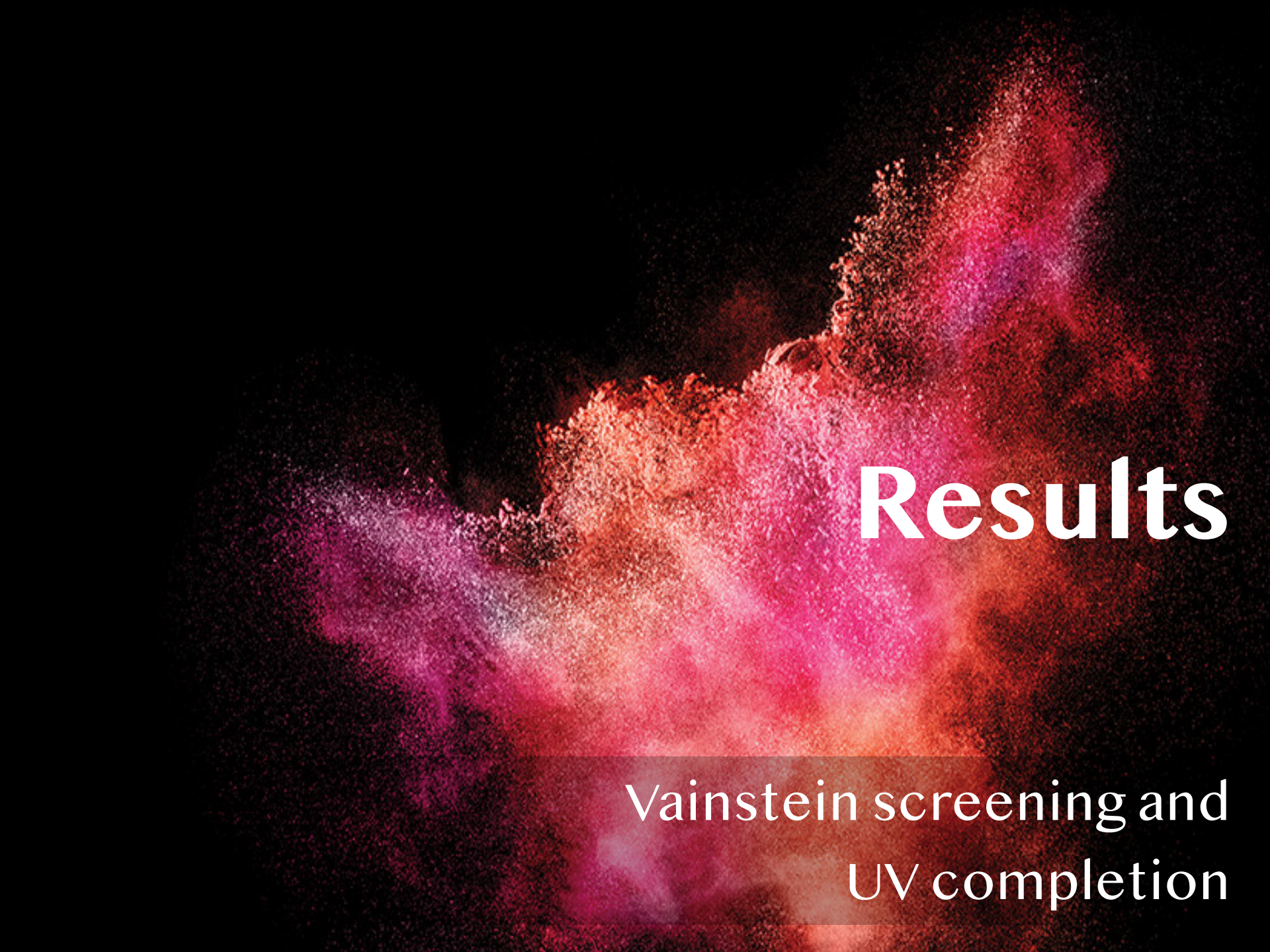


Light-field dominated

- source dominated
- · - · source-vacuum transition
- Yukawa suppressed

Heavy-field dominated

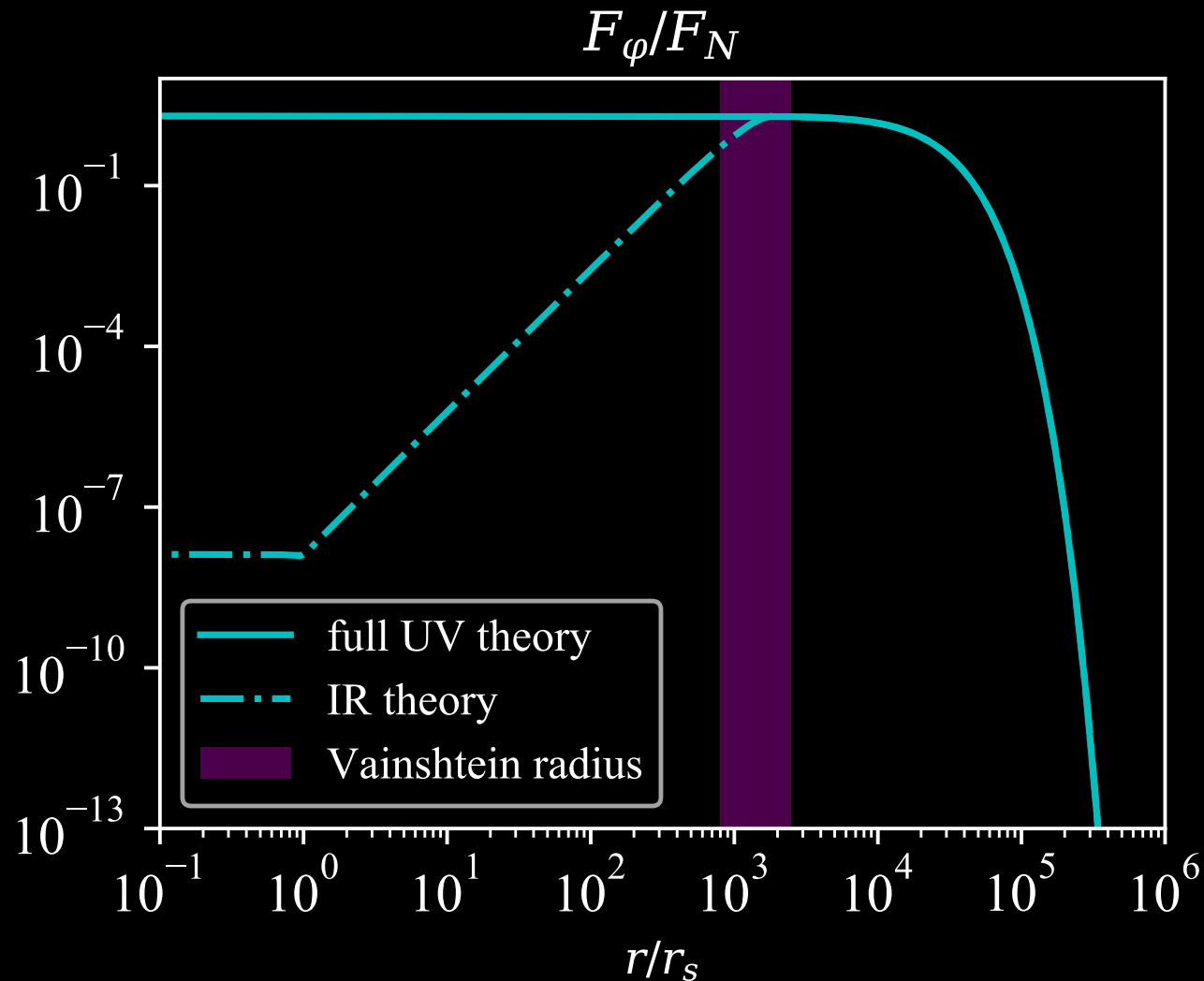
- · - · $\nabla^2 \phi \sim -\frac{\lambda}{6\alpha} H^3$
- $\nabla^2 \phi \sim -\frac{\lambda}{6\alpha} H^3 + \frac{1}{\alpha} \nabla^2 H$



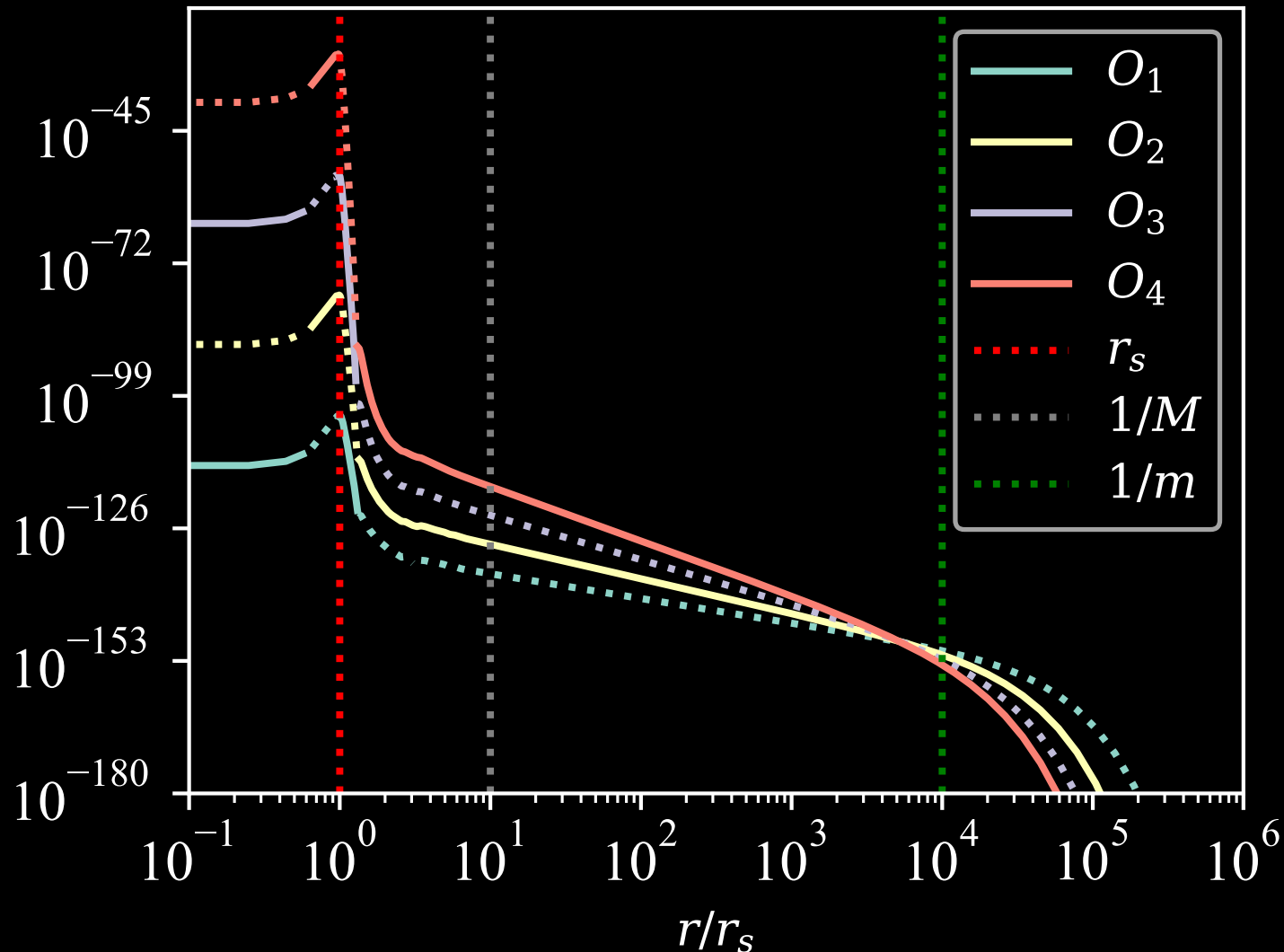
Results

Vainstein screening and
UV completion

Vainshtein screening – is it viable?



Vainshtein screening – is it viable?





Summary and conclusions

Summary and conclusions

- Vainshtein screening is an effective mechanism to hide the fifth force in high-density regions; it is characterised by complicated equations that are impossible to solve analytically and difficult to solve numerically;
- We developed **ϕ enics**, finite-element code for the full numerical solution of equations of screening
- We compared a theory of massive Galileons against its UV completion: the IR theory screens, **the UV theory doesn't!**
- It is unclear Vainshtein screening can be invoked within the limits of validity of the theories that make use of it

The background of the image is a dynamic, abstract composition of fine, glowing particles in shades of magenta, pink, and orange. These particles are concentrated in the center and right side, creating a sense of movement and energy, while the left side fades into a solid black background. The overall effect is reminiscent of a cosmic nebula or a microscopic view of a chemical reaction.

**Thank you for
your attention!**