Some applications of Q balls in the early Universe

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Based on arXiv:2105.11655 with Lauren Pearce, Daniel Vagie and Alex Kusenko and arXiv 2111.xxxx with Lauren Pearce, Volodymyr Takhistov and Alex Kusenko



Q balls are macroscopic objects that can exist in the early Universe

Conditions

1. A potential that is flatter than quadratic

$$\frac{V(\phi)}{\phi^2} \supset minimum$$

2. A conserved charge

When both are satisfied the energy of a bunch of scalar quanta forming an extended macroscopic object is lower than the same number of quanta being free particles

$$E = UV + A\sigma + \frac{Q^2}{\phi^2 V}$$

In the thin wall approximation, energy of a Q ball:

$$E = UV + A\sigma + \frac{Q^2}{\phi^2 V}$$

Radius of a Q-ball

$$R \sim \left(\frac{Q}{\omega v^2}\right)^{1/3}$$

Energy of each quanta in a Q ball

$$\omega = \sqrt{\frac{2U}{\phi^2}}$$

Part 1: Higgs Balls

Higgs example:

- 1. Early enough in the Universe, Higgs decays are slow compared to H
- 2. Naively the Higgs potential is much faster than quadratic

$$V = -\mu^2 \left| H^{\dagger} H \right| + \lambda \left| H^{\dagger} H \right|^2$$

However, the potential flattens out in the early Universe for two reasons

1.
$$\lambda \to \lambda(\mu)$$

2. $V \to V + cT^2h^2 - ETh^3$

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When is Higgs number a good conserved quantum number?

Naively quite high temperatures! $O(10^{13})$ GeV!

Not impossible for preheating for two reasons

- 1. Preheating temperature can be about 2 orders of magnitude higher
- 2. The instability scale grows with temperature



3. In any case, within experimental error, the Instability scale can be higher than this

When is Higgs number a good conserved quantum number?

Naively quite high temperatures! $O(10^{13})$ GeV!

For Higgs balls though, stability looks even rosier

- 1. Fermions can only decay at the surface due to Pauli Blocking inside
- 2. Decay rate to gluons would be dominant but Higgs quanta is below the transverse gluon Debye mass



Note on hypercharge:

Higgs balls with a large H number require a local hypercharge asymmetry to cancel the hypercharge, otherwise there is a maximum size Higgs ball

To minimize the energy, the system produces fermions at the same time the phase starts spinning We have a new macroscopic object in the Standard Model!

Now what to do with it?

Weinberg operator
$$\frac{v_H^2}{\Lambda^2}LL$$
 unsuppressed for huge vev

Can have CP violating interactions between neutrinos and the Higgs wall

Results in an (up to) O(1) asymmetry in the Higgs ball which is suppressed by the volume fraction within the balls

$$R \sim \left(\frac{Q}{\omega v^2}\right)^{1/3} \sim Q^{1/3}/T < < R_H$$

Stay tuned for details in upcoming paper

Summary part 2:

New macroscopic object in the standard model

Higgs charge is a conserved quantity in the early Universe

The Standard Model Higgs potential at finite temperature allows a Higgs ball solution

Can be used to facilitate Baryogenesis

No doubt many other applications (e.g. very heavy freeze out dark matter)

Part 2: Q balls formed by Affleck Dine Baryogenesis **Baryon asymmetry of the Universe:**

A triumph and a mystery

$$Y_B = \frac{n_B - \overline{n}_B}{s} \approx \frac{n_B}{s} = \begin{cases} (7.3 \pm 2.5) \times 10^{-11}, \text{BBN} \\ (9.2 \pm 1.1) \times 10^{-11}, \text{WMAP} \\ (8.59 \pm 0.11) \times 10^{-11}, \text{Planck}. \end{cases}$$

Affleck Dine Mechanism

- SUSY has many flat directions
- Some violate B or L



Figure from review by Rouzbeh Allahverdi and Anupam Mazumdar



Elegant but difficult to test!!





Form Q balls if decays are sufficiently slow, and the potential allows

$$\operatorname{Min}\left[\frac{V(\phi)}{\phi^2}\right] \neq 0$$

Simulations show Q balls have a large symmetric component. So initial energy density in Q balls is large

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 $\rho_Q = Q_0 \omega n_0$

Where

 $n_0 \sim N_Q H_0^3$

Simulations show $N_Q \sim 10^3$ gravity mediation and $N_Q \sim O(1)$ for gauge mediation. Q_0 can be written in terms of the initial baryon asymmetry

$$Q_0 = \frac{3Y_{B_0}M_{\rm Pl}^3}{800\sqrt{5}\pi^{5/2}g_*rT_0^3}$$

B(L) balls are usually long lived

- Gauge singlet (so no decay into gauge bosons)



Forbidden either kinematically or due to Pauli blocking

Decay rate has surface to volume ratio suppression

Early period of matter domination followed by a sudden transition to radiation

Resonantly enhanced gravitational waves

Initial power spectrum

$$P_{\zeta}(k) = \Theta(k_{\inf} - k)A_s \left(\frac{k}{k_*}\right)^{n_s}$$

- During matter domination perturbations grow
- The Q-balls decay faster than exponential
- Perturbations cannot melt away in this case but produce sound waves

$$c_s = \sqrt{1/3}$$
 Constructive interference

Specific scenario

Choose gauge mediated scenario

$$V(\Phi) = m^4 \log\left(1 + \frac{|\Phi|^2}{m^2}\right) + \frac{1}{\Lambda^2} |\Phi|^6$$

With Λ in between the inflaton and GUT scale and $N_Q = 1000$

Low reheating temperature motivated by gravitino problem(s)

 $T_R \lesssim 10^7 {
m GeV}$

Benchmarks $y = (y_e, y_\mu, y_b)$



Comments about scan:

- 1. It is very common for Q-balls to be long lived enough to produce a period of matter domination
- 2. Frequently the matter domination lasts so long that we enter the non-linear regime

Distinguishing from other cosmological scenarios

- Many causes of matter domination do not work as the transition to radiation domination is too slow to resonantly enhance the GWs

- Pbhs also have surface to volume ratio suppression in decay rate
- A possible method to distinguish is GWs is sensitive to the mass spectrum which could be quite different for PBHs and Q-balls
- Need detailed simulations and scans.

Summary part 2:

Affleck Dine Baryogenesis is an elegant paradigm explaining the matterantimatter asymmetry

Very difficult to test

We argue it generically produces long lived Q-balls that leave a gravitational wave signature

There are limited cosmological scenarios that do this, and there might be methods to distinguish signals

Might be complimentary to other probes of Affleck Dine