## 3d Supersymmetric Theories on an Elliptic Curve

Mathew Bullimore

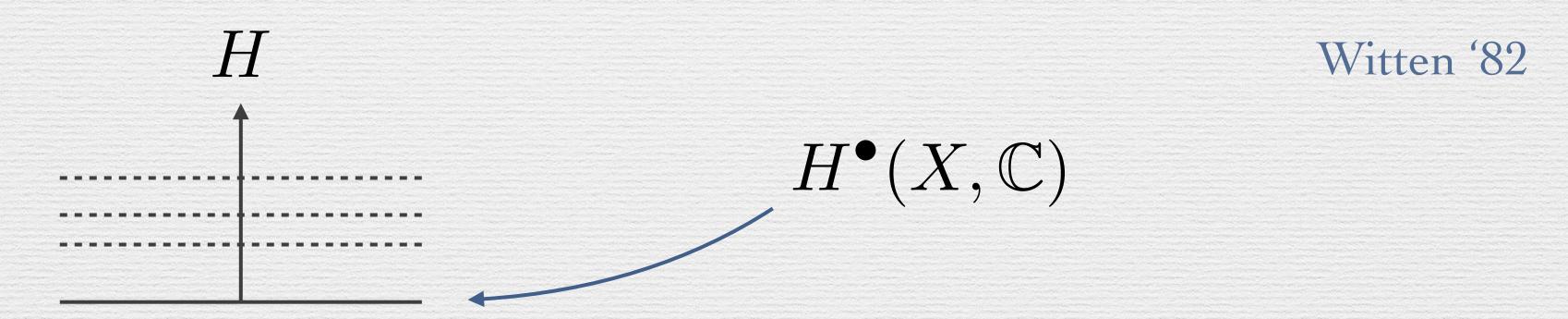


2109.10907 with Daniel Zhang

## Introduction

#### Introduction

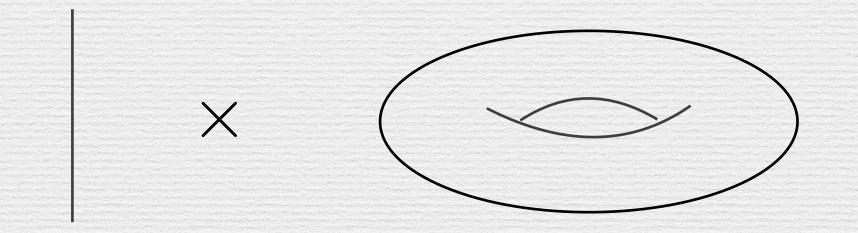
- $\bullet$  Consider a supersymmetric quantum mechanics with target X.
- The supersymmetric ground states are de Rham cohomology.



- Generalisations to 2d / 3d related to K-theory / elliptic cohomology
- The generalisation is more interesting in the equivariant case when *X* has symmetries.

#### Introduction

- What physical setup realises equivariant elliptic cohomology?
- Study 3d supersymmetric theories on an elliptic curve.



- Study supersymmetric ground states of this system.
- Physical construction of elliptic stable envelopes and R-matrices.

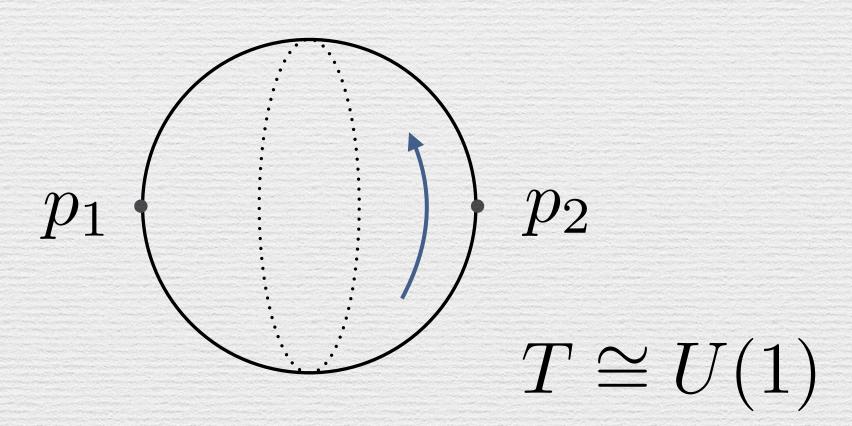
## Part 1:Background

#### Supersymmetry

- A three-dimensional QFT with at least N=2 supersymmetry.
- ullet Example: choose a smooth Kahler manifold X.
  - $\star$  A supersymmetric sigma model to X.
  - $\star$  A supersymmetric gauge theory with Higgs branch X.
- I will assume X is compact to simplify the presentation.

## Flavour Symmetry

- ullet I will assume there is an abelian symmetry T.
- This acts by Hamiltonian isometries of X.
- I will assume isolated fixed points  $X^T = \{p_{\alpha}\}$   $\alpha = 1, \dots, N$ .
- Example:  $X = \mathbb{CP}^1$

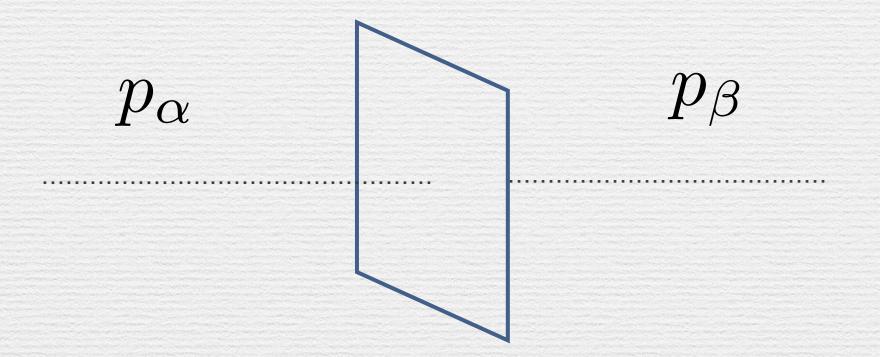


#### Massive Vacua

- Introduce real mass parameters  $m \in \mathfrak{t} := \mathrm{Lie}(T)$ .
- This generates a real super-potential  $h: X \to \mathbb{R}$ .
- It is the moment map for IPS  $T_m \subset T$  generated by m.
- For generic mass parameters  $\operatorname{Crit}(h) = \{p_{\alpha}\}$ .
- The critical points are isolated massive vacua.

#### Domain Walls

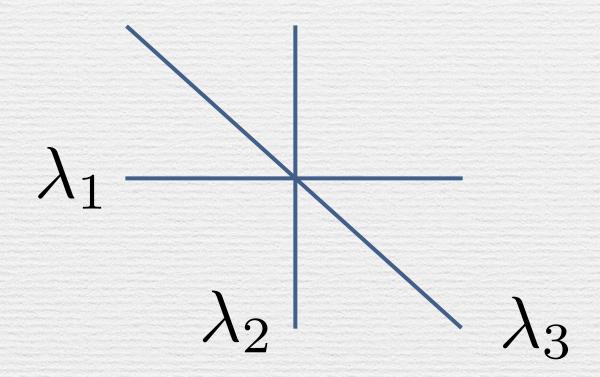
• There are BPS domain walls interpolating between vacua  $~p_{lpha} 
ightarrow p_{eta}$  .



- They correspond to gradient flow for h with tension  $|h_{\alpha} h_{\beta}|$ .
- Define linear hyperplanes in  $\mathfrak{t}$  where  $|h_{\alpha} h_{\beta}| = 0$ .
- Hyperplanes are loci where  $\operatorname{Crit}(h) \neq X^T$  and a moduli space containing  $p_{\alpha}, p_{\beta}$  opens up.

## Hyperplane Arrangement

• This forms a hyperplane arrangement in  $\mathfrak{t} \cong \mathbb{R}^{\mathrm{rk}(T)}$ .



• Each hyperplane is labelled by a  $\,T$ -weight  $\,\lambda\,$  such that

$$h_{\alpha} - h_{\beta} \propto \langle \lambda, m \rangle$$

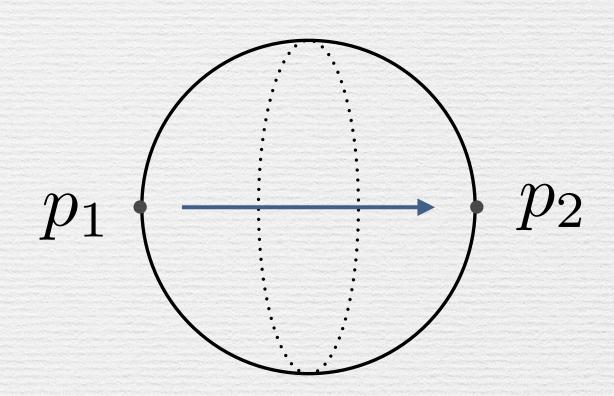
• It is a common tangent weight  $\lambda \in T_{\alpha}X$  and  $\lambda \in -T_{\beta}X$ .

## Example

- Consider again  $X = \mathbb{CP}^1$ .
- There are two massive vacua  $p_1, p_2$ .



- Here  $\zeta$  is the Kahler parameter of  $\mathbb{CP}^1$ .
- The hyperplane arrangement is:



$$\operatorname{Crit}(h) = \mathbb{CP}^1$$
 $m \in \mathbb{R}$ 

$$Crit(h) = \{p_1, p_2\}$$

# Part II: Ground States on an Elliptic Curve

### The Setup

• Place the theory on the real line times an elliptic curve  $\,E_{ au}\,$  .



- Ramond-Ramond boundary conditions for fermions.
- This preserves a quantum mechanics with 4 supercharges.

#### Flat Connections

- We have already introduced mass parameters  $m \in \mathfrak{t} \cong \mathbb{R}^{\mathrm{rk}\,T}$ .
- We can introduce a background flat connection on  $\,E_{ au}\,$  ,

$$a \sim a + \mu + \tau \nu$$

- They parametrise are  $\mathrm{rk}(T)$  -dimensional torus  $E_T \cong (E_\tau)^{\mathrm{rk}(T)}$ .
- The total moduli space of parameters m, a is

$$M_T \cong (\mathbb{R} \times E_\tau)^{\operatorname{rk}(T)}$$

#### Supersymmetric Ground States

• We are interested in supersymmetric ground states.



• For generic m, a, they are in 1-1 correspondence with fixed points.

$$\Psi_{\alpha}$$
  $\alpha = 1, \dots, N$ 

• What is the general dependence on the parameters m,a?

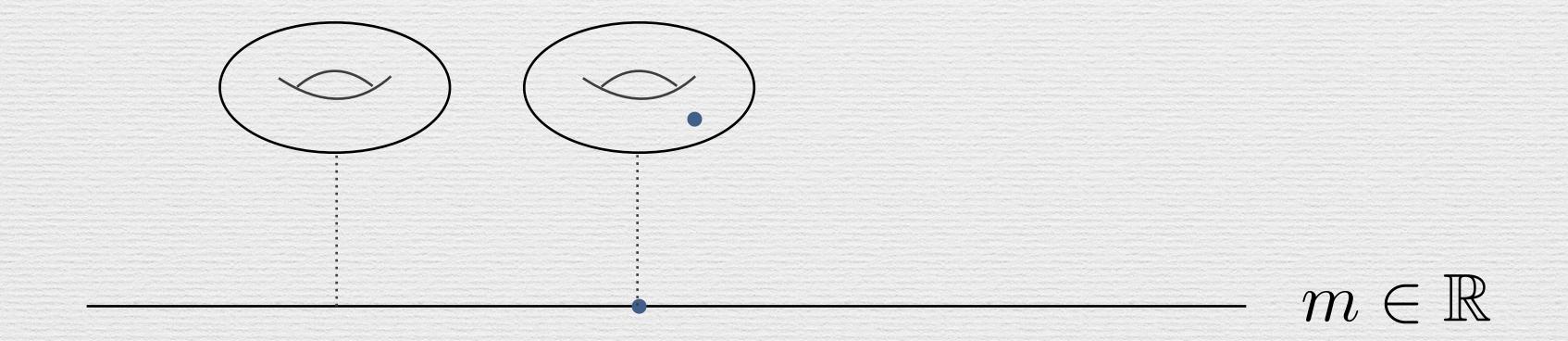
#### The Berry Connection

- The dependence is controlled by a supersymmetric Berry Connection:
  - $\star$  A connection on a principle  $\mathrm{SU}(N)$  bundle P.
  - \* A  $\mathfrak{t}^*$ -valued section  $\phi$  of  $\mathrm{Ad}(P)$ .
  - \* They solve a set of generalised Bogomolny equations on

$$(\mathbb{R} \times E_{\tau})^{\operatorname{rk} T}$$

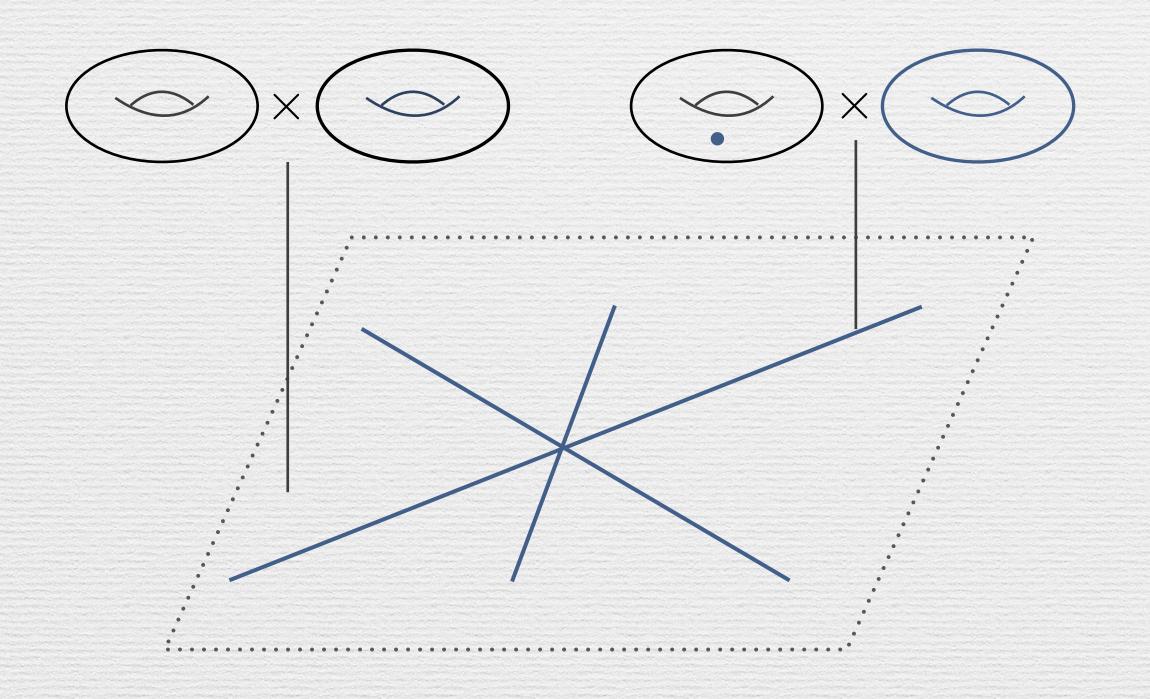
### Example

- Consider again  $X = \mathbb{CP}^1$ .
- There are two supersymmetric ground states  $\Psi_1,\Psi_2$  .
- The Berry connection is a smooth  $\mathrm{SU}(2)$  monopole on  $\mathbb{R} \times E_{ au}$  .



#### General Structure

• Smooth monopoles centred on codimension-3 loci in  $(\mathbb{R} \times E_{\tau})^{\mathrm{rk} T}$ .



$$\lambda \cdot m = 0$$

$$\lambda \cdot a \in \mathbb{Z} + \tau \mathbb{Z}$$

• They project onto the hyperplane arrangement in  $\mathfrak{t} \cong \mathbb{R}^{\mathrm{rk}\,T}$ .

#### Algebraic Picture

- The Bogomolny equations include  $[D_m + \phi, D_{\bar{a}}] = 0$ .
- A rank N holomorphic vector bundle  $\mathcal{E}$  on each fiber

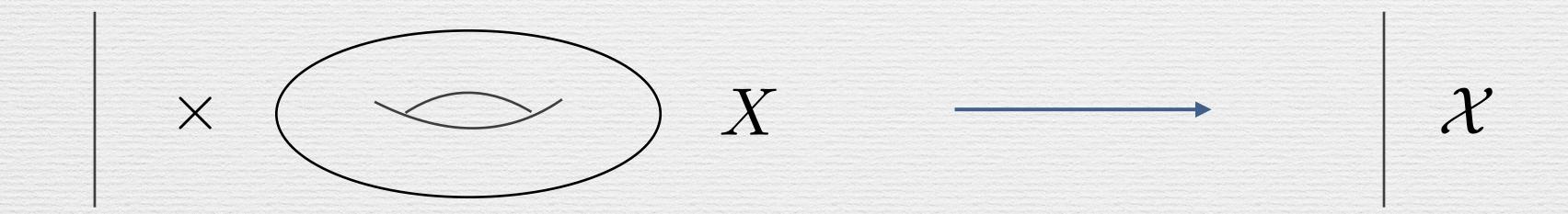
$$\{m\} \times E_T \cong (E_\tau)^{\operatorname{rk}(T)}$$

- The holomorphic bundle  ${\mathcal E}$  is piecewise constant in m .
- This suggests a more algebraic approach.

# Part III: Algebraic Description

## Supersymmetric QM

• Reduce three-dimensional theory on  $E_{ au}$  to obtain a supersymmetric quantum mechanics on  $\mathbb R$  .



- Quantum mechanics has infinite-dimensional target  $\mathcal{X} = \operatorname{Map}(E_{\tau}, X)$
- Coupled to background vectormultiplets for the  $S^1 \times S^1$  and T actions on  $\mathcal X$  .

#### Supersymmetric Ground States

- Alternative presentation of supersymmetric ground states.
- They are representatives of cohomology classes of the supercharge

$$Q = e^{-h}(d + \iota_V)e^h$$

Witten '82

• Acts on forms on  $\mathcal{X}$  where

- \* h is the moment map for the 1PS  $T_m$ -action on  $\mathcal X$ .
- \* V is a complex vector field generating  $S^1 \times S^1$  action and the T action with parameter a .

#### Localisation

- Supersymmetric localisation leads to a simpler description.
- Supersymmetric ground state are representatives of cohomology classes of the equivariant differential  $d + \iota_V$  on  $\mathrm{Crit}(h) \subset \mathcal{X}$ .
- The outcome depends on a face of the hyperplane arrangement.
- I will focus on the origin m = 0.

#### Zero Mass

- In this case,  $\operatorname{Crit}(h) = \mathcal{X}$ .
- For generic flat connection a , the fixed points of the vector field V are constant maps  $E_{ au} \to p_{\alpha}$  .
- A basis of supersymmetric ground states from equivariant Euler class of tangent space to constant maps in  $\mathcal{X}$ :

$$\Psi_{\alpha} \sim \prod_{\lambda \in T_{\alpha}X} \prod_{n,m \in \mathbb{Z}} (n + m\tau + \lambda \cdot a) \sim \prod_{\lambda \in T_{\alpha}X} \frac{\vartheta_{1}(\lambda \cdot a, \tau)}{\eta(\tau)}$$

#### Spectral Data

- They are sections of holomorphic line bundles  $L_1, \ldots, L_N$  on  $E_T$ .
- Combine into section of holomorphic line bundle on

$$E_T(X) := \bigsqcup_{\alpha=1}^{N} E_T^{(\alpha)} / \Delta$$

- \* Identical copies of  $E_T$  associated to each vacuum  $\alpha$ .
- + Pairs  $E_T^{(\alpha)}$  and  $E_T^{(\beta)}$  are identified at loci  $\lambda \cdot a \in \mathbb{Z} + \tau \mathbb{Z}$ .
- \* Here  $\lambda$  labels the hyperplane where domain walls  $p_{\alpha} \to p_{\beta}$  are tensionless.

### Example

- Consider again  $X = \mathbb{CP}^1$ .
- Supersymmetric ground states at m = 0:

$$\Psi_1 \sim \frac{\vartheta_1(2a,\tau)}{\eta(\tau)} \qquad \Psi_2 \sim \frac{\vartheta_1(-2a,\tau)}{\eta(\tau)}$$

• They glue to section of holomorphic line bundle on

$$E_T(X) = E_T^{(1)} \sqcup E_T^{(2)} / \Delta$$

where  $\Delta$  identifies the two copies at  $a \in \mathbb{Z} + \tau \mathbb{Z}$ .

#### General Picture

- $\bullet$  On a general face of the hyperplane arrangement containing m .
- Supersymmetric ground states transform as section of a line bundle on the equivariant elliptic cohomology variety:

$$\mathrm{Ell}_T(X^{T_m})$$

• Is this a type of spectral data for the supersymmetric Berry connection?

## Part IV: Boundary Conditions

### Boundary Conditions

• Boundary condition preserving at least  $\mathcal{N} = (0, 2)$  supersymmetry.



- $\bullet$  Assume it preserves the flavour symmetry T.
- It may have 't Hooft anomalies for T.
- It may or may not be compatible with mass parameters m.

## Boundary Amplitudes

• Now consider boundary amplitudes:



- They behave like  $\mathcal{N}=(0,2)$  elliptic genera.
- They may be computed exactly by supersymmetric localisation.

Sugiyama-Yoshida '20

· Quasi-periodicities fixed by boundary 't Hooft anomalies.

## Elliptic Cohomology Classes

- $\bullet$  Suppose the boundary condition is compatible with mass m.
- The boundary amplitudes glue to a section of a line bundle on.

$$\mathrm{Ell}_T(X^{T_m})$$

• Boundary conditions produce classes in equivariant elliptic cohomology.

## N = 4 Supersymmetry

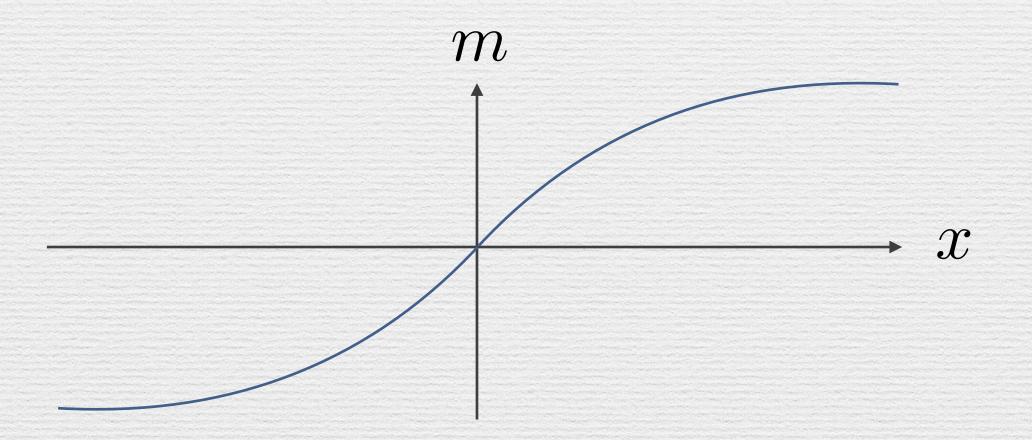
- Special interest in N = 4 supersymmetry, broken to N = 2.
- Interesting classes of boundary conditions  $B_{lpha}$  labelled by vacua lpha .

  MB-Dimofte-Gaiotto-Hilburn '16
- We constructed boundary conditions corresponding to
  - \* Attracting sets.
  - \* Stable envelopes.

    Aganagic-Okounkov '16
- They are exchanged under 3d mirror symmetry.

#### Janus Interfaces

- We also studied correlation functions of Janus interfaces.
- This is a position dependent mass m(x) along  $\mathbb R$  that interpolates between two faces of the hyperplane arrangement at  $x \to \pm \infty$ .



• This reproduces the construction of elliptic R-matrices.

## Thank you!