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String Theory and the Very Early Universe

Robert Brandenberger,
McGill University and IHEP (Beijing)

February 9, 2009

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Current Paradigm for Early Universe Cosmology

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The **Inflationary Universe Scenario** is the current paradigm of early universe cosmology.

Successes:

- Solves horizon problem
- Solves flatness problem
- Solves size/entropy problem
- Provides a causal mechanism of generating **primordial cosmological perturbations** (Chibisov & Mukhanov, 1981).

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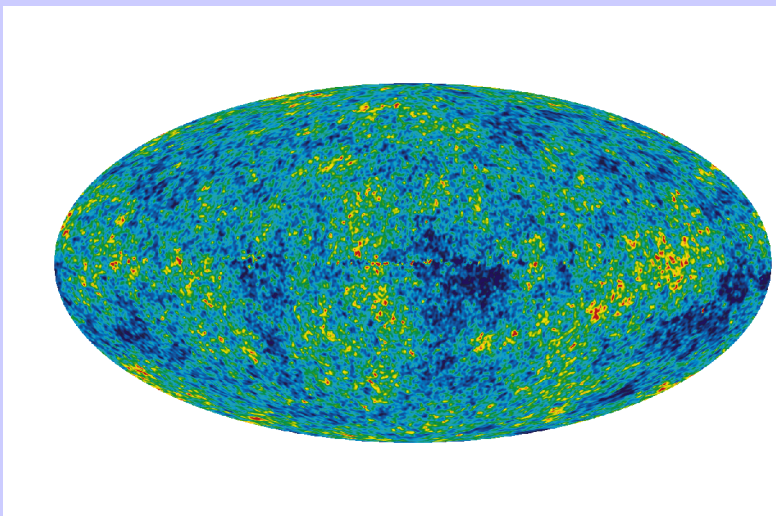
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Credit: NASA/WMAP Science Team

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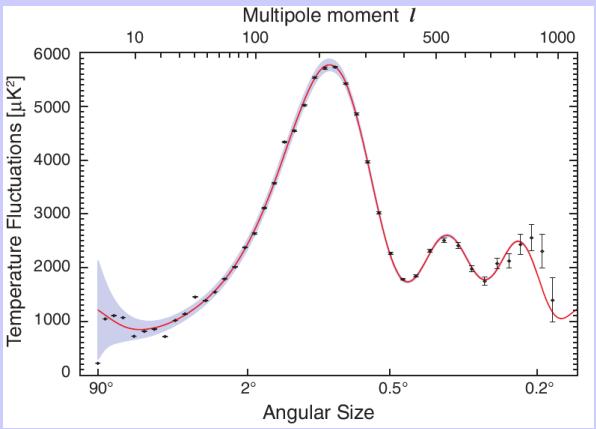
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- In spite of the phenomenological successes, the inflationary scenario suffers from several **conceptual problems**.
- In light of these problems we need to look for input from new fundamental physics to construct a new theory which will overcome these problems.
- Question: Can **Superstring theory** lead to a new and improved paradigm?
- Question: Can this new paradigm be **tested** in cosmological observations?

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Review of Inflationary Cosmology

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Context:

- General Relativity
- Scalar Field Matter

$$\text{Metric : } ds^2 = dt^2 - a(t)^2 d\mathbf{x}^2 \quad (1)$$

Inflation:

- phase with $a(t) \sim e^{tH}$
- requires matter with $p \sim -\rho$
- requires a slowly rolling scalar field φ
- - in order to have a potential energy term
- - in order that the potential energy term dominates sufficiently long

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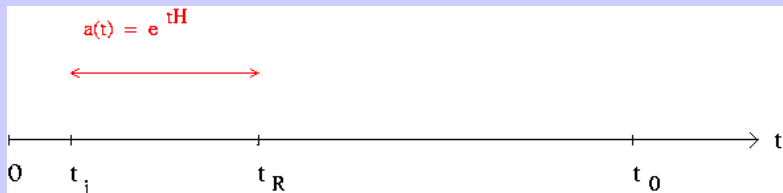
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Time line of inflationary cosmology:



- t_i : inflation begins
- t_R : inflation ends, reheating

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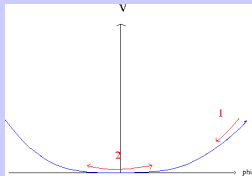
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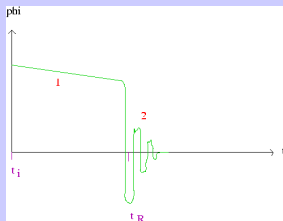
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Matter scalar field:

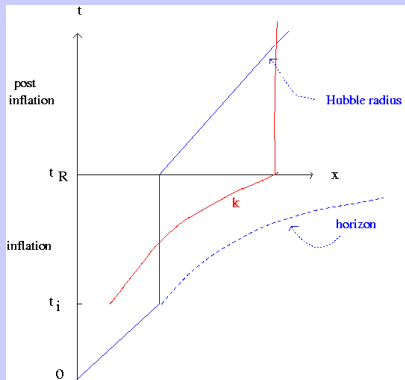


Scalar field evolution:



Review of Inflationary Cosmology II

Space-time sketch of inflationary cosmology:



Note:

- $H = \frac{\dot{a}}{a}$

- curve labelled by k : wavelength of a fluctuation

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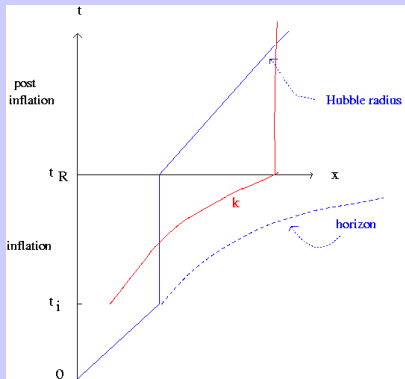
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- inflation renders the universe large, homogeneous and spatially flat
- classical matter redshifts \rightarrow matter vacuum remains
- quantum vacuum fluctuations: seeds for the observed structure [Chibisov & Mukhanov, 1981]
- sub-Hubble \rightarrow locally causal

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Conceptual Problems of Inflationary Cosmology

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- Nature of the scalar field φ (the “inflaton”)
- Conditions to obtain inflation (initial conditions, slow-roll conditions, graceful exit and reheating)
- Amplitude problem
- **Trans-Planckian problem**
- **Singularity problem**
- Cosmological constant problem
- **Applicability of General Relativity**

Trans-Planckian Problem

R.B., hep-ph/9910410

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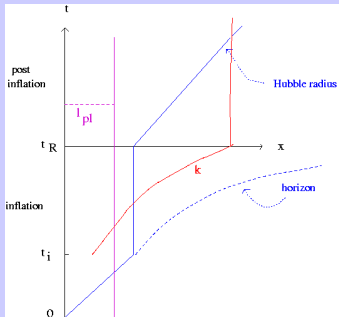
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- **Success of inflation:** At early times scales are inside the Hubble radius \rightarrow causal generation mechanism is possible.
- **Problem:** If time period of inflation is more than $70H^{-1}$, then $\lambda_p(t) < l_{pl}$ at the beginning of inflation
- \rightarrow new physics **MUST** enter into the calculation of the fluctuations.

Trans-Planckian Problem

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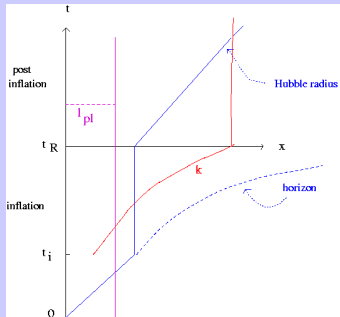
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Trans-Planckian Window of Opportunity

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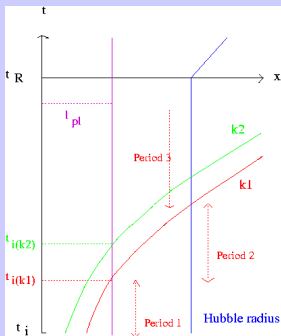
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- If evolution in Period I is non-adiabatic, then scale-invariance of the power spectrum will be lost [J. Martin and RB, 2000]
- → Planck scale physics testable with cosmological observations!

Singularity Problem

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- Standard cosmology: Penrose-Hawking theorems → initial singularity → incompleteness of the theory.
- Inflationary cosmology: In scalar field-driven inflationary models the initial singularity persists [Borde and Vilenkin] → **incompleteness of the theory**.

Penrose-Hawking theorems:

- Ass: i) Einstein action, 2) weak energy conditions
 $\rho > 0, \rho + 3p \geq 0$
- → space-time is geodesically incomplete.

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Penrose-Hawking theorems:

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Cosmological Constant Problem

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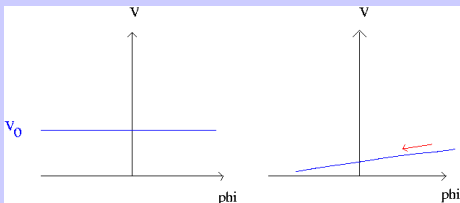
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- Quantum vacuum energy does not gravitate.
- Why should the almost constant $V(\varphi)$ gravitate?

$$\frac{V_0}{\Lambda_{obs}} \sim 10^{120} \quad (2)$$

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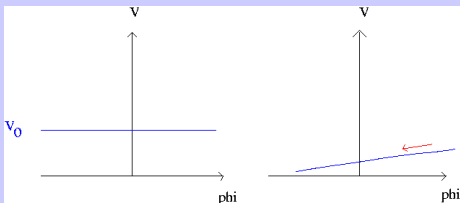
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Applicability of GR

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- In all approaches to quantum gravity, the Einstein action is only the leading term in a low curvature expansion.
- Correction terms may become dominant at much lower energies than the Planck scale.
- Correction terms will dominate the dynamics at high curvatures.
- The energy scale of inflation models is typically $\eta \sim 10^{16} \text{GeV}$.
- $\rightarrow \eta$ too close to m_{pl} to trust predictions made using GR.

Zones of Ignorance

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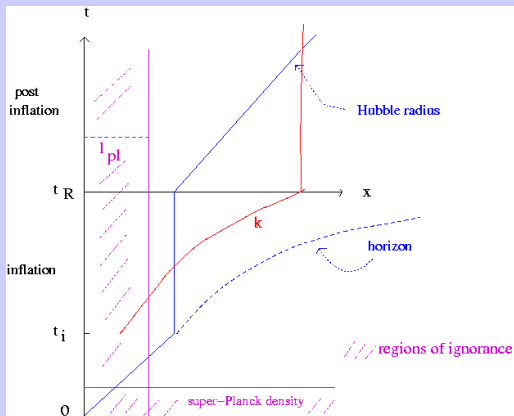
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- Current realizations of inflation have serious conceptual problems.
- We need a new paradigm of very early universe cosmology based on new fundamental physics.
- **Hypothesis:** New paradigm based on **Superstring Theory**.
- The new paradigm of early universe cosmology may **not** involve inflation.
- New cosmological model motivated by superstring theory: **String Gas Cosmology** (SGC) [R.B. and C. Vafa, 1989]
- **New structure formation scenario** emerges from SGC [A. Nayeri, R.B. and C. Vafa, 2006].

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String Gas Cosmology makes **testable predictions** for cosmological observations

- **Blue tilt** in the spectrum of **gravitational waves** [R.B., A. Nayeri, S. Patil and C. Vafa, 2006]
- **Line discontinuities** in **CMB anisotropy maps** [N. Kaiser and A. Stebbins, 1984]
- Line discontinuities may have **junctions**
- Line discontinuities can be detected using the **CANNY edge detection algorithm** [S. Amsel, J. Berger and R.B., 2007, A. Stewart and R.B., 2008, R. Danos and R.B., 2008]

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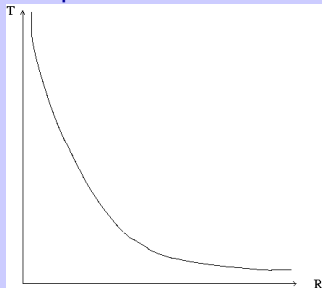
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Temperature-size relation in standard cosmology



Principles

R.B. and C. Vafa, *Nucl. Phys. B316:391 (1989)*

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Idea: make use of the **new symmetries** and **new degrees of freedom** which string theory provides to construct a new theory of the very early universe.

Assumption: Matter is a gas of fundamental strings

Assumption: Space is compact, e.g. a torus.

Key points:

- **New degrees of freedom:** string oscillatory modes
- Leads to a **maximal temperature** for a gas of strings, the Hagedorn temperature
- **New degrees of freedom:** string winding modes
- Leads to a **new symmetry:** physics at large R is equivalent to physics at small R

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T-Duality

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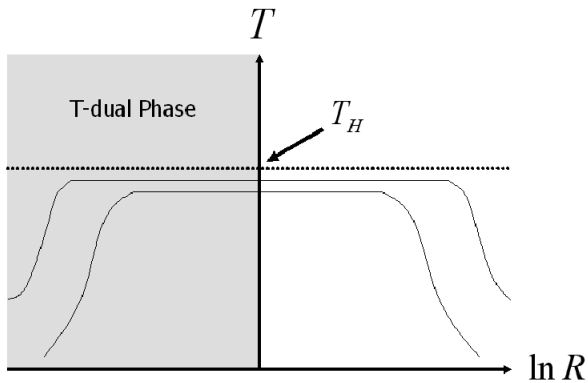
T-Duality

- Momentum modes: $E_n = n/R$
- Winding modes: $E_m = mR$
- Duality: $R \rightarrow 1/R$ $(n, m) \rightarrow (m, n)$
- Mass spectrum of string states unchanged
- Symmetry of vertex operators
- Symmetry at non-perturbative level \rightarrow existence of D-branes

Adiabatic Considerations

R.B. and C. Vafa, *Nucl. Phys. B316:391 (1989)*

Temperature-size relation in string gas cosmology



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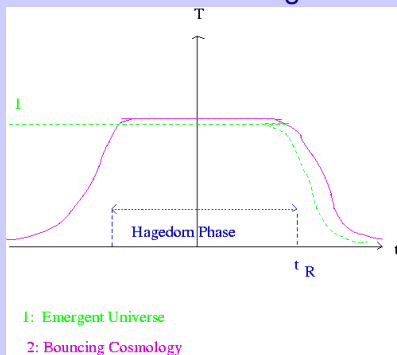
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Assume some action gives us $R(t)$



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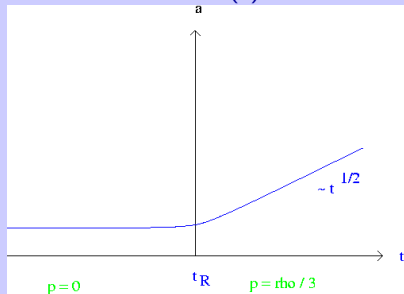
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We will thus consider the following background dynamics for the scale factor $a(t)$:



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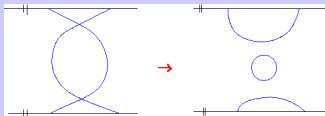
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Conclusions

- Begin with all 9 spatial dimensions small, initial temperature close to T_H → winding modes about all spatial sections are excited.
- Expansion of any one spatial dimension requires the annihilation of the winding modes in that dimension.



- Decay only possible in three large spatial dimensions.
- → dynamical explanation of why there are exactly three large spatial dimensions.

Note: this argument assumes constant dilaton [R. Danos, A. Frey and A. Mazumdar]

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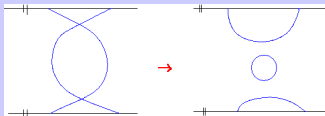
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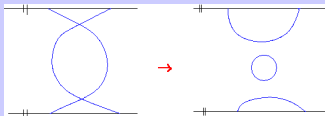
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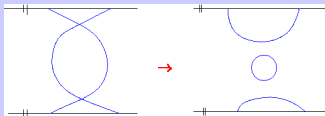
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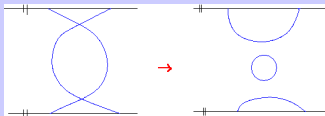


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Moduli Stabilization in SGC

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Size Moduli [S. Watson, 2004; S. Patil and R.B., 2004, 2005]

- winding modes prevent expansion
- momentum modes prevent contraction
- $\rightarrow V_{\text{eff}}(R)$ has a minimum at a finite value of R , $\rightarrow R_{\text{min}}$
- in heterotic string theory there are **enhanced symmetry states** containing both momentum and winding which are massless at R_{min}
- $\rightarrow V_{\text{eff}}(R_{\text{min}}) = 0$
- \rightarrow **size moduli stabilized** in Einstein gravity background

Shape Moduli [E. Cheung, S. Watson and R.B., 2005]

- enhanced symmetry states
- \rightarrow harmonic oscillator potential for θ
- \rightarrow shape moduli stabilized

Dilaton stabilization in SGC

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Conclusions

- The only remaining modulus is the dilaton
- Make use of **gaugino condensation** to give the dilaton a potential with a unique minimum
- → dilaton is stabilized
- Dilaton stabilization is consistent with size stabilization [R. Danos, A. Frey and R.B., 2008]

Dilaton stabilization in SGC

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Theory of Cosmological Perturbations: Basics

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Cosmological fluctuations connect early universe theories with observations

- Fluctuations of **matter** → large-scale structure
- Fluctuations of **metric** → CMB anisotropies
- N.B.: Matter and metric fluctuations are coupled

Key facts:

- 1. Fluctuations are small today on large scales
- → fluctuations were very small in the early universe
- → can use **linear perturbation theory**
- 2. Sub-Hubble scales: matter fluctuations dominate
- Super-Hubble scales: metric fluctuations dominate

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Quantum Theory of Linearized Fluctuations

V. Mukhanov, H. Feldman and R.B., *Phys. Rep.* 215:203 (1992)

Step 1: Metric including fluctuations

$$ds^2 = a^2[(1 + 2\Phi)d\eta^2 - (1 - 2\Phi)d\mathbf{x}^2] \quad (3)$$

$$\varphi = \varphi_0 + \delta\varphi \quad (4)$$

Note: Φ and $\delta\varphi$ related by Einstein constraint equations

Step 2: Expand the action for matter and gravity to second order about the cosmological background:

$$S^{(2)} = \frac{1}{2} \int d^4x ((v')^2 - v_{,i}v^{,i} + \frac{z''}{z}v^2) \quad (5)$$

$$v = a(\delta\varphi + \frac{z}{a}\Phi) \quad (6)$$

$$z = a\frac{\varphi_0'}{\mathcal{H}} \quad (7)$$

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Step 3: Resulting equation of motion (Fourier space)

$$v_k'' + \left(k^2 - \frac{z''}{z}\right)v_k = 0 \quad (8)$$

Features:

- **oscillations** on sub-Hubble scales
- **squeezing** on super-Hubble scales $v_k \sim z$

Quantum vacuum initial conditions:

$$v_k(\eta_i) = (\sqrt{2k})^{-1} \quad (9)$$

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Requirements for a model which agrees with observations:

- **Scale-invariant spectrum** of fluctuations.
- **Generation on sub-Hubble** scales.
- **Frozen propagation on super-Hubble** scales.

Structure formation in inflationary cosmology

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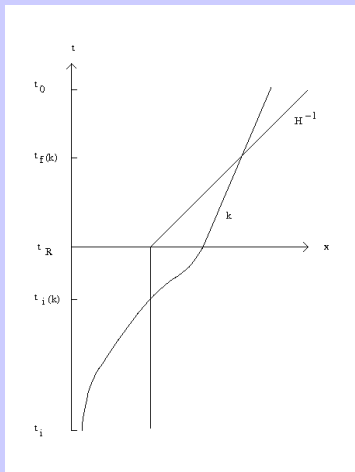
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N.B. Perturbations originate as quantum vacuum fluctuations.

Background for string gas cosmology

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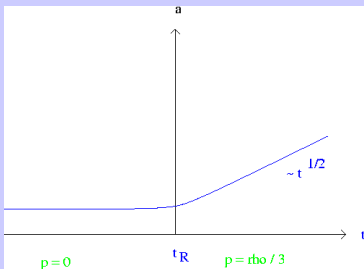
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Structure formation in string gas cosmology

A. Nayeri, R.B. and C. Vafa, *Phys. Rev. Lett.* 97:021302 (2006)

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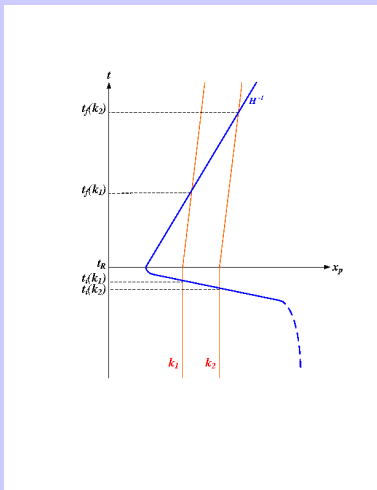
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N.B. Perturbations originate as thermal string gas fluctuations.

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1. Calculate matter correlation functions in the Hagedorn phase (neglecting the metric fluctuations)
2. For fixed k , convert the matter fluctuations to metric fluctuations at Hubble radius crossing $t = t_i(k)$
3. Evolve the metric fluctuations for $t > t_i(k)$ using the usual theory of cosmological perturbations

Extracting the Metric Fluctuations

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Ansatz for the metric including cosmological perturbations and gravitational waves:

$$ds^2 = a^2(\eta) \left((1 + 2\Phi) d\eta^2 - [(1 - 2\Phi)\delta_{ij} + h_{ij}] dx^i dx^j \right). \quad (10)$$

Inserting into the perturbed Einstein equations yields

$$\langle |\Phi(k)|^2 \rangle = 16\pi^2 G^2 k^{-4} \langle \delta T^0_0(k) \delta T^0_0(k) \rangle, \quad (11)$$

$$\langle |h(k)|^2 \rangle = 16\pi^2 G^2 k^{-4} \langle \delta T^i_j(k) \delta T^i_j(k) \rangle. \quad (12)$$

Power Spectrum of Cosmological Perturbations

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Key ingredient: For **thermal fluctuations**:

$$\langle \delta\rho^2 \rangle = \frac{T^2}{R^6} C_V. \quad (13)$$

Key ingredient: For **string thermodynamics** in a compact space

$$C_V \approx 2 \frac{R^2/\ell_s^3}{T(1 - T/T_H)}. \quad (14)$$

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Power spectrum of cosmological fluctuations

$$P_{\Phi}(k) = 8G^2 k^{-1} \langle |\delta\rho(k)|^2 \rangle \quad (15)$$

$$= 8G^2 k^2 \langle (\delta M)^2 \rangle_R \quad (16)$$

$$= 8G^2 k^{-4} \langle (\delta\rho)^2 \rangle_R \quad (17)$$

$$= 8G^2 \frac{T}{\ell_s^3} \frac{1}{1 - T/T_H} \quad (18)$$

Key features:

- **scale-invariant** like for inflation
- **slight red tilt** like for inflation

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- 1. Evolution for $t > t_i(k)$: $\Phi \simeq \text{const}$ since the equation of state parameter $1 + w$ stays the same order of magnitude **unlike in inflationary cosmology**.
- 2. Squeezing of the fluctuation modes takes place on super-Hubble scales **like in inflationary cosmology** \rightarrow **acoustic oscillations** in the CMB angular power spectrum
- 3. In a dilaton gravity background the dilaton fluctuations dominate \rightarrow different spectrum [R.B. et al, 2006; Kaloper, Kofman, Linde and Mukhanov, 2006]

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Spectrum of Gravitational Waves

R.B., A. Nayeri, S. Patil and C. Vafa, *Phys. Rev. Lett.* (2007)

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$$P_h(k) = 16\pi^2 G^2 k^{-1} \langle |T_{ij}(k)|^2 \rangle \quad (19)$$

$$= 16\pi^2 G^2 k^{-4} \langle |T_{ij}(R)|^2 \rangle \quad (20)$$

$$\sim 16\pi^2 G^2 \frac{T}{\ell_s^3} (1 - T/T_H) \quad (21)$$

Key ingredient for **string thermodynamics**

$$\langle |T_{ij}(R)|^2 \rangle \sim \frac{T}{\ell_s^3 R^4} (1 - T/T_H) \quad (22)$$

Key features:

- scale-invariant (like for inflation)
- slight blue tilt (unlike for inflation)

Spectrum of Gravitational Waves

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Conclusions

- Static Hagedorn phase (including static dilaton) → new physics required.
- $C_V(R) \sim R^2$ obtained from a thermal gas of strings provided there are winding modes which dominate.
- Cosmological fluctuations in the IR are described by Einstein gravity.

Note: Specific higher derivative toy model: T. Biswas, R.B., A. Mazumdar and W. Siegel, 2006

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Network of cosmic superstrings

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- Remnant of the Hagedorn phase: **network of cosmic superstrings**
- This string network will be present at **all times** and will achieve a **scaling solution** like cosmic strings forming during a phase transition.
- **Scaling Solution**: The network of strings looks statistically the same at all times when scaled to the Hubble radius.

Kaiser-Stebbins Effect

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Space perpendicular to a string is **conical** with **deficit angle**

$$\alpha = 8\pi G\mu, \quad (23)$$

Photons passing by the string undergo a **relative Doppler shift**

$$\frac{\delta T}{T} = 8\pi\gamma(v)vG\mu, \quad (24)$$

→ network of **line discontinuities** in CMB anisotropy maps

.

*N.B. characteristic scale: comoving Hubble radius at the time of recombination → need **good angular resolution** to detect these edges.*

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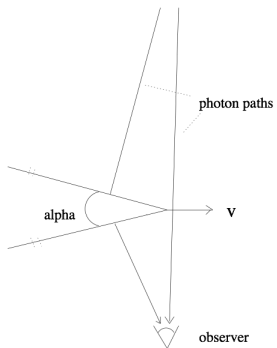
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Gaussian temperature map

$10^0 \times 10^0$ map of the sky at 1.5' resolution (South Pole Telescope specifications)

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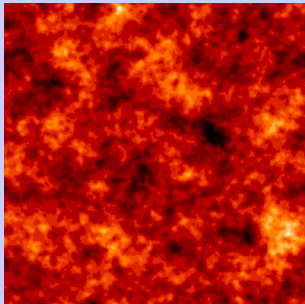
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Cosmic string temperature map

$10^0 \times 10^0$ map of the sky at 1.5' resolution

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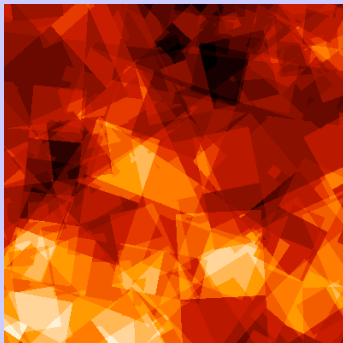
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This signal is superimposed on the Gaussian map. The relative power of the string signature depends on $G\mu$ and is bound to contribute less than 10% of the power (L. Pogosian & M. Wyman).

CANNY edge detection algorithm

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- Challenge: pick out the string signature from the Gaussian "noise" which has a much larger amplitude
- New technique: use **CANNY edge detection algorithm** [Canny, 1986]
- Idea: find edges across which the gradient is in the correct range to correspond to a Kaiser-Stebbins signal from a string
- Step 1: generate "Gaussian" and "Gaussian plus strings" CMB anisotropy maps: size and angular resolution of the maps are free parameters, flat sky approximation, cosmic string toy model in which a fixed number of straight string segments is laid down at random in each Hubble volume in each Hubble time step between t_{rec} and t_0 .

Temperature map Gaussian + strings

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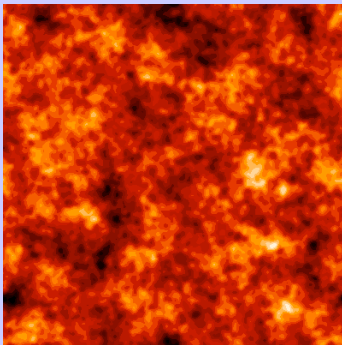
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CANNY algorithm II

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- Step 2: run the CANNY algorithm on the temperature maps to produce **edge maps**.
- Step 3: Generate histogram of edge lengths
- Step 4: Use Fisher combined probability test to check for difference compared to a Gaussian distribution.

Edge map

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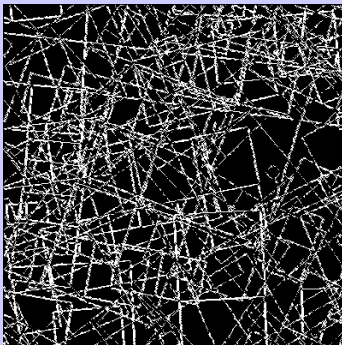
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- For South Pole Telescope (SPT) specification: limit $G\mu < 2 \times 10^{-8}$ can be set [A. Stewart and R.B., 2008, R. Danos and R.B., 2008]
- Anticipated SPT instrumental noise only insignificantly effects the limits [A. Stewart and R.B., 2008]
- WMAP data: limit $G\mu < 2 \times 10^{-7}$ can be set [E. Thewalt, in prep.]

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Conclusions :

- The **Inflationary scenario**, the current paradigm of early universe cosmology, has serious **conceptual problems**.
- This motivated the search for a new paradigm based on improved fundamental physics.
- **String Gas Cosmology**: Model of cosmology of the very early universe based on new degrees of freedom and new symmetries of superstring theory.
- SGC \rightarrow **nonsingular cosmology**
- SGC \rightarrow natural explanation of the number of large spatial dimensions.
- SGC \rightarrow **new scenario of structure formation**
- Scale invariant spectrum of cosmological fluctuations (like in inflationary cosmology).
- **Spectrum of gravitational waves** has a **small blue tilt** (unlike in inflationary cosmology).

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- Two **consistency relations** between the four observables (amplitudes and slopes of the spectra of cosmological perturbations and gravitational waves).
- SGC leaves behind a network of **cosmic superstrings**
- These cosmic superstrings give rise to **line discontinuities** in CMB anisotropy maps which can be probed using a **CANNY** edge detection algorithm.
- A specific signature are **string junctions**.

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Requirements on the Structure Formation Scenario Revisited

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The proposed alternative to inflation relies on:

- **New physics phase** of background cosmology with thermal equilibrium over 1 mm scales.
- **Holographic scaling** of the specific heat capacity.
- Applicability of linearized Einstein equations after the new physics phase on infrared scales.

Challenge to string theorists: construct an improved new physics phase with holographic scaling of C_V .

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- Find a **Background effective field theory** which yields quasi-static Hagedorn phase.
- Some progress on this issue: R.B., A. Frey and S. Kanno, 2007 (in the context of a model including a dynamical tachyon field)
- **Non-perturbative understanding** of the Hagedorn phase.
- Does String Gas Cosmology solve the horizon, flatness, entropy and size problems of standard cosmology?
- Is the model free of small-scale instabilities (Jeans instabilities)?

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