A simulation of a cosmic magnetic field, showing a complex, filamentary structure. The field is represented by a color map where blue indicates low field strength and yellow/red indicates high field strength. The structure is highly irregular and interconnected, with many small loops and filaments. The background is black, making the colored structures stand out.

Magnetic Field Generation in Cosmic Structure

Francesco Miniati
ETH-Zurich

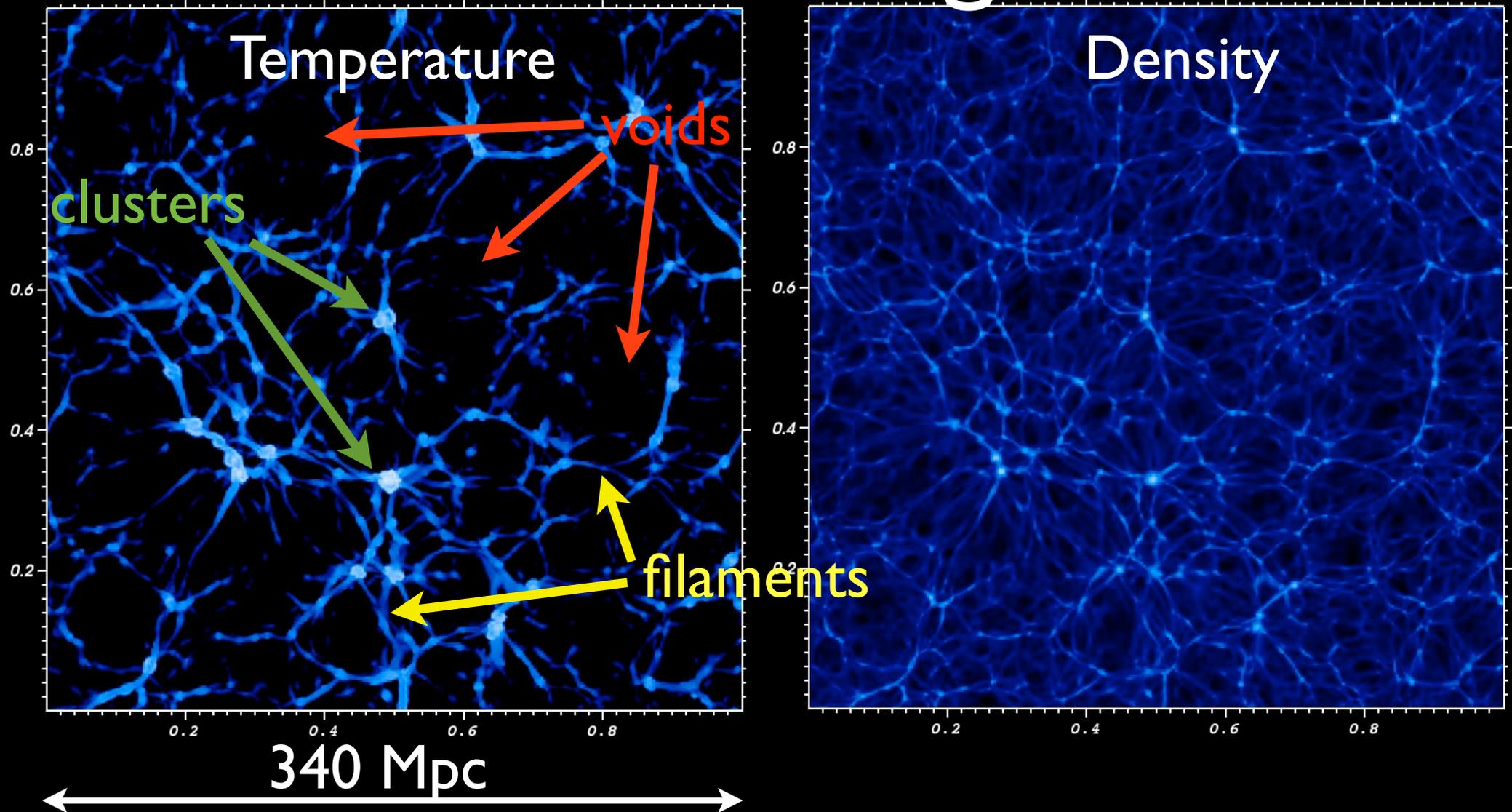
November 24 2011, IPMU (Kashiwa), Japan

Cosmic Magnetism

Most astrophysical bodies are magnetized:

- Main sequence stars 1-100 G.
- Galaxy, nearby galaxies and “high-redshift” galaxies: $B \sim 1-10 \mu\text{G}$ (Beck 1996, Bernet et al 2008)
- Clusters of Galaxies: $B \sim 0.1-10 \mu\text{G}$
- Filament of Galaxies: $B \sim \text{nG}$ (?)
- Cosmic Voids: $B \sim 0.1-0.01 \text{ fG}$

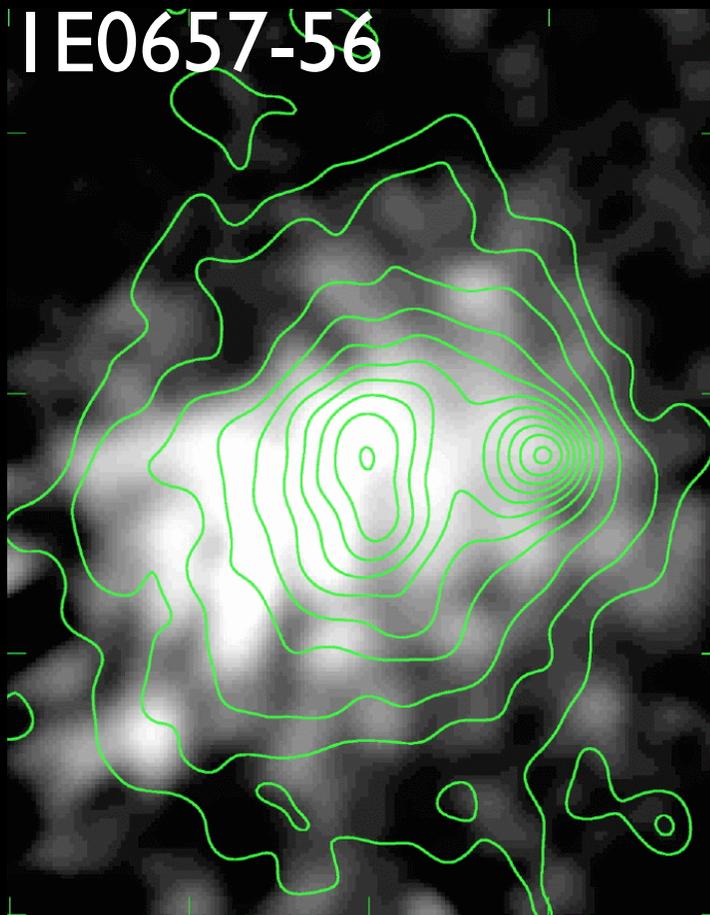
The LSS: voids, filaments and clusters of galaxies



Outline

- measurement of magnetic fields in the LSS, i.e. galaxy clusters, filaments and voids
- resistive generation mechanism
- conclusions

Probes of Magnetism in Galaxy Clusters



Radio Halo

Synchrotron Radiation

relativistic electrons

$$\frac{dN}{dE}(E) = N_0 E^{-s}$$

synchrotron flux

$$F_\nu \propto N_0 B^{\frac{1+s}{2}} \nu^{-\frac{s-1}{2}}$$

Probes of Magnetism

Rotation Measure

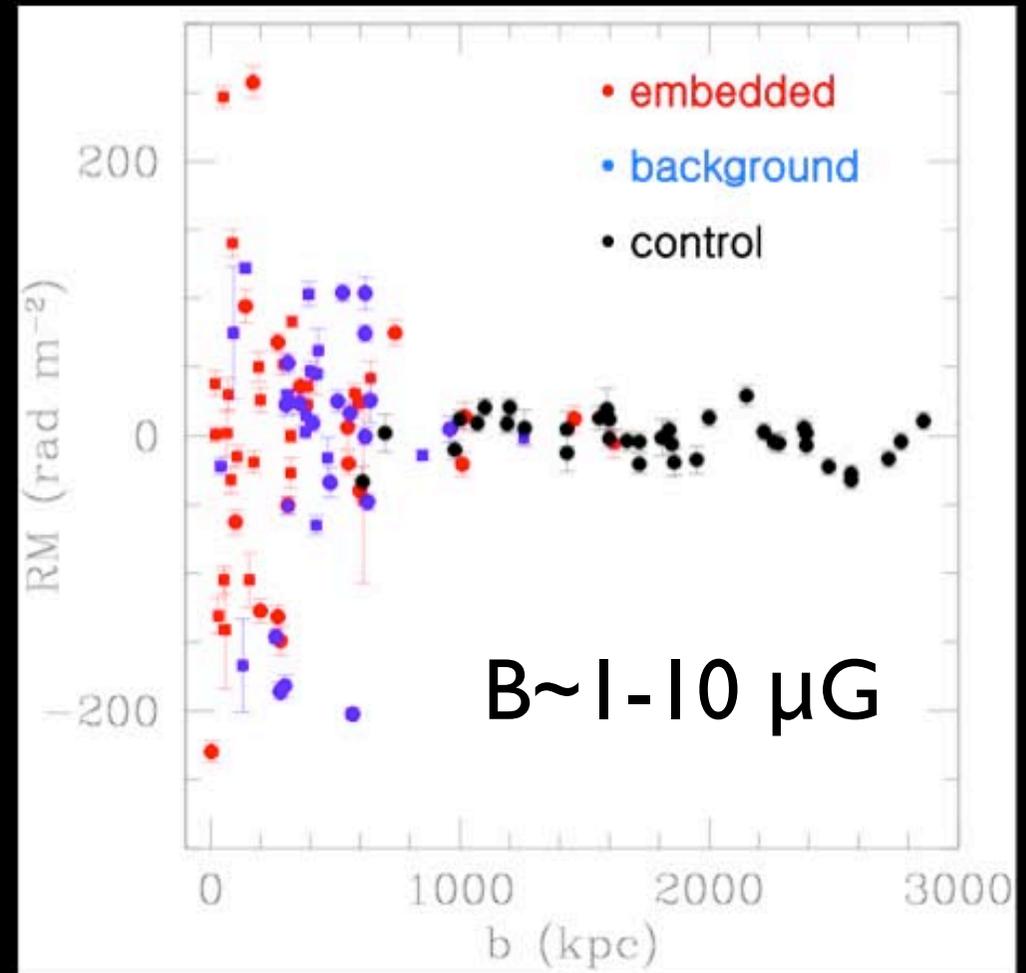
$$\chi = \chi_0 + \text{RM} \lambda^2$$

$$\text{RM} = 0.8 \int n_e \vec{B} \cdot d\vec{l}$$



$$\langle \text{RM} \rangle = 0$$

$$\sigma_{\text{RM}} \propto B \ell_c^{1/2}$$



Clarke et al (2001, 2004)

E.M. cascade in Voids

$$\gamma_{Blz} \gamma_{EBL} \rightarrow e^+ e^-$$

$$\Gamma_{e^\pm} = E_{\gamma_{Blz}} / 2m_e c^2 \approx 10^6 E_{\gamma_{Blz}, TeV}$$

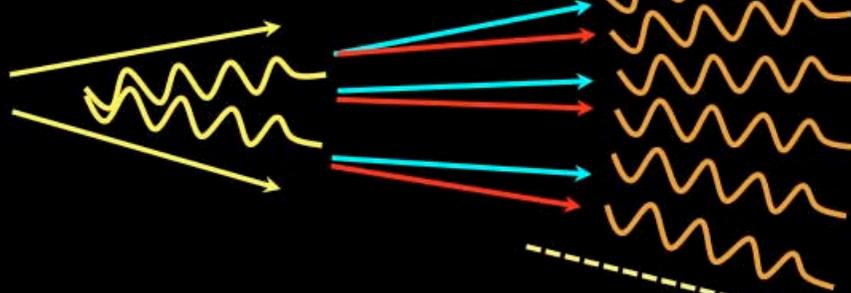
$$\ell_{\gamma} \sim 100 \text{ Mpc}$$

$$e^\pm \gamma_{CMB} \rightarrow e^\pm \gamma$$

$$E_\gamma \approx 1.3 \epsilon_{CMB} \Gamma_{e^\pm}^2 \approx 88 \left(E_{\gamma, Blz} / 10 \text{ TeV} \right)^2 \text{ GeV}$$

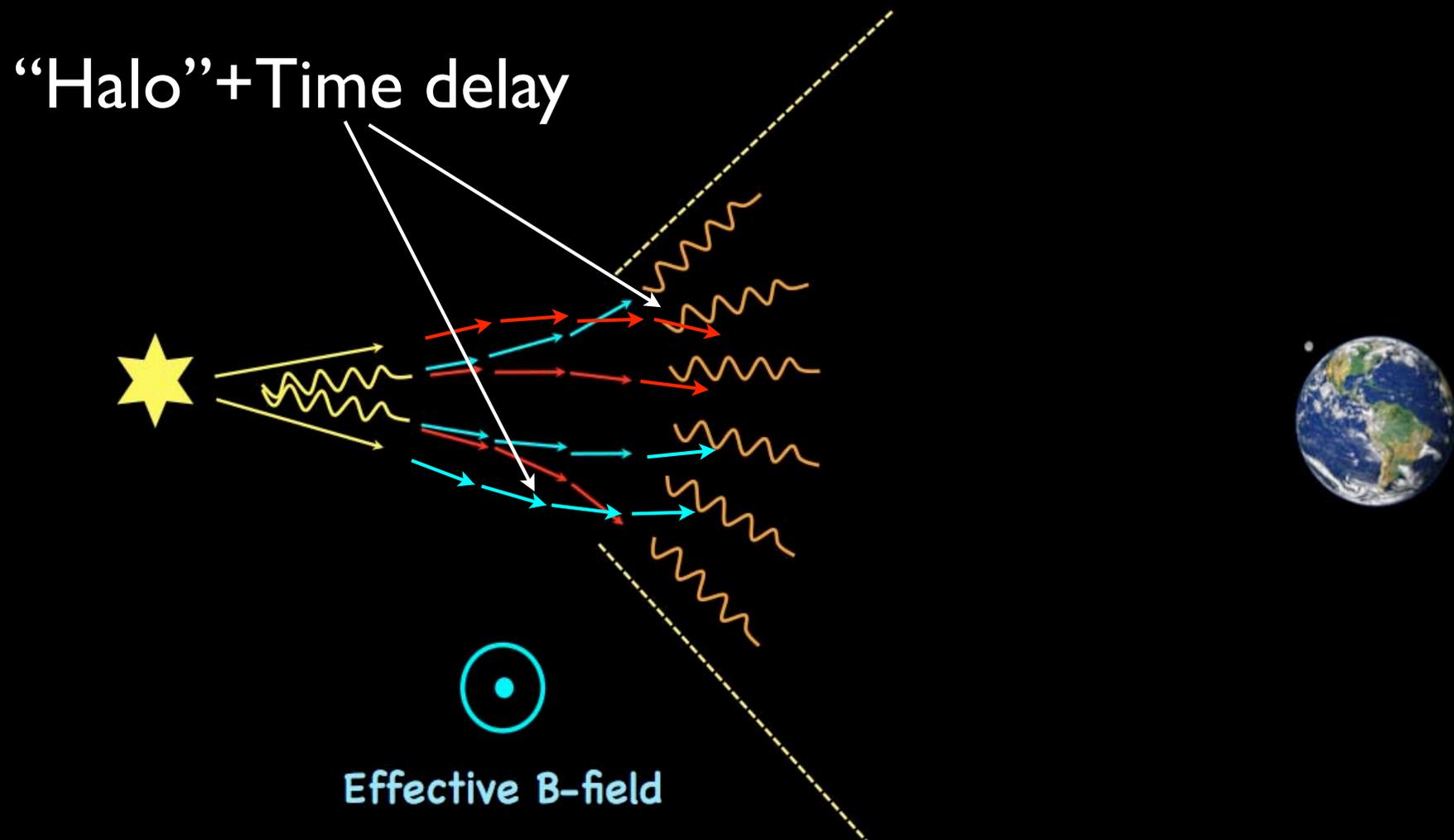
$$\ell_{IC} \sim 30 \text{ kpc}$$

Blazar



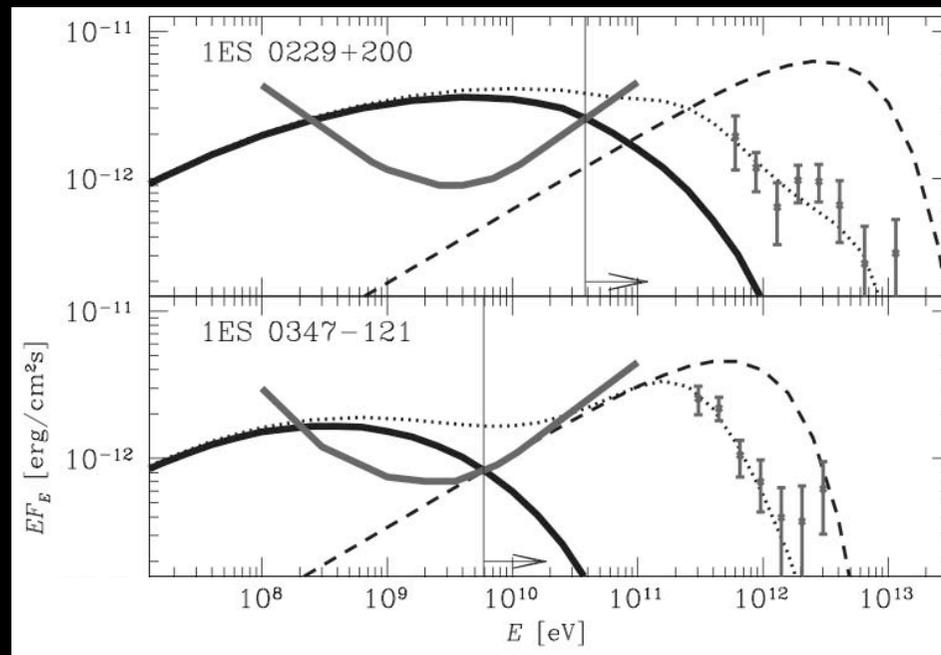
Typical energies of
reprocessed photons
1 – 100 GeV

Non vanishing B-field deflects secondary e^\pm



Magnetism in Voids

$$B \geq 10^{-18} - 10^{-17} \text{ Gauss}$$



Neronov and Vovk, Science 328 73 (2010), see also Tavecchio et al (2010),
Dermer et al. 2010, Taylor et al. 2011

Evolution of B

Electrostatic fields are weak and negligible!

$$\vec{E}' \approx 0 \Rightarrow \vec{E} = -\frac{\vec{u}}{c} \times \vec{B}$$

Induction equation

$$\frac{\partial \vec{B}}{\partial t} = -c \vec{\nabla} \times \vec{E} = \vec{\nabla} \times (\vec{u} \times \vec{B}) + \dot{\vec{B}}_{Source}$$

Growth timescale of B set by flow (\vec{u}) geometry.

B evolution in various environments

- In Galaxy Clusters the requirement on the initial seeds is not well constrained because the amplification is not well known
- In filaments the amplification due to turbulence is probably poor
- In voids the flow is divergent, so turbulence, if any, decays and we are most likely seeing are the fields directly seeded by the initial mechanism
- Initial seeds $B_0 \leq 1$ nG (e.g. Schleicher & Miniati 2011)

Generation Mechanisms for B

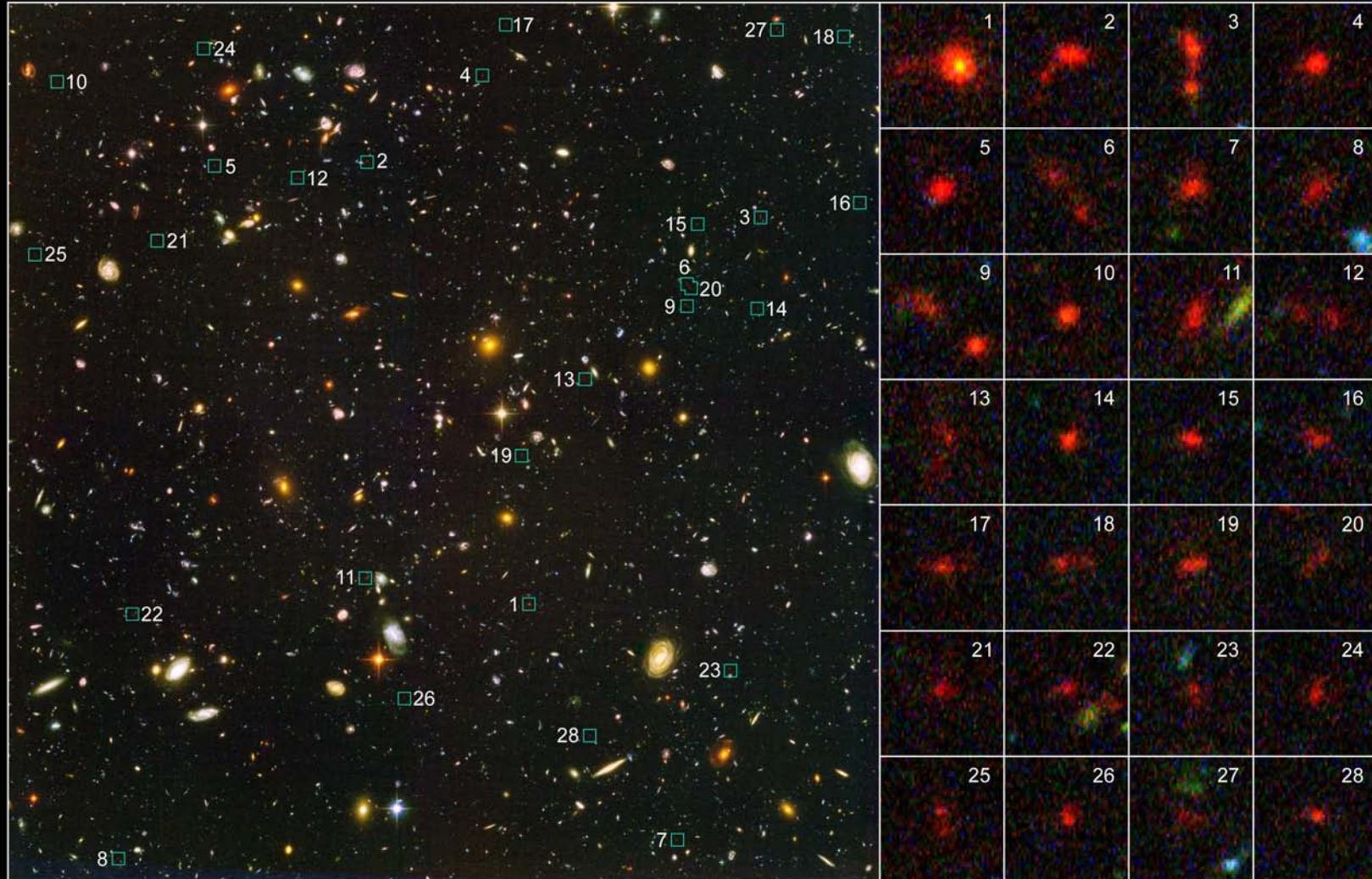
$$\frac{\partial \vec{B}}{\partial t} = -c \vec{\nabla} \times \vec{E} = \vec{\nabla} \times (\vec{u} \times \vec{B}) + \frac{\eta c^2}{4\pi} \nabla^2 \vec{B} + \dot{\vec{B}}_{Source}$$

- Plasma processes:
 - (i) resistive mechanism (Miniati & Bell 2011)
 - (ii) Biermann's battery (Subramanian et al. 1992, Kulsrud et al. 1997, Gnedin et al. 2000),
 - (iii) Weibel's instability (Schlickeiser & Shukla 2003, Medvedev et al 2004).
- Outflows:
 - (i) Galactic (Kronberg et al. 1999, Bertone et al. 2006, Donnert et al. 2009, Dubois & Teyssier 2010)
 - (ii) Jets from radiogalaxies (e.g., Furlanetto & Loeb 2001, Kronberg et al 2001)
- Early Universe (Ichiki et al. 2006)
- Inflationary processes (see work by, e.g., Kanishvili, Rathra, Jedamzik, Sigl)

Remarks

- Weibel can generate large fields in collapsed structures.
- Biermann mostly generates magnetic seeds in collapsed structures which need large amplification (turbulence + adiabatic compression)
- galactic outflows are characterized by strong B field, but their filling factor is quite difficult to predict
- same is true for outflows from AGN

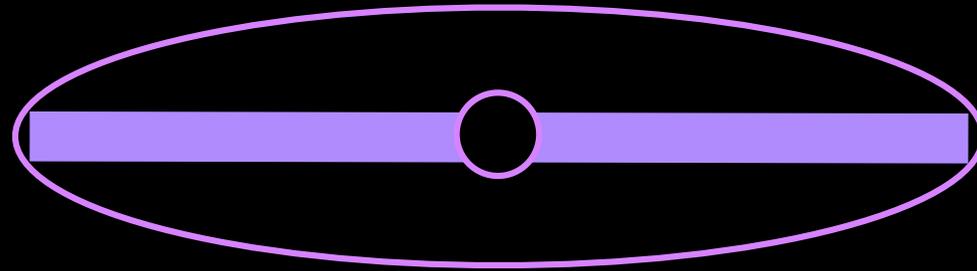
High redshift galaxies reionize the universe



Distant Galaxies in the Hubble Ultra Deep Field
Hubble Space Telescope • Advanced Camera for Surveys

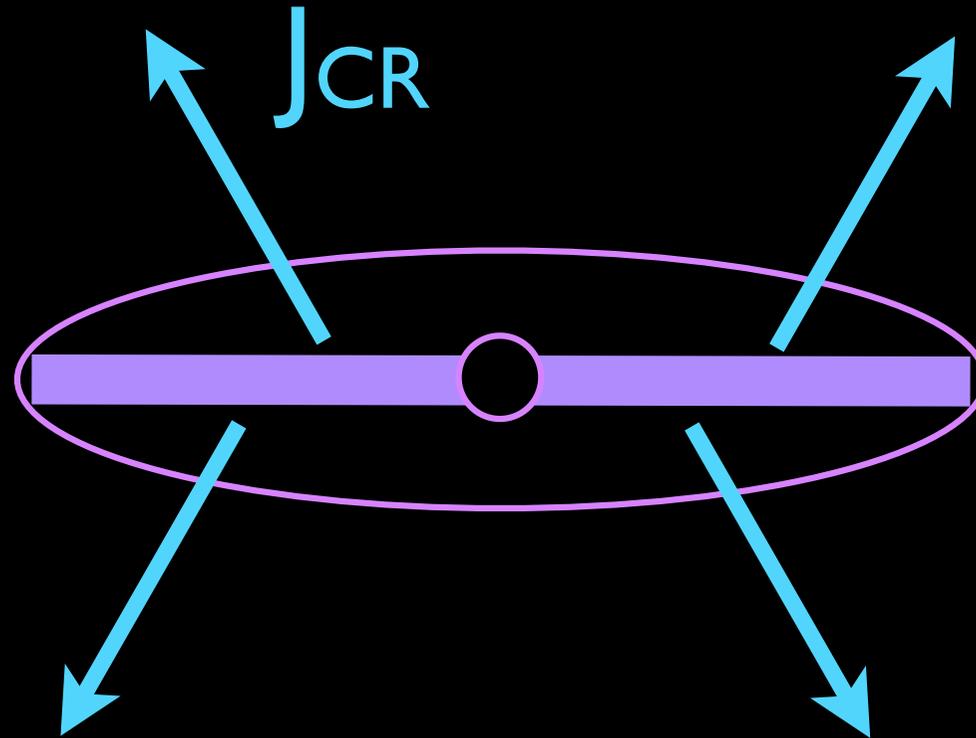
Resistive Mechanism

(FM & Bell, ApJ 2011, 729, 73; arXiv:1001.2011)



Resistive Mechanism

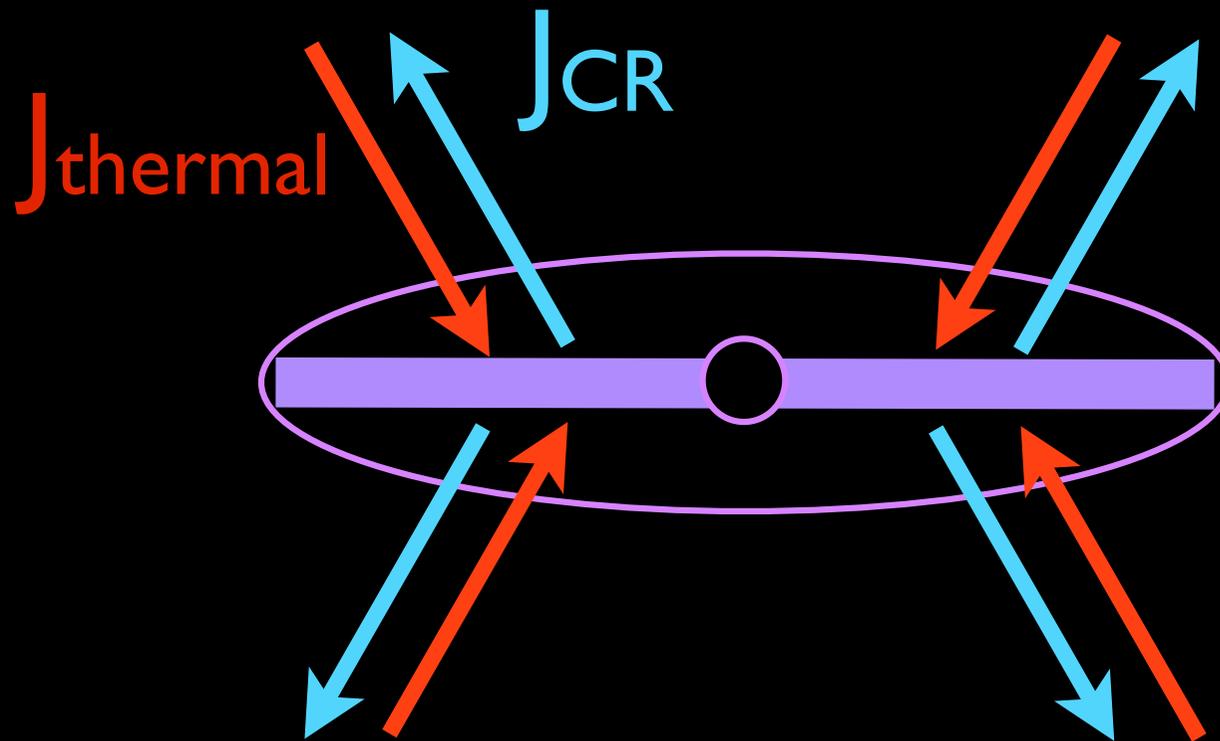
(FM & Bell, ApJ 2011, 729, 73; arXiv:1001.2011)



High-redshift ($z > 6$) star forming galaxies produce copious amount of cosmic-rays which escape into the intergalactic medium.

Resistive Mechanism

(FM & Bell, ApJ 2011, 729, 73; arXiv:1001.2011)



High-redshift ($z > 6$) star forming galaxies produce copious amount of cosmic-rays which escape into the intergalactic medium.

Basics

- the CR current, j_{cr} , drives a return current in the plasma, j_{th} , that tends to cancel j_{cr} itself, i.e.

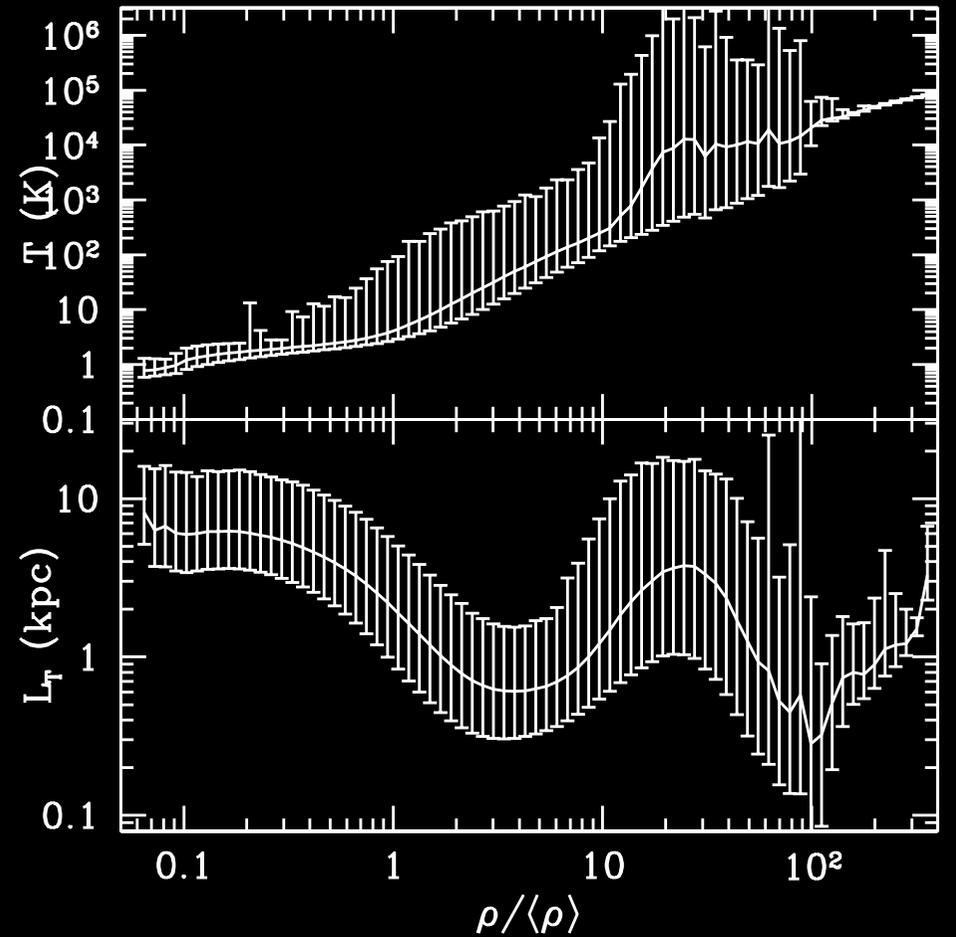
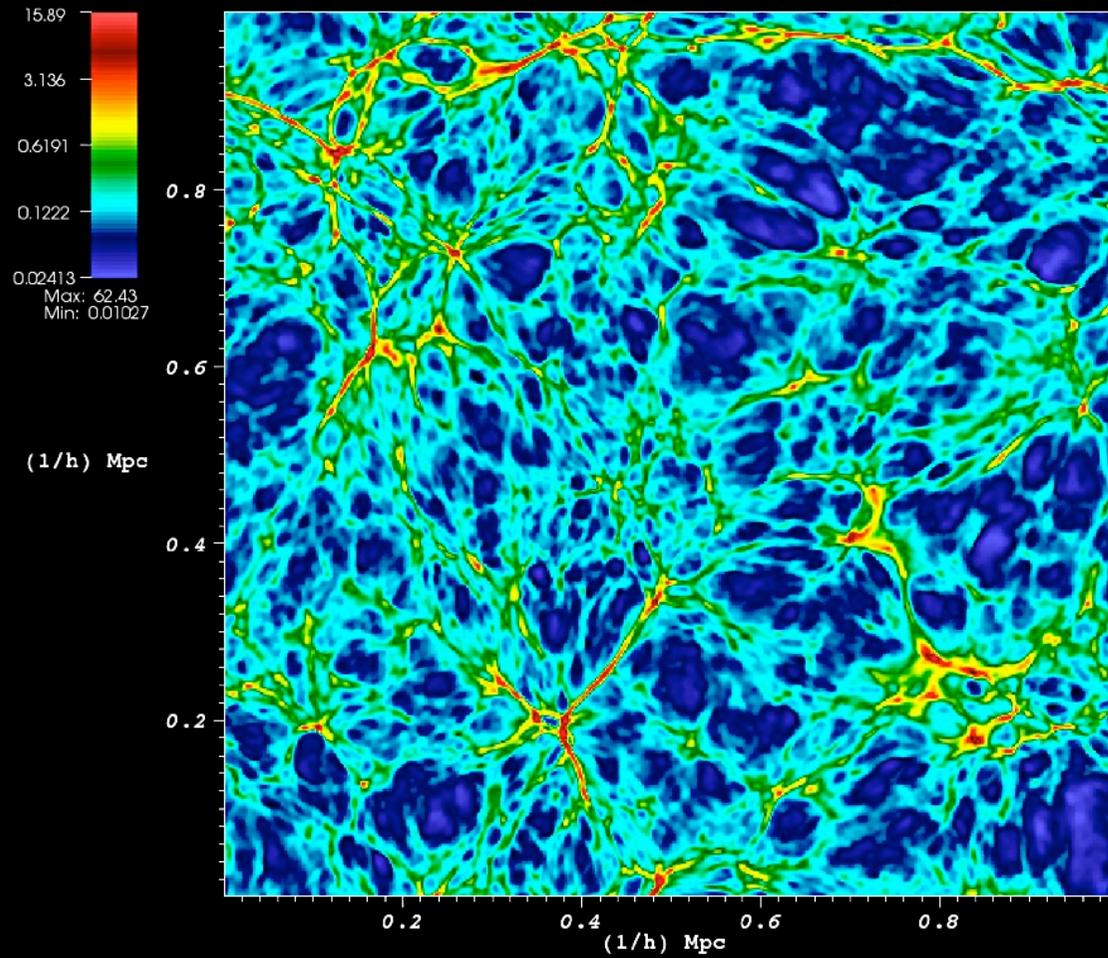
$$j_{th} \approx -j_{cr}$$

- the return current is associated with an electric field:

$$\vec{E}' = \frac{\vec{j}_{th}}{\sigma}, \quad \text{where (Spitzer)} \quad \sigma \simeq 10^7 \left(\frac{T}{K} \right)^{3/2} s^{-1}$$

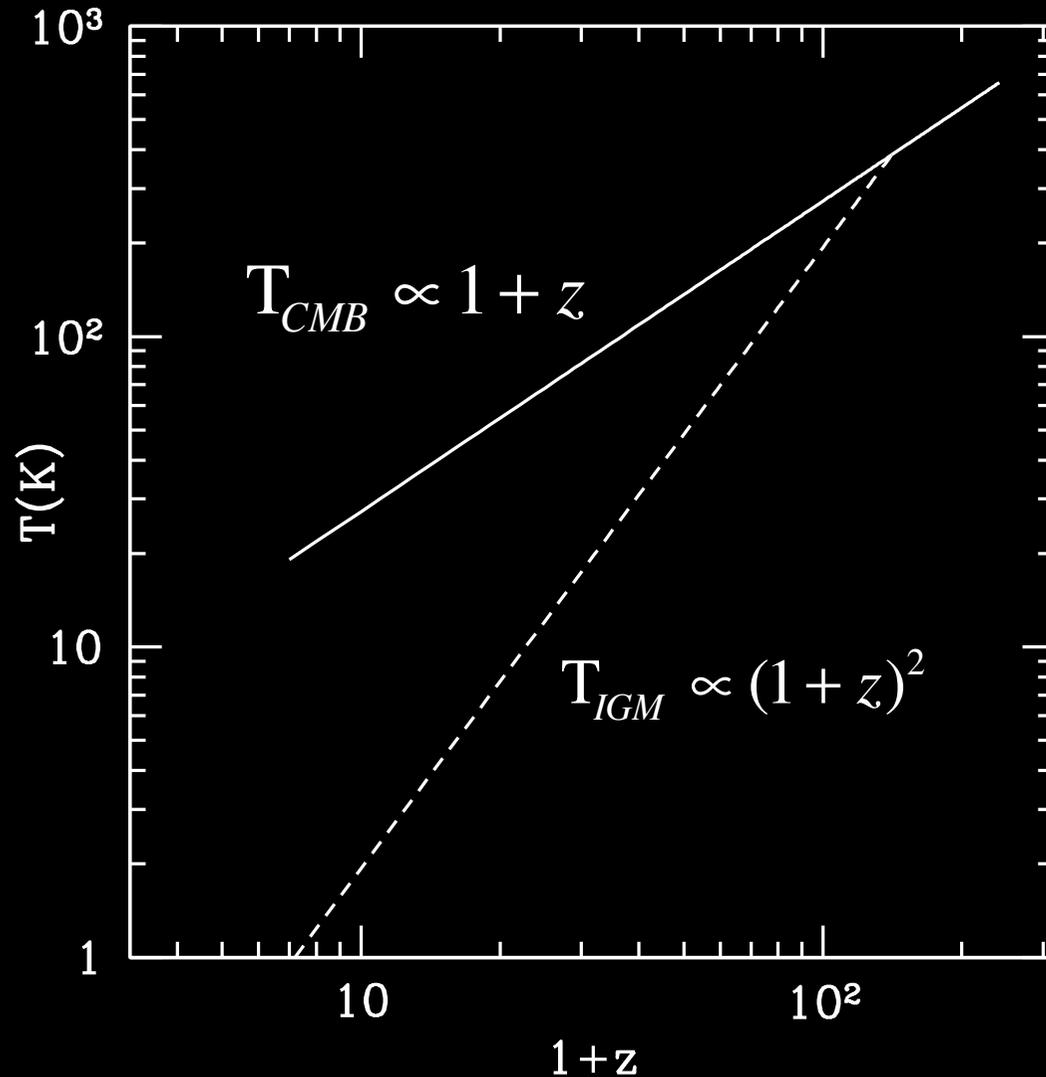
$$\frac{\partial \vec{B}}{\partial t} = -c \vec{\nabla} \times \vec{E} \simeq c \frac{\vec{j}_{cr}}{\sigma} \times \frac{\vec{\nabla} T}{T}$$

IGM inhomogeneities at $z \approx 6$



Miniati & Bell (2011)

IGM Temperature



Compton scattering
efficiently couples
T_{IGM} and T_{CMB}
only for z > 140.

Just prior to cosmic reionization the temperature of the IGM
was at its lowest point (~1K).

Growth rate around L_* galaxies

Temperature
scale-length

$$\ell_T \equiv \frac{T}{|\nabla T|} \sim 1 \text{ kpc}$$

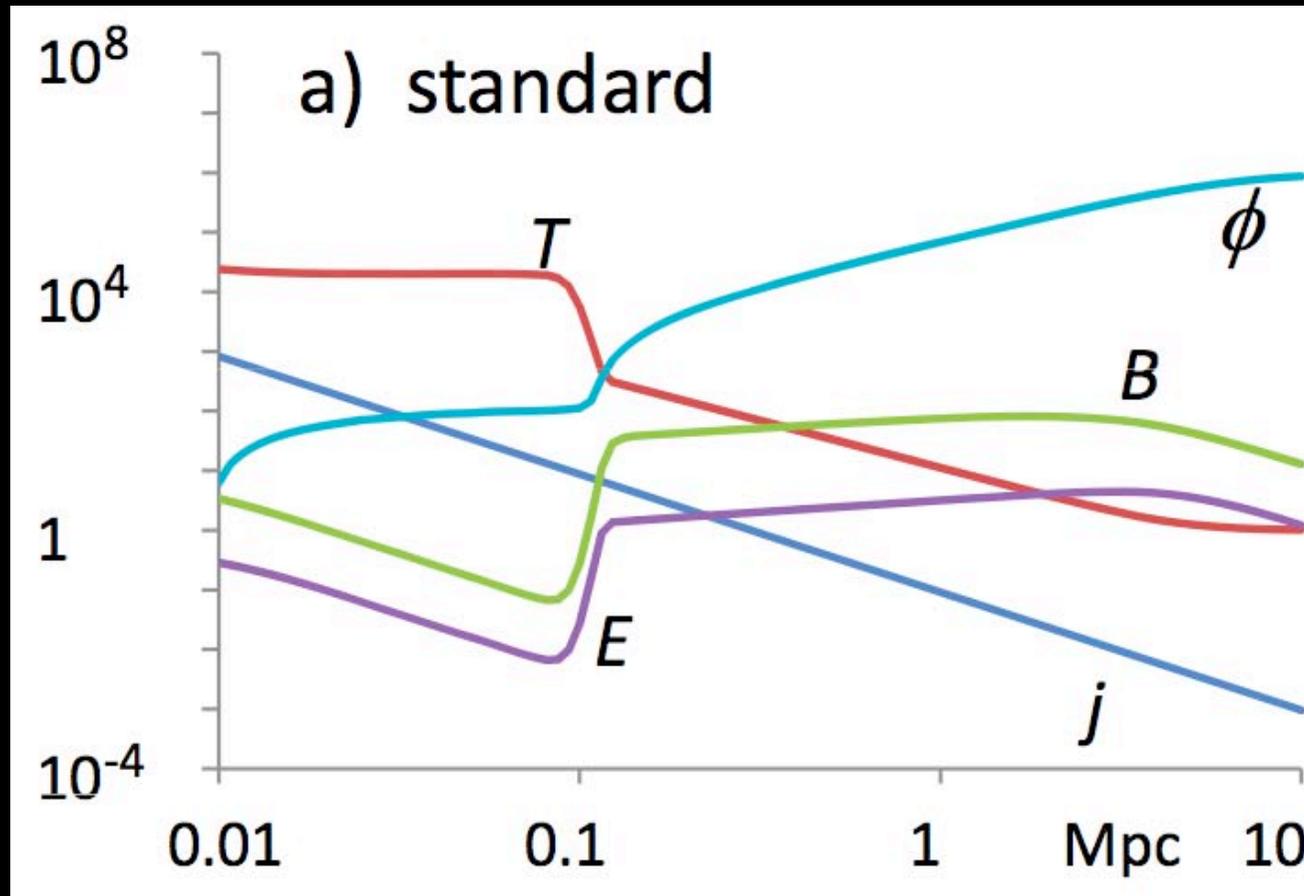
CR-current
density

$$j_{CR} \approx \frac{e\varepsilon_{CR}L}{2\pi\theta p_{\min}c\Lambda_{CR}} d^{-2}$$

$$\frac{\partial B}{\partial t} \approx \frac{cj_{CR}}{\sigma\ell_T} \approx 10^{-15} \left(\frac{\ell_T}{\text{kpc}}\right)^{-1} \left(\frac{T}{\text{K}}\right)^{-3/2} \left(\frac{L}{L_*}\right) \left(\frac{d}{\text{Mpc}}\right)^{-2} \frac{\text{Gauss}}{\text{Gyr}}$$

+ Ohmic heating: $\frac{3}{2}nk_B \frac{dT}{dt} = j_{cr}^2 / \sigma$

Solution's radial profile



$$[j] = 10^{-18} \text{ Am}^{-2}$$

$$[T] = \text{K}$$

$$[B] = 10^{-18} \text{ G}$$

$$[E] = 10^{-18} \text{ Vm}^{-1}$$

$$[\phi] = \text{Volt}$$

Growth rate in the IGM

Luminosity
Function

$$\Phi(L) = \Phi_* \left(\frac{L}{L_*} \right)^{-\alpha} e^{-L/L_*} : n(L) \approx \Phi(L) \frac{L}{L_*}$$

Mean distance
between L-galaxies

$$\langle d_L \rangle = \left[\frac{L}{L_*} \Phi(L) \right]^{-1/3} \propto L^{(\alpha-1)/3} \approx L^{1/4}$$

Bowens et al. 2009

Magnetization
around L-galaxies

$$\dot{B} \propto L \langle d_L \rangle^{-2} \propto L^{1/2}$$

Magnetization of the IGM is dominated by
the most luminous galaxies

MHD + Resistive Source

$$\frac{\partial U}{\partial t} + \vec{\nabla} \cdot \vec{F}_{MHD} = \dot{U}_{Cosmol} + \begin{pmatrix} 0 \\ 0 \\ j_{CR}^2 / \sigma \end{pmatrix} \begin{matrix} \text{mass} \\ \text{momentum} \\ \text{energy} \end{matrix}$$

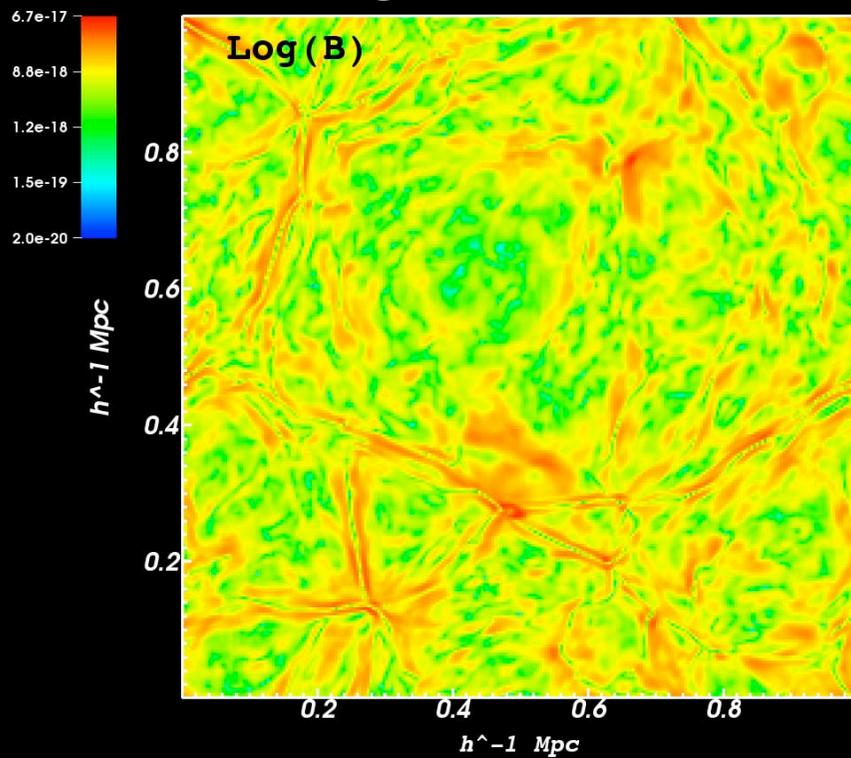
$$\frac{\partial \vec{B}}{\partial t} + \vec{\nabla} \cdot \vec{F}_B = \dot{\vec{B}}_{Cosmol} + c \vec{\nabla} \times (\vec{j}_{CR} / \sigma)$$

Cosmological MHD code (Miniati & Colella 2007; Miniati & Martin 2011). ‘ $\nabla \cdot B = 0$ ’ constraint enforced through Constrained-Transport, hence ‘B’ is not artificially generated (Brackbill & Barnes, 1980).

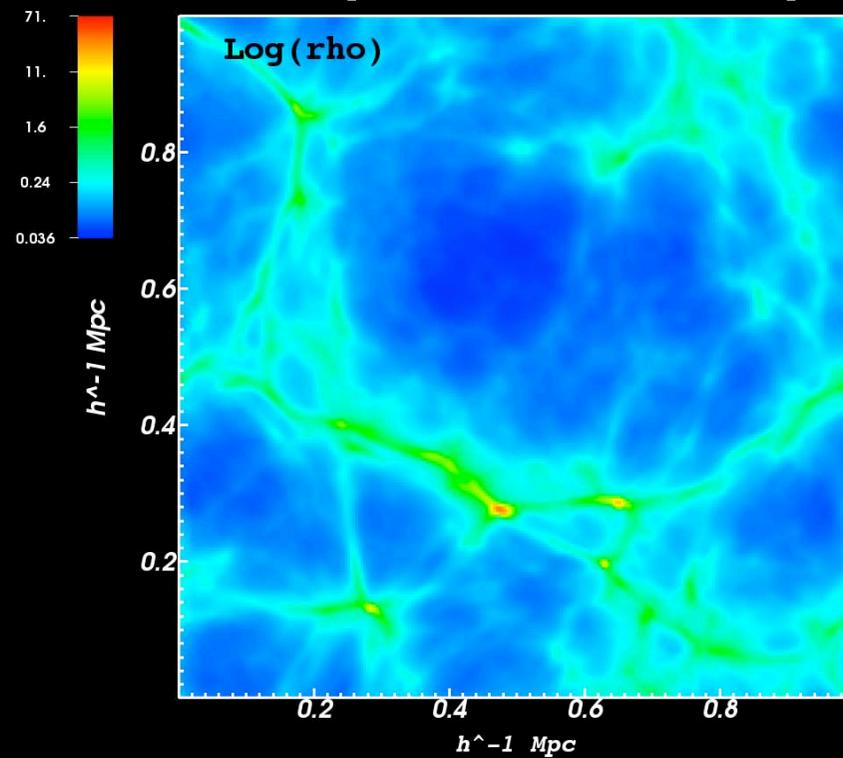
$Z \sim 10$

Box = $100 h^{-1} \text{ kpc}$
 ΛCDM (WMAP7)

Magnetic Field

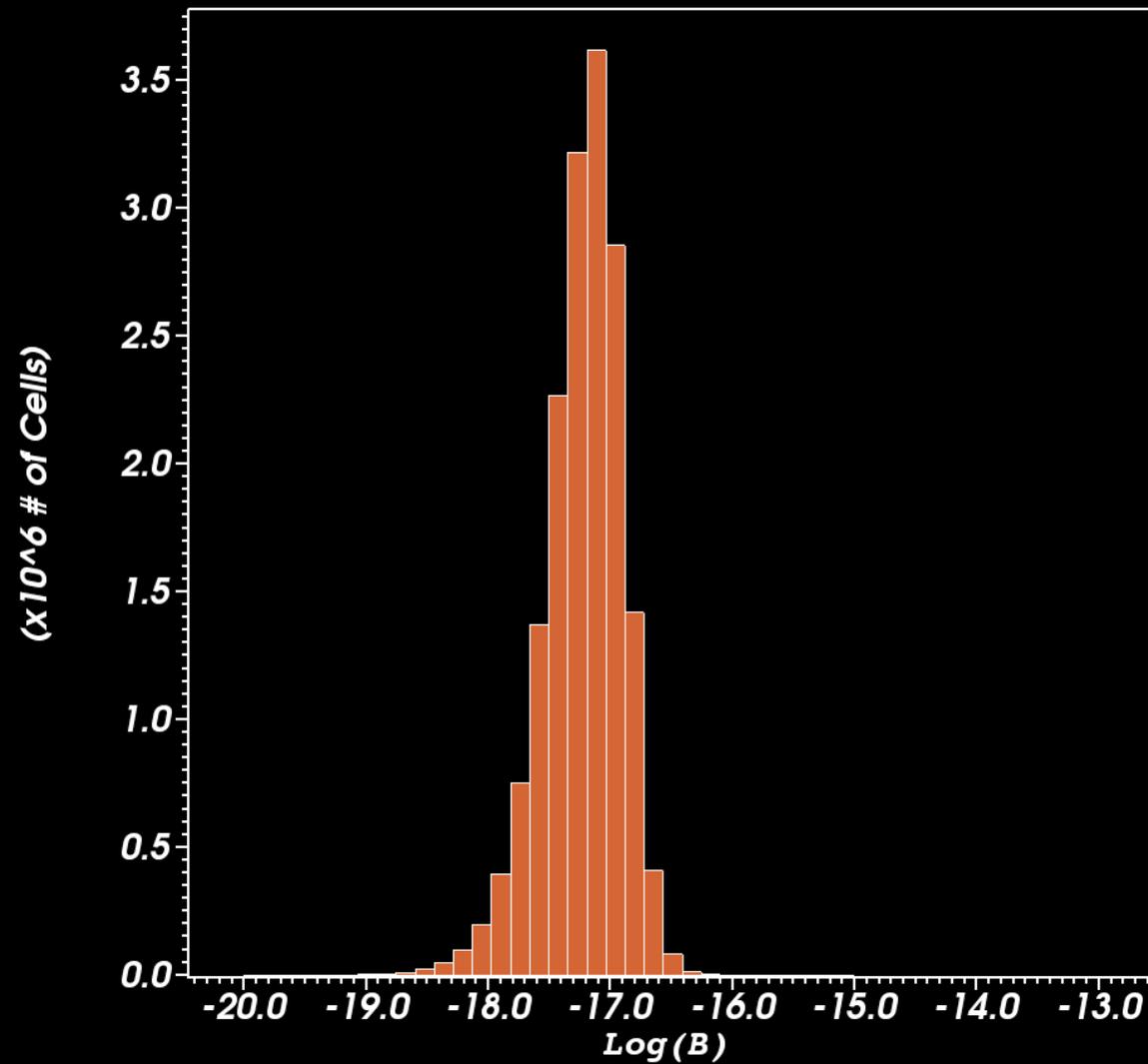


Baryonic Density



$Z \sim 10$

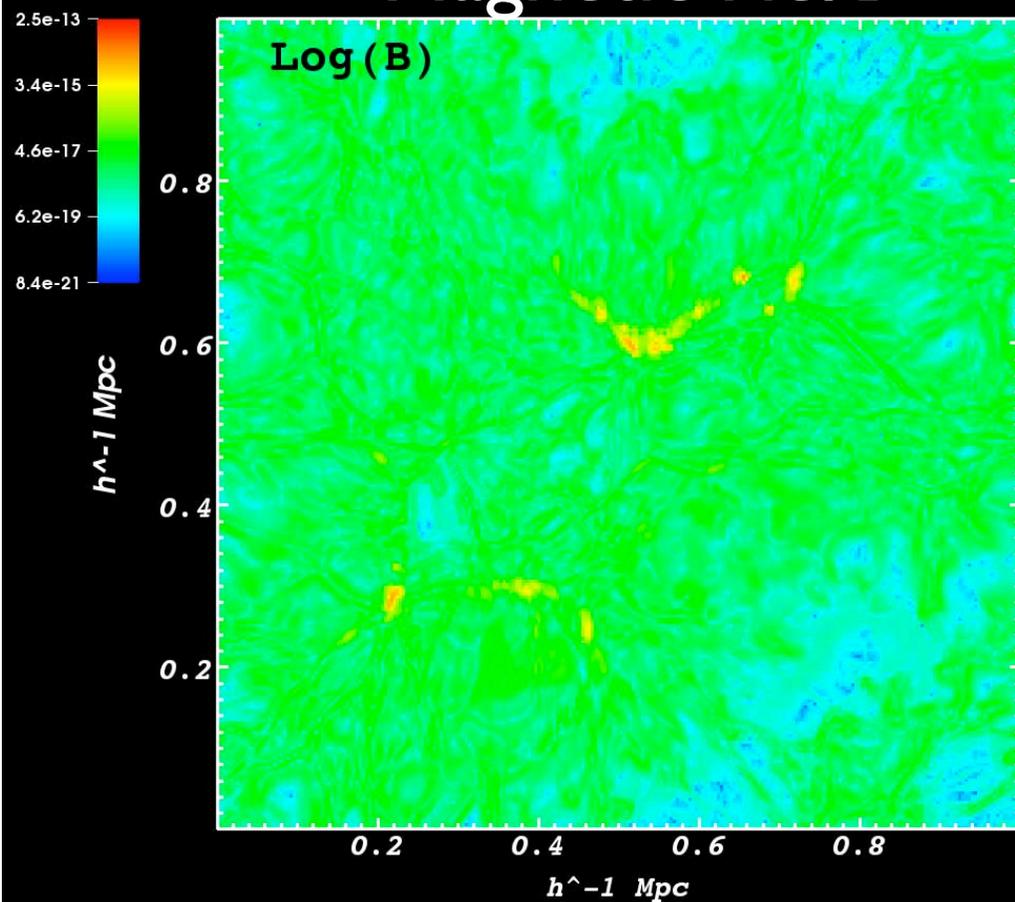
Box = $100 h^{-1} \text{ kpc}$
 ΛCDM (WMAP7)



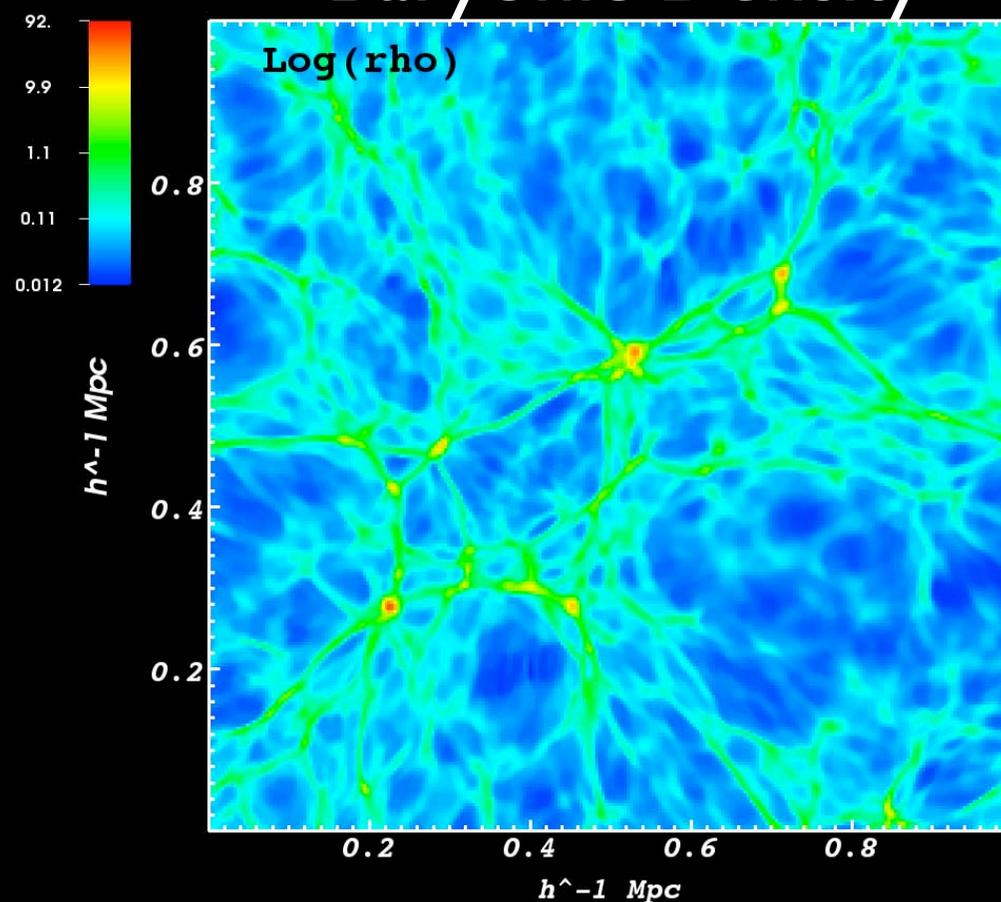
$Z \sim 6$

Box = 1 h^{-1} Mpc
 Λ CDM (WMAP7)

Magnetic Field



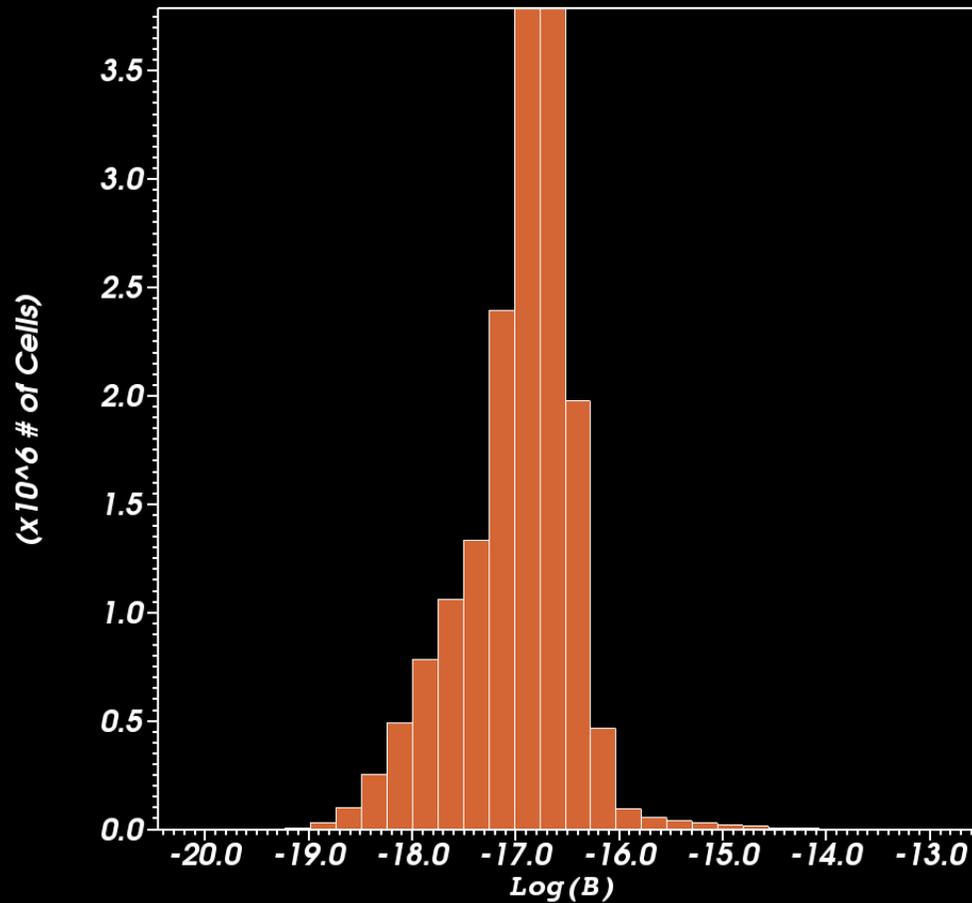
Baryonic Density



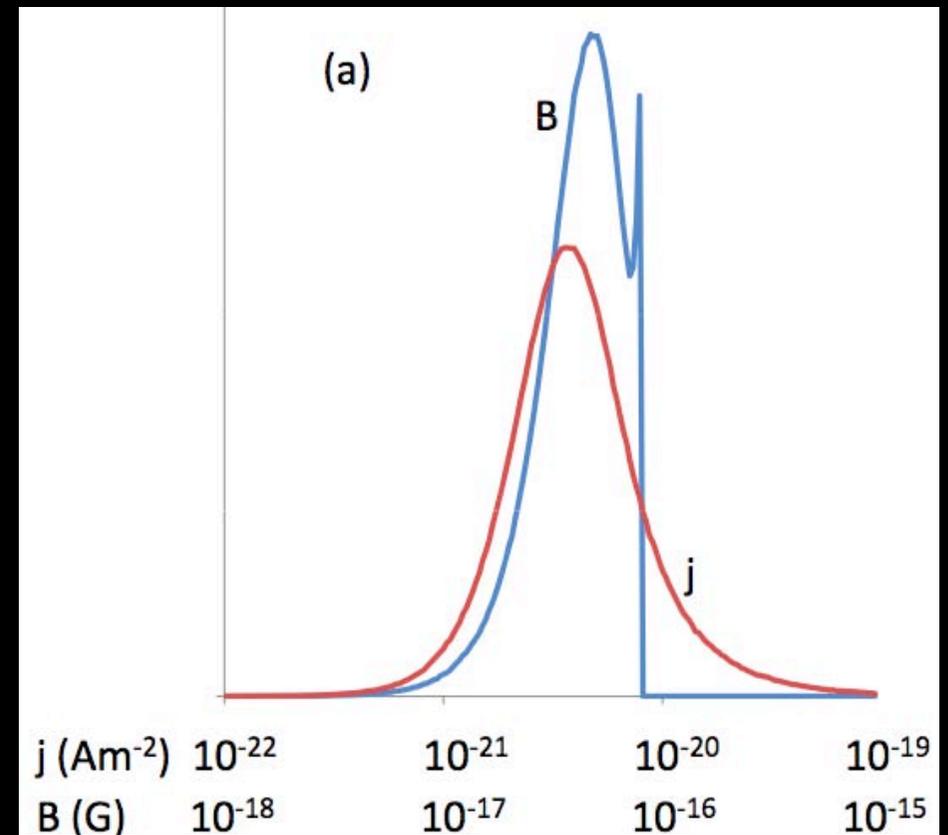
$Z \sim 6$

Box = $1 \text{ h}^{-1} \text{ Mpc}$
 ΛCDM (WMAP7)

Numerical Simulation



Monte Carlo solution



Non-resonant term

$$\frac{\partial U}{\partial t} + \vec{\nabla} \cdot \vec{F}_{MHD} = S_{Cosmol+Resistive} - \begin{pmatrix} 0 \\ \vec{j}_{CR} \times \vec{B} \\ \vec{u} \cdot (\vec{j}_{CR} \times \vec{B}) \end{pmatrix} \begin{matrix} \text{mass} \\ \text{momentum} \\ \text{energy} \end{matrix}$$

whence an instability with growth rate (Bell 2005)

$$\gamma = \left(\frac{kBj_{CR}}{\rho} \right)^{\frac{1}{2}} \approx 0.5 \text{ Gyr}^{-1} \left(\frac{\lambda}{kpc} \right)^{-\frac{1}{2}} \left(\frac{B}{5 \times 10^{-17} G} \right)^{\frac{1}{2}} \left(\frac{d}{Mpc} \right)^{-1}$$

important, may be, on 10 Gyr time scales ($z \sim 1$)

Summary

- Magnetic fields are ubiquitous and play several different roles in astrophysical plasma
- Recently discovered in cosmic voids!
- Resistive mechanism provides suitable seed fields $B \sim 10^{-16} - 10^{-17}$ G at $z \sim 10^{-6}$
- Degree of amplification by the non-resonant instability is not clear yet