

Neutrinos: Kage-Musha in nature which
however have a key to understand her
fundamental structure

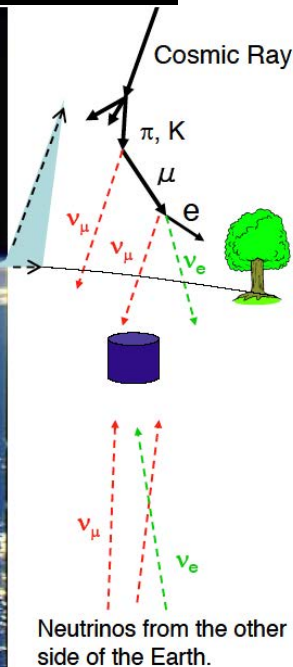
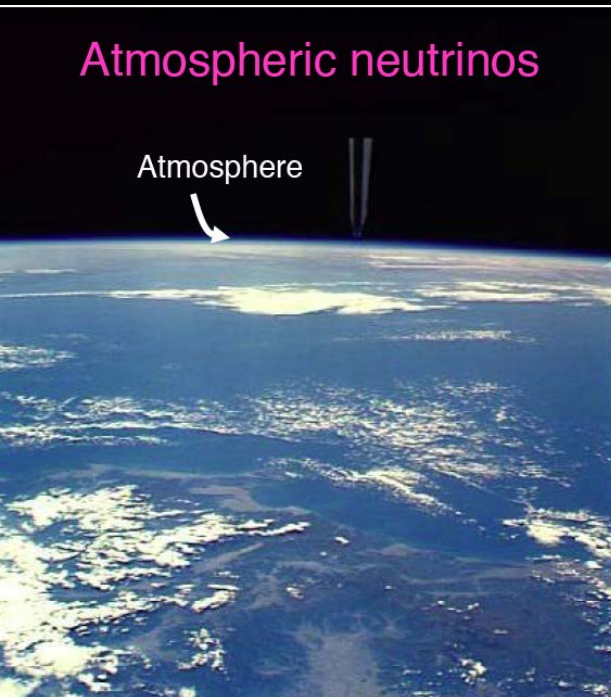


Hisakazu Minakata



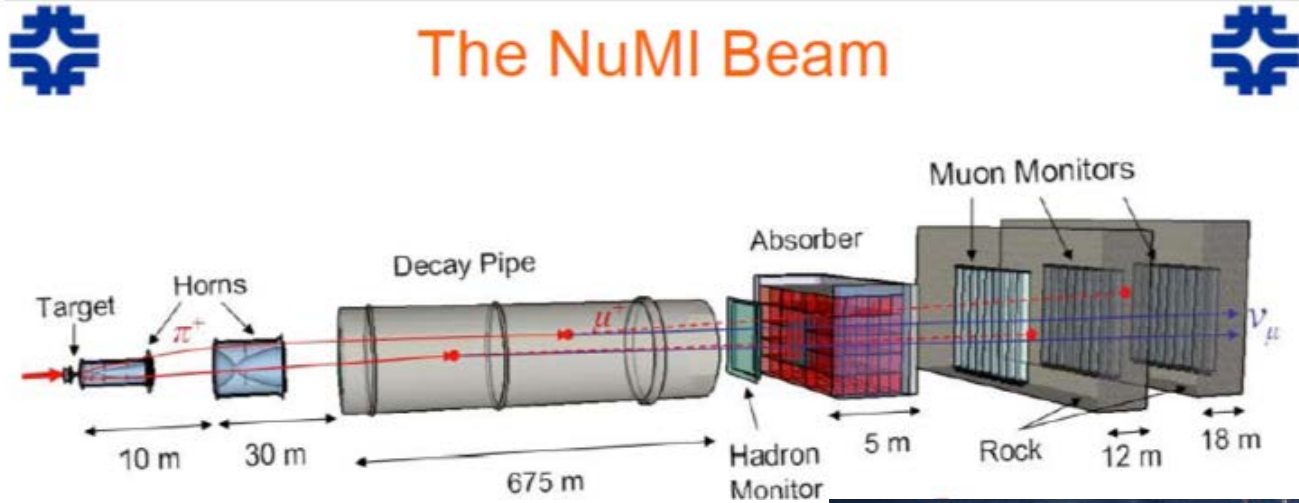
Neutrinos: interdiscipli nary science

Neutrinos come from natural sources

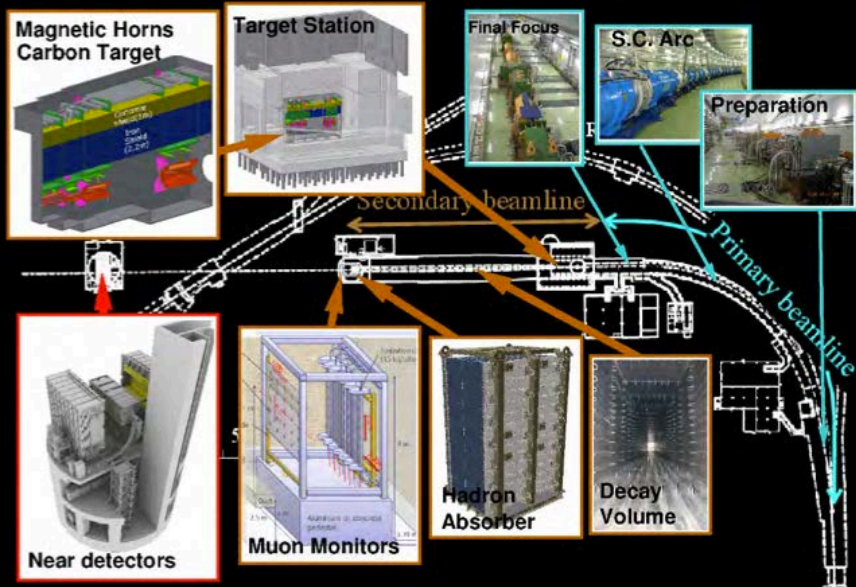


Nature **436**, 499-503 (28 July 2005)

Neutrinos come from artificial sources



Tokai Site



Two aspects of neutrino physics

Understanding phenomena
in nature through neutrinos

- core collapse SN explosion: cannot occur without neutrinos
- Sun does not shine without neutrinos
- BBN not OK without ν
- Geo-neutrinos
- Ex-high energy ν ?

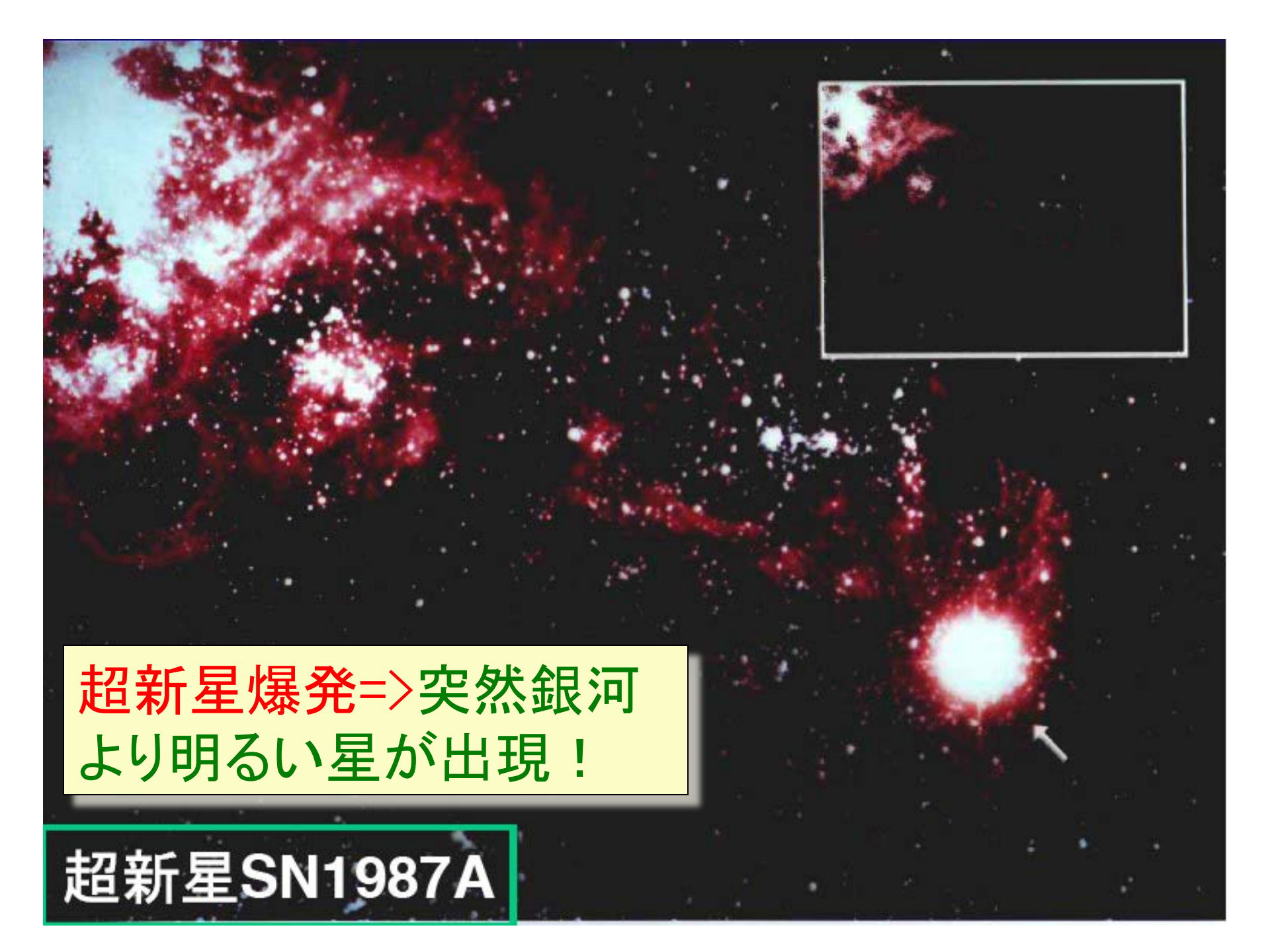
Understanding nature
using neutrinos as probe

- Pure chirality of ν leads to chiral picture of fundamental particles
- Neutrino mass probes high-energy worlds
- Through- ν discovered lepton flavor mixing shed light on quark-lepton relation



Neutrinos
arrived
from
outside our
galaxy !



This is a deep-field astronomical photograph showing a large, irregularly shaped galaxy with a complex, filamentary structure. The galaxy is primarily red, with numerous bright white and yellow stars scattered throughout. In the lower right quadrant, a single, exceptionally bright white star stands out from the rest. A small white arrow points directly at this star. In the upper right corner, there is a rectangular inset box containing a zoomed-in view of a portion of the galaxy's structure. The background is a deep black, filled with distant, faint stars.

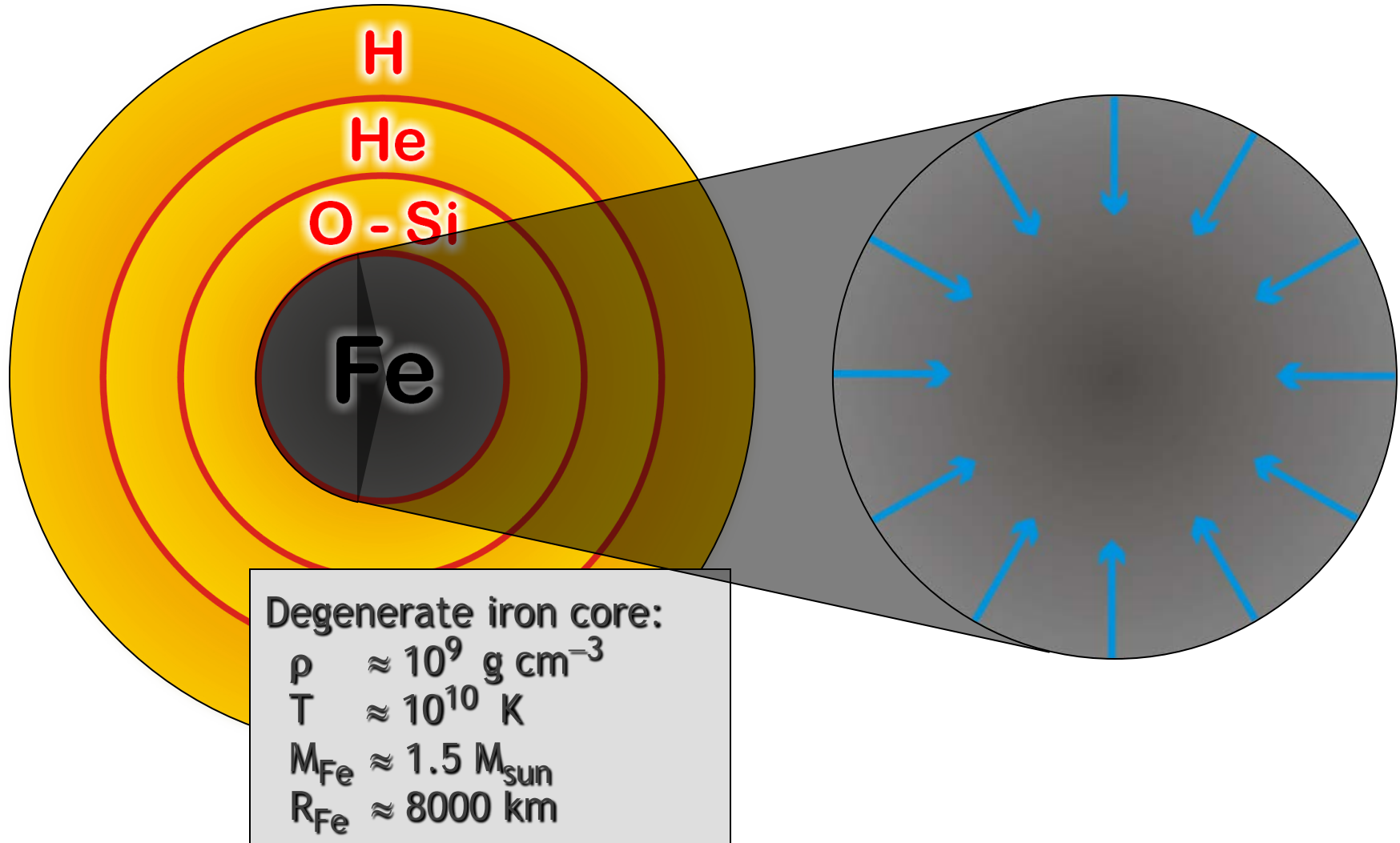
超新星爆発=>突然銀河
より明るい星が出現！

超新星SN1987A

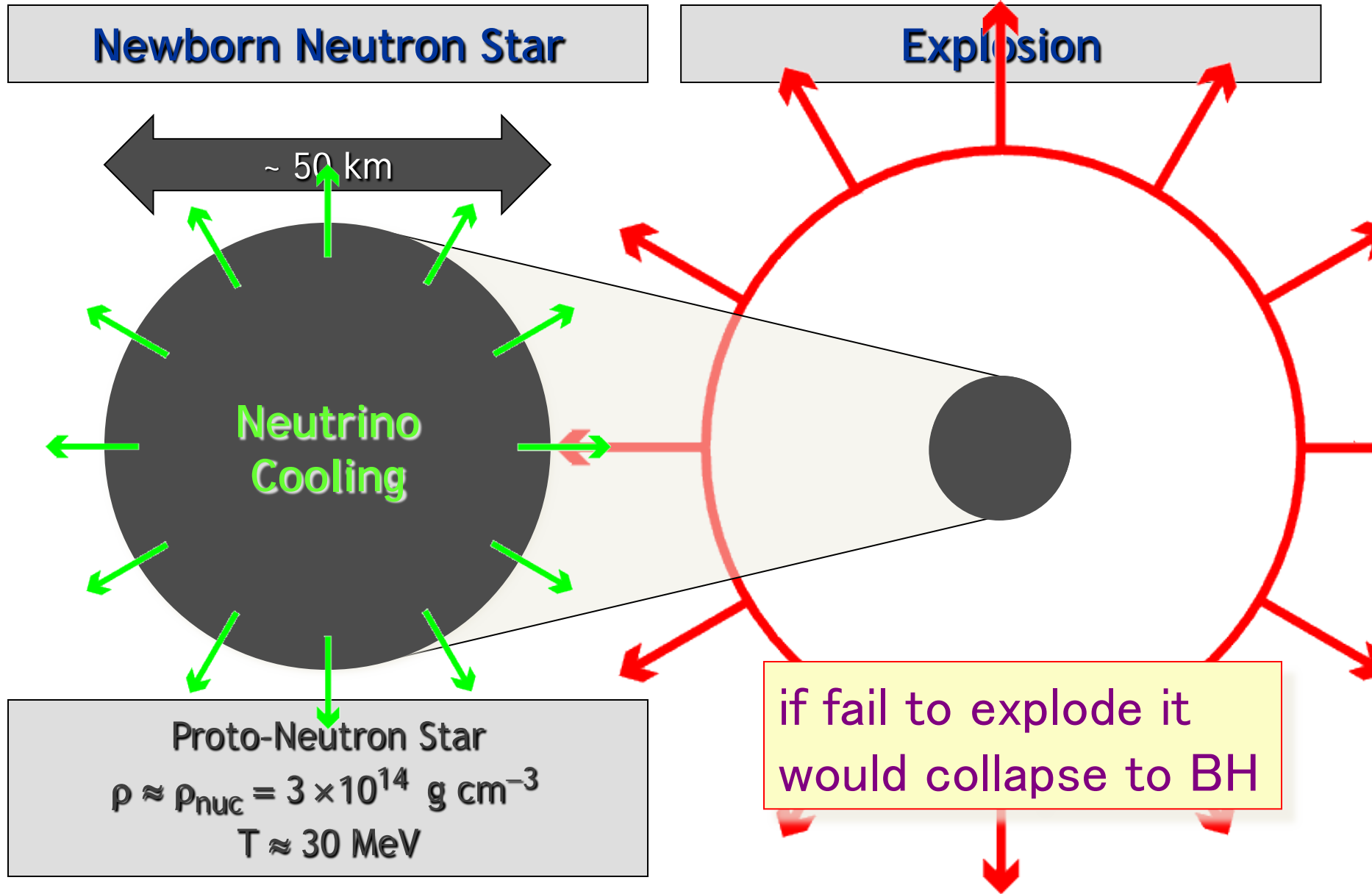
Stellar Collapse and Supernova Explosion

Onion structure

Collapse (implosion)




Stellar Collapse and Supernova Explosion



Supernova energy budget

- How big amount of energy is released ?

 simple to compute!

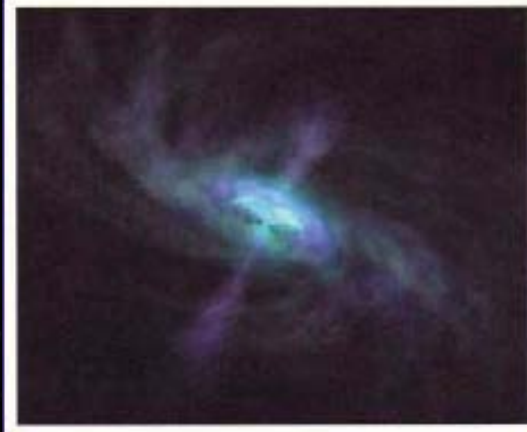
- To form neutron star how much energy should be extracted? $A = \text{grav. binding energy} = 10^{46} \text{ J}$
- How can it be transported to outside star in 10 s?  Neutrinos!
- 99% of 10^{46} J are carried away by ν
- Only remaining 1% explodes outer progenitor

Neutrino IS the major player in core collapse supernova !

SN1054 now: crab



Ejected material
is expanding
with velocity \sim
600km/s

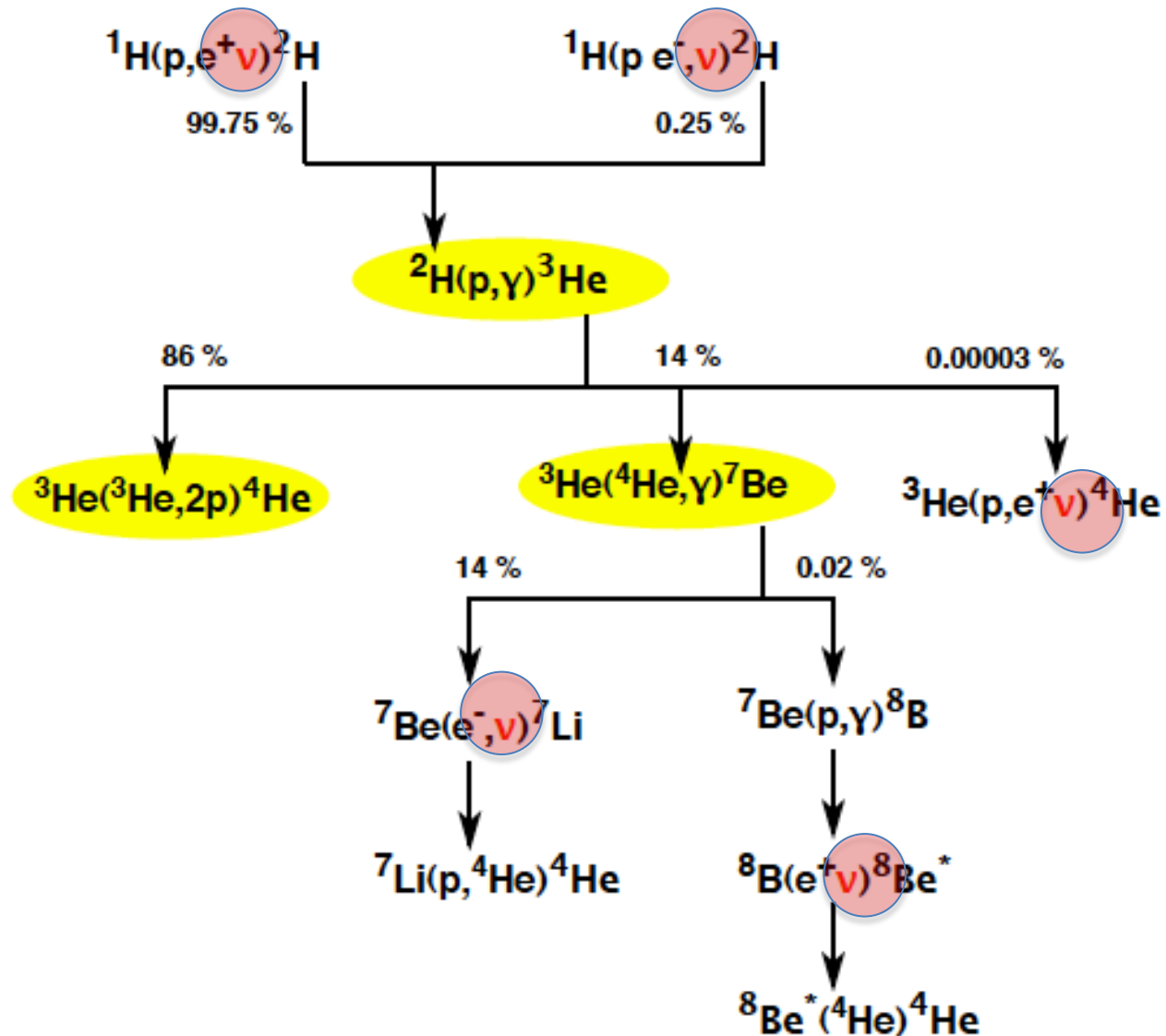


中心部には高速回転するパルサーがある。(毎秒30回、回転している)



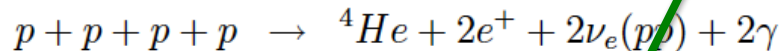
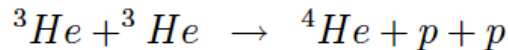
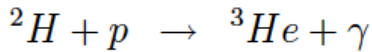
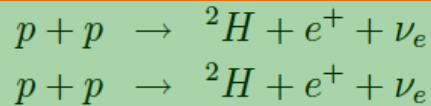
Solar neutrino problem

Sun does not shine without neutrinos: Nuclear chain reaction in the Sun

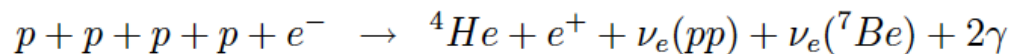
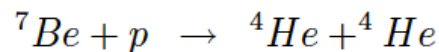
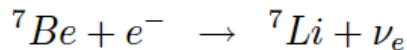
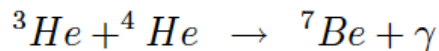
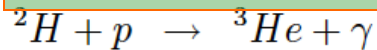
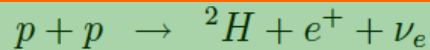


Net reaction: $pppp \rightarrow {}^4\text{He} + 2e^+ + 2\nu_e + \sim 25 \text{ MeV}$

pp I chain (termination 85%)

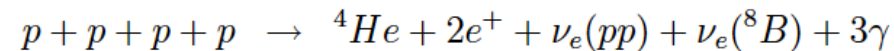
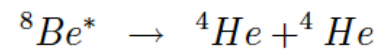
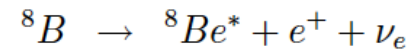
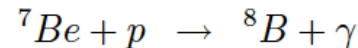
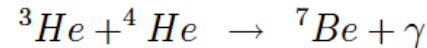
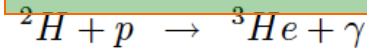
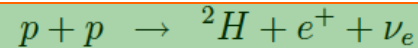


pp II chain (termination 15%)



ppl: main engine of the Sun

pp III chain (termination 0.02%)



Solar neutrino problem since 60' s



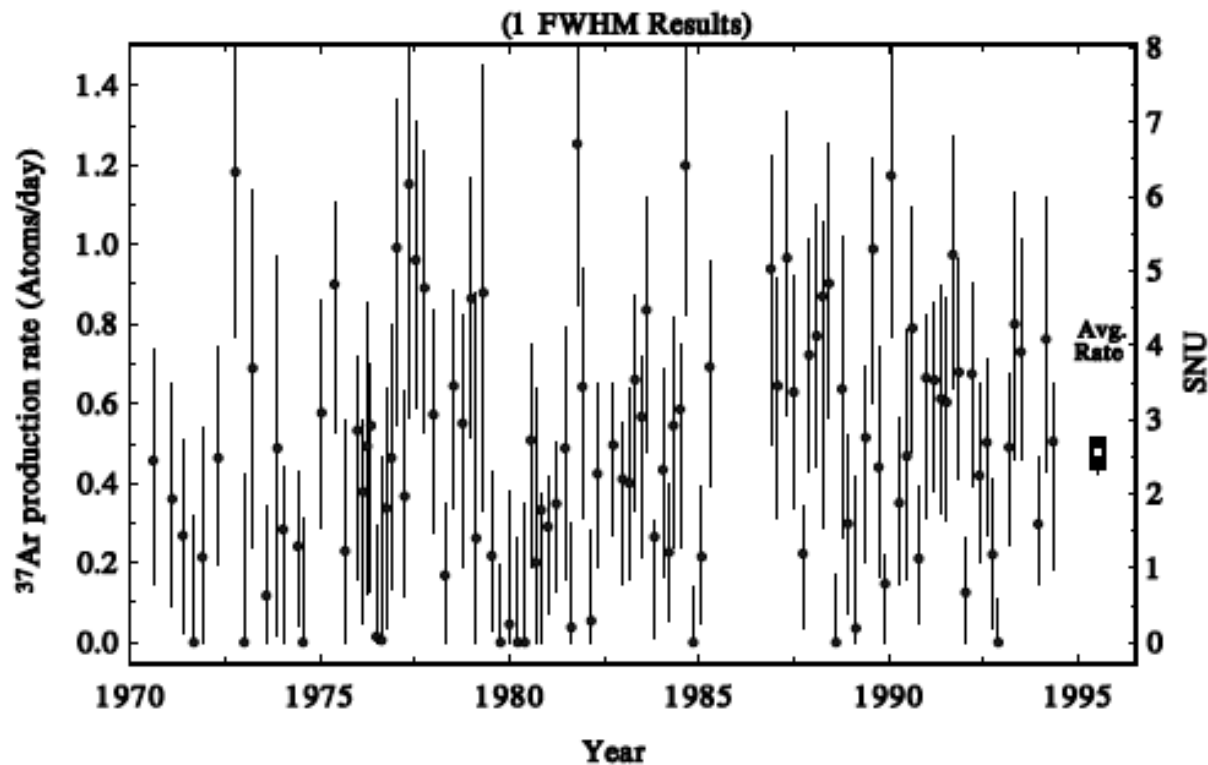
Raymond David Jr. (left) and John Bahcall in miner's clothing and protective hats. The photograph was taken in 1967 about a mile underground in the Homestake Gold Mine in Lead, South Dakota, USA. Davis is pictured showing Bahcall his newly constructed steel tank (6 meters in diameter, 15 meters long), which contained a large amount of cleaning fluid (40,000 liters) and was used to capture neutrinos from the Sun.

October 3, 2012

Colloquium@Kavli-IPMU



Result of ^{37}Cl experiment: ^{37}Ar production rate

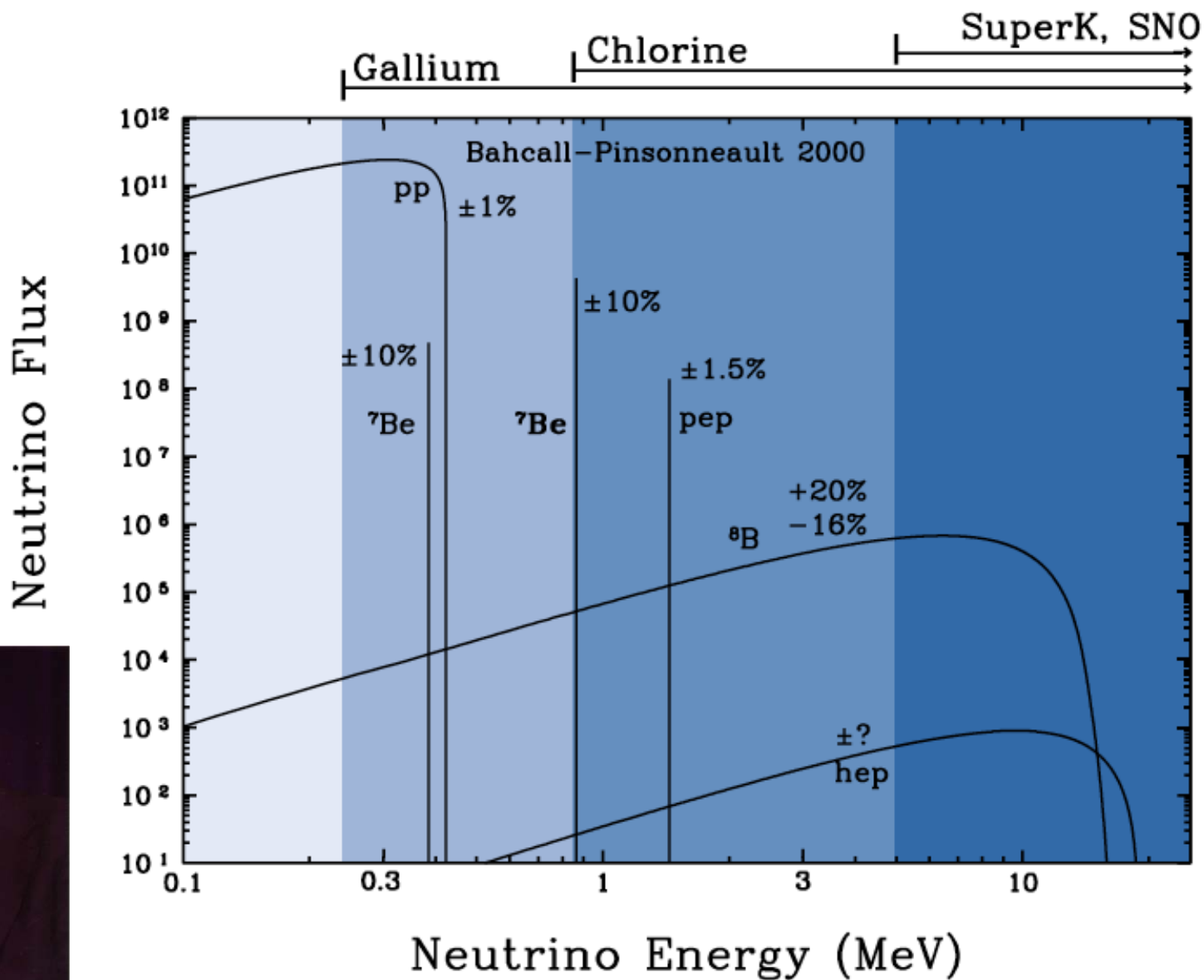


1 SNU = 10^{-36}
events / (No. of
target atom sec)

~1/3 of SSM
prediction

FIG. 13.—Homestake Experiment—one FWHM results. Results for 108 individual solar neutrino observations made with the Homestake chlorine detector. The production rate of ^{37}Ar shown has already had all known sources of nonsolar ^{37}Ar production subtracted from it. The errors shown for individual measurements are statistical errors only and are significantly non-Gaussian for results near zero. The error shown for the cumulative result is the combination of the statistical and systematic errors in quadrature.

計算された太陽ニュートリノ強度 (flux)



2039m underground

SNO

6000 mwe
overburden

1000 tonnes D₂O

12 m Diameter
Acrylic Vessel

1700 tonnes Inner
Shield H₂O

Support Structure
for 9500 PMTs,
60% coverage

5300 tonnes Outer
Shield H₂O

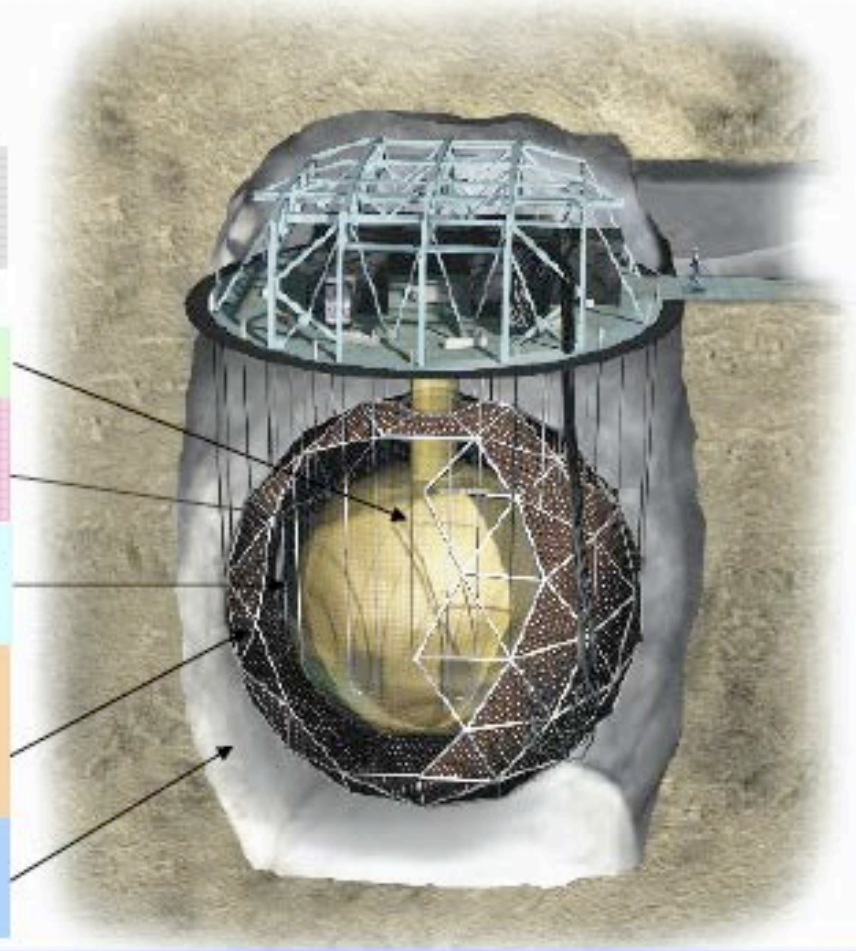
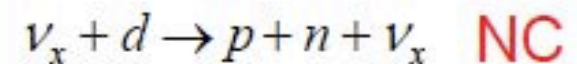
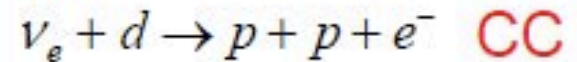
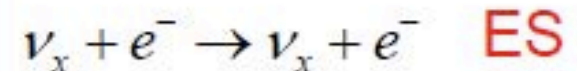
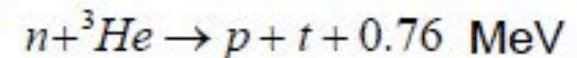
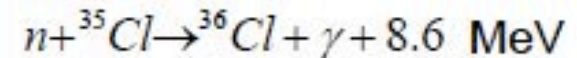
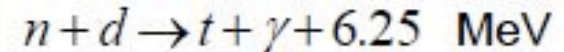


Image courtesy National Geographic

3 Reactions:



3 neutron detection methods:



3 Phases:

- Just D₂O
- D₂O + 2 tonnes NaCl
- D₂O + ³He Proportional Counters (“NCDs”)

391 days salt data - in numbers

$$\phi_{CC} = 1.68^{+0.06}_{-0.06}(\text{stat.})^{+0.08}_{-0.09}(\text{syst.})$$

$$\phi_{NC} = 4.94^{+0.21}_{-0.21}(\text{stat.})^{+0.38}_{-0.34}(\text{syst.})$$

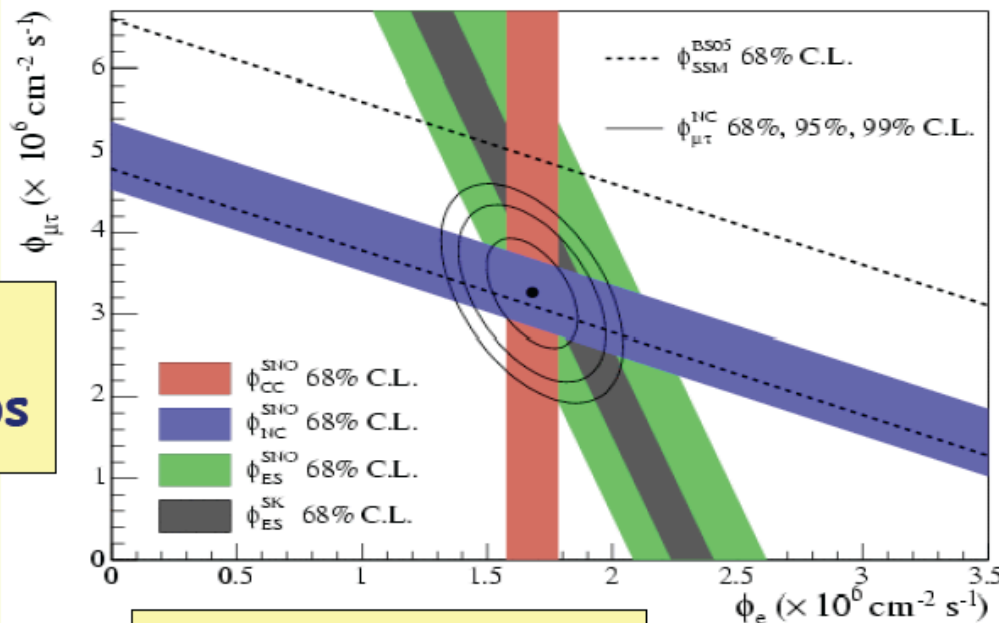
$$\phi_{ES} = 2.35^{+0.22}_{-0.22}(\text{stat.})^{+0.15}_{-0.15}(\text{syst.})$$

$$5.25 \pm 3.7\% \\ [10^6 \text{ cm}^{-2} \text{ s}^{-1}]$$

$$\frac{\phi_{CC}}{\phi_{NC}} = 0.340 \pm 0.023(\text{stat.})^{+0.029}_{-0.031}$$

(In units of
 $10^6 \text{ cm}^{-2} \text{ s}^{-1}$)

μ, τ
neutrinos



electron ν

fluxes
for all
neutrinos

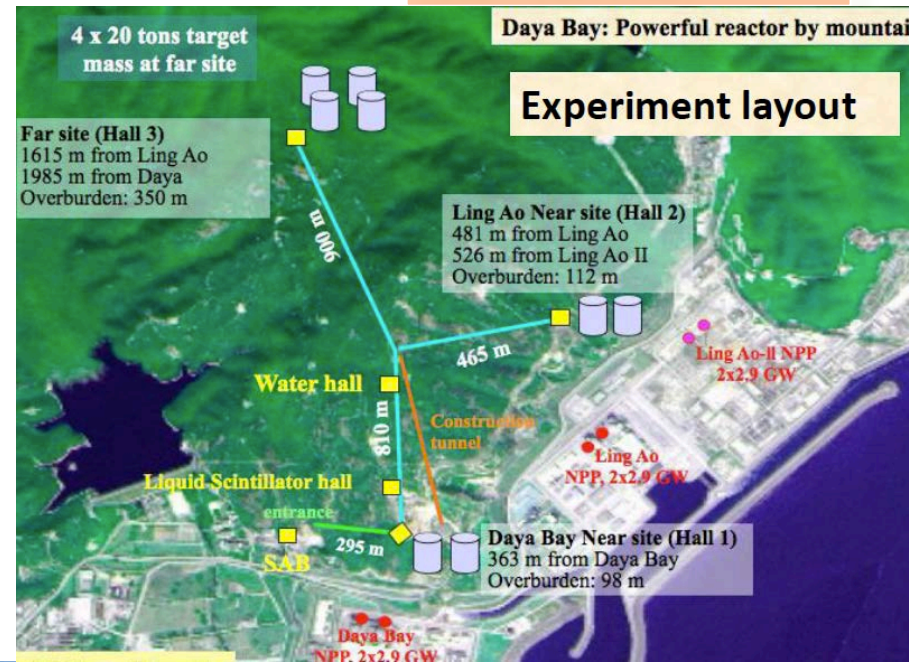
SNO solves the solar neutrino problem



The same
physics was
seen by
reactor v

Reactor is powerful ν source

Daya-Bay, Ling-Ao (China)



Kashiwazaki-Kariwa
(Japan)

YongGwang (Korea)

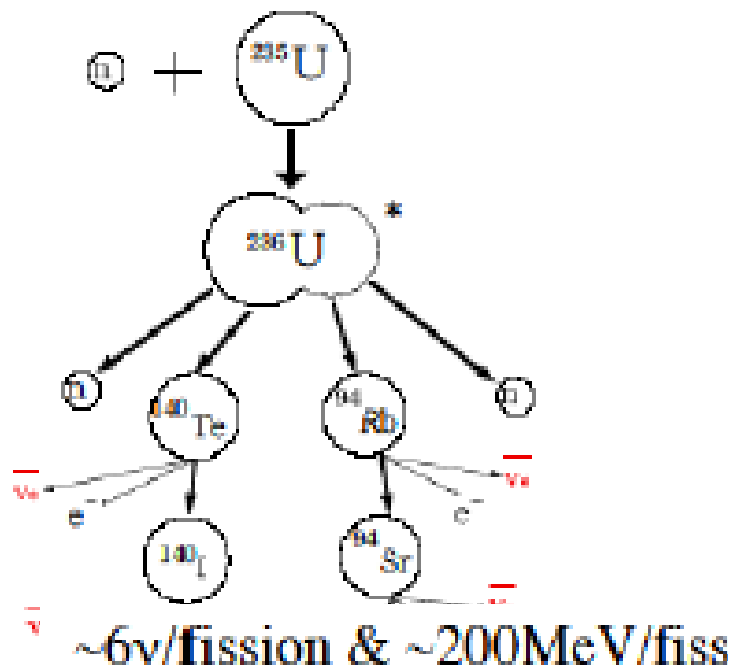


原子炉ニュートリノ実験

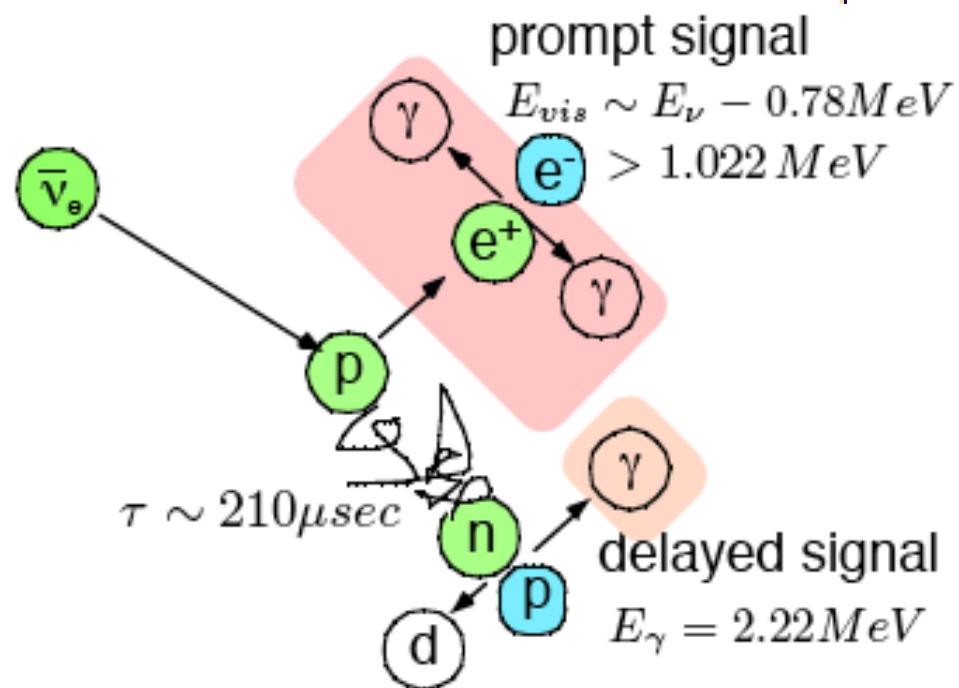
Reactor $\bar{\nu}_e$

$\bar{\nu}_e$ Detection

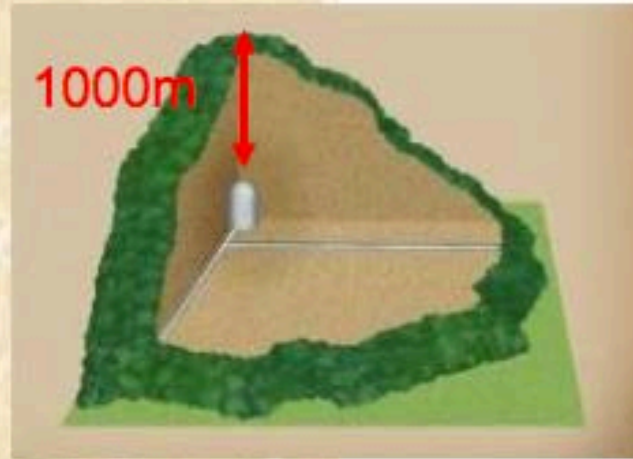
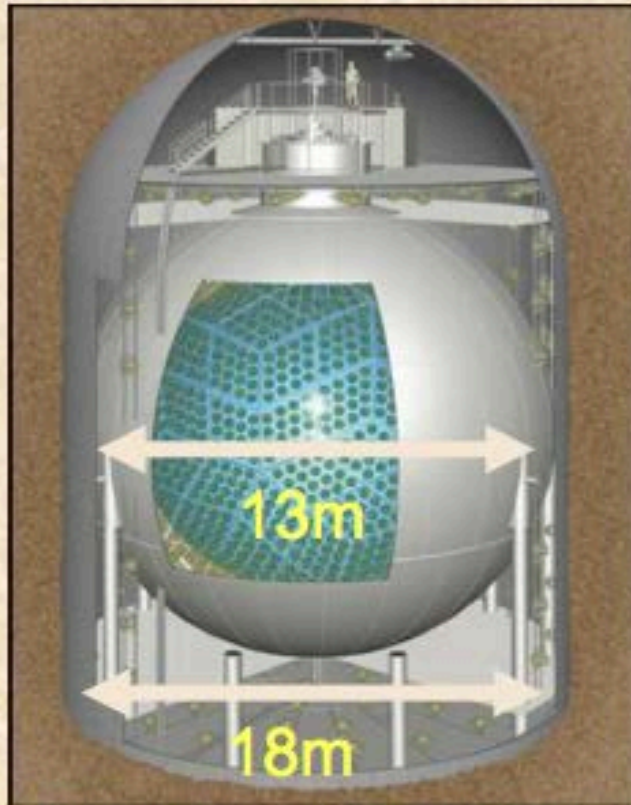
$\bar{\nu}_e + p \rightarrow e^+ + n$



$\sim 6 \times 10^{20} \bar{\nu}_e / s / \text{reactor (1GWe)}$



KamLAND : calorimeter type detector



Shield of cosmic
ray background
 $\sim 10^{-5}$ than ground



1,000 tons pure liquid scintillator (LS)

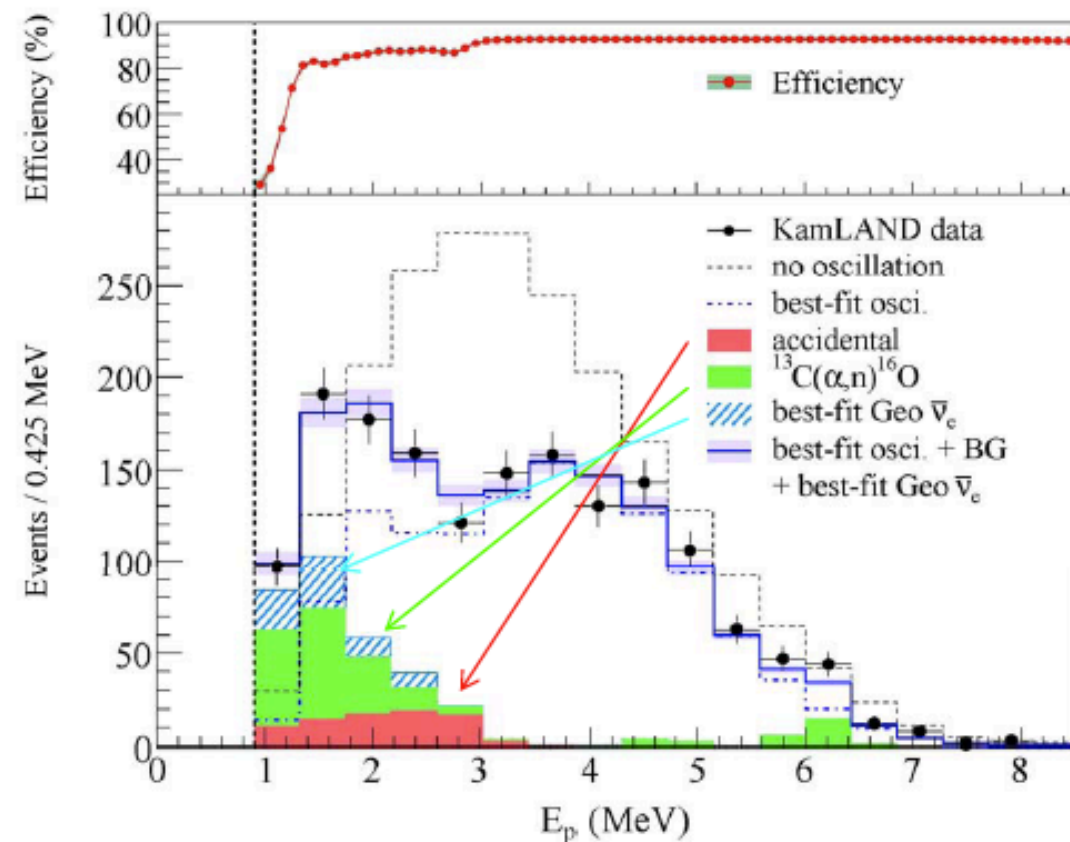
Buffer oil : for environmental radiation

PMT : 17inch : 1325 + 20inch : 554

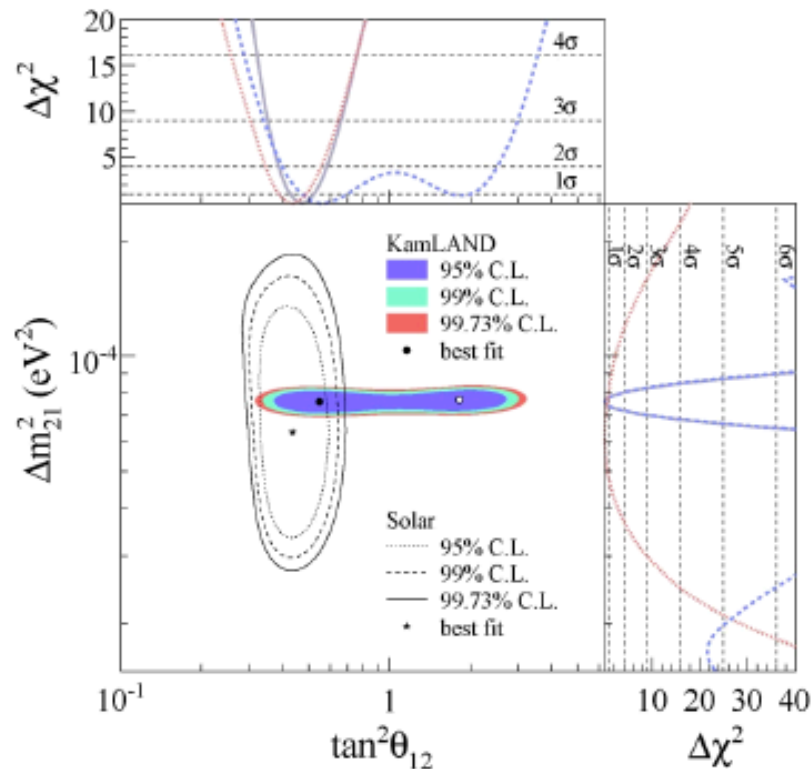
Water cherenkov anti counter
225 20inch PMT with water

Resolution : $\sim 12\text{cm} / \sqrt{E(\text{MeV})}$
 $\sim 6.4\% / \sqrt{E(\text{MeV})}$

Reactor neutrino energy spectrum



Fit to scaled no-oscillation spectrum
: exclude at 5.1σ



$$\Delta m^2 = 7.58^{+0.21}_{-0.20} \times 10^{-5} \text{eV}^2$$

$$\tan^2 \theta = 0.56^{+0.14}_{-0.09}$$

Cleanest signature of neutrino oscillation by KamLAND

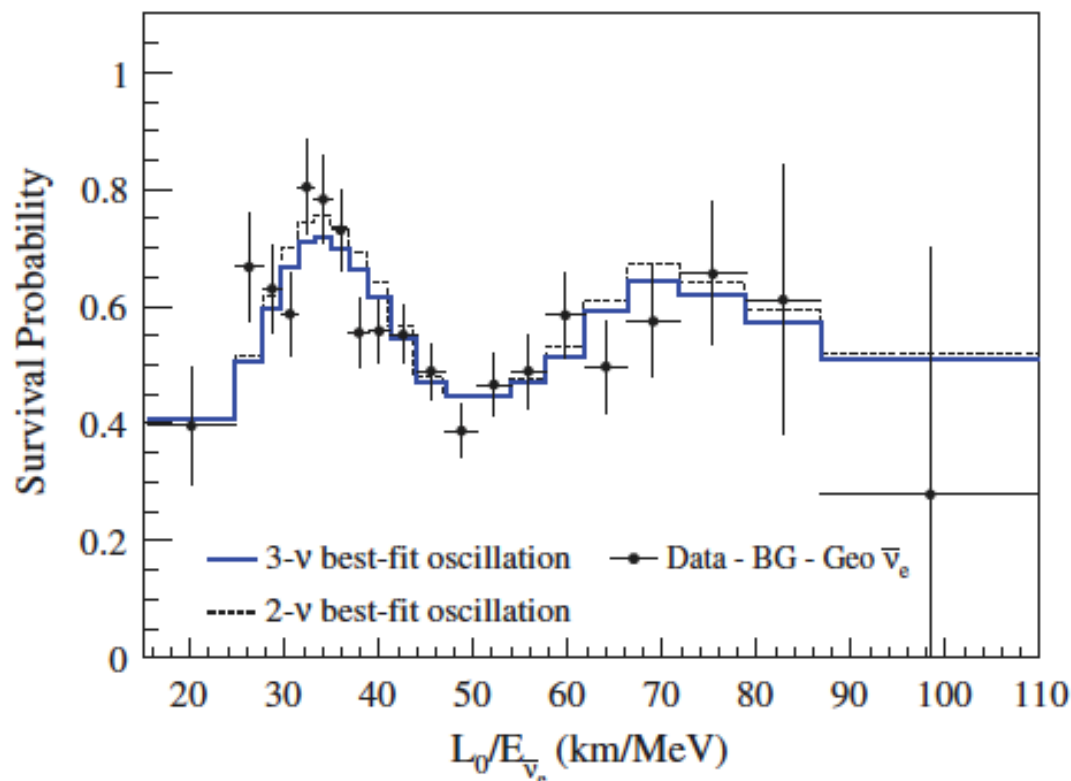
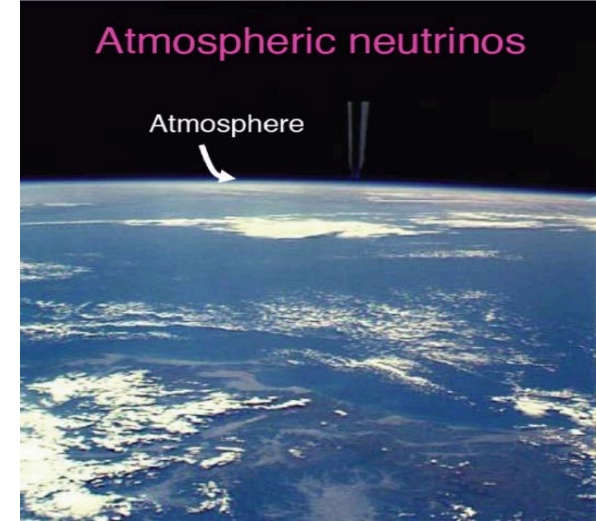
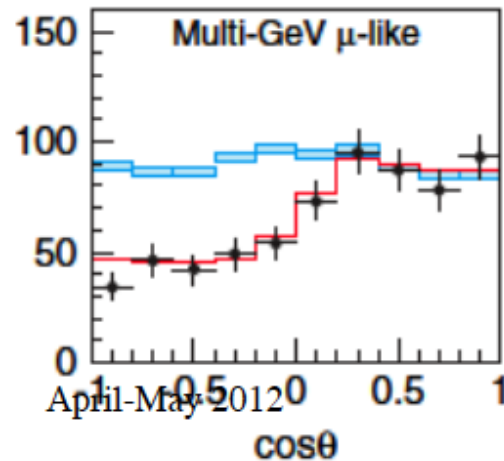
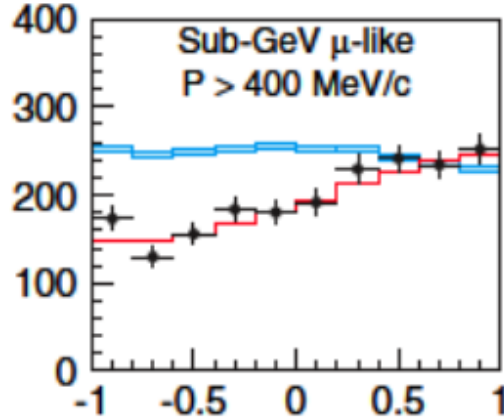
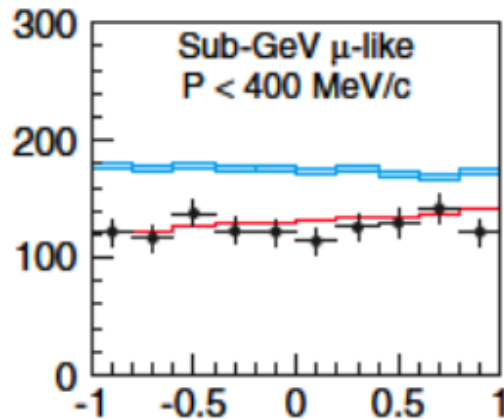


FIG. 5 (color). Ratio of the observed $\bar{\nu}_e$ spectrum to the expectation for no-oscillation versus L_0/E for the KamLAND data. $L_0 = 180$ km is the flux-weighted average reactor baseline. The 2- ν and 3- ν histograms are the expected distributions based on the best-fit parameter values from the two- and three-flavor unbinned maximum-likelihood analyses of the KamLAND data.

Atmospheric neutrinos



Atmospheric ν in a nutshell



- μ -like events at medium-high energies have strong zenith angle dependence (not in MC)
- μ -like events at low energies do not have strong zenith angle dependence



consistent with ν oscillation

Neutrino 1998 @Takayama

1998, @Takayama
June 1998

Atmospheric neutrino results
from Super-Kamiokande & Kamiokande

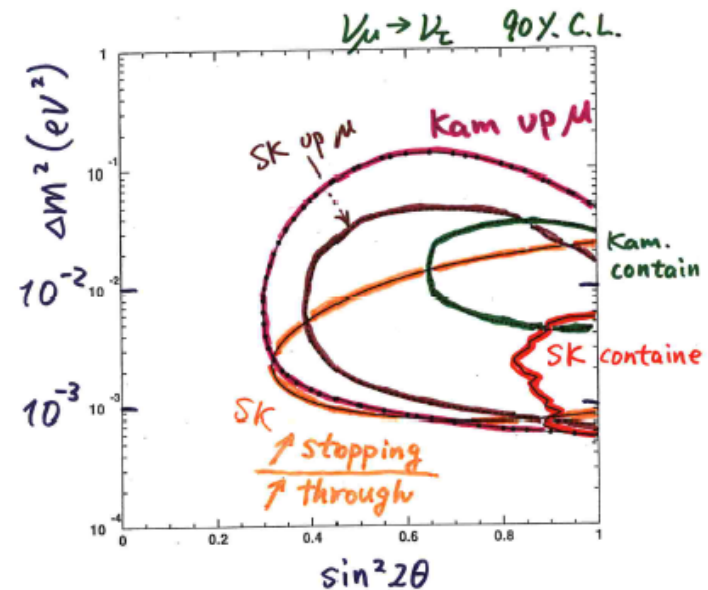
— Evidence for ν_μ oscillations —

T. Kajita
Kamioka observatory, Univ. of Tokyo
for the { Kamiokande
Super-Kamiokande } Collaborations



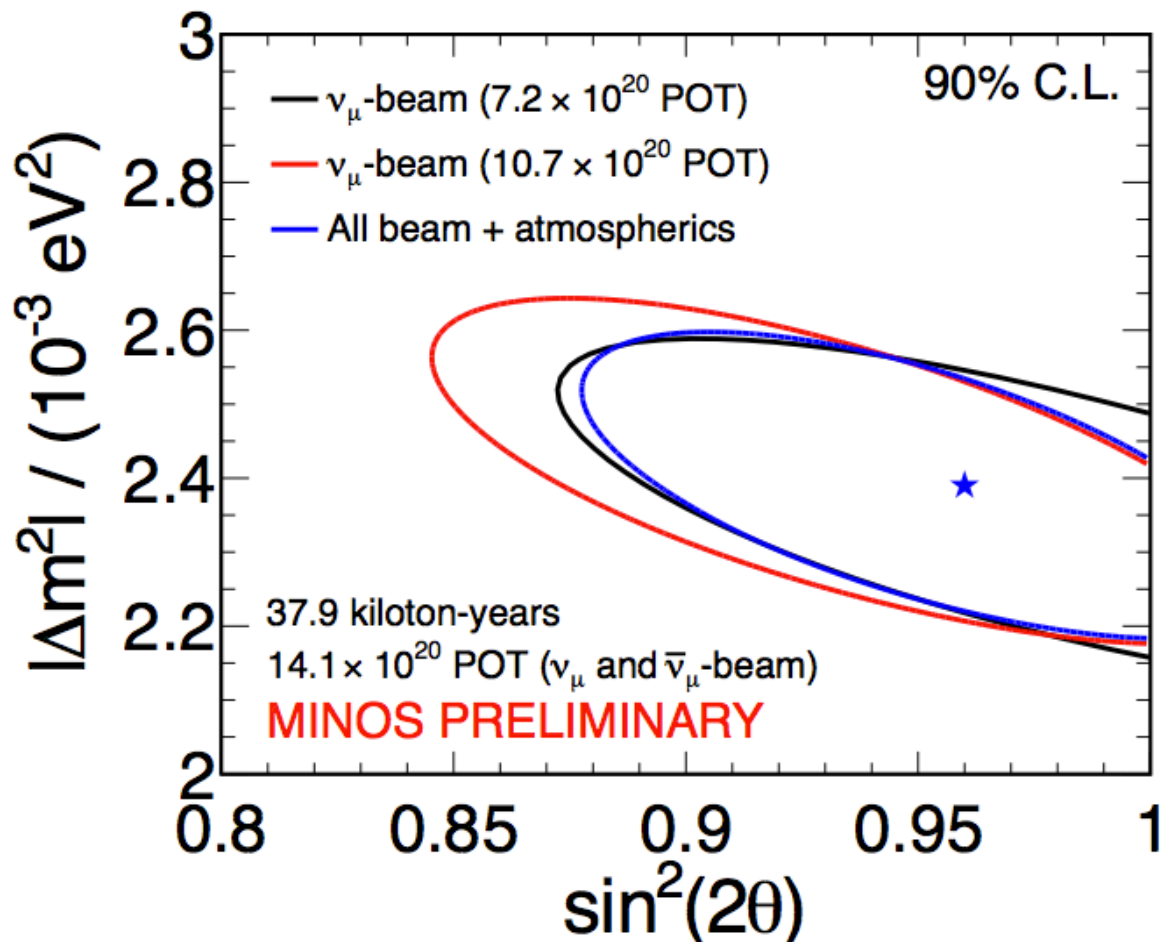
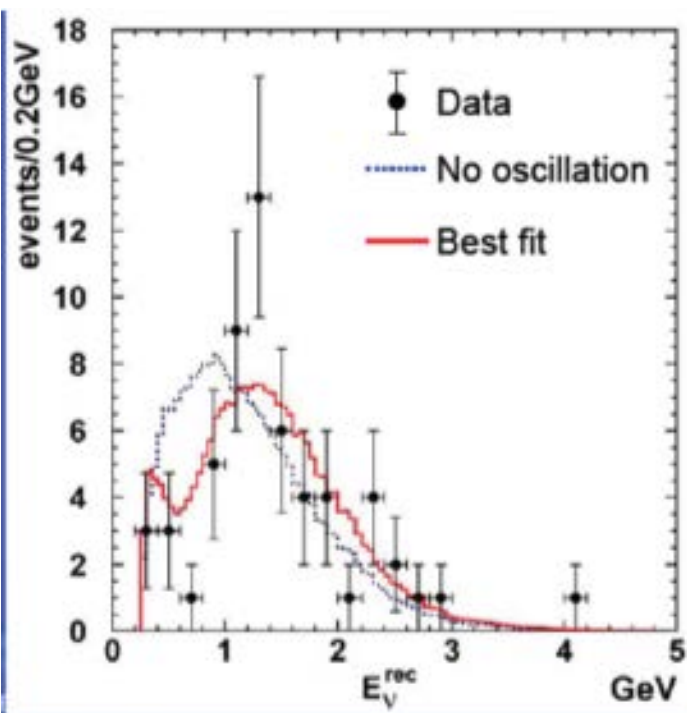
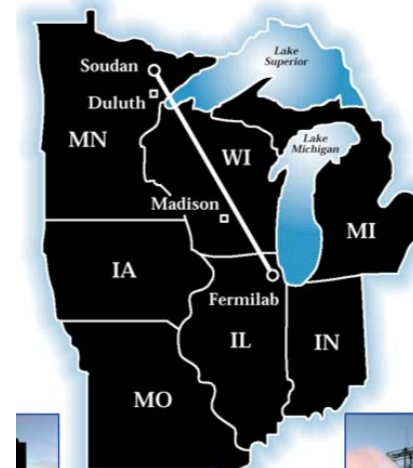
Summary

Evidence for ν_μ oscillations



$$\bullet \begin{cases} \sin^2 2\theta > 0.8 \\ \Delta m^2 \sim 10^{-3} \sim 10^{-2} \end{cases}$$

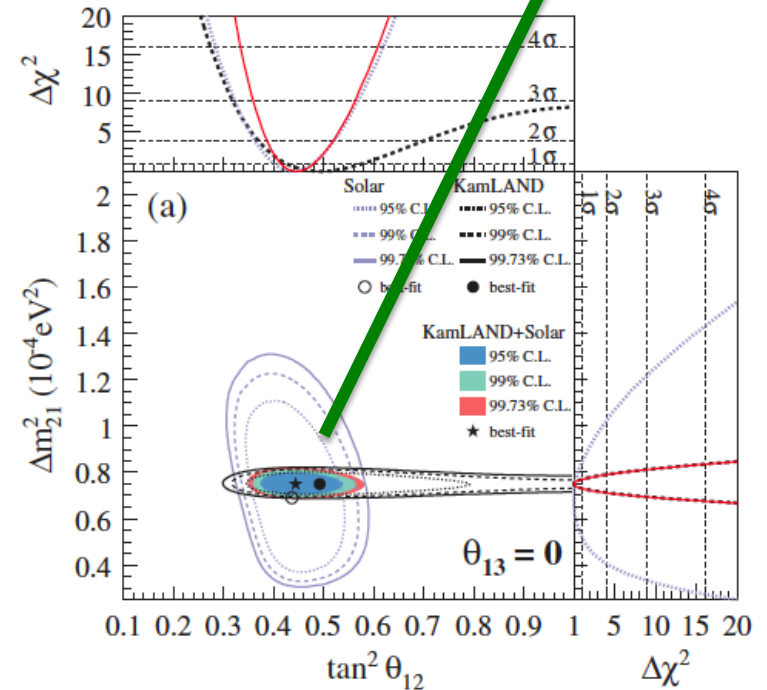
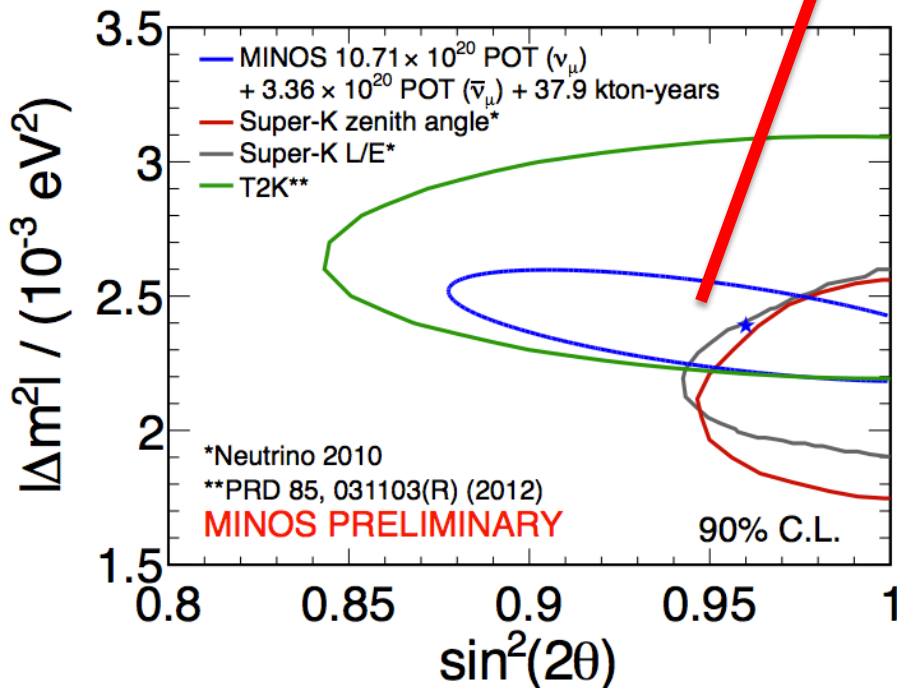
($\bullet \nu_\mu \rightarrow \nu_e$ or $\nu_\mu \rightarrow \nu_s$?)



We enjoy redundancy: Each of θ_{12} and θ_{23} is measured by two ways

$$\nu_\alpha = U_{\alpha i} \nu_i$$

$$U_{MNS} = U_{23}U_{13}U_{12} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{bmatrix} \begin{bmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{bmatrix} \begin{bmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{bmatrix}$$



SK-atm+T2K+MINOS

October 3, 2012

Colloquium@Kavri-IPMU

solar+KamLAND

Anthropic principle?

- Two neutrino mass² difference turn out to be comparable to sizes in the solar system:
- $\Delta m^2_{\text{atm}} / E \sim (\text{earth radius})^{-1}$
- $\Delta m^2_{\text{solar}} / E \sim \text{solar density} / G_F$

Nature's kindness or something behind?

A year
from June
2011-June
2012:
Year of θ_{13}



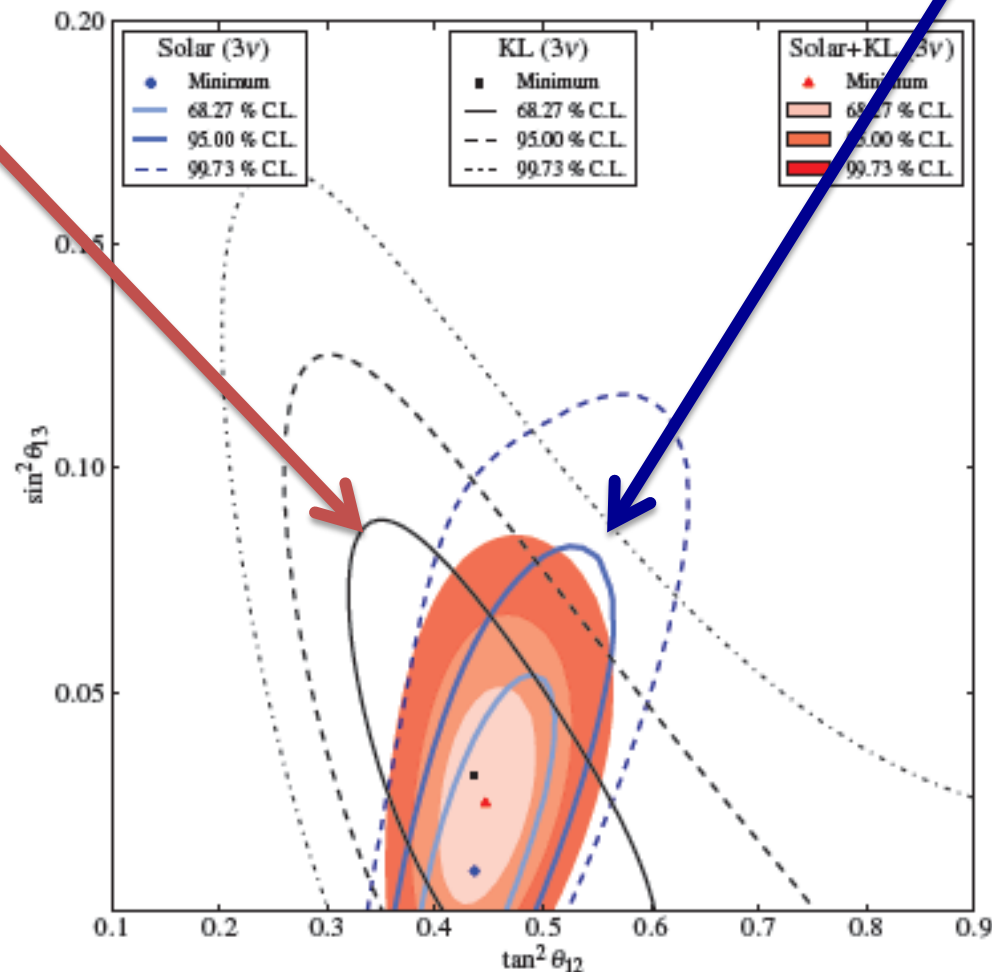
θ_{13} before direct measurement (2010-2011)

KamLAND

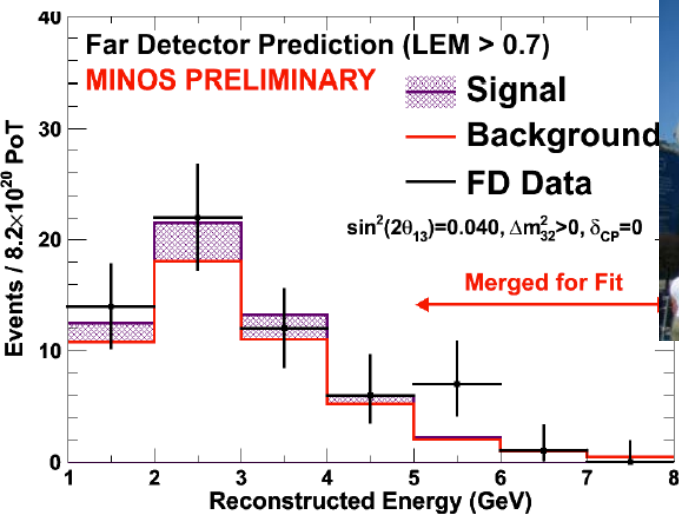
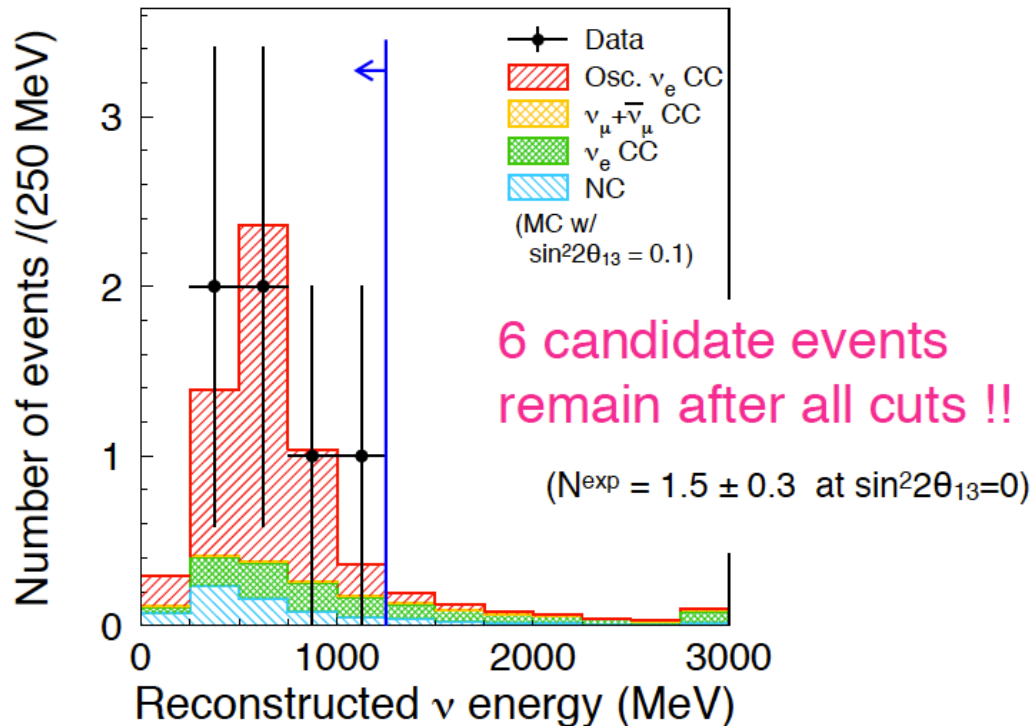
solar

$$P_{ee} \approx \cos^4 \theta_{13} \left(1 - \frac{1}{2} \sin^2 2\theta_{12} \right)$$

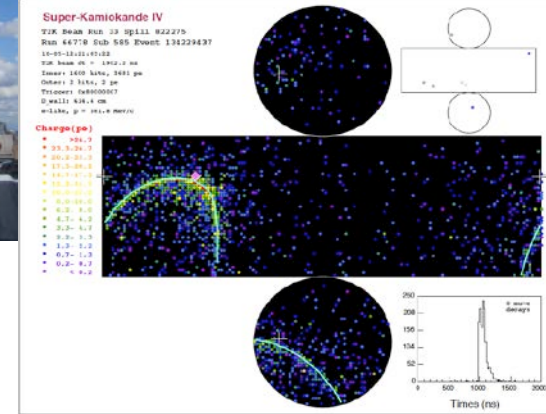
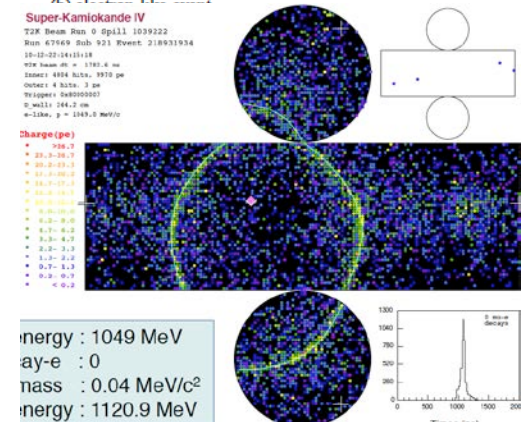
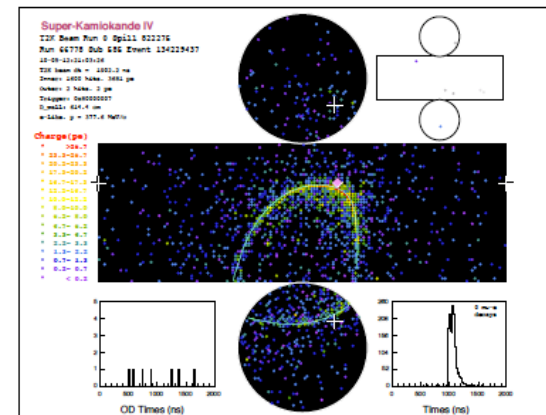
$$P_{ee} \approx \cos^4 \theta_{13} \sin^2 \theta_{12}$$



June 2011: T2K 6 events: start of year of θ_{13}



Auror @ Saint Petersburg

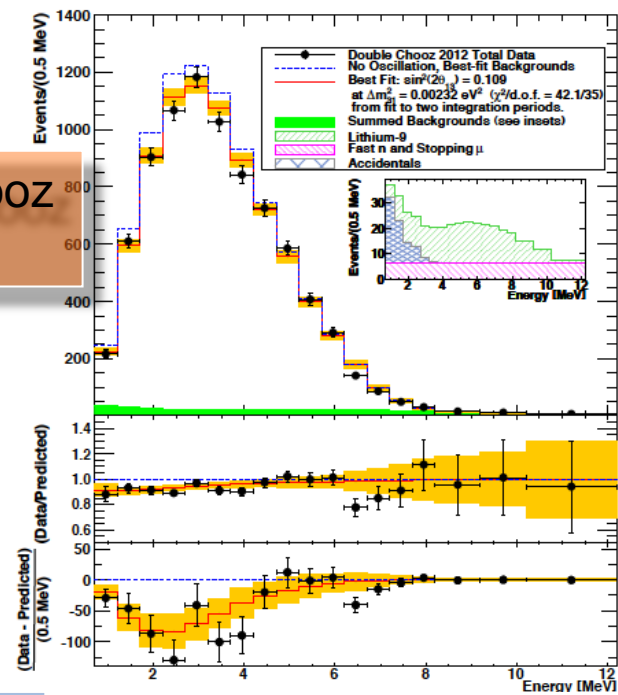
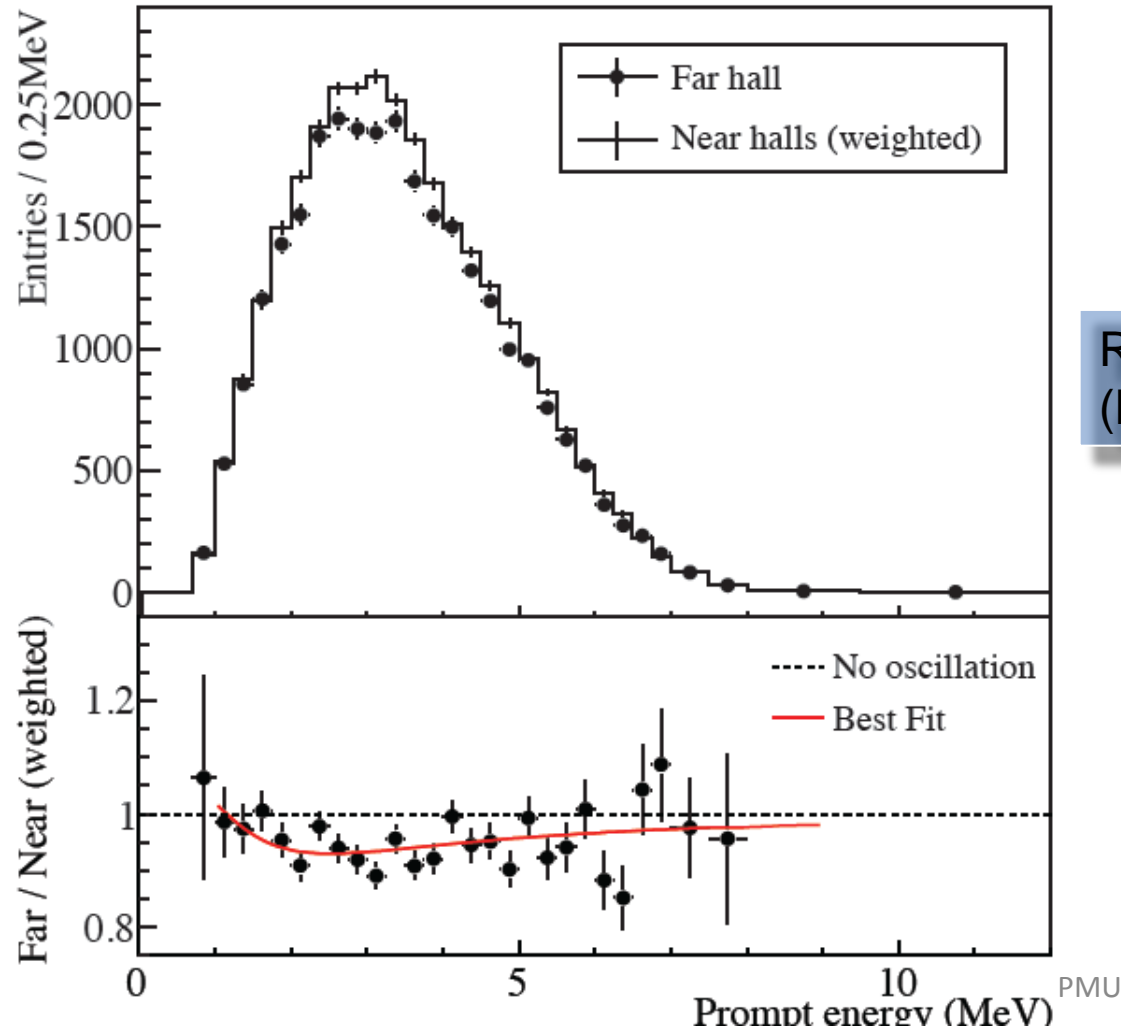


Reactor θ_{13} experiments

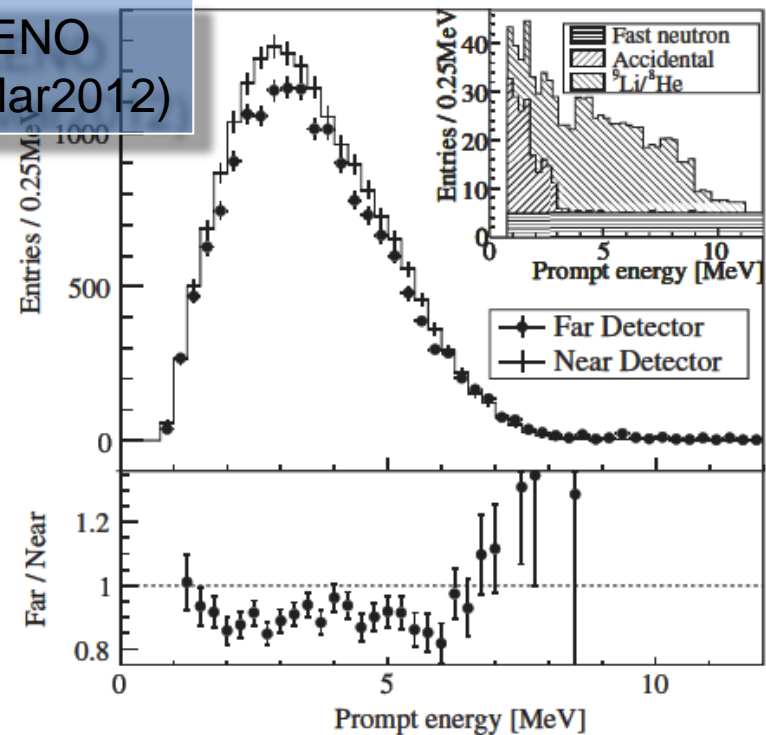
Daya Bay (Nu2012)

Double Chooz
(Jul.2012)

7.7 σ evidence for nonzero θ_{13} !



RENO
(Mar2012)

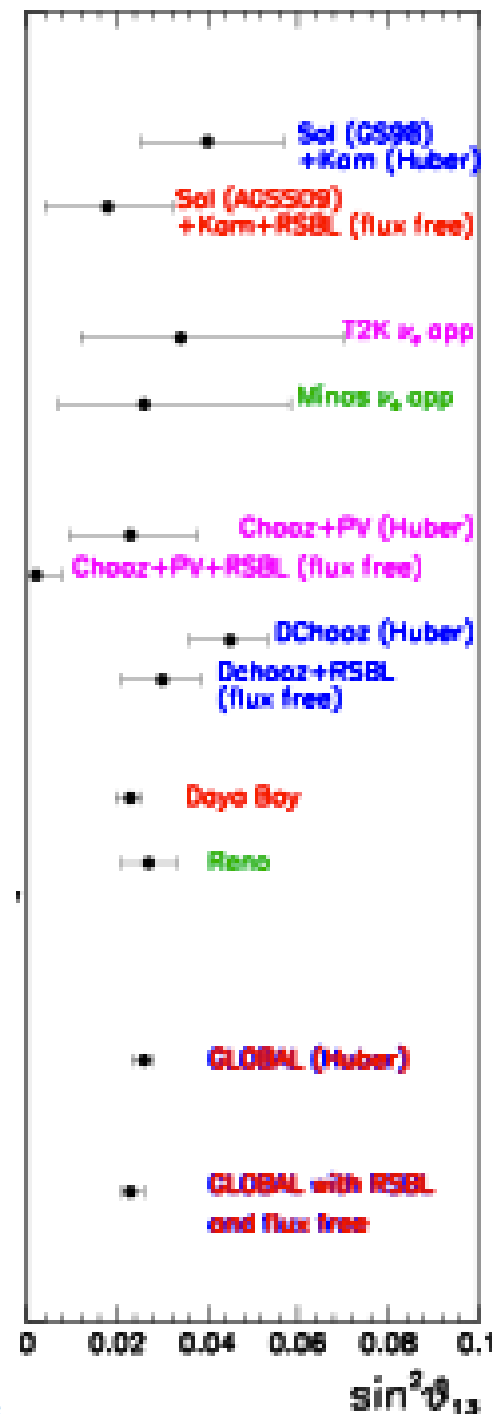
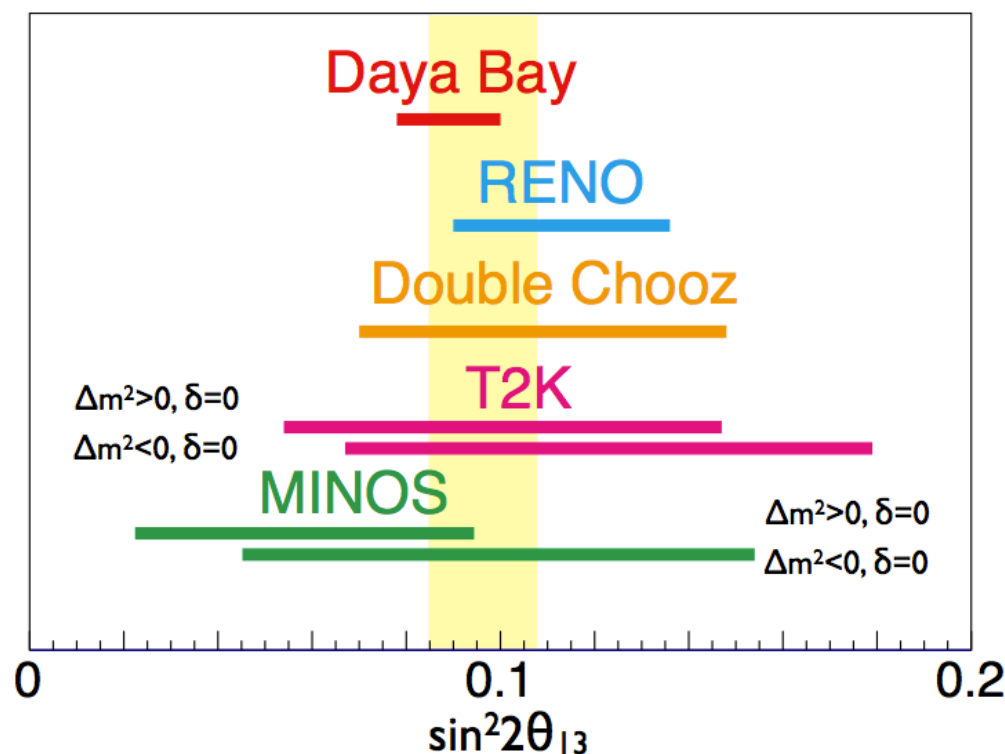


$$\sin^2 2\theta_{13} = 0.090^{+0.008(0.027)}_{-0.010(0.029)}$$

M.C.Gonzalez-Garcia et al. 2012



$\theta_{13} \neq 0$ established...



Now is the time to move forward to the next step!

What large θ_{13} means?

- It is natural because $U_{MNS}=U_l^\dagger U_\nu$ and θ_{12} and θ_{23} are large it is unlikely that only θ_{13} is extremely small
HM in Nu2008 at Christchurch
- This argument is invalidated if there is a symmetry which enforce $\theta_{13}=0$
- So my interpretation is: there is no such symmetry (sorry, negative statement)
- Anarchy ? Director's choice
- I have my own favorite (see later)

We now know
all the 3
lepton mixing
angles!



What we know

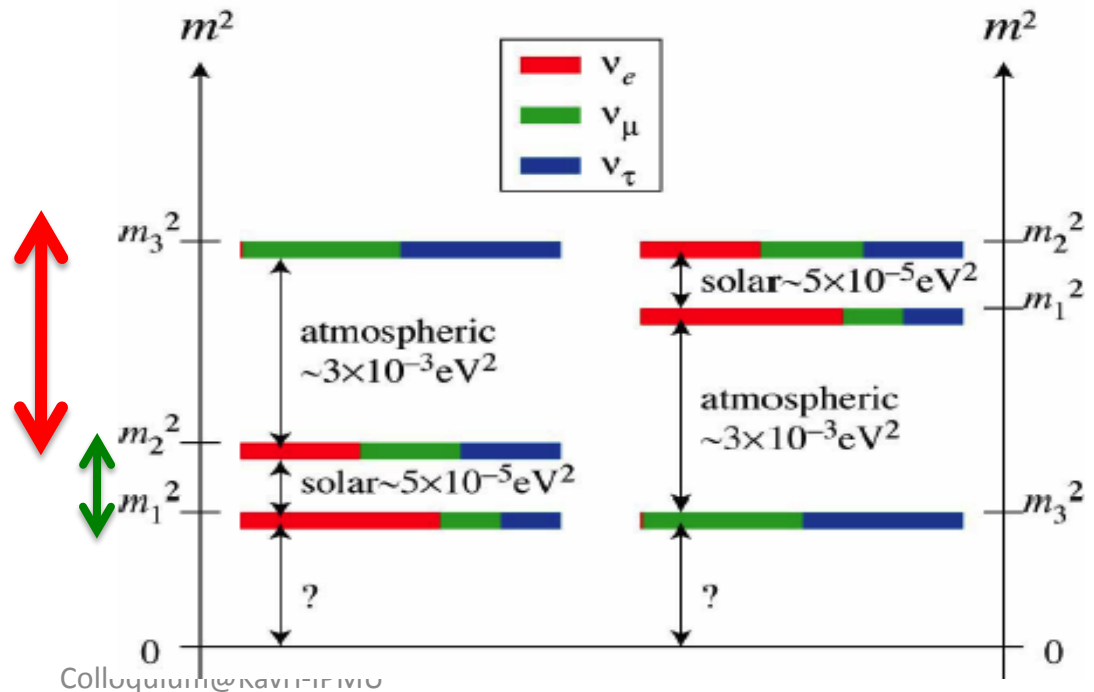
$$\nu_{\alpha} = U_{\alpha i} \nu_i$$

$$U_{MNS} = U_{23}U_{13}U_{12} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{bmatrix} \begin{bmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{bmatrix} \begin{bmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

$$\theta_{12} = 34^{\circ}, \theta_{23} \sim 45^{\circ}, \theta_{13} = 9^{\circ}$$

Atmospheric mass²
splitting = $2.4 \times 10^{-3} \text{ eV}^2$

Solar mass² splitting =
 $7.5 \times 10^{-5} \text{ eV}^2$



What we do not know

Lepton KM phase

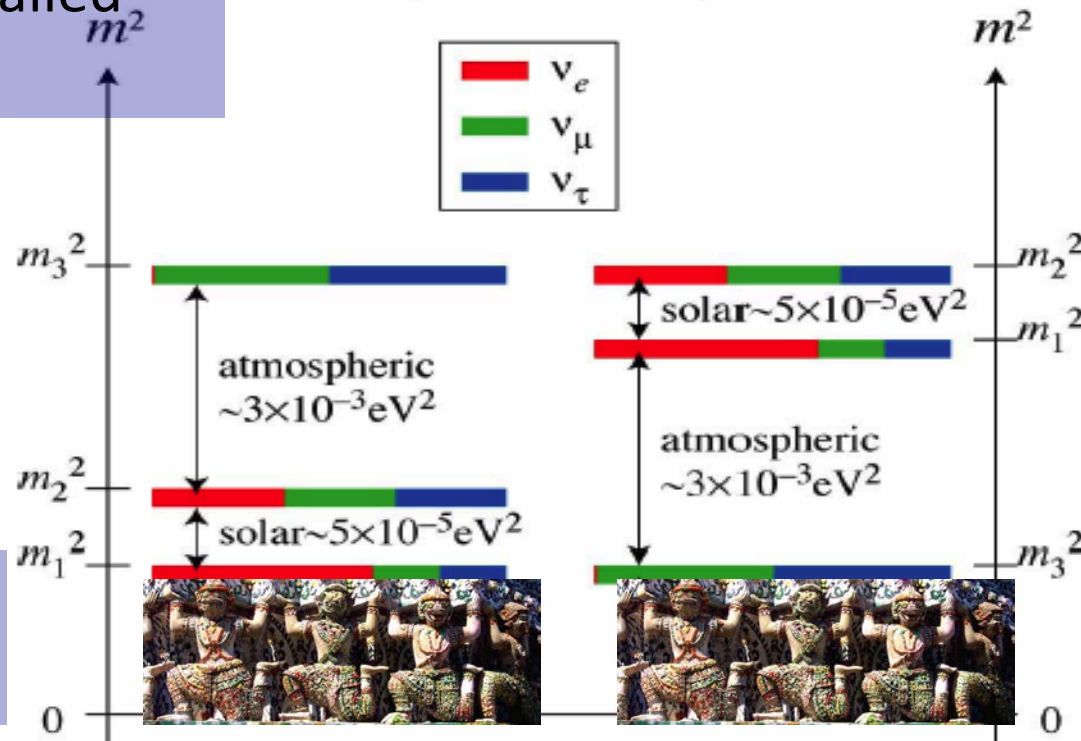
$$U = \begin{bmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{bmatrix} \begin{bmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{bmatrix} \begin{bmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} e^{i\alpha_1} & 0 & 0 \\ 0 & e^{i\alpha_2} & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

Another CP phases:
If Majorana

Which mass pattern?
(sometimes called
“hierarchy”)

Normal
ordering

Absolute
ν masses



inverted
ordering



Then what is
remain to be
understood?

There are 2 aspects

Experimentally

- We now know 3 mixing angles and $\Delta m^2_{\text{solar}}$, Δm^2_{atm}
- What's remain to be explored experimentally?
- Two “disjoint” frontier

1. Lepton KM phase, mass pattern

2. Absolute ν mass, Majorana vs. Dirac, Majorana phase

Theoretically

- We roughly know ν mass scale
- What the mass scale means?
- Lepton mixing pattern is quite different from quarks, what that means?



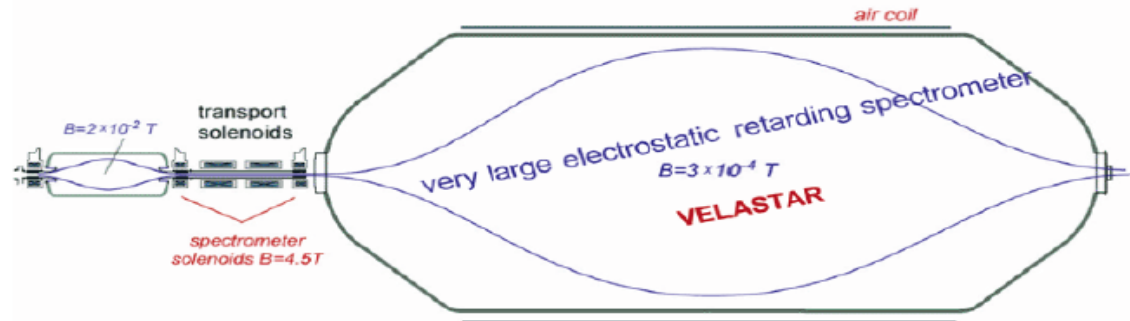
Absolute v mass, Majorana vs. Dirac

KATRIN: endpoint of beta decay spectrum..

electrostatic spectrometers: tandem design

electrostatic pre-filtering & analysis of tritium β -decay electrons

Sensitivity reach
 $m_\nu \sim 0.2$ eV



pre-spectrometer

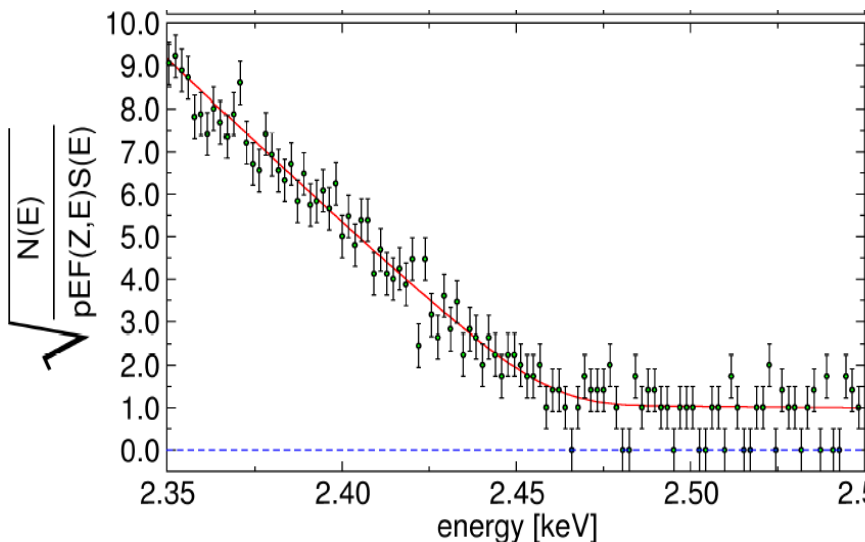
fixed retarding potential 18.5-18.6 kV
 $\varnothing = 1.7$ m / $L = 3.5$ m
 $\Delta E = 70$ eV

main spectrometer

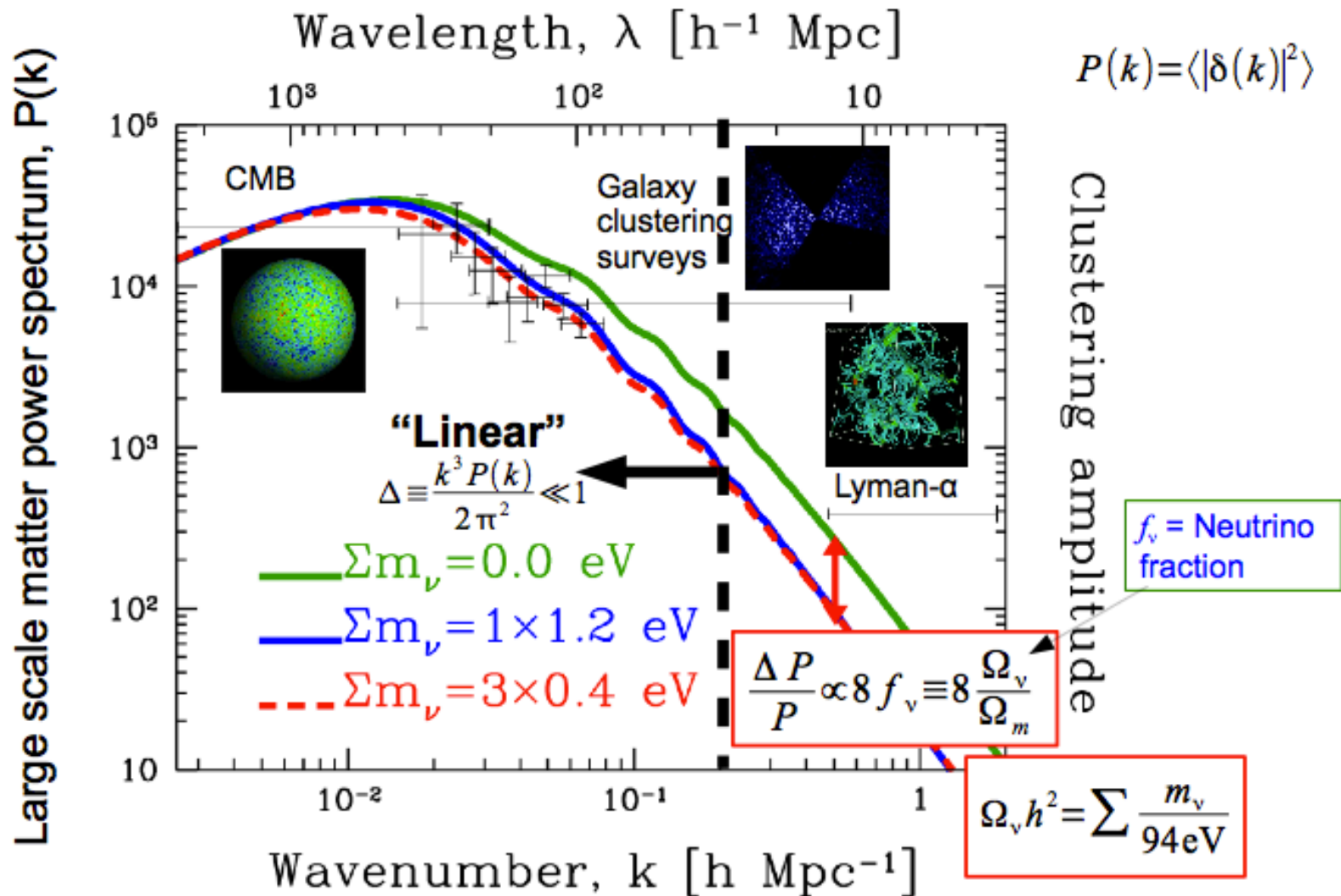
variable retarding potential 18.5-18.6 kV
 $\varnothing = 10$ m / $L = 22$ m
 $\Delta E = 1$ eV

XHV conditions $p < 10^{-11}$ mbar : main challenge

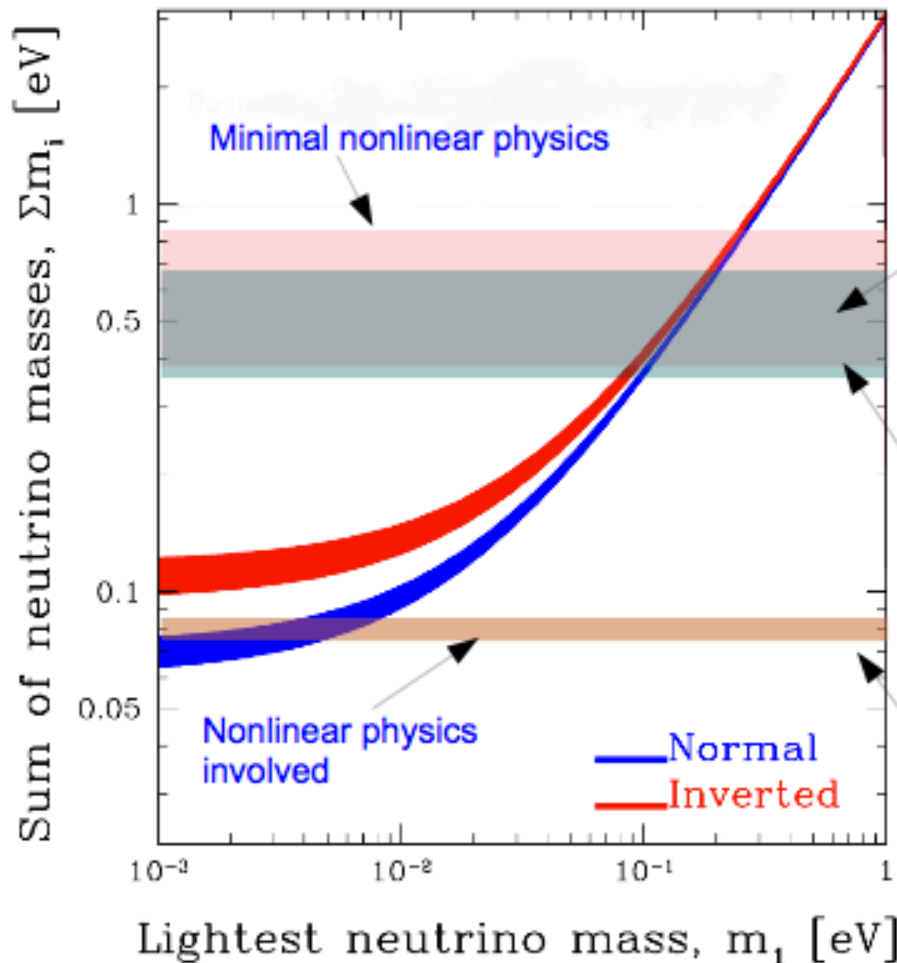
Figure: Pre-Spectrometer and Main Spectrometer



Cosmology: ν suppresses small scale structure



Present constraints and future sensitivities...



CMB (WMAP7+ACBAR+BICEP+QuaD)
+ LSS (SDSS-HPS)
+ H_0 +SNIa

$$\Sigma m_\nu < 0.36 \rightarrow 0.76 \text{ eV (95\% CI)}$$

depending on the model complexity

Hannestad, Mirizzi, Raffelt & Y³W 2010
Gonzalez-Garcia et al. 2010
de Putter et al. 2011, etc.

Planck alone (1 year) **2013**

$$\Sigma m_\nu < 0.38 \rightarrow 0.84 \text{ eV (95\% CI)}$$

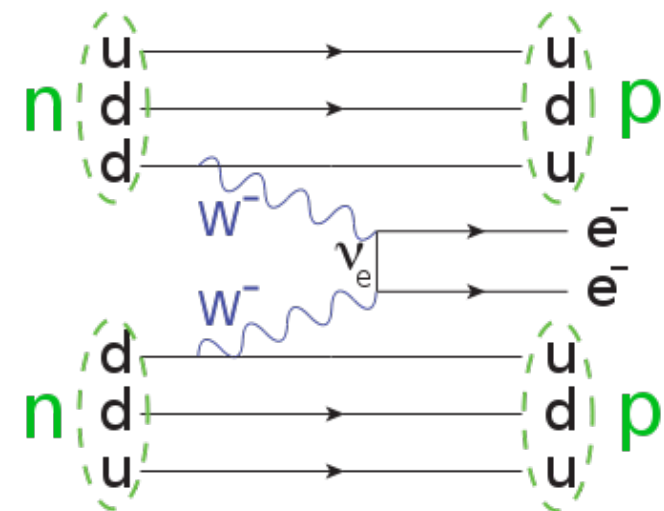
Perotto et al. 2006

Planck+Weak lensing (**Euclid**) **2020+**

$$\Sigma m_\nu < 0.074 \rightarrow 0.086 \text{ eV (95\% CI)}$$

Hannestad, Tu & Y³W 2006

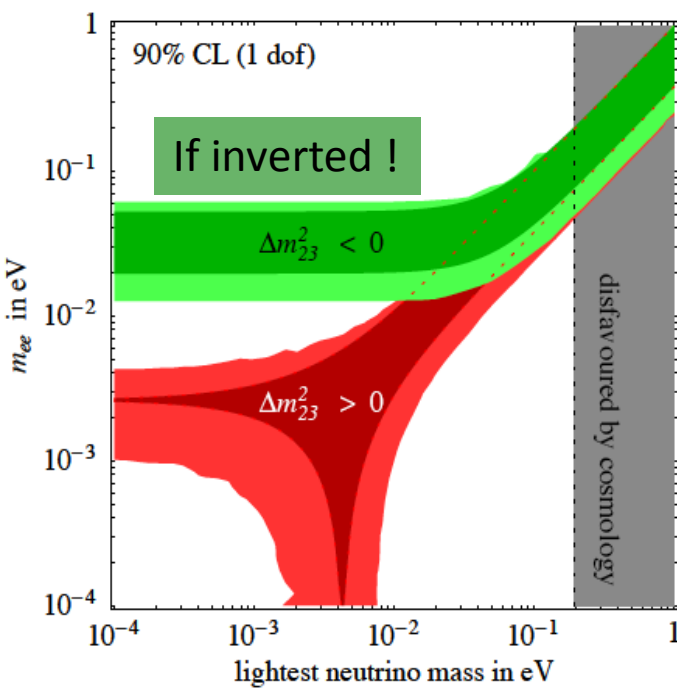
Majorana vs. Dirac: double β decay



$$(T_{1/2}^{0\nu})^{-1} = G^{0\nu}(Q, Z) |M^{0\nu}|^2 \langle m_{ee} \rangle^2$$

$$\langle m_{ee} \rangle = \left| \sum_i U_{ei}^2 m_i \right|$$

Simplified List of Limits for $\beta\beta 0\nu$ decay



Candidate nucleus	Detector type	(kg yr)	Present $T_{1/2}^{0\nu\beta\beta}$ (yr)	$\langle m \rangle$ (eV)
^{48}Ca	Ge diode	~47.7	$> 5.8 \cdot 10^{22}$ (90%CL)	< 0.35
^{76}Ge			$> 1.9 \cdot 10^{25}$ (90%CL)	
^{82}Se			$> 2.1 \cdot 10^{23}$ (90%CL)	
^{96}Zr			$> 9.2 \cdot 10^{21}$ (90%CL)	
^{100}Mo			$> 5.8 \cdot 10^{23}$ (90%CL)	
^{116}Cd	Foil.Geiger tubes		$> 1.7 \cdot 10^{23}$ (90%CL)	$< 0.19 - 0.68$
^{128}Te			$> 1.1 \cdot 10^{23}$ (90%CL)	
^{130}Te			$> 3 \cdot 10^{24}$ (90%CL)	
^{136}Xe	Xe scint	~4.5	$> 1.2 \cdot 10^{24}$ (90%CL)	$< 1.1 - 2.9$
^{150}Nd			$> 1.8 \cdot 10^{22}$ (90%CL)	
^{160}Gd			$> 1.3 \cdot 10^{21}$ (90%CL)	

With regard
to mass
hierarchy &
CP we have
clearer view



Is large θ_{13} good news? Matter vs. CP

- $P_{\mu e} = P_{\text{atm}}(\text{matter}) + P_{\text{interfere}}(\delta) + P_{\text{solar}}$

$(\theta_{13})^2$

Merit for mass hierarchy

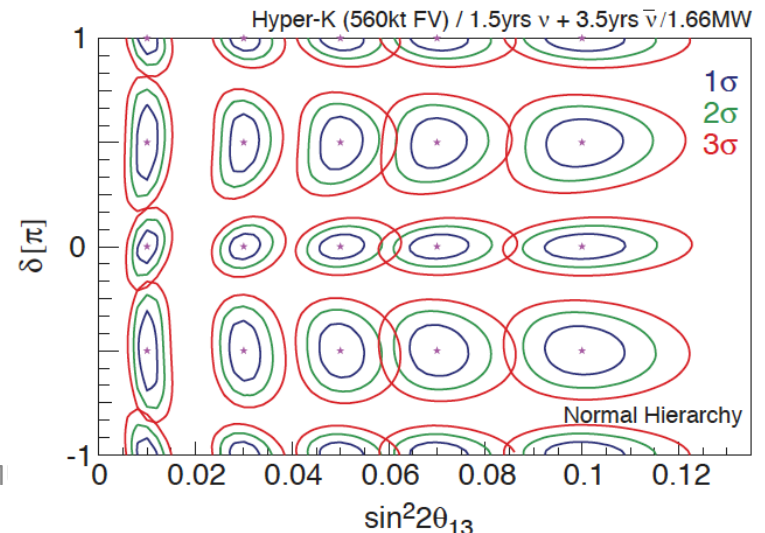
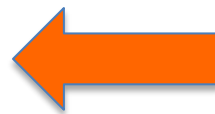
$(\theta_{13})^1$

Merit / demerit for δ

θ_{13} -indep

- IS a good news for mass hierarchy
- Need not be a good news for CP because:
- $A = (P_{\mu e} - \text{anti-}P_{\mu e}) / (P_{\mu e} + \text{anti-}P_{\mu e}) \sim 1/\theta_{13}$

Sensitivity to δ is roughly independent of θ_{13} (HK-LOI)

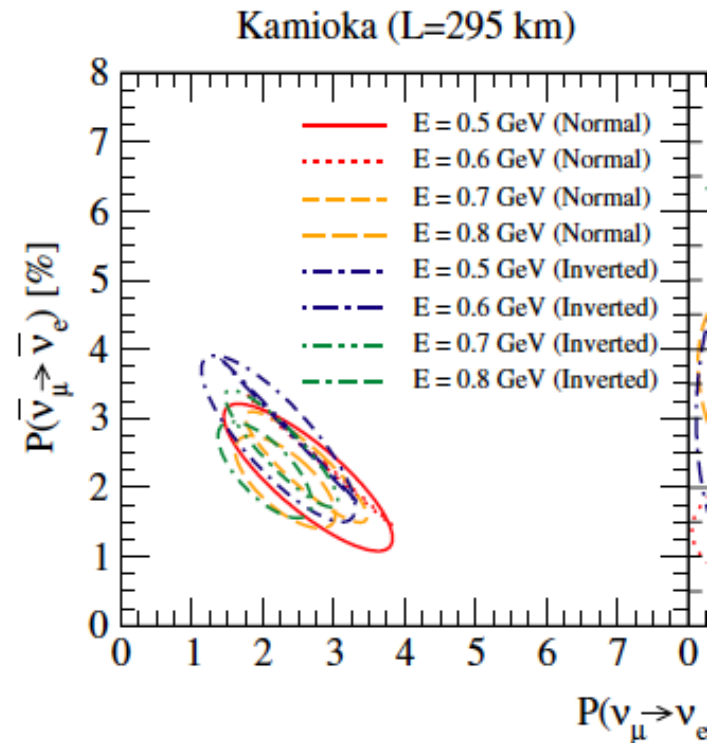
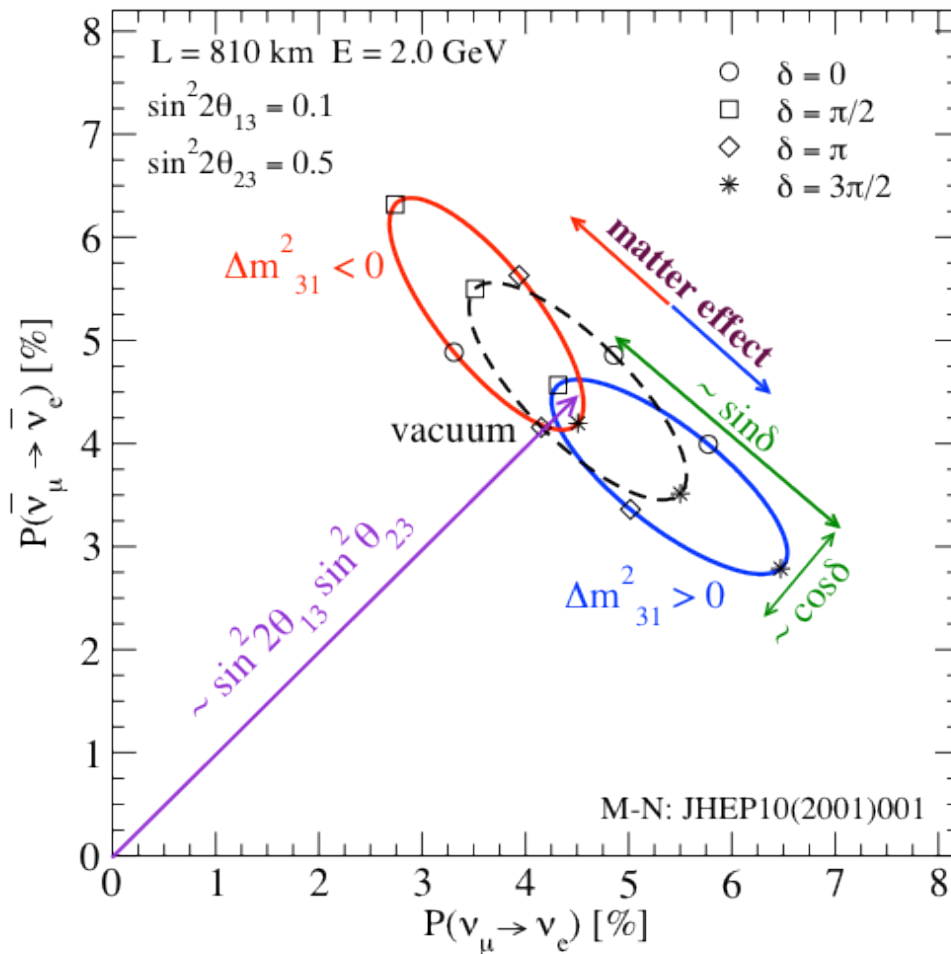


Mass hierarchy



Mass hierarchy resolution and CP: understanding principle

NOVA



bi-probability plot in $P_{\mu e}$ – anti- $P_{\mu e}$ space

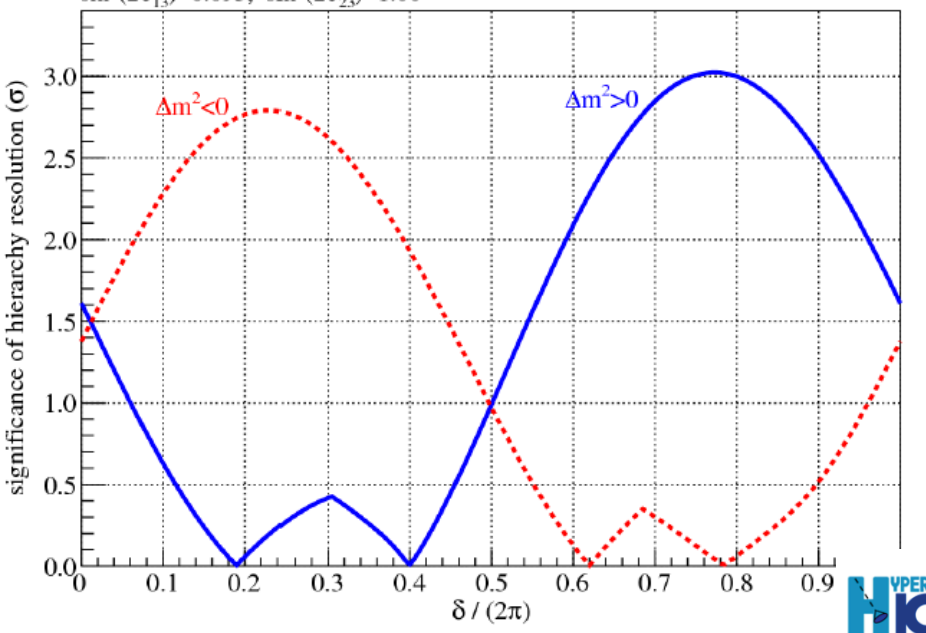
HM-H.Nunokawa,JHEP01

Mass hierarchy: a shopping list

- NOVA, if lucky
- NOVA+INO (up to 3σ)
- JPARC-HK, if lucky
- LBNE (but ?? if on surface..)
- T2KK (Tokai-Kamioka-Korea)
- JPARC-Okinoshima
- CERN-Pihasarmi
- Atmospheric neutrinos (earth matter effect)
- Reactor neutrinos (atm wiggle at solar scale oscillation)

NOvA hierarchy resolution, 3+3 yr ($\nu + \bar{\nu}$)

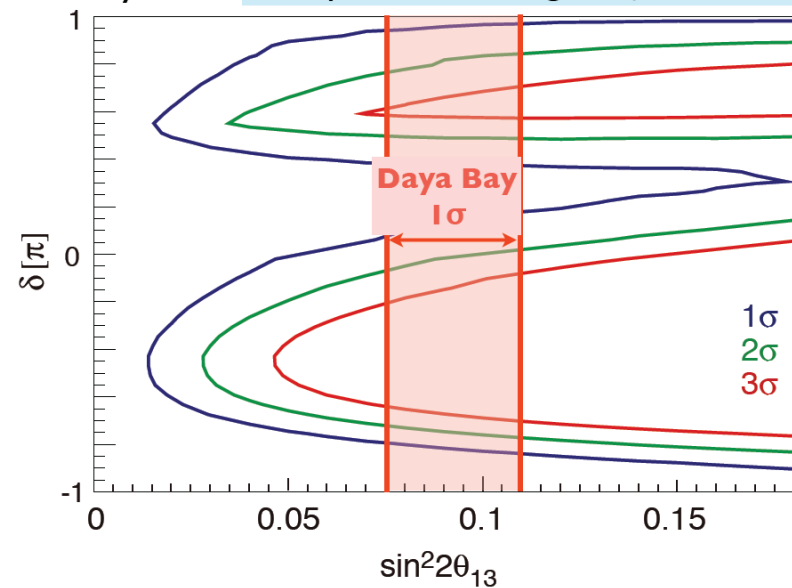
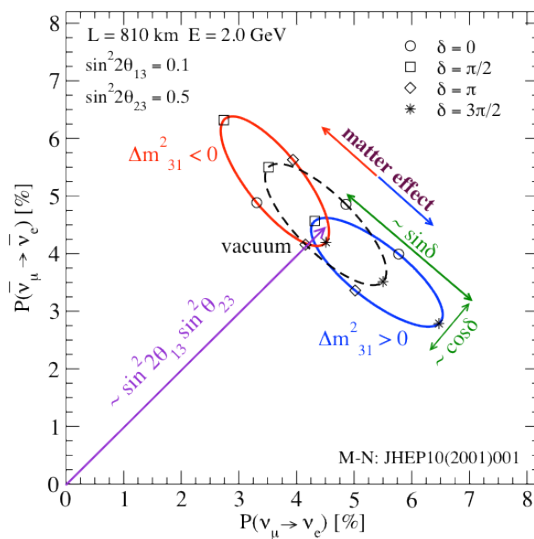
$\sin^2(2\theta_{13})=0.095$, $\sin^2(2\theta_{23})=1.00$



Mass hierarchy

0.75MW×10yrs

5% systematics on signal, ν_μ BG, ν_e BG, $\nu/\bar{\nu}$



Chance to determine MH! See also later talk (atm ν)

October 3, 2012

Col

Masashi Yokoyama (U.Tokyo)

1st open meeting for Hyper-K project, Aug.22-23 2012

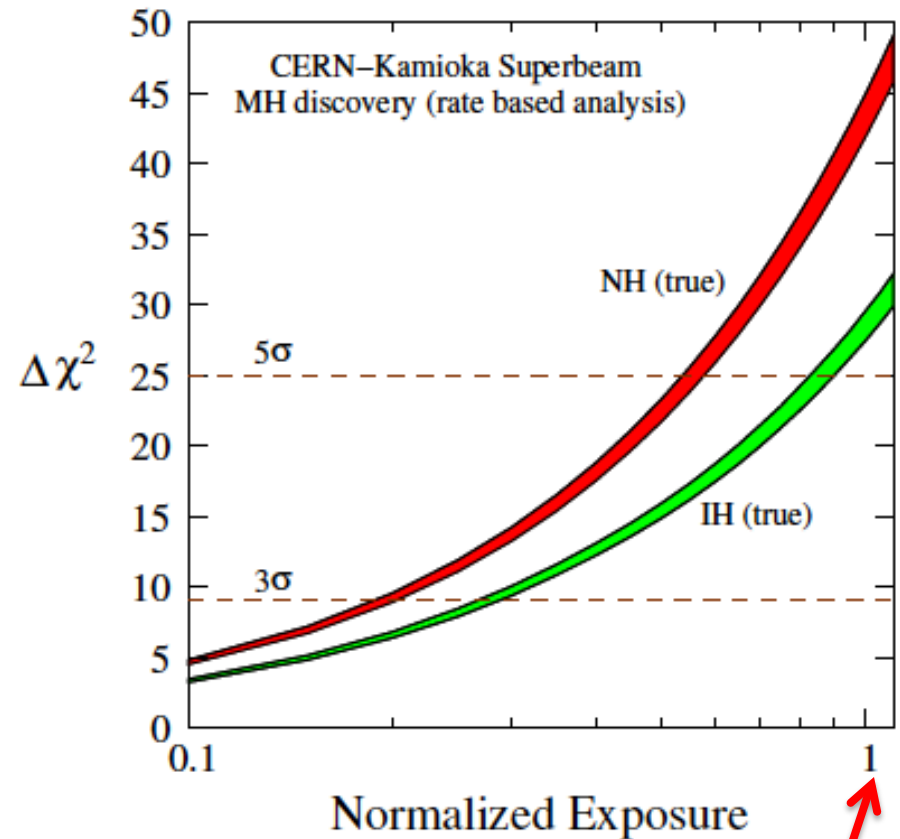
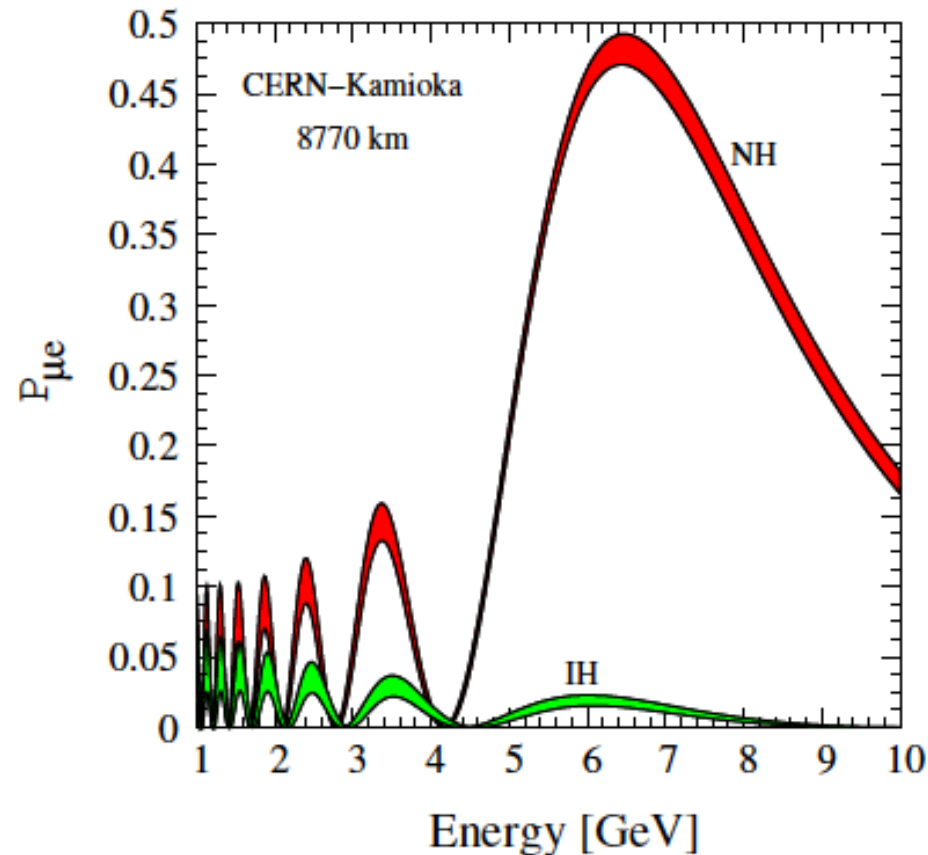
LBL and p-decay with Hyper-K

19

Wednesday, August 22, 12

CERN-Super-K (8870 km)

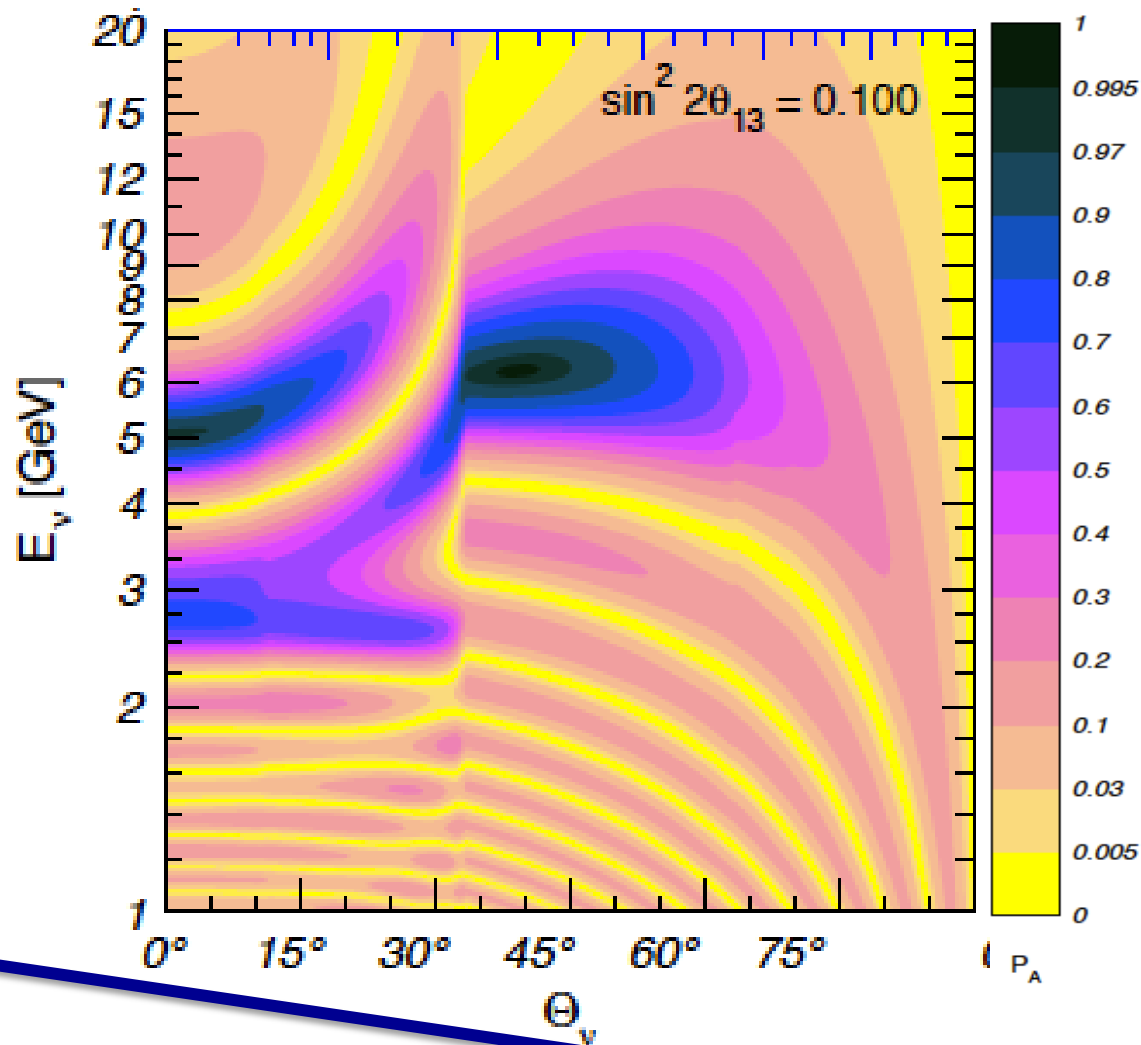
Agarwalla-Hernandez April 12



Channel	CERN-Kamioka (8870 km)	
	Signal	Background
	CC-1 ring	Int+Mis-id+NC = Total
$\nu_{\mu} \rightarrow \nu_e$ (NH)	44	1+2+16=19
$\nu_{\mu} \rightarrow \nu_e$ (IH)	2	1+3+16=20
$\nu_{\mu} \rightarrow \nu_{\mu}$ (NH)	83	2
$\nu_{\mu} \rightarrow \nu_{\mu}$ (IH)	91	2

5×10^{21} pot.

Use of atmospheric ν for mass hierarchy



Akhmedov et al
JHEP07

$$P_{ee} = 1 - P_A,$$

$$P_{\mu e} = P_{e\mu} = s_{23}^2 P_A,$$

$$P_{e\tau} = c_{23}^2 P_A,$$

$$P_{\mu\mu} = 1 - \frac{1}{2} \sin^2 2\theta_{23} - s_{23}^4 P_A + \frac{1}{2} \sin^2 2\theta_{23} \sqrt{1 - P_A} \cos \phi_X,$$

$$P_{\mu\tau} = \frac{1}{2} \sin^2 2\theta_{23} - s_{23}^2 c_{23}^2 P_A - \frac{1}{2} \sin^2 2\theta_{23} \sqrt{1 - P_A} \cos \phi_X, \quad \text{-IPMU}$$

P_A

Atmospheric ν @PINGU

Doug Cowen, NuSky, ICTP, June 2011

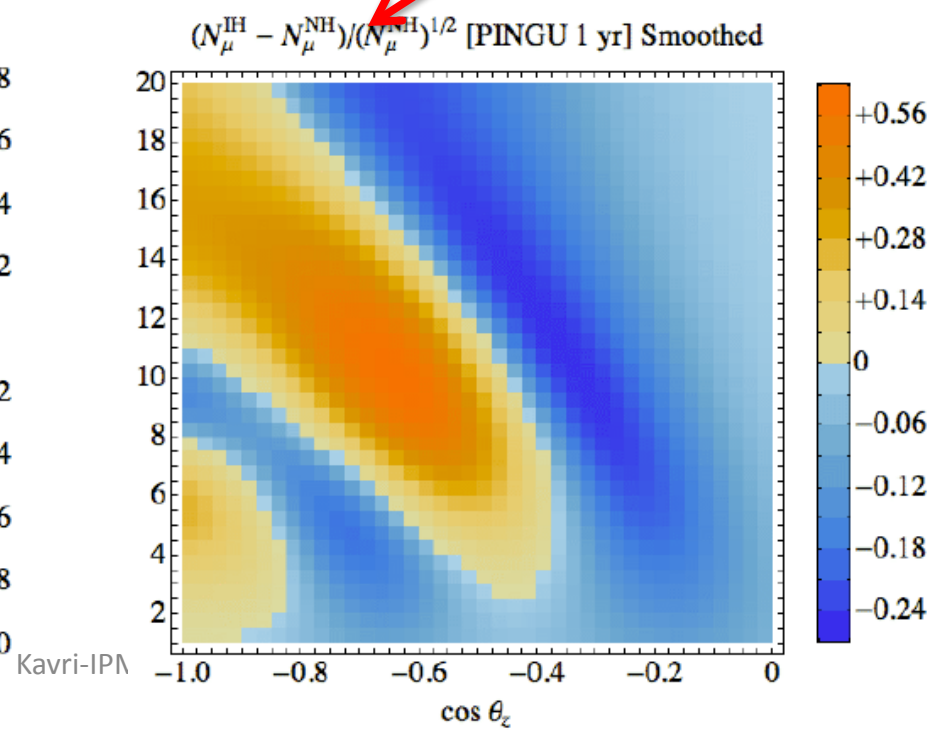
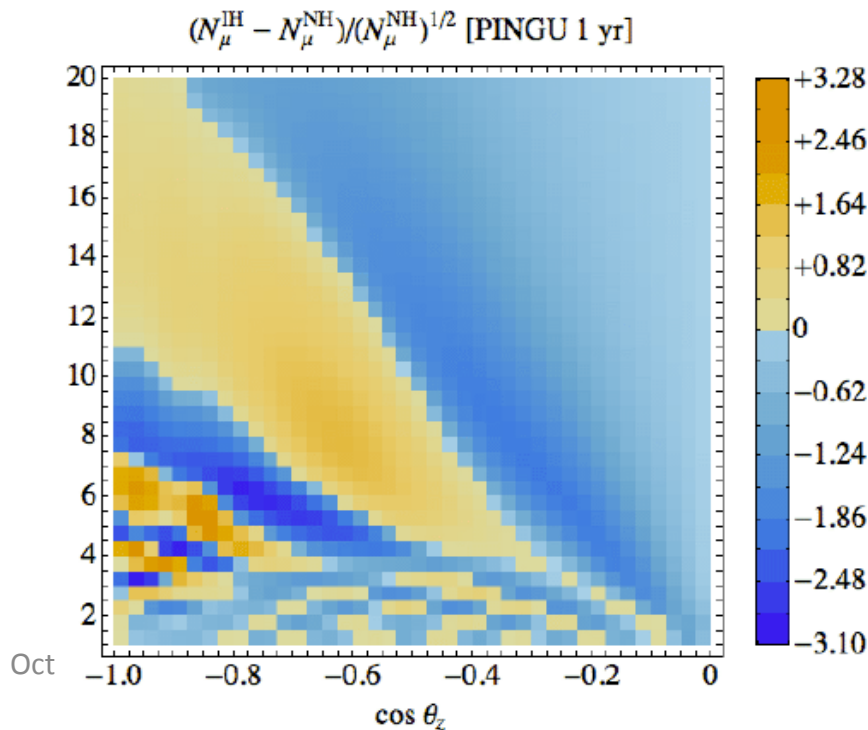
IceCube \rightarrow DeepCore \rightarrow **PINGU**

- ~ 20 additional strings within DeepCore
- lower threshold to few GeV
- ~ 10 Mt effective volume
- construction within 1 yr, $\sim \$25$ M

Akhmedov-Razzaque-Smirnov June 12

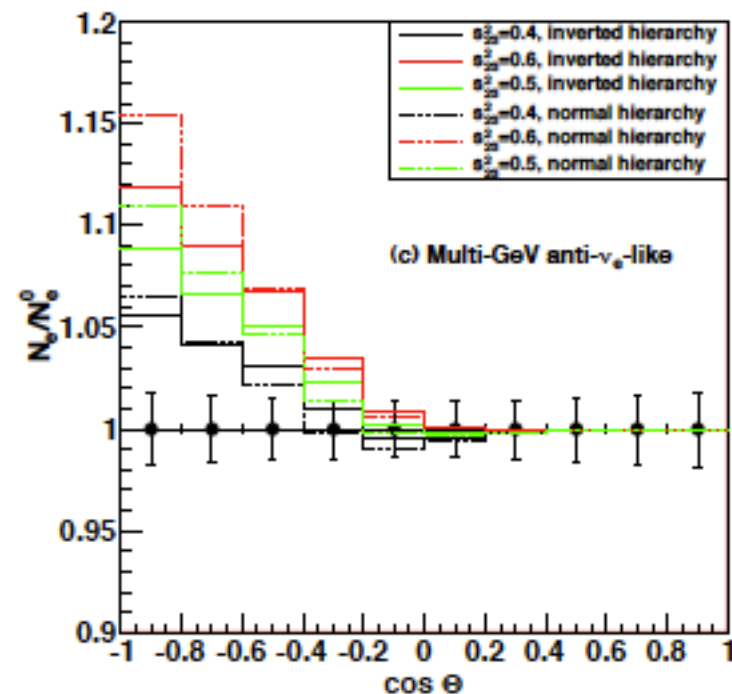
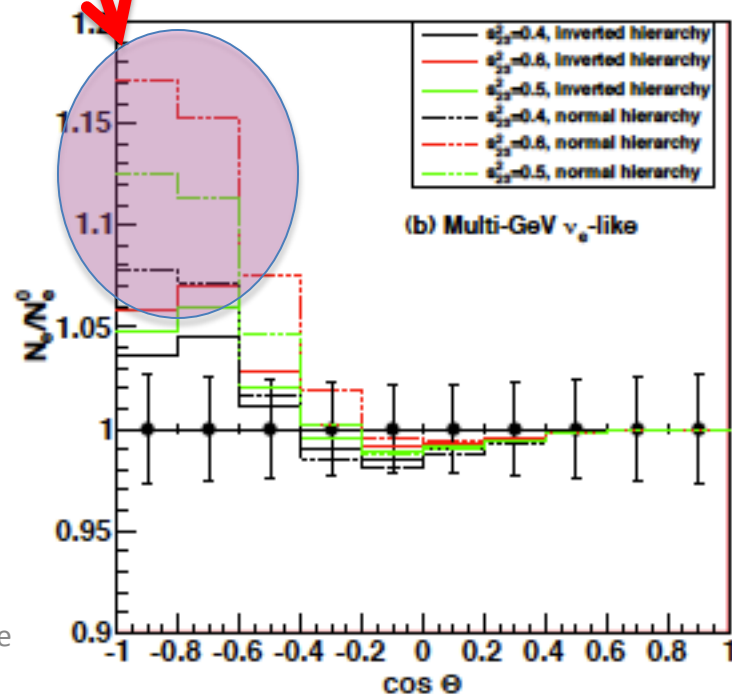
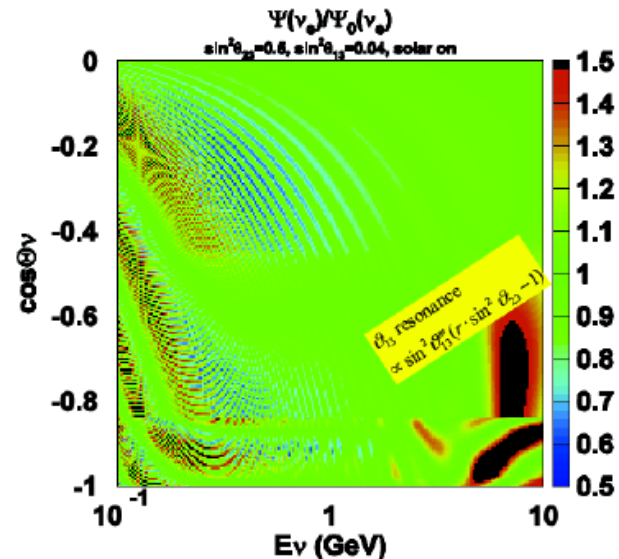
MH resolution 4σ -
 11σ in 5 years !

$\sigma_E = 2$ GeV
 $\sigma_\theta = 11.25^\circ$

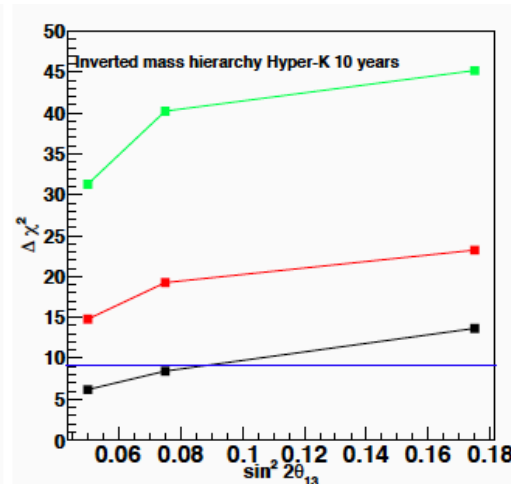
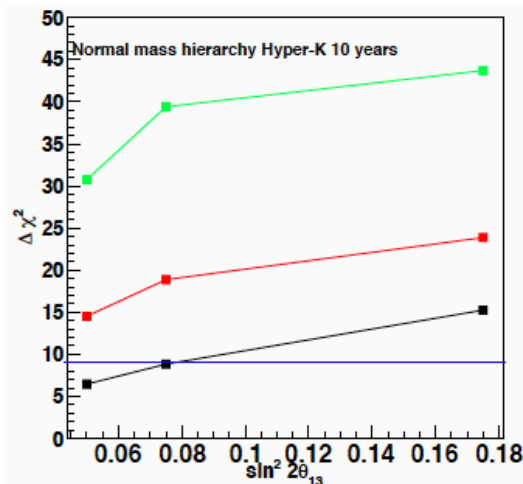
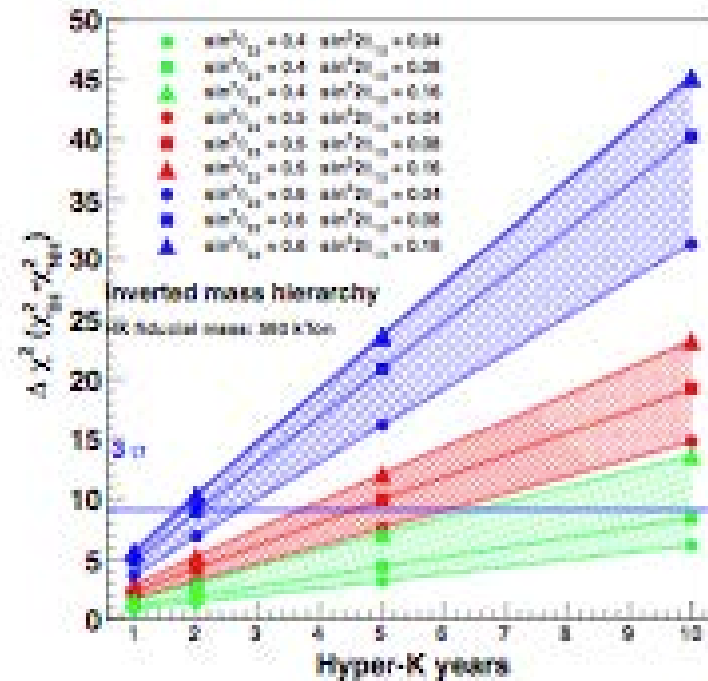
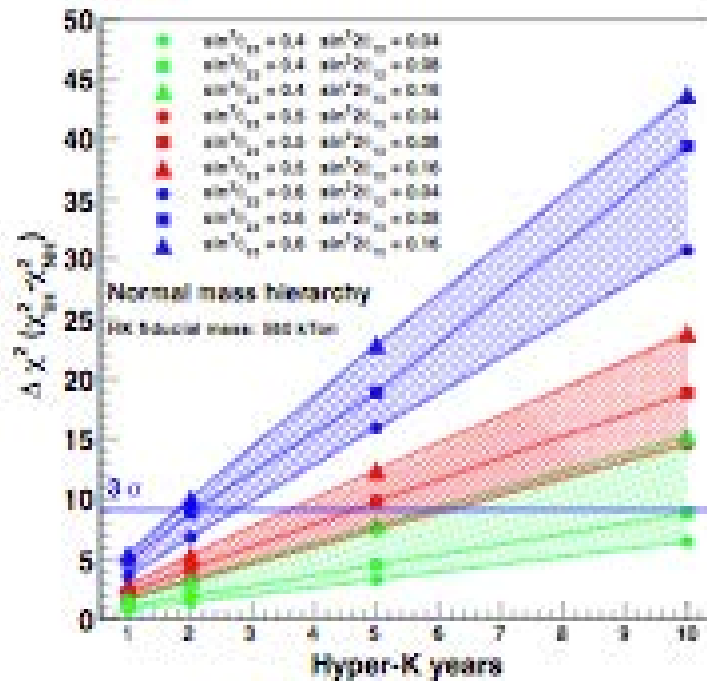


Use of atmospheric ν for mass hierarchy: SK & HK

e-like larger
for normal
hierarchy



Hyper-K; Good MH sensitivity, depend on θ_{23}



$$s_{23}^2 = 0.6$$

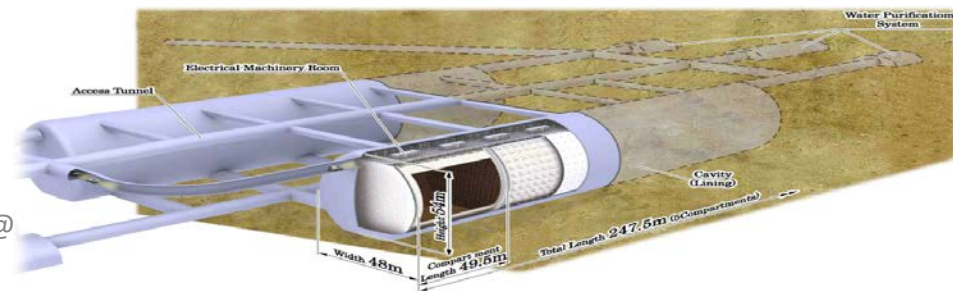
$$s_{23}^2 = 0.5$$

$$s_{23}^2 = 0.4$$

What is really good with large θ_{13} :

“All in one” approach:

- With large θ_{13} there arises an exciting possibility that CP and Mass hierarchy can be determined in situ in a single apparatus (concrete example: Hyper-K)
- With intense ν and $\bar{\nu}$ beam it can measure δ
- With gigantic atmospheric ν events it could determine the mass hierarchy
- \sim megaton scale water Cherenkov can do
- ~ 100 kt scale Liquid Ar detector can do
- It can do proton decay, interesting astrophysics ..



CP Violation: more difficult to see



Need for dedicated machine for CP

- Effect of δ (cosine and sine) has to be tiny because it is suppressed by two small numbers:

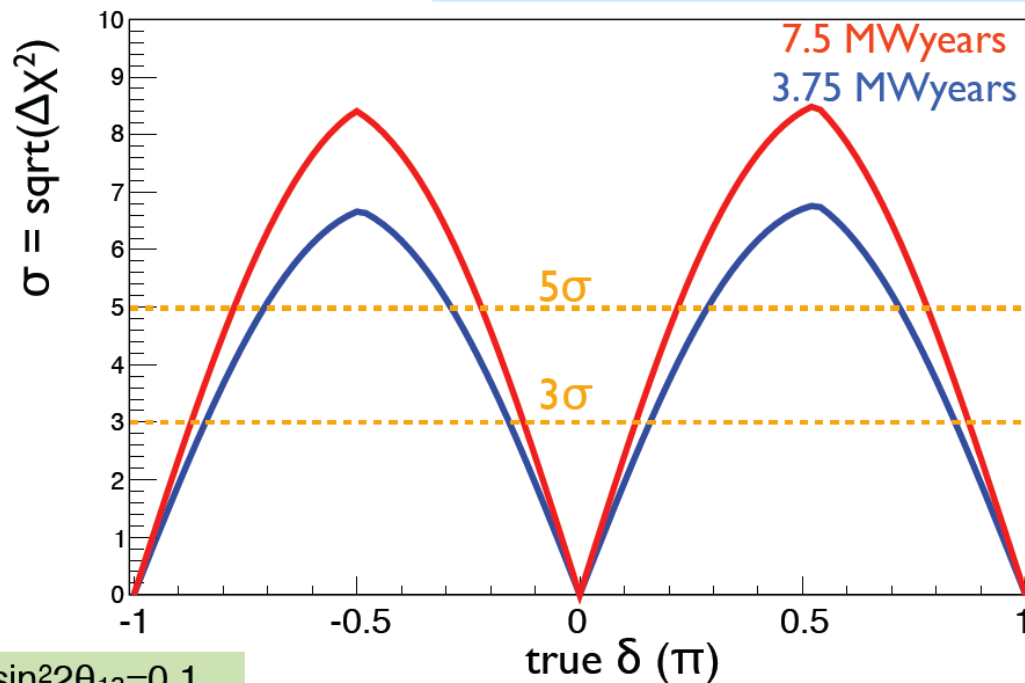
$$\Delta m_{21}^2 / \Delta m_{31}^2 \sim 0.031, J_r = c_{12}s_{12}c_{23}s_{23}s_{13} \sim 0.035$$



Hyper-K CPV sensitivity

(Exclusion of $\delta=0, \pi$)

5% systematics on signal, ν_μ BG, ν_e BG, $\nu/\bar{\nu}$



$\sin^2 2\theta_{13}=0.1$

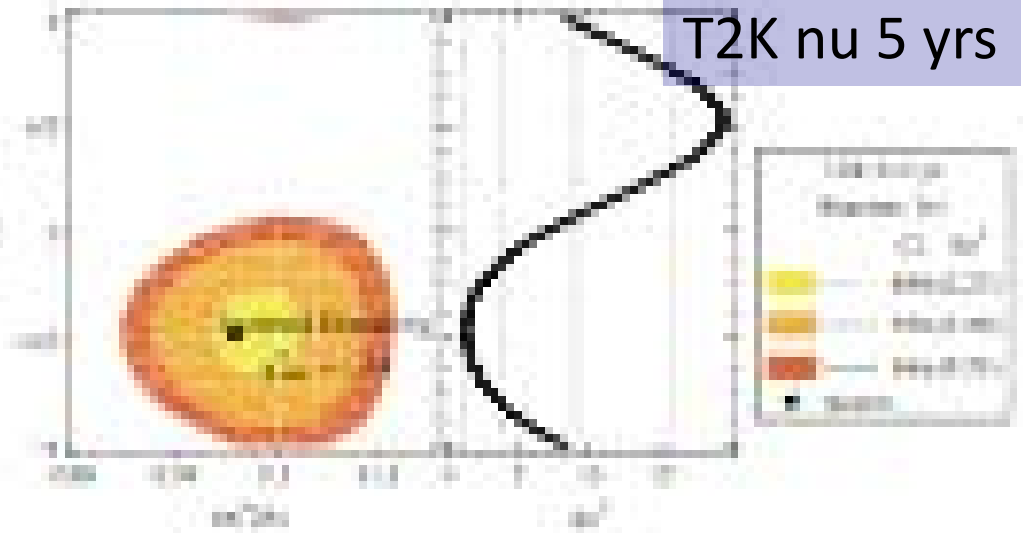
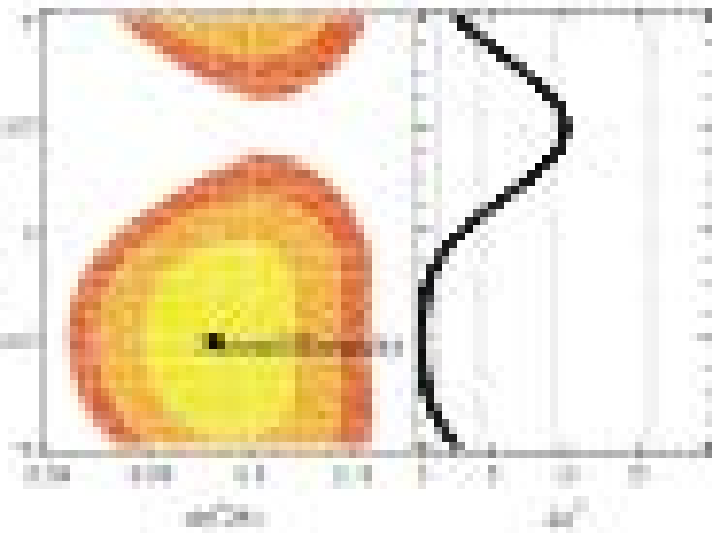
Normal hierarchy

For 74(55)% of δ , $>3(5)\sigma$ with 7.5MWyrs

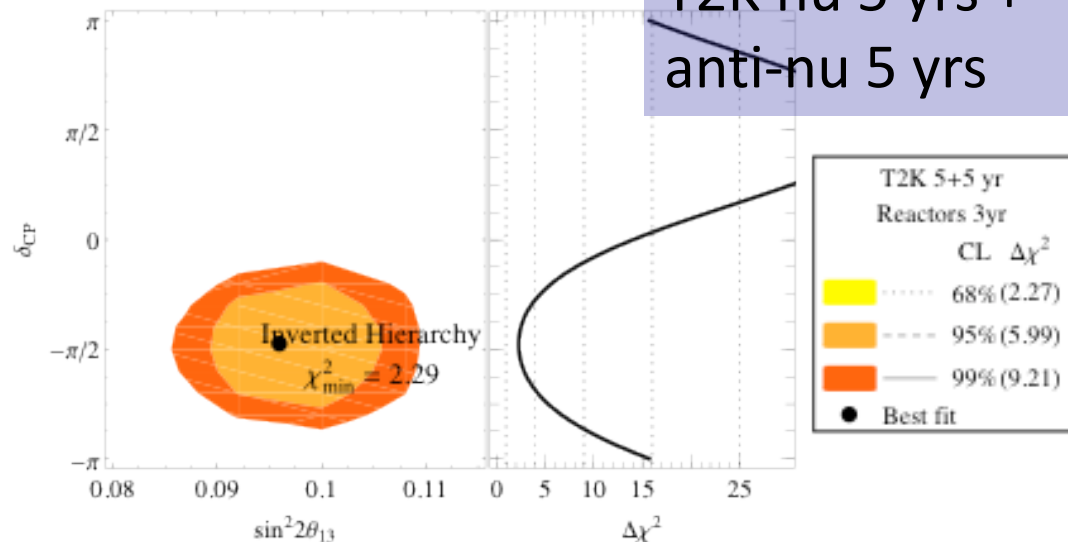
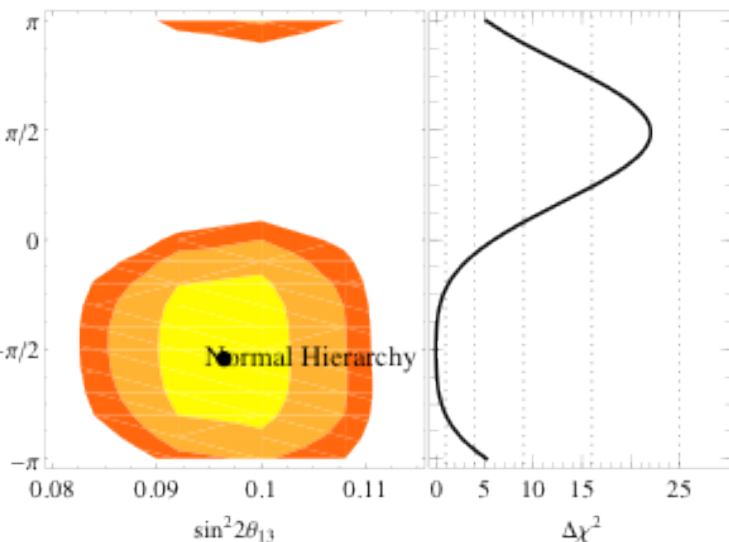
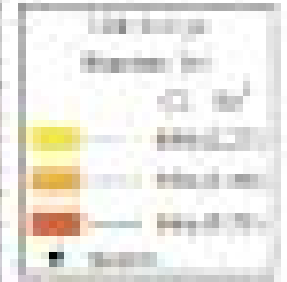
Reactor- accelerator method for CP



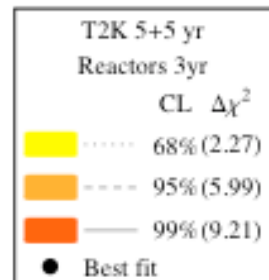
Reactor 3 years +T2K 5/10 years



T2K ν 5 yrs



T2K ν 5 yrs +
anti- ν 5 yrs

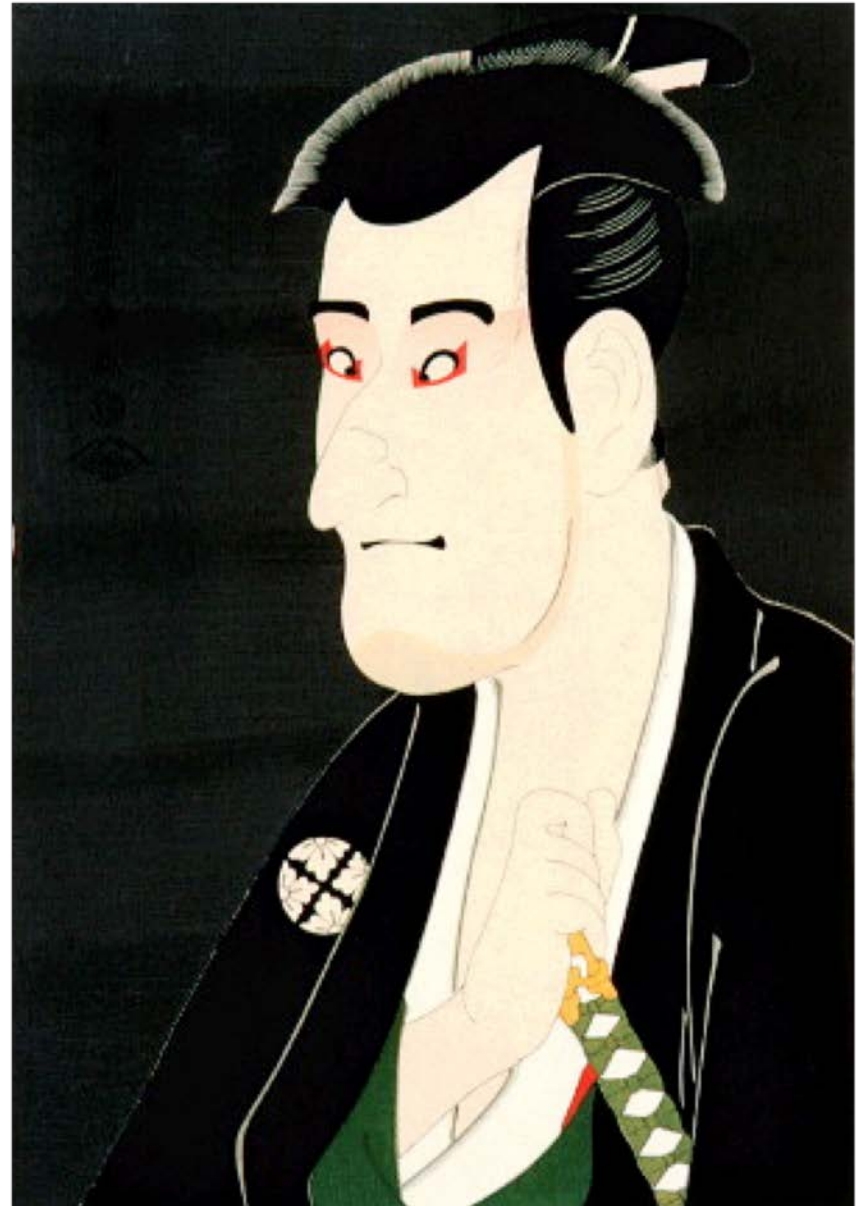


Why do you want to measure δ and mass hierarchy so much?

My prejudice !

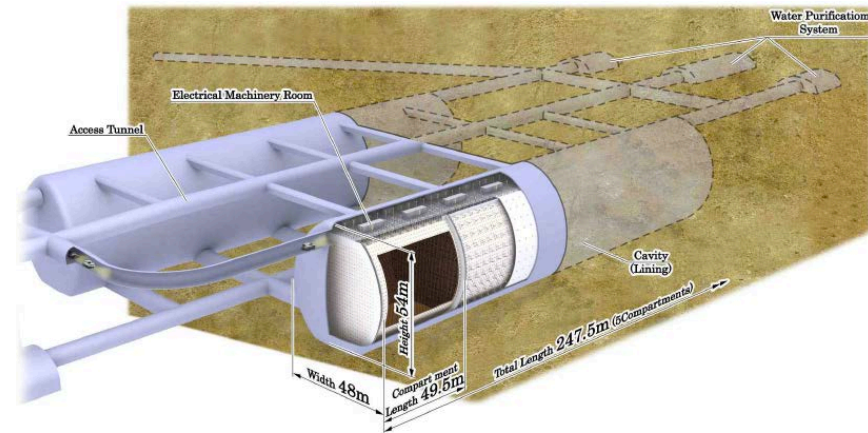
- In SM quarks and leptons are related through anomaly (short distance phenomena)
- It can be interpreted as indication of quark-lepton relation in a much deeper level than the SM scale
- Leptons and quarks have similarity (quark & charged lepton mass spectra..) and dissimilarity (small vs. large mixing angles..)
- So we want to know every feature of quark and lepton correspondence

ν mass and
large lepton
mixing:
what we are
seeing



Fermi told us about the meaning of higher dimensional operators

- $H_W = (1/v_{\text{higgs}})^2 \psi\psi\psi\psi$
- $L_{\text{mass}} = (1/M_{\text{NP}}) \langle \phi \rangle^2 \nu\nu$
- $M_{\text{NP}} \sim 10^{15} \text{ GeV} \rightarrow$ what the scale means?
- If GUT, $L_{\text{BV}} = (1/M_{\text{NP}})^2 uude$
- Simplest model for L_{mass} (seesaw) predict lepton# violation \Rightarrow leptogenesis



From where large mixing come?

$$\text{Since } U_{\text{MNS}} = U_l^\dagger U_\nu$$

- It can come from neutrino sector
- Example: Seesaw enhancement
- $L = N^c Y_\nu LH - E^c Y_\nu LH + (1/2) N^c M N$
- $m_\nu = Y_\nu^T (M)^{-1} Y_\nu$
- Natural because there is no N sector in charged leptons
- It can come from charged lepton sector
- Example: lopsided mass matrix (see next sheet)
- Interesting implications: large mixing can propagate to other lepton flavor violation (e.g., $\mu \rightarrow e \gamma$)

Lopsided lepton mass matrix (explanatory sheet)

- Consider SU(5) GUT

$$5^* = [d^c, (\nu, e)_L]$$

$$10 = [u^c, (u, d)_L, e^c]$$

- Quark mass

$$m_{LR}^{\text{down}} = 10 \cdot 5^* \langle H_1 \rangle$$

- Charged lepton mass

$$m_{LR}^{\text{lepton}} = 5^* \cdot 10 \langle H_2 \rangle$$

$$\Rightarrow m^{\text{lepton}} = (m^{\text{down}})^T$$

\Rightarrow lopsided structure; left-handed mixing of m^{lepton}
 = right-handed mixing of m^{down}

$$m^{d,l} m^{d,l+} = S m_i^2 S^+$$

$$U_{MNS} = S^{(\text{lepton})} + S^{(\nu)}$$

$$m^{\text{down}} = c \begin{bmatrix} \lambda^4 & \lambda^3 & \lambda^4 \\ x & \lambda^2 & \lambda^2 \\ y & z & 1 \end{bmatrix}, \quad m^{\text{down}} (m^{\text{down}})^\dagger = c^2 \begin{bmatrix} \lambda^6 & \lambda^5 & \lambda^4 \\ \lambda^5 & \lambda^4 & \lambda^2 \\ \lambda^4 & \lambda^2 & 1 \end{bmatrix}$$

$$m^{\text{lepton}} (m^{\text{lepton}})^\dagger = c'^2 \begin{bmatrix} x^2 + y^2 & yz + \lambda^2 x & y + \lambda^2 x \\ yz + \lambda^2 x & z^2 & z \\ y + \lambda^2 x & z & 1 \end{bmatrix}$$

$\lambda=0.2, x, y, z = O(1) \Rightarrow$ Large lepton mixing arises from quark mass

Quark-lepton complementarity

$$\theta_{\text{Cabibbo}} + \theta_{\text{solar}} = \pi / 4$$



QLC embedded into GUTs

$$U_{\text{MNS}} = U_{\text{lepton}}^\dagger U_n, \quad V_{\text{CKM}} = V_{\text{up}}^\dagger V_{\text{down}}$$

- Neutrino-induced bimaximal

$$\begin{array}{ll}
 U_\nu = U_{\text{bimaximal}} & \text{Seesaw enhance?} \\
 U_{\text{lepton}} = V_{\text{CKM}} & \text{<= GUTs =>}
 \end{array}
 \quad
 \begin{array}{l}
 V_{\text{up}} = I \\
 V_{\text{down}} = V_{\text{CKM}}
 \end{array}$$

- Lepton-induced bimaximal

$$\begin{array}{ll}
 U_\nu = V_{\text{CKM}}^\dagger & \text{<= GUTs =>} \\
 U_{\text{lepton}} = U_{\text{bimaximal}} & \text{<= Lopsided}
 \end{array}
 \quad
 \begin{array}{l}
 V_{\text{up}} = V_{\text{CKM}}^\dagger \\
 V_{\text{down}} = I
 \end{array}$$


Large θ_{13} in QLC context

“bimaximal minus CKM mixing.”

HM-A.Smirnov 04

Bi-maximal from neutrinos


$$U_\nu = R_{23}^m R_{12}^m, \quad U_l = V^{\text{CKM}},$$
$$U_{\text{MNS}} = V^{\text{CKM}\dagger} \Gamma_\delta R_{23}^m R_{12}^m$$
$$= R_{12}^{\text{CKM}\dagger} R_{13}^{\text{CKM}\dagger} R_{23}^{\text{CKM}\dagger} \Gamma_\delta R_{23}^m R_{12}^m,$$


$$\sin \theta_{13} \simeq \frac{1}{\sqrt{2}} \sin \theta_C$$
$$\sin^2 \theta_{13} = 0.026 \pm 0.008$$

Bi-maximal from charge leptons

$$V_\nu = V^{\text{CKM}\dagger}, \quad V_l = R_{12}^{m\dagger} R_{23}^{m\dagger}.$$

$$U_{\text{MNS}} = R_{23}^m \Gamma_\delta R_{12}^m V^{\text{CKM}\dagger}$$


$$\sin \theta_{13} \approx -\sin \theta_{\text{sun}} |V_{cb}|,$$
$$\sin^2 2\theta_{13} = 1.9 \times 10^{-3}$$

IN QLC context, large θ_{13} implies that large mixing comes from neutrino sector!

Conclusion

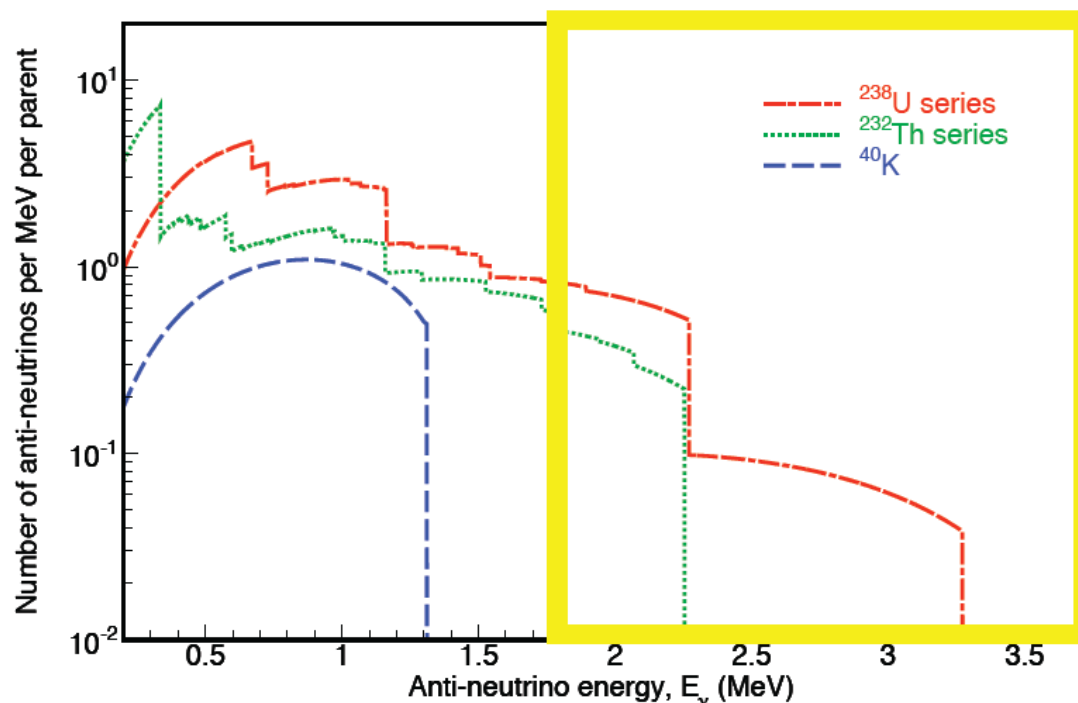
- We know know all the three lepton mixing angles and after this our next goal is well defined: δ and ν mass hierarchy
- We need "guaranteeing machine" to meet the goal but what is good with large θ_{13} is "all in one" approach becomes possible
- Large θ_{13} triggers "hundred flowers" situation for method for determining MH
- Absolute n mass, Majorana phase ..
- How we can make solid step-by-step progress in understanding physics?



Supplementa ry slides

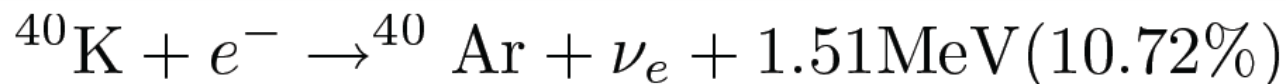
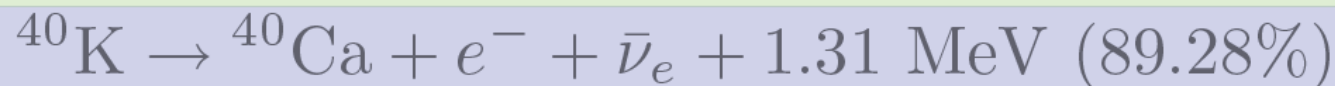
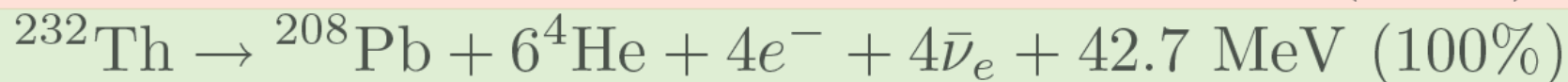
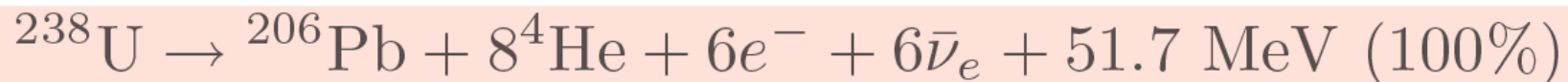
地球内部起源ニュートリノ

反ニュートリノフラックス



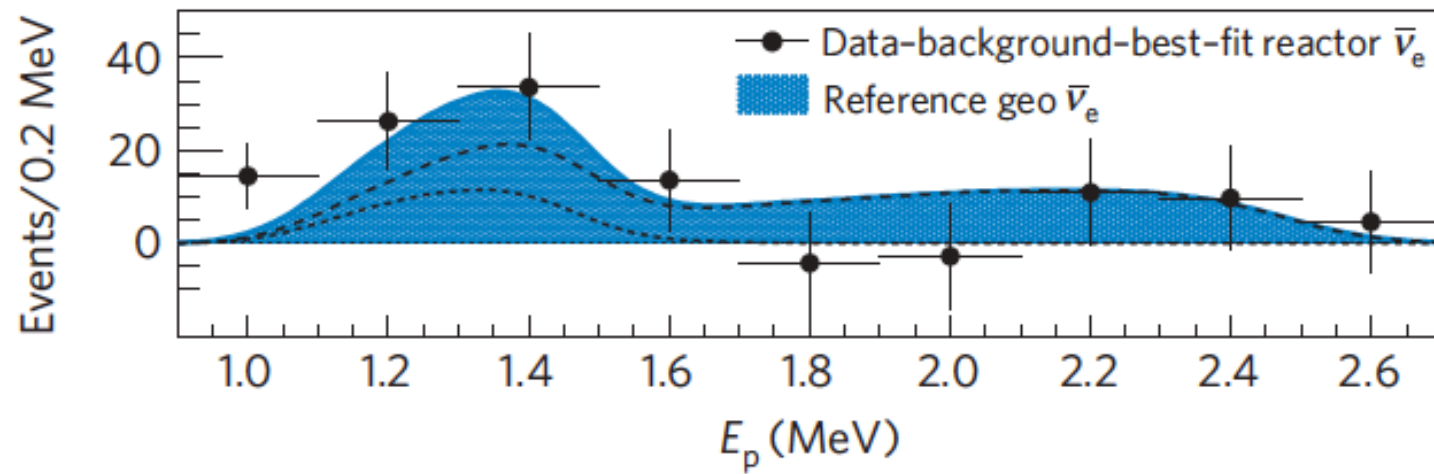
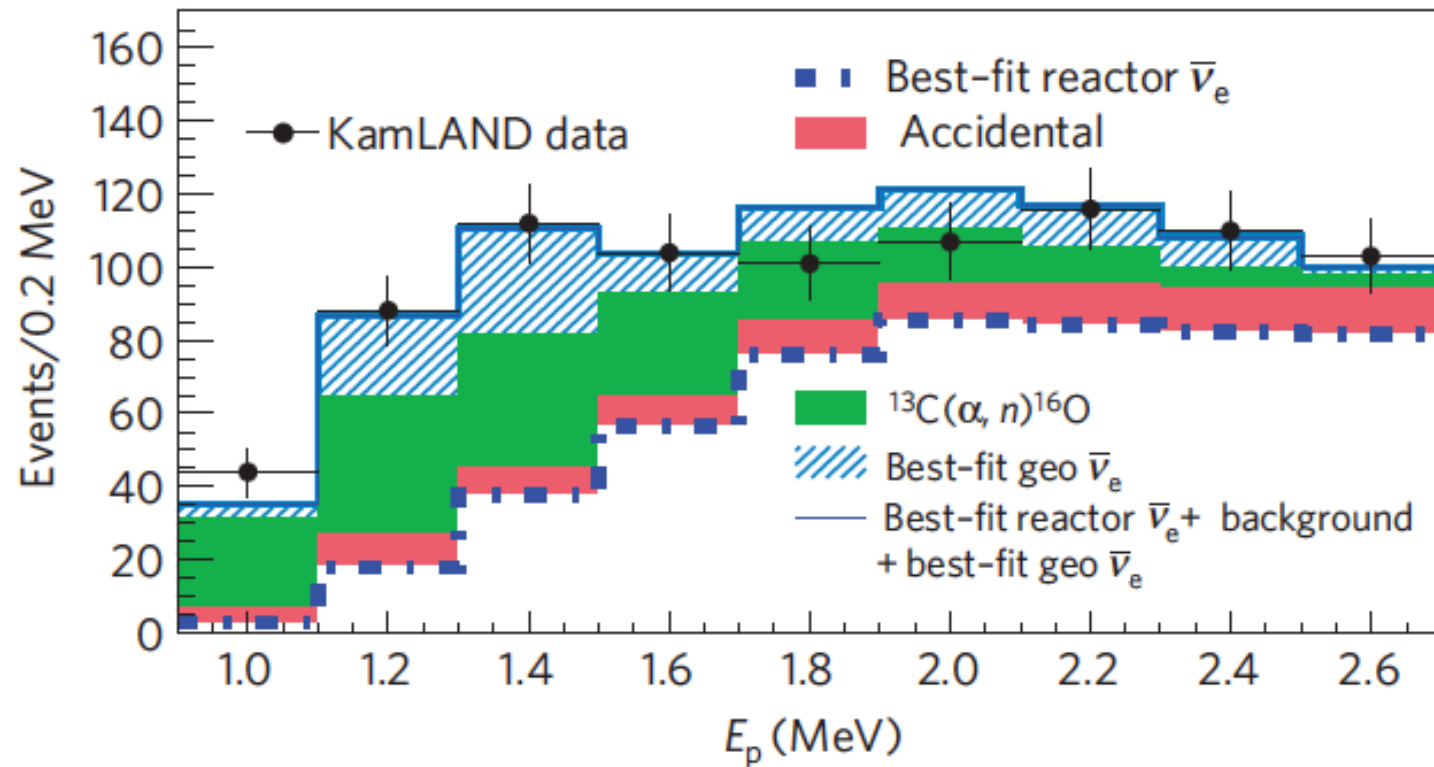
KamLAND
での観測領域

β崩壊



KamLAND result 2011

NATURE GEOSCIENCE DOI:10.1038/NNGEO1205



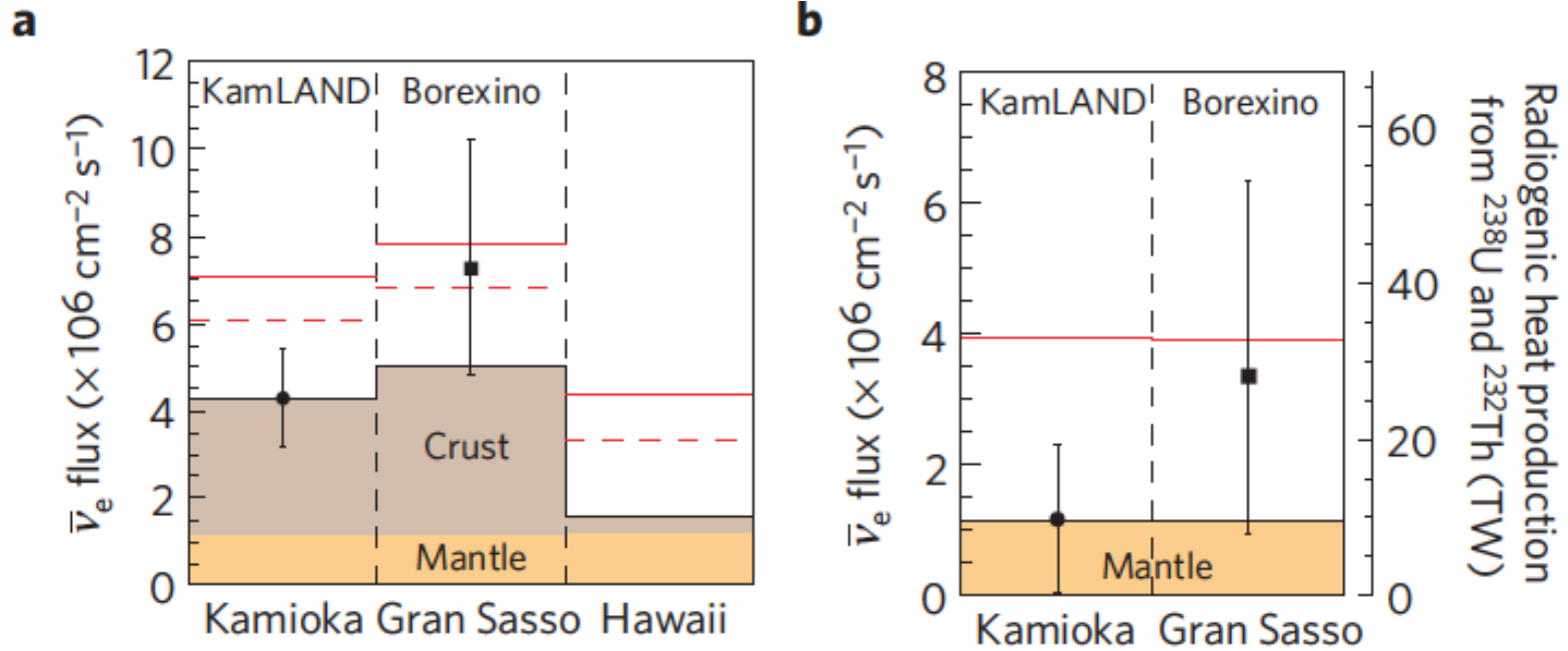
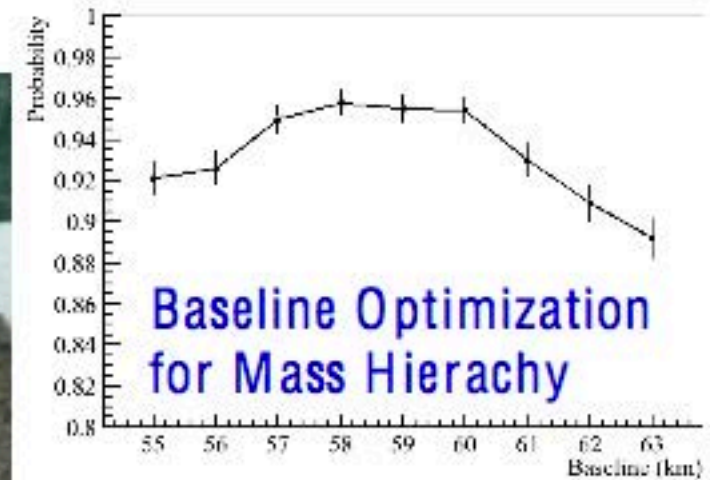
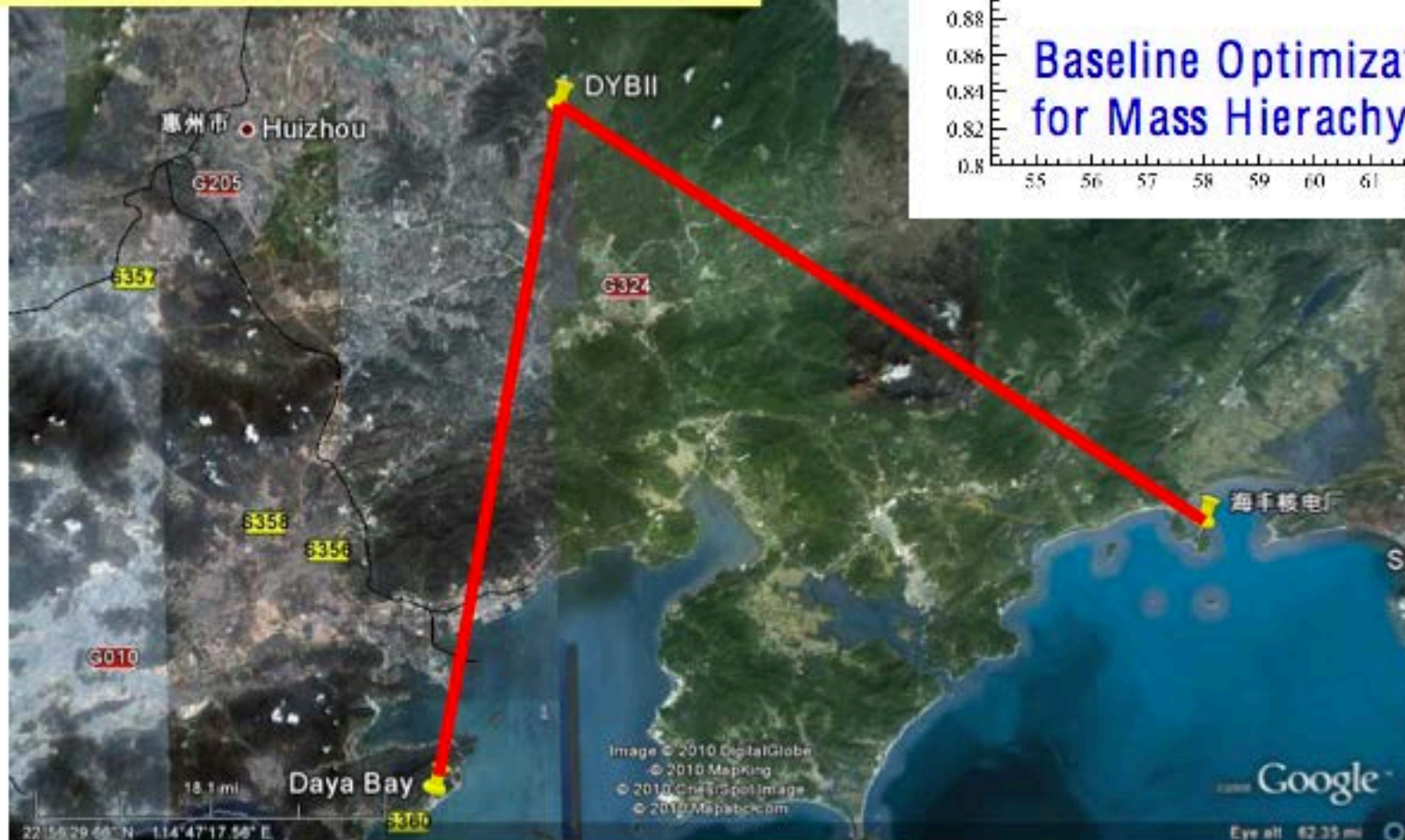


Figure 4 | Measured geoneutrino flux and models. a, Measured geoneutrino flux at Kamioka and Gran Sasso, and expected fluxes at these sites and Hawaii⁴. The solid and dashed red lines represent, respectively, the fluxes for a fully radiogenic model assuming the homogeneous and sunken-layer hypotheses. **b,** Measured geoneutrino flux after subtracting the estimated crustal contribution. No modelling uncertainties are shown. The right axis shows the corresponding radiogenic heat production assuming a homogeneous mantle. The solid red line indicates the fully radiogenic model where the contributions from the crust (7.0 TW) and the other isotopes^{6,24} (4.3 TW) are subtracted from the total heat flow⁷ (44.2 TW). Error bars, see text.

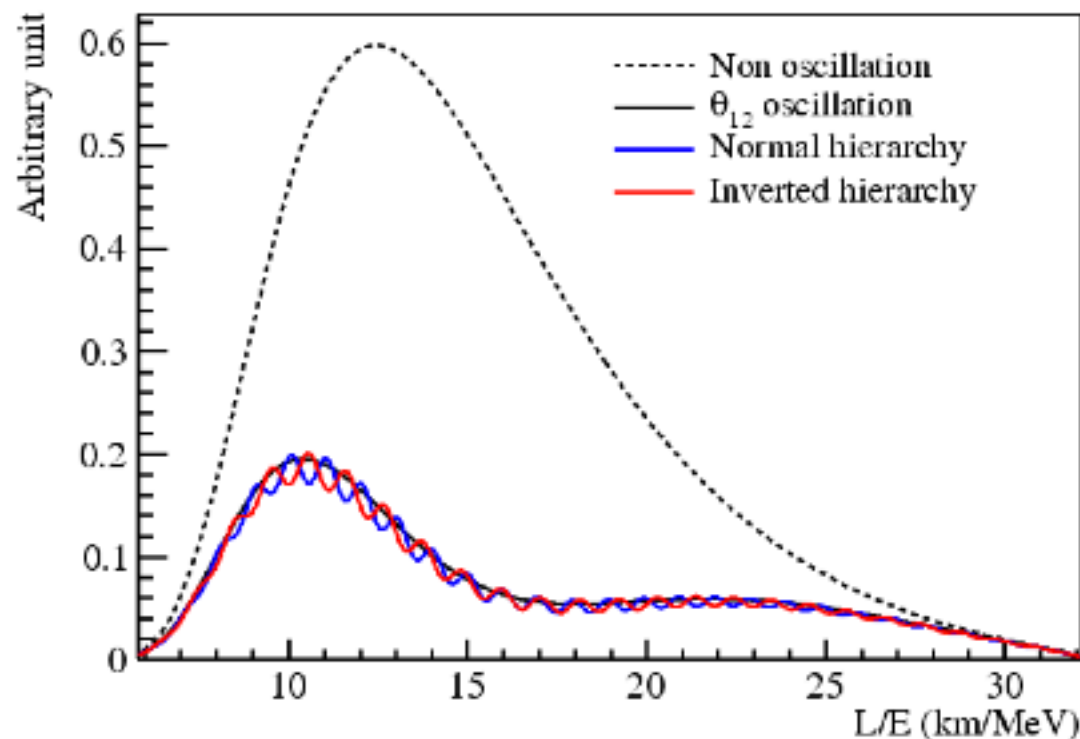
Daya Bay II Site Investigation

J.Cao@ NuTURN12

~60km to Daya Bay and to Haifeng
Thermal Power (17.4 GW + 17.4 GW)
Overburden > 1000 m.w.e



Reactor Exp. to determine MH



$$P_{ee}(L/E) = 1 - P_{21} - P_{31} - P_{32}$$

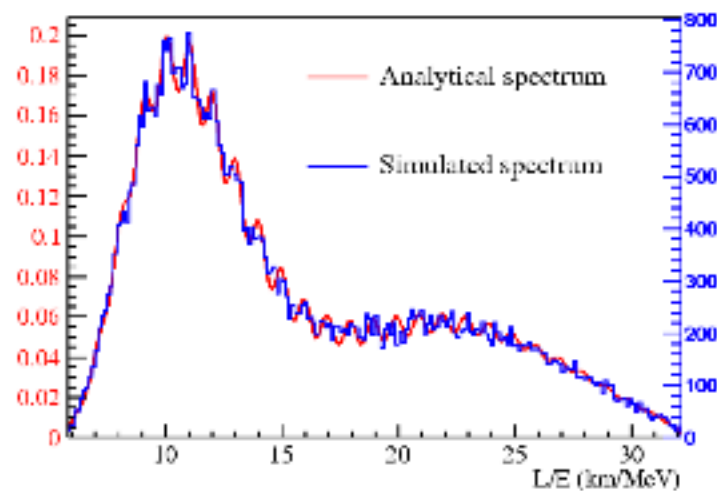
$$P_{21} = \cos^4(\theta_{13}) \sin^2(2\theta_{12}) \sin^2(\Delta_{21})$$

$$P_{31} = \cos^2(\theta_{12}) \sin^2(2\theta_{13}) \sin^2(\Delta_{31})$$

$$P_{32} = \sin^2(\theta_{12}) \sin^2(2\theta_{13}) \sin^2(\Delta_{32})$$

S.T. Petcov et al., PLB533(2002)94
 S.Choubey et al., PRD68(2003)113006
 J. Learned et al., hep-ex/0612022

L. Zhan, Y. Wang, J. Cao, L. Wen,
 PRD78:111103, 2008
 PRD79:073007, 2009



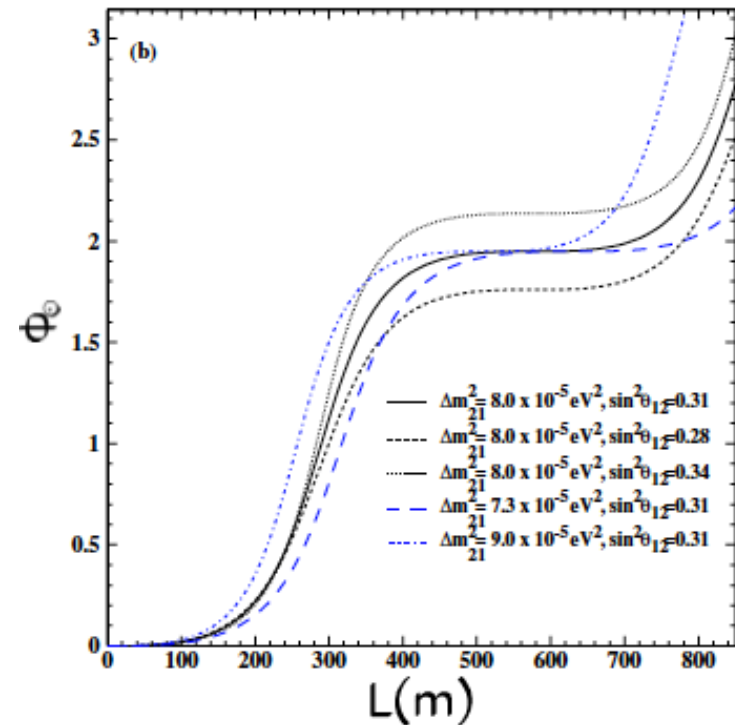
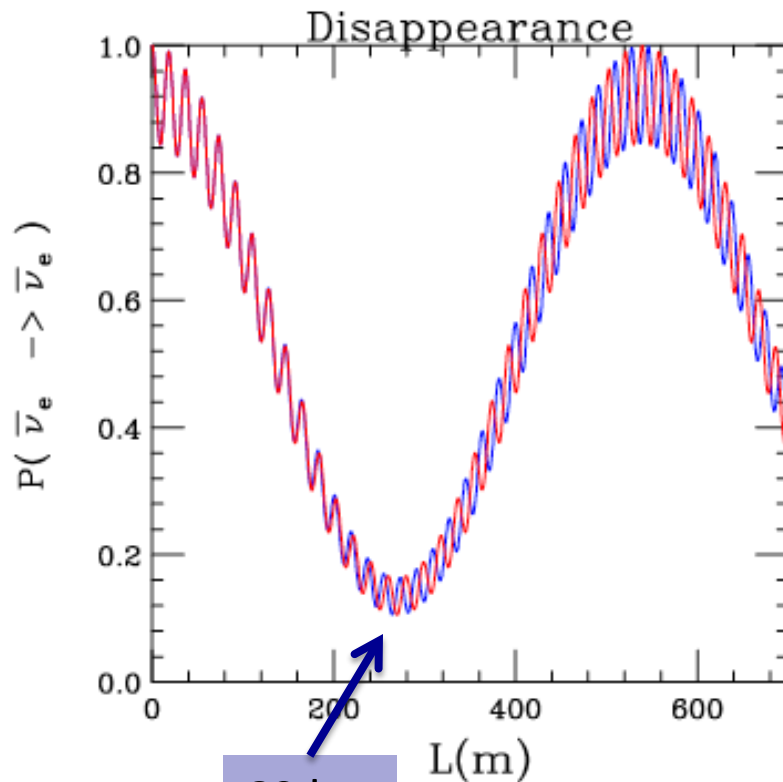
Phase advancement vs. retardation

$$P(\nu_e \rightarrow \nu_e) = 1 - \frac{1}{2} \sin^2 2\theta_{13} [1 - \sqrt{1 - \sin^2 2\theta_{12} \sin^2 \Delta_{21}} \times \cos(2\Delta_{ee} \pm \phi_\odot)] - P_\odot,$$

Hierarchy sign

$$\Delta m_{ee}^2 \equiv \cos^2 \theta_{12} |\Delta m_{31}^2| + \sin^2 \theta_{12} |\Delta m_{32}^2|.$$

$$\sin \phi_\odot = \frac{c_{12}^2 \sin(2s_{12}^2 \Delta_{21}) - s_{12}^2 \sin(2c_{12}^2 \Delta_{21})}{\sqrt{1 - \sin^2 2\theta_{12} \sin^2 \Delta_{21}}}$$

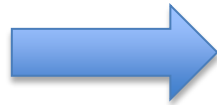


Mass hierarchy confusion

$$P(\nu_e \rightarrow \nu_e) = 1 - \frac{1}{2} \sin^2 2\theta_{13} [1 - \sqrt{1 - \sin^2 2\theta_{12} \sin^2 \Delta_{21}} \times \cos(2\Delta_{ee} \pm \phi_\odot)] - P_\odot, \quad ($$

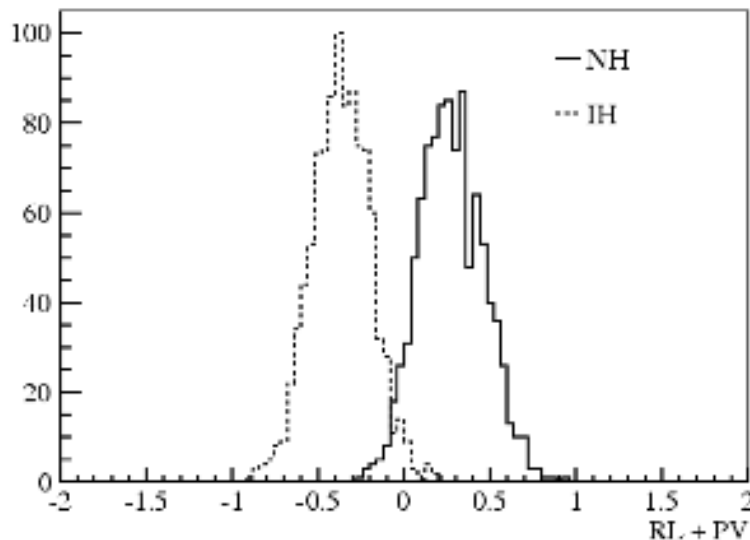
If accuracy of Δm^2_{atm} is not good enough

$$2\Delta_{ee}|_{\text{NH}} + \phi_\odot = 2\Delta_{ee}|_{\text{IH}} - \phi_\odot$$



complete confusion of mass hierarchy

One need to watch over many wiggles → Fourier transform

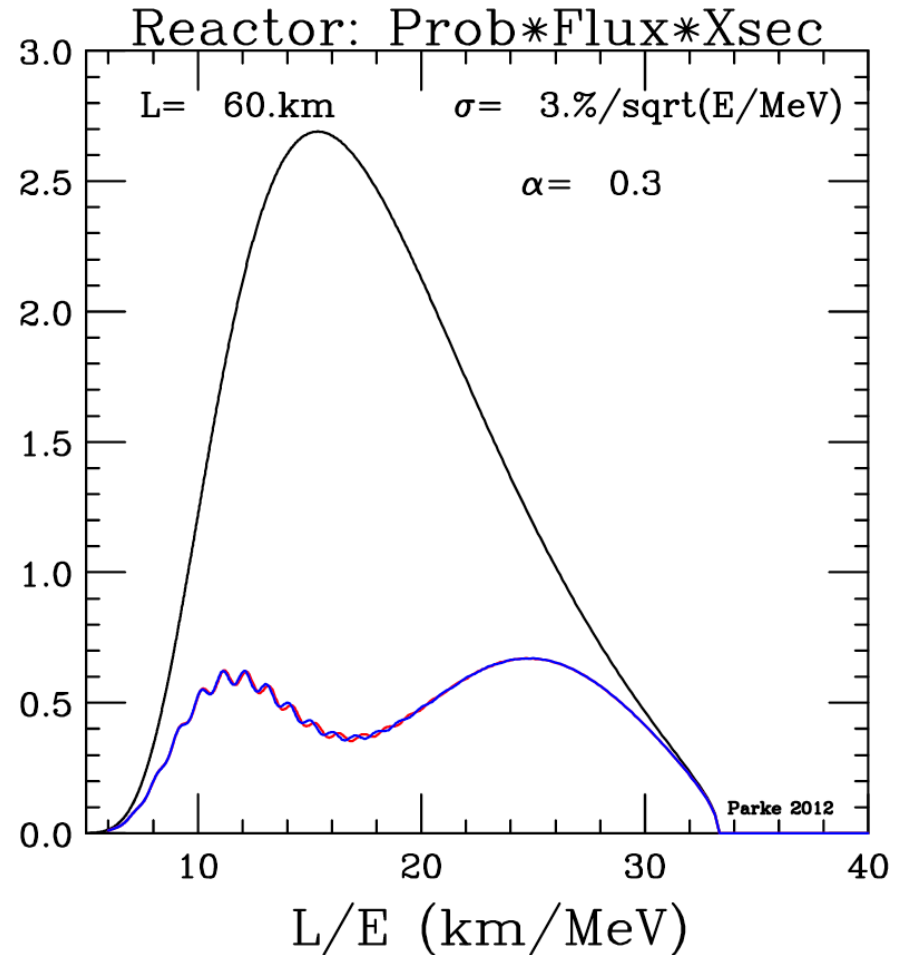
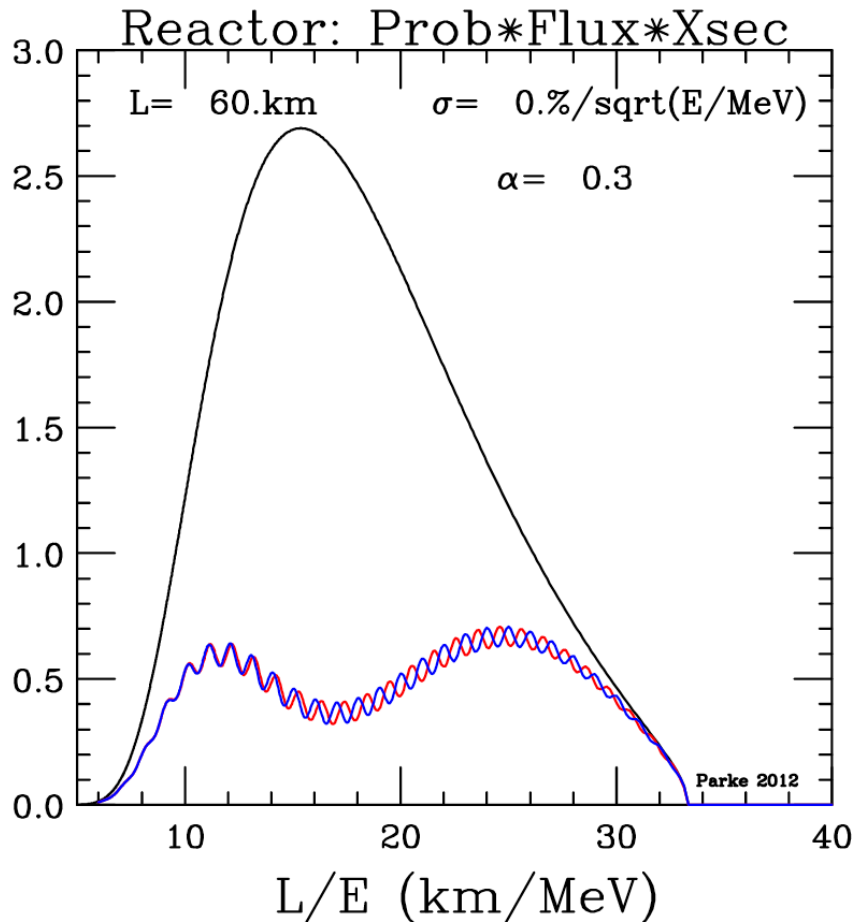


Energy resolution	3%/sqrt(E)
Baseline	58 km
Thermal Power	35 GW

50k events = 20k tons X 3 years

50k events: 96% probability

Then, energy resolution (absolute & relative) is the issue



See X.~Qian, et al. arXiv:1208.1551 [hep-ex] for experimental requirements

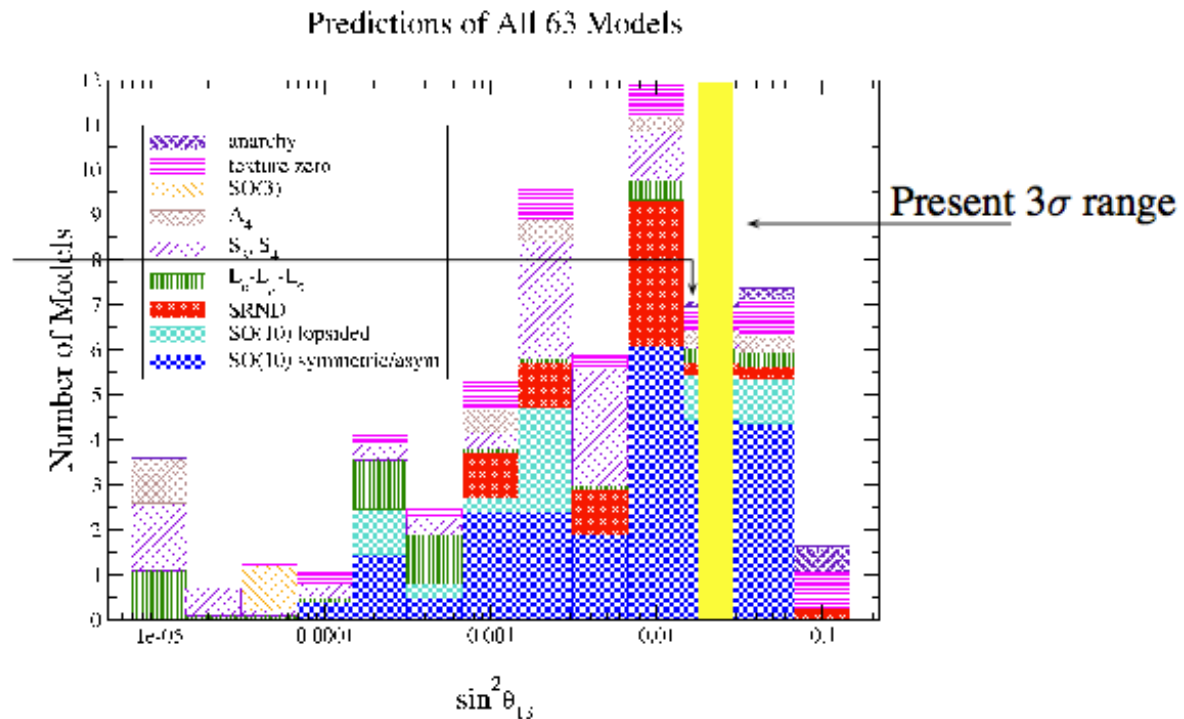
Implication of large θ_{13} to the theory is NOT that clear..

Concha@ICHEP2012

Implications

- Survey of 63 ν mass models in 2006 (Albright, M-C Chen, hep-ph/0608136)

Only 7
got it right !



S. Antusch