

MSSM Neutralino signature

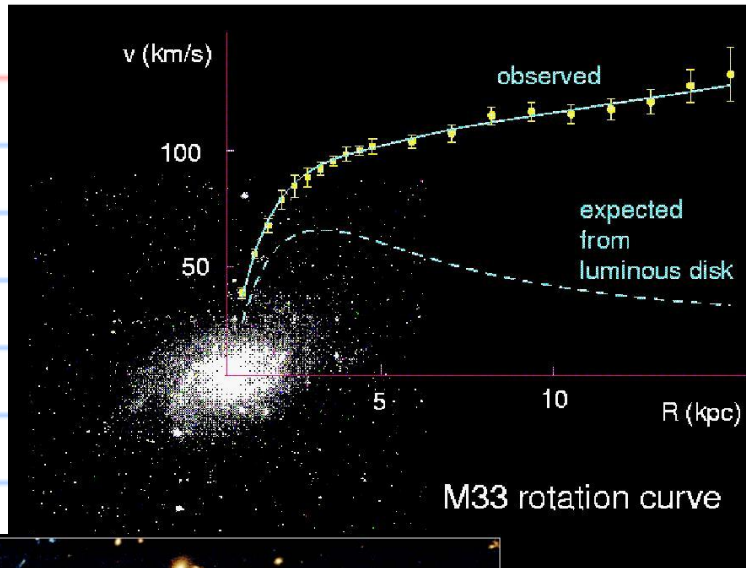
IPMU

1 November 2013

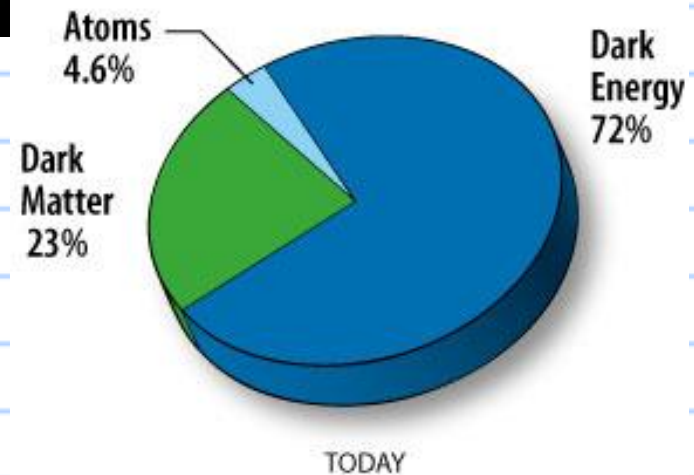
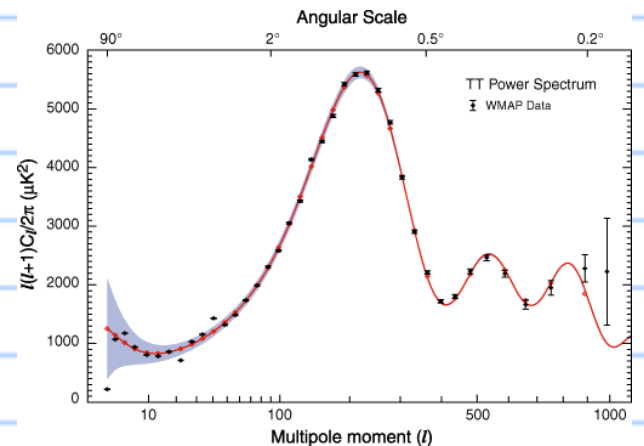
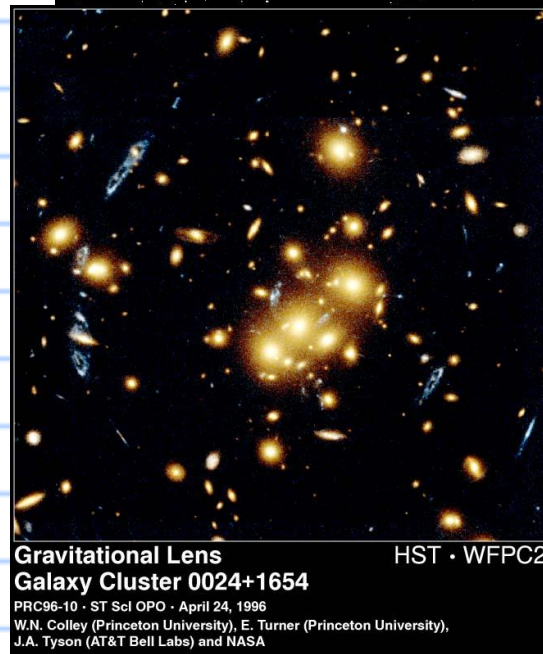
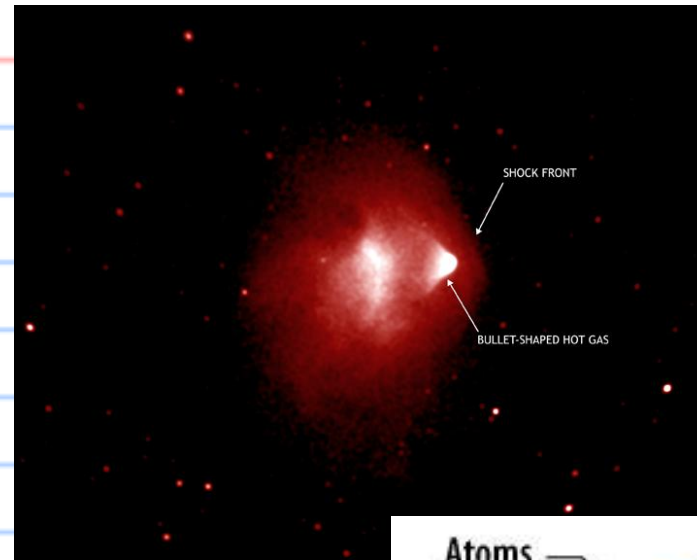
Y-L Sming Tsai (蔡岳霖・駟名)

Based on 1306.1567
Phys.Rev. D88 (2013) 055012
in Collaboration with BayesFITS group

Dark Matter evidence

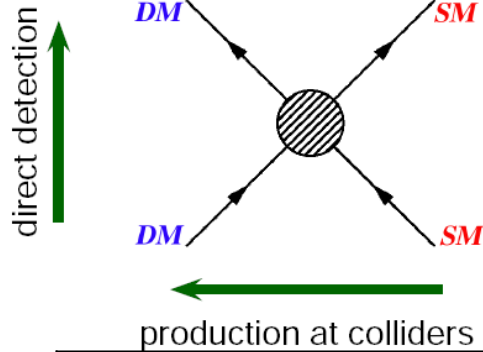


M33 rotation curve



It will be difficult to explain the universe without dark matter assumption.

thermal freeze-out (early Univ.)
indirect detection (now)



Dark Matter search



Signals at	Experiments	DM Hints
Coillders	LHC, LEP, Tevatron, ...	None
Direct detection	XENON100, CDMS, LUX ...	DAMA, CoGeNT, CRESST at low DM mass region.
Cosmic rays <ol style="list-style-type: none"> 1. Positrons 2. antiprotons 3. neutrinos 	<ol style="list-style-type: none"> 1. PAMELA, Fermi-LAT, AMS02... 2. PAMELA... 3. IceCube... 	<ol style="list-style-type: none"> 1. High energy positron excess 2. None 3. PeV neutrinos
Gamma rays	Fermi-LAT, ...	FERMI bubbles, Fermi Gamma ray line at 130 GeV...
Radio	WMAP, Planck	WMAP (Planck) haze

Outline

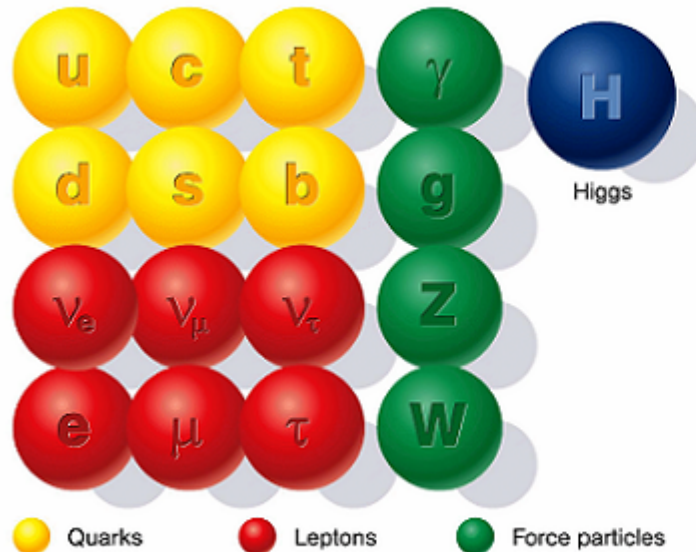
1. MSSM neutralino dark matter
2. SUSY Dark Matter search at the LHC
3. Impact of the XENON100 (2012) result
4. Neutralino indirect detection
5. Summary



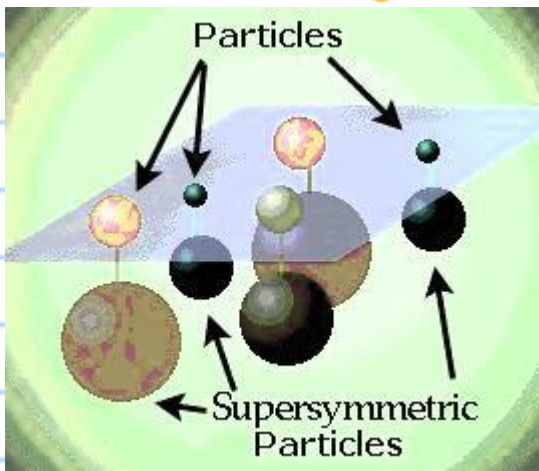
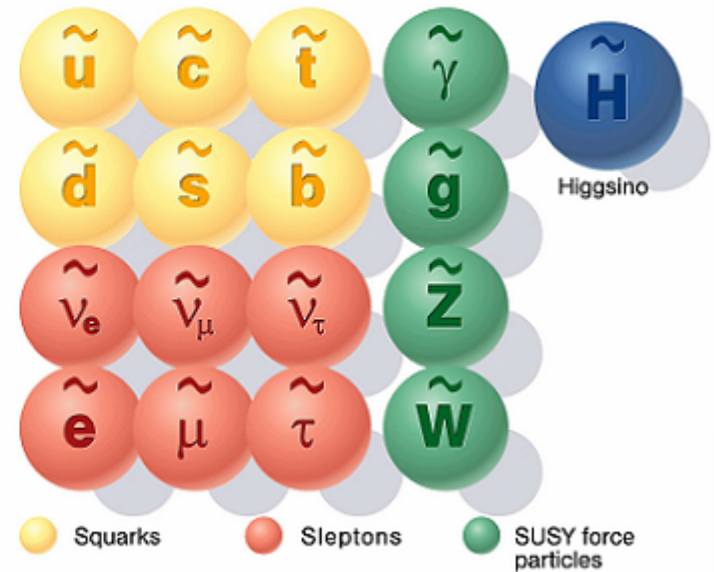
Neutralino Dark Matter

Supersymmetry

Standard particles



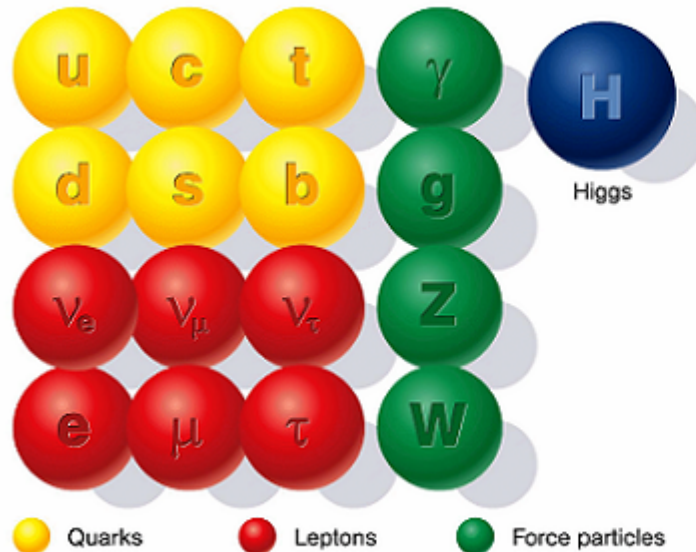
SUSY particles



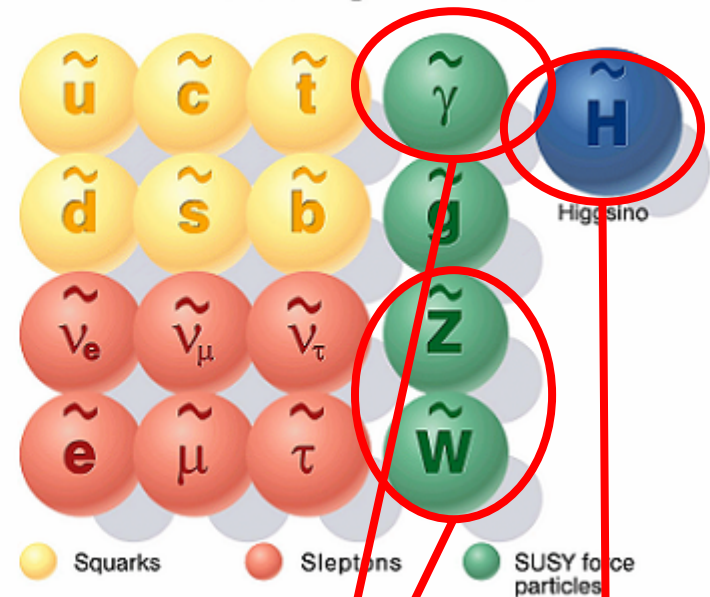
Supersymmetry

No DM candidate !!

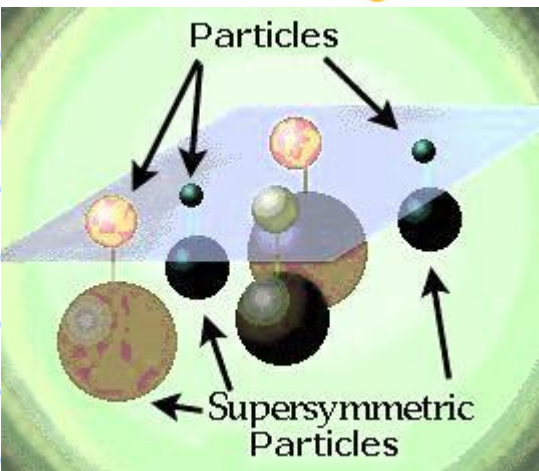
Standard particles



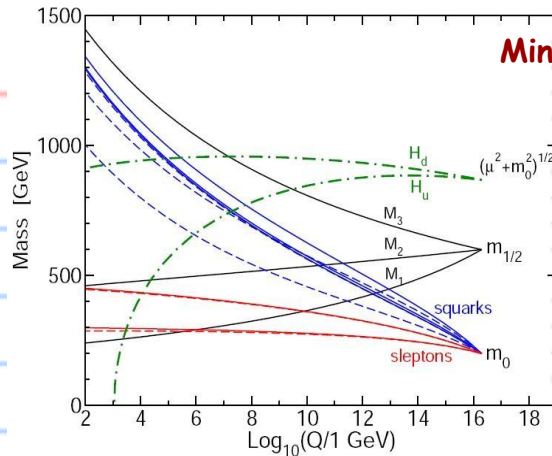
SUSY particles



"neutralino" χ : lightest mass state of neutral gauginos and higgsinos
stable WIMP: excellent DM candidate



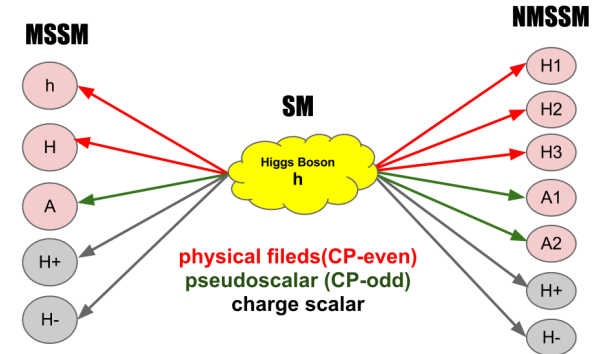
Some popular supersymmetric models:



Minimal and Next Minimal Supersymmetric Standard Model

G. L. Kane, C. F. Kolda, L. Roszkowski and J. D. Wells, Phys. Rev. D 49 (1994) 6173

Higgs sector



Models	Input parameters	Boundary conditions
CMSSM	$m_0, m_{1/2}, \tan\beta, A_0, \text{sgn}(\mu)$	GUT ($Q/\text{GeV} \sim 2 \times 10^{16}$)
CMSSM +non-universal Higgs mass	CMSSM+ (m_{Hu}, m_{Hd})	GUT ($Q/\text{GeV} \sim 2 \times 10^{16}$)
MSSM	all trilinear couplings, masses, and Higgs parameters	SUSY ($Q \sim \text{stop mass}$)
CNMSSM	CMSSM+ Singlet higgs mass	GUT ($Q/\text{GeV} \sim 2 \times 10^{16}$)
NMSSM	all trilinear couplings, masses, and Higgs parameters	SUSY ($Q \sim \text{stop mass}$)

Choice of parameters (p9MSSM)

All ranges in TeV

Parameter	Range
gluino mass	$0.7 < M_3 < 8$
wino mass	$0.01 < M_2 < 4$
bino mass	$M_1 = 0.5M_2$
stop trilinear coupl.	$-7 < A_t < 7$
τ trilinear coupl.	$-7 < A_\tau < 7$
sbottom trilinear coupl.	$A_b = -0.5$
pseudoscalar mass	$0.2 < m_A < 4$
μ parameter	$0.01 < \mu < 4$
3rd gen. soft squark mass	$0.3 < m_{\tilde{Q}_3} < 4$
3rd gen. soft slepton mass	$0.1 < m_{\tilde{L}_3} < 2$
1st/2nd gen. soft slepton mass	$m_{\tilde{L}_{1,2}} = M_1 + 50 \text{ GeV}$
1st/2nd gen. soft squark mass	$m_{\tilde{Q}_{1,2}} = 2.5$
ratio of Higgs doublet VEVs	$3 < \tan \beta < 62$
Nuisance parameter	Central value, error
Bottom mass $m_b(m_b)^{\overline{MS}}$ (GeV)	(4.18, 0.03)
Top mass at pole M_t (GeV)	(173.5, 1.0)

Bino DM candidate
(Fit relic density and ID)

Choice of parameters (p9MSSM)

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Bino DM candidate
(Fit relic density and ID)

SMUON light
(Fit g-2)

1st 2nd gen squarks heavy
(LHC disfavored)

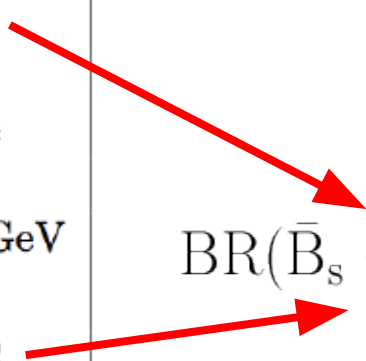
Choice of parameters (p9MSSM)

All ranges in TeV

$M_2, M_3, m_{\tilde{Q}_3}, m_{\tilde{L}_3}, A_t, A_\tau, m_A, \mu, \tan \beta$

1.8×10^6 points
collected
subject to constraints

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$$\text{BR}(\bar{B}_s \rightarrow \mu^+ \mu^-) \propto \frac{\tan^6 \beta}{m_A^4}$$

Constraints in Likelihood

Measurement	Mean or Range	Error: (Exp., Th.)	Distribution
CMS α_T 11.7/fb , $\sqrt{s} = 8$ TeV	See text	See text	Poisson
m_h by CMS	125.8 GeV	0.6 GeV, 3 GeV	Gaussian
$\Omega_\chi h^2$	0.1199	0.0027, 10%	Gaussian
$\text{BR}(\bar{B} \rightarrow X_s \gamma) \times 10^4$	3.43	0.22, 0.21	Gaussian
$\text{BR}(B_u \rightarrow \tau \nu) \times 10^4$	1.66	0.33, 0.38	Gaussian
ΔM_{B_s}	17.719 ps ⁻¹	0.043 ps ⁻¹ , 2.400 ps ⁻¹	Gaussian
$\sin^2 \theta_{\text{eff}}$	0.23146	0.00012, 0.00015	Gaussian
M_W	80.399 GeV	0.023 GeV, 0.015 GeV	Gaussian
$\text{BR}(B_s \rightarrow \mu^+ \mu^-) \times 10^9$	3.2	+1.5 - 1.2, 10%	Gaussian
$m_b(m_b)^{\overline{MS}}$	4.18 GeV	0.03 GeV, 0	Gaussian
M_t	173.5 GeV	1.0 GeV, 0	Gaussian
$\delta(g-2)_\mu^{\text{SUSY}} \times 10^{10}$	28.7	8.0, 1.0	Gaussian
XENON100 (2012)	See text	See text	Poisson
CMS $3l + E_T^{\text{miss}}$ 9.2/fb, $\sqrt{s} = 8$ TeV	See text	See text	Poisson

PLUS LEP CONSTRAINTS:

NEW!!! XENON100 (2012)
included.

$$\begin{aligned}
 m_\chi &> 46 \text{ GeV} , \\
 m_{\tilde{e}} &> 107 \text{ GeV} , \\
 m_{\chi_1^\pm} &> 94 \text{ GeV if } m_{\chi_1^\pm} - m_\chi > 3 \text{ GeV and } \tan \beta < 40 , \\
 m_{\tilde{\mu}} &> 94 \text{ GeV if } m_{\tilde{\mu}} - m_\chi > 10 \text{ GeV and } \tan \beta < 40 , \\
 m_{\tilde{\tau}} &> 81.9 \text{ GeV if } m_{\tilde{\tau}_1} - m_\chi > 15 \text{ GeV} , \\
 m_{\tilde{b}_1} &> 89 \text{ GeV if } m_{\tilde{b}_1} - m_\chi > 8 \text{ GeV} , \\
 m_{\tilde{t}_1} &> 95.7 \text{ GeV if } m_{\tilde{t}_1} - m_\chi > 10 \text{ GeV} .
 \end{aligned}$$

Constraints in Likelihood

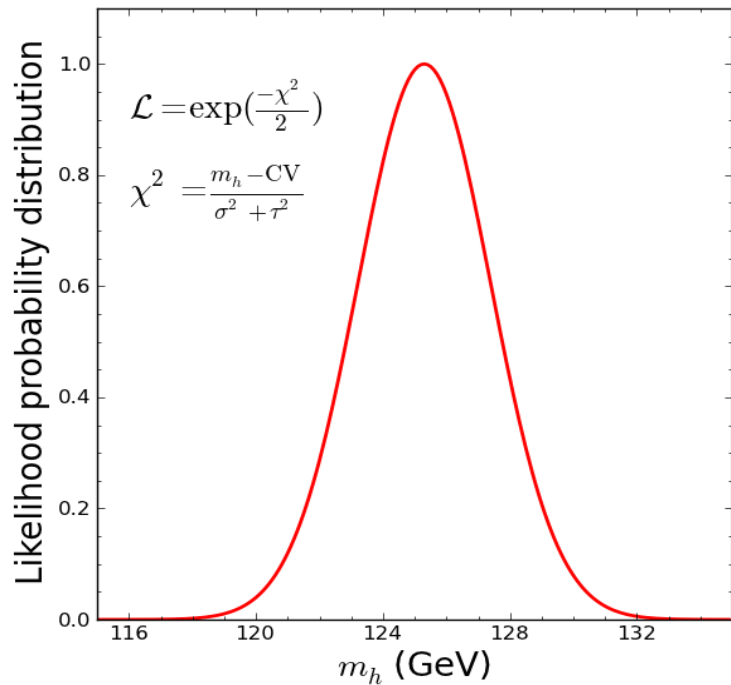
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Basic set

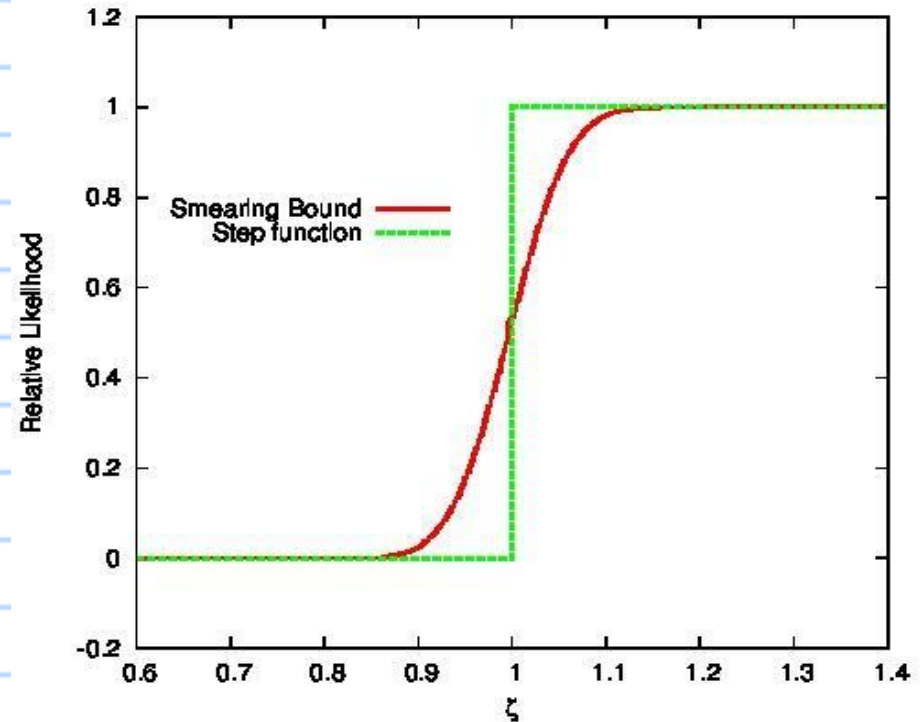
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 \end{aligned}$$

Likelihood types



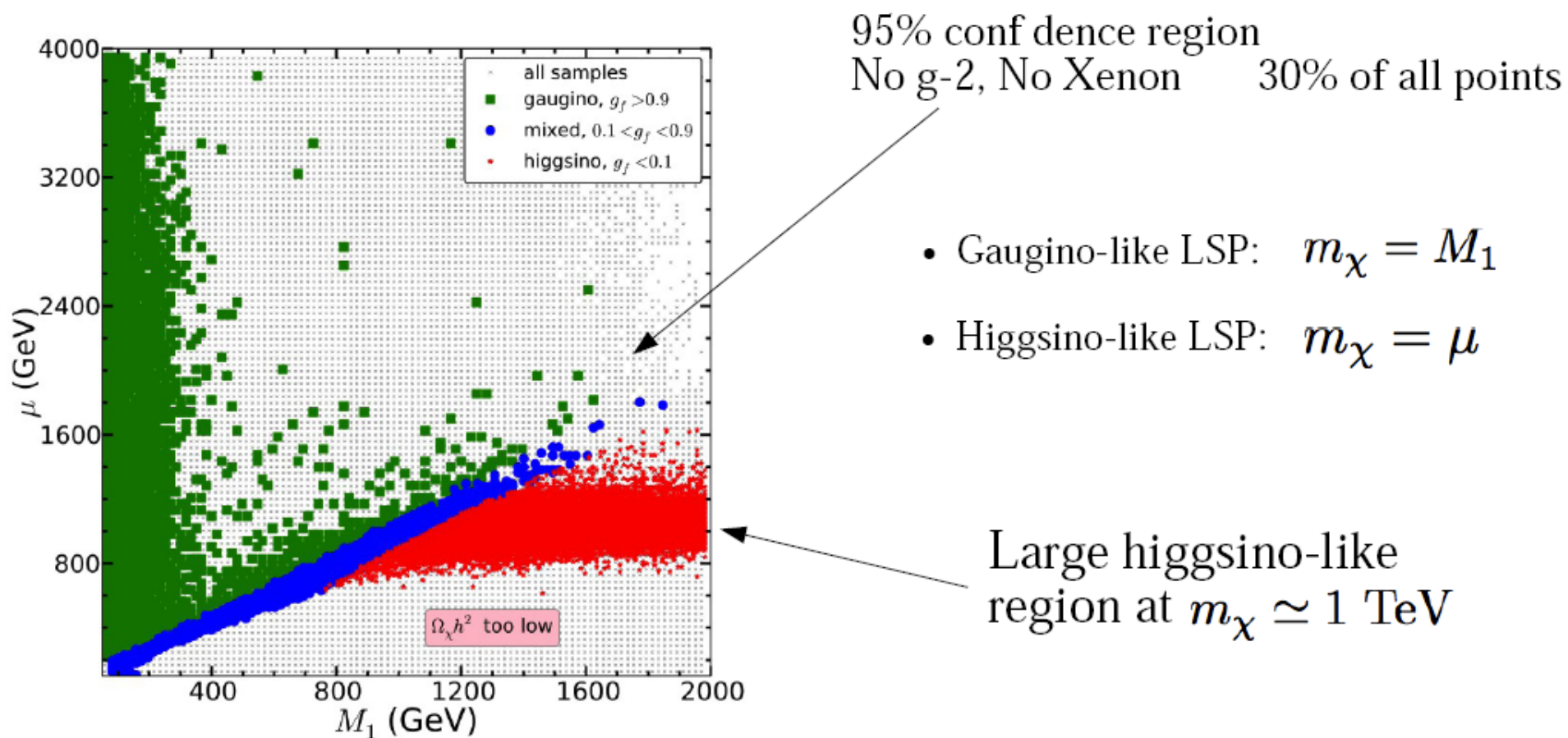
Gaussian Likelihood distribution



Lower-limit Likelihood distribution

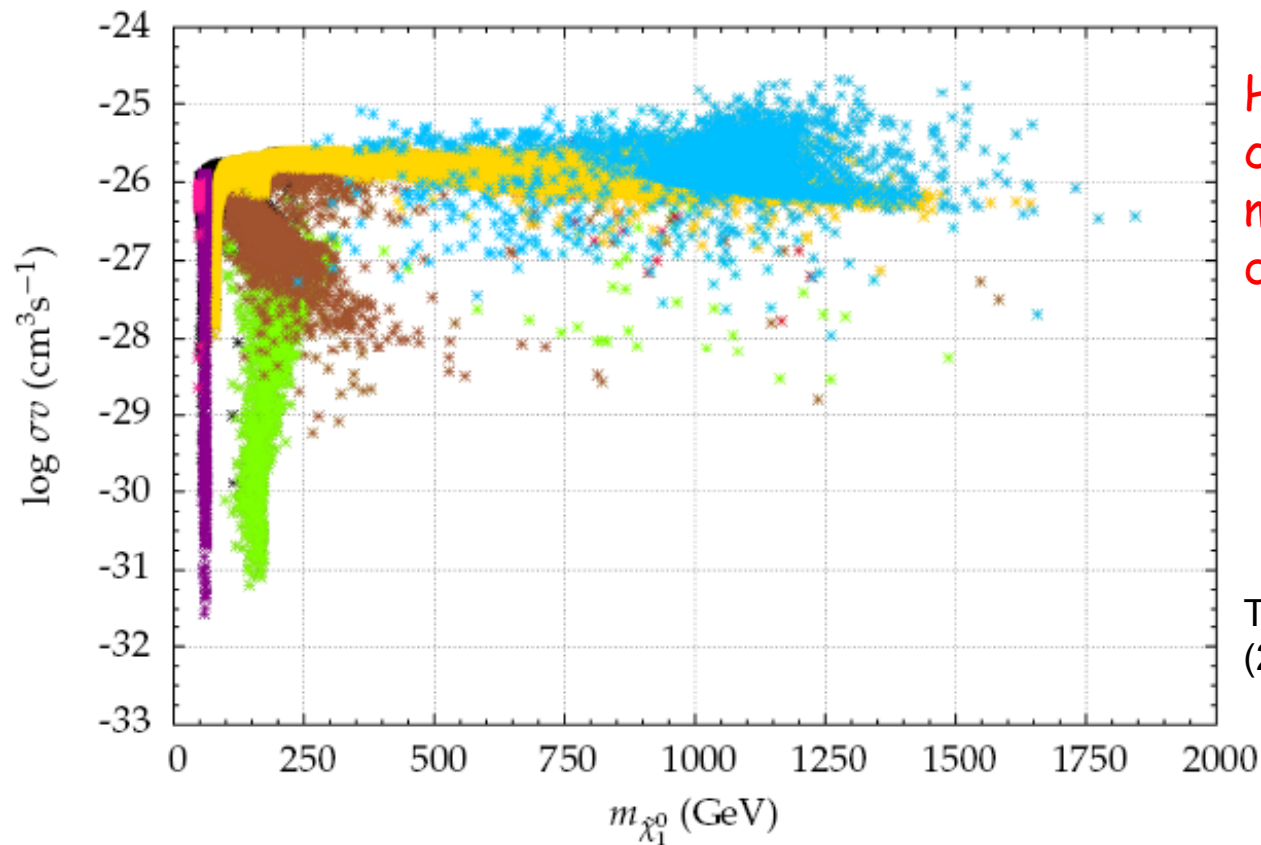
Poisson Likelihood distribution will be shown later...

Impact of the relic density



Bulk \times
 \tilde{f} -coannihilation \times
 $\tilde{\tau}_1$ -coannihilation \times
 $\tilde{\chi}_1^\pm$ -coannihilation \times
 Mixed neutralino \times
 A-funnel \times
 h-funnel \times
 Z-funnel \times

All samples with $\Delta\chi^2 < 5.99$



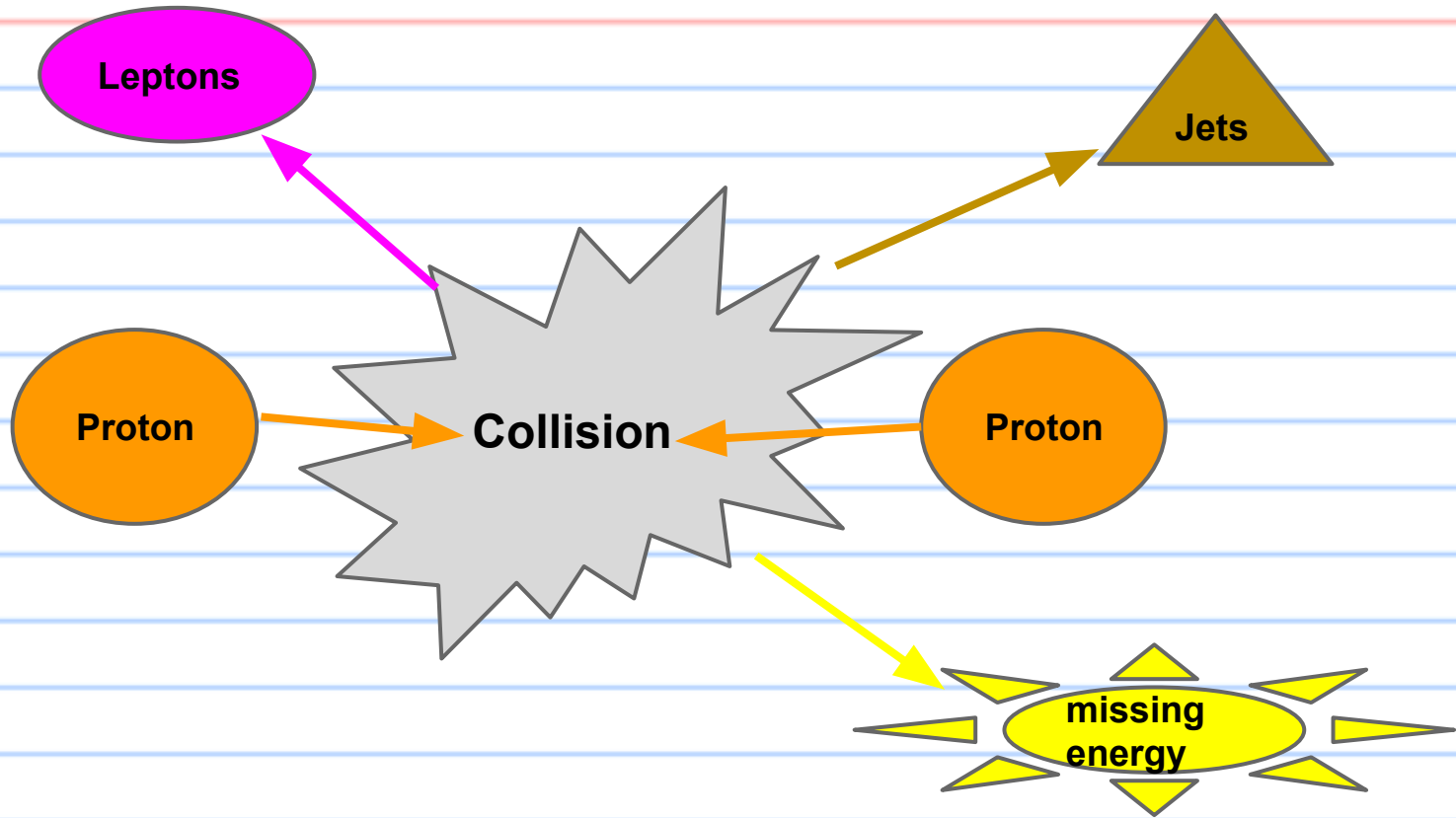
High probability region
can be grouped by their
main contribution to relic
density reduction.

Taken from Fowlie's PhD thesis
(2013)



SUSY Dark Matter search at the LHC

Dark Matter production at the LHC



Dark Matter studies at the LHC are highly model-dependent: one can identify DM candidate.
This will need a confirmation from DM DD and ID search.

315 Physicists Report Failure In Search for Supersymmetry

The negative result illustrates the risks of Big Science, and its often sparse pickings.

By MALCOLM W. BROWNE

Three hundred and fifteen physicists worked on the experiment.

Their apparatus included the Tevatron, the world's most powerful particle accelerator, as well as a \$65 million detector weighing as much as a warship, an advanced new computing system and a host of other innovative gadgets.

But despite this arsenal of brains and technological brawn assembled at the Fermilab accelerator laboratory, the participants have failed to find their quarry, a disagreeable reminder that as science gets harder, even Herculean efforts do not guarantee success.

That was reported 10 years ago about Tevatron...

What has been changed since then?

315 Physicists Report Failure

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27 August 2011 Last updated at 06:41 GMT

7.4K Share



LHC results put supersymmetry theory 'on the spot'



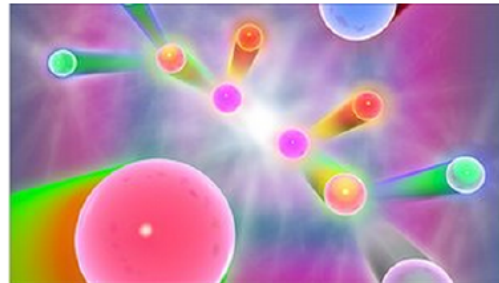
By Pallab Ghosh

Science correspondent, BBC News

Results from the Large Hadron Collider (LHC) have all but killed the simplest version of an enticing theory of sub-atomic physics.

Researchers failed to find evidence of so-called "supersymmetric" particles, which many physicists had hoped would plug holes in the current theory.

Theorists working in the field have told BBC News that they may have to come up with a completely new idea.



Supersymmetry predicts the existence of mysterious super particles.

Data were presented at the Lepton Photon science meeting in Mumbai.

Related Stories

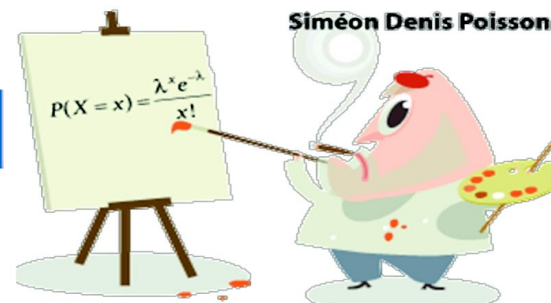
Increasing energy, luminosity and number of physicist failed to find SUSY have increased by factor of 10...

Herculean efforts do not guarantee success.

Let's look at the
situation in the MSSM.

LHC likelihood

Siméon Denis Poisson



Poisson likelihood
 $p(o | s; b)$

CMS boxes for **Observed** and **BG** yields

$$\mathcal{L}(s) = \int p(o | s; b') \exp \left[-\frac{(b - b')^2}{2\delta b^2} \right] db'$$

Simulate "signal"

Table 1: CMS preliminary 2012, $L_{\text{int}} = 11.7 \text{ fb}^{-1}$, $\sqrt{s} = 8 \text{ TeV}$. Comparison of the observed yields in the different H_T bins when requiring $2 \leq n_{\text{jet}} \leq 3$ and $n_b^{\text{reco}} = 0$ for the signal region and data control samples with the SM expectations and combined statistical and systematic uncertainties given by the simultaneous fit.

H_T Bin (GeV)	275–325	325–375	375–475	475–575	575–675	675–775	775–875	875– ∞
SM hadronic	6235^{+100}_{-67}	2900^{+60}_{-54}	1955^{+34}_{-39}	558^{+14}_{-15}	186^{+11}_{-10}	$51.3^{+3.4}_{-3.8}$	$21.2^{+2.3}_{-2.2}$	$16.1^{+1.7}_{-1.7}$
Data hadronic	6232	2904	1965	552	177	58	16	25
SM μ +jets	9696^{+102}_{-83}	5044^{+69}_{-81}	4655^{+57}_{-73}	1806^{+43}_{-34}	766^{+27}_{-24}	307^{+14}_{-16}	147^{+10}_{-11}	195^{+13}_{-13}
Data μ +jets	9698	5039	4653	1808	779	294	150	193
SM $\mu\mu$ +jets	1334^{+28}_{-37}	707^{+26}_{-27}	594^{+18}_{-22}	238^{+10}_{-10}	116^{+9}_{-10}	$46.2^{+3.8}_{-4.8}$	$22.4^{+3.0}_{-2.7}$	$28.6^{+4.3}_{-4.0}$
Data $\mu\mu$ +jets	1336	708	623	205	120	44	21	26
SM γ +jets	–	–	2637^{+68}_{-42}	818^{+24}_{-20}	261^{+16}_{-15}	$85.1^{+5.3}_{-5.7}$	$31.3^{+4.0}_{-4.0}$	$25.4^{+3.3}_{-3.4}$
Data γ +jets	–	–	2601	854	252	94	35	21

SUSY
Spectrum
(model)

Decay BR's

Event generation
X-sec storage
Hadronization

Detector simulation
Object reconstruction
Cards, b -tagging

Kinematical cuts

SoftSUSY

SUSYHIT

PYTHIA

PGS4

Follow CMS

In our scan:

Inclusive prod. for squarks + gluinos

EW production

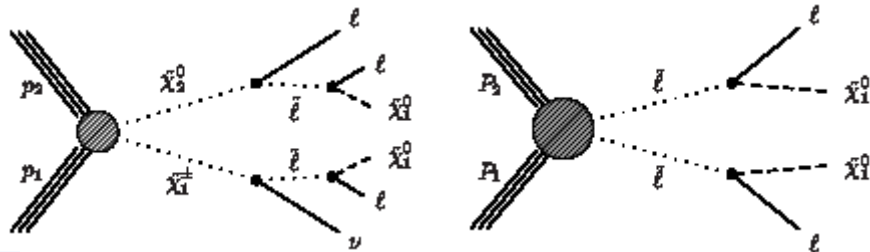
$\alpha_T, \sqrt{s} = 8 \text{ TeV}, 11.7/\text{fb}$

$3l + E_T^{\text{miss}}, \sqrt{s} = 8 \text{ TeV}, 9.2/\text{fb}$

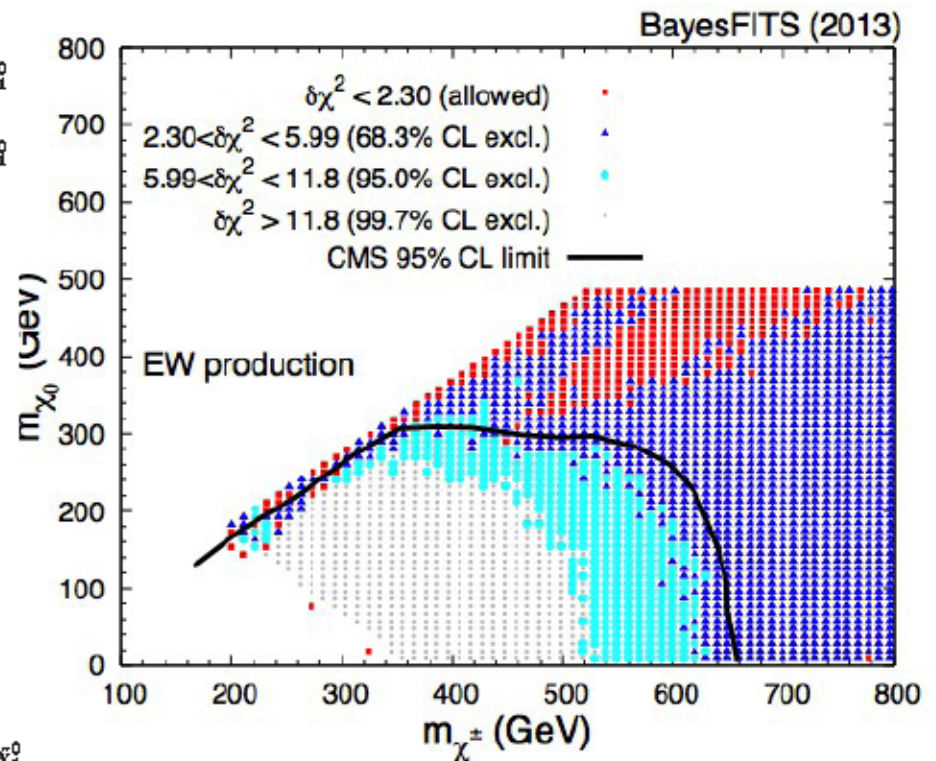
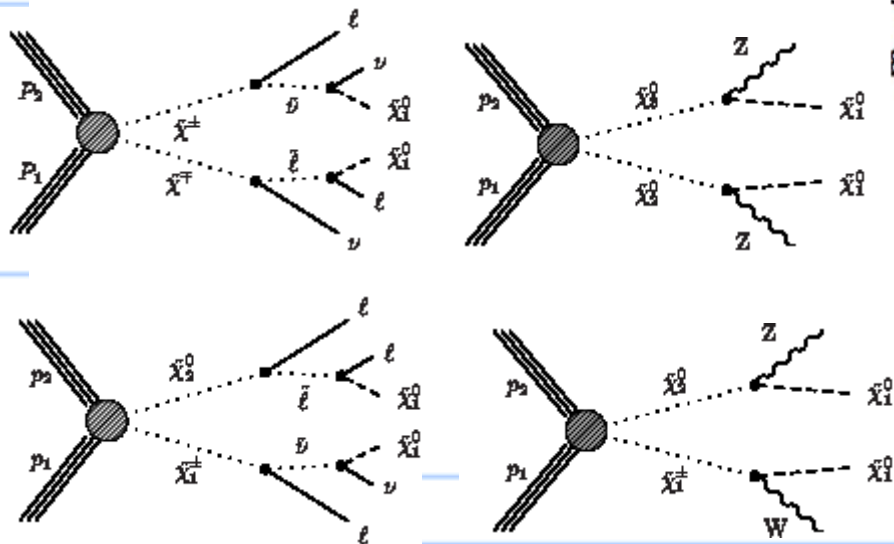
chargino-neutralino pair production (EW production)

$3l + E_T^{\text{miss}}, \sqrt{s} = 8 \text{ TeV}, 9.2/\text{fb}$

(SMS: $m_{\tilde{l}} = 0.5m_{\chi_1^0} + 0.5m_{\chi_1^\pm}$)

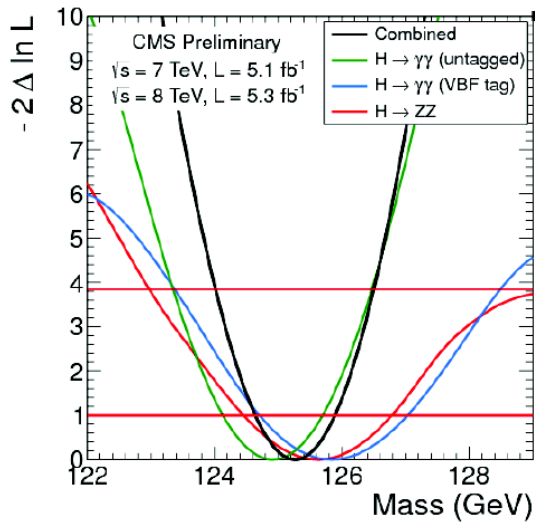


Very Good!



The impact of 125 GeV Higgs mass

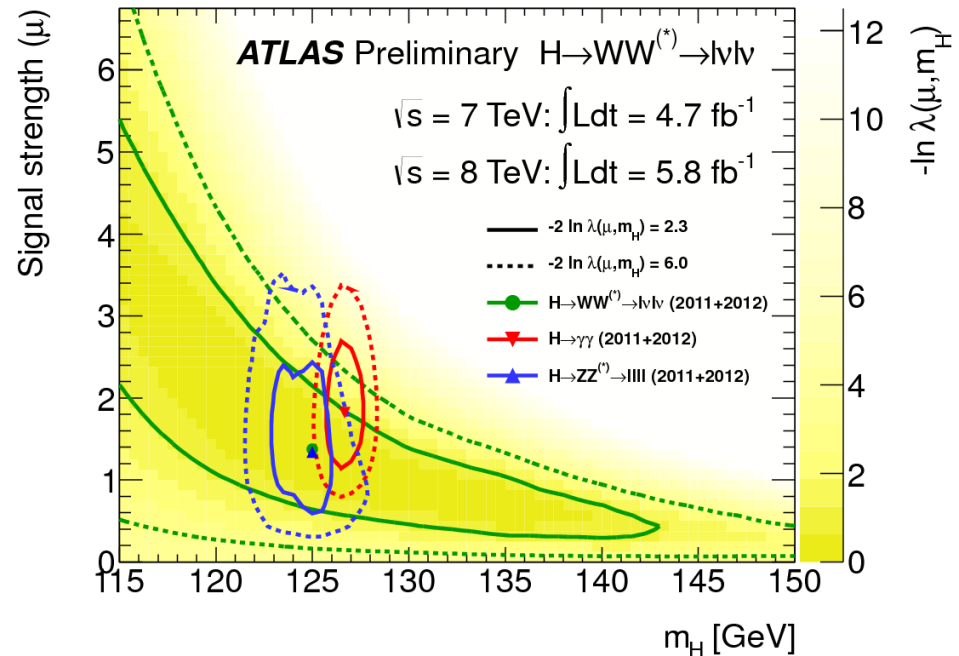
Characterization of the excess: mass



To reduce model dependence,
allow for free cross sections
in three channels
and fit for the common mass:

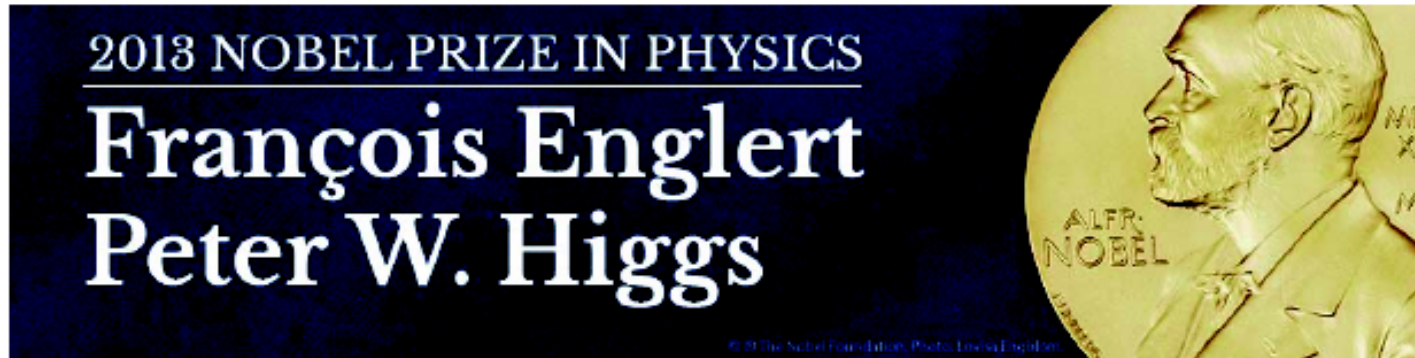
$$m_X = 125.3 \pm 0.6 \text{ GeV}$$

04th July 2012, CMS and
ATLAS



The impact of 125 GeV Higgs mass

Characterization of the excess: **mass**



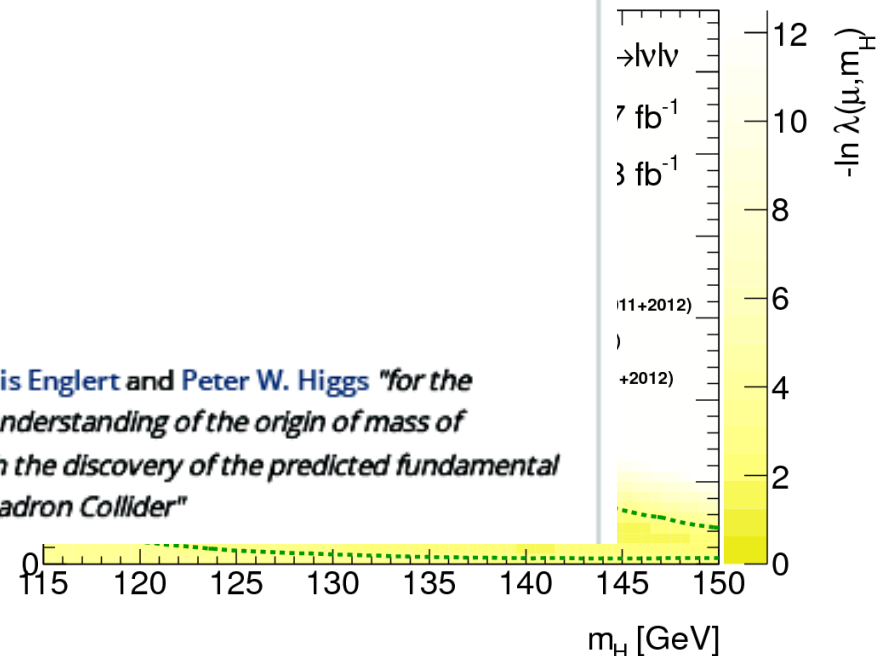
and



F. Englert and P. Higgs
Photo: Wikimedia Commons

2013 Nobel Prize in Physics

The [Nobel Prize in Physics 2013](#) was awarded jointly to [François Englert](#) and [Peter W. Higgs](#) "for the theoretical discovery of a mechanism that contributes to our understanding of the origin of mass of subatomic particles, and which recently was confirmed through the discovery of the predicted fundamental particle, by the ATLAS and CMS experiments at CERN's Large Hadron Collider"



The impact of 125 GeV Higgs mass

Characterization of the excess: **mass**



12th 2012 The Status of the Higgs Search J. Incandella for the CMS COLLABORATION



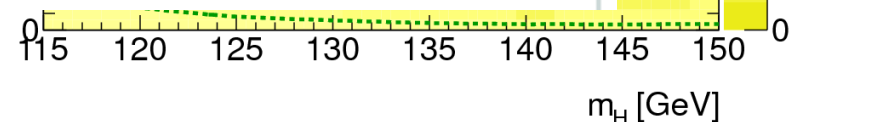
and



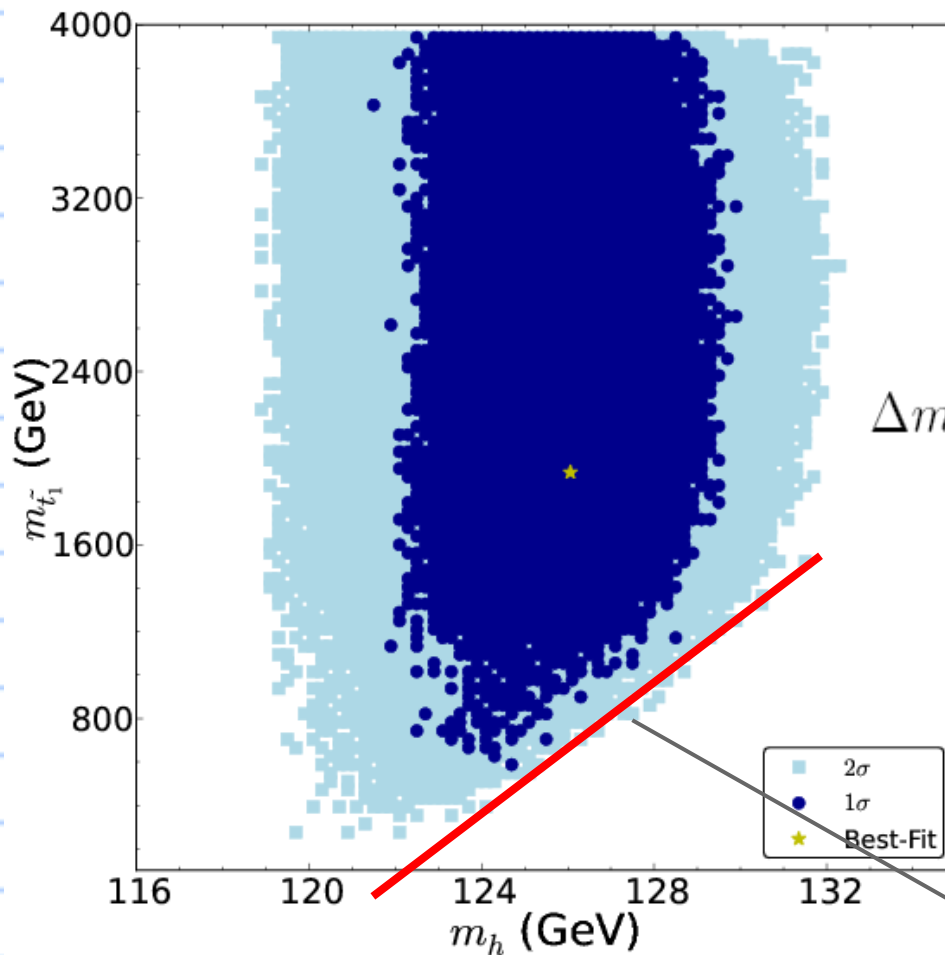
Time to celebrate?
Hold on a moment, is Higgs Standard Model like?

... recently was confirmed through the discovery of the predicted fundamental experiments at CERN's Large Hadron Collider"

for the
mass of



The impact of 126 GeV Higgs mass



input parameters

$$X_t = A_t - \mu \cot \beta$$

$$\Delta m_h^2 \propto \ln \frac{M_{\text{SUSY}}^2}{m_t^2} + \frac{X_t^2}{M_{\text{SUSY}}^2} \left(1 - \frac{X_t^2}{12M_{\text{SUSY}}^2} \right)$$

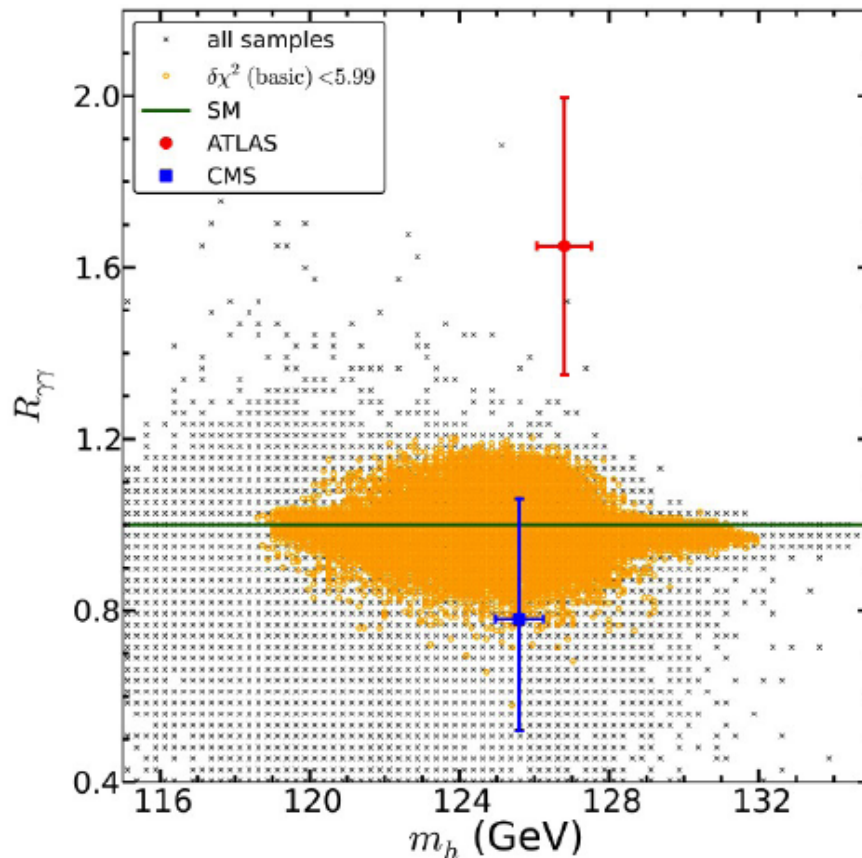
No difficulty to fit Higgs mass at 126 GeV!!!

M_SUSY Scale

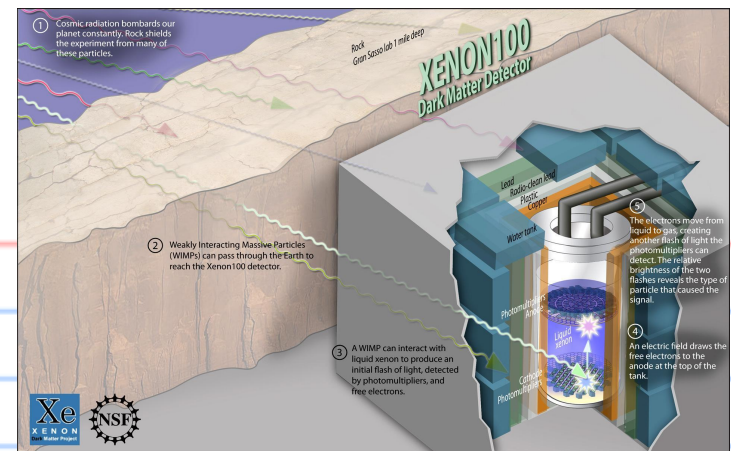
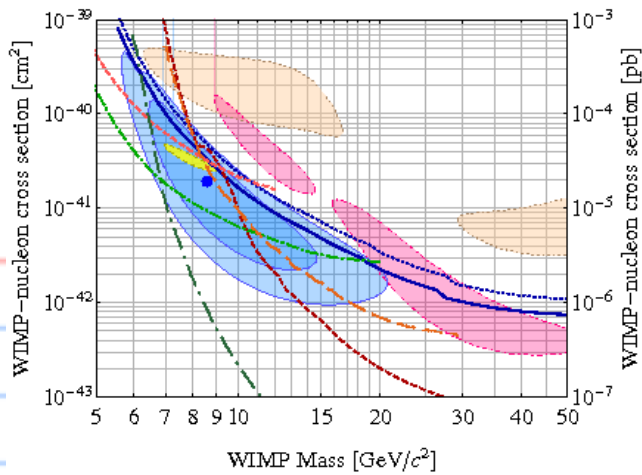
The impact of 126 GeV Higgs mass

Gamma-gamma rate

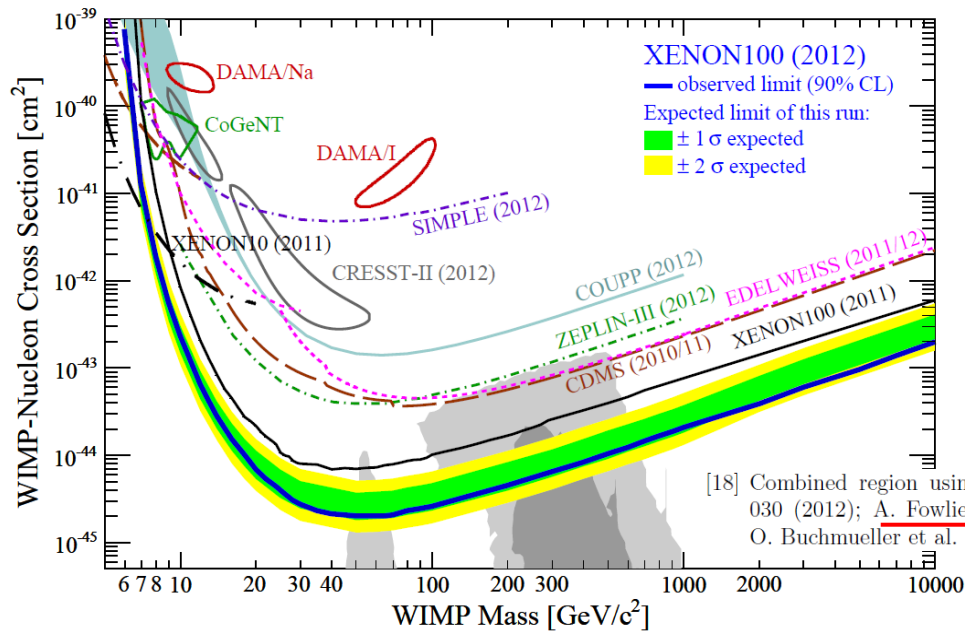
$$R_h(\gamma\gamma) = \frac{\sigma(pp \rightarrow h)}{\sigma(pp \rightarrow h_{\text{SM}})} \times \frac{BR(h \rightarrow \gamma\gamma)}{BR(h_{\text{SM}} \rightarrow \gamma\gamma)} = \sum_{Y \in \text{prod}} \frac{\sigma(pp \rightarrow Y \rightarrow h_{\text{SM}})}{\sigma(pp \rightarrow h_{\text{SM}})} \times \frac{\sigma(Y \rightarrow h)}{\sigma(Y \rightarrow h_{\text{SM}})}$$

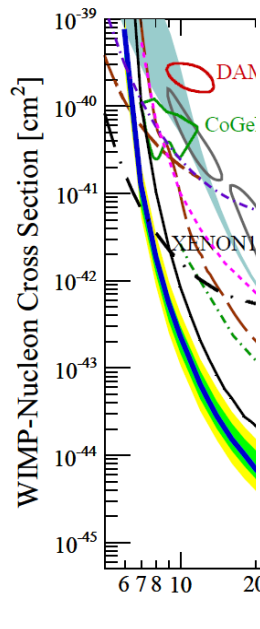
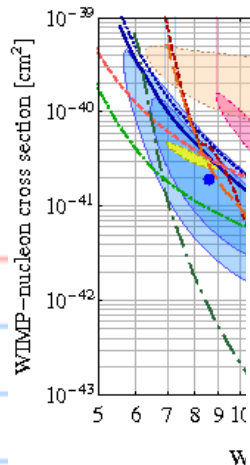


Agree with SM !



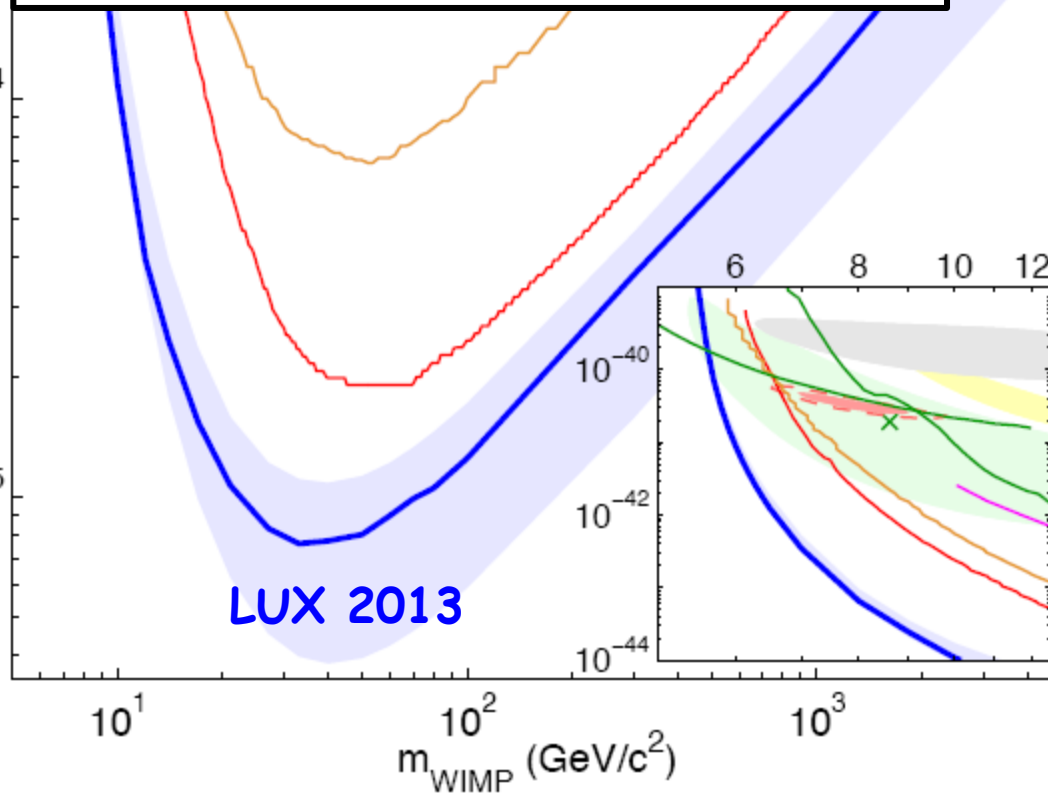
Dark matter direct detection





WIMP-nucleon cross section (cm²)

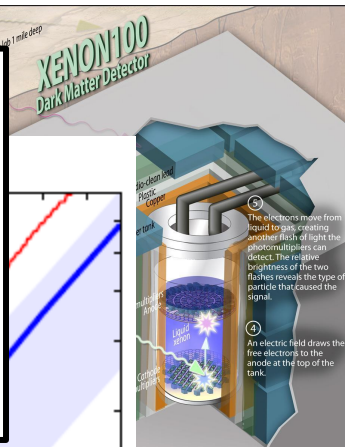
"We find that the LUX data are in strong disagreement with low-mass WIMP signal interpretations of the results from several recent direct detection experiments"



LUX 2013

U. Buchmueller et al. (2011), [arXiv:1112.3564](https://arxiv.org/abs/1112.3564).

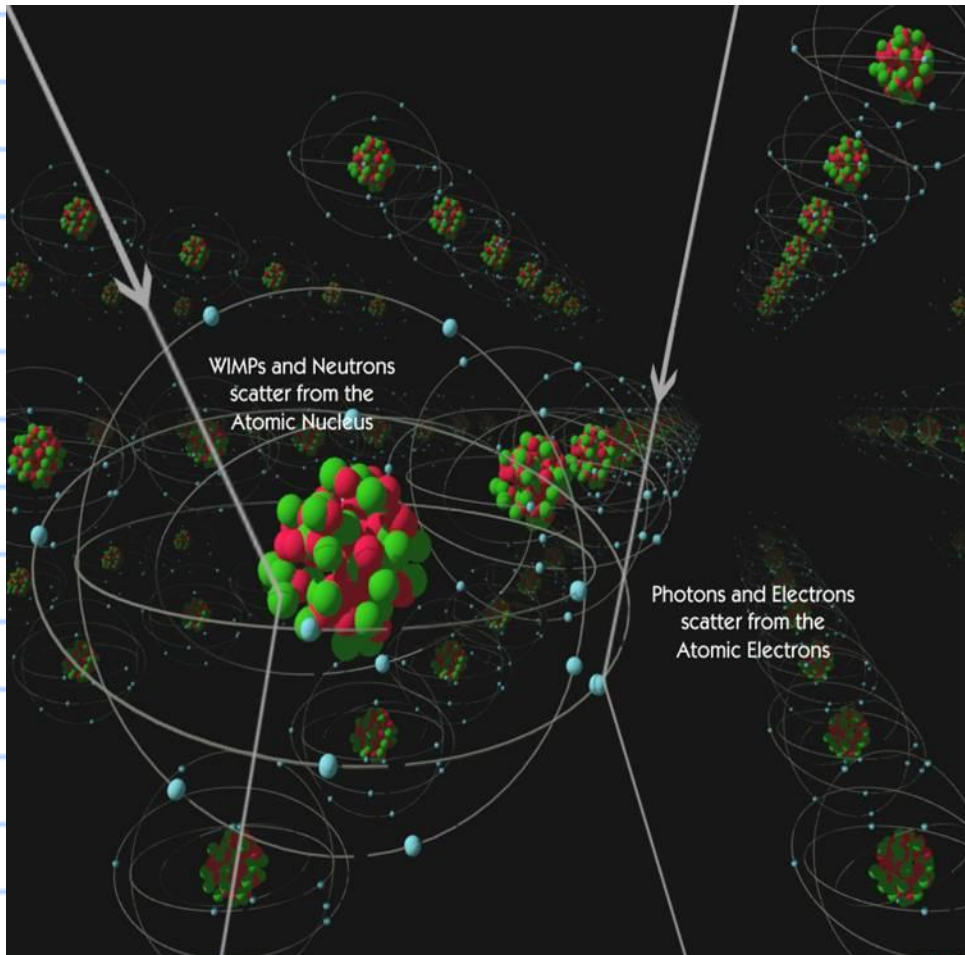
① Cosmic radiation bombards our planet constantly. Rock shields the experiment from many of



on



Detect Dark Matter elastic scattering



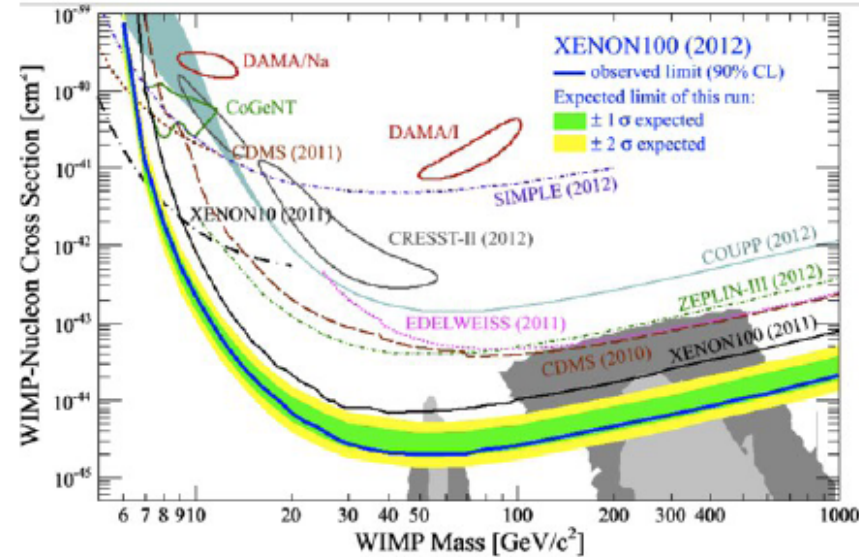
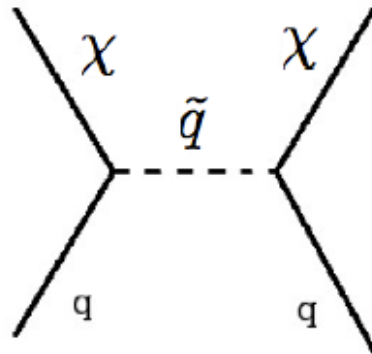
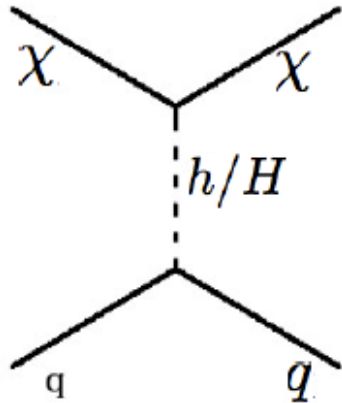
Undetermined direct detection parameters:

- Local dark matter density
- Dark matter velocity distribution
- Dark matter mass
- Form Factor
- effective couplings of DM to protons and neutrons.

Large theoretical uncertainties from lattice physics

Astrophysical uncertainties ~ 2 (1111.0292, 1012.3458)
Form Factor uncertainties ~ 1.2 (hep-ph/0608035v2)

XENON100 likelihood



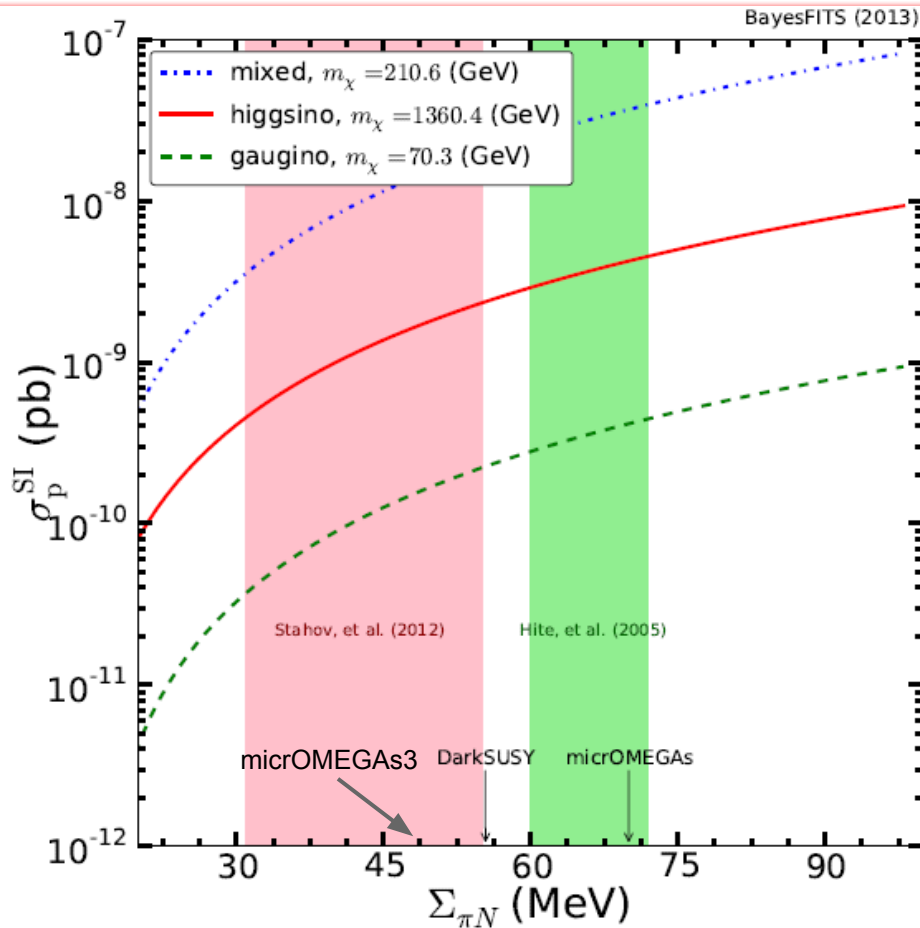
Experimental part:

$o = 2, b = 1, \delta b = 0.2 \longrightarrow$ Simulate signal with micrOMEGAs

$$\mathcal{P}(s + b|o) = \int_0^\infty \frac{e^{-(s+b')}(s+b')^o}{o!} \exp\left[-\frac{(b' - b)^2}{2\delta b^2}\right] db'$$

**Kingman Cheung, Yue-Lin Sming Tsai,
Po-Yan Tseng, Tzu-Chiang Yuan, A. Zee**
Published in JCAP 1210 (2012) 042
e-Print: arXiv:1207.4930

Large theoretical error from lattice calculations



$$\sigma_0 = \frac{m_u + m_d}{2} \langle N | \bar{u}u + \bar{d}d - 2\bar{s}s | N \rangle,$$

$$\Sigma_{\pi N} = \frac{m_u + m_d}{2} \langle N | \bar{u}u + \bar{d}d | N \rangle.$$

From quark level to parton level

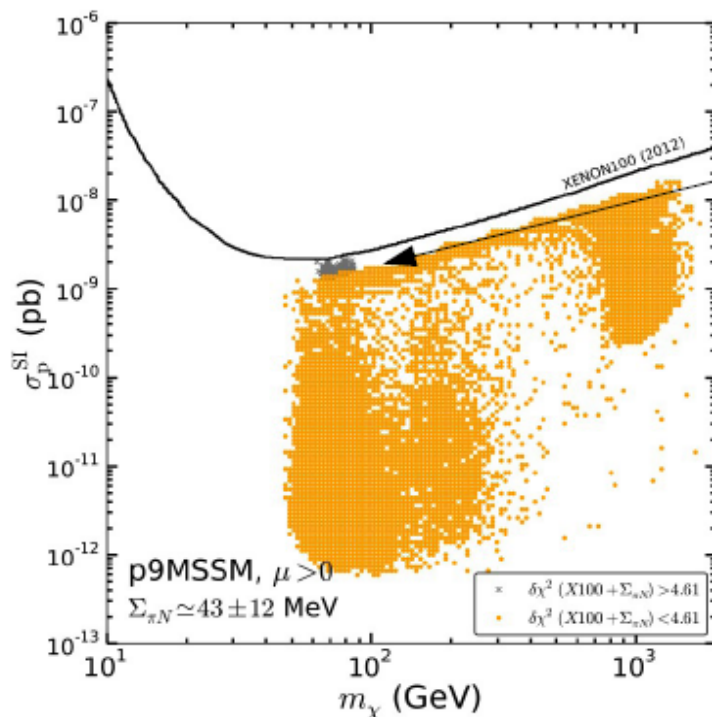
Varying sigma-term between two scenarios (30 MeV and 80 MeV),
spin-independent cross-section can vary by ~one order of magnitude.

The effect of theoretical error

Our likelihood:

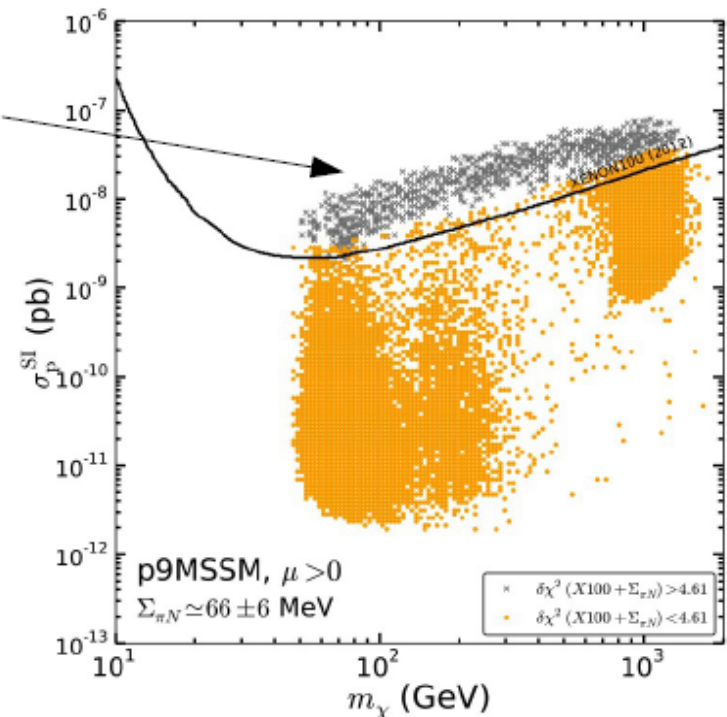
$$\mathcal{L}[m_\chi, \sigma_p^{\text{SI}}(\Sigma_{\pi N})] = \max_{\Sigma'_{\pi N}} \mathcal{P}[\varepsilon s_{\text{mo}}(m_\chi, \sigma_p^{\text{SI}}(\Sigma'_{\pi N})) + b|o] \cdot \exp \left[-\frac{(\Sigma'_{\pi N} - \Sigma_{\pi N})^2}{2\sigma_{\Sigma_{\pi N}}^2} \right]$$

$$\Sigma_{\pi N} = 43 \pm 12 \text{ MeV}$$



90% CL
exclusion

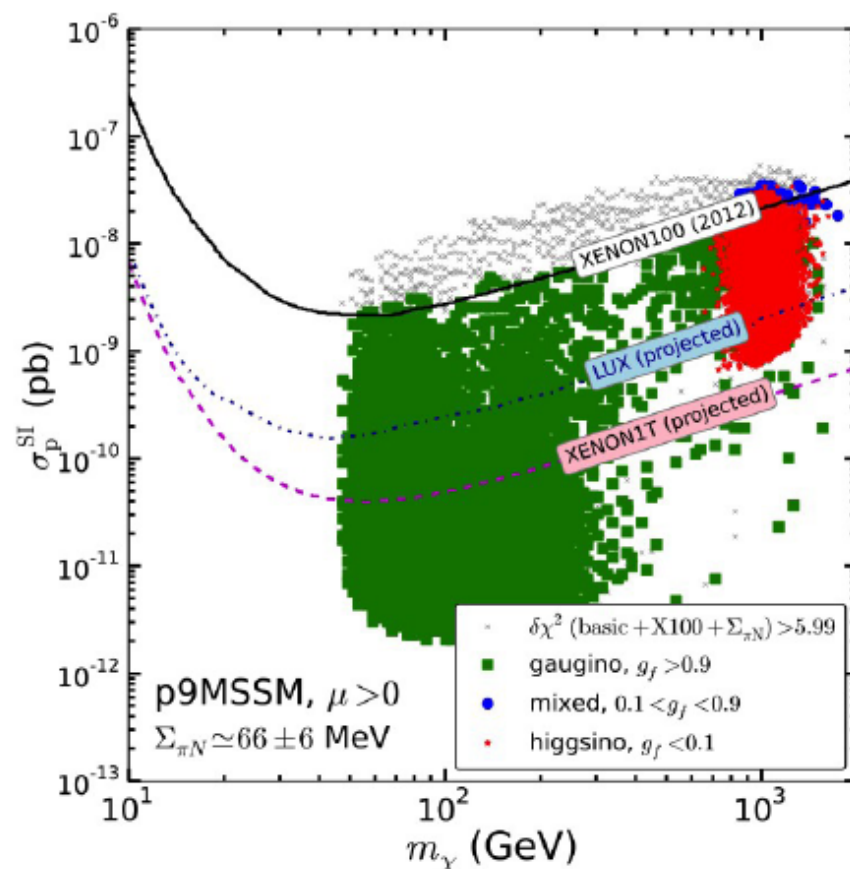
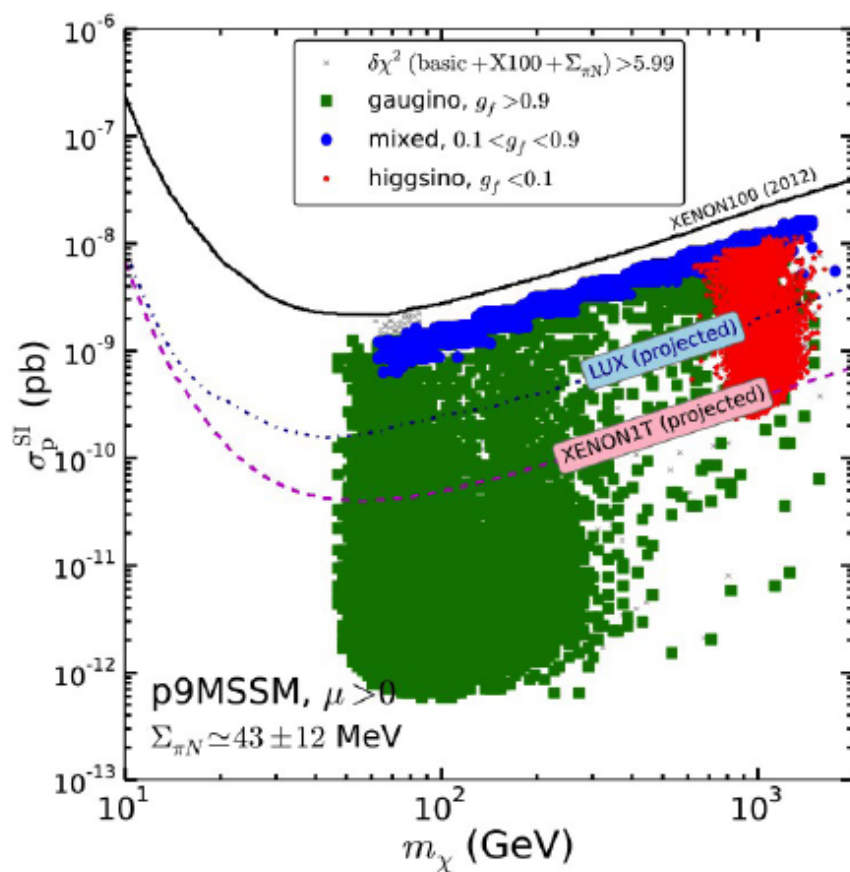
$$\Sigma_{\pi N} = 66 \pm 6 \text{ MeV}$$



Impact of the XENON100

Impact of the XENON100 and future sensitivities at LUX and XENON1T

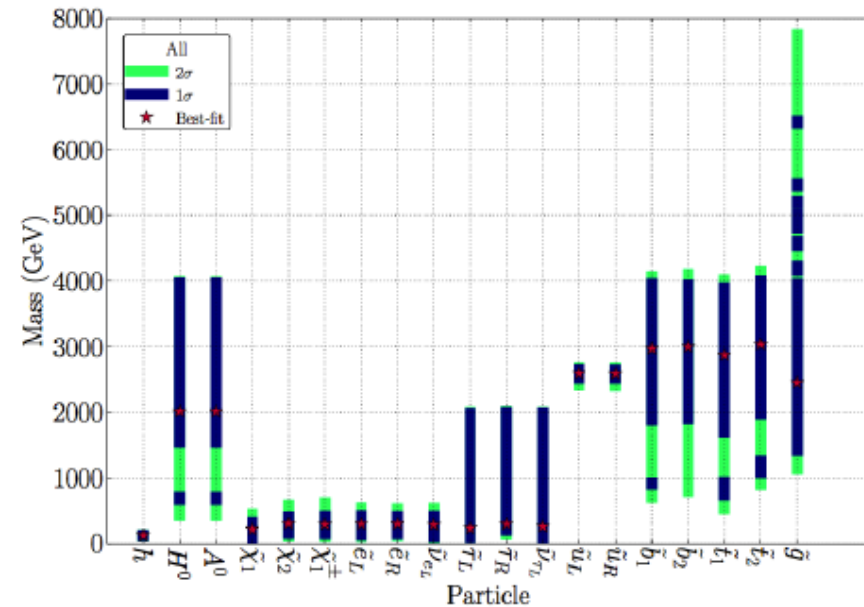
95% confidence region
No g-2



Results

Constraints

Mass spectrum and BF points

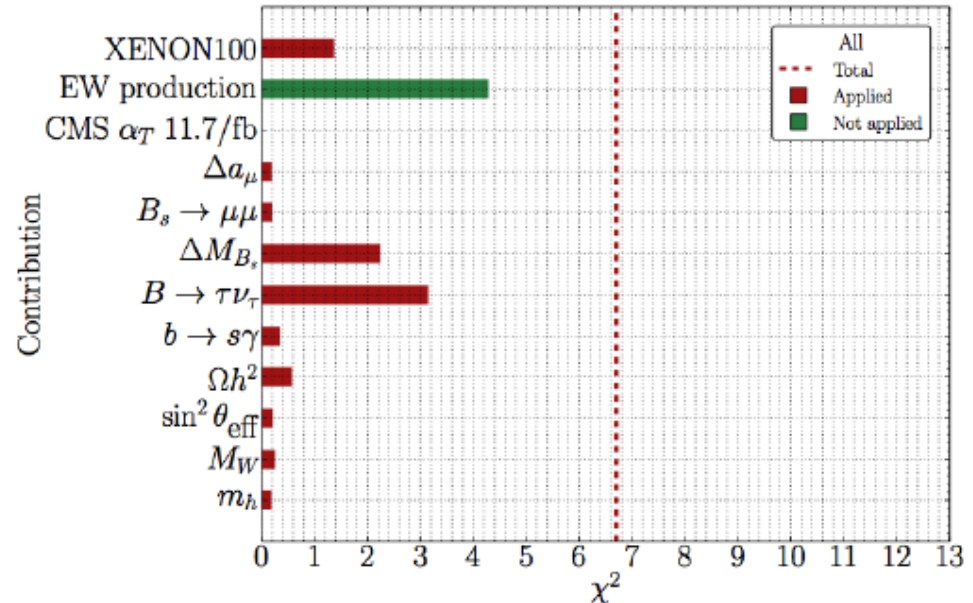


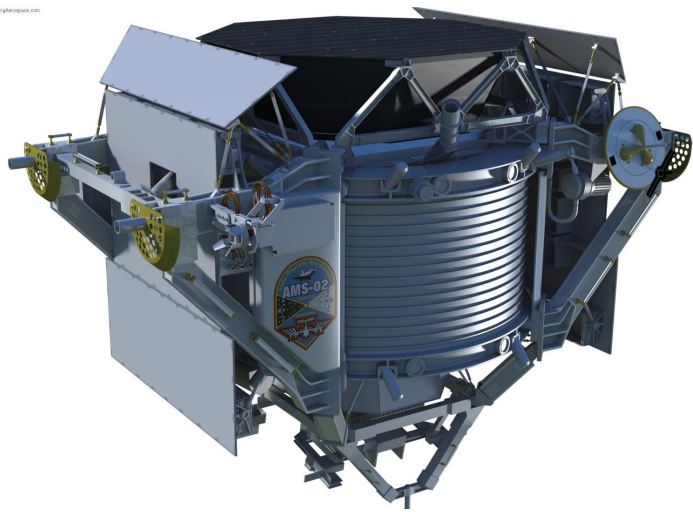
200 GeV < m_X < 500 GeV

Measurement	Mean or Range	Error: (Exp., Th.)	Distribution
CMS α_T 11.7/fb, $\sqrt{s} = 8$ TeV	See text	See text	Poisson
m_h by CMS	125.8 GeV	0.6 GeV, 3 GeV	Gaussian
$\Omega_\chi h^2$	0.1199	0.0027, 10%	Gaussian
BR ($B \rightarrow X_s \gamma$) $\times 10^4$	3.43	0.22, 0.21	Gaussian
BR ($B_u \rightarrow \tau \nu$) $\times 10^4$	1.66	0.33, 0.38	Gaussian
ΔM_{B_s}	17.719 ps ⁻¹	0.043 ps ⁻¹ , 2.400 ps ⁻¹	Gaussian
$\sin^2 \theta_{\text{eff}}$	0.23146	0.00012, 0.00015	Gaussian
M_W	80.399 GeV	0.023 GeV, 0.015 GeV	Gaussian
BR ($B_s \rightarrow \mu^+ \mu^-$) $\times 10^9$	3.2	+1.5 - 1.2, 10%	Gaussian
$m_b(m_b)^{\overline{MS}}$	4.18 GeV	0.03 GeV, 0	Gaussian
M_t	173.5 GeV	1.0 GeV, 0	Gaussian
$\delta(g - 2)_\mu^{\text{SUSY}} \times 10^{10}$	28.7	8.0, 1.0	Gaussian
XENON100 (2012)	See text	See text	Poisson
CMS $3l + E_T^{\text{miss}}$ 9.2/fb, $\sqrt{s} = 8$ TeV	See text	See text	Poisson

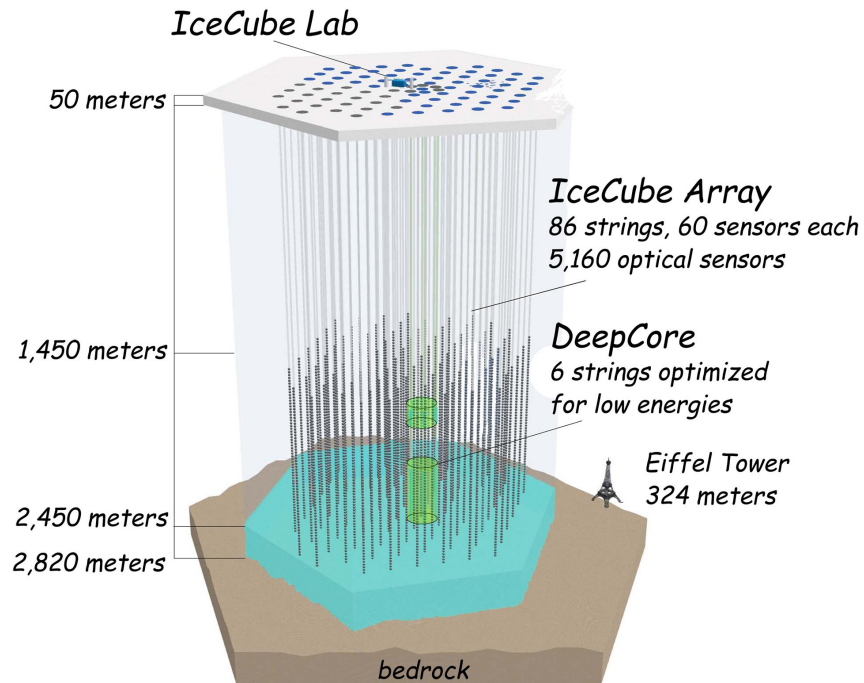
PLUS LEP CONSTRAINTS:

- $m_\chi > 46$ GeV,
- $m_e > 107$ GeV,
- $m_{\chi_1^\pm} > 94$ GeV if $m_{\chi_1^\pm} - m_\chi > 3$ GeV and $\tan \beta < 40$,
- $m_{\tilde{\mu}} > 94$ GeV if $m_{\tilde{\mu}} - m_\chi > 10$ GeV and $\tan \beta < 40$,
- $m_{\tilde{\tau}} > 81.9$ GeV if $m_{\tilde{\tau}} - m_\chi > 15$ GeV,
- $m_{\tilde{b}_1} > 89$ GeV if $m_{\tilde{b}_1} - m_\chi > 8$ GeV,
- $m_{\tilde{t}_1} > 95.7$ GeV if $m_{\tilde{t}_1} - m_\chi > 10$ GeV.

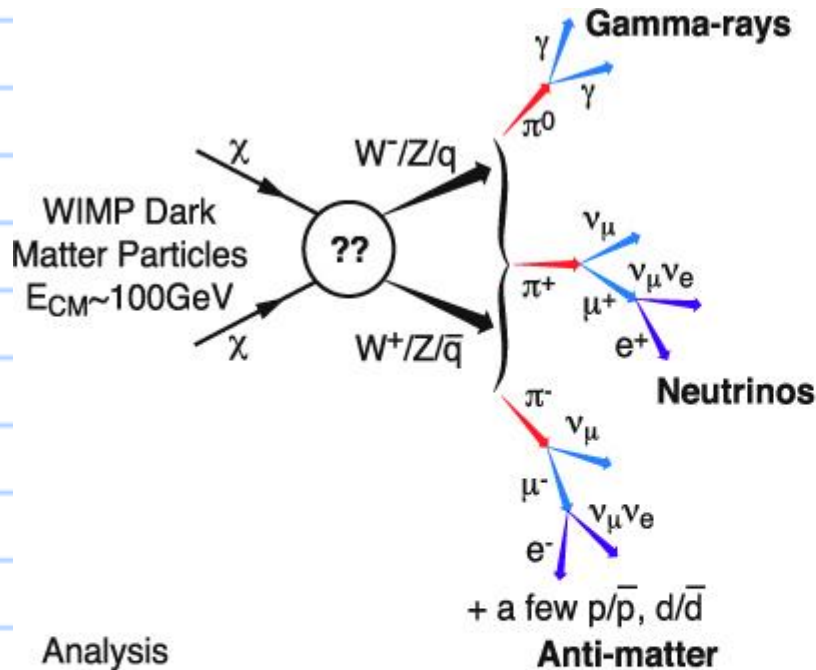




Neutralino indirect detection

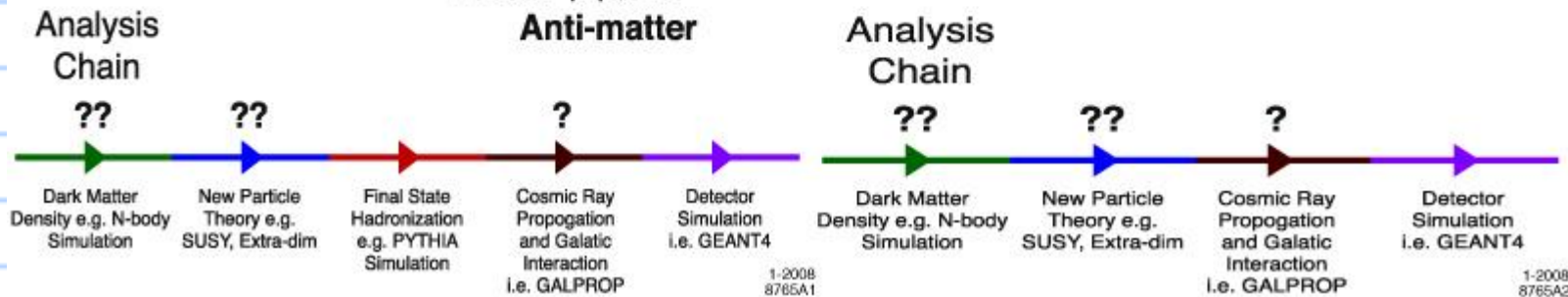
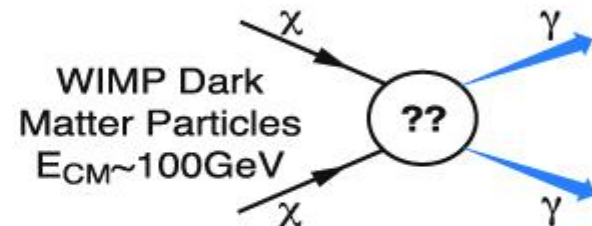


Detect Dark Matter annihilation



Undetermined indirection detection inputs:

- Astrophysics background
- Halo model
- Dark matter mass
- annihilation cross-sections



Fermi DM gamma ray search



Source	Advantage	Disadvantage
Galactic Center	Strong DM signal, good statistics	High astrophysics background, unclear source
Milky Way Halo	Large statistics	High astrophysics background
dSphs	Low astrophysics background	Low statistics
Gamma-ray line	No similar astrophysical signal	Low statistics
Extragalactic gamma-ray background	Large statistics	Huge astrophysical uncertainties
Galaxy Clusters	Low astrophysics background	Low statistics and astrophysical uncertainties

neutralino DM gamma ray search



Fermi dSphs data can provide sensitive bound for CMSSM.

See: Roszkowski, Sessolo, Tsai (1202.1503)

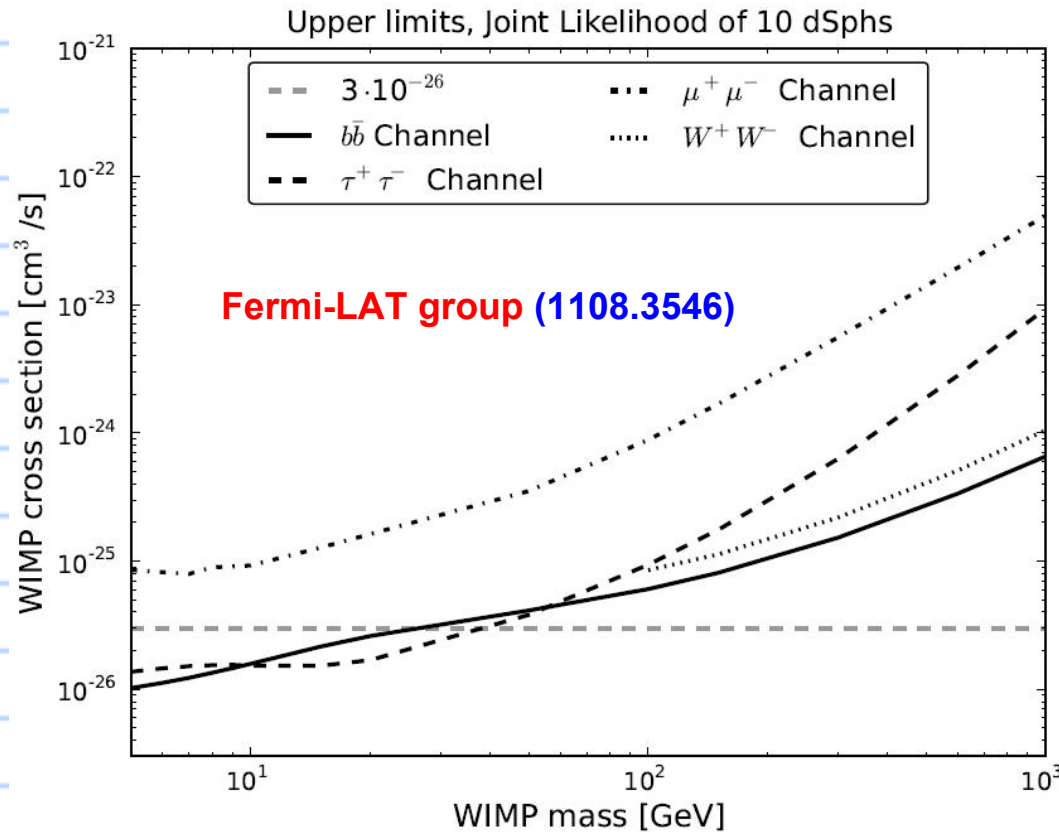
Source	Advantage	Disadvantage
Galactic Center	Strong DM signal, good statistics	High astrophysics background, unclear source
Milky Way Halo	Large statistics	High astrophysics background
dSphs	Low astrophysics background	Low statistics
Gamma-ray line	No similar astrophysical signal	Low statistics
Extragalactic gamma-ray background	Large statistics	Huge astrophysical uncertainties
Galaxy Clusters	Low astrophysics background	Low statistics and astrophysical uncertainties

MSSM neutralino flux has $O(1\sim 2)$ lower than Fermi data.

NMSSM: one can have larger fluxes.

See: Guillaume Chalons (1204.4591) and Das, Ellwanger, Mitropoulos (1206.2639)

DM gamma rays from dSphs



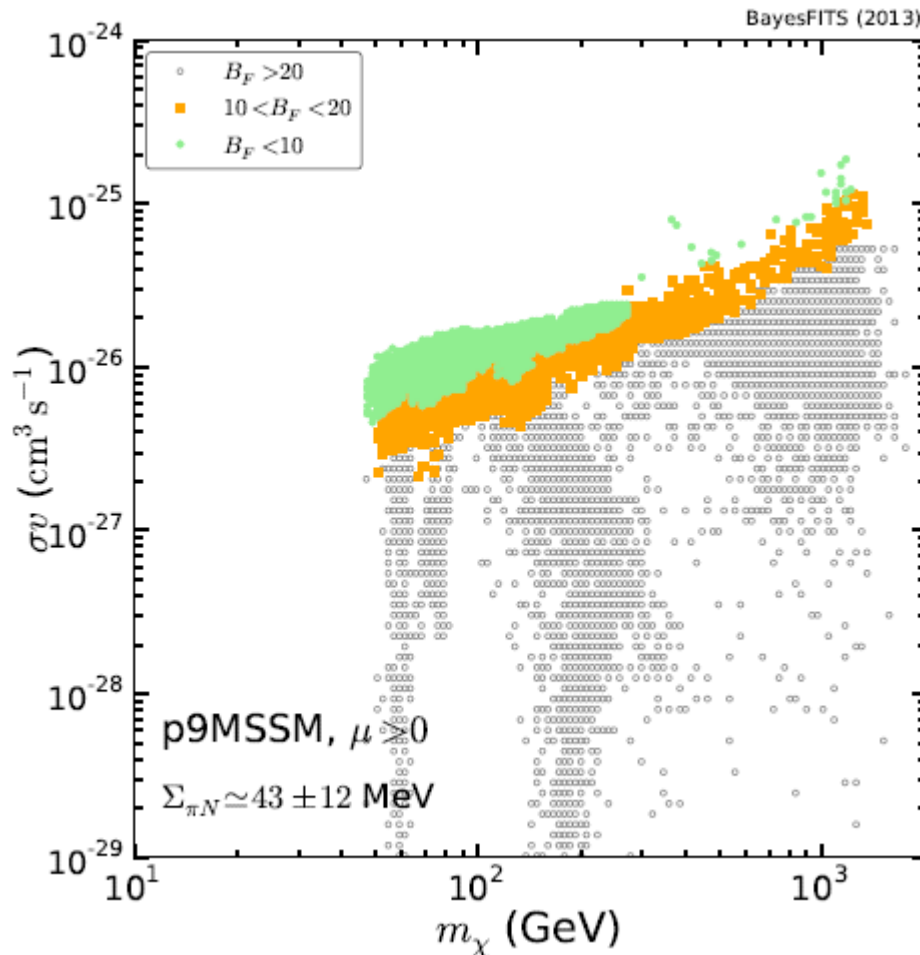
$$\frac{d\Phi_\gamma}{dE_\gamma}(E_\gamma, \psi) = \sum_i \frac{\sigma_i v}{8\pi m_\chi^2} \frac{dN_\gamma^i}{dE_\gamma} \int_{\text{l.o.s.}} dl \rho_\chi^2(r(l, \psi))$$

TABLE I. Position, distance, and J factor (under assumption of a Navarro-Frenk-White profile) of each dSph. The 4th column shows the mode of the posterior distribution of $\log_{10} J$, and the 5th column indicates its 68% C.L. error. See the text for further details. The J factors correspond to the pair annihilation flux coming from a cone of solid angle $\Delta\Omega = 2.4 \cdot 10^{-4}$ sr. The final column indicates the reference for the kinematic dataset used.

Name	l deg.	b deg.	d kpc	$\log_{10}(J)$ $\log_{10}[\text{GeV}^2 \text{cm}^{-5}]$	σ	ref.
Bootes I	358.08	69.62	60	17.7	0.34	[15]
Carina	260.11	-22.22	101	18.0	0.13	[16]
Coma Berenices	241.9	83.6	44	19.0	0.37	[17]
Draco	86.37	34.72	80	18.8	0.13	[16]
Fornax	237.1	-65.7	138	17.7	0.23	[16]
Sculptor	287.15	-83.16	80	18.4	0.13	[16]
Segue 1	220.48	50.42	23	19.6	0.53	[18]
Sextans	243.4	42.2	86	17.8	0.23	[16]
Ursa Major II	152.46	37.44	32	19.6	0.40	[17]
Ursa Minor	104.95	44.80	66	18.5	0.18	[16]

Fermi-LAT put the upper limit on WIMP annihilation cross section.

DM gamma rays from dSphs



Parameter	gaugino	mixed	higgsino
$\tan \beta$	7.94	52.63	4.76
M_2 (GeV)	148.56	457.82	3810.98
M_3 (GeV)	1847.66	2785.12	2281.29
μ (GeV)	620.95	275.72	1345.21
m_A (GeV)	1454.49	3648.84	2716.89
A_t (GeV)	2086.52	3607.58	5277.32
A_τ (GeV)	-2786.17	5256.12	-4519.05
$m_{\tilde{Q}_3}$ (GeV)	3335.38	2330.03	3577.35
$m_{\tilde{L}_3}$ (GeV)	144.74	1613.40	1500.82
m_χ (GeV)	70.29	210.61	1360.36
σv ($\text{cm}^3 \text{s}^{-1}$)	1.33×10^{-26}	2.64×10^{-26}	3.75×10^{-25}
B_F	4.49	6.11	3.14

Fermi LAT data taken from 4 August 2008 to 2 August 2012 with the pass 7 photon selection, and energy from 200 MeV to 500 GeV.

Fermi-dSphs data is more sensitive to low neutralino mass region.

Uncertainties of DM annihilation to gamma rays

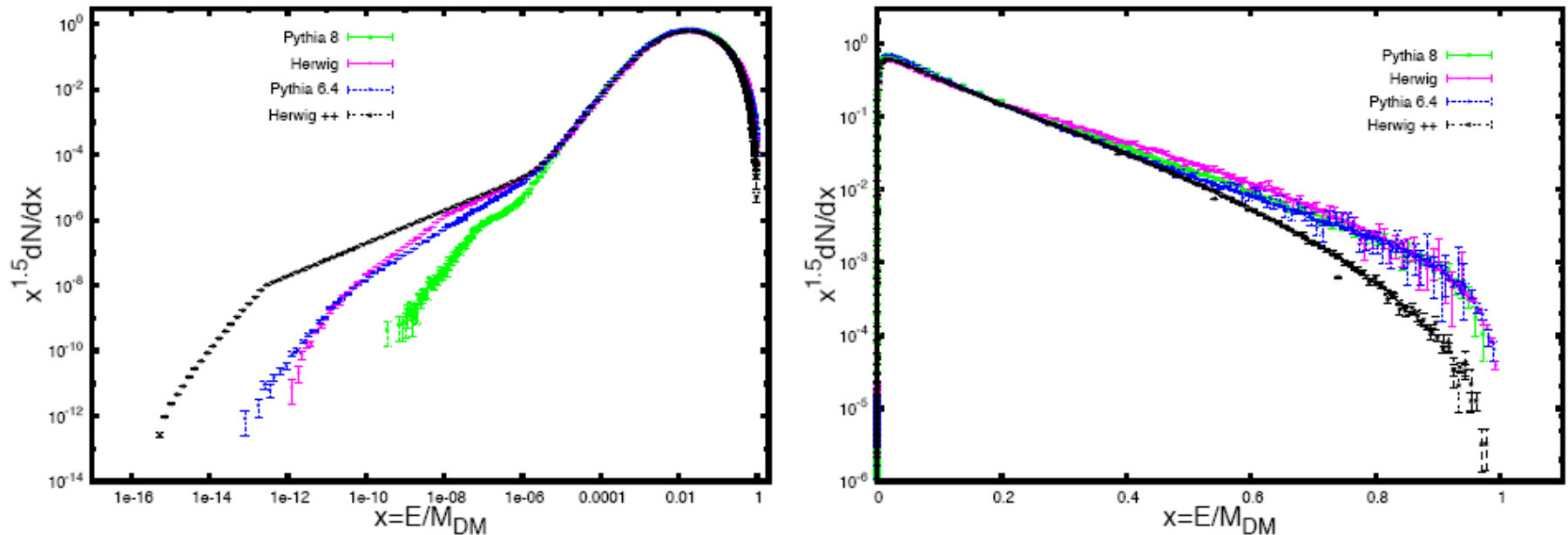
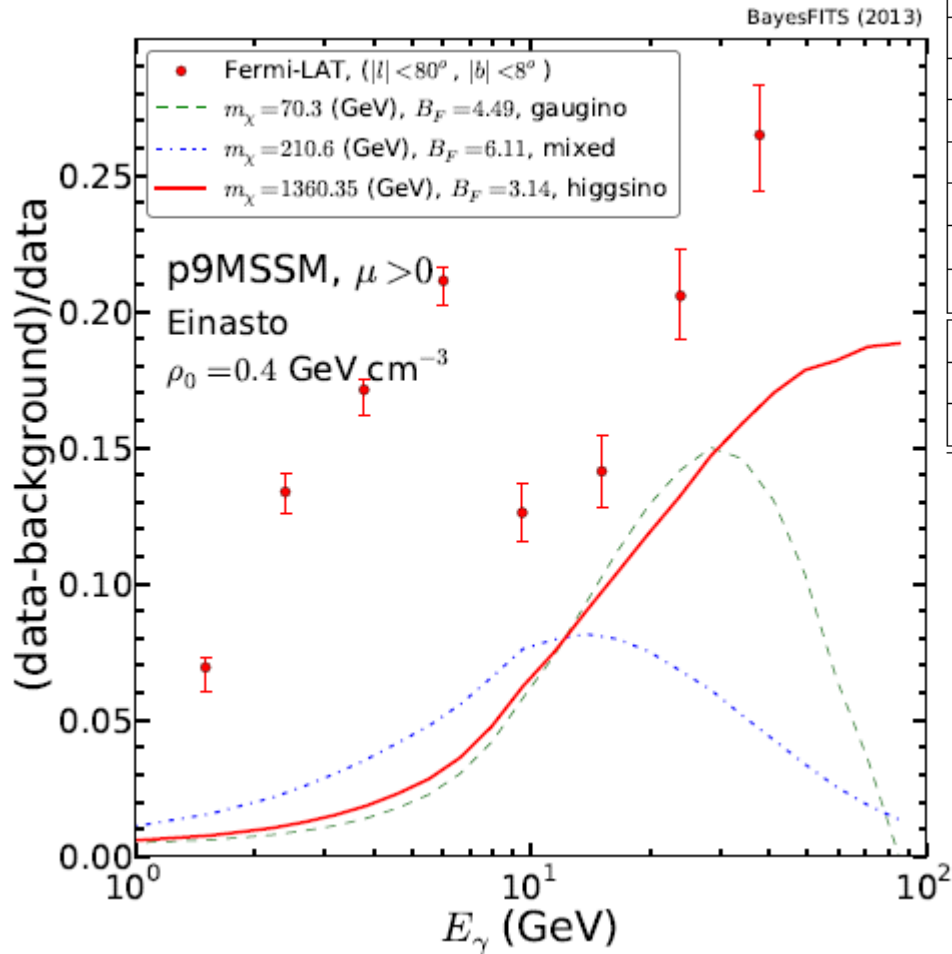


Figure 2. (*Left-panel*) W^+W^- annihilation channel with $M_{\text{DM}} = 1$ TeV in logarithmic scale. As in Fig. 1, the simulations are consistent down to a value of x , that is 10^{-6} in the case of $M_{\text{DM}} = 1$ TeV (a factor ten lower in x with respect to the case with $M_{\text{DM}} = 100$ GeV). Similar behaviors of the lower energy cuts-off are also observed, with a general shift of x cut-off value of order 10^{-2} . (*Right-panel*) W^+W^- annihilation channel with $M_{\text{DM}} = 1$ TeV in linear scale. All the simulations except for HERWIG++ exhibit the same behavior as in Fig. 1, but within $x \simeq 0.3$ and $x \simeq 0.7$ and a maximum discrepancy at $x \simeq 0.5$. The shift with respect to Fig. 1 can be simply explained by the increment of the WIMP mass.

Can be more than one order of magnitude

Taken from
1305.2124

DM gamma rays from GC

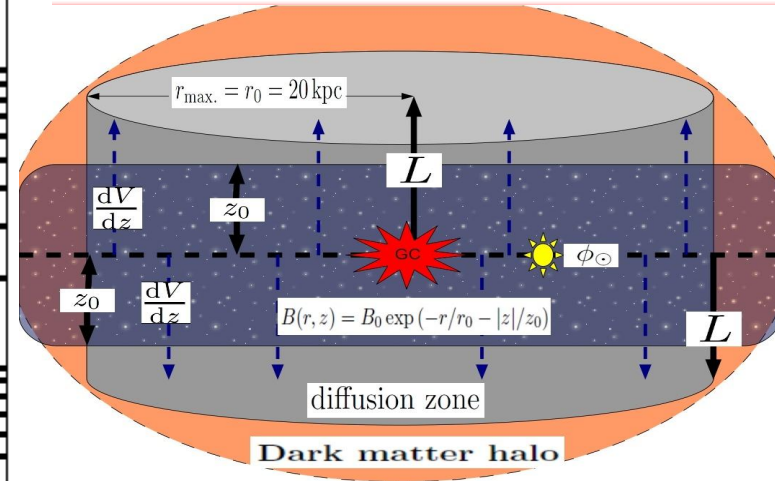
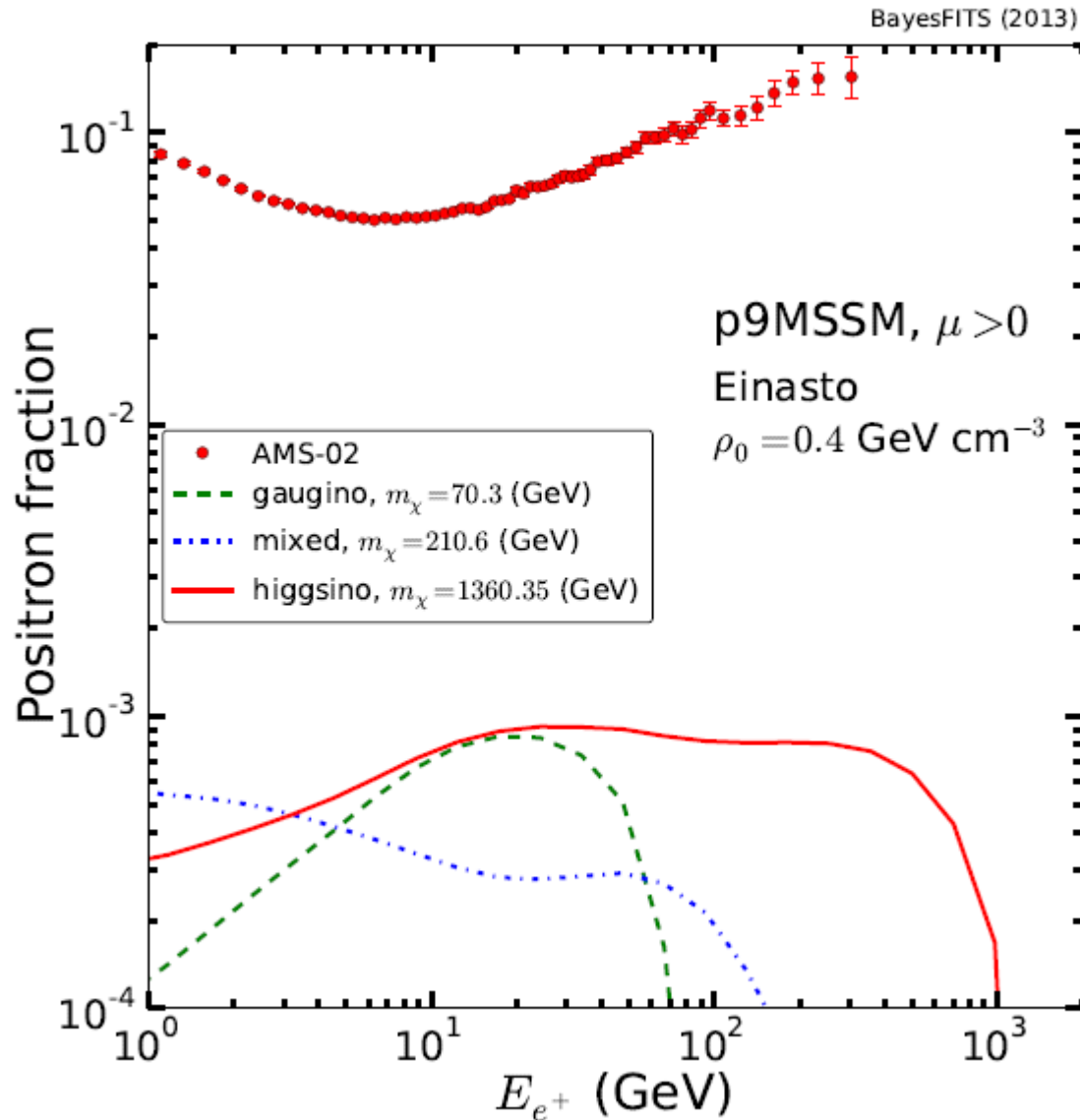


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B_F	4.49	6.11	3.14

1. Dark matter π^0 -decay
2. Inverse Compton
3. bremsstrahlung

SUSY can be also testable in Fermi-GC search.

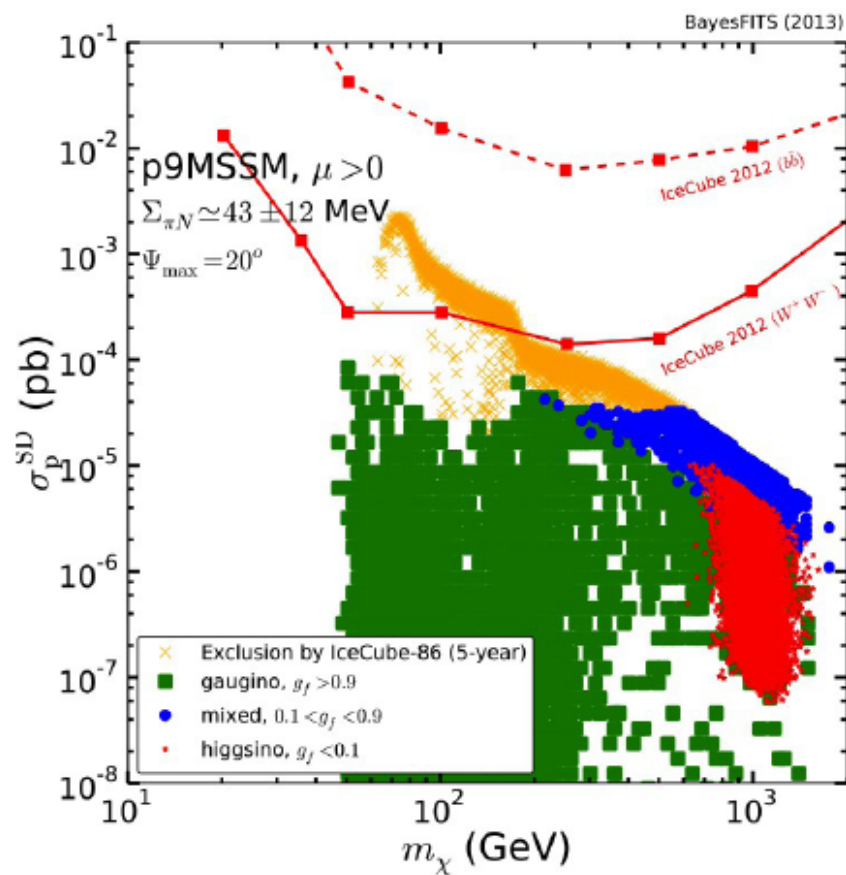
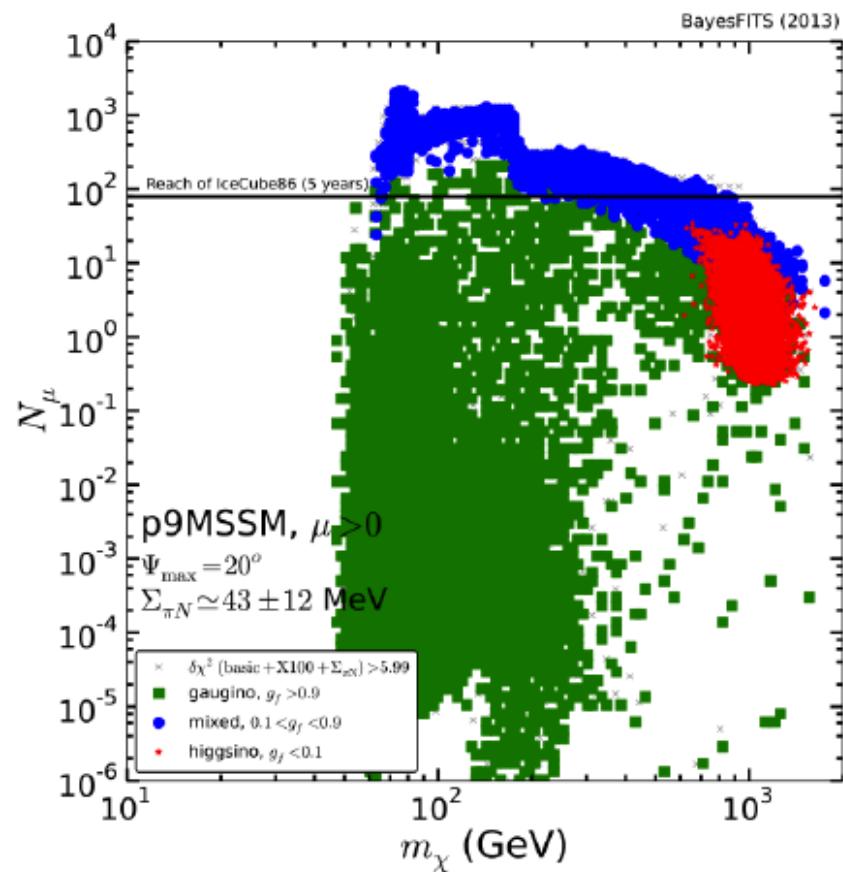
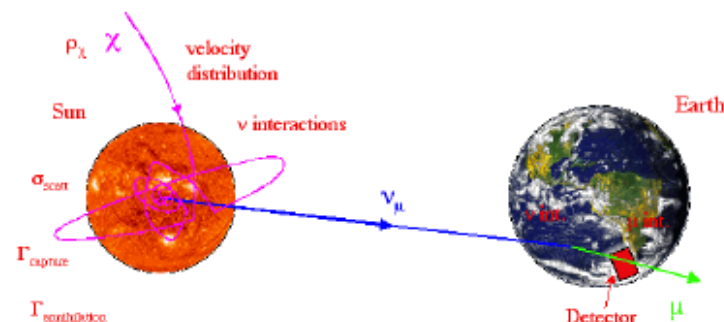
Positron signatures of neutralino DM



SUSY does not reproduce AMS-02 e^+ fluxes (pulsars).

Impact of DM indirect detection

Fermi-LAT, AMS02, IceCube
(see arXiv:1306.1567)



Model
↓
FR Model File



Galileo

- Models from symmetry
 - Christensen, Salmon, Setzer, Stefanus

↓
FeynRules

↓
Validation

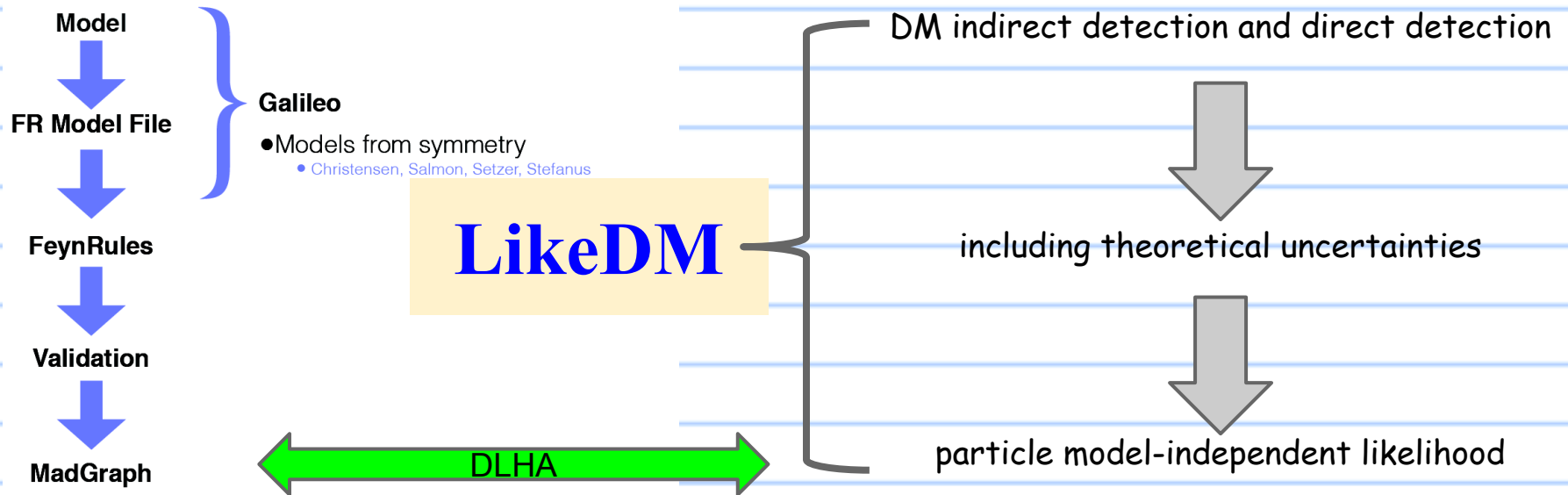
↓
MadGraph
(micromegas)

How about the experimental constraints?

Slide taken from MadGraph School 2013 Taipei,
Neil Christensen

LikeDM code

in Collaboration with Q. Yuan and X. Huang



- We can more confidently and efficiently check every dark matter model.
- Can be extend to cosmology constraints.
- Similar to “DMFIT” but starting from data level

(Tesla E. Jeltema, Stefano Profumo)

Fitting DM gamma rays by using FermiTools

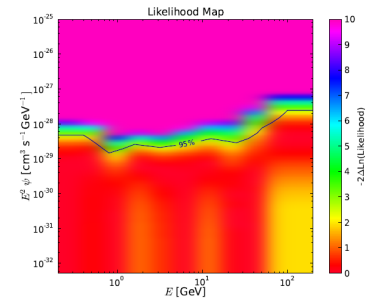
1. Halo models
2. DM model information
3. Astrophysics sources
4. Source locations
5. Background

FermiTools

Likelihoods

Too much CPU time consuming to do particle model fitting

Fitting DM gamma rays by using LikeDM



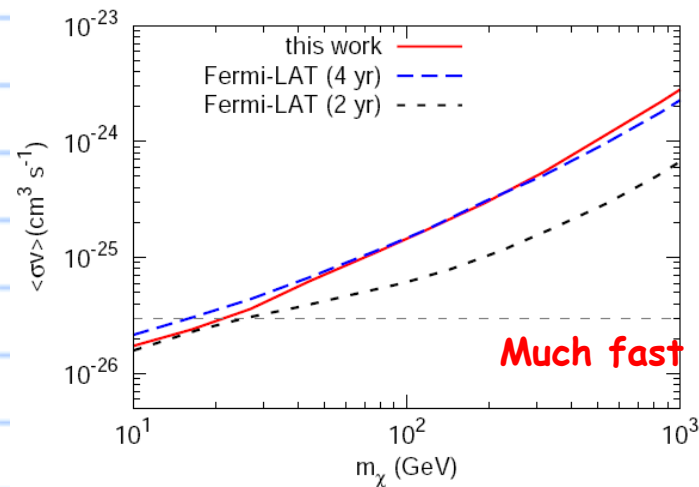
1. Halo models
2. Astrophysics sources
3. Source locations
4. Background

FermiTools

Energy-Residual
likelihood map

Likelihoods

Y-L Sming Tsai, Qiang Yuan, Xiaoyuan Huang
Published in JCAP 1303 (2013) 018
e-Print: arXiv:1212.3990



Much fast to do particle model fitting

DM particle model information

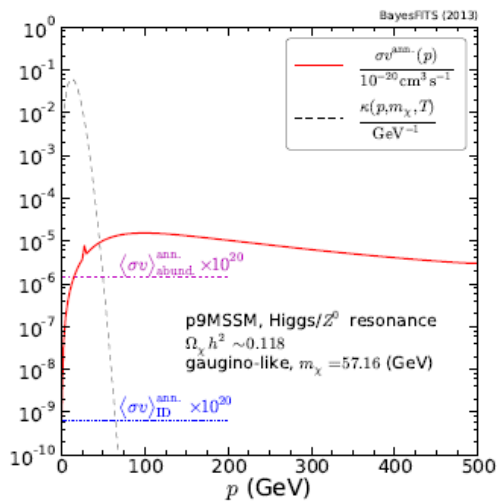
Summary

1. A global Statistical Analysis of the MSSM with a minimum set of 9 parameters.
2. Better numerical simulation of LHC, XENON100, and Fermi gamma-ray likelihoods.
3. Theoretical error strongly limits power of XENON100 on MSSM (Not so for XENONIT)
4. With all the constraints, especial $g-2$ and LHC, the limit $200 \text{ GeV} < m_{\tilde{\chi}} < 500 \text{ GeV}$.
5. Interesting prospects for indirect detection
6. Our method is completely general. It can be applied to other models (SUSY or not).

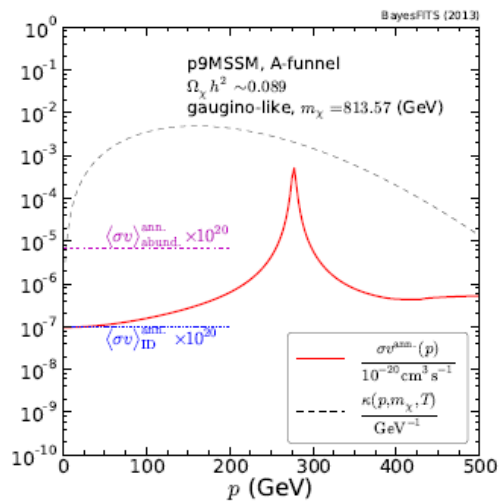


**Thank you very much
for your attention.**

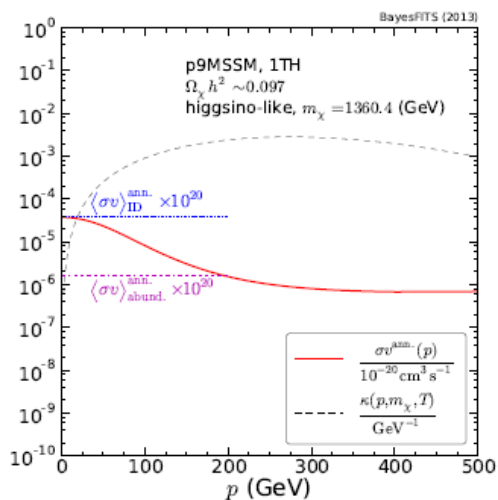
Backup



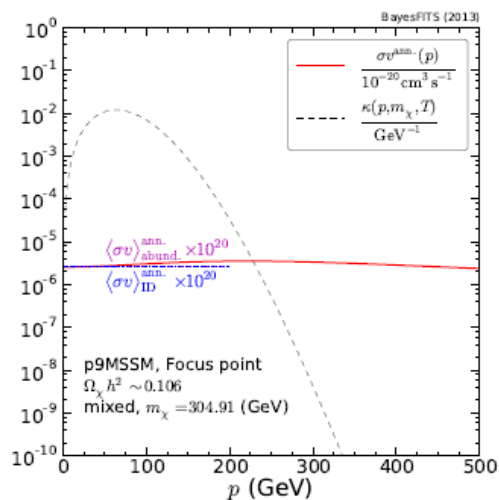
(a)



(b)



(c)



(d)

$$\langle \sigma v \rangle_{\text{abund}}^{\text{ann+co}} \sim \int \sigma v^{\text{ann+co}}(p) \kappa(p, m_\chi, T) dp.$$

Boltzmann distribution