GTM TK

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Isolating blocks

An *isolating block* is a special isolating neighborhood such that every point on its boundary leaves the neighborhood immediately in one or another time direction.



Useful Lemmas

- An isolating block always exists
- For an isolating block N, define its **exit set** n^- to be the set of points on the boundary that leave N immediately.

$$n^{-} := \{ x \in \partial N \, | \, \exists \, \delta_0 > 0 \text{ such that } x \cdot (0, \delta_0) \cap N = \emptyset \},\$$

- $\bullet\,$ Then, (N,n^-) is an index pair and $\mathcal{I}(N)=[N/n^-]$
- If Y is contractible, then Y/Z is homotopy equivalent to the unreduced suspension SZ of Z.

Tools

Example

Consider a flow on \mathbb{R}^2 given by $(x, y) \cdot t = (2^{-t}x, 2^t y)$. The disk around the origin (0, 0), which is a hyperbolic fixed point, is an isolating block.





Computation

Homology

Example (Continued)



The Conley index is (homotopy type of) the circle S^1 .

Linear flow

Consider a flow given by a bounded, self-adjoint, linear L with no kernel on a Hilbert space.

- Let X be a unit ball in H. Then X is an isolating block.
- Pick a finite-dimensional eigenspace V of L. Similarly, the Conley index $\mathcal{I}(X \cap V)$ is S^{V^-} the sphere with dimension equal to the number of negative eigenvalues in V.

Stable Conley index

It turns out that we can define E(X) as

$$\Sigma^{-V^-}\mathcal{I}(X\cap V)\cong\Sigma^{-W^-}\mathcal{I}(X\cap W)\cong S^0$$

Remark

- Conley theory can be considered as a generalization of Morse theory.
- If a group G acts on M and the flow is G-equivariant, then we can define an equivariant version of Conley indices.
- Seiberg-Witten Floer theory has U(1)-symmetry.

2. Seiberg-Witten Theory

Reference

Kronheimer and Mrowka, Monopoles and three-manifolds.

Let Y be a closed, oriented, Riemannian 3-manifold equipped with a spin^c structure \mathfrak{s} .

• This also means Y is given a unitary rank-two "spinor" bundle $S \to Y$ with Clifford multiplication

 $\rho: TY \to Hom(S,S)$

Computation

Homology

Configuration space

Definition

- The configuration space $C(Y, \mathfrak{s}) := \mathcal{A}(Y) \oplus \Gamma(S)$ consists of pairs of a spin^c connection and a section of the spinor.
- The gauge group $\mathcal{G} := \operatorname{Map}(Y, S^1)$ acts on $\mathcal{C}(Y, \mathfrak{s})$ by

$$u \cdot (B, \Psi) = (B - u^{-1}du, u\Psi).$$

• The quotient configuration space is $\mathcal{B}(Y, \mathfrak{s}) := \mathcal{C}(Y, \mathfrak{s})/\mathcal{G}$

Seiberg-Witten vector field

• The Chern-Simons-Dirac functional is given by

$$CSD(B,\Psi) = -\frac{1}{8} \int_Y (B^t - B_0^t) \wedge (F_{B^t} + F_{B_0^t}) + \frac{1}{2} \int_Y \langle D_B \Psi, \Psi \rangle$$

 $\bullet\,$ The Seiberg-Witten vector field is the gradient of the CSD

$$\nabla CSD(B,\Psi) = (\frac{1}{2} * F_{B^t} + \tau(\Psi), D_B\Psi)$$

- Roughly speaking, monopole Floer homology is a semi-infinite S¹-equivariant Morse homology of B(Y, s) using CSD as the Morse function.
- The equation $\nabla CSD = 0$ is the (3D) Seiberg-Witten equation. Critical points are moduli space of the solutions (and they are compact!).

The Coulomb slice

In general, $\mathcal{B}(Y, \mathfrak{s})$ is not a vector space, but a Hilbert bundle over the Picard torus

$$\mathbb{I}^{b_1} = H^1(Y; i\mathbb{R})/(2\pi i H^1(Y; \mathbb{Z})).$$

To apply finite dimensional approximation, we instead consider

Definition

The Coulomb slice $\mathcal{K} := \ker d^* \oplus \Gamma(S) \subset \Omega^1(Y; i\mathbb{R}) \oplus \Gamma(S)$

• This is a universal cover

F

$$\mathbb{Z}^{b_1} \to \mathcal{K} \to \mathcal{B}(Y,\mathfrak{s}).$$

Outline of the construction of SWF

- \mathcal{K} is a vector space, but we lose compactness.
- We can consider a certain collection \mathcal{R}_k of enlarging closed and bounded subsets of \mathcal{K} .
- On each \mathcal{R}_k , we can find its stable Conley index $E(\mathcal{R}_k)$.
- Finally, we take limit

$$SWF(Y) := \underset{\longrightarrow}{\lim} E(\mathcal{R}_k).$$

Previous works

- (1992) Cohen, Jones, and Segal started studying about existence of "Floer homotopy type".
- (1999) The stable Conley index on Hilbert spaces was developed by Geba, Izydorek, Pruszko
- (2003) In the Seiberg-Witten case, Manolescu successfully constructed Floer homotopy type for 3-manifolds with $b_1 = 0$.
- (2003) Kronheimer and Manolescu extended the construction to 3-manifolds with $b_1 = 1$ with nontorsion spin^c structure.

3. Computation

- We will give an explicit computation for when $Y = T^3$ with a trivial spin^c structure.
- The computation applies to 3-manifolds with nonnegative scalar curvature.
- Using Hodge decomposition, we have ker $d^* = \Omega_h \oplus \Omega_{\perp}$, where Ω_h is a space of harmonic 1-forms isomorphic to $H^1(Y; i\mathbb{R}) \cong \mathbb{R}^{b_1}$.

Tubular neighborhoods as isolating neighborhoods

- $\bullet\,$ In this case, all critical points and flow lines lie on Ω_h
- The subset we consider will be a tubular neighborhood $\nu(\epsilon)$ of a subset $I_k \subset \Omega_h \cong \mathbb{R}^3$, i.e. $\mathcal{R}_k = I_k \times B(\Omega_\perp \oplus \Gamma(S), \epsilon)$



Trick: linearizing a vector field

- We can freely scale the radius of tubular neighborhoods. This, in turn, scales the vector field.
- We also use the fact that the (perturbed) vector field can be written as L + Q where L is linear and Q is quadratic.

$$\begin{array}{cccc} \text{Isolating neighborhood} : & \nu(\epsilon) & \xrightarrow{\text{Scaling}} & \nu(1) \\ & \text{Vector field} : & L(h)v + \pi Qv & \xrightarrow{\text{Scaling}} & L(h)v + \epsilon \pi Qv \\ & \text{Approximating} & & & \downarrow \\ & L(h)v + Qv & & L(h)v \end{array}$$

The linearized vector field

• The linearized vector field on $\Omega_h \times \Gamma(S)$ has a form

$$(h,\psi)\mapsto (\epsilon\nabla f, (D_{B_0+b_h}-\delta)\psi),$$

where δ, ϵ are sufficiently small positive numbers (the choice of ϵ depends on δ) and f is a Morse function on Ω_h .

• We pick $f(\theta_1, \theta_2, \theta_3) = -\cos \theta_1 - \cos \theta_2 - \cos \theta_3$.

The flow on Ω_h

- ∇f induces a flow on Ω_h .
- For subsets on Ω_h , we will consider the cube $I_k = [-(2k + \frac{1}{2})\pi, (2k + \frac{1}{2})\pi]^3$, which is an isolating block with no exit set.



A tubular neighborhood as an isolating block.

Claim

The tubular neighborhood $\nu(I_k)$ is an isolating block.



- One can check directly that the norm $||v(t)||^2$ has no local maximum using the second derivative test.
- Consequently, we only need to understand the exit set.

Homology

Describing the exit set

From the first derivative of the norm $||v(t)||^2$, we can deduce

Exit set of the tubular neighborhood

The exit set of $\nu(I_k)$ is given by the set

 $\{(h,v)\in I_k\times S(1)\,|\,\langle L(h)v,v\rangle\leq 0\}.$



• We can view this as a union of unit nonpositive cones $\operatorname{Cone}_{\leq}(V, L(h)) = \{v \in S(1) \mid \langle L(h)v, v \rangle \leq 0\}$ along fibers.

The spinor bundle

- The spinor bundle is the trivial bundle $T^3 \times \mathbb{C}^2$.
- Using Fourier transform, we have a decomposition

$$\Gamma(S) = \bigoplus_{n_1, n_2, n_3 \in \mathbb{Z}} V_{n_1, n_2, n_3},$$

each V_{n_1,n_2,n_3} is a 2-dimensional complex vector space spanned by $(e^{i(n_1\theta_1+n_2\theta_2+n_3\theta_3)}, 0)$ and $(0, e^{i(n_1\theta_1+n_2\theta_2+n_3\theta_3)})$.

• Note that the S^1 -action comes from multiplication on $\Gamma(S)$ by unit complex numbers.

The Dirac operator

• For $\vec{b} = (b_1, b_2, b_3) \in \mathbb{R}^3$, the operator $D_{\vec{b}}$ on V_{n_1, n_2, n_3} is given by a matrix

$$\begin{bmatrix} -(n_1 + \frac{b_1}{2\pi}) - \delta & -(n_3 + \frac{b_3}{2\pi}) - (n_2 + \frac{b_2}{2\pi})i \\ -(n_3 + \frac{b_3}{2\pi}) + (n_2 + \frac{b_2}{2\pi})i & n_1 + \frac{b_1}{2\pi} - \delta \end{bmatrix}$$

- The eigenvalues of the above matrix is $-\delta \pm \sqrt{(n_1 + \frac{b_1}{2\pi})^2 + (n_2 + \frac{b_2}{2\pi})^2 + (n_3 + \frac{b_3}{2\pi})^2}$ with eigenvectors $\begin{bmatrix} -b_1 + R\\ -b_3 + b_2i \end{bmatrix}$ and $\begin{bmatrix} -b_1 - R\\ -b_3 + b_2i \end{bmatrix}$ respectively.
- The gauge symmetry can be seen as the matrix of $D_{\vec{b}+2\pi\vec{u}}$ on $V_{\vec{n}}$ is the same as the matrix of $D_{\vec{b}}$ on $V_{\vec{n}+\vec{u}}$.

Computation

Homology

The kernel of $D_{\vec{b}}$

• $D_{\vec{b}}$ has kernel on a small sphere centered at centered at $2\pi(b_1, b_2, b_3)$, which is a critical point of index 0.



Warm-up (1-dim down)

- Consider a tubular neighborhood of a square $\left[-\frac{1}{2}, 2\pi + \frac{1}{2}\right]^2 \times \{0\}$ with respect to a subspace $V = \bigoplus_{i,j=0,1} V_{-i,-j,0}.$
- Look at signature of Dirac operators on V.



Conclusion

• After desuspension, the stable Conley index $E(\nu(I_k))$ is an "S⁰-sum" of $(2k+1)^3$ copies of S².



• Equivariantly, the 2-sphere is $S^{\mathbb{C}}$

Induced map on Conley indices

• The inclusion $\nu(I_k) \subset \nu(I_{k+1})$ induces a map between Conley indices as well as their associated spectra.

$$E(\nu(I_k)) \rightarrow E(\nu(I_{k+1})).$$

 The inclusion ν(I_k) ⊂ ν(I_{k+1}) also induces an inclusion of their exit sets

The colimit of spectra

 $SWF(T^3)$ is the S^0 -sum of \mathbb{Z}^3 copies of $S^{\mathbb{C}}$.

4. Homology

• Monopole Floer homology comes with three flavors $\overrightarrow{HM}, \overrightarrow{HM}, \overrightarrow{HM}$ related by a long exact sequence

$$\ldots \to \overline{HM}(Y) \to \widetilde{HM}(Y) \to \widehat{HM}(Y) \to \ldots$$

• There are three S^1 -equivariant homology theories: the Borel homology $H_*^{S^1}$, the coBorel homology $\hat{H}_*^{S^1}$, and the Tate homology $\bar{H}_*^{S^1}$.

Conjecture

These S^1 -equivariant homology groups of SWF(Y) agree with the monopole Floer groups of Y. Tools

Example

- As $SWF(S^3) = S^0$, the groups agree.
- But for T^3

	-4	-3	-2	-1	0	1	2	3	4
$\overline{HM}_*(T^3)$	\mathbb{Z}^3	\mathbb{Z}^3	\mathbb{Z}^3	\mathbb{Z}^3	\mathbb{Z}^3	\mathbb{Z}^3	\mathbb{Z}^3	\mathbb{Z}^3	\mathbb{Z}^3
$\widecheck{HM}_*(T^3)$	0	0	0	0	0	\mathbb{Z}^3	\mathbb{Z}^3	\mathbb{Z}^3	\mathbb{Z}^3
$\widehat{HM}_*(T^3)$	\mathbb{Z}^3	\mathbb{Z}^3	\mathbb{Z}^3	\mathbb{Z}^3	\mathbb{Z}^3	\mathbb{Z}^3	0	0	0
$\bar{H}_*^{S^1}(SWF(T^3))$	\mathbb{Z}	0	\mathbb{Z}	0	\mathbb{Z}	0	\mathbb{Z}	0	\mathbb{Z}
$H_*^{S^1}(SWF(T^3))$	0	0	0	0	0	\mathbb{Z}^{N-1}	\mathbb{Z}	0	\mathbb{Z}
$\widehat{H}_*^{S^1}(SWF(T^3))$	0	\mathbb{Z}	0	\mathbb{Z}	0	\mathbb{Z}^N	0	0	0
where $N = \mathbb{Z}^3$									

• This is because there is an action by \mathbb{Z}^3 left.

Twisted parametrized spectrum

The construction is based on a concept of twisted parametrized spectra introduced by Douglas.

- Roughly speaking, a twisted parametrized spectrum is a bundle of spectra twisted by automorphisms of the category of spectra, e.g. suspensions.
- One simple way to describe this is to use open cover and transition functions.

Twisted Manolescu-Floer spectrum

• There is a isomorphism between Conley indices induced by the action of a harmonic gauge transformation u

$$\mathcal{I}(\mathcal{R} \cap V) \to \ \mathcal{I}((u \cdot \mathcal{R}) \cap (u \, V)).$$

• However, when passing to spectra, we have

$$S^{V^-} \wedge SWF(Y) \to S^{(uV)^-} \wedge SWF(Y),$$

Twisted Manolescu-Floer spectrum

We define a twisted Manolescu-Floer spectrum $\widetilde{SWF}(Y)$ as a twisted parametrized spectrum over the Picard torus.



• We use the standard cover of the torus and the transition functions g_i is the map given by u_i earlier.

The manifold $S^1 \times S^2$

- Consider $S^1 \times S^2$ with torsion spin^c structure.
- We can find that $SWF(S^1 \times S^2, \mathfrak{s}) \simeq S^0$ using the same method as T^3 .
- Because the Dirac operator has no kernel, it turns out that $\widetilde{SWF}(S^1 \times S^2, \mathfrak{s}) \simeq S^1 \times SWF(S^1 \times S^2, \mathfrak{s}).$

Homology



There is a spectral sequence for Tate homology.

