Core-Collapse Supernova Explosions: Some Observables and Diagnostics

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Diagnostics of Explosion

- Morphology of Blasts: Spatial Element/Debris Distribution
- Pulsar Proper Motions ("Kicks")
- Pulsar Spin Periods
- Neutrinos
- Gravitational Waves
- Nucleosynthesis R-process ?

Cas A Remnant



Fesen/Milisavljevic et al. 2011

Cas A Remnant



Fesen/Milisavljevic et al. 2011

Cas A Remnant



Fesen/Milisavljevic et al. 2011

NuSTAR will map the remnants of recent supernova explosions, testing theories of where the elements are born





NuSTAR will measure and map the ⁴⁴Ti lines at 68 and 78 keV in historic remnants: Tycho, Kepler, Cas A and SN1987A





Pulsar Kicks

 Hydrodynamic Mechanism Seems Natural - "Simple" Recoil

Net Neutrino momentum component may be small
 No need for exotic mechanisms (Harrison-Tademaru, parity violation, B-fields,)

Pulsar Recoil: A Generic Feature

Pulsar Kicks: Pulsar B2224+65 and Bow Shock

Cordes, Romani, Lundgren '93

Guitar Nebula



Nordhaus, Burrows, & Ott 2010

Acceleration ~ 500 km/

See also pioneering work of Scheck et al., for an extensive suite of related simulations



Nordhaus, Brandt, & Burrows 2012





Induced Rotation?

Found and Pioneered by J. Blondin et al.







Integrated Angular Momentum from 3D simulations: Induced Spins







CASTRO 3D AMR Neutrino-Driven Explosion Model: Induced Spin (1.7 vs. 1.9)?



Neutrino Bursts/Signatures



Core-Collapse Neutrinos Detected



Kamioka II



Breakout Burst of Neutrinos: Precision Boltzmann Transfer







Multi-Angle, Multigroup, Time-Dependent Transport in 2D SN Sigheletrogh

Modifications to the Solver

[Ott et al. 2008, arxiv:0804.239, ApJ 682, 1277, 2008]

S_n – MGFLD hybrid scheme:

- S_n solver in VULCAN converges slowly at high optical depth.
 time step limitation.
- Idea: Use MGFLD at high optical depths and transition to Sn at intermediate optical depth.
- Set up boundary data using to Eddington approximation:

$$I(\vec{n}) = J_{\text{MGFLD}} + 3(\vec{n} \cdot \vec{H}_{\text{MGFLD}})$$



 Efficient at high optical depths and accurate in semi-transparent regions.



- s20.nr: Little difference between MGFLD and S_n at 160 ms after bounce.
- s20. π : Large (factor ~3) polar differences in specific heating rates.

• (only $\approx 2\%$ difference; S_n gain < MGFLD gain!)

Gravitational Radiation from Supernovae

Ott et al. (2004,2006)















A Boltzmann Formalism for Oscillating Neutrinos

P. Strack and A. Burrows Phys.Rev.D 71:093004,2005 (hep-ph/0504035) & hep-ph/0505056

Y. Zhang & A. Burrows 2013

I. Quasi-classical Boltzmann equations

Quantum analog to the classical phase-space density

$$\rho(\mathbf{r}, \mathbf{p}, t) = \int \frac{d^3 \mathbf{p}'}{(2\pi)^3} e^{-i\mathbf{p}'\mathbf{r}} a^{\dagger}(\mathbf{p} - \frac{1}{2}\mathbf{p}', t) a(\mathbf{p} + \frac{1}{2}\mathbf{p}', t)$$

Make operators to numbers

 $\mathcal{F} = \langle n_i | \rho | n_j \rangle = \begin{pmatrix} f_{\nu_e} & f_{e\mu} \\ f_{e\mu}^* & f_{\nu_{\mu}} \end{pmatrix}$

Diagonal elements: real numbers: phasespace densities

Off-diagonal elements: complex numbers: macroscopic overlap

$$f_{e\mu} = \langle n_{\nu_e} | \rho(\mathbf{r}, \mathbf{p}, t) | n_{\nu_{\mu}} \rangle$$

=
$$\int \frac{d^3 \mathbf{p}'}{(2\pi)^3} e^{i\mathbf{p'r}} \langle n_{\nu_e} | a^{\dagger}(\mathbf{p} - \frac{1}{2}\mathbf{p'}, t) a(\mathbf{p} + \frac{1}{2}\mathbf{p'}, t) | n_{\nu_{\mu}} \rangle$$

II. Quasi-classical Boltzmann equations

Merging the Boltzmann equation with the Heisenberg equation yields

$$\frac{\partial \mathcal{F}}{\partial t} + \frac{1}{2} \left\{ \mathbf{v}, \frac{\partial \mathcal{F}}{\partial \mathbf{r}} \right\} + \frac{1}{2} \left\{ \dot{\mathbf{p}}, \frac{\partial \mathcal{F}}{\partial \mathbf{p}} \right\} = -i \left[\Omega, \mathcal{F} \right] + C$$

Flavor Oscillation physics is implemented via

$$\Omega(\mathbf{r}, \mathbf{p}, \varepsilon, t) = \Omega_{\text{vac}}(\varepsilon) + \Omega_{\text{mat}}(\mathbf{r}) + \Omega_{\nu\nu}(\mathbf{r}, \mathbf{p}, t) - \tilde{\Omega}_{\tilde{\nu}\nu}(\mathbf{r}, \mathbf{p}, t)$$

$$\Omega(\varepsilon, \mathbf{r}) = \frac{\pi c}{L} \begin{pmatrix} -\cos 2\theta + 2A & \sin 2\theta \\ \sin 2\theta & \cos 2\theta \end{pmatrix}$$
Nonlinearities through self-interactions.
$$L = \frac{4\pi\hbar c\varepsilon}{\Delta m^2 c^4} \qquad \mathbf{\&} \qquad A = \left(\frac{L}{\pi c}\right) \frac{2\sqrt{2}G_F}{\hbar} n_e(\mathbf{r})$$

III. Quasi-classical Boltzmann equations

Separation of scales (small de Broglie wavelength of the v's when compared to density gradients & no density fluctuations)
 neglect off-diagonal Liouville terms

New "customized" variables Macroscopic <u>Overlap</u> densities:

$$f_r = \frac{1}{2} \left(f_{e\mu} + f_{e\mu}^* \right)$$
$$f_i = \frac{1}{2i} \left(f_{e\mu} - f_{e\mu}^* \right)$$

Self-interactions

$$\begin{aligned} \frac{\partial f_{v_e}}{\partial t} + \mathbf{v} \cdot \frac{\partial f_{v_e}}{\partial \mathbf{r}} + \dot{\mathbf{p}} \cdot \frac{\partial f_{v_e}}{\partial \mathbf{p}} &= -f_i \left(\frac{2\pi c}{L} \sin 2\theta + 2\beta \int (1 - \cos \theta^{\mathbf{pq}}) \left(f_r + \tilde{f}_r \right) d^3 \mathbf{q} \right) + 2\beta f_r \int (1 - \cos \theta^{\mathbf{pq}}) \left(f_i + \tilde{f}_i \right) d^3 \mathbf{q} + C_{v_e} \\ \frac{\partial f_{v_\mu}}{\partial t} + \mathbf{v} \cdot \frac{\partial f_{v_\mu}}{\partial \mathbf{r}} + \dot{\mathbf{p}} \cdot \frac{\partial f_{v_\mu}}{\partial \mathbf{p}} &= f_i \left(\frac{2\pi c}{L} \sin 2\theta + 2\beta \int (1 - \cos \theta^{\mathbf{pq}}) \left(f_r + \tilde{f}_r \right) d^3 \mathbf{q} \right) - 2\beta f_r \int (1 - \cos \theta^{\mathbf{pq}}) \left(f_i + \tilde{f}_i \right) d^3 \mathbf{q} + C_{v_\mu} \\ \frac{\partial f_r}{\partial t} + \mathbf{v} \cdot \frac{\partial f_r}{\partial \mathbf{r}} + \dot{\mathbf{p}} \cdot \frac{\partial f_r}{\partial \mathbf{p}} &= f_i \left[\frac{2\pi c}{L} \left(A - \cos 2\theta \right) + \beta \int (1 - \cos^{\mathbf{pq}}) \left(f_{v_e} - \tilde{f}_{v_e} - f_{v_\mu} + \tilde{f}_{v_\mu} \right) d^3 \mathbf{q} \right] + \\ &\qquad \left(f_{v_e} - f_{v_\mu} \right) \beta \int (1 - \cos \theta^{\mathbf{pq}}) \left(\tilde{f}_i - f_i \right) d^3 \mathbf{q} \\ \frac{\partial f_i}{\partial t} + \mathbf{v} \cdot \frac{\partial f_i}{\partial \mathbf{r}} + \dot{\mathbf{p}} \cdot \frac{\partial f_i}{\partial \mathbf{p}} &= \left(f_{v_e} - f_{v_\mu} \right) \left[\frac{\pi c}{L} \sin 2\theta + \beta \int (1 - \cos^{\mathbf{pq}}) \left(f_r - \tilde{f}_r \right) d^3 \mathbf{q} \right] - f_r \left[\frac{2\pi c}{L} \left(A - \cos 2\theta \right) + \beta \int (1 - \cos^{\mathbf{pq}}) \left(f_r - \tilde{f}_r \right) d^3 \mathbf{q} \right] \\ + \beta \int (1 - \cos^{\mathbf{pq}}) \left(f_{v_e} - \tilde{f}_{v_e} - f_{v_\mu} + \tilde{f}_{v_\mu} \right) d^3 \mathbf{q} \right] \end{aligned}$$

Quasi-classical Boltzmann Wigner density matrix, ensemble-averaging

$$\mathcal{F} = \langle n_i | \rho | n_j \rangle = \begin{pmatrix} f_{\nu_e} & f_{e\mu} \\ f_{e\mu}^* & f_{\nu_{\mu}} \end{pmatrix}$$

Diagonal elements: real numbers: Phase-Space densities Off-diagonal elements: complex numbers: Macroscopic Overlap

$$f_r = \frac{1}{2} \left(f_{e\mu} + f_{e\mu}^* \right)$$
$$f_i = \frac{1}{2i} \left(f_{e\mu} - f_{e\mu}^* \right)$$

(Real part)

(Imaginary part)

$$L = \frac{4\pi\hbar c\varepsilon}{\Delta m^2 c^4}$$

$$A = \left(\frac{L}{\pi c}\right) \frac{2\sqrt{2}G_F}{\hbar} n_e(\mathbf{r})$$

$$\frac{\partial f_{\nu_e}}{\partial t} + \mathbf{v} \cdot \frac{\partial f_{\nu_e}}{\partial \mathbf{r}} + \dot{\mathbf{p}} \cdot \frac{\partial f_{\nu_e}}{\partial \mathbf{p}} = -\frac{2\pi c}{L} f_i \sin 2\theta + C_{\nu_e}$$

$$\frac{\partial f_{\nu_{\mu}}}{\partial t} + \mathbf{v} \cdot \frac{\partial f_{\nu_{\mu}}}{\partial \mathbf{r}} + \dot{\mathbf{p}} \cdot \frac{\partial f_{\nu_{\mu}}}{\partial \mathbf{p}} = \frac{2\pi c}{L} f_i \sin 2\theta + C_{\nu_{\mu}}$$

$$\frac{\partial f_r}{\partial t} + \mathbf{v} \cdot \frac{\partial f_r}{\partial \mathbf{r}} + \dot{\mathbf{p}} \cdot \frac{\partial f_r}{\partial \mathbf{p}} = -\frac{2\pi c}{L} f_i (\cos 2\theta - A)$$

$$\frac{\partial f_i}{\partial t} + \mathbf{v} \cdot \frac{\partial f_i}{\partial \mathbf{r}} + \dot{\mathbf{p}} \cdot \frac{\partial f_i}{\partial \mathbf{p}} = \frac{2\pi c}{L} \left(\frac{f_{\nu_e} - f_{\nu_{\mu}}}{2} \sin 2\theta + f_r (\cos 2\theta - A)\right)$$

"R-process" Nucleosynthesis?

No natural site for 2nd or 3rd Peak in CCSNe (~ Solar Metallicity)?

R-Process Nucleosynthesis

Nucleosynthesis in the r-process



Ejecta mass versus Entropy and Y_e for Acoustic Mechanism: R-process?



M (s > 300): $1.25 \times 10^{-4} M$?

M (s > 100): 1.07×10^{-5} M?

Preliminary R-process Calculations for the Long-term Acoustic Mechanism



Summary: Diagnostics of Explosion

- Morphology of Blasts: Spatial Element/Debris Distribution
- Explosion Energies and Nickel Yields
- Pulsar Proper Motions ("Kicks")
- Pulsar Spin Periods
- Neutrinos !!
- Gravitational Waves
- Nucleosynthesis R-process ? The first peak?