

# B Physics: New Physics and The Next Generation

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Complex phases in the weak interaction:  $V_{td}$  and  $V_{ts}$  and associated CPV asymmetries

Excitement in Flavor Physics:

- Connections to the charged Higgs
- Rare B Decays + NP

Flavor Physics, The Next Generation:  
Belle II and the LHCb upgrade

Apologies: In the limited time, I cannot cover all the recent results from BaBar, Belle, LHCb, CMS, ATLAS, Tevatron ... I have borrowed slides from many excellent physicists and will aim for the “big picture” but skip most details.

# Amplitudes and Phases in the Weak Interaction

N. Cabibbo



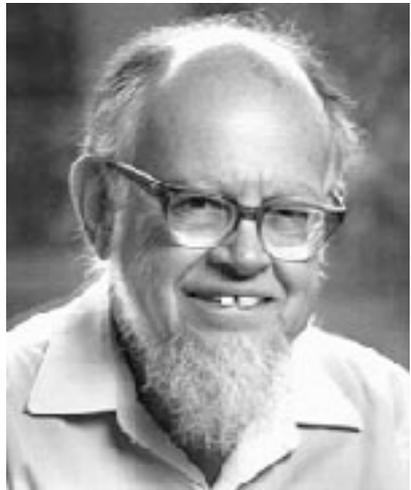
M. Kobayashi



T. Maskawa



$$V_{\text{CKM}} \equiv V_L^u V_L^{d\dagger} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ \underline{V_{td}} & V_{ts} & \underline{V_{tb}} \end{pmatrix}$$

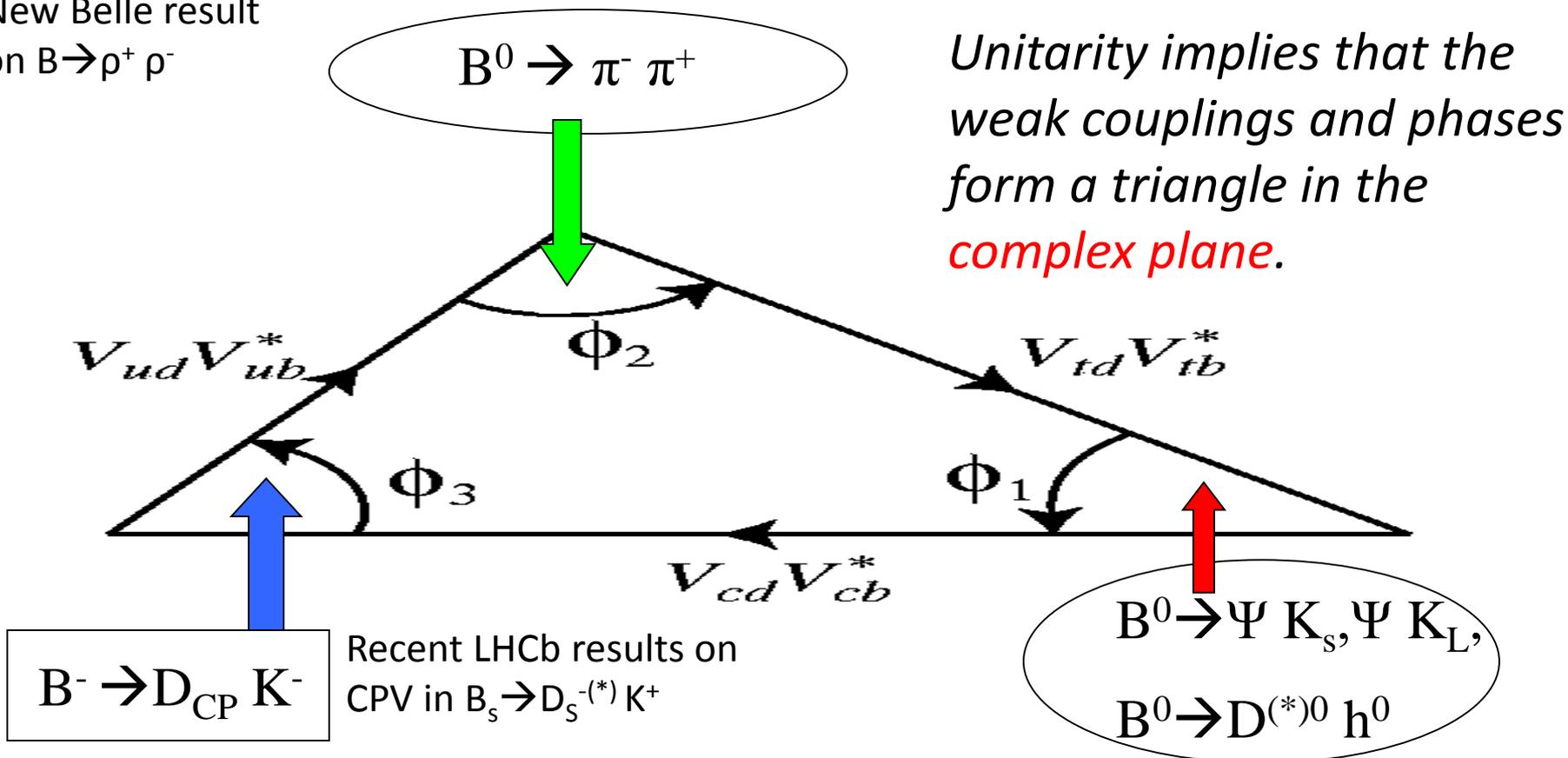


to  $O(\lambda^3)$

$$V = \begin{pmatrix} \begin{Bmatrix} 1 - \frac{1}{2}\lambda^2 & \lambda \\ -\lambda & 1 - \frac{1}{2}\lambda^2 \end{Bmatrix} & \underline{A\lambda^3(\rho - i\eta)} \\ \underline{A\lambda^3(1 - \rho - i\eta)} & \underline{-A\lambda^2} & 1 \end{pmatrix}$$

# Three Angles: $(\varphi_1, \varphi_2, \varphi_3)$ or $(\beta, \alpha, \gamma)$

New Belle result  
on  $B \rightarrow \rho^+ \rho^-$



Recent LHCb results on  
CPV in  $B_s \rightarrow D_s^{(*)} K^+$

$B^- \rightarrow D_{CP} K^-$

$B^0 \rightarrow \Psi K_s, \Psi K_L,$

$B^0 \rightarrow D^{(*)0} h^0$

Big Questions: *Are determinations of angles consistent with determinations of the sides of the triangle? Are angle determinations from **loop** and **tree** decays consistent?*

# Time-dependent $CP$ violation is “A Double-Slit experiment” with particles and antiparticles

## QM interference between two diagrams

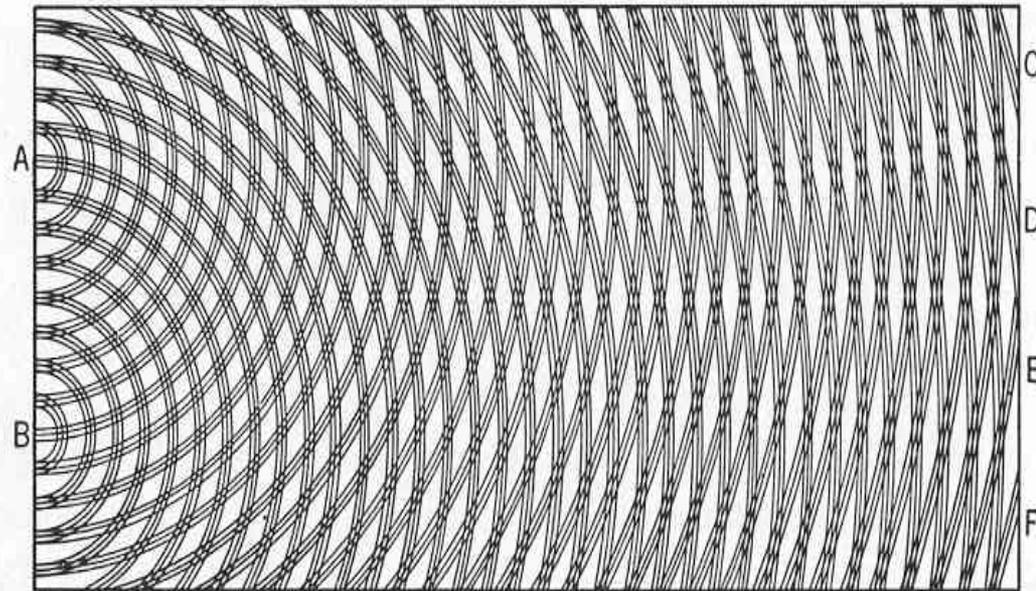
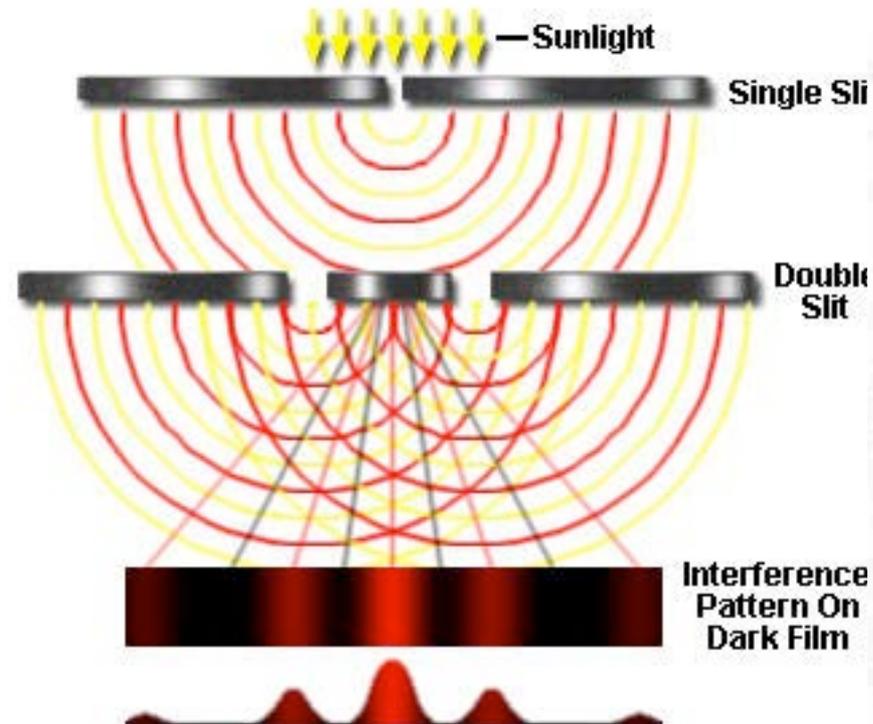
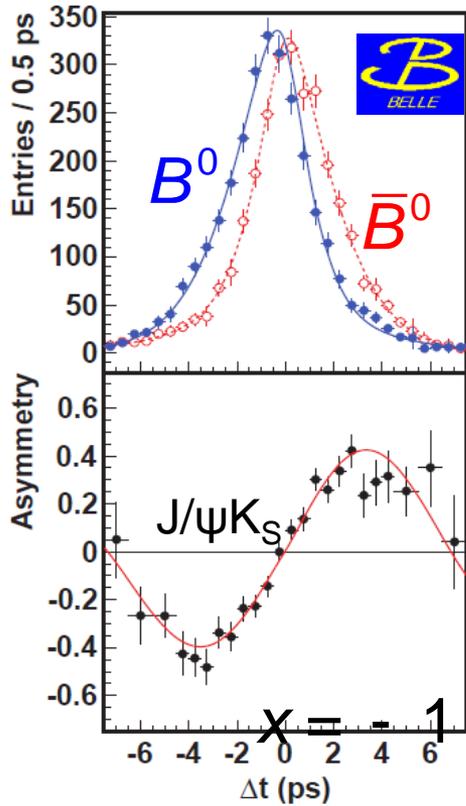


FIG. 1

Measures the phase of  $V_{td}$  or equivalently the phase of  $B_d$ -anti  $B_d$  mixing.

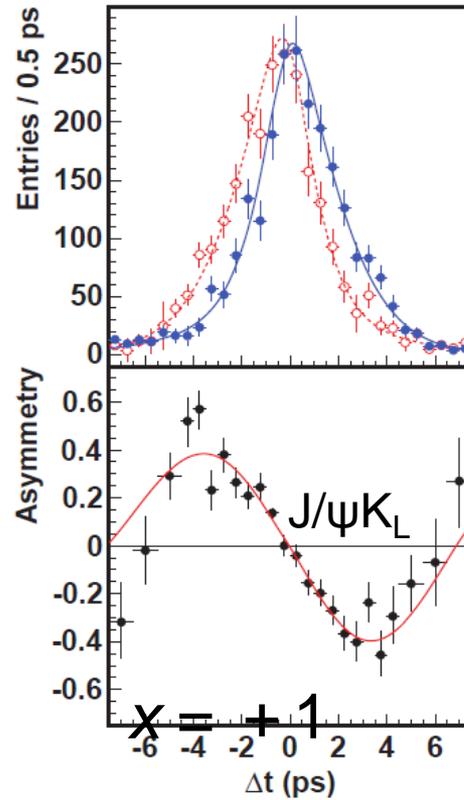
# Measurement of $\sin(2\phi_1)/\sin(2\beta)$ in $B \rightarrow \text{Charmonium } K^0$ modes



$$\sin 2\phi_1 = 0.667 \pm 0.023 \pm 0.012$$

$$A_f = 0.006 \pm 0.016 \pm 0.012$$

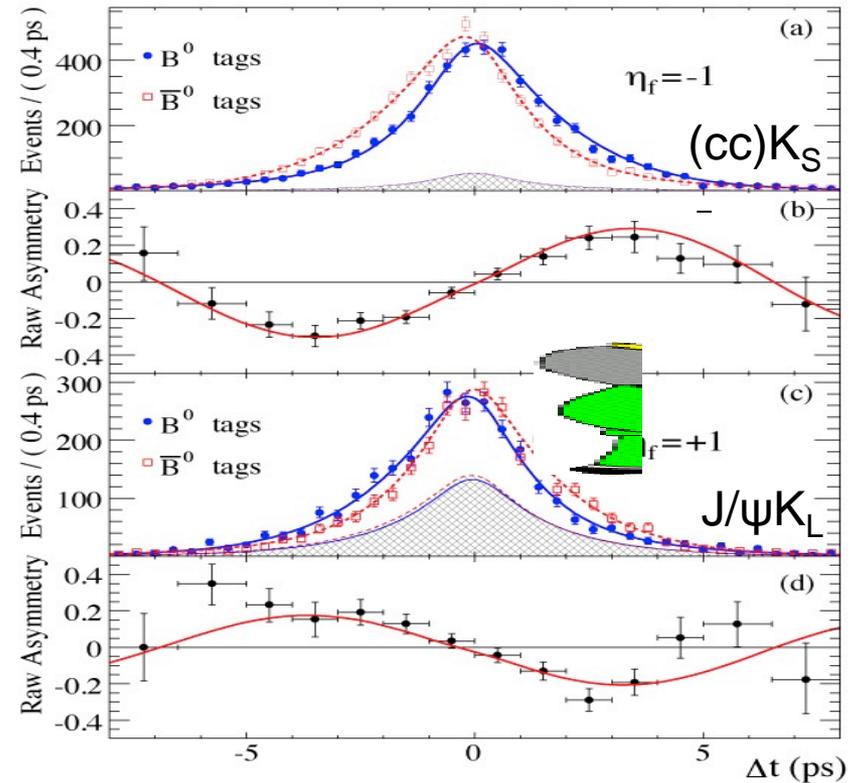
PRL108,171802 (2012)



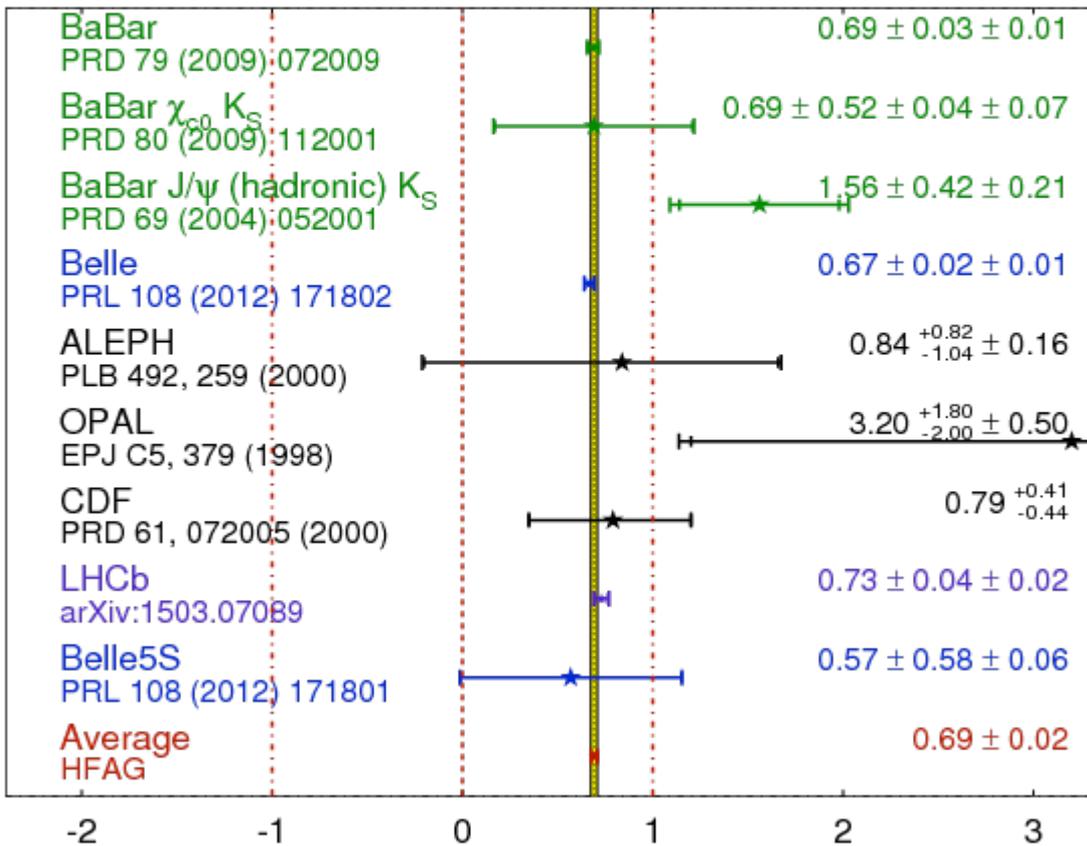
$$\sin 2\phi_1 = 0.687 \pm 0.028 \pm 0.012$$

$$A_f = -0.024 \pm 0.020 \pm 0.016$$

PRD79,072009 (2009)

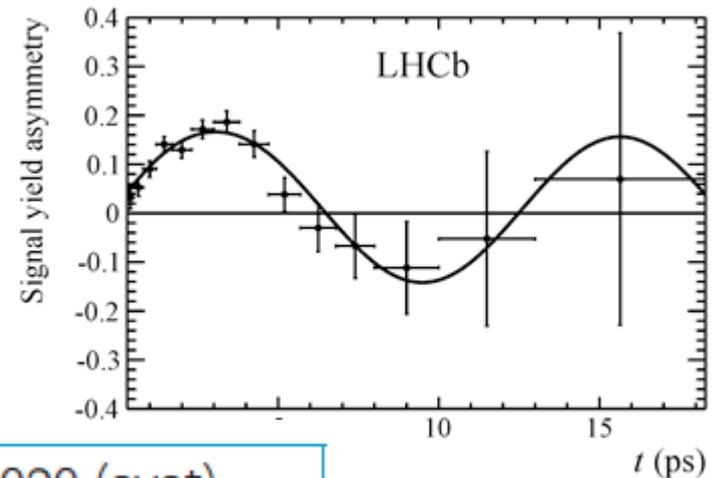


Overpowering evidence for CP violation (matter-antimatter asymmetries). >>>> The phase of  $V_{td}$  is in good agreement with Standard Model expectations. *This is the phase of  $B_d$  mixing.*



B factories: High precision CPV measurement and a calibration for NP.

2013: **LHCb** joins the game [ $\pm 0.07$ (stat)]  
 2015: Latest LHCb measurement is comparable in precision to B factories (now uses same-side B tagging)



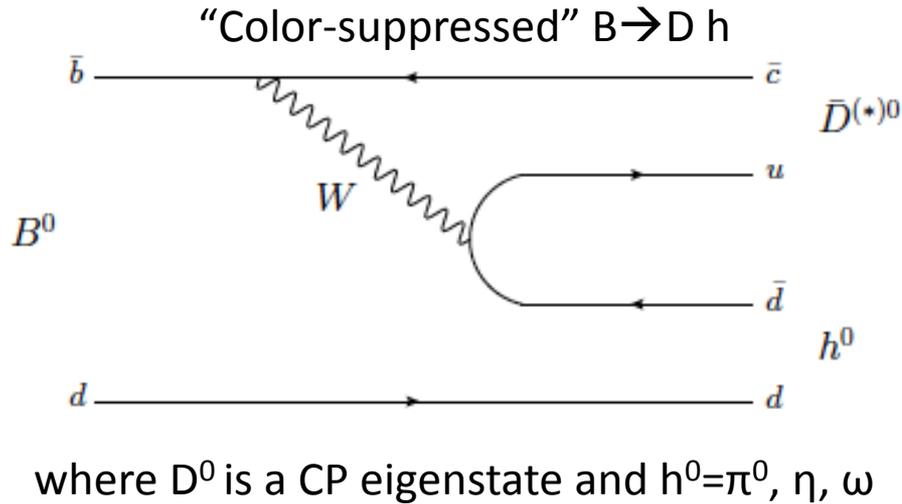
$$S(B^0 \rightarrow J/\psi K_S^0) = 0.731 \pm 0.035 \text{ (stat)} \pm 0.020 \text{ (syst)},$$

$$C(B^0 \rightarrow J/\psi K_S^0) = -0.038 \pm 0.032 \text{ (stat)} \pm 0.005 \text{ (syst)}.$$

# B factories: *Check CP violation in $b \rightarrow c$ [ $u\bar{b} d$ ] processes*

## 2015: First joint BaBar-Belle data analysis

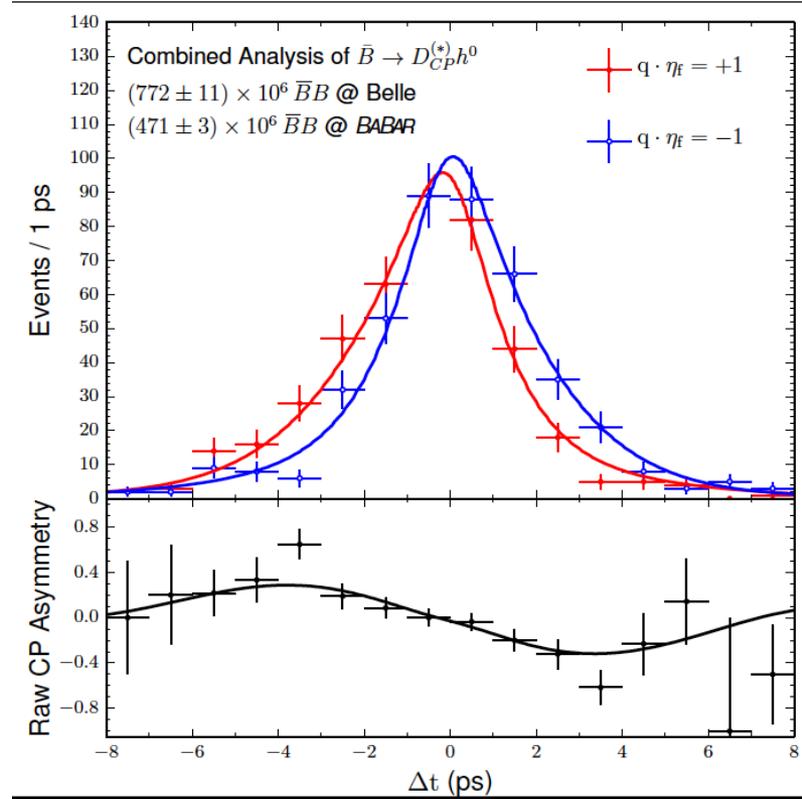
M. Rohrken et al



Combining Belle and BaBar datasets,  $\sim 1260$  signal events, obtain a  $5.4\sigma$  CP violation signal  $\rightarrow$  First observation  
 $\sin(2\beta_{\text{eff}}) = 0.66 \pm 0.10(\text{stat}) \pm 0.06(\text{sys})$

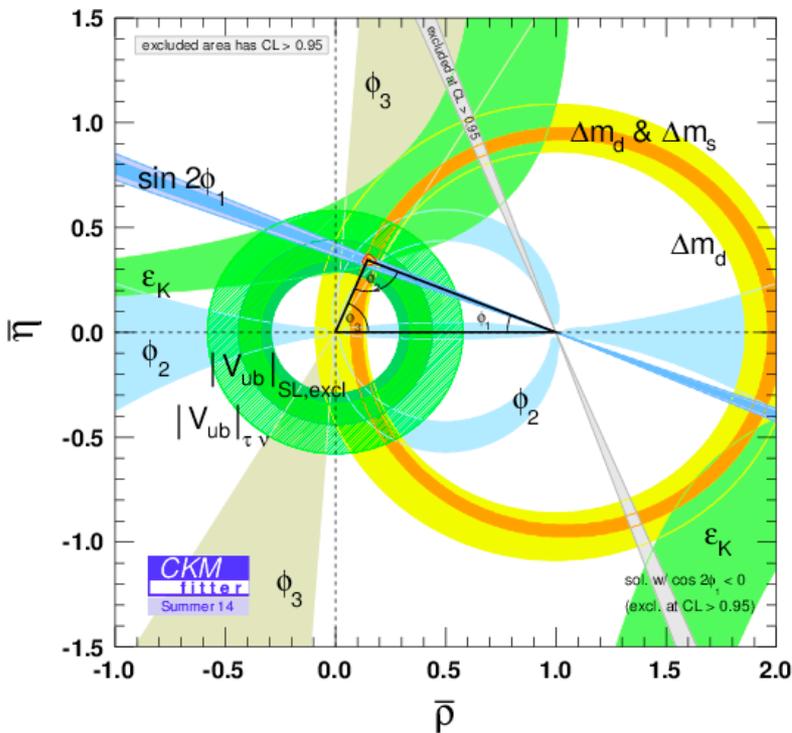
Phase of  $V_{td}$  again

Conclusion: CP violation in  $b \rightarrow c$   $u\bar{b} d$  modes is the same as in  $b \rightarrow c$   $c\bar{b} s$  modes (e.g.  $B \rightarrow J/\psi K_S$ )



# Results from Global Fits to Data (CKMFitter Group)

Great progress on  $\varphi_3$  or  $\gamma$  (first from B factories and now in the last two years from LHCb). These measure the phase of  $V_{ub}$  [CKM2014, K. Trabelsi's review:  $\pm 7^\circ$ ]

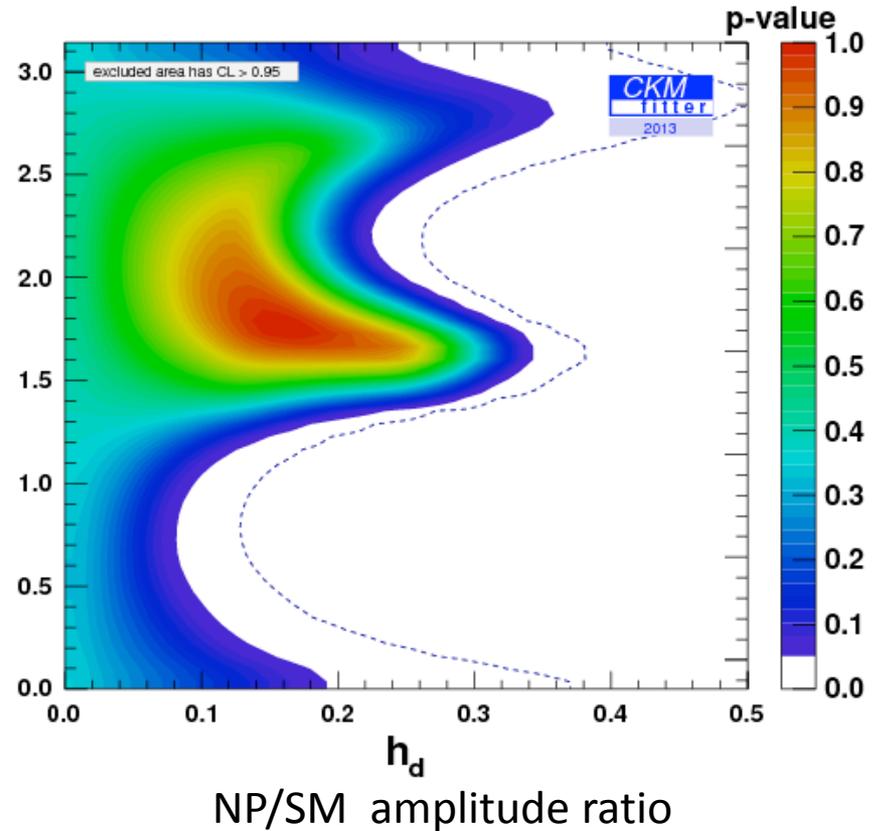


Looks good  
(except for an issue with  $|V_{ub}|$ )

Similar results from UTFIT as well from G. Eigen et al.

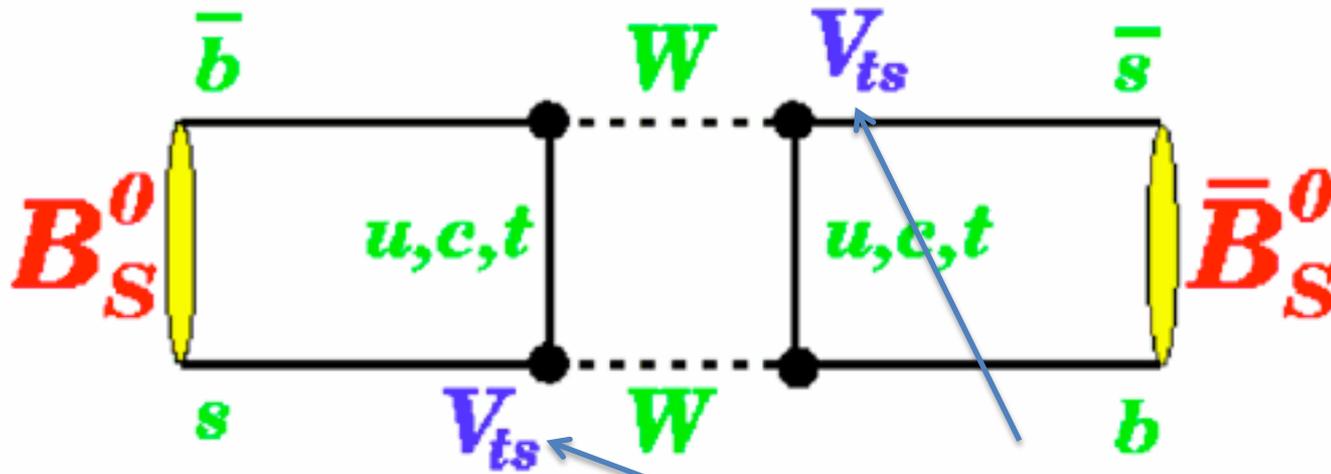
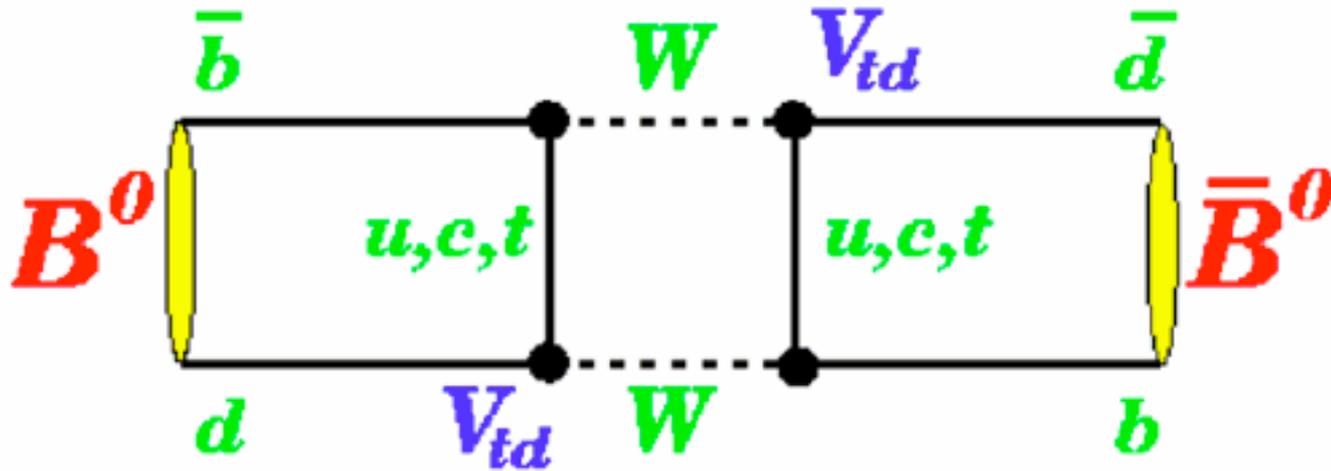
NP  
Phase

$\sigma_d$



But a 10-20% NP amplitude in  $B_d$  mixing is perfectly compatible with all current data.

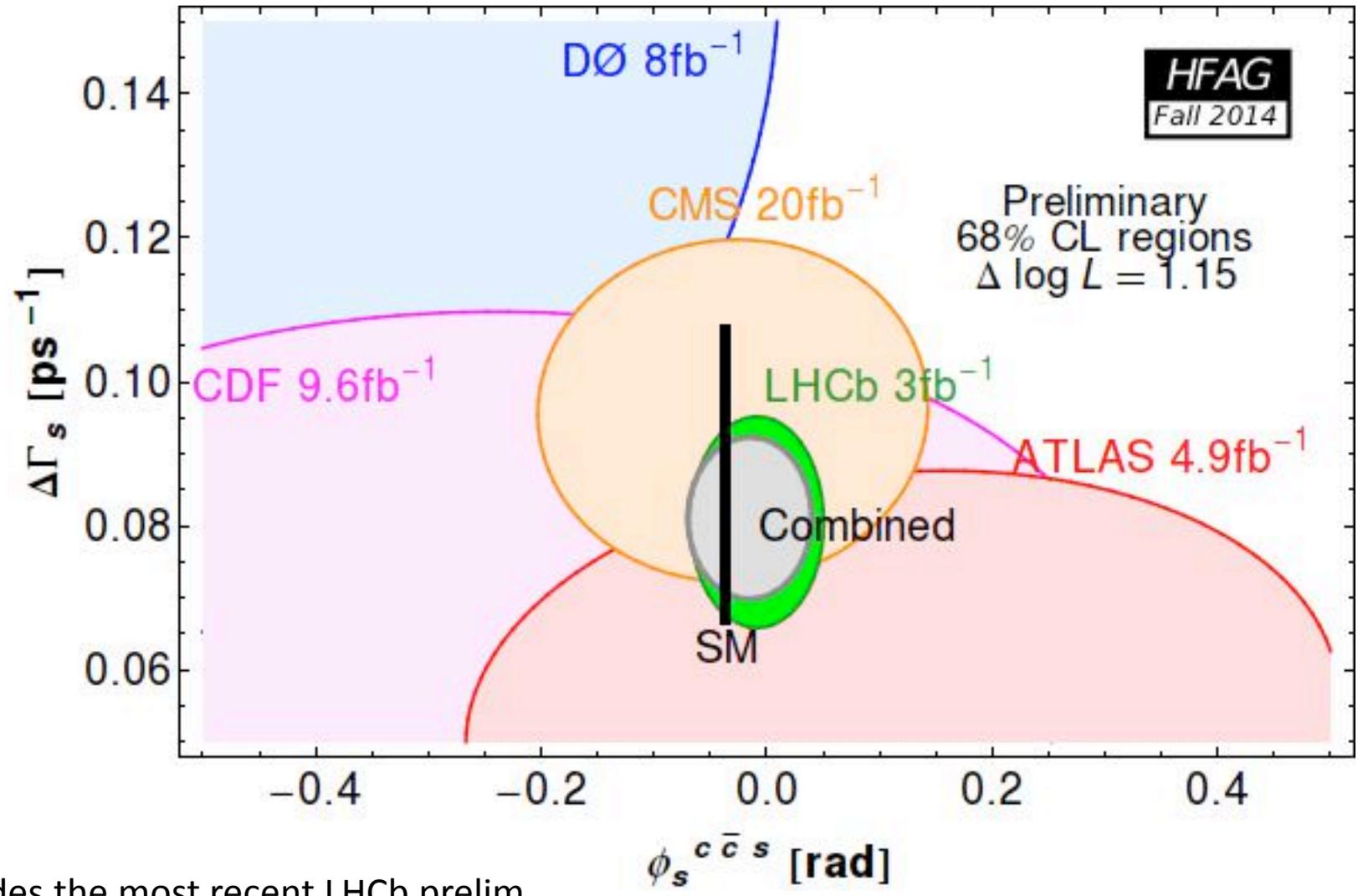
# Boxes



Although B factories can run on the Upsilon(5S), LHCb dominates here

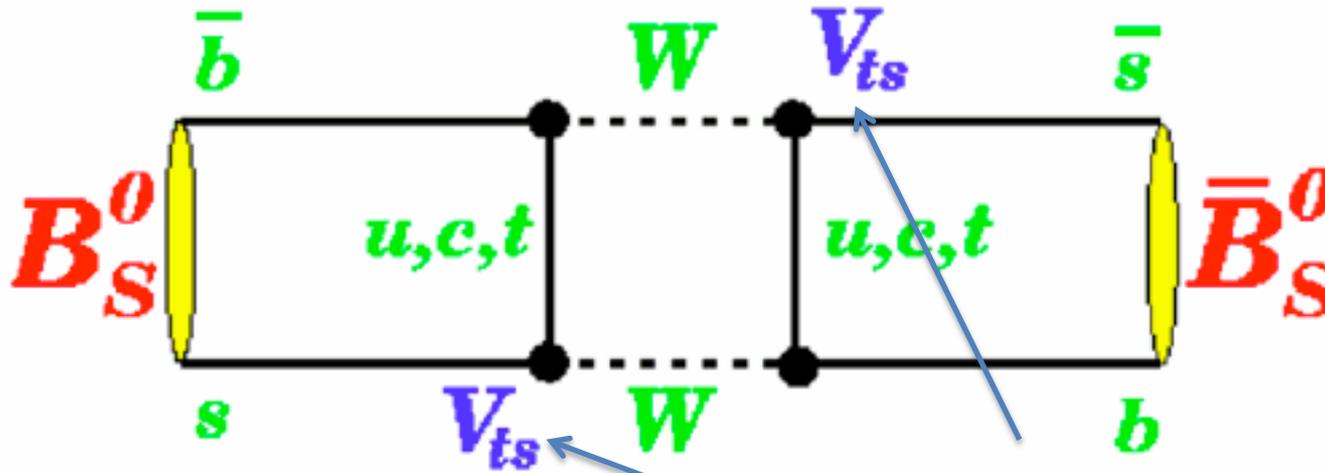
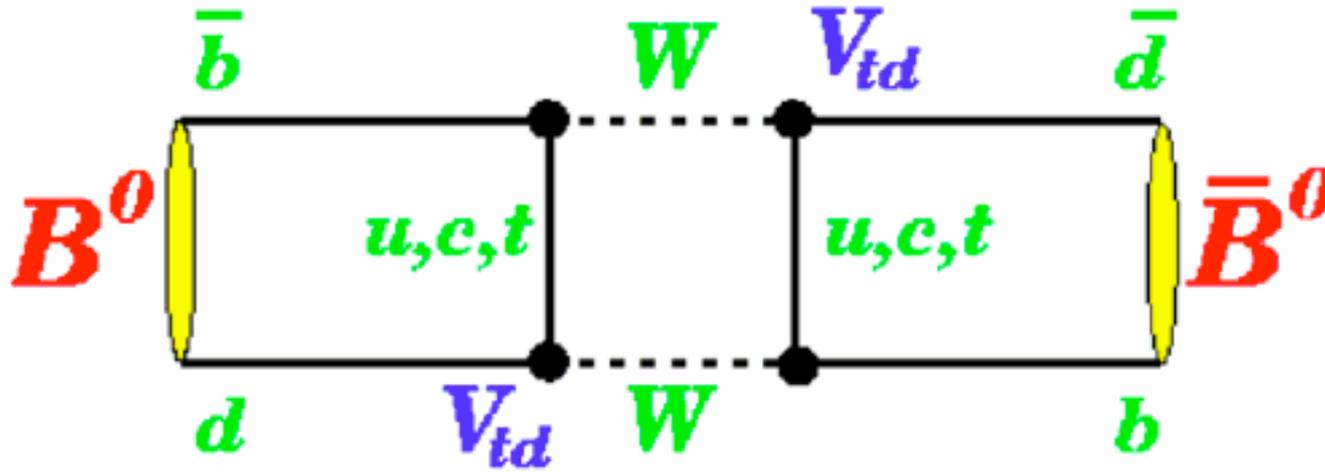
No phase expected from SM but possible from NP particles

# Results on the phase of $B_s$ -anti $B_s$ mixing (i.e. phase of $V_{ts}$ ) [use $B_s \rightarrow J/\psi\varphi$ ; $J/\psi\pi\pi$ modes]

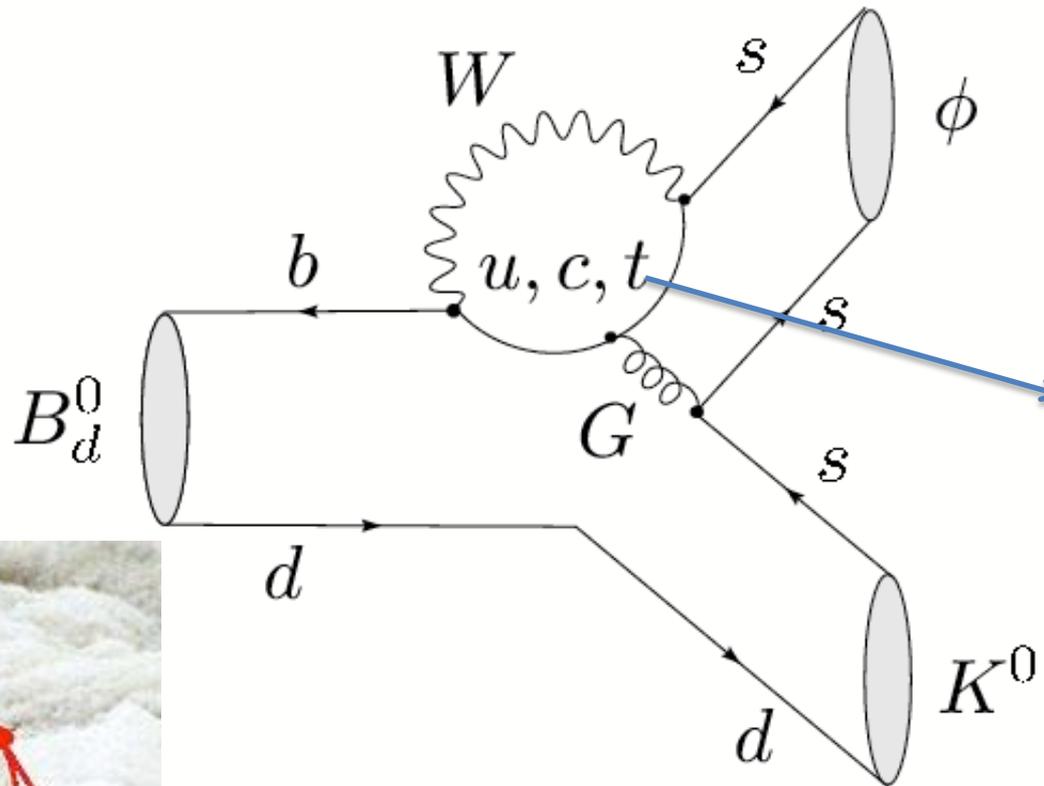


(Includes the most recent LHCb prelim result, giving  $-10 \pm 40$  mrad)

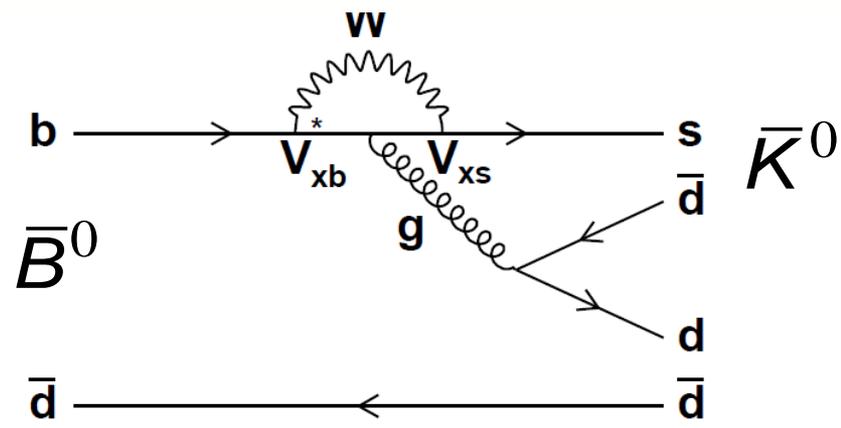
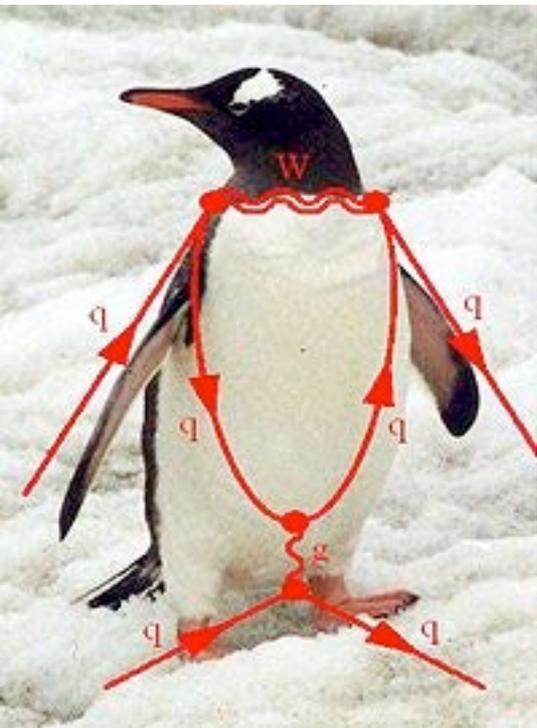
# Boxes



No phase expected from SM but possible from NP particles



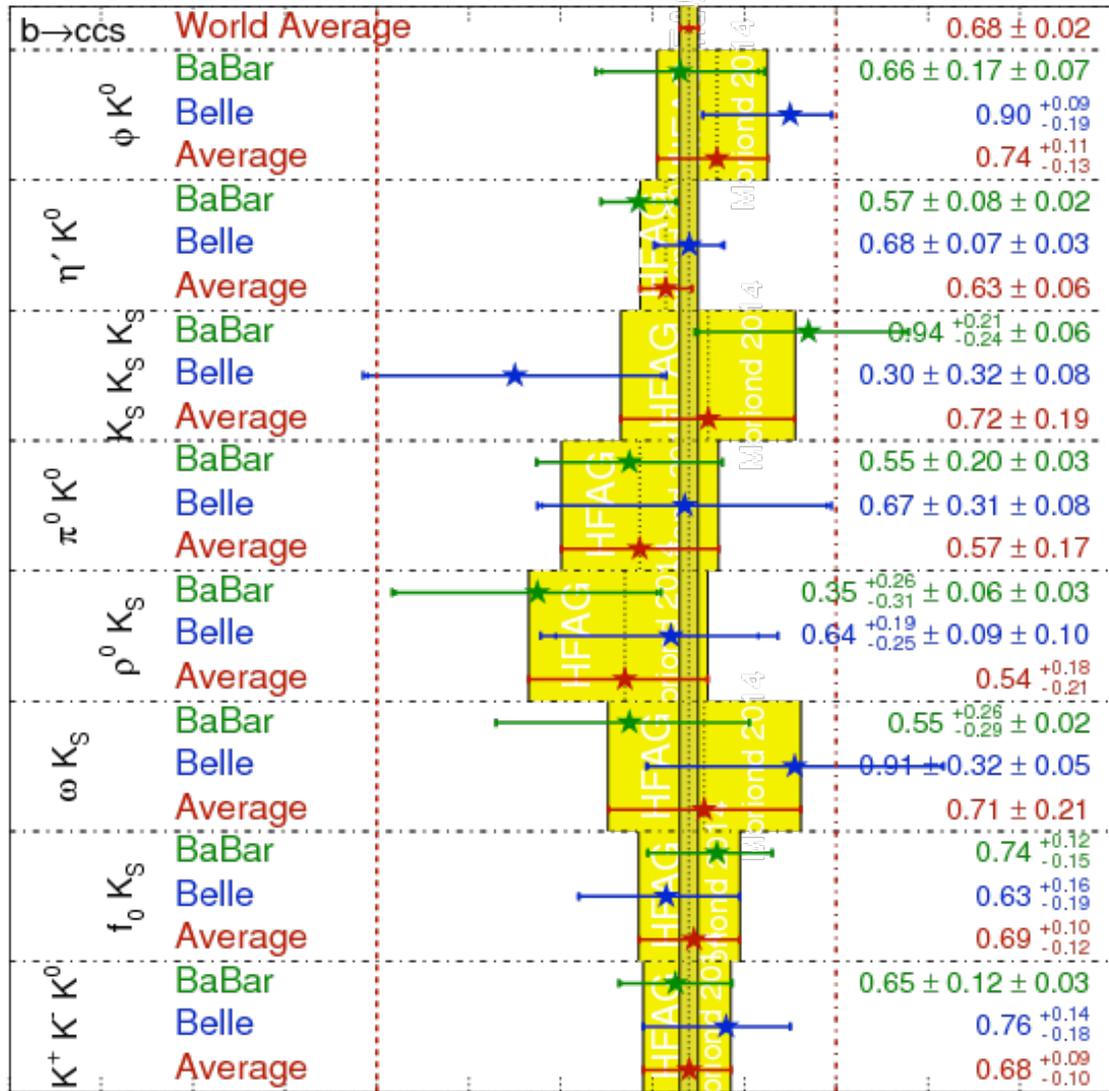
Phase of  $V_{ts}$



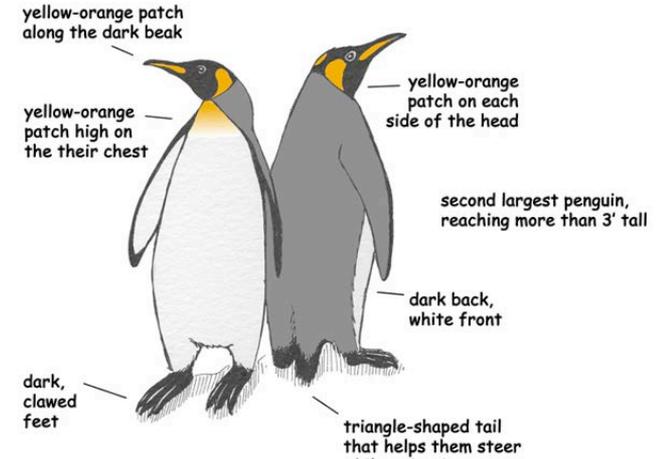
Recent Belle results on  $B \rightarrow \omega K_S$ ,  $B \rightarrow \eta' K_S$

$$\sin(2\beta^{\text{eff}}) \equiv \sin(2\phi_1^{\text{eff}})$$

**HFAG**  
Morioud 2014  
PRELIMINARY



# New Physics Phases in Penguin $b \rightarrow s$ decays



No evidence for NP at current level of sensitivity

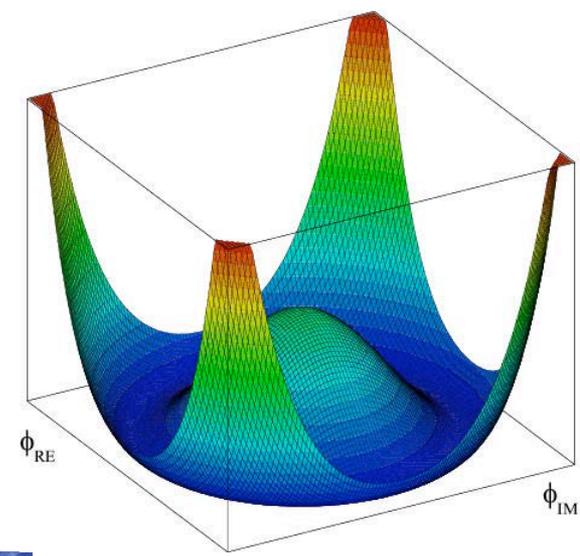
LHCb is absent from this game (lower  $K_S$  eff and flavor tagging eff) but contributes in  $B_S$  modes (Nakada's talk)

# “Missing Energy” Decays

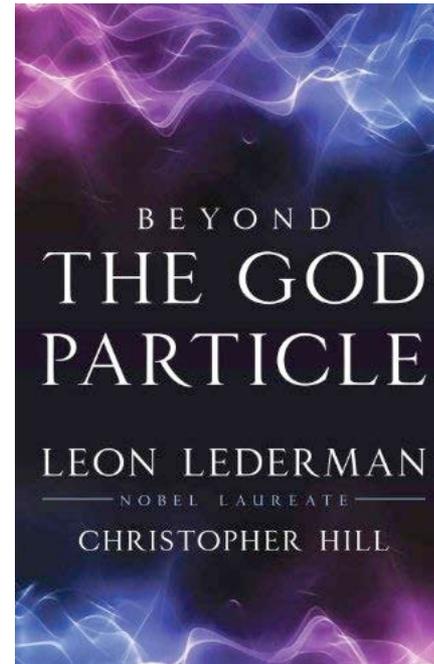


Tokyo Sky Tree

The BEH boson is now firmly established by experimental results from ATLAS and CMS. *Now planning for Higgs flavor factory facilities (e.g ILC, FCC, CEPC, H-LHC).*



Does the GP (Brout-Englert-Higgs particle) have a “brother” i.e. the charged Higgs ?



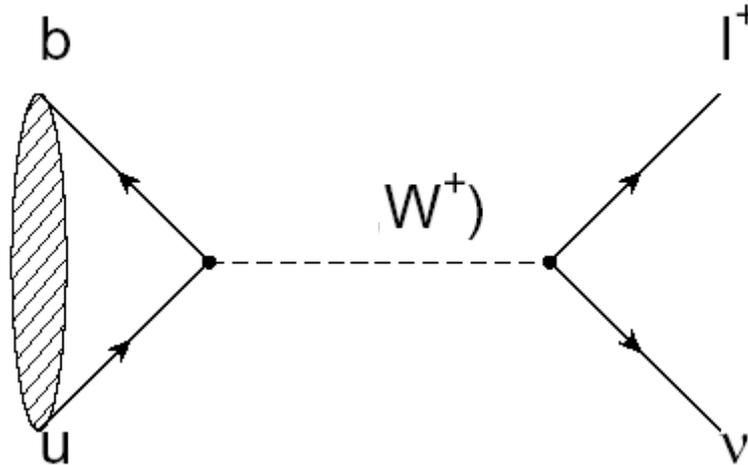
Y. Nambu, 1921-2015

Measurements at Belle II and direct searches at hadron colliders take *complementary* approaches to this important question.

# $B \rightarrow \tau \nu$

(Decay with *Large Missing Energy*)

Sensitivity to new physics from a charged Higgs



$$\mathcal{B}(B^+ \rightarrow \tau^+ \nu_\tau) = \frac{G_F^2 m_B}{8\pi} m_\tau^2 \left(1 - \frac{m_\tau^2}{m_B^2}\right)^2 f_B^2 |V_{ub}|^2 \tau_B$$

$$\mathcal{B}_{(B \rightarrow \tau \nu)} = \mathcal{B}_{SM} \times \left(1 - \tan^2 \beta \frac{m_{B^\pm}^2}{m_{H^\pm}^2}\right)$$



W.S.Hou, PRD 48, 2342 (1993)

The  $B$  meson decay constant, determined by the  $B$  wavefunction at the origin

( $|V_{ub}|$  taken from indep. measurements.)

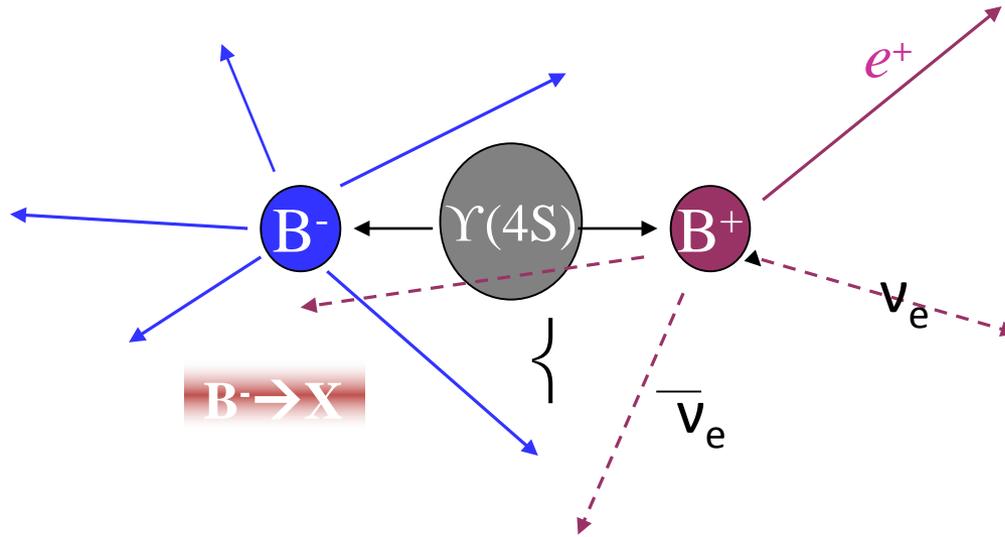
# Consumer's guide to charged Higgs

- Higgs doublet of type I ( $\phi_1$  couples to upper (u-type) and lower (d-type) generations. No fermions couple to  $\phi_2$ )

- Higgs doublet of type II ( $\phi_u$  couples to u type quarks,  $\phi_d$  couples to d-type quarks, u and d couplings are different;  $\tan(\beta) = v_u/v_d$ ) [avored NP scenario e.g. MSSM, generic SUSY]

- Higgs doublet of type III (not type I or type II; anything goes. "FCNC hell"  $\rightarrow$  many FCNC signatures)

# Why measuring $B^+ \rightarrow \tau^+ \nu$ is non-trivial

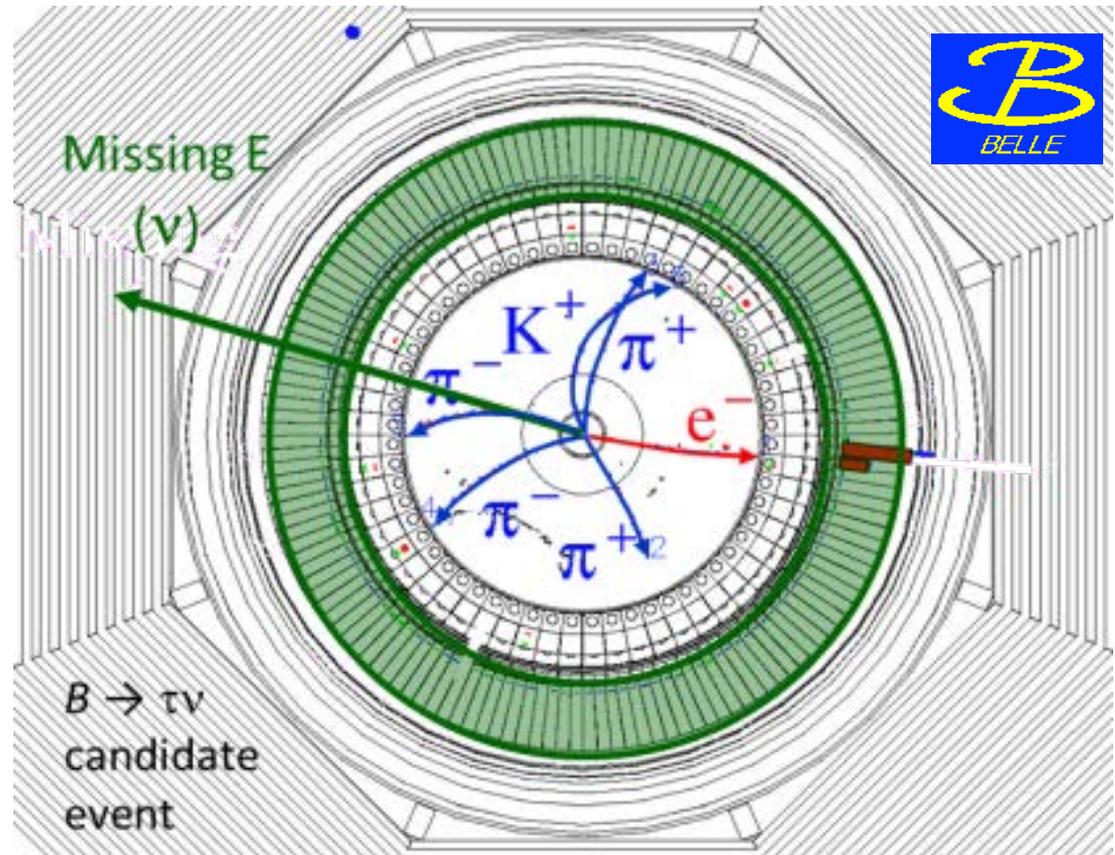
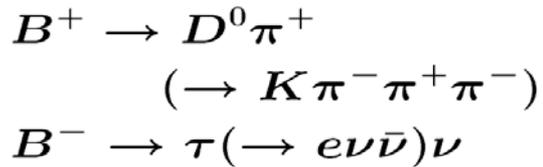


Most of the sensitivity is from tau modes with 1-prongs.

*The experimental signature is rather difficult:  
B decays to a **single charged track + nothing***

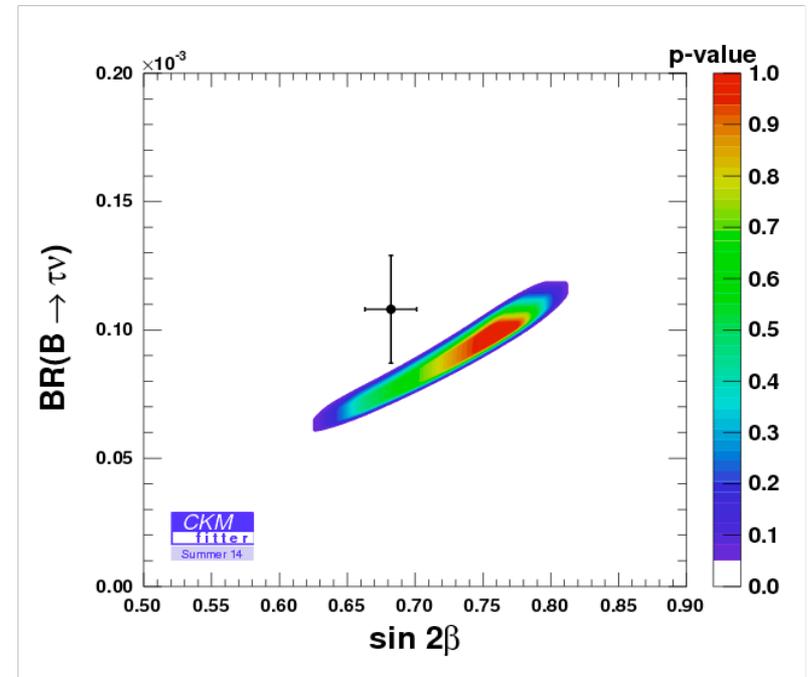
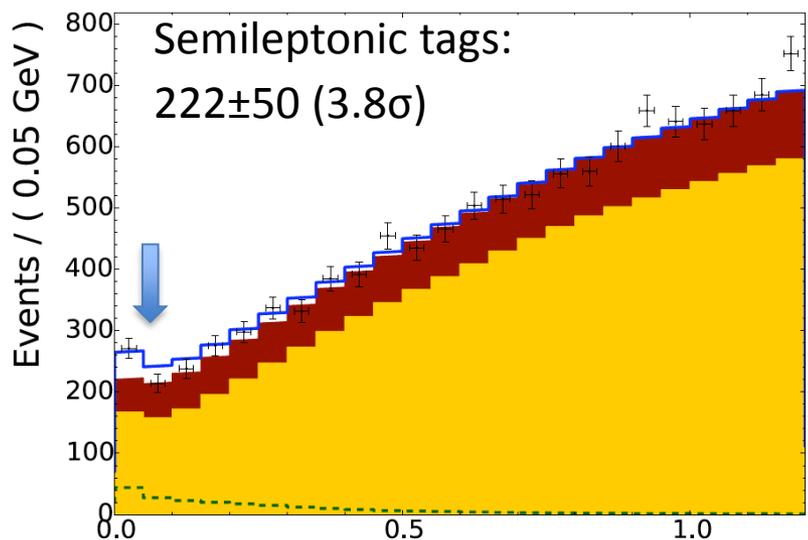
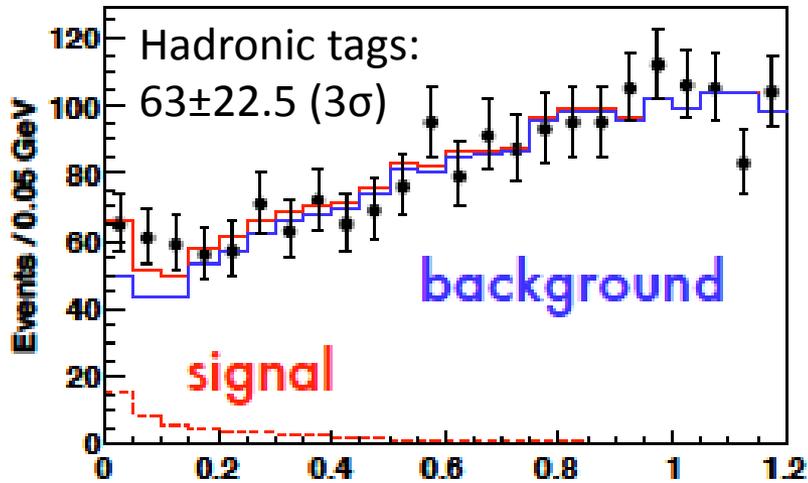
(This may be hard at a hadron collider)

# Example of a Missing Energy Decay ( $B \rightarrow \tau \nu$ ) in Data



The clean  $e^+e^-$  environment makes this possible

Example: Belle  $B \rightarrow \tau \nu$  results with full *reprocessed* data sample and either hadronic or semileptonic tags (**new** arXiv: 1409.5269)



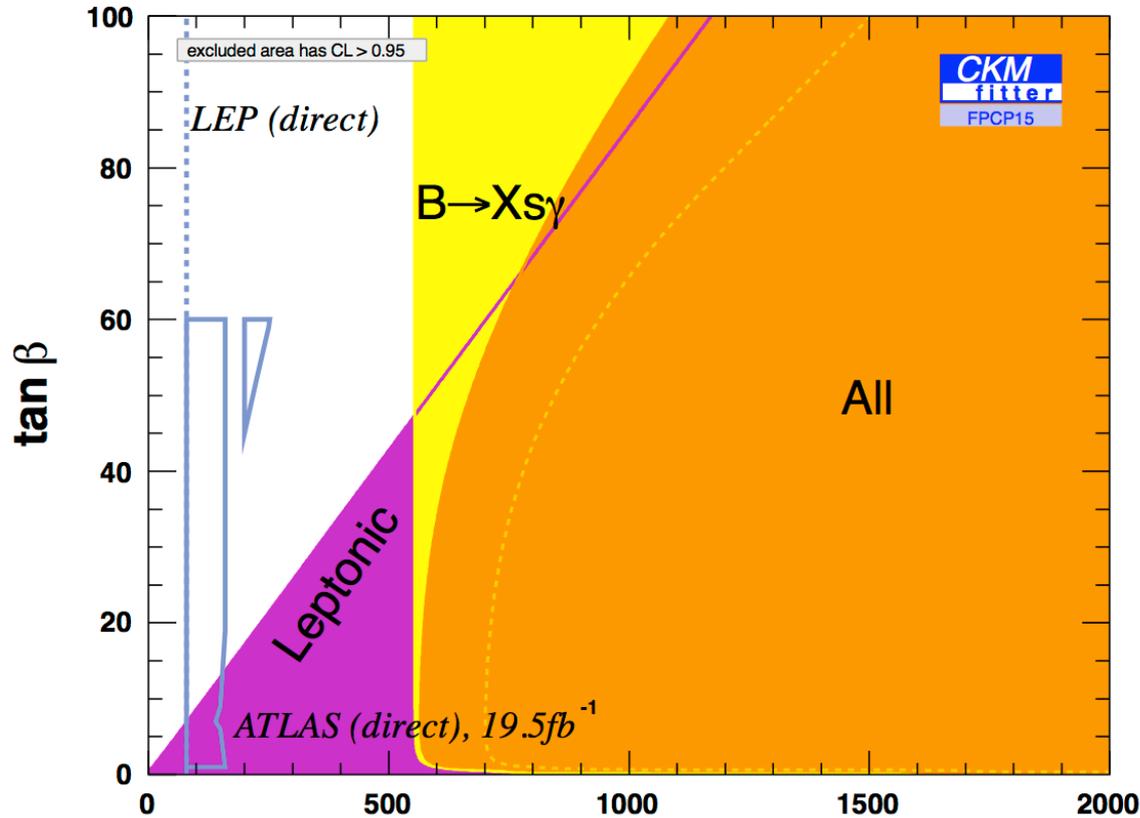
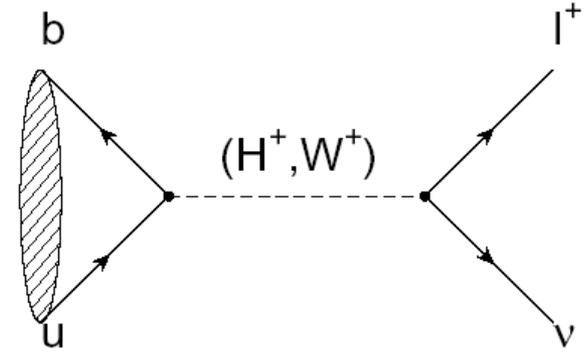
With the full B factory statistics only “evidence”.  
 No single observation from either Belle or BaBar.

The horizontal axis is the “Extra Calorimeter Energy”

# Complementarity of $e^+ e^-$ factories and LHC

(Slide adapted from A. Bevan)

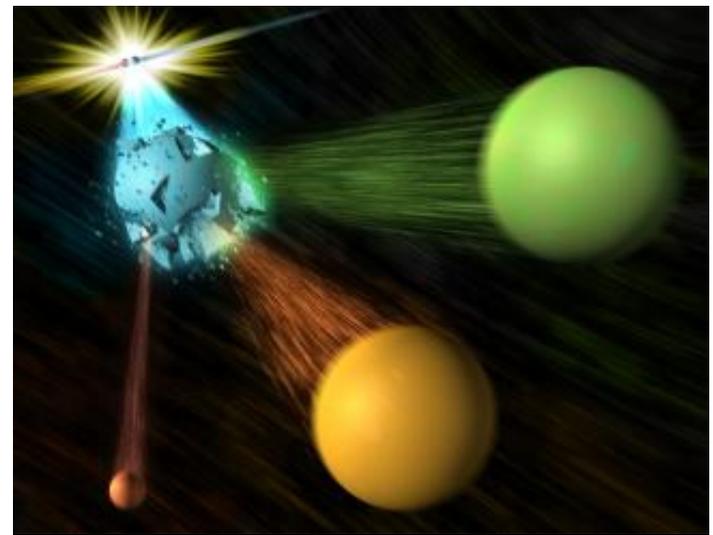
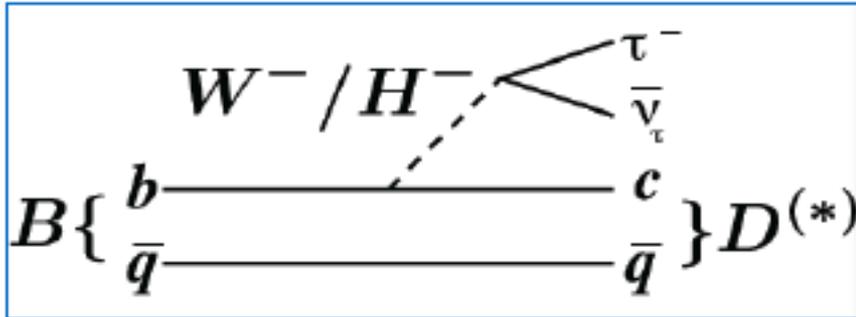
The current combined  $B \rightarrow \tau u$  limit places a stronger constraint than direct searches from LHC exps. for the next few years.



$$r_H = \left( 1 - \frac{m_B^2}{m_H^2} \tan^2 \beta \right)^2$$

Currently **inclusive b to sy** rules out  $m_{H^+}$  below  $\sim 480 \text{ GeV}/c^2$  range at 95% CL (independent of  $\tan\beta$ ), M. Misiak et al.

<http://arxiv.org/abs/1503.01789>



$$\mathcal{R}(D^{(*)})_{2\text{HDM}} = \mathcal{R}(D^{(*)})_{\text{SM}} + A_{D^{(*)}} \frac{\tan^2 \beta}{m_{H^+}^2} + B_{D^{(*)}} \frac{\tan^4 \beta}{m_{H^+}^4}$$

	$D\tau\nu$	$D^*\tau\nu$
$A_{D^{(*)}} \text{ (GeV}^2\text{)}$	$-3.25 \pm 0.32$	$-0.230 \pm 0.029$
$B_{D^{(*)}} \text{ (GeV}^4\text{)}$	$16.9 \pm 2.0$	$0.643 \pm 0.085$

$$R(D^{(*)}) = \frac{\mathcal{B}(\bar{B} \rightarrow D^{(*)} \tau^- \bar{\nu}_\tau)}{\mathcal{B}(\bar{B} \rightarrow D^{(*)} \ell^- \bar{\nu}_\ell)} \begin{matrix} \longrightarrow & \text{Signal} \\ \longrightarrow & \text{Normalization } (l = e \text{ or } \mu) \end{matrix}$$

Example from a BaBar paper

Signals in  $B \rightarrow D^{(*)} \tau \nu$  ( $489 \pm 63$ ,  $888 \pm 63$ )

Missing mass variable:

$$m_{\text{miss}}^2 = p_{\text{miss}}^2 = (p[e^+e^-] - p_{\text{tag}} - p_{D^{(*)}} - p_l)^2$$

$P_l^*$  = momentum of lepton in B rest frame

**But wait !!! Now possible at LHCb.**

*Production of B meson pairs at threshold is critical to the separation of backgrounds from the missing energy/ momentum signal.*

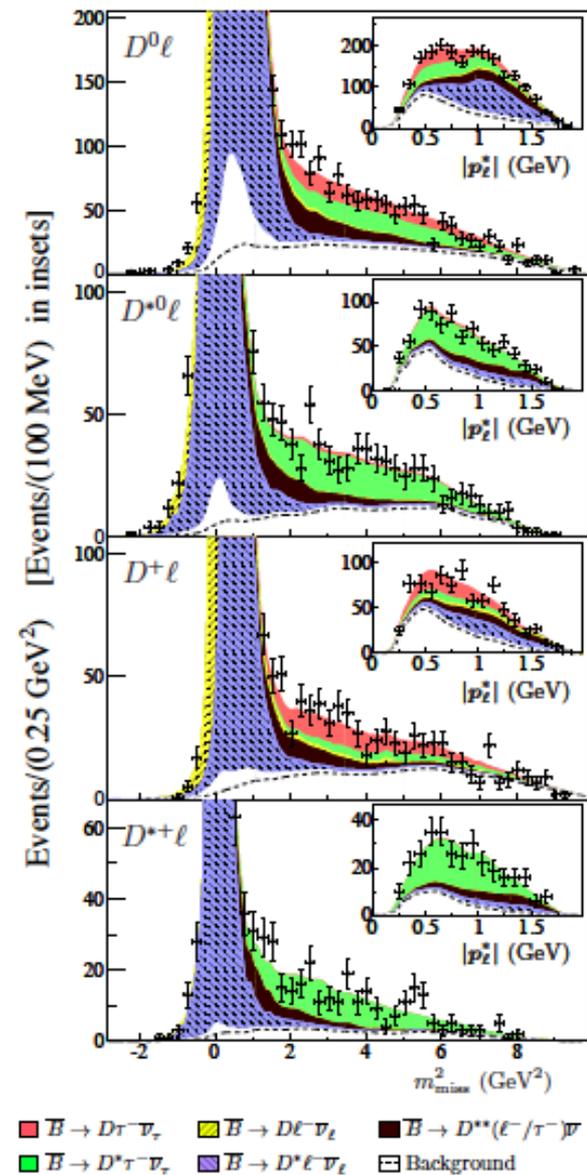
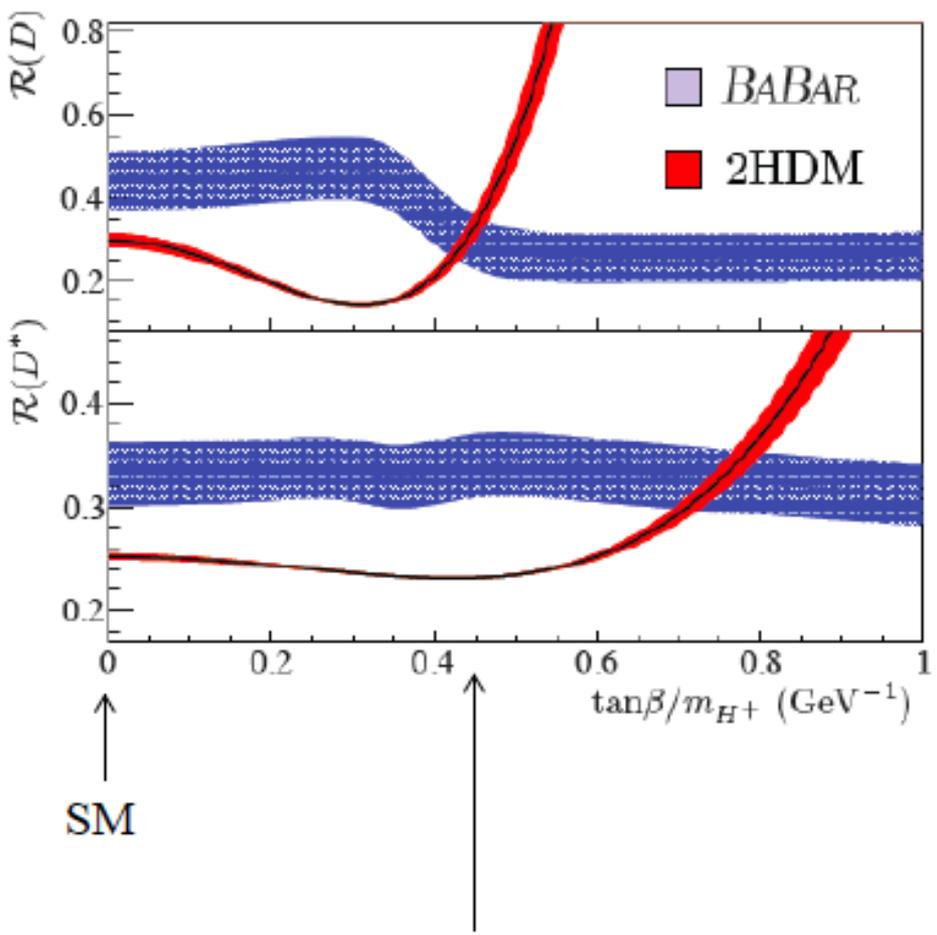


FIG. 1. (Color online) Comparison of the data and the fit projections for the four  $D^{(*)} \ell$  samples. The insets show the  $|p_l^*|$  projections for  $m_{\text{miss}}^2 > 1 \text{ GeV}^2$ , which excludes most of the normalization modes. In the background component, the region above the dashed line corresponds to charge cross-feed, and the region below corresponds to continuum and  $B\bar{B}$ .

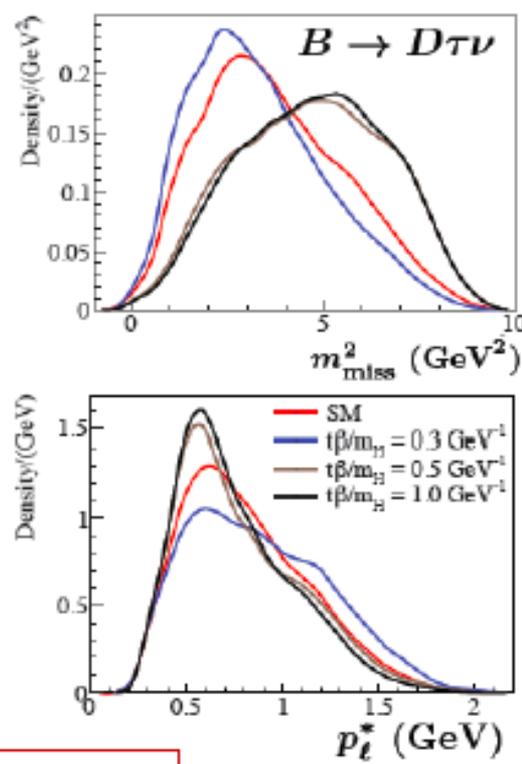
# Limits on type-II 2HDM



$\tan\beta/m_{H^+} =$   
 $0.44 \pm 0.02$  GeV $^{-1}$

$\tan\beta/m_{H^+} =$   
 $0.75 \pm 0.04$  GeV $^{-1}$

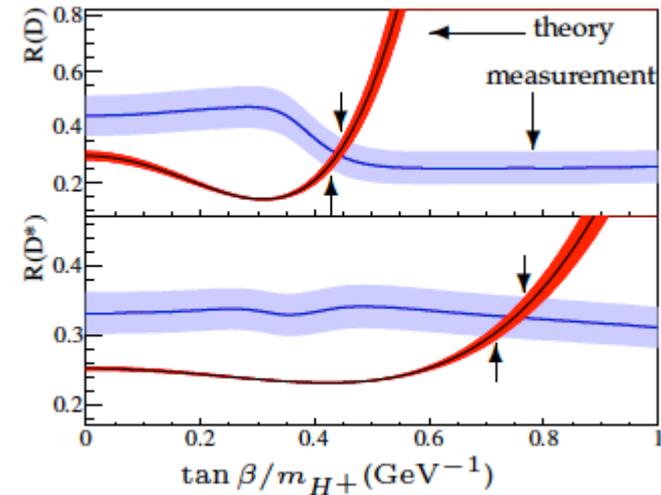
2HDM modifies fit-variable distribution and hence the efficiency



Best point is  $\tan\beta/m_{H^+} = 0.45$  GeV $^{-1}$ , excluded at 99.8% CL ( $3.1 \sigma$ ).  
 All other values (with  $m_{H^+} > 15$  GeV) are worse.

BaBar collaboration, Phys. Rev. Lett. 109, 101802 (2012)

“However, the combination of  $R(D)$  and  $R(D^*)$  excludes the type II 2HDM charged Higgs boson with a 99.8% confidence level for any value of  $\tan(\beta)/m_{H^+}$ ”



In other words, found NP but *killed* the 2HDM NP model.

Warning: color-coding different from BaBar

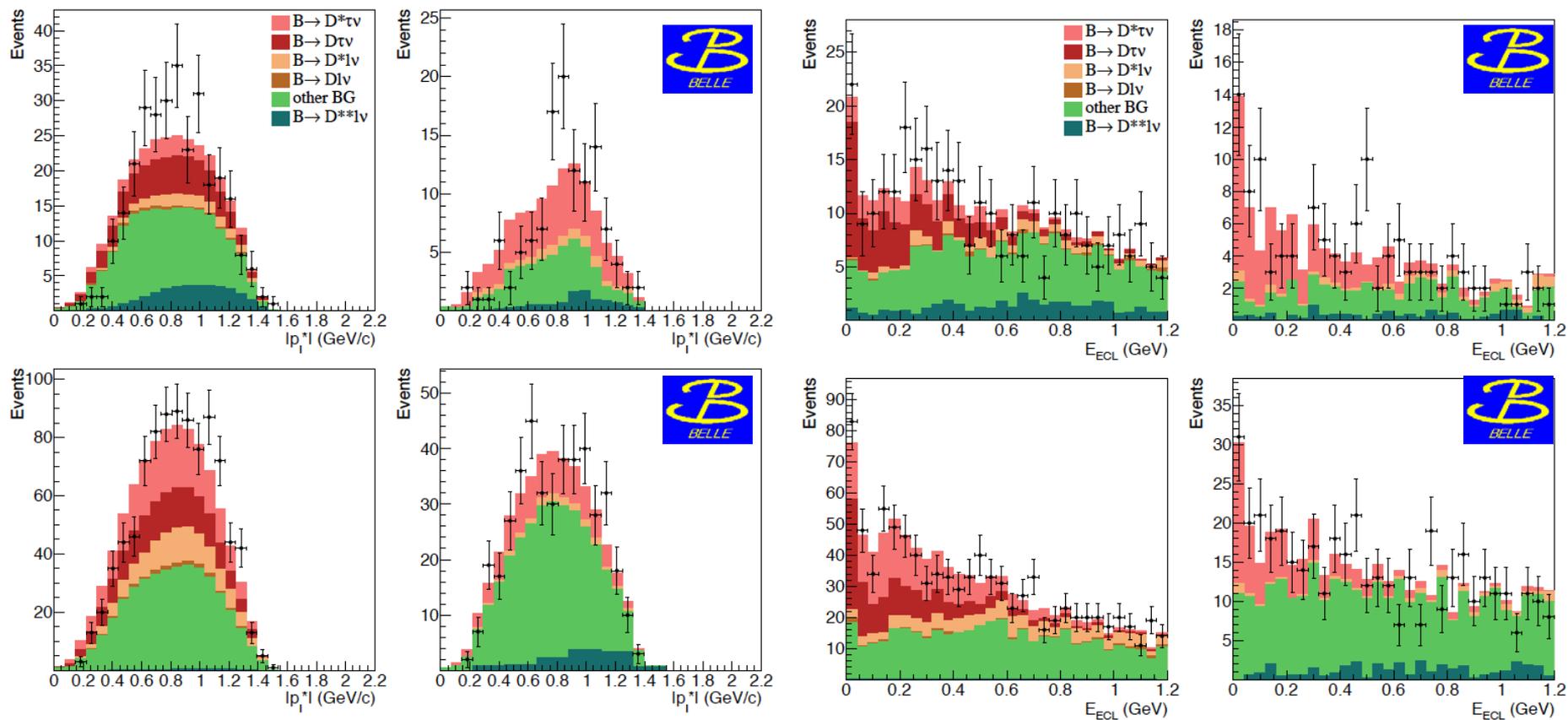


FIG. 6. Projections of the fit results and data points with statistical uncertainties in a signal enhanced region of  $M_{miss}^2 > 2.0 \text{ GeV}^2/c^4$  in the  $p_l^*$  dimension. Top left:  $D^+\ell^-$ ; top right:  $D^{*+}\ell^-$ ; bottom left:  $D^0\ell^-$ ; bottom right:  $D^{*0}\ell^-$ .

FIG. 5. Projections of the fit results and data points with statistical uncertainties in a signal enhanced region of  $M_{miss}^2 > 2.0 \text{ GeV}^2/c^4$  in the  $E_{ECL}$  dimension. Top left:  $D^+\ell^-$ ; top right:  $D^{*+}\ell^-$ ; bottom left:  $D^0\ell^-$ ; bottom right:  $D^{*0}\ell^-$ .

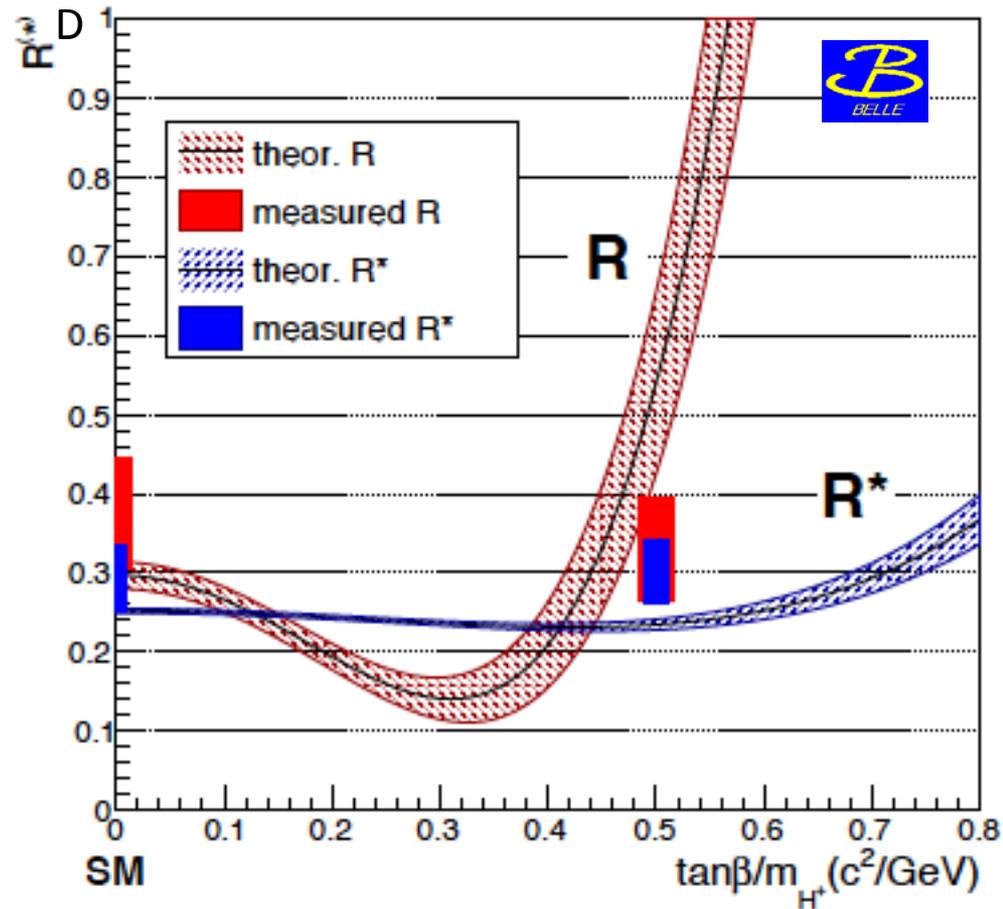
Signal enhanced projections  
of lepton momenta in the  
high  $M_{miss}^2$  region

Signal enhanced projections of  
 $E_{ECL}$  momenta in the high  $M_{miss}^2$   
region

# New Belle result with hadronic tags

May 25 2015, Nagoya FPCP

<http://xxx.lanl.gov/abs/1507.03233>; submitted to Phys Rev



Compatible with both  
BaBar and the 2HDM  
model (and SM !).

Need more data and  
more Belle analyses to  
resolve the issue.

FIG. 8. Theoretical predictions with  $1\sigma$  error ranges for  $R$  (red) and  $R^*$  (blue) for different values of  $\tan\beta/m_{H^+}$  in the 2HDM of type II. This analysis' fit results for  $\tan\beta/m_{H^+} = 0.5 c^2/\text{GeV}$  and SM are shown with their  $1\sigma$  ranges as red and blue bars with arbitrary width for better visibility.

# New LHCb result

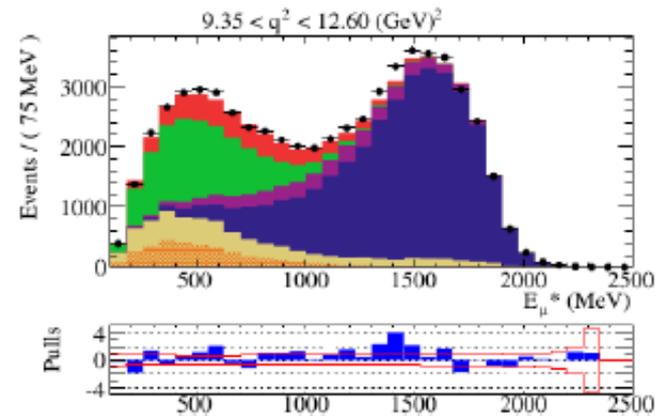
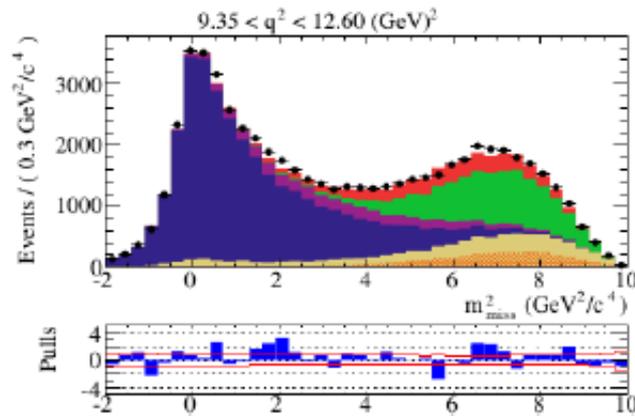
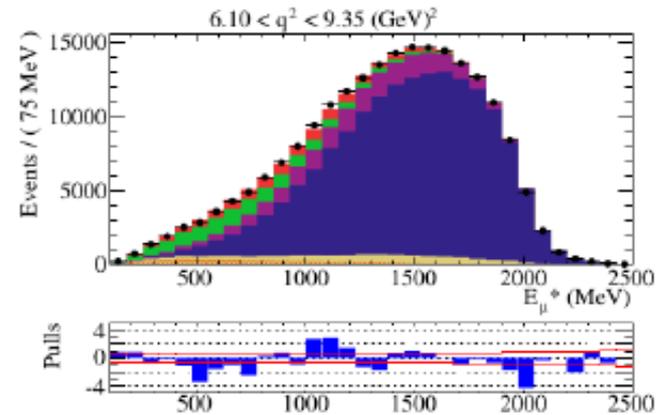
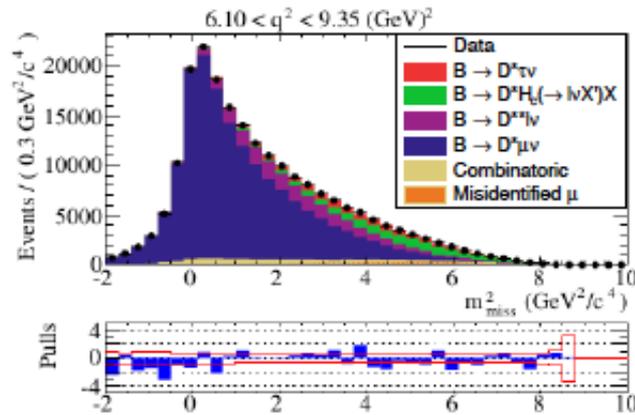
May 25 2015, Nagoya FPCP

Compatible with BaBar  $D^* \tau \nu$   
BF ( $B \rightarrow D \tau \nu$  coming soon)

## 3. $B \rightarrow D^* \tau \nu$ (LHCb-PAPER-2015-025)

### Signal fit

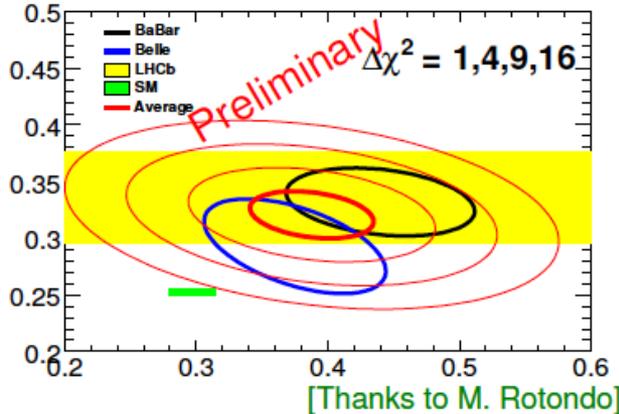
Mass resolution  
is poor but  
vertex isolation  
is very powerful



# Apres Nagoya: New World Averages for $R(D)$ and $R(D^*)$

	$R(D)$	$R(D^*)$
BaBar	$0.440 \pm 0.058 \pm 0.042$	$0.332 \pm 0.024 \pm 0.018$
Belle	$0.375^{+0.064}_{-0.063} \pm 0.026$	$0.293^{+0.039}_{-0.037} \pm 0.015$
LHCb		$0.336 \pm 0.027 \pm 0.030$
Average	$0.388 \pm 0.047$	$0.321 \pm 0.021$
SM expectation	$0.300 \pm 0.010$	$0.252 \pm 0.005$
Belle II, 50/ab	$\pm 0.010$	$\pm 0.005$

Now  $\sim 4\sigma$  from SM



A few points:

- (0) Need additional Belle results (e.g. semileptonic tags) and LHCb  $B \rightarrow D^0 \tau \nu$  result
- (1) Still tension with *NP from type II charged Higgs*
- (2) It is *obvious* that we need two orders of magnitude of data to solve these issues related to the charged Higgs.

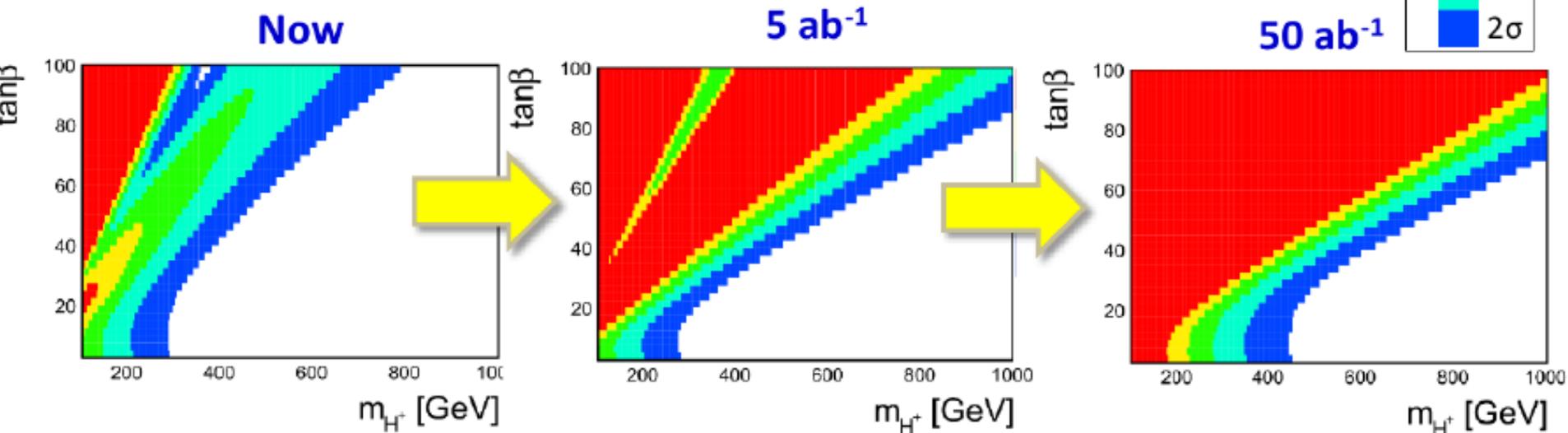
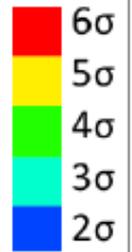
# Initial Belle II projections for charged Higgs sensitivity

	Exp.			Th.
	Now	5 ab <sup>-1</sup>	50 ab <sup>-1</sup>	Now
$B \rightarrow \tau\nu$	25%	10%	3%	-7+14%
$B \rightarrow D\tau\nu$	30%	11%	4%	4%
$B \rightarrow D^*\tau\nu$	19%	7%	2%	2%
$B \rightarrow X_s\gamma$	7%	5%	4%	7%

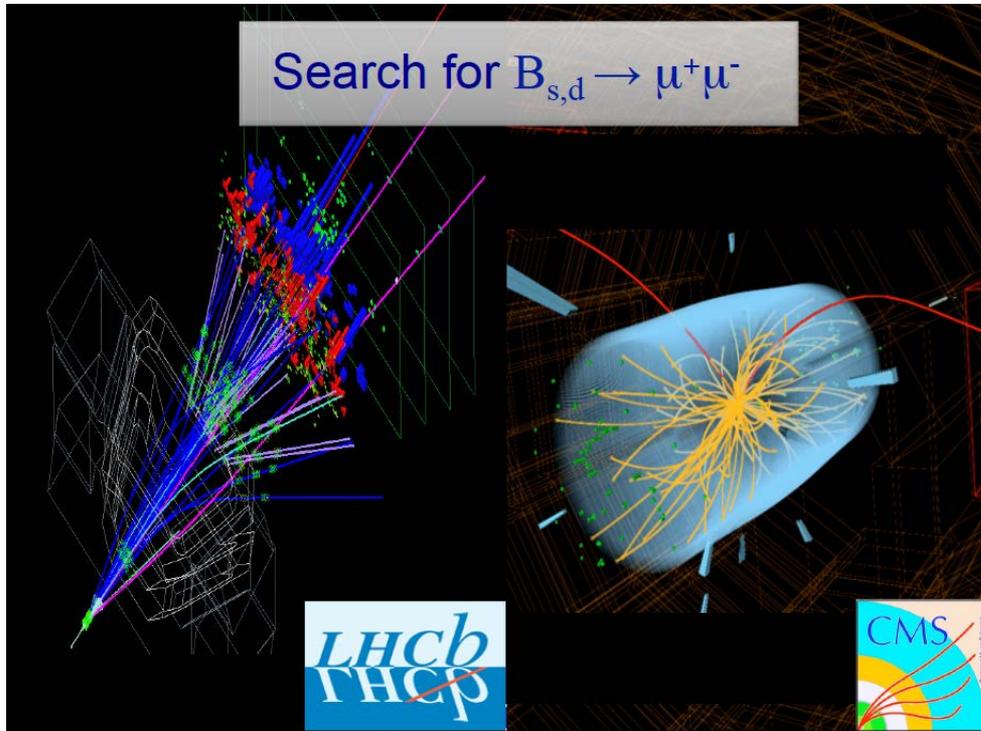
Will improved by precise  $V_{ub}$  measurements.  
 My naive estimation assuming  $\sigma_{fB} \sim 1\%$  :  
 $\sim 5\%$  @ Belle II era

My naive estimation

Excl. at



# Rare B Decays



J. Albrecht



Goa, India

# LHC found the rarest B decay; $B_s \rightarrow \mu^+ \mu^-$

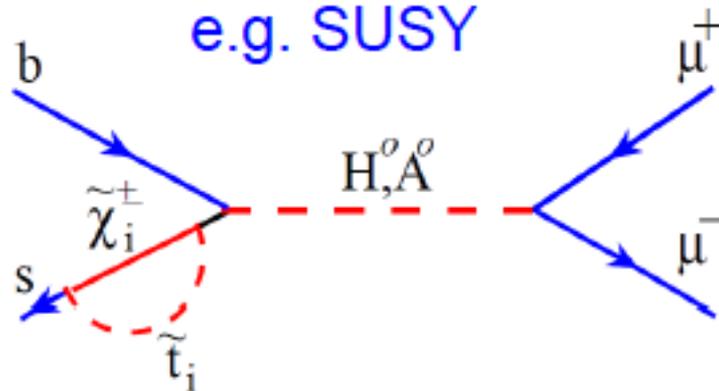
Standard Model



$BF \sim O(10^{-9})$

Left handed couplings  
 $\rightarrow$  helicity suppressed

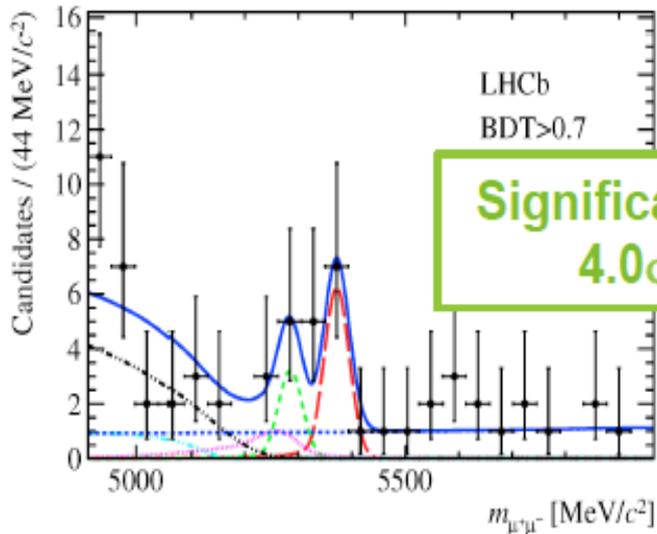
e.g. SUSY



N. B. Here and in  $b \rightarrow s l^+ l^-$  all the heavy particles of the SM enter as virtual particles in the Feynman diagrams

# LHCb

- Update: full dataset:  $3\text{fb}^{-1}$ 
  - Improved BDT
  - Expected sensitivity:  $5.0\sigma$

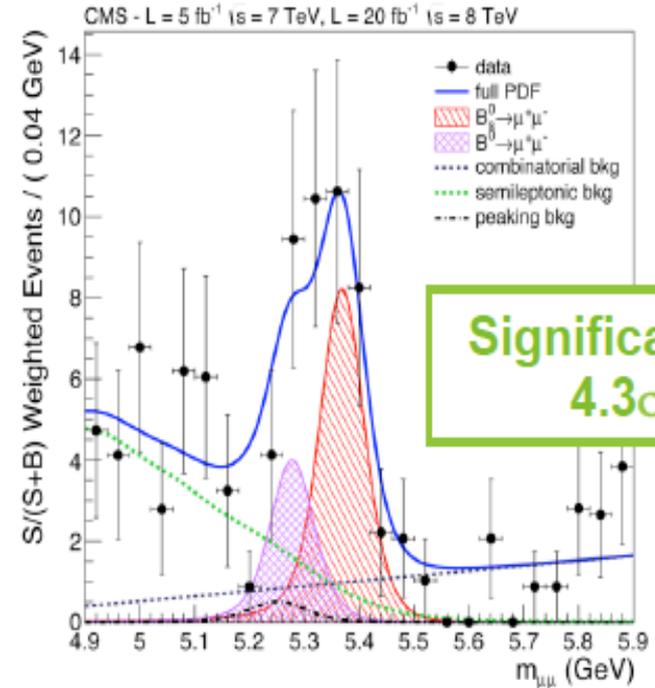


$$BR(B_s \rightarrow \mu^+ \mu^-) = (2.9^{+1.1}_{-1.0}) \times 10^{-9}$$

$$BR(B^0 \rightarrow \mu^+ \mu^-) = (3.7^{+2.4}_{-2.1}) \times 10^{-10}$$

$$BR(B^0 \rightarrow \mu^+ \mu^-) < 7 \times 10^{-10} @ 95\%CL$$

# CMS

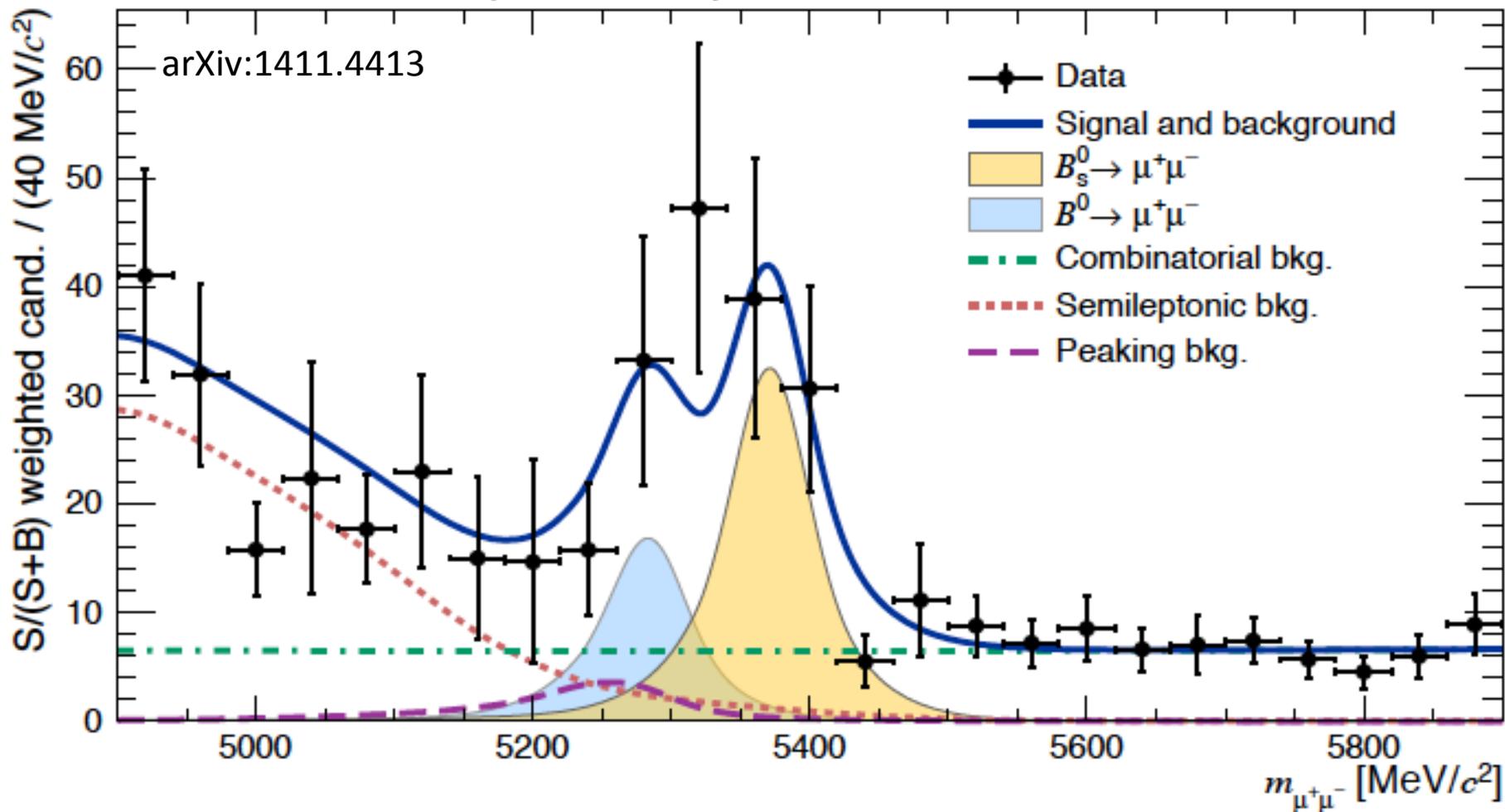


$$BR(B_s \rightarrow \mu^+ \mu^-) = (3.0^{+1.0}_{-0.9}) \times 10^{-9}$$

$$BR(B^0 \rightarrow \mu^+ \mu^-) = (3.5^{+2.1}_{-1.8}) \times 10^{-10}$$

$$BR(B^0 \rightarrow \mu^+ \mu^-) < 11 \times 10^{-10} @ 95\%CL$$

# CMS and LHCb (LHC run I)



Published in Nature: June 4, 2015

$$\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-) = (2.9 \pm 0.7) \times 10^{-9}$$

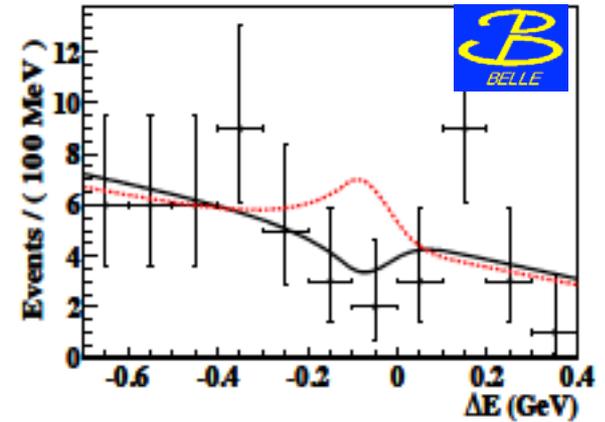
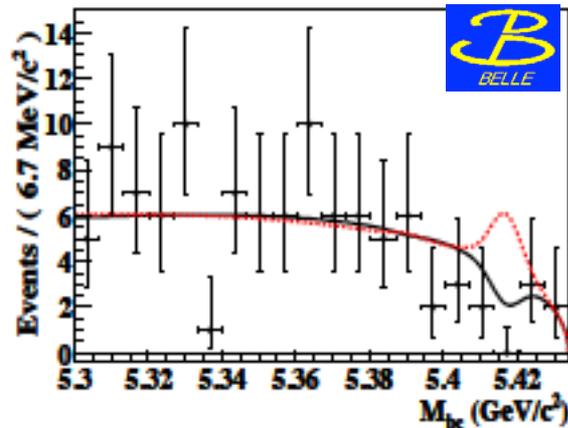
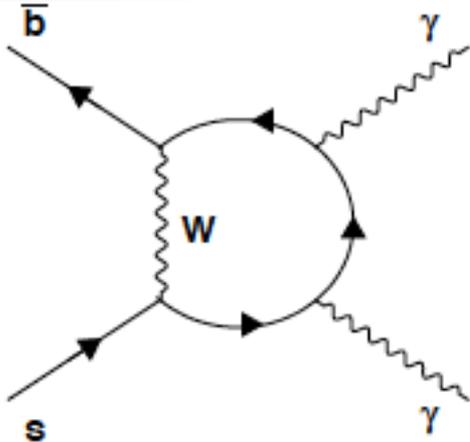
$$\mathcal{B}(B^0 \rightarrow \mu^+ \mu^-) = (3.6 \pm_{-1.4}^{+1.6}) \times 10^{-10}$$



?

SM:  $\text{BR}(B_c) = (3.65 \pm 0.23) \times 10^{-9}$   
 $\text{BR}(B^0) = (1.1 \pm 0.1) \times 10^{-10}$   
 PRL 112 101801 (2014)

Combining evidence from two LHC experiments (LHCb and CMS),  $B_s \rightarrow \mu^+ \mu^-$  is observed with  $6.2\sigma$  significance. The corresponding  $B_d$  decay is not clearly seen yet.



$$\text{BF}(B_s \rightarrow \gamma\gamma) < 3.1 \times 10^{-6}$$

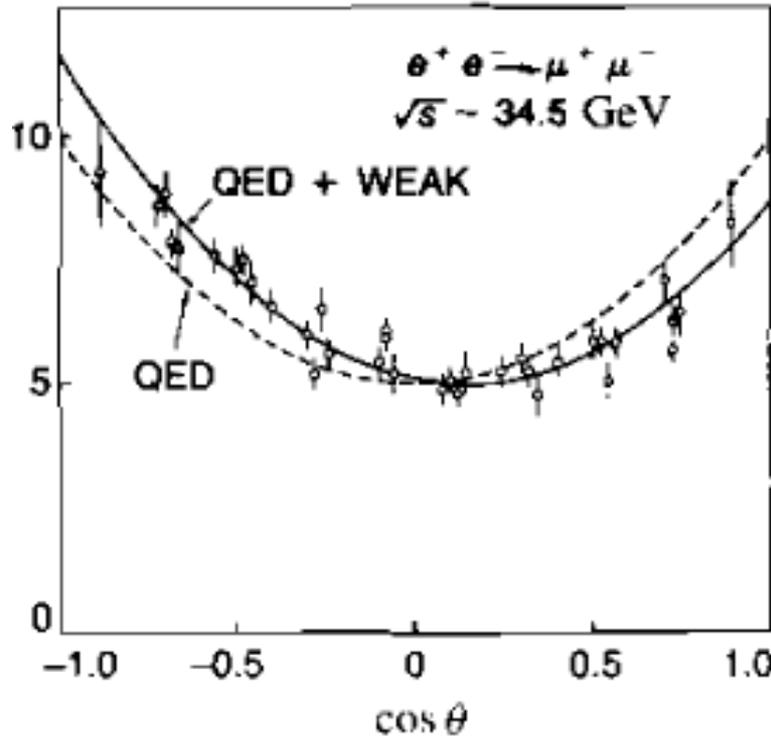
D. Dutta et al (Belle)  
 Phys. Rev. D 91, 011101(R)

Complementarity [uses and requires Upsilon(5S) data]

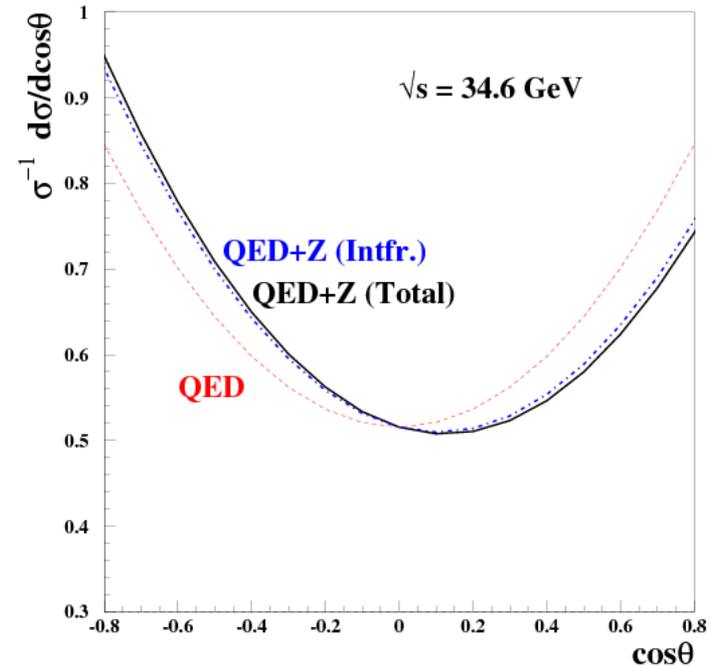
# Red Hot Flavor Physics



# High Energy Physics History: finding NP in $A_{FB}$ (using interference)



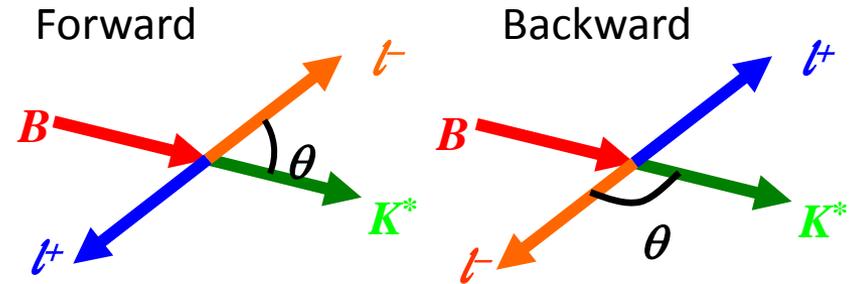
a fit including the weak interaction (solid line).



*Conclusion: There is a Z boson at higher energy even though colliders of the time did not have enough  $\sqrt{s}$  to produce it*

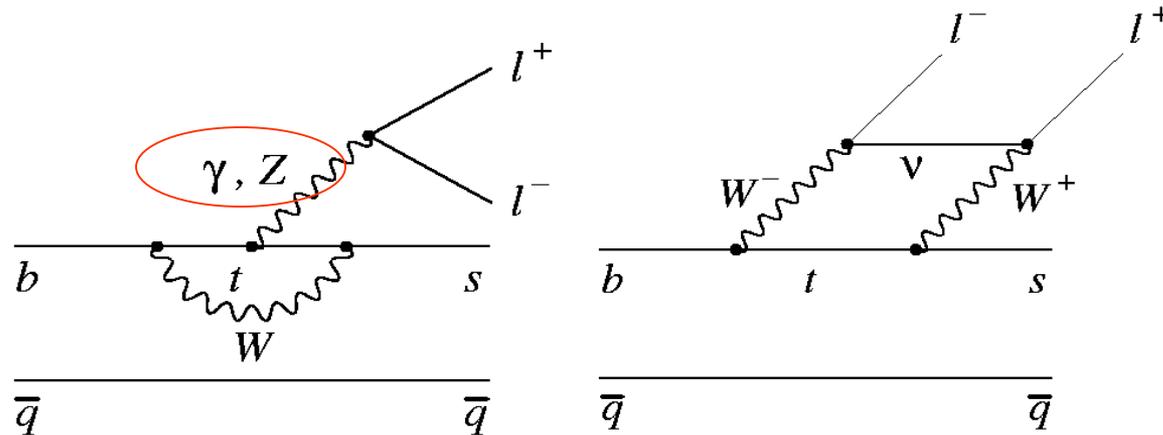
# $A_{FB}(B \rightarrow K^* l^+ l^-)(q^2)$

The SM forward-backward asymmetry in  $b \rightarrow s l^+ l^-$  arises from the interference between  $\gamma$  and  $Z^0$  contributions.



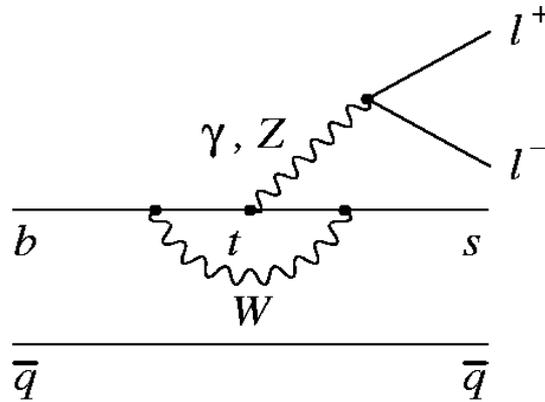
$$A_{FB}(B \rightarrow K^* l^+ l^-) = -C_{10} \xi(q^2) \left[ \text{Re}(C_9) F_1 + \frac{1}{q^2} C_7 F_2 \right]$$

Ali, Mannel, Morozumi, PLB273, 505 (1991)



Note that all the heavy particles of the SM (W, Z, top) enter in this decay.

# More on $A_{FB}(B \rightarrow K^* l^+ l^-)(q^2)$



Can in effect vary  $v_s$  for NP

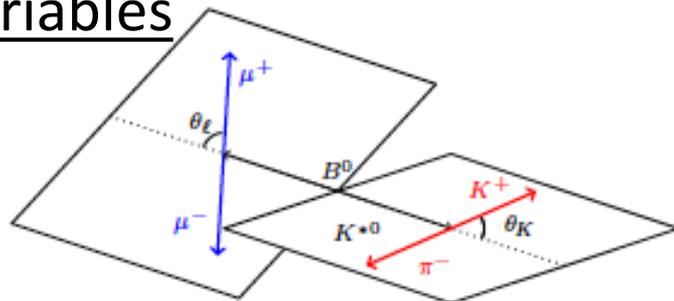
$A_{FB}$  depends on  $q^2 = M^2(l^+ l^-)$

$$A_{FB}(B \rightarrow K^* l^+ l^-) = -C_{10} \xi(q^2) \left[ \text{Re}(C_9) F_1 + \frac{1}{q^2} C_7 F_2 \right]$$

Ali, Mannel, Morezumi, PLB273, 505 (1991)

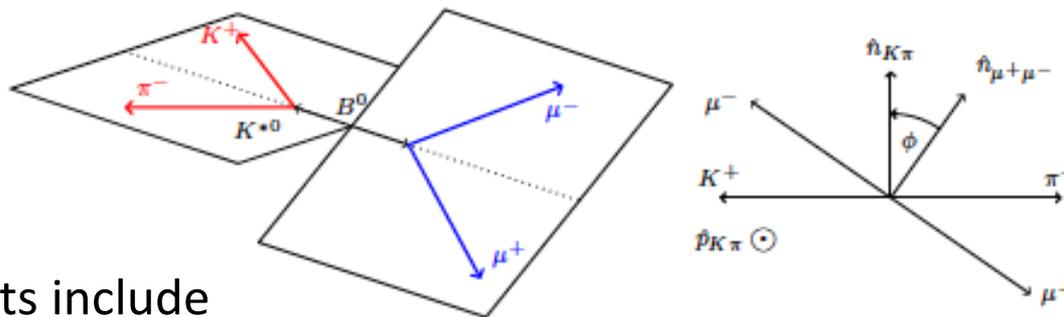
The “zero-crossing” of  $A_{FB}$  depends only on a ratio of form factors and is a *clean* observable.

# $B \rightarrow K^* \ell \ell$ angular variables



$K^*$  and  $\ell^+ \ell^-$  helicity angles

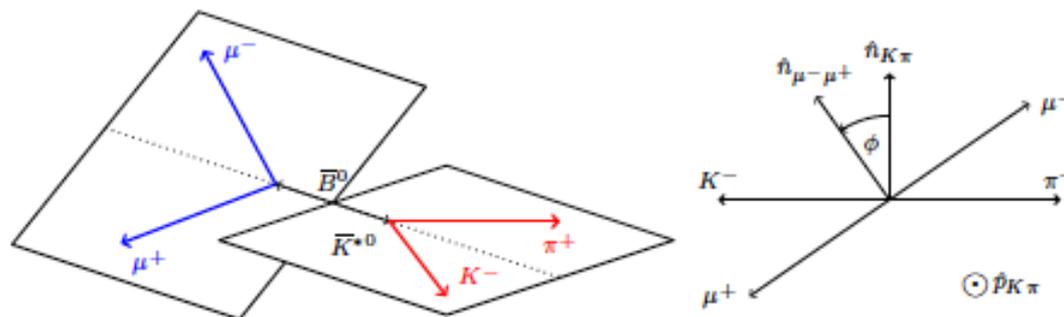
(a)  $\theta_K$  and  $\theta_\ell$  definitions for the  $B^0$  decay



(b)  $\phi$  definition for the  $B^0$  decay

Angle between the normals to the two decay planes.

N.B. Recent measurements include  $\phi$  angle data

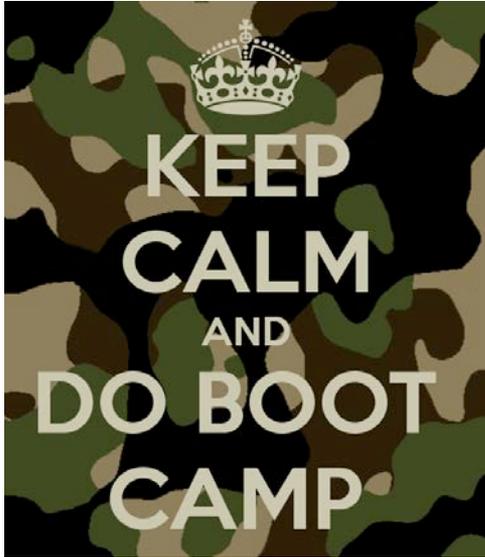


(c)  $\phi$  definition for the  $\bar{B}^0$  decay

From the 2013 LHCb paper

# B → K\* 1+ 1-(q<sup>2</sup>) bootcamp

Angular dependence



(-) means the term is only in  $\Gamma - \bar{\Gamma}$

$$\frac{1}{d(\Gamma + \bar{\Gamma}) / dq^2} \frac{d^3(\Gamma + \bar{\Gamma})}{d\bar{\Omega}} =$$

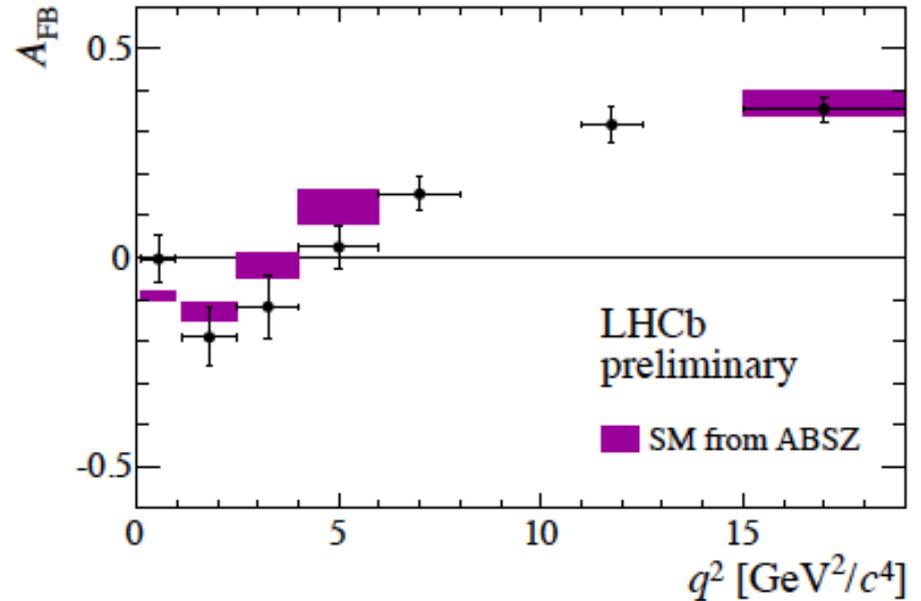
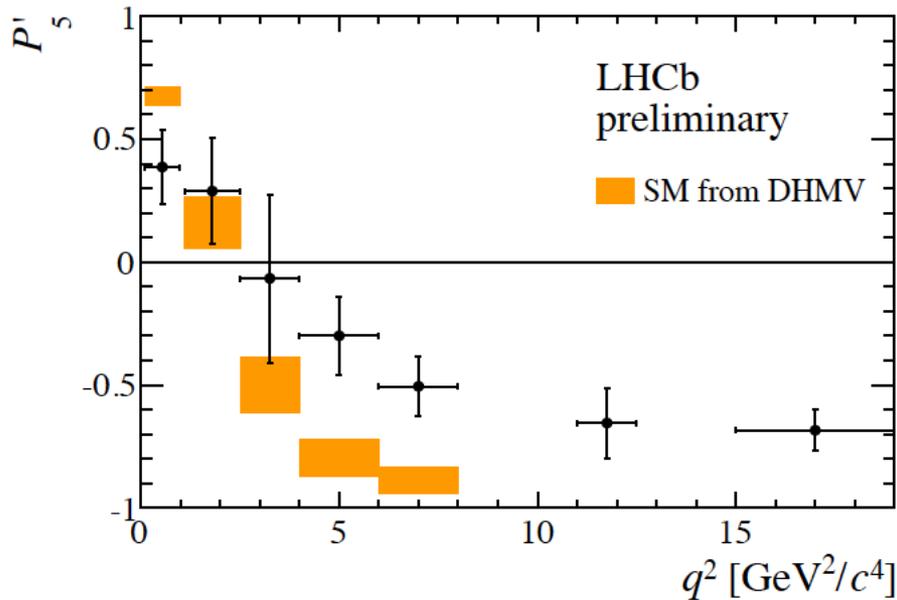
$F_L$  is the longitudinal polarization fraction.

$$\frac{9}{32\pi} \left[ \begin{aligned} & \frac{3}{4}(1 - F_L) \sin^2 \vartheta_K + F_L \cos^2 \vartheta_K \\ & + \frac{1}{4}(1 - F_L) \sin^2 \vartheta_K \cos 2\vartheta_L \\ & - F_L \cos^2 \vartheta_K \cos 2\vartheta_L + S_3 \sin^2 \vartheta_K \sin^2 \vartheta_L \cos 2\phi \\ & + S_4 \sin 2\vartheta_K \sin 2\vartheta_L \cos \phi + \boxed{\phantom{S_5 \sin 2\vartheta_K \sin 2\vartheta_L \sin \phi}} \\ & + \boxed{\phantom{S_5 \sin 2\vartheta_K \sin 2\vartheta_L \sin \phi}} + S_7 \sin 2\vartheta_K \sin \vartheta_L \sin \phi \\ & + \boxed{\phantom{S_5 \sin 2\vartheta_K \sin 2\vartheta_L \sin \phi}} \end{aligned} \right]$$

*Introduce  $P_{4,5} = S_{4,5} / \text{sqrt}[F_L(1 - F_L)]$  to reduce dependence on form factors*

# New LHCb $3fb^{-1}$ results on $B \rightarrow K^* \mu^+ \mu^- (q^2)$

Angular Asymmetries based on  $2398 \pm 57$  signal events



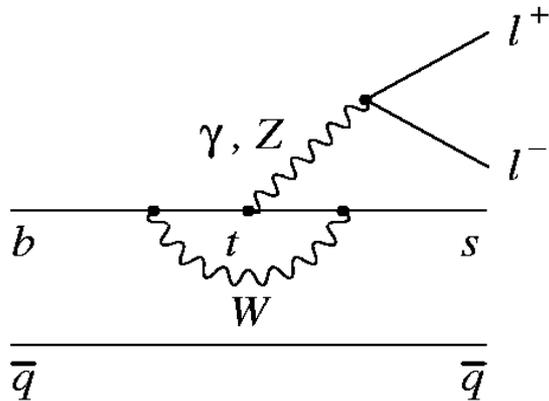
“The  $P_5'$  measurements are only compatible with the SM prediction at a level of  $3.7\sigma$ .....A mild tension can also be seen in the  $A_{FB}$  distribution, where the measurements are systematically  $\leq 1\sigma$  below the SM prediction in the region  $1.1 < q^2 < 6.0$  GeV<sup>2</sup>”

Blank regions are the  $J/\psi$  and  $\psi'$  vetoes

# *New LHCb results on $B \rightarrow K^* \mu^+ \mu^- (q^2)$*

*Is HEP History repeating itself?* [*Make sure this is not a tricky SM form factor effect.*]

Why does NP appear first in this mode (and not others) ?



Possible answer: All the heavy particles of the SM ( $t$ ,  $W$ ,  $Z$ ) and maybe NP (except the Higgs) appear here. Sensitive to NP via interference (linear effects).

# Theory issues on $B \rightarrow K^* \mu^+ \mu^- (q^2)$

*“To better understand the compatibility of the data with the SM a global analysis of the data, taking into account correlations between observables is necessary”-LHCb conference paper.*

- ➔ Check dependence on light-cone form factors (some checks already done by Lattice QCD group, or *fit form factors from data a la Mandal and Sinha and check for consistency*)
- ➔ Check binning effects (do an unbinned analysis of NP)
- ➔ Can tails of large  $B \rightarrow K^* [c\text{-}c\text{-bar}]$  produce the anomalies found in the angular distributions ?



# $B \rightarrow K^* l^+ l^-$ form factor ratios determined from data disagree with theory

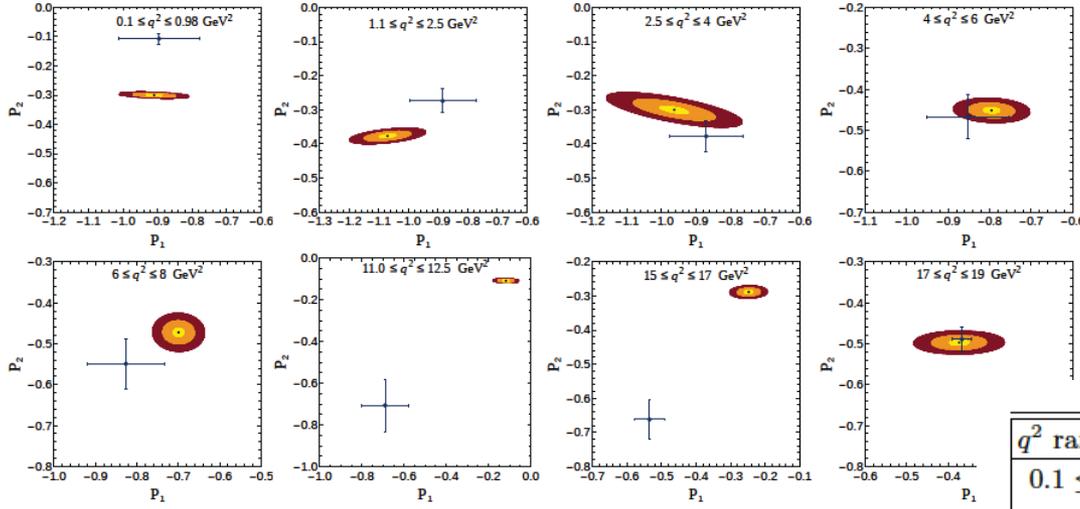


FIG. 1. (color online). The allowed region for  $P_1$  versus  $P_2$  plane. The innermost yellow (lightest), the middle orange and outer most red (dark) contours represent  $1\sigma$ ,  $3\sigma$  and  $5\sigma$  regions, respectively. The theoretical estimates Ref. [7] for  $q^2 \leq 8 \text{ GeV}^2$  and Ref. [10] for  $q^2 \geq 11 \text{ GeV}^2$  are shown as points with error bars. In most cases, the agreement between the theoretical values and those obtained from data. However, for the ranges  $0.1 \leq q^2 \leq 11.0 \leq q^2 \leq 12.5 \text{ GeV}^2$  and  $15 \leq q^2 \leq 17 \text{ GeV}^2$  there are significant disagreements.

Three form factors here

$q^2$ range in $\text{GeV}^2$	$V(q^2)$	$A_1(q^2)$	$A_{12}(q^2)$
$0.1 \leq q^2 \leq 0.98$	$0.704 \pm 0.404$ ( $0.81\sigma$ )	$0.538 \pm 0.309$ ( $0.79\sigma$ )	$0.246 \pm 0.141$ ( $1.27\sigma$ )
$1.1 \leq q^2 \leq 2.5$	$0.624 \pm 0.081$ ( $2.48\sigma$ )	$0.384 \pm 0.051$ ( $1.42\sigma$ )	$0.331 \pm 0.052$ ( $0.72\sigma$ )
$2.5 \leq q^2 \leq 4.0$	$0.318 \pm 0.185$ ( $0.70\sigma$ )	$0.204 \pm 0.119$ ( $0.89\sigma$ )	$0.270 \pm 0.177$ ( $1.56\sigma$ )
$4.0 \leq q^2 \leq 6.0$	$0.556 \pm 0.026$ ( $1.42\sigma$ )	$0.398 \pm 0.020$ ( $2.02\sigma$ )	$0.359 \pm 0.032$ ( $1.28\sigma$ )
$6.0 \leq q^2 \leq 8.0$	$0.597 \pm 0.017$ ( $0.83\sigma$ )	$0.437 \pm 0.014$ ( $2.74\sigma$ )	$0.394 \pm 0.022$ ( $2.18\sigma$ )
$11.0 \leq q^2 \leq 12.5$	$0.172 \pm 0.006$ ( $5.65\sigma$ )	$0.539 \pm 0.027$ ( $2.43\sigma$ )	$0.462 \pm 0.028$ ( $2.82\sigma$ )
$15.0 \leq q^2 \leq 17.0$	$0.713 \pm 0.004$ ( $6.25\sigma$ )	$0.638 \pm 0.026$ ( $3.36\sigma$ )	$0.505 \pm 0.016$ ( $4.64\sigma$ )
$17.0 \leq q^2 \leq 19.0$	$1.936 \pm 0.007$ ( $4.38\sigma$ )	$0.678 \pm 0.025$ ( $3.82\sigma$ )	$0.498 \pm 0.014$ ( $4.64\sigma$ )

TABLE I. The form factor values obtained from fit to  $3 \text{ fb}^{-1}$  of LHCb data [4]. Round brackets indicate the standard deviation between fitted values and theoretical estimates [7, 10]. We find significant discrepancies for several values, especially for the large  $q^2$  region.

$B \rightarrow K^* l^+ l^-$  form factor ratios determined from data are inconsistent: violate HQET equalities at  $q^2_{\max}$

$q^2$ range in $\text{GeV}^2$	$u_0$	$u_{\parallel}$
$15 \leq q^2 \leq 17$	$0.001 \pm 0.015$	$0.013$
$17 \leq q^2 \leq 19$	$0.137 \pm 0.013$	$0.002$
$15 \leq q^2 \leq 19$	$0.068 \pm 0.005$	$0.002$

TABLE II. The values of  $u_0$ ,  $u_{\parallel}$  and  $u_{\perp}$  for the  $B \rightarrow K^* l^+ l^-$  data at  $3 \text{ fb}^{-1}$  of LHCb data [4]. In large  $q^2$  region, the equality  $u_0 = u_{\parallel} = u_{\perp}$  is expected to hold in the absence of charm loop contributions with negligible value of  $u_{\parallel}$  for the larger  $q^2$  bin is unexpected. Significant discrepancies are observed between the values of  $u_{\perp}$  and  $u_0$ .

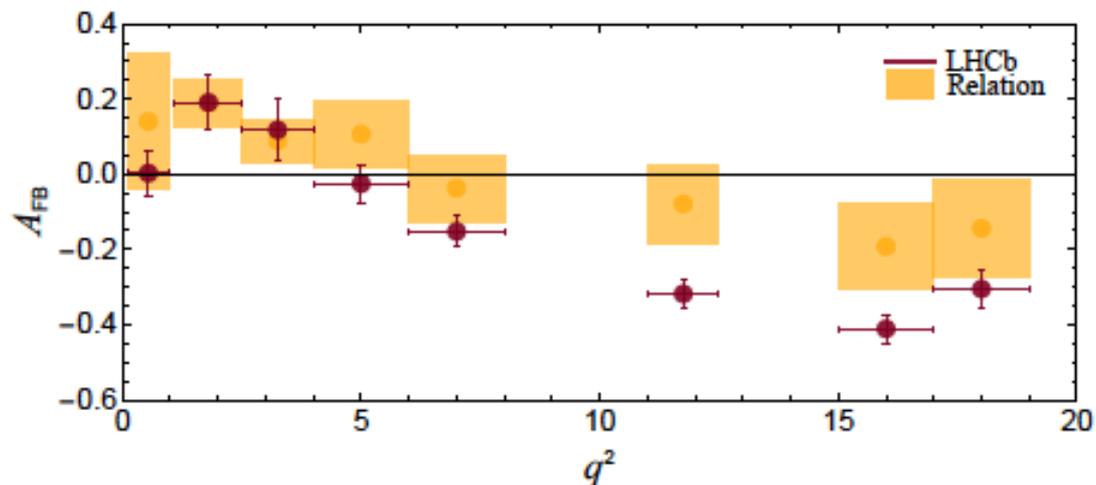


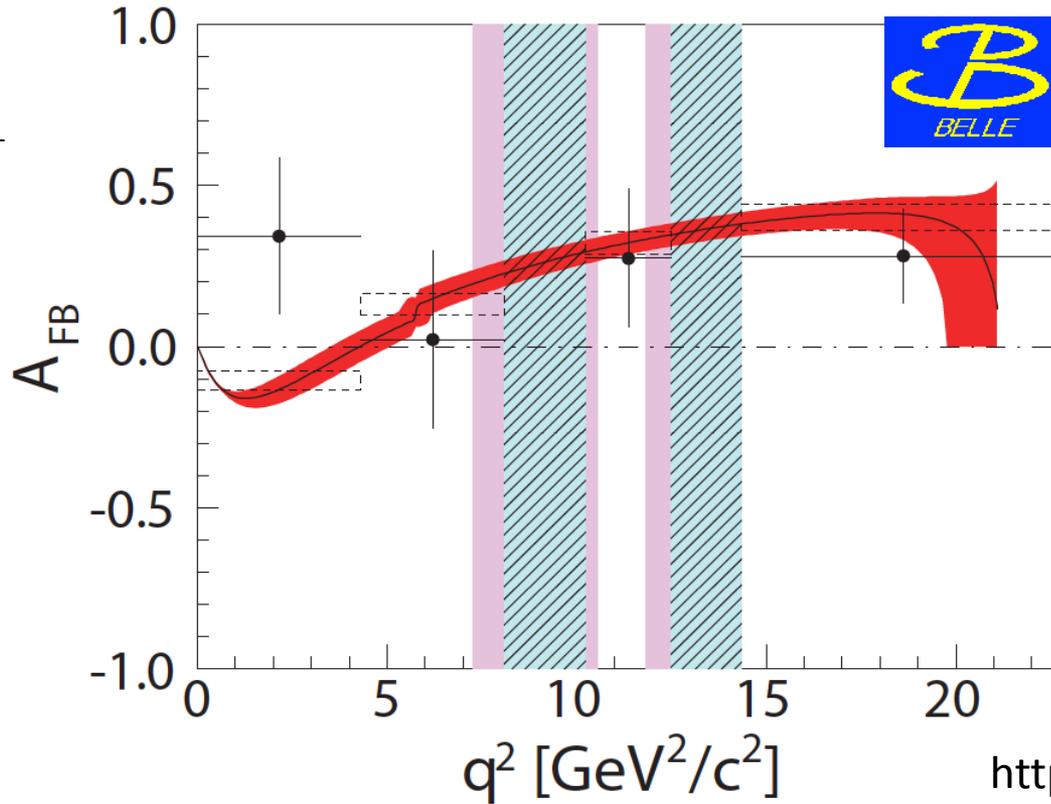
FIG. 3. (color online). The mean values and  $1\sigma$  uncertainties for  $A_{\text{FB}}$  calculated using the ‘Relation’ (Eq. (17)) are shown by yellow (light) bands. We emphasize that these bands are derived using only experimentally measured observables and do not depend on any theoretical input. The error bars in red (dark) correspond to the experimentally measured central values and errors in  $A_{\text{FB}}$  for the respective  $q^2$  bins. See text for details.

*Cannot recover by adding resonances or large non-perturbative effects. But NP works.*

Confirmation and more data is needed to close the case

# $A_{FB}(q^2)$ for Inclusive $b \rightarrow s l^+ l^-$

~301  $b \rightarrow s l^+ l^-$   
signal events



No form factors

Where is the zero crossing ?

Precise result useful  
for NP diagnosis

<http://arxiv.org/abs/1402.7134>

TABLE II. Fit results for the four  $q^2$  bins. For  $\mathcal{A}_{FB}$ , the first uncertainty is statistical and the second uncertainty is systematic.  $\mathcal{A}_{FB}$  values predicted by the SM [4, 7] are also shown with systematic uncertainties. For the signal yields, only statistical uncertainties are shown. The uncertainties of  $\alpha$  and  $\beta$  are due to the statistical uncertainties of the MC.

	1st bin	2nd bin	3rd bin	4th bin
$q^2$ range [ $\text{GeV}^2/c^2$ ]	[0.2,4.3]	[4.3,7.3]	[10.5,11.8]	[14.3, 25.0]
$(B \rightarrow X_s e^+ e^-)$		[4.3,7.3]	[10.5,11.8]	[14.3, 25.0]
$(B \rightarrow X_s \mu^+ \mu^-)$		[4.3,8.1]	[10.2,12.5]	
$\mathcal{A}_{FB}$	$0.34 \pm 0.24 \pm 0.02$	$0.04 \pm 0.31 \pm 0.05$	$0.28 \pm 0.21 \pm 0.01$	$0.28 \pm 0.15 \pm 0.01$
$\mathcal{A}_{FB}$ (theory)	$-0.11 \pm 0.03$	$0.13 \pm 0.03$	$0.32 \pm 0.04$	$0.40 \pm 0.04$
$N_{sig}^{ee}$	$45.6 \pm 10.9$	$30.0 \pm 9.2$	$25.0 \pm 7.0$	$39.2 \pm 9.6$
$N_{sig}^{\mu\mu}$	$43.4 \pm 9.2$	$23.9 \pm 10.4$	$30.7 \pm 9.9$	$62.8 \pm 10.4$
$\alpha^{ee}$	$1.289 \pm 0.004$	$1.139 \pm 0.003$	$1.063 \pm 0.003$	$1.121 \pm 0.003$
$\alpha^{\mu\mu}$	$2.082 \pm 0.010$	$1.375 \pm 0.003$	$1.033 \pm 0.003$	$1.082 \pm 0.003$
$\beta$	1.000	$1.019 \pm 0.003$	$1.003 \pm 0.000$	1.000

TABLE I: Projections for the statistical uncertainties on the  $B \rightarrow K^{(*)}\nu\bar{\nu}$  branching fractions.

Mode	$\mathcal{B}$ [ $10^{-6}$ ]	Efficiency	$N_{\text{Backg.}}$	$N_{\text{Sig-exp.}}$	$N_{\text{Backg.}}$	$N_{\text{Sig-exp.}}$	Statistical error	Total Error
		Belle [ $10^{-4}$ ]	711 $\text{fb}^{-1}$ Belle	711 $\text{fb}^{-1}$ Belle	50 $\text{ab}^{-1}$ Belle II	50 $\text{ab}^{-1}$ Belle II		
$B^+ \rightarrow K^+\nu\bar{\nu}$	3.98	5.68	21	3.5	2960	245	23%	24%
$B^0 \rightarrow K_S^0\nu\bar{\nu}$	1.85	0.84	4	0.24	560	22	110%	110%
$B^+ \rightarrow K^{*+}\nu\bar{\nu}$	9.91	1.47	7	2.2	985	158	21%	22%
$B^0 \rightarrow K^{*0}\nu\bar{\nu}$	9.19	1.44	5	2.0	704	143	20%	22%
$B \rightarrow K^*\nu\bar{\nu}$ combined							15%	17%

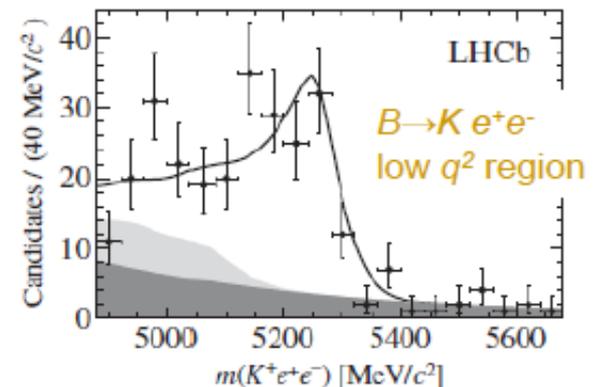
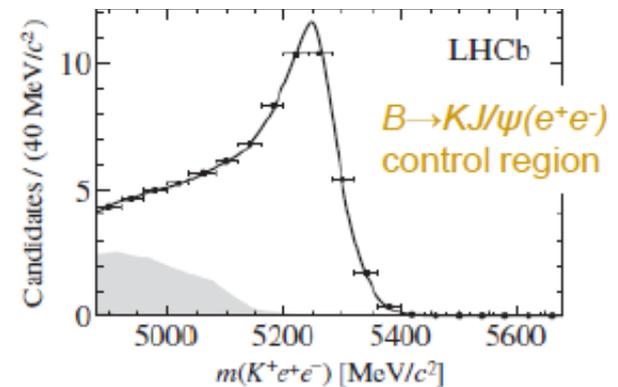
Ans: Verify hint of lepton universality breakdown at **Belle II** (good electron eff)

Control region gives  $R_K$  consistent with unity. Interesting, low  $q^2$  region gives:

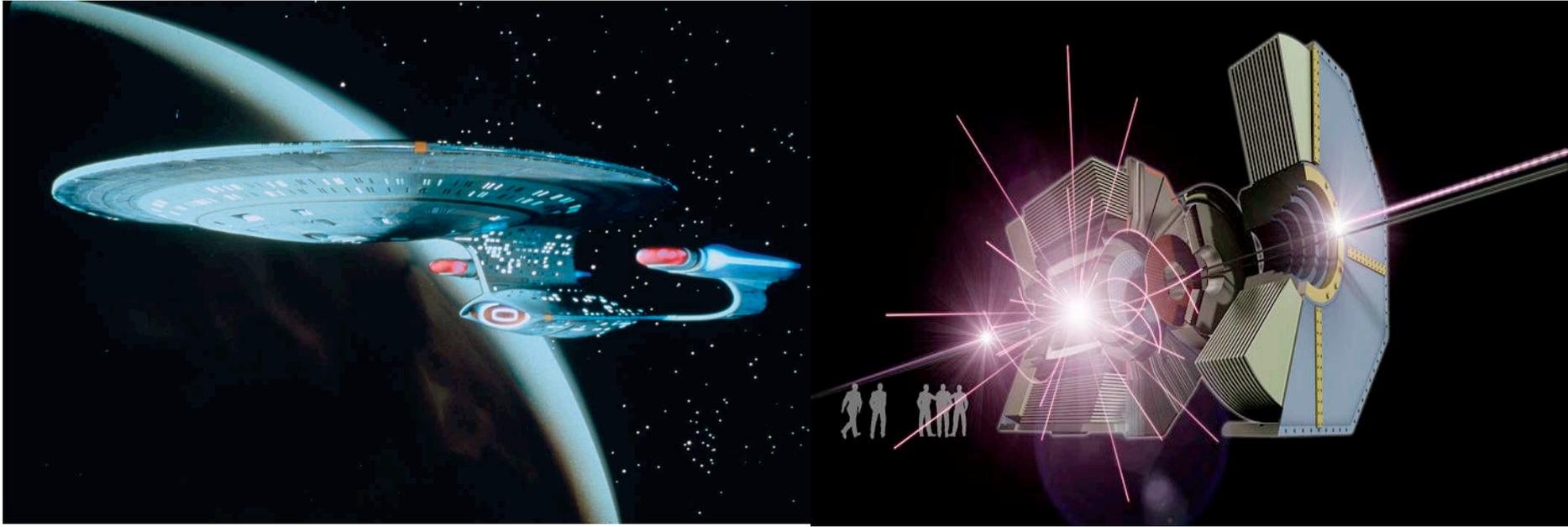
$$R_K = 0.745^{+0.090}_{-0.074}(\text{stat}) \pm 0.036(\text{syst})$$

which is  $2.6\sigma$  from unity,  $3\sigma$  if BaBar included.

R. Aaij et al. (LHCb collab); PRL 113, 151601 (2014)



2014 was the 50<sup>th</sup> anniversary of the discovery of CP violation in the kaon sector [see <http://pprc.qmul.ac.uk/research/50-years-cp-violation>



## The Next Generation

### *Belle II and the LHCb upgrade*

US P5 report (p. v): “Explore the unknown: new particles, interactions, and physical principles”

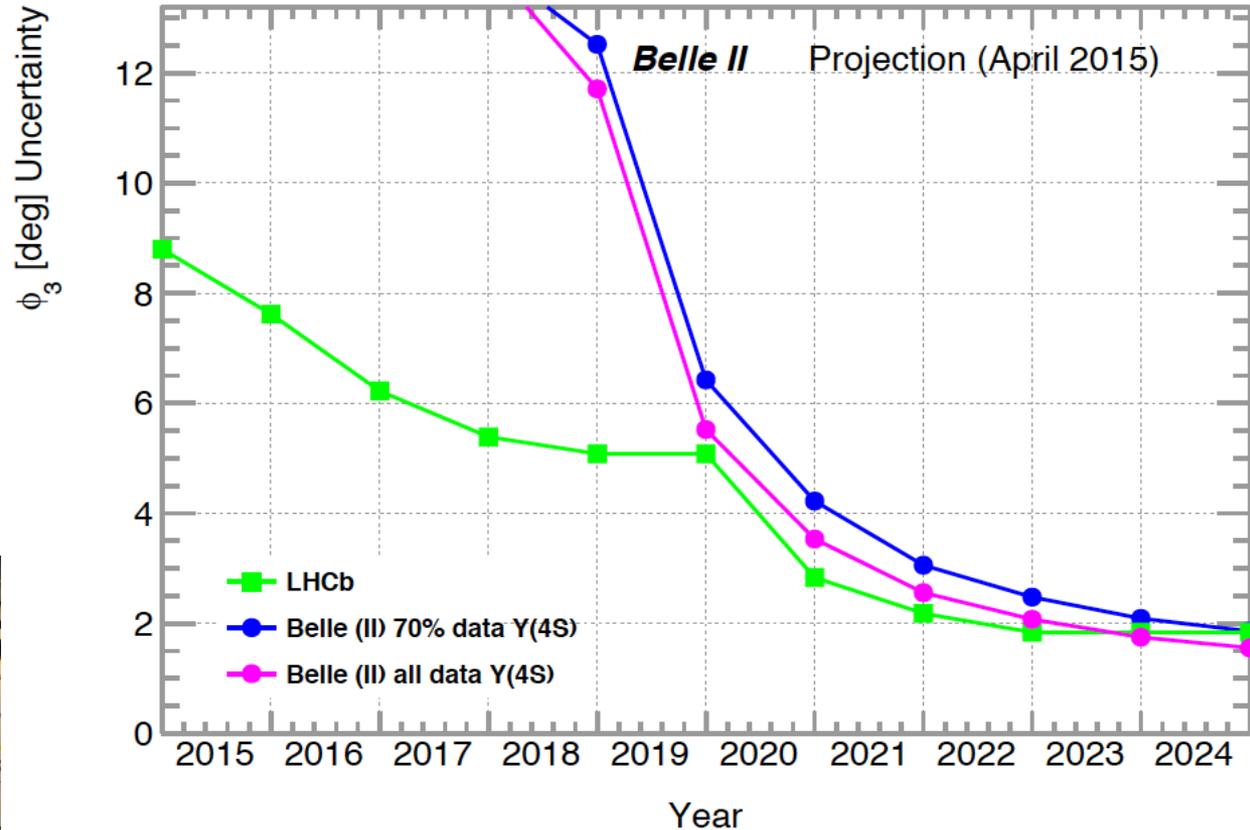
# Physics Reach of Belle II and the LHCb upgrade

## Competition and complementarity



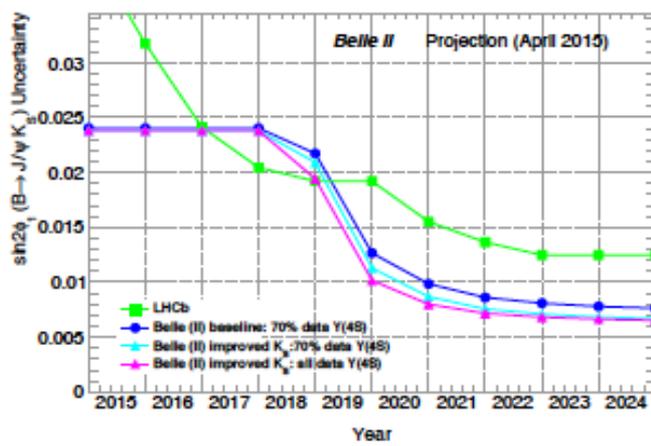
Tofu Gelato ?

Observable	Expected th. accuracy	Expected exp. uncertainty	Facility
CKM matrix			
$ V_{us} $ [ $K \rightarrow \pi \ell \nu$ ]	**	0.1%	<i>K</i> -factory
$ V_{cb} $ [ $B \rightarrow X_c \ell \nu$ ]	**	1%	Belle II
$ V_{ub} $ [ $B \rightarrow X_u \ell \nu$ ]	*	10%	Belle II

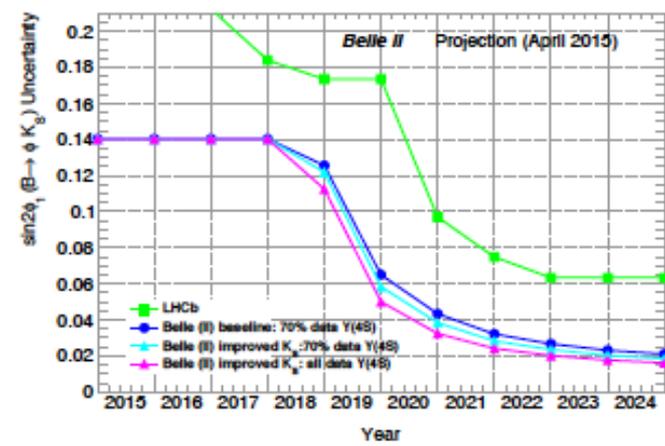


$B(K \rightarrow e \nu) / B(K \rightarrow \mu \nu)$	***	0.1%	<i>K</i> -factory
charm and $\tau$			
$B(\tau \rightarrow \mu \gamma)$	***	$3 \cdot 10^{-9}$	Belle II
$ q/p _D$	***	0.03	Belle II
$arg(q/p)_D$	***	$1.5^\circ$	Belle II

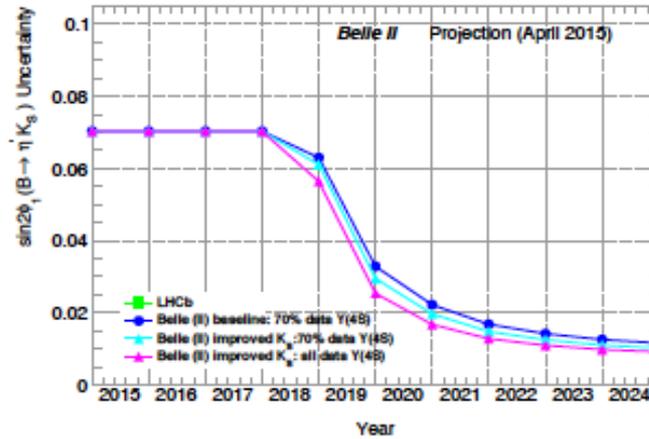
Tight race



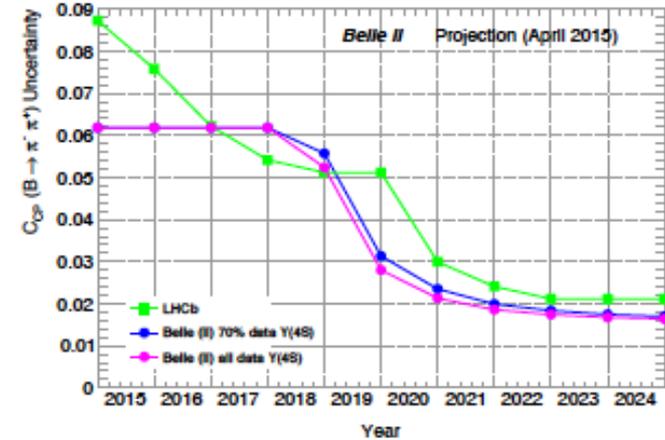
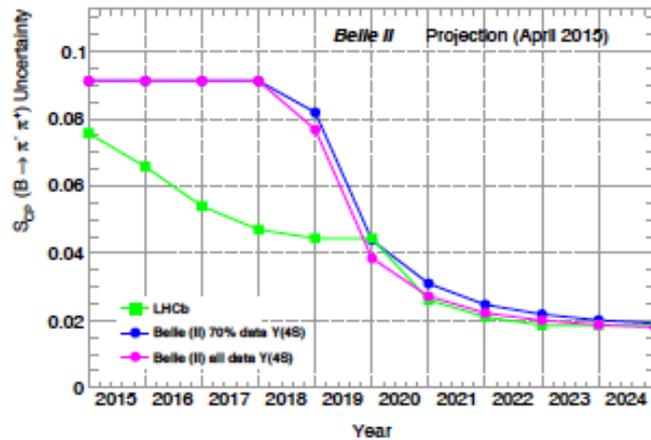
Belle II ahead



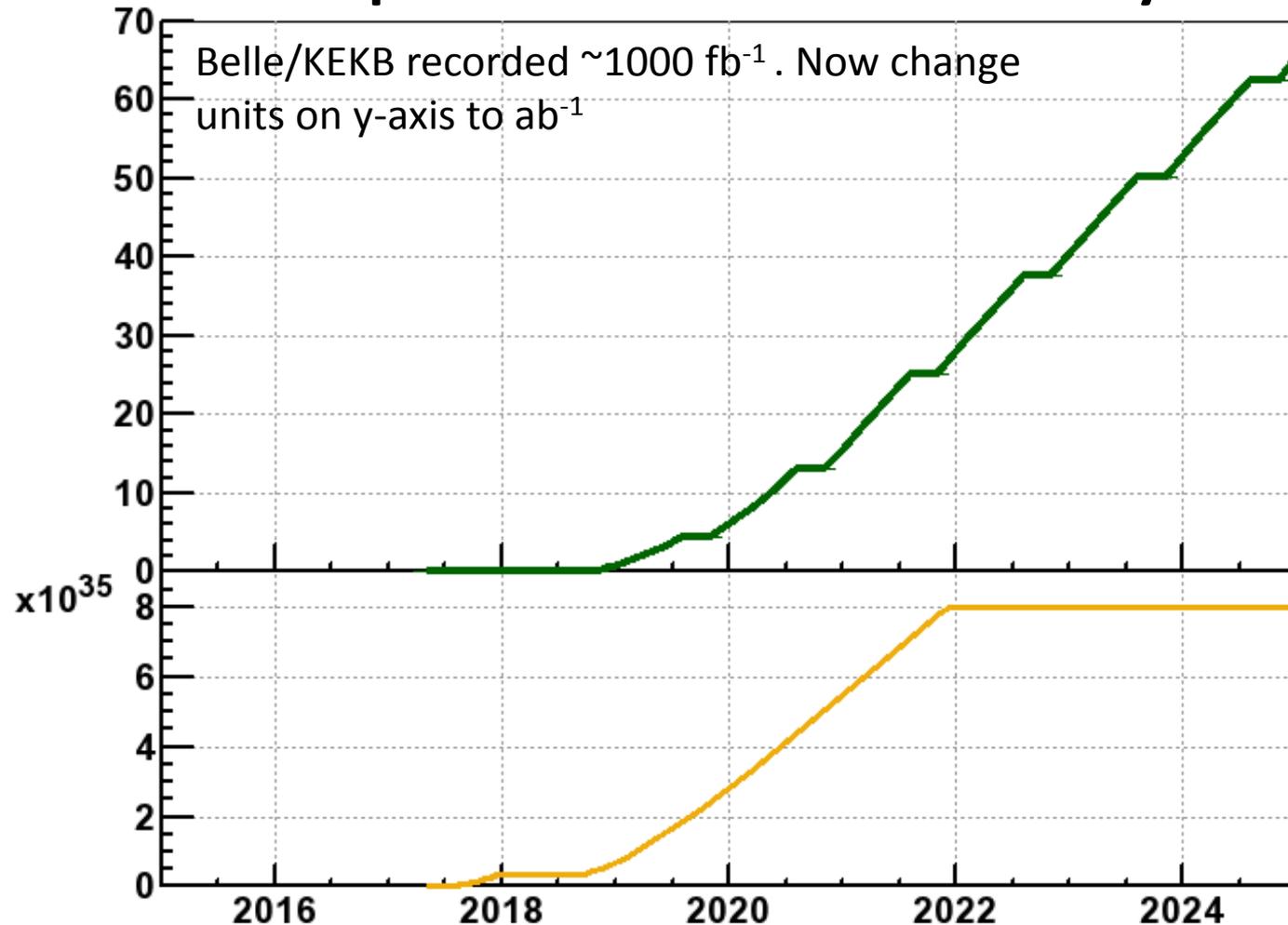
Belle II dominates here



Tight race



# Latest SuperKEKB Luminosity Profile

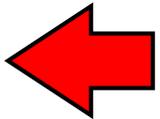


N.B. To realize this steep turn-on, requires close cooperation between Belle II and SuperKEKB [and *international collaboration* on the accelerator].

Also assumes full operation funding profile.

# Compare the Parameters for KEKB and SuperKEKB

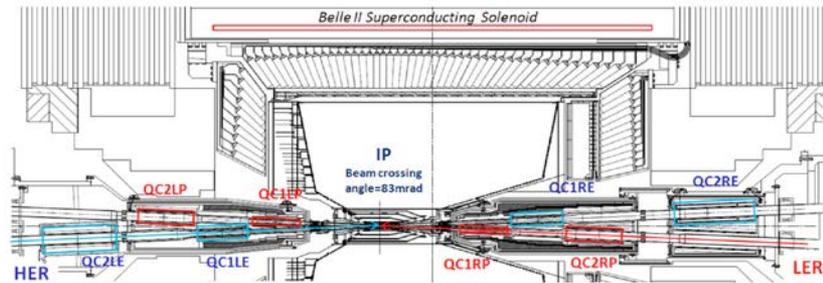
	KEKB Design	KEKB Achieved : with crab	SuperKEKB Nano-Beam
Energy (GeV) (LER/HER)	3.5/8.0	3.5/8.0	4.0/7.0
$\beta_y^*$ (mm)	10/10	5.9/5.9	0.27/0.30
$\beta_x^*$ (mm)	330/330	1200/1200	32/25
$\epsilon_x$ (nm)	18/18	18/24	3.2/5.3
$\epsilon_y / \epsilon_x$ (%)	1	0.85/0.64	0.27/0.24
$\sigma_y$ (mm)	1.9	0.94	0.048/0.062
$\sigma_y$	0.052	0.129/0.090	0.09/0.081
$\sigma_z$ (mm)	4	6 - 7	6/5
$I_{\text{beam}}$ (A)	2.6/1.1	1.64/1.19	3.6/2.6
$N_{\text{bunches}}$	5000	1584	2500
Luminosity ( $10^{34} \text{ cm}^{-2} \text{ s}^{-1}$ )	1	2.11	80



Nano-beams are the key (vertical spot size is  $\sim 50\text{nm}$  !!)

This is not a typo

2015: Basic hardware (except final focus) now in place



New superconducting final focusing magnets near the IP



$e^+$  3.6A

$e^-$  2.6A

## KEKB to SuperKEKB

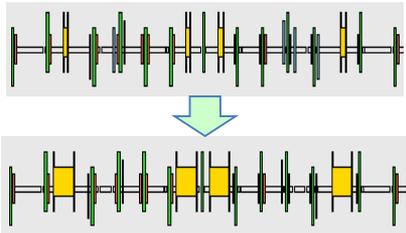
- ◆ Nano-Beam scheme  
extremely small  $\beta_y^*$   
low emittance
- ◆ Beam current X 2

$$L = \frac{\gamma_{\pm}}{2er_e} \left( 1 + \frac{\sigma_y^*}{\sigma_x^*} \left( \frac{I_{\pm} \xi_{\pm y}}{\beta_y^*} \right) \left( \frac{R_L}{R_y} \right) \right)$$

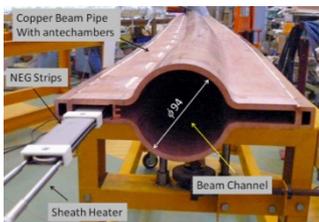
40 times higher luminosity  
 $2.1 \times 10^{34} \rightarrow 8 \times 10^{35} \text{ cm}^{-2} \text{ s}^{-1}$



Redesign the lattice to reduce the emittance (replace short dipoles with longer ones, increase wiggler cycles) (*all magnets installed 8/2014*)



Replace beam pipes with TiN-coated beam pipes with antechambers (*installed*)



DR tunnel

*New  $e^+$  Damping Ring constructed*



Reinforce RF systems for higher beam currents



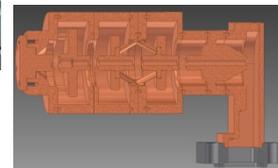
Improve monitors and control system

Injector Linac upgrade

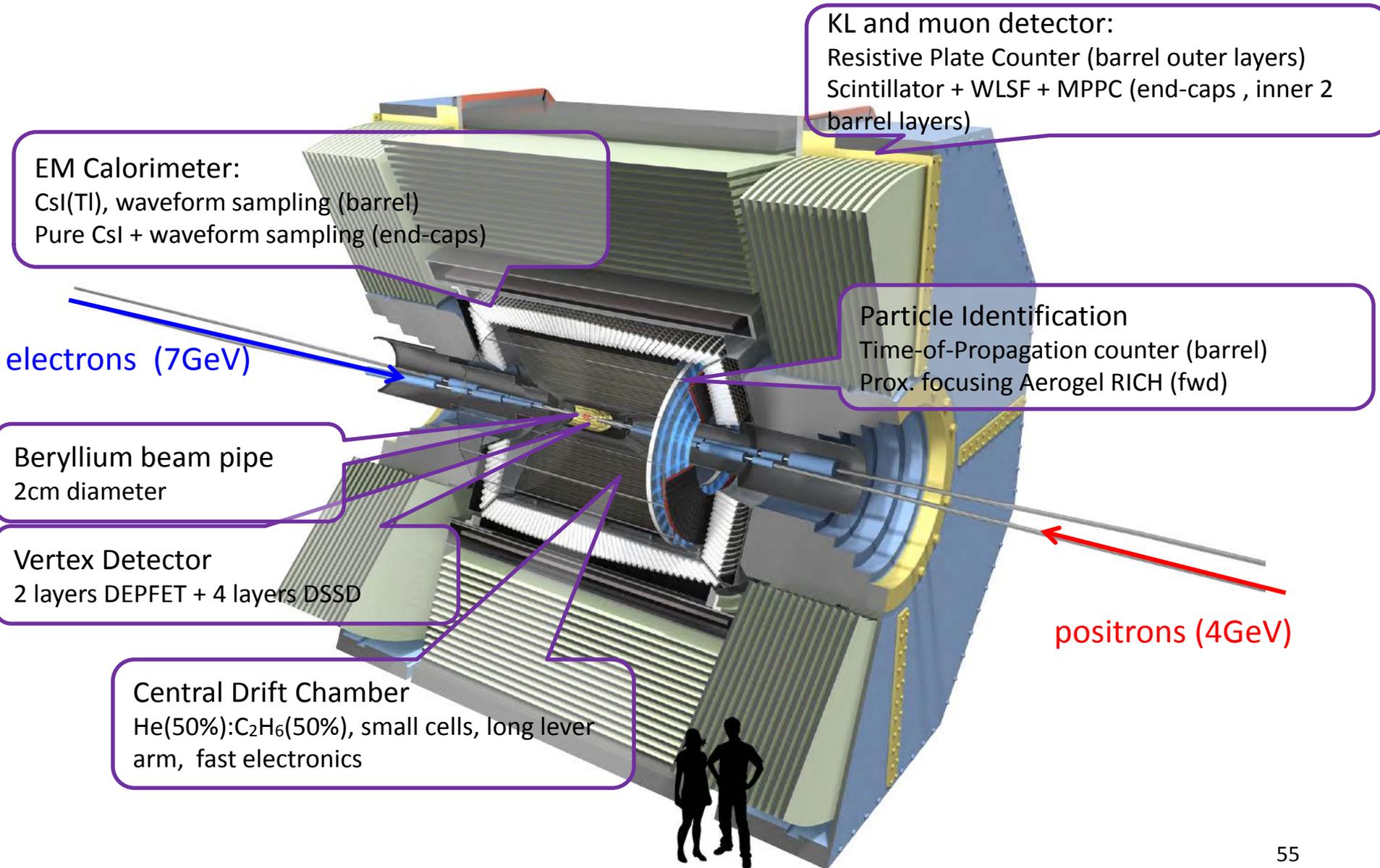
Upgrade positron capture section



Low emittance RF electron gun

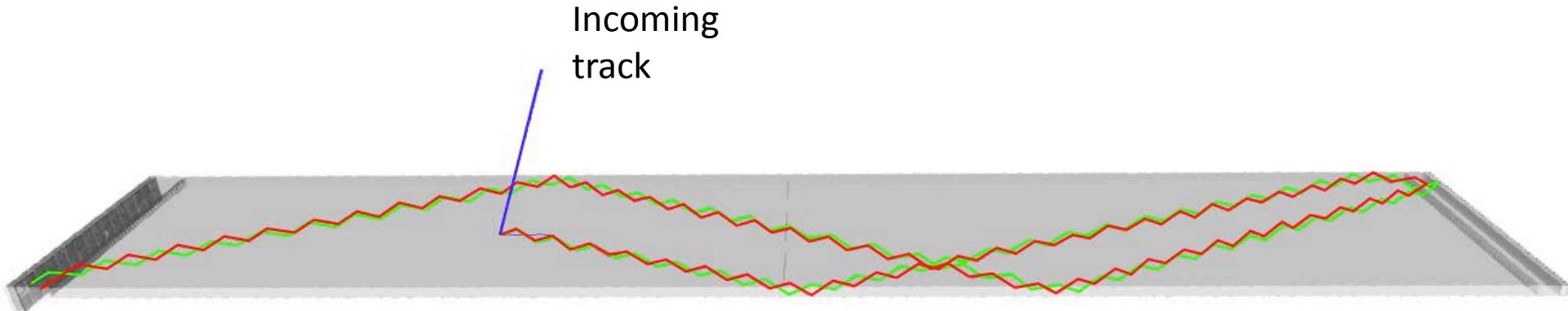


# Belle II Detector

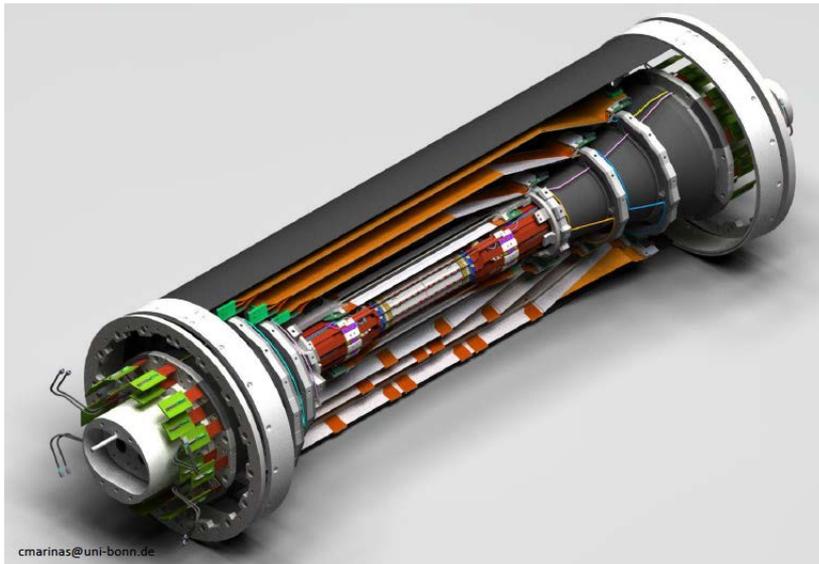


# Barrel PID

A GEANT4 event display of a 2 GeV pion and kaon interacting in a TOP quartz bar. (Japan, US, Slovenia, Italy)



## Vertexing/Inner Tracking



cmarinas@uni-bonn.de

- Beampipe  $r = 10$  mm
- DEPFET pixels (Germany, Czech Republic...)
  - Layer 1  $r = 14$  mm
  - Layer 2  $r = 22$  mm
- DSSD (double sided silicon detectors) FWD/BWD
  - Layer 3  $r = 38$  mm (Australia)<sup>Italy</sup>
  - Layer 4  $r = 80$  mm (India)
  - Layer 5  $r = 115$  mm (Austria)
  - Layer 6  $r = 140$  mm (Japan)

+Poland, Korea

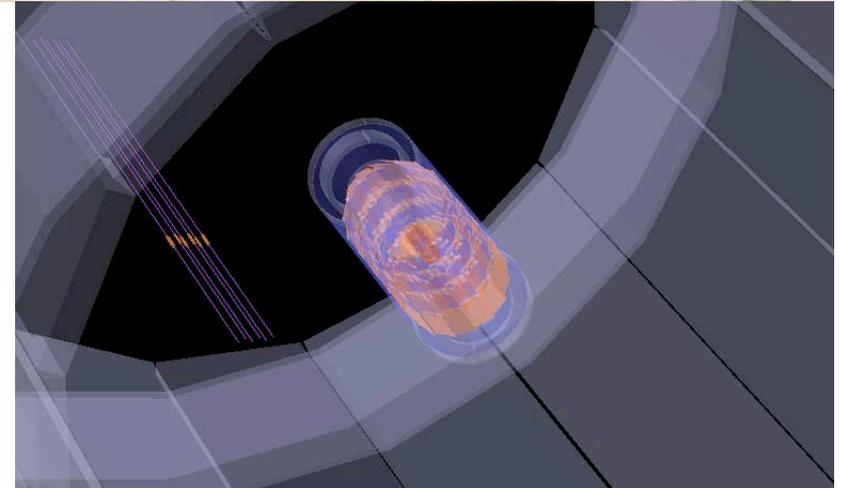


# Highlights of Belle II construction

SuperKEKB hardware is being finalized.



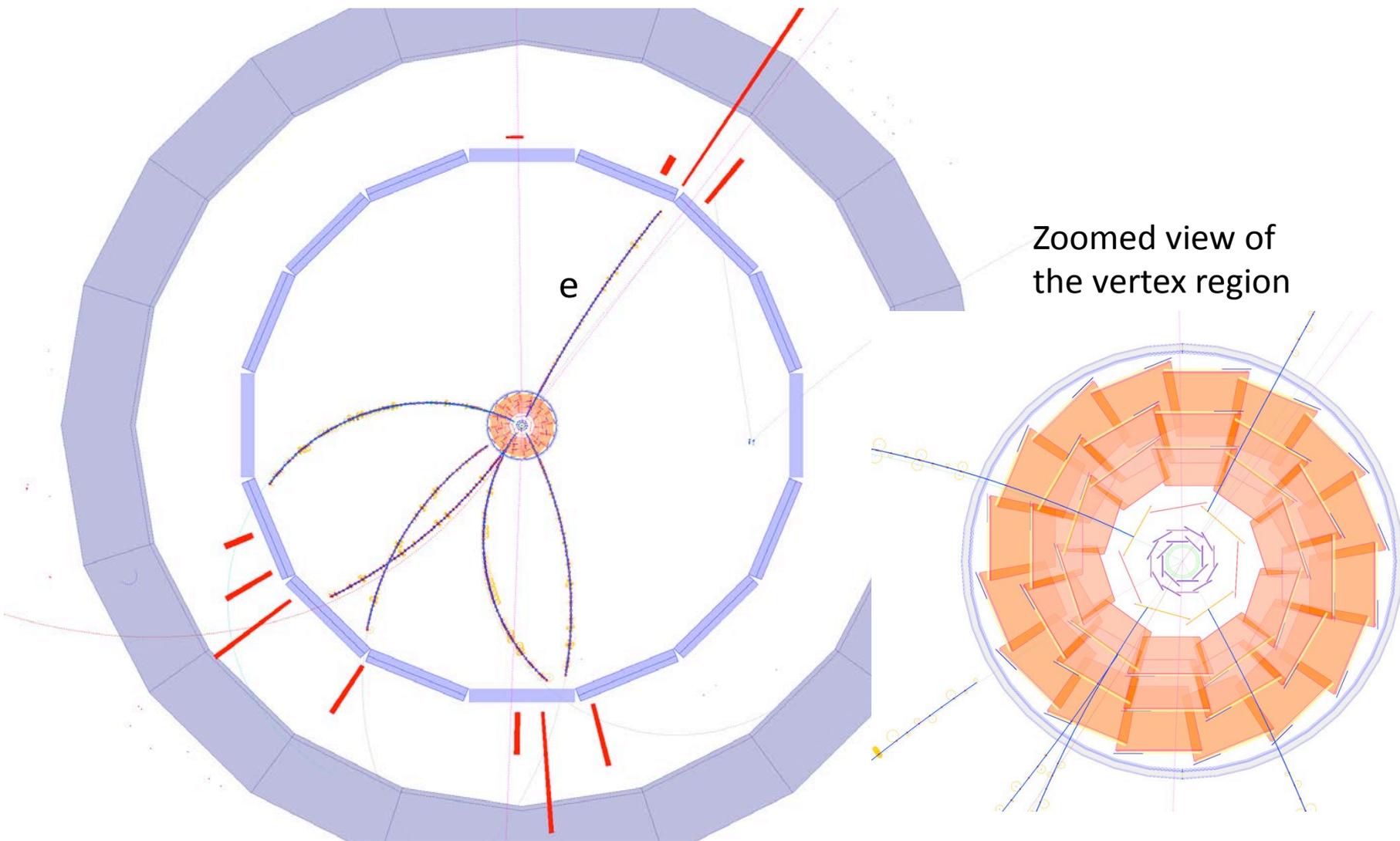
BEAST PHASE I beampipe installed



(a) CDC arriving at Tsukuba Hall; (b) first cosmics with partly instrumented electronics (6 layers)

# “Missing Energy Decay” in a Belle II GEANT4 MC simulation

$B \rightarrow \tau \nu$ ,  $\tau \rightarrow e \nu \nu$        $B \rightarrow D \pi$ ,  $D \rightarrow K \pi \pi \pi$

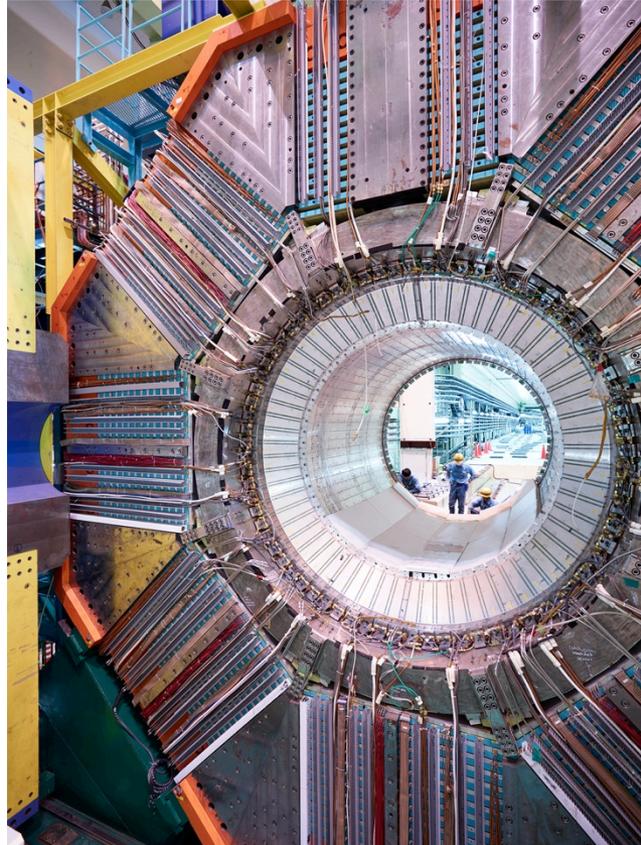


# Conclusion/Next Generation

- The  $e^+e^-$  B factories confirmed that the KM phase is responsible for most of the observed CPV [Physics Nobel Prize 2008]
- Nevertheless, 10-20% NP effects are consistent with all current flavor data.
- LHCb has ruled out large CPV phases from NP in the  $B_s$  sector.
- “Missing energy B decays” provide important high—mass sensitivity to the charged Higgs in the multi-TeV range.
- LHC exps (CMS, LHCb) measured  $B_s \rightarrow \mu^+ \mu^-$
- Angular anomalies in  $B \rightarrow K^* l^+ l^-$  from LHCb with  $3 \text{ fb}^{-1}$
- *Flavor physics is exciting and fundamental. (Did we just find NP ? Is the path for the future ?)*

*SuperKEKB commissioning starts in January. Belle II Physics runs in 2018 and the LHCb upgrade in  $\sim 2020$ . These facilities will inaugurate a new era of flavor physics and the study of CP violation.*

# Backup slides

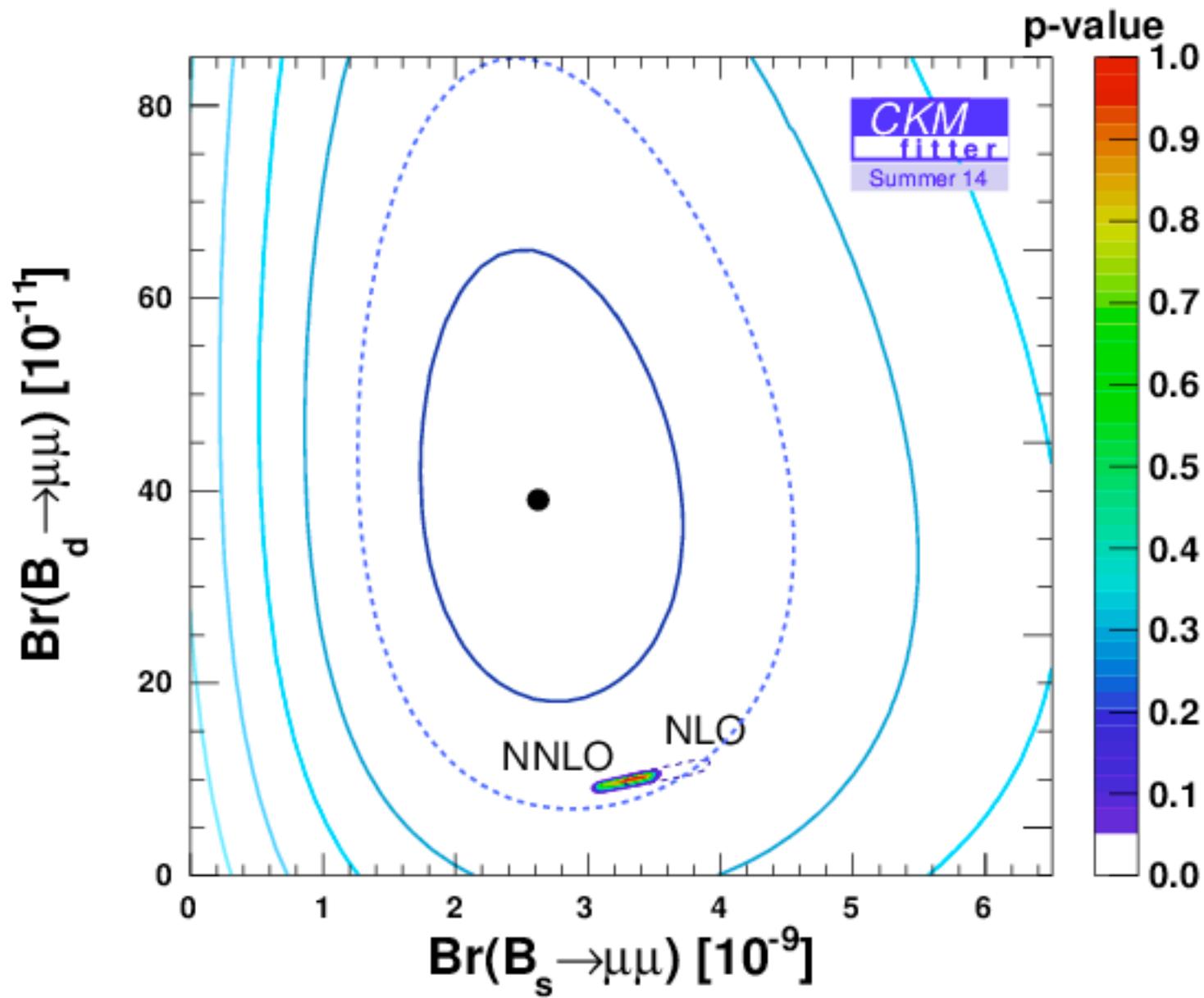


# Updated projections for $B \rightarrow K^{(*)} \nu \bar{\nu}$ modes

TABLE I: Projections for the statistical uncertainties on the  $B \rightarrow K^{(*)} \nu \bar{\nu}$  branching fractions.

Mode	$\mathcal{B} [10^{-6}]$	Efficiency Belle [ $10^{-4}$ ]	$N_{\text{Backg.}}$		$N_{\text{Sig-exp.}}$		Statistical error 50 $\text{ab}^{-1}$	Total Error
			711 $\text{fb}^{-1}$ Belle	711 $\text{fb}^{-1}$ Belle	50 $\text{ab}^{-1}$ Belle II	50 $\text{ab}^{-1}$ Belle II		
$B^+ \rightarrow K^+ \nu \bar{\nu}$	3.98	5.68	21	3.5	2960	245	23%	24%
$B^0 \rightarrow K_S^0 \nu \bar{\nu}$	1.85	0.84	4	0.24	560	22	110%	110%
$B^+ \rightarrow K^{*+} \nu \bar{\nu}$	9.91	1.47	7	2.2	985	158	21%	22%
$B^0 \rightarrow K^{*0} \nu \bar{\nu}$	9.19	1.44	5	2.0	704	143	20%	22%
$B \rightarrow K^* \nu \bar{\nu}$ combined							15%	17%

- [1] A. J. Buras, J. Girrbach-Noe, C. Niehoff and D. M. Straub, JHEP **1502**, 184 (2015) [arXiv:1409.4557 [hep-ph]].
- [2] O. Lutz *et al.* [Belle Collaboration], Phys. Rev. D **87**, no. 11, 111103 (2013) [arXiv:1303.3719 [hep-ex]].
- [3] T. Kuhr, “ $B \rightarrow h^{(*)} \nu \bar{\nu}$ ”, KEK-FF Workshop (2013).



Discussion Topic: What additional *Theoretical Work* is required to determine whether NP is present in B decays ?

Participants: Wolfgang Altmannshofer, Christoph Bobeth, Jorge Martin Camalich, Robert Fleischer, Zoltan Ligeti, Rahul Sinha

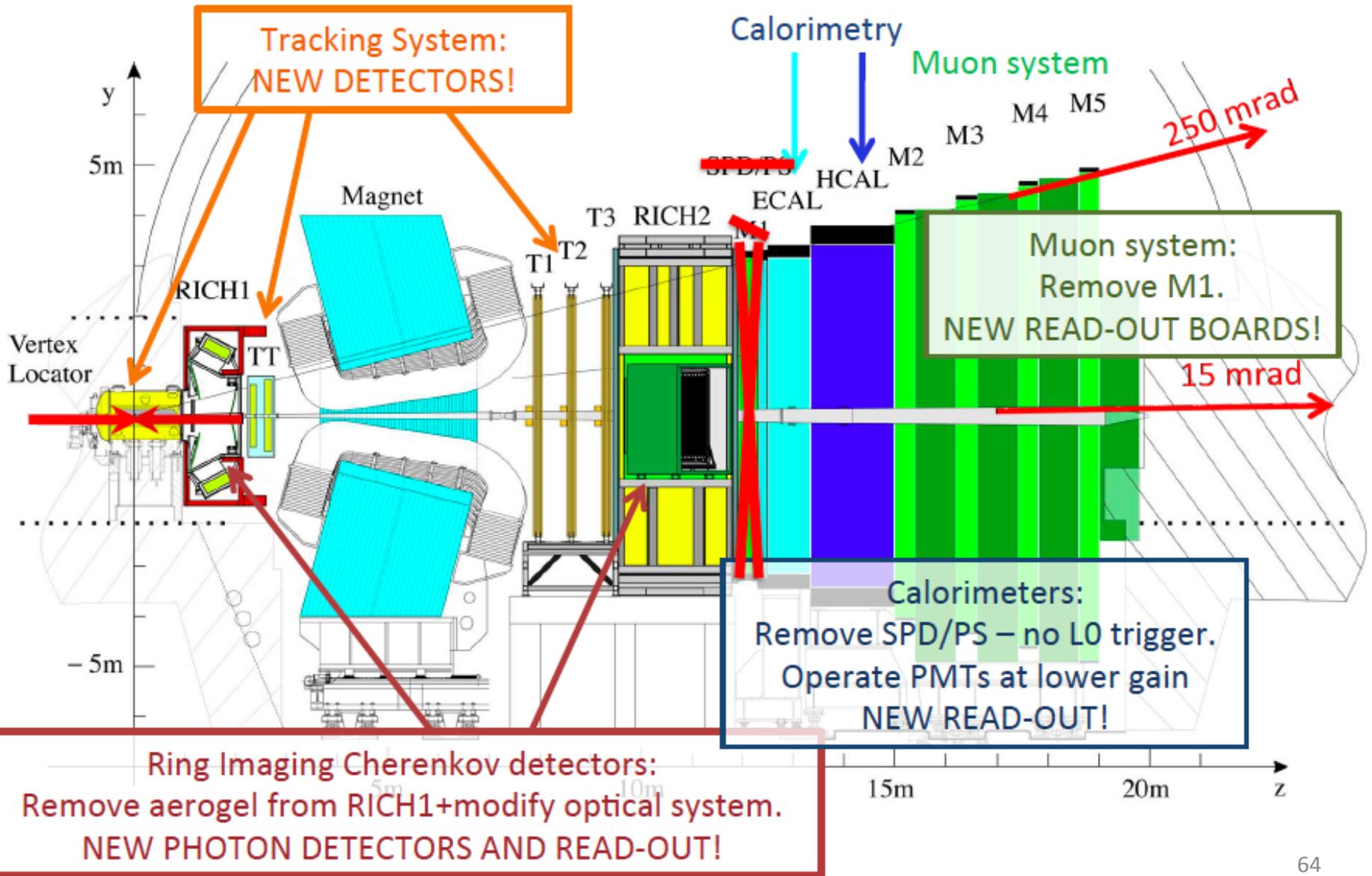


Nagoya  
FPCP15  
roundtable

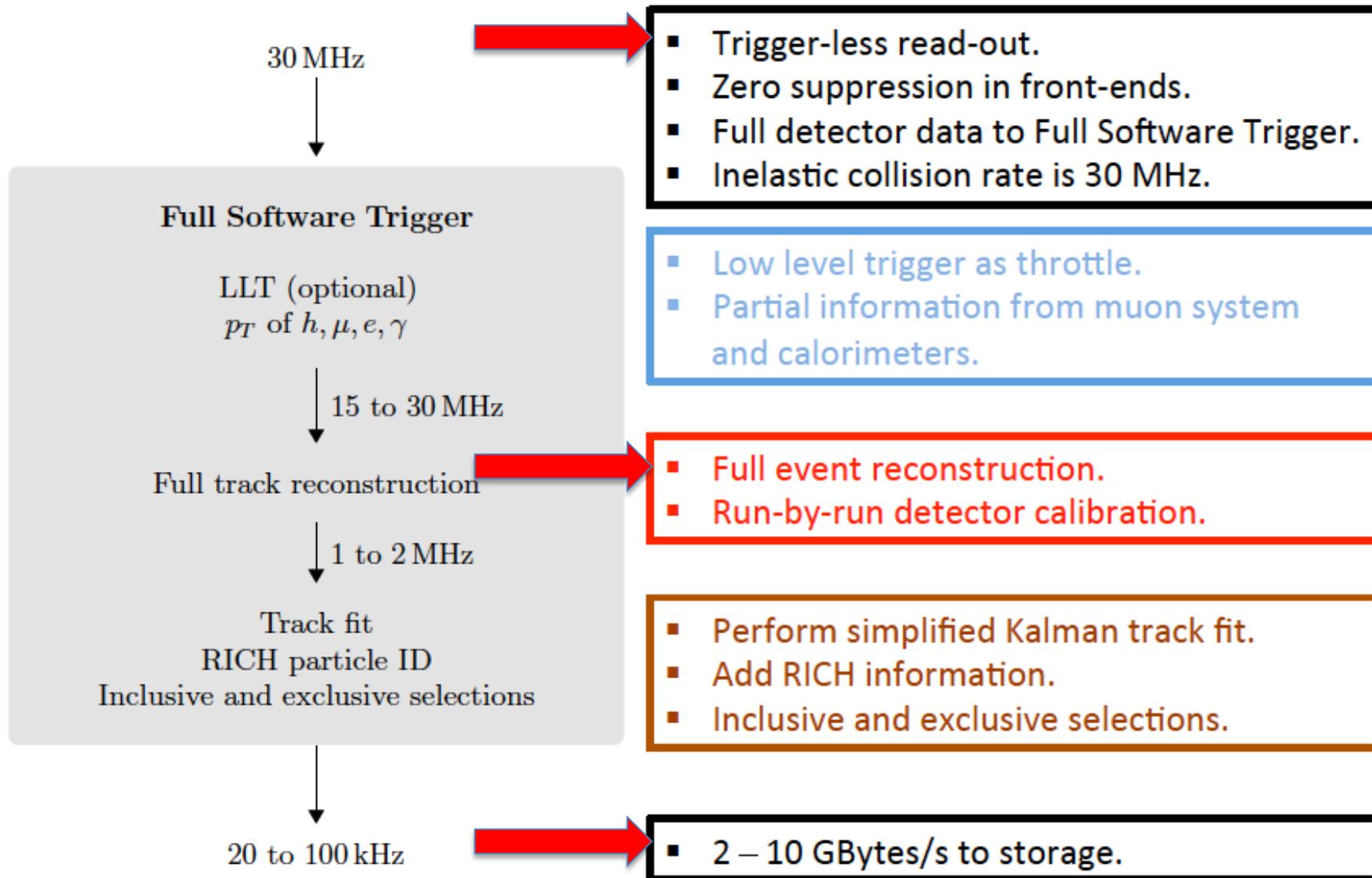
鵜飼  
“ukai”

Cormorant Fishing on the Nagara River during the Edo Period

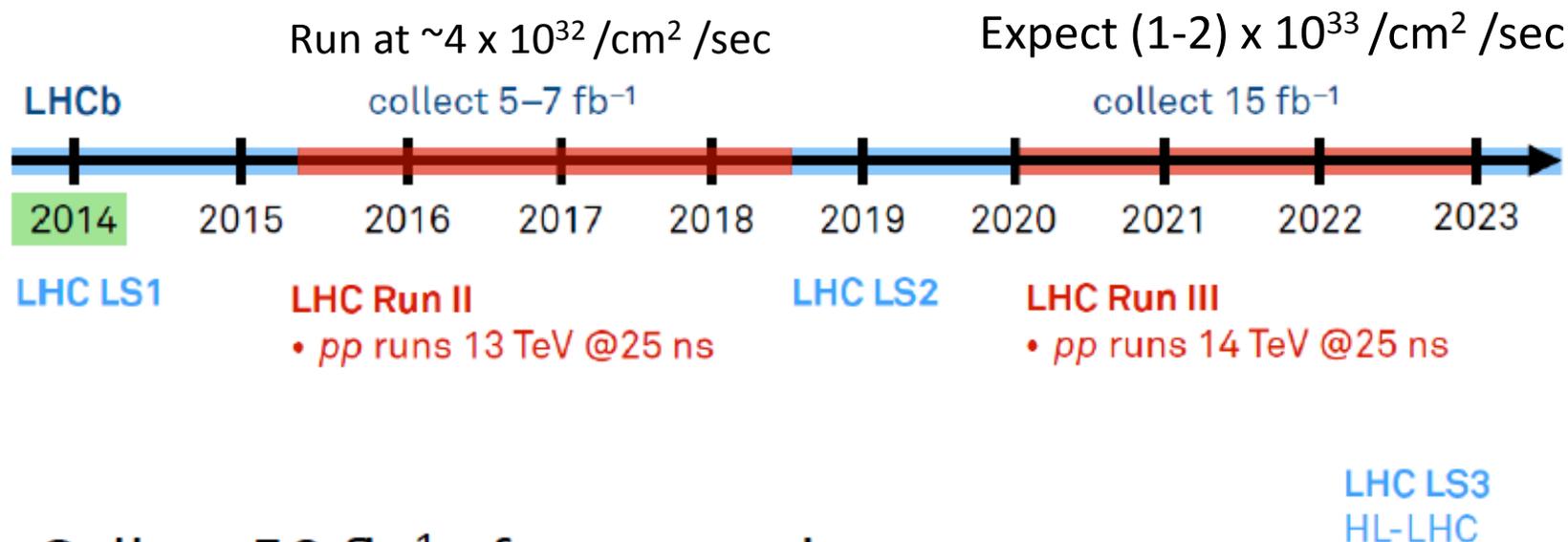
# Upgraded LHCb detector



# LHCb Upgrade: *Key Feature* is Trigger-less readout



# LHCb upgrade timeline



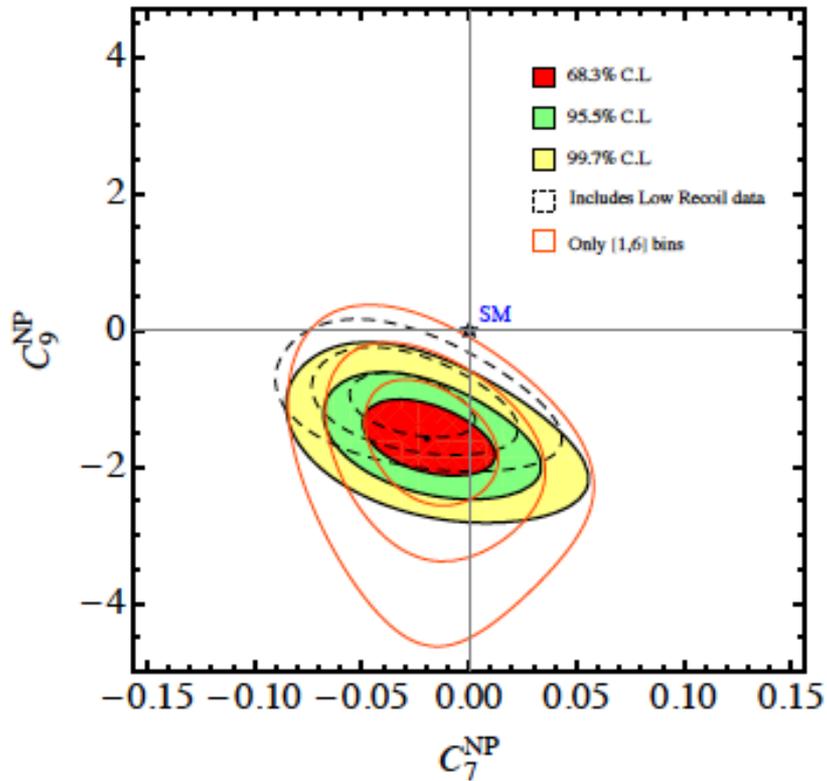
- Collect  $50 \text{ fb}^{-1}$  after upgrade.
- Continue taking data during HL-LHC.

*Upgraded trigger and DAQ is the key feature*

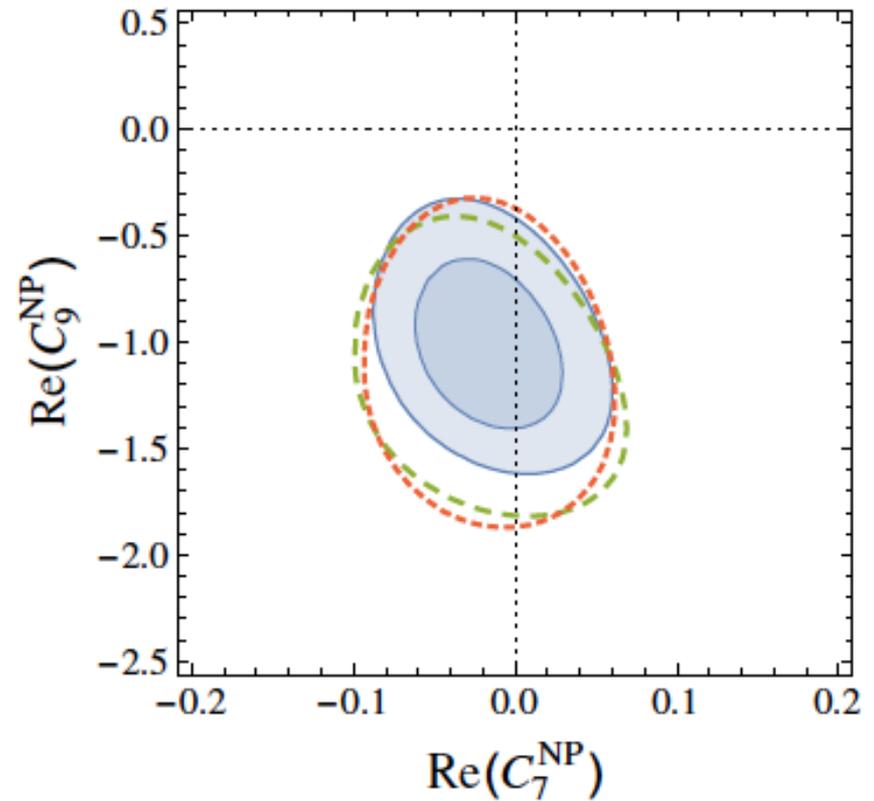
Discussed in Nakada-san's talk

# NP Fits

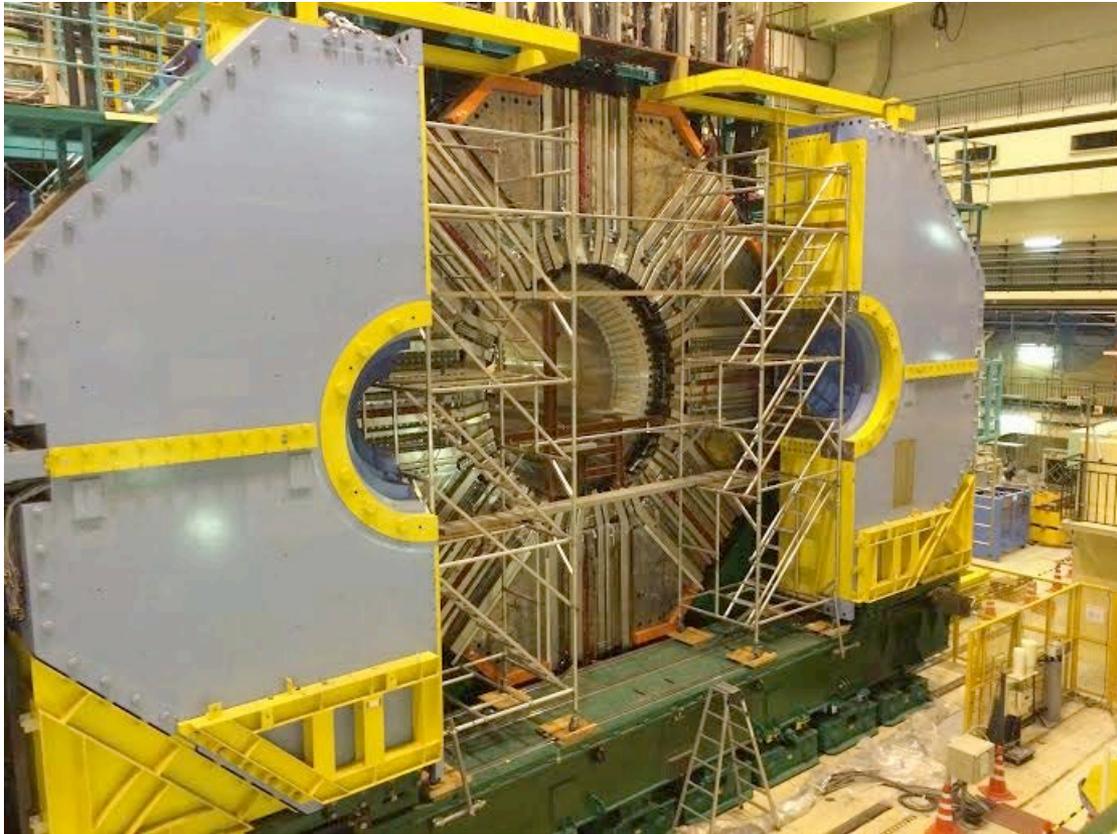
Descotes-Genon, Matias, JV 1307.5683



Altmannshofer, Straub 1503.06199



## *Belle II@Tsukuba Hall, KEK*



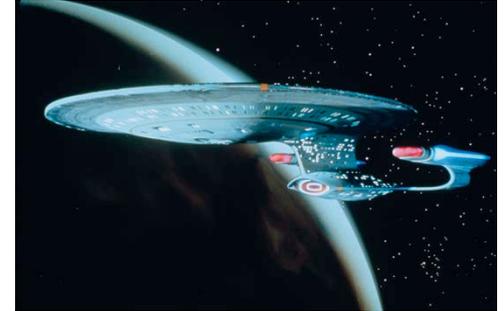
Belle II detector with barrel KLM upgrade as well as forward and backward muon endcap upgrades now installed.



# New Reference *for the Next Generation*

## **The Physics of the B Factories**

<http://arxiv.org/abs/1406.6311>



This work is on the Physics of the B Factories. Part A of this book contains a brief description of the SLAC and KEK B Factories as well as their detectors, BaBar and Belle, and data taking related issues. Part B discusses tools and methods used by the experiments in order to obtain results. The results themselves can be found in Part C.

Comments: 928 pages

Subjects: High Energy Physics - Experiment (hep-ex); High Energy Physics - Phenomenology (hep-ph)

Report number: SLAC-PUB-15968, KEK Preprint 2014-3

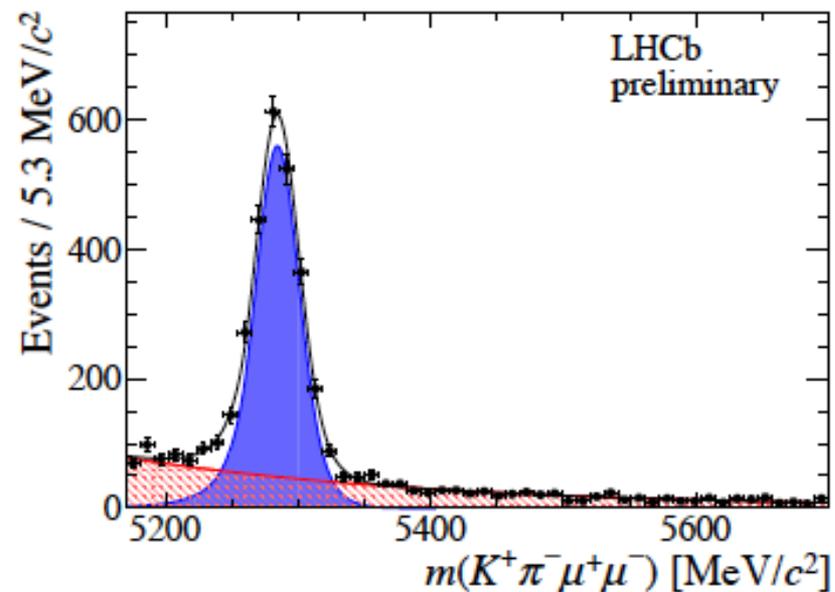
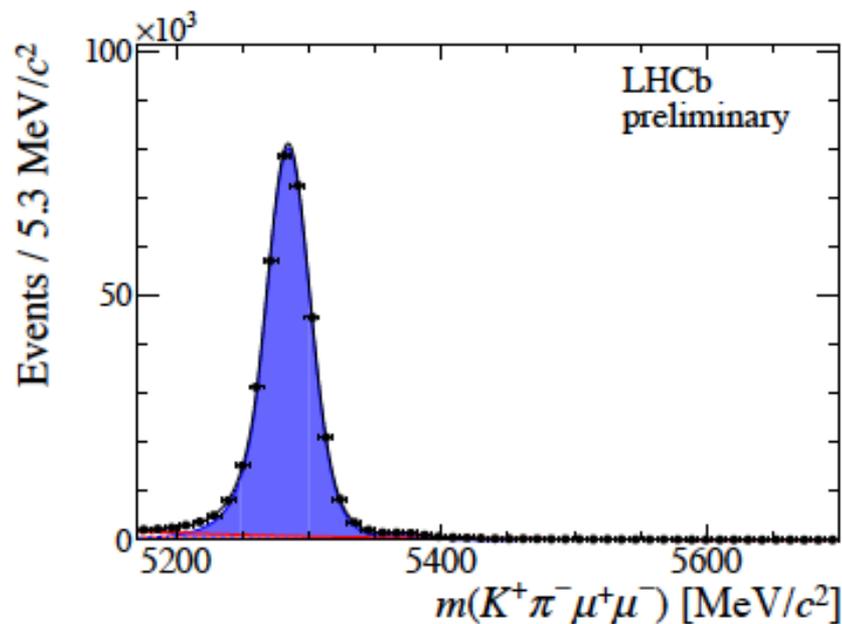
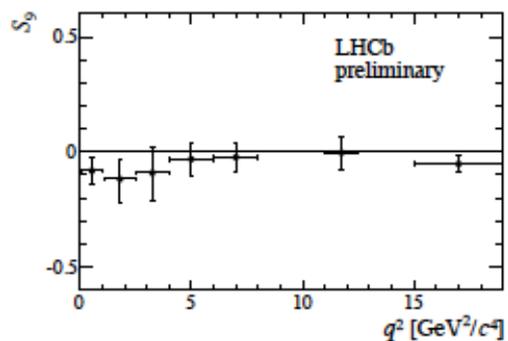
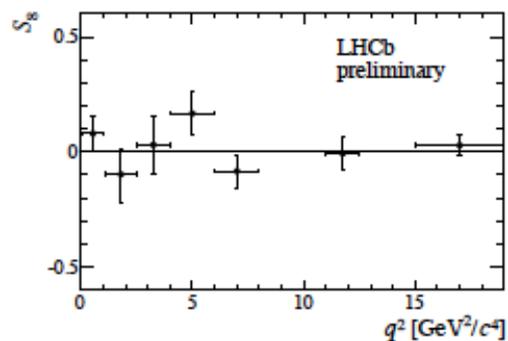
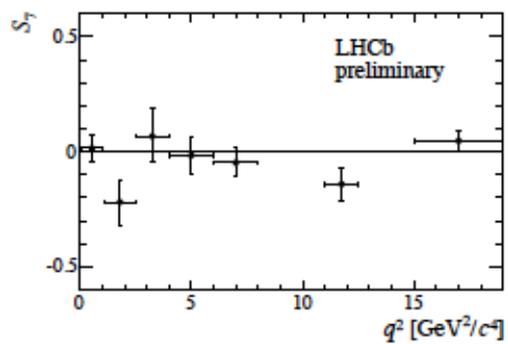
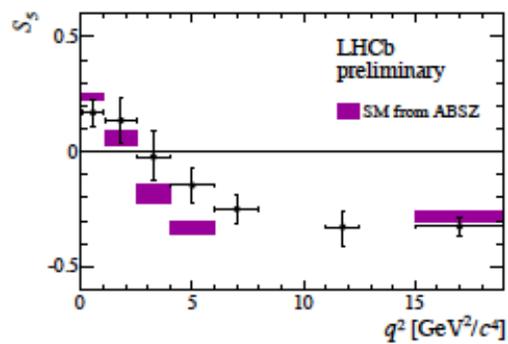
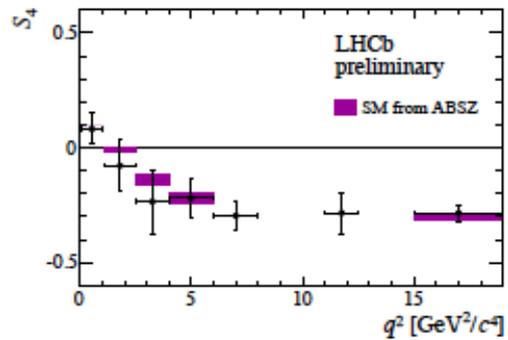
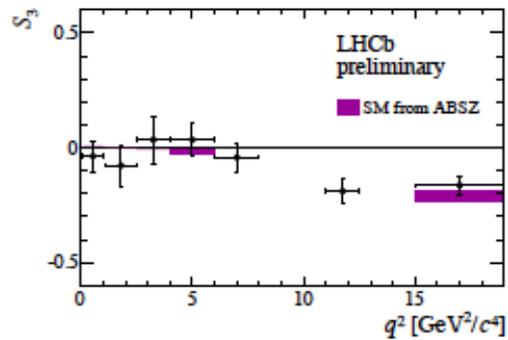
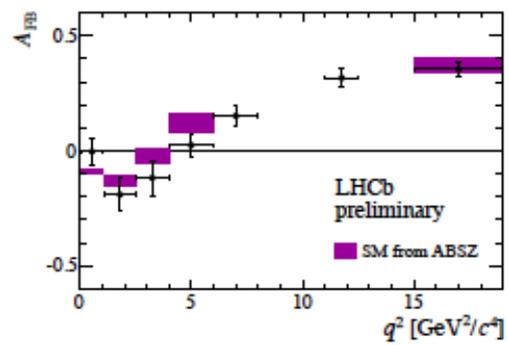
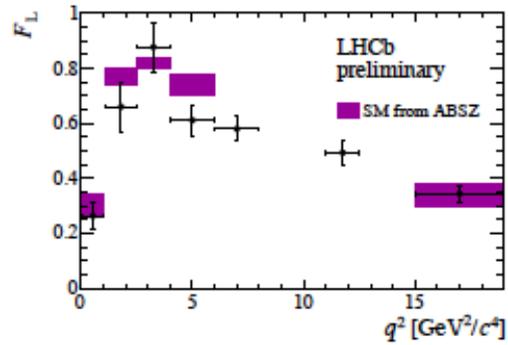
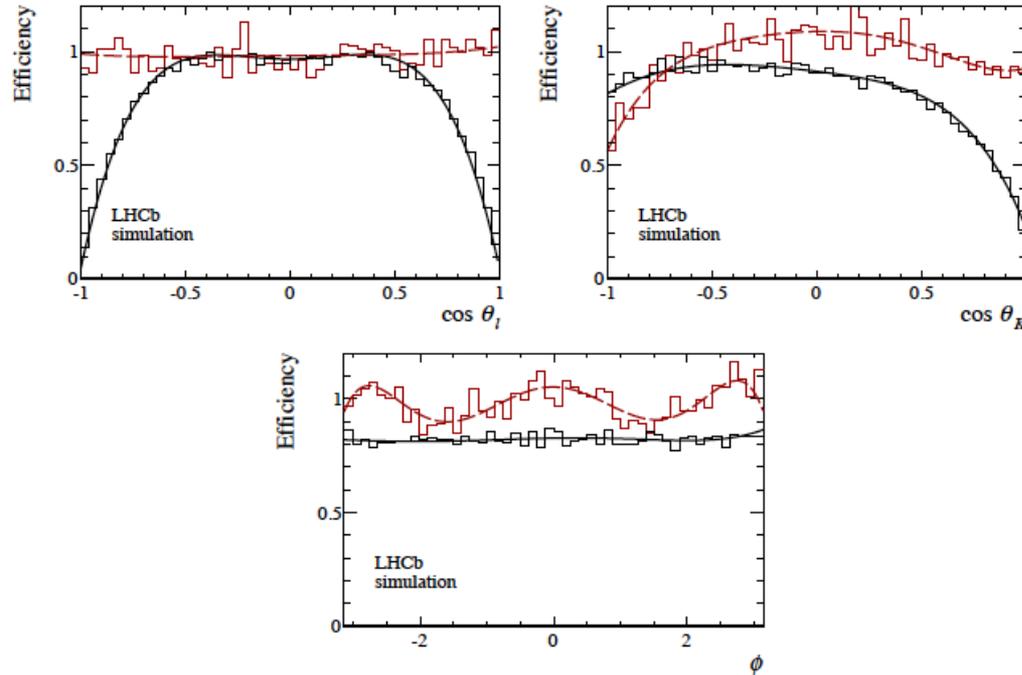


Figure 3: Invariant mass  $m(K^+\pi^-\mu^+\mu^-)$  for (left) the control decay  $B^0 \rightarrow J/\psi K^{*0}$  and (right) the signal decay  $B^0 \rightarrow K^{*0}\mu^+\mu^-$ , integrated over the full  $q^2$  range. The  $B^0 \rightarrow K^{*0}\mu^+\mu^-$  signal yield integrated over  $q^2$  is determined to be  $2398 \pm 57$ . Overlaid are the projections of the total fitted distribution (black line) and its different components. The signal is shown by the blue component and the background is shown by the red hatched component.



# Efficiency at low $q^2$ for Belle II ?



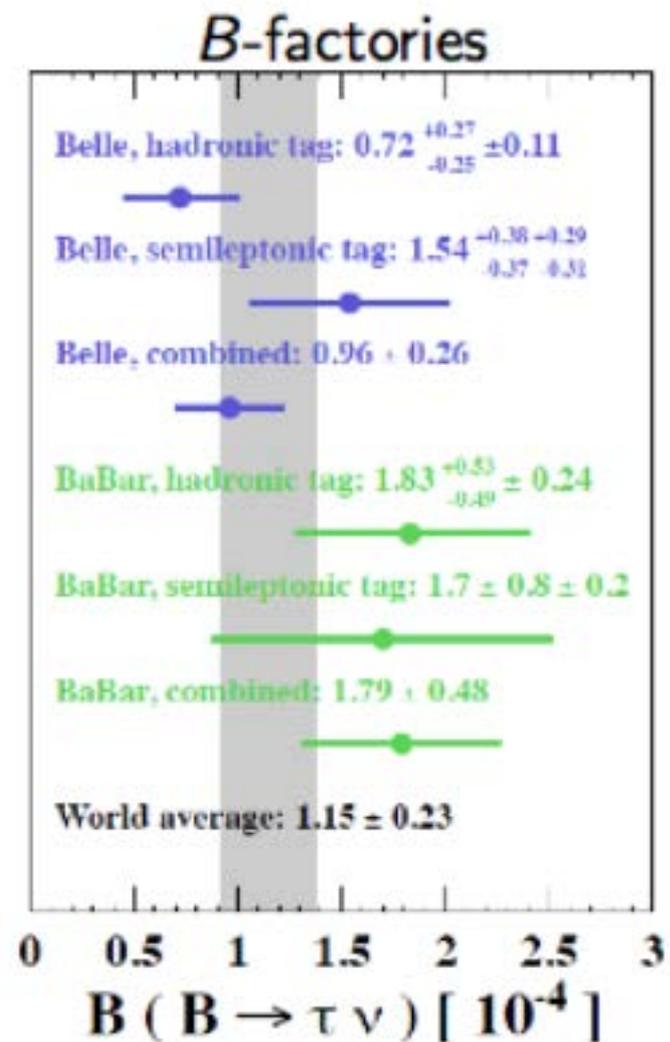
LHCb detection efficiency plots.

Figure 2: Angular efficiency in  $\cos \theta_\ell$ ,  $\cos \theta_K$  and  $\phi$ , as determined from a principal moment analysis of simulated three-body  $B^0 \rightarrow K^{*0} \mu^+ \mu^-$  phase-space decays. The efficiency is shown for the regions  $0.1 < q^2 < 0.98 \text{ GeV}^2/c^4$  (black solid line) and  $18.0 < q^2 < 19.0 \text{ GeV}^2/c^4$  (red dashed line). The histograms indicate the distribution of simulated three-body  $B^0 \rightarrow K^{*0} \mu^+ \mu^-$  phase-space decays used to determine the acceptance. The absolute normalisation of the distributions is arbitrary.

How important are the di-electron modes that Belle II does well ? (see the effect of the [photon pole](#) more clearly; NP from right handed currents ??)

## Current world average

$$\mathcal{B}(B^+ \rightarrow \tau^+ \nu_\tau) = (1.14 \pm 0.27) \times 10^{-4}$$



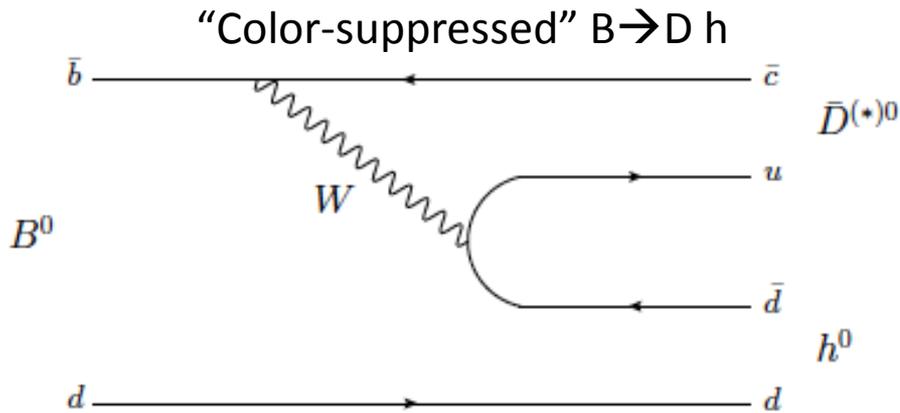
## New Belle result (semileptonic tags)

$$\mathcal{B}(B^+ \rightarrow \tau^+ \nu_\tau) = [1.25 \pm 0.28(\text{stat.}) \pm 0.27(\text{syst.})] \times 10^{-4}$$

# B factories: *Check CP violation in $b \rightarrow c$ [ $\bar{u} \bar{d}$ ] processes*

## 2015: First joint BaBar-Belle data analysis

M. Rohrken et al

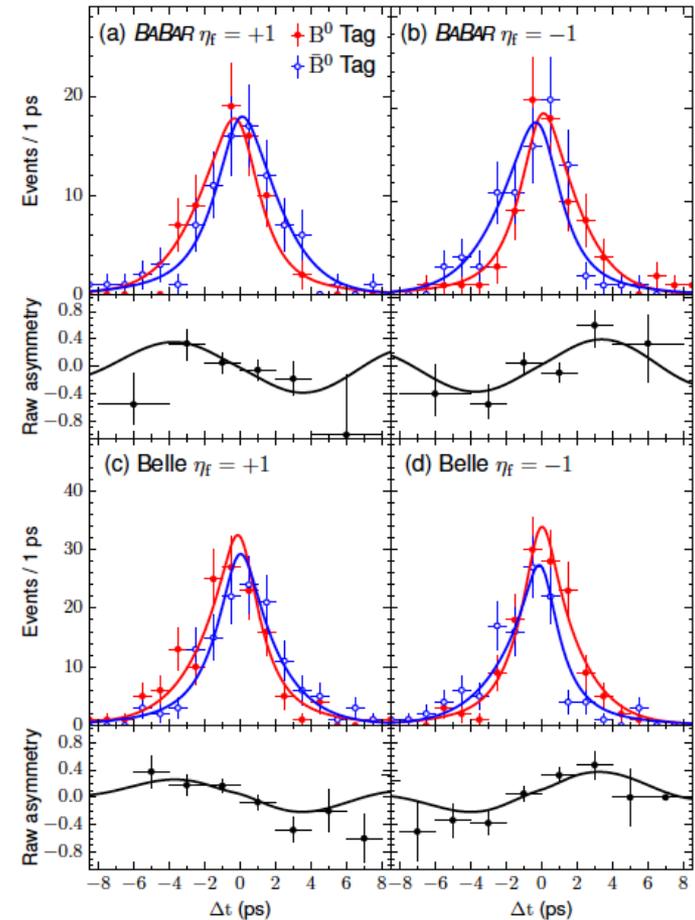


where  $D^0$  is a CP eigenstate and  $h^0 = \pi^0, \eta, \omega$

Combining Belle and BaBar datasets,  
 $\sim 1260$  signal events, obtain a  $5.4\sigma$  CP  
 violation signal  $\rightarrow$  First observation  
 $\sin(2\beta_{\text{eff}}) = 0.66 \pm 0.10(\text{stat}) \pm 0.06(\text{sys})$

Phase of  
 $V_{td}$  again

Conclusion: CP violation in  $b \rightarrow c$   $\bar{u} \bar{d}$  modes is  
 the same as in  $b \rightarrow c$   $\bar{c} \bar{s}$  modes (e.g.  $B \rightarrow J/\psi K_S$ )



More backup

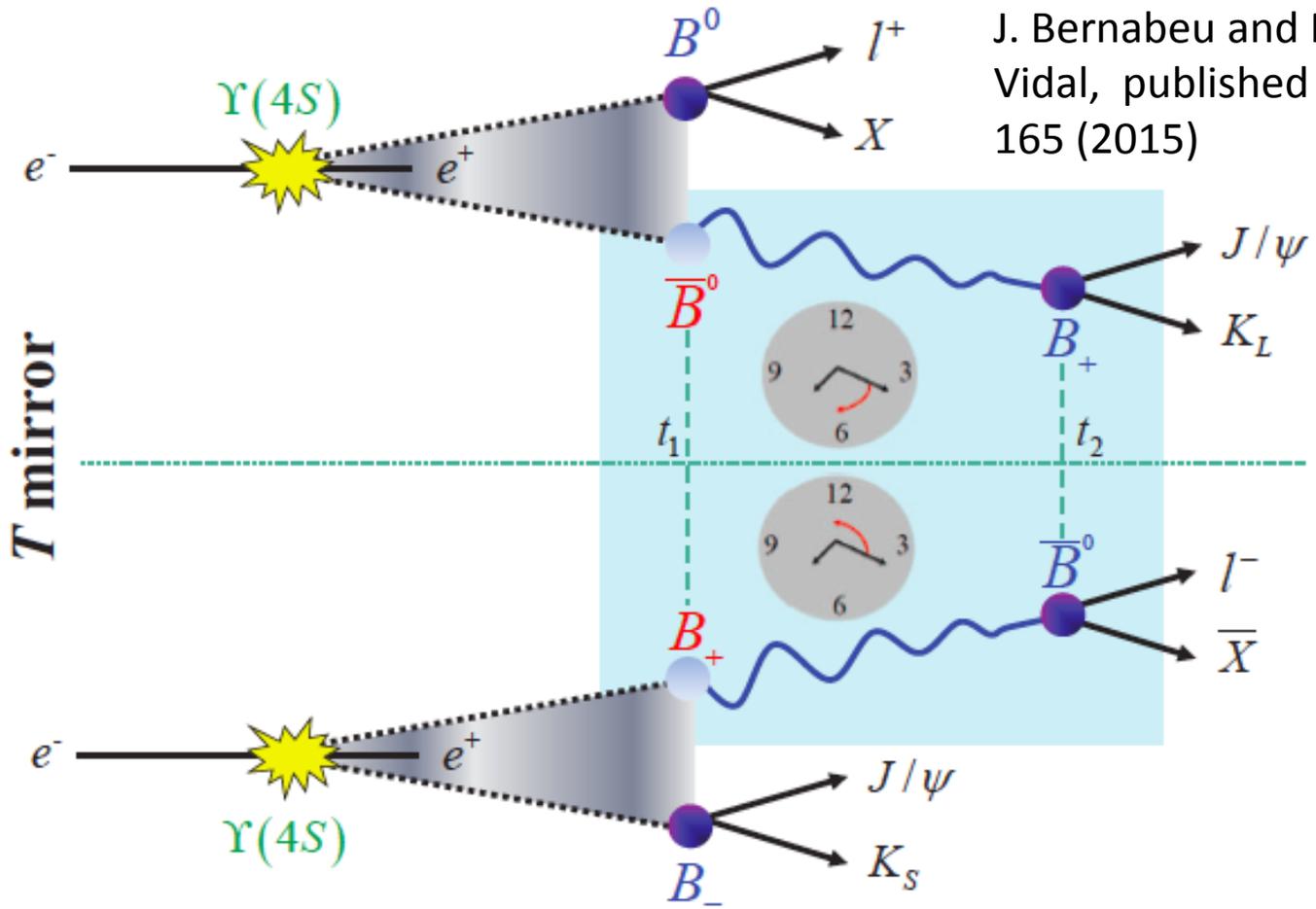
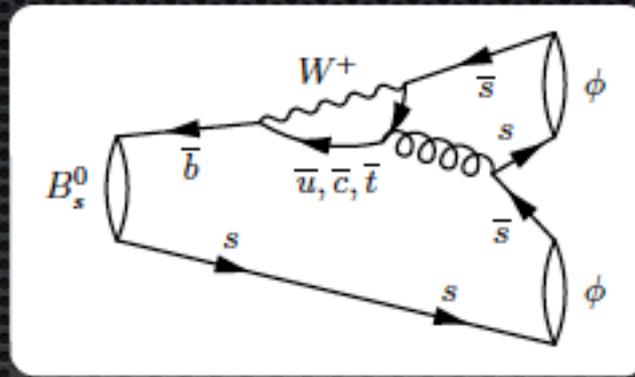
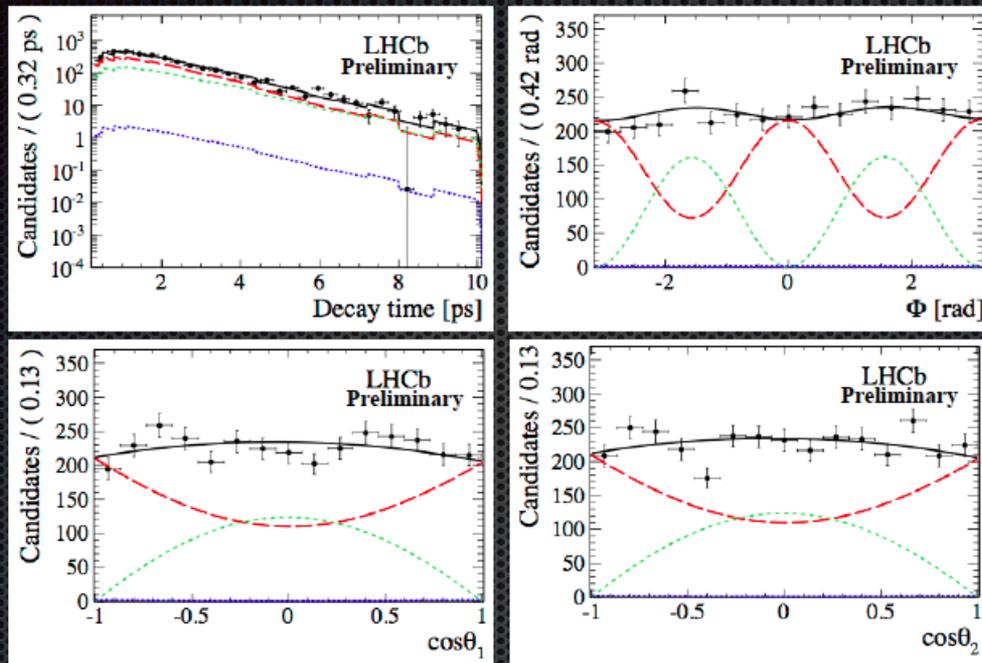


FIG. 11. Foundations of the time-reversal experiment. Electron-positron collisions at the asymmetric  $B$  factory produce  $\Upsilon(4S)$  resonances, each of which decays through strong interaction in an entangled pair of  $B$  mesons. When one  $B$  meson decays at  $t_1$ , the identity of the other is “tagged” without measuring it specifically. In the top panel, the  $B$  meson observed to decay to the final state  $l^+X$  at  $t_1$  transfers information to the (still living) partner meson and dictates that it is in a  $\bar{B}^0$  state. This surviving meson tagged as  $\bar{B}^0$  is observed later at  $t_2$ , encapsulating a time ordering, to decay into a final state  $J/\psi K_L^0$  that filters the  $B$  meson to be in a  $B_+$  state, a linear combination of  $B^0$  and  $\bar{B}^0$  states. This case corresponds to a transition  $\bar{B}^0 \rightarrow B_+$ . To study time reversal we have to compare the rate at which this transition occurs to the rate of the time-reversed transition,  $B_+ \rightarrow \bar{B}^0$  (bottom panel). Adapted from<sup>2</sup>.

But LHCb dominates on these  $B_s$  modes



## $B_s \rightarrow \phi\phi$ - Time-Dependent Results



Projections are s-weighted and include acceptances,  
 Decay time acceptance from  $B_s \rightarrow D_s \pi$  data,  
 Angular acceptance from simulated events.



# Signal model

We use the same notations as LHCb [arXiv:1304.2600]:

$$\frac{d^4\Gamma(B_s(t))}{d\Theta dt} = X(\Theta, \alpha, t) = \sum_{i=1}^{10} O_i(\alpha, t) \cdot g_i(\Theta),$$

$$O_i(\alpha, t) = N_i e^{-\Gamma_s t} \left[ a_i \cosh\left(\frac{1}{2}\Delta\Gamma_s t\right) + b_i \sinh\left(\frac{1}{2}\Delta\Gamma_s t\right) + c_i \cos(\Delta m_s t) + d_i \sin(\Delta m_s t) \right]$$

$i$	$g_i(\theta_T, \psi_T, \phi_T)$	$N_i$	$a_i$	$b_i$	$c_i$	$d_i$
1	$2 \cos^2 \psi_T (1 - \sin^2 \theta_T \cos^2 \phi_T)$	$ A_0(0) ^2$	1	$D$	$C$	$-S$
2	$\sin^2 \psi_T (1 - \sin^2 \theta_T \sin^2 \phi_T)$	$ A_{\parallel}(0) ^2$	1	$D$	$C$	$-S$
3	$\sin^2 \psi_T \sin^2 \theta_T$	$ A_{\perp}(0) ^2$	1	$-D$	$C$	$S$
4	$-\sin^2 \psi_T \sin 2\theta_T \sin \phi_T$	$ A_{\parallel}(0)A_{\perp}(0) $	$C \sin(\delta_{\perp} - \delta_{\parallel})$	$S \cos(\delta_{\perp} - \delta_{\parallel})$	$\sin(\delta_{\perp} - \delta_{\parallel})$	$D \cos(\delta_{\perp} - \delta_{\parallel})$
5	$\frac{1}{\sqrt{2}} \sin 2\psi_T \sin^2 \theta_T \sin 2\phi_T$	$ A_0(0)A_{\parallel}(0) $	$\cos(\delta_{\parallel} - \delta_0)$	$D \cos(\delta_{\parallel} - \delta_0)$	$C \cos(\delta_{\parallel} - \delta_0)$	$-S \cos(\delta_{\parallel} - \delta_0)$
6	$\frac{1}{\sqrt{2}} \sin 2\psi_T \sin 2\theta_T \sin \phi_T$	$ A_0(0)A_{\perp}(0) $	$C \sin(\delta_{\perp} - \delta_0)$	$S \cos(\delta_{\perp} - \delta_0)$	$\sin(\delta_{\perp} - \delta_0)$	$D \cos(\delta_{\perp} - \delta_0)$
7	$\frac{2}{3}(1 - \sin^2 \theta_T \cos^2 \phi_T)$	$ A_S(0) ^2$	1	$-D$	$C$	$S$
8	$\frac{1}{3}\sqrt{6} \sin \psi_T \sin^2 \theta_T \sin 2\phi_T$	$ A_S(0)A_{\parallel}(0) $	$C \cos(\delta_{\parallel} - \delta_S)$	$S \sin(\delta_{\parallel} - \delta_S)$	$\cos(\delta_{\parallel} - \delta_S)$	$D \sin(\delta_{\parallel} - \delta_S)$
9	$\frac{1}{3}\sqrt{6} \sin \psi_T \sin 2\theta_T \cos \phi_T$	$ A_S(0)A_{\perp}(0) $	$\sin(\delta_{\perp} - \delta_S)$	$-D \sin(\delta_{\perp} - \delta_S)$	$C \sin(\delta_{\perp} - \delta_S)$	$S \sin(\delta_{\perp} - \delta_S)$
10	$\frac{4}{3}\sqrt{3} \cos \psi_T (1 - \sin^2 \theta_T \cos^2 \phi_T)$	$ A_S(0)A_0(0) $	$C \cos(\delta_0 - \delta_S)$	$S \sin(\delta_0 - \delta_S)$	$\cos(\delta_0 - \delta_S)$	$D \sin(\delta_0 - \delta_S)$

$$C = \frac{1 - |\lambda|^2}{1 + |\lambda|^2},$$

$$S = -\frac{2|\lambda| \sin \phi_s}{1 + |\lambda|^2},$$

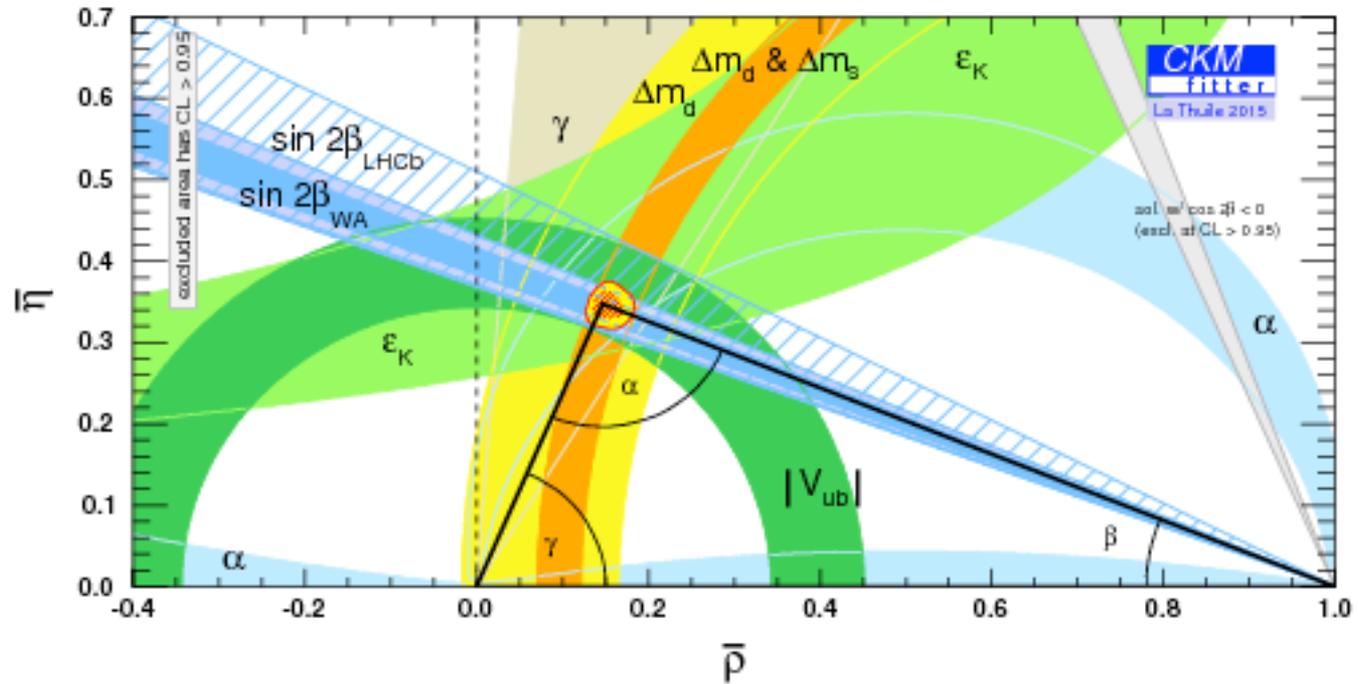
$$D = -\frac{2|\lambda| \cos \phi_s}{1 + |\lambda|^2}$$

$|\lambda|$  includes possible contribution from CP violation in direct decay, we assume  $|\lambda| = 1$  and we assign a systematic.

$\Delta\Gamma_S > 0$ : we use previous LHCb results.  $\alpha$  physics parameters ( $\Delta\Gamma_S, \phi_s, c\tau, |A_0|^2, |A_S|^2, |A_{\perp}|^2, \delta_{\parallel}, \delta_{S\perp}, \delta_{\perp}$ )



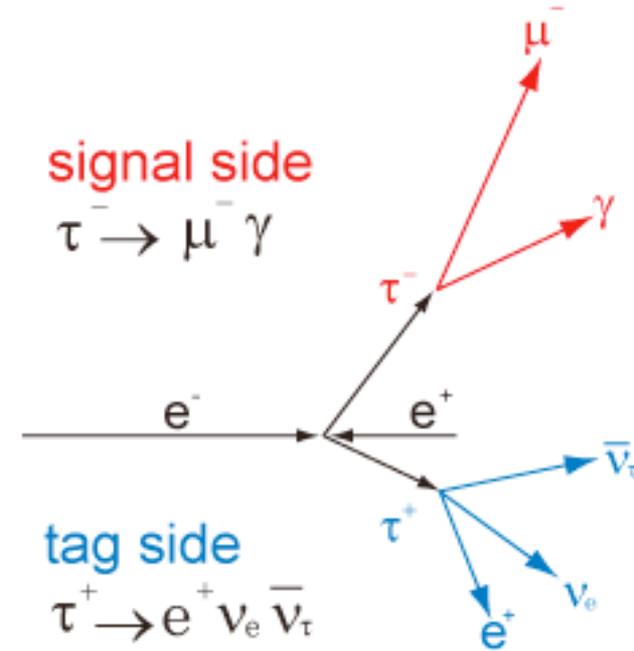
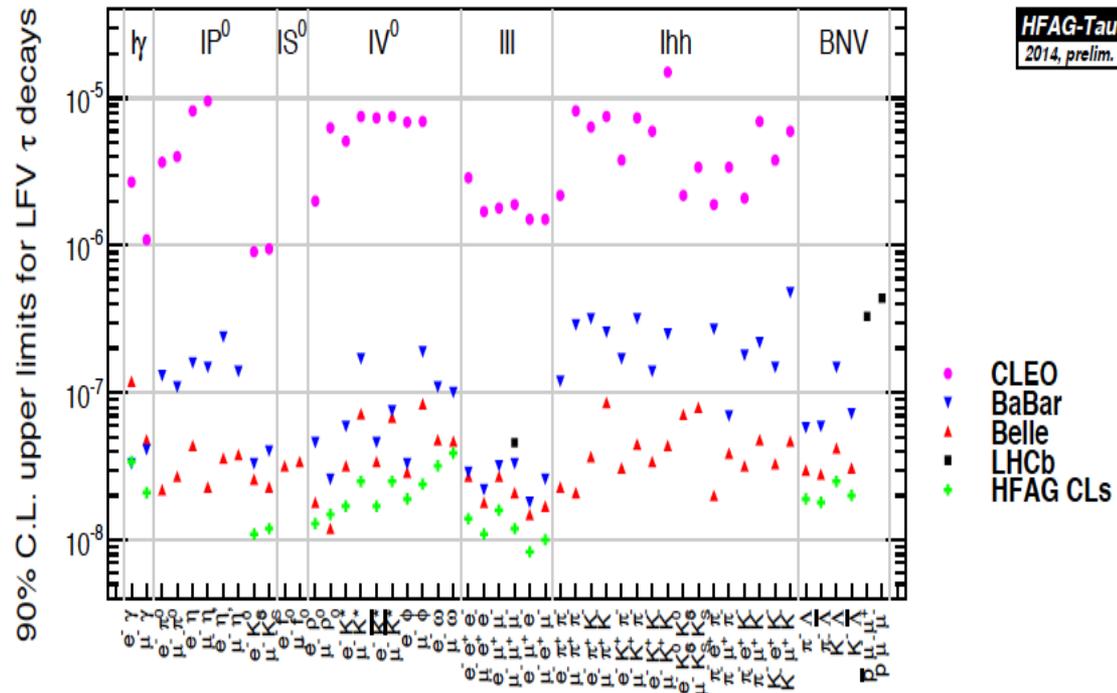
# CKMFitter with LHCb $\sin(2\beta)$ included



# Tau Lepton Flavor Violation

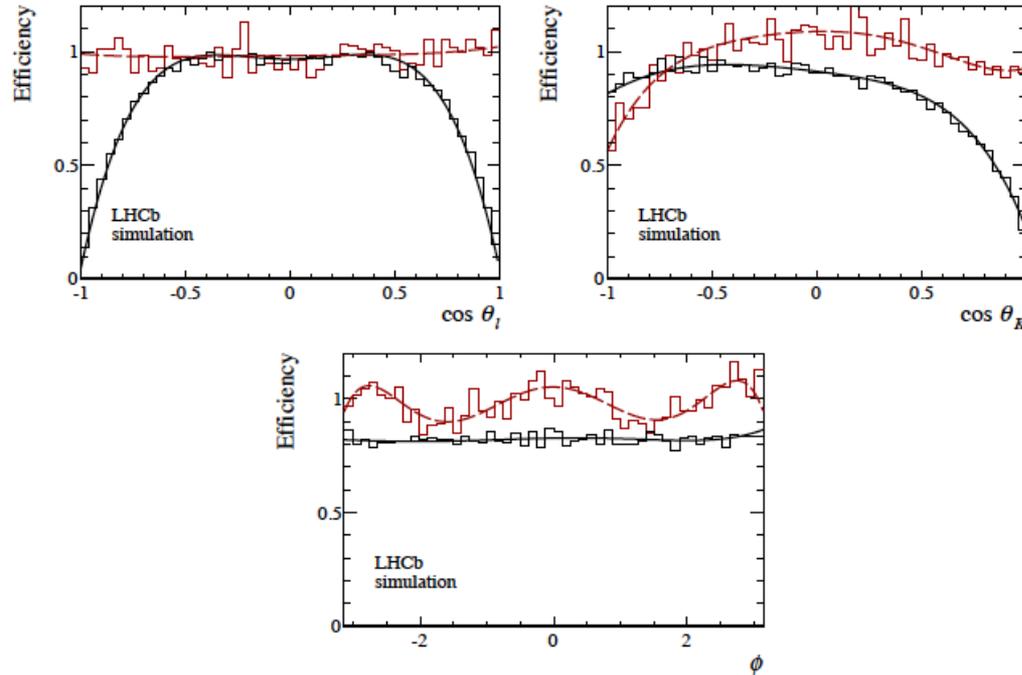
- Several processes:  $\tau \rightarrow l\gamma$ ,  $\tau \rightarrow l_\alpha \bar{l}_\beta l_\beta$ ,  $\tau \rightarrow lY$   
 $\leftarrow P, S, V, P\bar{P}, \dots$

Example of the decay topology



Belle II will push many limits below  $10^{-9}$  ; LHCb has very limited capabilities.

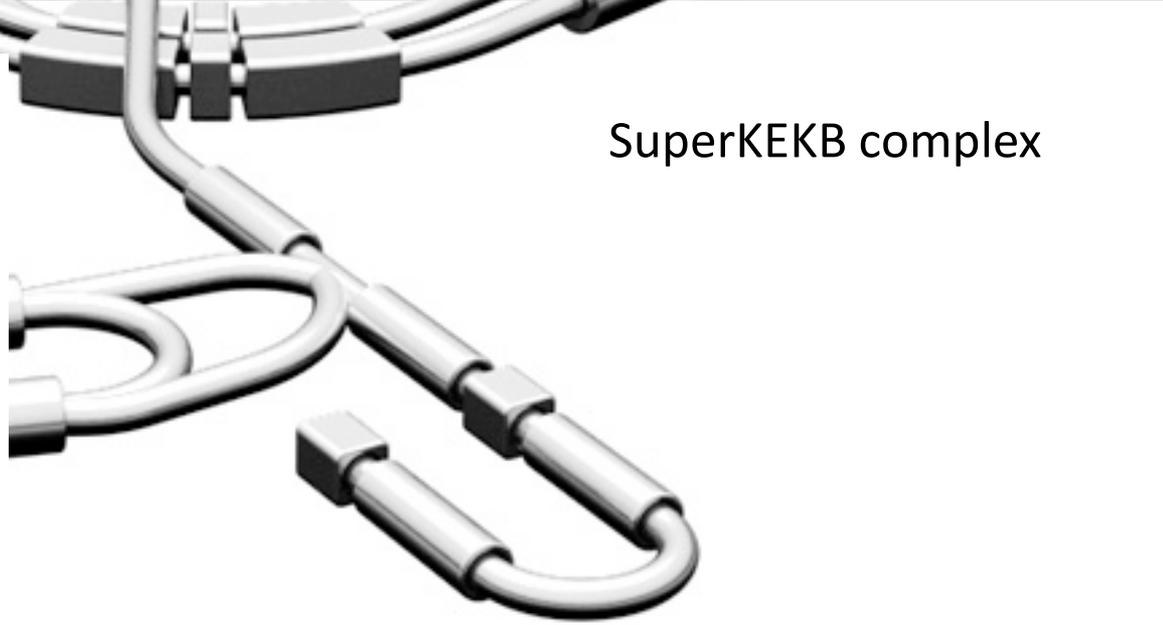
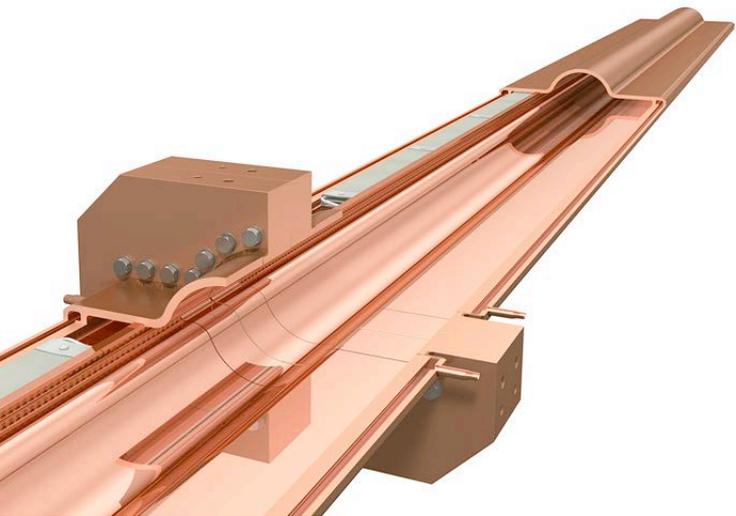
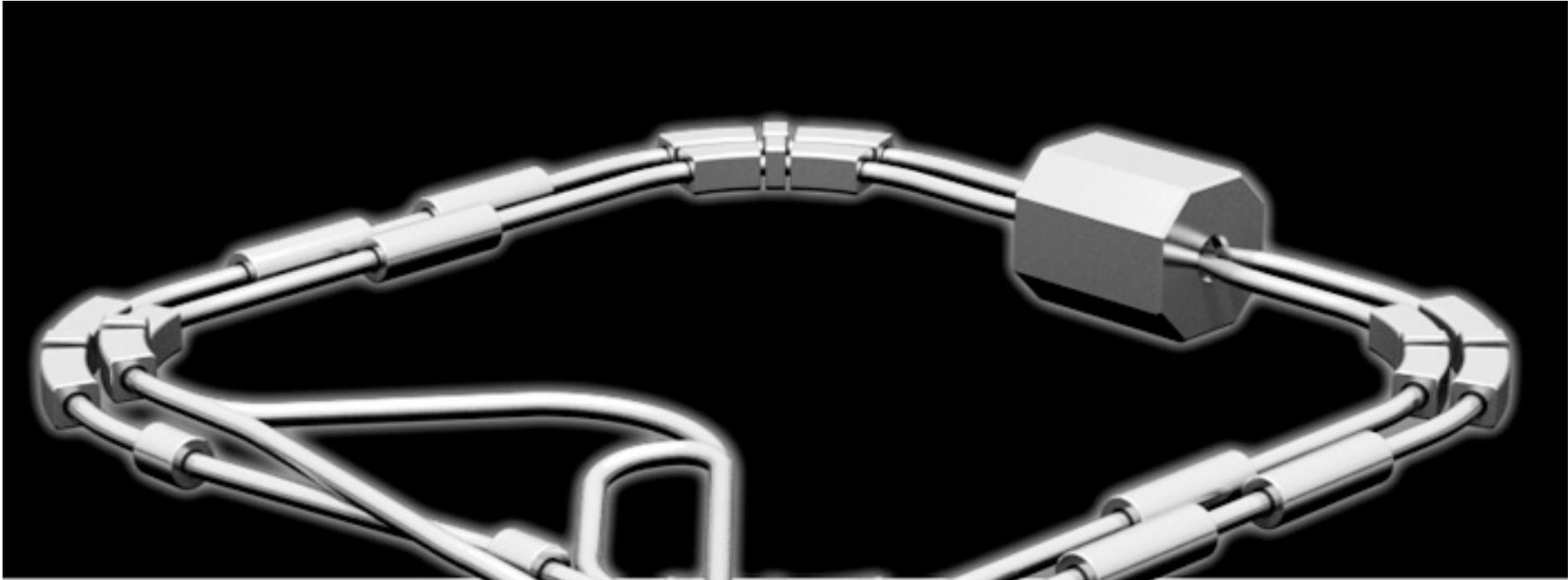
# Efficiency at low $q^2$ for Belle II ?



LHCb detection efficiency plots.

Figure 2: Angular efficiency in  $\cos \theta_\ell$ ,  $\cos \theta_K$  and  $\phi$ , as determined from a principal moment analysis of simulated three-body  $B^0 \rightarrow K^{*0} \mu^+ \mu^-$  phase-space decays. The efficiency is shown for the regions  $0.1 < q^2 < 0.98 \text{ GeV}^2/c^4$  (black solid line) and  $18.0 < q^2 < 19.0 \text{ GeV}^2/c^4$  (red dashed line). The histograms indicate the distribution of simulated three-body  $B^0 \rightarrow K^{*0} \mu^+ \mu^-$  phase-space decays used to determine the acceptance. The absolute normalisation of the distributions is arbitrary.

How important are the di-electron modes that Belle II does well ? (see the effect of the [photon pole](#) more clearly; NP from right handed currents ??)



SuperKEKB complex

# ATF2 nanobeams

