The one-loop matter bispectrum as a probe of gravity and dark energy

Ben Bose

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12th September 2018

in collaboration with Atsushi Taruya [1808.01120]



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Talk Outline:

Motivation.

- 2 Description of generalised perturbative approach to Large Scale Structure (LSS).
- 3 The Importance of theoretical **consistency** for upcoming surveys.
- 4 The 1-loop matter bispectrum in non-standard models of gravity and dark energy.

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5 Applications and Extensions.

Why Modify ACDM?

- Extrapolation of general relativity to cosmological scales.
- Dominant dark sector in standard picture.
- Tentative tensions in data sets.



[1409.2769]



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Generality is useful

We want a general picture of gravity

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Generality is useful ...

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... but must respect solar system tests \rightarrow screening.

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Comparisons of Approaches to LSS Modelling



[1607.03150,1704.05309,1804.05867]

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Core Ingredients:

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 Evolution of initial Gaussian perturbations to LSS in a flat FRLW background.

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- **2** Working in the Newtonian Regime, i.e. $\mathbf{v}, \Phi \ll 1$.
- **3** Loop corrections added to improve non-linear regime modelling.

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- Evolution of initial Gaussian perturbations to LSS in a flat FRLW background.
- **2** Working in the Newtonian Regime, i.e. $\mathbf{v}, \Phi \ll 1$.
- **3** Loop corrections added to improve non-linear regime modelling.
- 4 Additional modelling required to connect with observables.



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Eulerian Perturbation Theory

Conservation of energy and momentum lead to the **Continuity** and **Euler** equations:

$$a\frac{\partial\delta(\mathbf{k},a)}{\partial a} + \theta(\mathbf{k},a) = -\int \frac{d^3\mathbf{k}_1d^3\mathbf{k}_2}{(2\pi)^3} \delta_{\mathrm{D}}(\mathbf{k}-\mathbf{k}_1-\mathbf{k}_2)\alpha(\mathbf{k}_1,\mathbf{k}_2)\,\theta(\mathbf{k}_1,a)\delta(\mathbf{k}_2,a),$$

$$\begin{aligned} a\frac{\partial\theta(\mathbf{k},a)}{\partial a} + \left(2 + \left[\stackrel{???}{H}\right] + \frac{aH'}{H^2}\right)\theta(\mathbf{k},a) - \left(\frac{k}{aH}\right)^2 \Phi(\mathbf{k},a) = \\ &- \frac{1}{2}\int \frac{d^3\mathbf{k}_1 d^3\mathbf{k}_2}{(2\pi)^3}\delta_{\mathrm{D}}(\mathbf{k} - \mathbf{k}_1 - \mathbf{k}_2)\beta(\mathbf{k}_1,\mathbf{k}_2)\theta(\mathbf{k}_1,a)\theta(\mathbf{k}_2,a), \end{aligned}$$

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 $\cdots + \delta_q, \theta_q$ equations, etc.

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 $\cdots + \delta_q, \theta_q$ equations, etc.

* $\delta_{NL}(\mathbf{k}, \mathbf{a}) = \sum_{n=1}^{\infty} \delta_n(\mathbf{k}, \mathbf{a}), \qquad \theta_{NL}(\mathbf{k}, \mathbf{a}) = \sum_{n=1}^{\infty} \theta_n(\mathbf{k}, \mathbf{a}).$

* $\delta_n(\mathbf{k}; \mathbf{a}) = \frac{1}{(2\pi)^{3(n-1)}} \int d^3 \mathbf{k}_1 ... d^3 \mathbf{k}_n \delta_D(\mathbf{k} - \mathbf{k}_{1...n}) F_n(\mathbf{k}_1, ..., \mathbf{k}_n; \mathbf{a}) \delta_0(\mathbf{k}_1) ... \delta_0(\mathbf{k}_n).$

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Evolution equations solved **numerically** using MGCopter up to 4th order in the perturbations.

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Gravitational Modelling

We can keep gravity general through Poisson equation in the evolution equations

$$-\left(rac{k}{aH}
ight)^2\Phi=rac{3\Omega_m(a)}{2}\mu(k;a)\,\delta(\mathbf{k})+S(\mathbf{k};a),$$

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$$-\left(\frac{k}{aH}\right)^2 \Phi = \frac{3\Omega_m(a)}{2}\mu(k;a)\,\delta(\mathbf{k}) + S(\mathbf{k};a),$$

where the non-linear interaction term is given by

$$\begin{split} \mathcal{S}(\mathbf{k}, \mathbf{a}) &= \int \frac{d^3 \mathbf{k}_1 d^3 \mathbf{k}_2}{(2\pi)^3} \, \delta_{\mathrm{D}}(\mathbf{k} - \mathbf{k}_{12}) \gamma_2(\mathbf{k}_1, \mathbf{k}_2; \mathbf{a}) \delta(\mathbf{k}_1) \, \delta(\mathbf{k}_2) \\ &+ \int \frac{d^3 \mathbf{k}_1 d^3 \mathbf{k}_2 d^3 \mathbf{k}_3}{(2\pi)^6} \delta_{\mathrm{D}}(\mathbf{k} - \mathbf{k}_{123}) \gamma_3(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3; \mathbf{a}) \delta(\mathbf{k}_1) \, \delta(\mathbf{k}_2) \, \delta(\mathbf{k}_3) \\ &+ \int \frac{d^3 \mathbf{k}_1 d^3 \mathbf{k}_2 d^3 \mathbf{k}_3 d^3 \mathbf{k}_4}{(2\pi)^9} \delta_{\mathrm{D}}(\mathbf{k} - \mathbf{k}_{1234}) \gamma_4(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3, \mathbf{k}_4; \mathbf{a}) \delta(\mathbf{k}_1) \, \delta(\mathbf{k}_2) \, \delta(\mathbf{k}_3) \delta(\mathbf{k}_4). \end{split}$$

[1606.02520]

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This term is responsible for screening.

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Dark Sector Modelling

We can incorporate **general interactions** within the dark sector *easily*. Some examples:

Momentum exchange models:

$$\left(2+A(a)+\frac{aH'}{H}\right)\theta(k;a),$$

[1412.1080]

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Image: A math a math

Dark Sector Modelling

We can incorporate **general interactions** within the dark sector *easily*. Some examples:

Momentum exchange models:

$$\left(2+A(a)+\frac{aH'}{H}\right)\theta(k;a),$$

[1412.1080]

or NOT so easily ...

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Clustering quintessence systems :

$$\begin{aligned} a\theta_{q}' + \left[2 - 3w + \frac{aH'}{H}\right]\theta_{q} &- \frac{k^{2}}{a^{2}H^{2}}\Phi - \frac{c_{s}^{2}k^{2}\delta_{q}}{(1+w)a^{2}H^{2}} = \\ &\int d^{3}k_{1}d^{3}k_{2}\delta_{D}(k-k_{1}-k_{2})\left[-\frac{\beta(k_{1},k_{2})}{2}\theta_{q}(k_{1})\theta_{q}(k_{2}) + \frac{3(c_{s}^{2}-w)}{(1+w)}\delta_{q}(k_{1})\theta_{q}(k_{2})\alpha(k_{1})\right] \\ &- \frac{ac_{s}^{2}}{(1+w)}\delta_{q}'(k_{1})\theta_{q}(k_{2})\alpha(k_{1},k_{2}) + \frac{c_{s}^{2}(1+c_{s}^{2})}{a^{2}H^{2}(1+w)^{2}}\delta_{q}(k_{1})\delta_{q}(k_{2})(k_{1}k_{2}\mu - k_{2}^{2})\right]. \end{aligned}$$

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The **power spectrum** or **correlation function** are Fourier doubles of a 2-point correlation of the perturbations

$$\langle \delta(\mathbf{k})\delta(\mathbf{k}')\rangle = (2\pi)^3 \delta_{\mathrm{D}}(\mathbf{k}+\mathbf{k}') P(k).$$

- Has been well studied in perturbative regime. [ex. 1607.03150,1607.03148]
- High precision measurements demand accurate theory. There are many difficulties in this respect.
- Information rich scales difficult to model in a **useful** way.

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Example: Relevance of Consistent Modelling for Future Spectroscopic Surveys (z = 1)



* Screening effects on matter power spectrum in above analysis $\leq 0.5\%$.

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Information Extraction: 3-point Statistics

Galaxy Survey

$\times 10^7$ B_{000} 🖡 data 0.8 S3-wide ($f_{sky} = 0.5, \sigma_P = 6\mu K', \theta = 1'$) B_{110} MD-Patchy mocks B_{220} shot noise S3-deep ($f_{sky} = 0.05$, $\sigma_p = 3\mu K'$, $\theta = 1'$) $k^2\,B_{\ell_1\ell_2L=0}(k,k)$ 0.6 50 0.4Solid: iterative ∛ 40 - Dashed: quadratic $(S/N)_{<}$ 0.20.0 20 -0.210 "*B*/*B*" 500 1000 1500 2000 0.050.10 0.150.20 maximum multipole $k \left[h \, \mathrm{Mpc}^{-1} \right]$

[1803.02132,1604.08578]

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CMB Lensing

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The Matter Bispectrum

The **bispectrum** is a Fourier space 3-point correlation measurement of the perturbations:

 $\langle \delta(\mathbf{k}_1)\delta(\mathbf{k}_2)\delta(\mathbf{k}_3) \rangle = (2\pi)^3 \delta_{\mathrm{D}}(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3) B(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3).$

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To include non-linear information we can include loop corrections. The 1-loop bispectrum is then given by

$$B_{1-loop}(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3) = B_{112} + [B_{222} + B_{321} + B_{411}],$$

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$$B_{1-loop}(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3) = B_{112} + [B_{222} + B_{321} + B_{411}],$$

where

$$\langle \delta_2(\mathbf{k}_1) \delta_2(\mathbf{k}_2) \delta_2(\mathbf{k}_3) \rangle = (2\pi)^3 \delta_{\mathrm{D}}(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3) B_{222},$$

$$\langle \delta_1(\mathbf{k}_1) \delta_2(\mathbf{k}_2) \delta_3(\mathbf{k}_3) \rangle \sim (2\pi)^3 \delta_{\mathrm{D}}(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3) B_{321},$$

 $\langle \delta_1(\mathbf{k}_1) \delta_1(\mathbf{k}_2) \delta_4(\mathbf{k}_4)
angle \sim (2\pi)^3 \delta_{\mathrm{D}}(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3) B_{411}$, we have $\delta_1(\mathbf{k}_2) \delta_4(\mathbf{k}_4) > 0$

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1-loop Bispectrum: DGP



Fitting formula for (Beyond) Horndeski: [Namikawa et al. 1805.10567]

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1-loop Bispectrum: f(R)



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1-loop Bispectrum: Momentum Exchange



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Power of 3-point Statistics: **DGP** , $\Omega_{rc} = 0.438$.



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Power of 3-point Statistics: **DGP** , $\Omega_{rc} = 0.438$.



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Power of 3-point Statistics: f(R), $|f_{R0}| = 2.5 \times 10^{-6}$



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Applications and Extensions: Galaxy Bispectrum



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 Redshift space galaxy bispectrum can complement power spectrum analyses of galaxy data sets: improve constraints (factor of up to 5) and break degeneracies.

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[1705.04392, 1606.00439]

- * Redshift space galaxy bispectrum can complement power spectrum analyses of galaxy data sets: improve constraints (factor of up to 5) and break degeneracies. [1705.04392,1606.00439]
- * Velocity kernels already computed by algorithm but additional integral required for multipoles.

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- * Total model parameters for a redshift space, biased tracer bispectrum prediction at 1-loop order is $\geq 10!$ [1705.02574,1806.04015]

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- * Redshift space galaxy bispectrum can complement power spectrum analyses of galaxy data sets: improve constraints (factor of up to 5) and break degeneracies. [1705.04392,1606.00439]
- * Velocity kernels already computed by algorithm but additional integral required for multipoles.
- * Requires galaxy bias and redshift space modelling so constraints may be weaker and/or biased....
- * Total model parameters for a redshift space, biased tracer bispectrum prediction at 1-loop order is > 10![1705.02574,1806.04015]
- * Convergence spectra has great scope to constrain gravitational theories but non-linear modelling must be accurate + intrinsic alignments, z-distribution. (ロ) (四) (三) (三)

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[1801.01741,1805.10567]

Applications and Extensions: CMB Lensing Bispectrum

Direct application to CMB experiments:

- * CMB lensing very clean but currently offers weaker constraints than convergence spectra.
- Next generation CMB experiments may be able to offer competitive constraints on gravity and dark energy. 1805.10567





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 SPT at 1-loop order plus some redshift space model (ex. TNS) seems to be a promising framework for upcoming survey analyses (see Fonseca et al. 1805.12394 for example).

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- Code has been validated against analytic results and simulation results and is ready for application to survey data.
- Redshift space extension is available but introduces many free parameters and numerically challenging in general case.
- CMB lensing offers clean measurement but may not have strong enough signal even at stage 4 plus added complication of post-born modelling.

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Thanks for listening!



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