

Predicting Proton Decay from low energy SUSY spectrum

Kazuki Sakurai

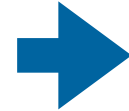
(University of Warsaw)

in collaboration with

Stefan Pokorski, Krzysztof Rolbiecki, Graham Ross

◎ Motivation of SUSY

- Fine-tuning Problem



tension with LHC;
but better than non-SUSY

- Dark Matter



well studied; consistent if its pure Higgsino ($\sim 1\text{TeV}$) or pure Wino ($\sim 3\text{TeV}$)

- Gauge Coupling Unification



not well studied compared to the other two

▶ How is the condition of GCU formulated?

▶ Is there an upper/lower limit on SUSY masses from GCU?

▶ Any relation between low energy SUSY and proton decay?

Contents

- Derivation of analytical formulae for GCU including the bulk 2-loop correction applicable for general SUSY and GUT spectra.
- Application:
 - Minimal SU(5)
coloured Higgs mass as a function of general SUSY mass spectrum
=> prediction of D=5 proton decay from SUSY spectrum
 - Orbifold SUSY SU(5) [Hall-Nomura model]
X,Y gauge boson mass as a function of general SUSY mass spectrum
=> prediction of D=6 proton decay from SUSY spectrum

lightest
sparticle

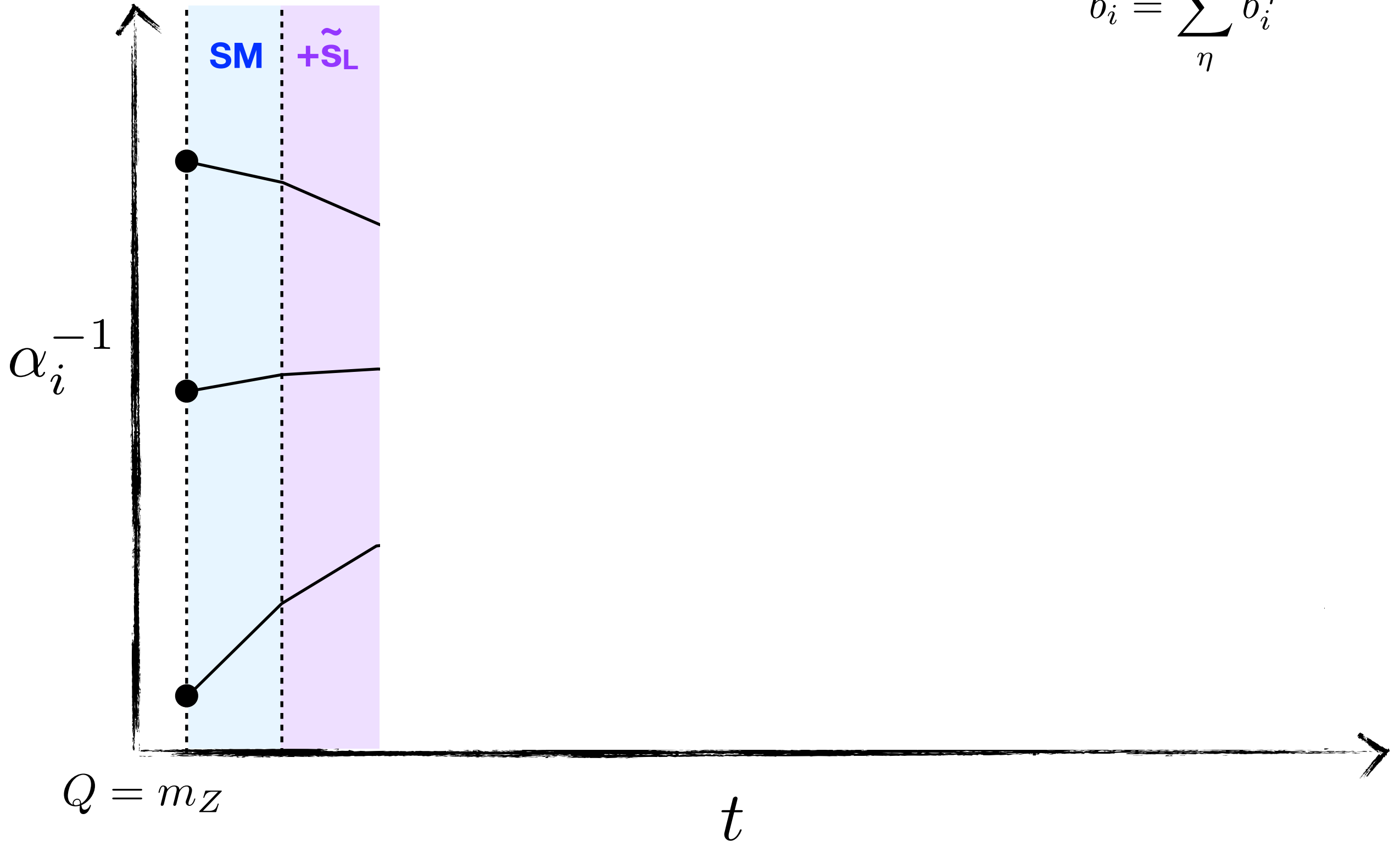
$\tilde{S}_L$

$$\frac{d}{dt} \tilde{\alpha}_i^{-1} = b_i$$

$$\tilde{\alpha}_i^{-1} \equiv 2\pi\alpha_i^{-1}$$

$$t \equiv \ln(Q/Q_0) > 0$$

$$b_i = \sum_{\eta} b_i^{\eta}$$



lightest
sparticle

heaviest
sparticle

$\tilde{S}_L$

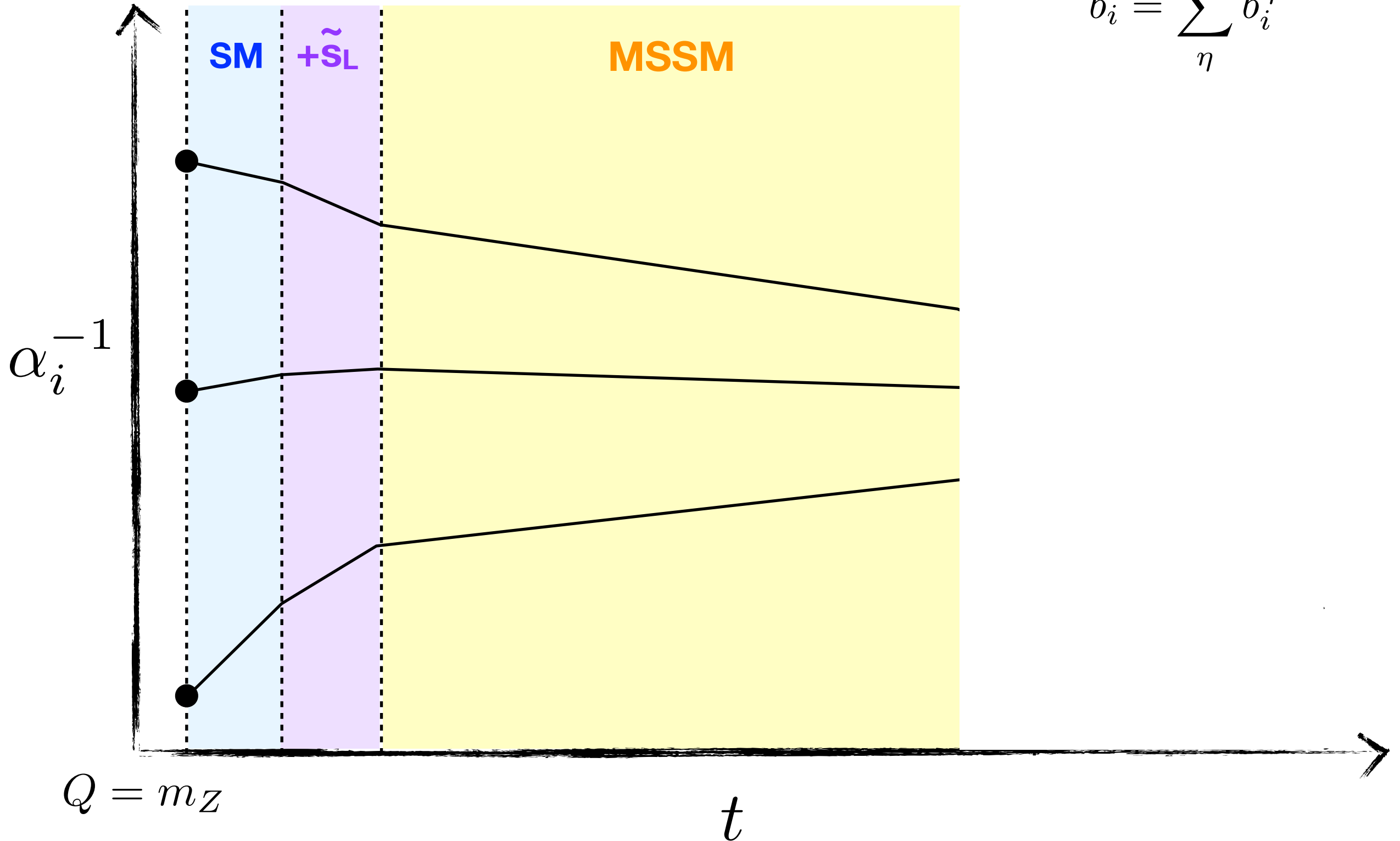
$\tilde{S}_H$

$$\frac{d}{dt} \tilde{\alpha}_i^{-1} = b_i$$

$$\tilde{\alpha}_i^{-1} \equiv 2\pi\alpha_i^{-1}$$

$$t \equiv \ln(Q/Q_0) > 0$$

$$b_i = \sum_{\eta} b_i^{\eta}$$



lightest
sparticle

heaviest
sparticle

lightest
GUT
particle

heaviest
GUT
particle

$$\frac{d}{dt} \tilde{\alpha}_i^{-1} = b_i$$

$\tilde{S}_L$

$\tilde{S}_H$

ξ_1

ξ_2

ξ_H

SM

$+\tilde{S}_L$

MSSM

$+\xi_1$

$+\xi_2$

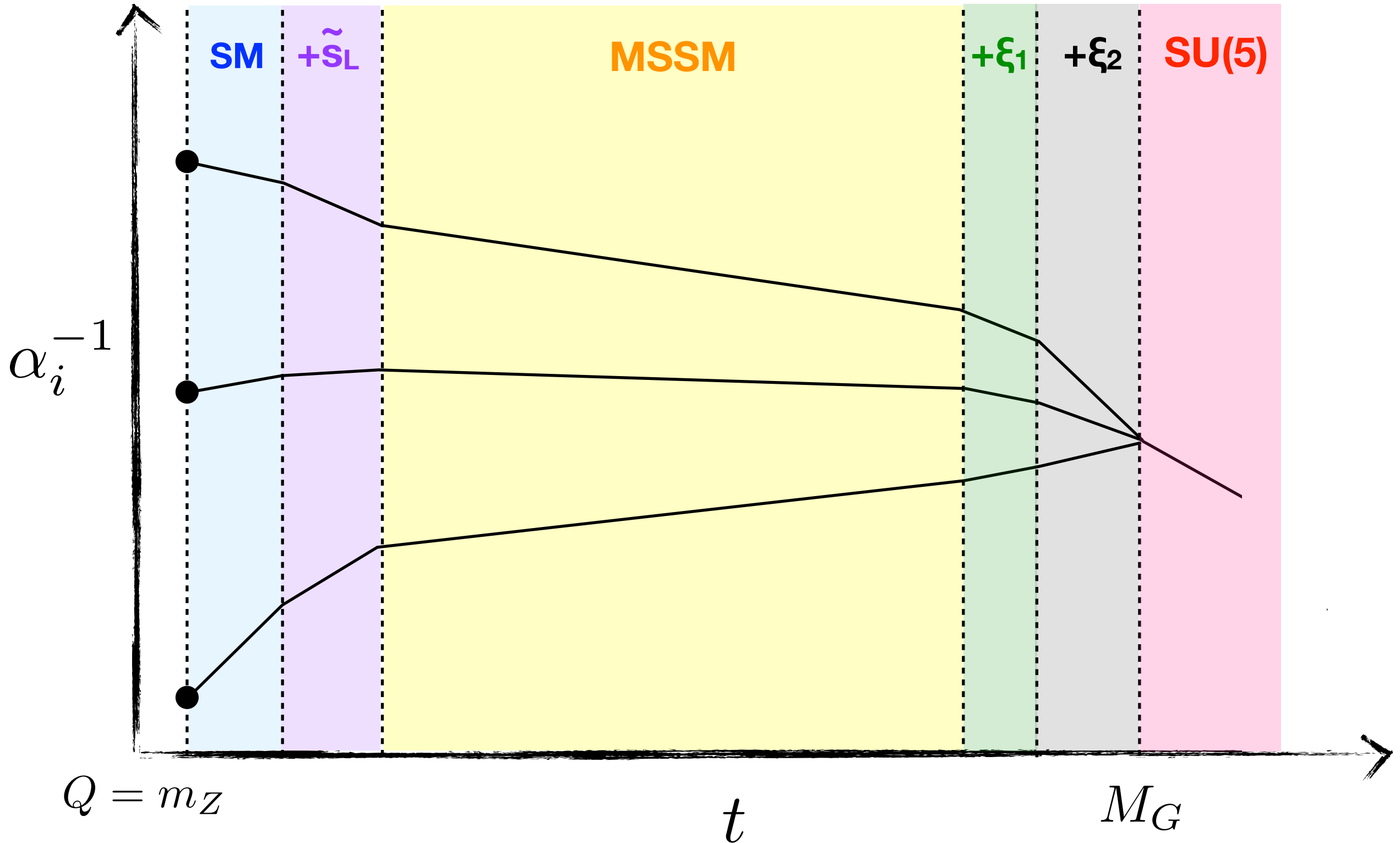
SU(5)

α_i^{-1}

$Q = m_Z$

t

M_G



lightest
sparticle

heaviest
sparticle

$\tilde{S}_L$

$\tilde{S}_H$

$$\frac{d}{dt} \tilde{\alpha}_i^{-1} = b_i$$

lightest
GUT
particle

ξ_1

ξ_2

heaviest
GUT
particle

ξ_H

SM

$+\tilde{S}_L$

MSSM

$+\xi_1$

$+\xi_2$

SU(5)

α_i^{-1}

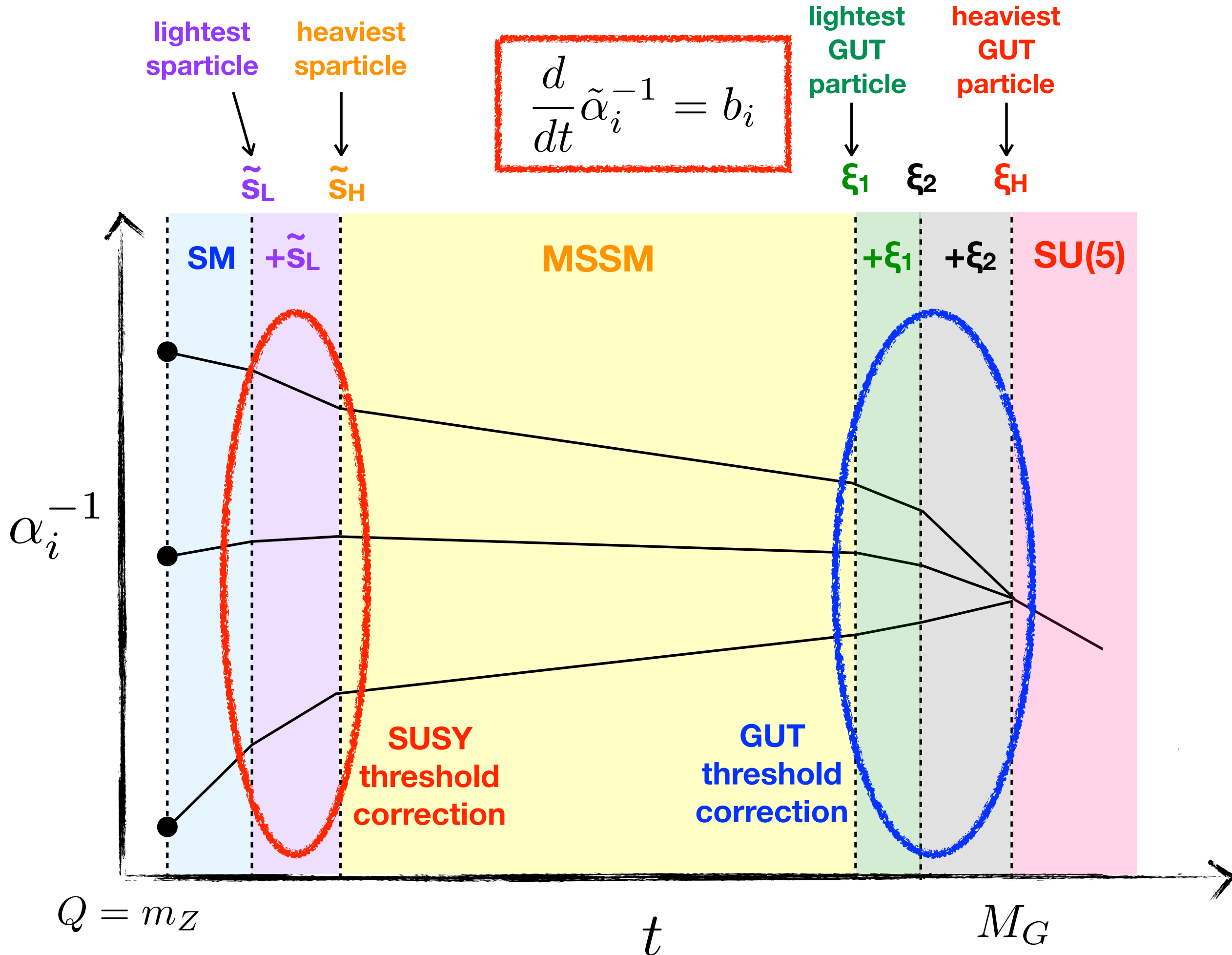
SUSY
threshold
correction

GUT
threshold
correction

$Q = m_Z$

t

M_G



General solution to RGE

$$\frac{2\pi}{\alpha(\Lambda)}$$

↑
unified coupling

$$= \frac{2\pi}{\alpha_i(m_Z)}$$

↑
experimental input

$$- b_i \ln \left(\frac{\Lambda}{m_Z} \right)$$

1-loop MSSM ↓

$$+ s_i$$

SUSY threshold ↓

$$+ r_i$$

GUT threshold ↑

$$+ \gamma_i$$

bulk 2-loop contribution ↓

$$+ \Delta_i$$

top-quark threshold, MSbar-DRbar conversion ↑

$$s_i = \sum_{\eta} b_i^{\eta} \ln \left(\frac{m_{\eta}}{m_Z} \right)$$

$$r_i = \sum_{\xi} b_i^{\xi} \ln \left(\frac{m_{\xi}}{\Lambda} \right)$$

[bulk 2-loop contribution]

$$b_i = \left(\frac{33}{5}, 1, -3 \right)$$

$$\gamma_i = -\frac{1}{2} \sum_j \frac{b_{ij}}{b_j} \ln \left(\frac{\alpha_j(\Lambda)}{\alpha_j(m_Z)} \right)$$

$$\simeq -\frac{1}{2} \sum_j \frac{b_{ij}}{b_j} \ln \left(1 + \frac{b_j \alpha(\Lambda)}{2\pi} \ln \frac{\Lambda}{m_Z} \right)$$

$$b_{ij} = \begin{pmatrix} 199/25 & 27/5 & 88/5 \\ 9/5 & 25 & 24 \\ 11/5 & 9 & 14 \end{pmatrix}$$

one can find γ_i by iteratively updating $\alpha(\Lambda)$ and Λ

General solution to RGE

$$\frac{2\pi}{\alpha(\Lambda)}$$

↑
unified coupling

$$= \frac{2\pi}{\alpha_i(m_Z)}$$

↑
experimental input

1-loop MSSM

SUSY threshold

bulk 2-loop contribution

$$- b_i \ln \left(\frac{\Lambda}{m_Z} \right) + s_i + r_i + \gamma_i + \Delta_i$$

↑ SUSY threshold
 ↑ GUT threshold
 ↑ top-quark threshold, MSbar-DRbar conversion

$$s_i = \sum_{\eta} b_i^{\eta} \ln \left(\frac{m_{\eta}}{m_Z} \right)$$

$$r_i = \sum_{\xi} b_i^{\xi} \ln \left(\frac{m_{\xi}}{\Lambda} \right)$$

- We trade 3 exp inputs with the 3 constants: $[\alpha_1(m_Z), \alpha_2(m_Z), \alpha_3(m_Z)] \rightarrow [M_s^*, M_G^*, \alpha_G^*]$

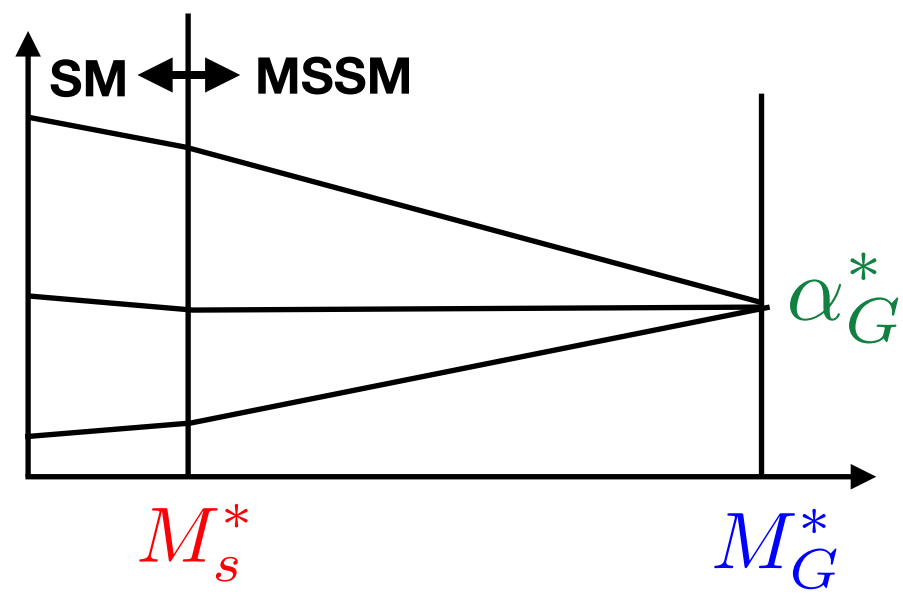
$$\frac{2\pi}{\alpha_i(m_Z)} = \frac{2\pi}{\alpha_G^*} + b_i \ln \left(\frac{M_G^*}{m_Z} \right) - \delta_i \ln \left(\frac{M_s^*}{m_Z} \right) - \gamma_i - \Delta_i$$

degenerate SUSY without GUT thres

$$b_i = \left(\frac{33}{5}, 1, -3 \right)$$

$$\delta_i \equiv \sum_{\eta} b_i^{\eta} = b_i - b_i^{\text{SM}}$$

$$= \left(\frac{2}{5}, \frac{25}{6}, 4 \right)$$



$$M_s^* = 2.13 \text{ TeV}$$

$$M_G^* = 1.26 \cdot 10^{16} \text{ GeV}$$

$$\alpha_G^{*-1} = 25.5$$

$$\boxed{\frac{2\pi}{\alpha_i(m_Z)}} = \frac{2\pi}{\alpha_G^*} + b_i \ln \left(\frac{M_G^*}{m_Z} \right) - \delta_i \ln \left(\frac{M_s^*}{m_Z} \right) - \gamma_i - \Delta_i$$

**degenerate SUSY
without GUT thres**

↓

$$\frac{2\pi}{\alpha(\Lambda)} = \boxed{\frac{2\pi}{\alpha_i(m_Z)}} - b_i \ln \left(\frac{\Lambda}{m_Z} \right) + s_i + r_i + \gamma_i + \Delta_i.$$

general solution

$$= \frac{2\pi}{\alpha_G^*} + b_i \ln \left(\frac{M_G^*}{\Lambda} \right) - \delta_i \ln \left(\frac{M_s^*}{m_Z} \right) + \overset{\substack{\uparrow \\ \text{SUSY threshold}}}{s_i} + \overset{\substack{\uparrow \\ \text{GUT threshold}}}{r_i}$$

$$M_s^* = 2.08 \text{ TeV}$$

$$M_G^* = 1.27 \cdot 10^{16} \text{ GeV}$$

$$\alpha_G^{*-1} = 25.5$$

$$\begin{aligned}
& \boxed{\frac{2\pi}{\alpha_i(m_Z)}} = \frac{2\pi}{\alpha_G^*} + b_i \ln \left(\frac{M_G^*}{m_Z} \right) - \delta_i \ln \left(\frac{M_s^*}{m_Z} \right) - \gamma_i - \Delta_i \quad \boxed{\text{degenerate SUSY}} \\
& \quad \downarrow \\
& \frac{2\pi}{\alpha(\Lambda)} = \boxed{\frac{2\pi}{\alpha_i(m_Z)}} - b_i \ln \left(\frac{\Lambda}{m_Z} \right) + s_i + r_i + \gamma_i + \Delta_i. \quad \boxed{\text{general solution}} \\
& = \frac{2\pi}{\alpha_G^*} + b_i \ln \left(\frac{M_G^*}{\Lambda} \right) - \delta_i \ln \left(\frac{M_s^*}{m_Z} \right) + \underset{\substack{\uparrow \\ \text{SUSY threshold}}}{s_i} + \underset{\substack{\uparrow \\ \text{GUT threshold}}}{r_i}
\end{aligned}$$

Any 3D vector can be decomposed into a sum of 3 independent vectors: $1, b_i, \delta_i$

$$\vec{1} = (1, 1, 1) \quad \vec{b} = \left(\frac{33}{5}, 1, -3 \right) \quad \vec{\delta} \equiv \vec{b} - \vec{b}_{\text{SM}} = \left(\frac{2}{5}, \frac{25}{6}, 4 \right)$$

$$s_i = \sum_{\eta} b_i^{\eta} \ln \left(\frac{m_{\eta}}{m_Z} \right) = C_S + b_i \ln \Omega_S + \delta_i \ln \left(\frac{T_S}{m_Z} \right)$$

$$r_i = \sum_{\xi} b_i^{\xi} \ln \left(\frac{m_{\xi}}{\Lambda} \right) = C_G - b_i \ln \left(\frac{T_G}{\Lambda} \right) - \delta_i \ln \Omega_G$$

$$\boxed{\frac{2\pi}{\alpha_i(m_Z)}} = \frac{2\pi}{\alpha_G^*} + b_i \ln \left(\frac{M_G^*}{m_Z} \right) - \delta_i \ln \left(\frac{M_s^*}{m_Z} \right) - \gamma_i - \Delta_i \quad \boxed{\text{degenerate SUSY}}$$

↓

$$\frac{2\pi}{\alpha(\Lambda)} = \boxed{\frac{2\pi}{\alpha_i(m_Z)}} - b_i \ln \left(\frac{\Lambda}{m_Z} \right) + s_i + r_i + \gamma_i + \Delta_i. \quad \boxed{\text{general solution}}$$

i-independent

$$= \frac{2\pi}{\alpha_G^*} + b_i \ln \left(\frac{M_G^*}{\Lambda} \right) - \delta_i \ln \left(\frac{M_s^*}{m_Z} \right) + s_i + r_i$$

$$= \left[\frac{2\pi}{\alpha_G^*} + C_S + C_G \right] + b_i \ln \left(\frac{M_G^* \Omega_S}{T_G} \right) + \delta_i \ln \left(\frac{T_S}{M_s^* \Omega_G} \right)$$

Any 3D vector can be decomposed into a sum of 3 independent vectors:

$$\vec{1} = (1, 1, 1) \quad \vec{b} = \left(\frac{33}{5}, 1, -3 \right) \quad \vec{\delta} \equiv \vec{b} - \vec{b}_{\text{SM}} = \left(\frac{2}{5}, \frac{25}{6}, 4 \right)$$

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$$\boxed{\frac{2\pi}{\alpha_i(m_Z)}} = \frac{2\pi}{\alpha_G^*} + b_i \ln \left(\frac{M_G^*}{m_Z} \right) - \delta_i \ln \left(\frac{M_s^*}{m_Z} \right) - \gamma_i - \Delta_i \quad \boxed{\text{degenerate SUSY}}$$

$$\frac{2\pi}{\alpha(\Lambda)} = \boxed{\frac{2\pi}{\alpha_i(m_Z)}} - b_i \ln \left(\frac{\Lambda}{m_Z} \right) + s_i + r_i + \gamma_i + \Delta_i. \quad \boxed{\text{general solution}}$$

$$= \frac{2\pi}{\alpha_G^*} + b_i \ln \left(\frac{M_G^*}{\Lambda} \right) - \delta_i \ln \left(\frac{M_s^*}{m_Z} \right) + s_i + r_i$$

must vanish

$$= \left[\frac{2\pi}{\alpha_G^*} + C_S + C_G \right] + b_i \ln \left(\frac{M_G^* \Omega_S}{T_G} \right) + \delta_i \ln \left(\frac{T_S}{M_s^* \Omega_G} \right)$$

The condition of gauge coupling unification:

$$T_S = M_s^* \Omega_G \cap T_G = M_G^* \Omega_S$$

$$M_s^* = 2.13 \text{ TeV}$$

$$M_G^* = 1.26 \cdot 10^{16} \text{ GeV}$$

$$\alpha_G^{*-1} = 25.5$$

The unified coupling at Λ

$$\alpha^{-1}(\Lambda) = \alpha_G^{*-1} + \frac{1}{2\pi} (C_S + C_G)$$

$$s_i = \sum_{\eta} b_i^{\eta} \ln \left(\frac{m_{\eta}}{m_Z} \right) = \sum_{\eta=\tilde{g}, \tilde{q}, \dots} C_S b_i^{\eta} \ln \left(\frac{m_{\eta}}{m_Z} \right) + C$$

$$b_i = \left(\frac{33}{5}, 1, -3 \right)$$

$$\delta_i = \left(\frac{2}{5}, \frac{25}{6}, 4 \right)$$

Acknowledgments

We thank Hans Peter Nilles^{b1} for useful discussion and comments on the manuscript. KS thanks Zackaria Chacko, Kiwoon Choi and Shigeki Matsumoto for helpful discussion. KS thanks the TU Munich for hospitality during the final stages of this work and has been partially supported by the DFG cluster of excellence EXC 153 ‘Origin and Structure of the Universe’, by the Collaborative Research Center SFB1258. The work of SP and KS is partially supported by the National Science Centre, Poland, under research grants DEC-2014/15/B/ST2/02157 and DEC-2015/18/M/ST2/00054. The work of KR and KS is supported by the National Science Centre (Poland) under Grant 2015/19/D/ST2/03136.

A

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$$T_S = \left[M_3^{-28} M_2^{32} \mu^{12} m_A^3 X_T \right]^{\frac{1}{19}}$$

$$X_T \equiv \prod_{i=1\dots 3} \left(\frac{m_{\tilde{l}_i}^3}{m_{\tilde{d}_{Ri}}^3} \right) \left(\frac{m_{\tilde{q}_i}^7}{m_{\tilde{e}_{Ri}}^2 m_{\tilde{u}_{Ri}}^5} \right)$$

$$\Omega_S = \left[M_3^{-100} M_2^{60} \mu^{32} m_A^8 X_{\Omega} \right]^{\frac{1}{288}}$$

$$X_{\Omega} \equiv \prod_{i=1\dots 3} \left(\frac{m_{\tilde{l}_i}^8}{m_{\tilde{d}_{Ri}}^8} \right) \left(\frac{m_{\tilde{q}_i}^6 m_{\tilde{e}_{Ri}}}{m_{\tilde{u}_{Ri}}^7} \right)$$

$$C_S = \frac{125}{19} \ln M_3 - \frac{113}{19} \ln M_2 - \frac{40}{19} \ln \mu - \frac{10}{19} \ln m_A$$

$$+ \sum_{i=1\dots 3} \left[\frac{79}{114} \ln m_{\tilde{d}_{Ri}} - \frac{10}{19} \ln m_{\tilde{l}_i} - \frac{121}{114} \ln m_{\tilde{q}_i} + \frac{257}{228} \ln m_{\tilde{u}_{Ri}} + \frac{33}{76} \ln m_{\tilde{e}_{Ri}} \right]$$

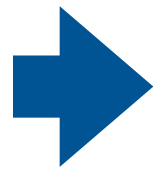
are the
tributi

ues of β_a are shown in Table 1. The β_a 's are to be understood as the components of a three

dimensional vector and expanded in terms of the three independent vectors δ_a , b_a and $(1, 1, 1)$,

where $\delta \equiv b - b^{\text{SM}} = (\frac{2}{5}, \frac{25}{6}, 4)$ is the difference between b and the SM β -function coefficient

$$r_i = \sum_{\xi} b_i^{\xi} \ln \left(\frac{m_{\xi}}{\Lambda} \right) = C_G - b_i \ln \left(\frac{T_G}{\Lambda} \right) - \delta_i \ln \Omega_G$$



$$\ln \left(\frac{T_G}{\Lambda} \right) = \sum_{\xi} \left(-\frac{5}{288} b_1^{\xi} - \frac{15}{76} b_2^{\xi} + \frac{25}{114} b_3^{\xi} \right) \ln \left(\frac{m_{\xi}}{\Lambda} \right)$$

$$\ln \Omega_G = \sum_{\xi} \left(\frac{10}{19} b_1^{\xi} - \frac{24}{19} b_2^{\xi} + \frac{14}{19} b_3^{\xi} \right) \ln \left(\frac{m_{\xi}}{\Lambda} \right)$$

$$C_G = \sum_{\xi} \left(\frac{165}{76} b_1^{\xi} - \frac{339}{76} b_2^{\xi} + \frac{125}{38} b_3^{\xi} \right) \ln \left(\frac{m_{\xi}}{\Lambda} \right)$$

The condition of gauge coupling unification:

$$T_S = M_s^* \Omega_G \cap T_G = M_G^* \Omega_S$$

$$M_s^* = 2.08 \text{ TeV}$$

$$M_G^* = 1.27 \cdot 10^{16} \text{ GeV}$$

$$\alpha_G^{*-1} = 25.5$$

The unified coupling at Λ

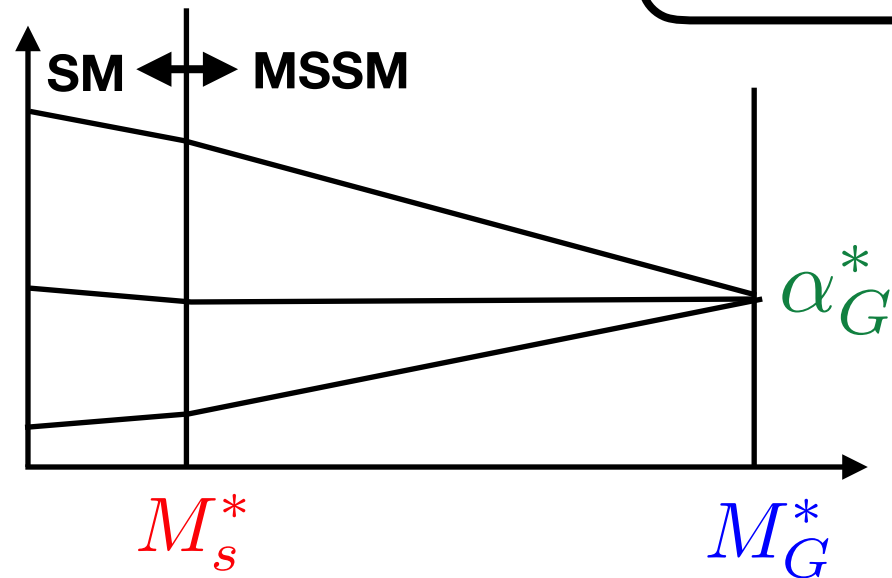
$$\alpha^{-1}(\Lambda) = \alpha_G^{*-1} + \frac{1}{2\pi} (C_S + C_G)$$

Uncertainty of $\alpha_s(m_Z)$

$$\alpha_s^0 \quad \Delta\alpha_s$$

$$\alpha_s(m_Z) = 0.1183 \pm 0.0008$$

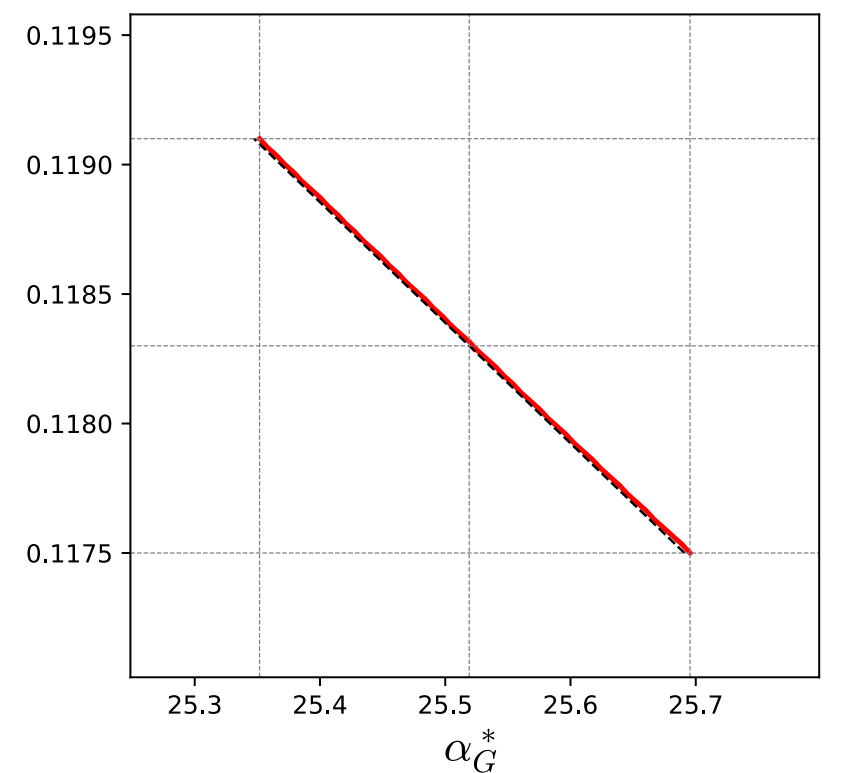
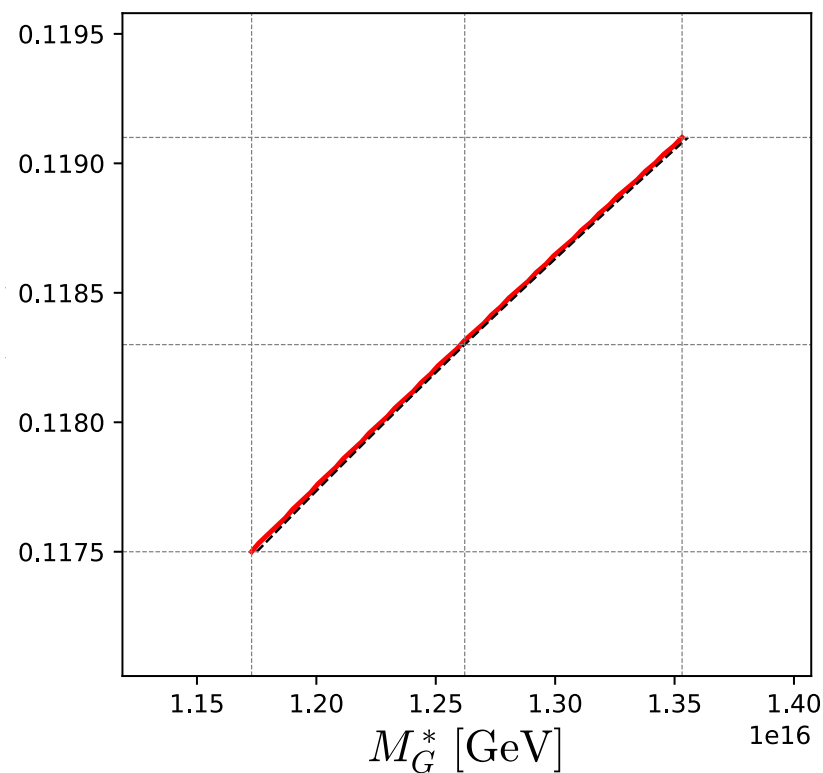
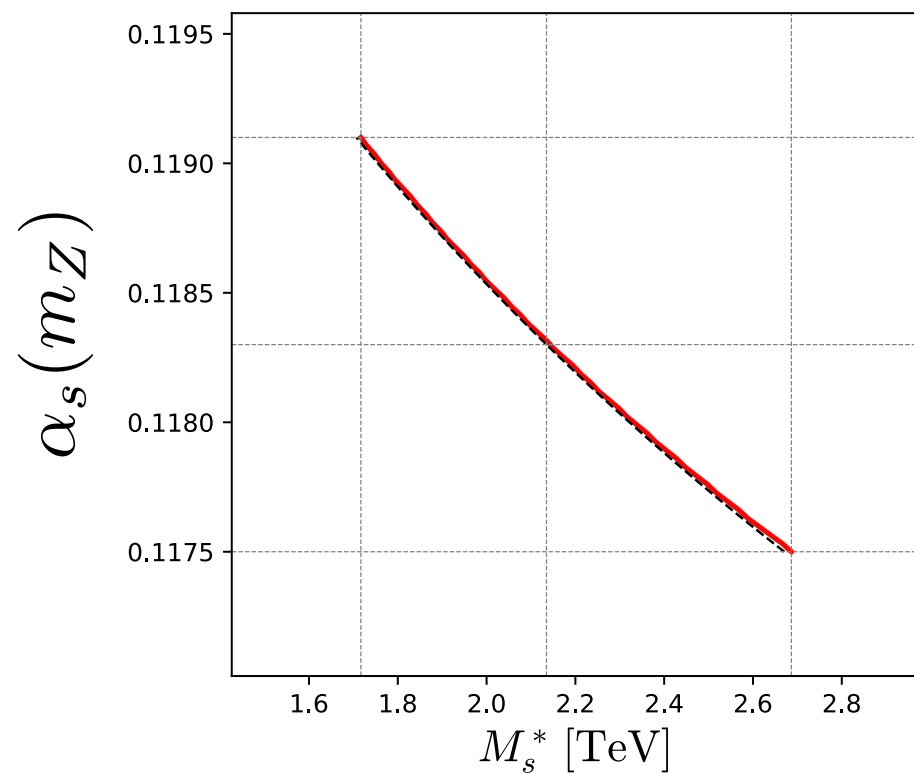
D. d'Enterria
[1806.06156]



$$\frac{M_s^*}{\text{TeV}} = \frac{2.13}{\text{TeV}} \cdot \exp \left[-0.224 \left(\frac{\alpha_s - \alpha_s^0}{\Delta\alpha_s} \right) \right],$$

$$\frac{M_G^*}{\text{GeV}} = \frac{1.26 \cdot 10^{16}}{\text{GeV}} \cdot \exp \left[0.0715 \left(\frac{\alpha_s - \alpha_s^0}{\Delta\alpha_s} \right) \right]$$

$$\alpha_G^{*-1} = 25.5 - 0.172 \left(\frac{\alpha_s - \alpha_s^0}{\Delta\alpha_s} \right).$$



$$M_s^* \in [2.69, 1.72] \text{ TeV}$$

$$M_G^* \in [1.17, 1.35] \cdot 10^{16} \text{ GeV}$$

$$\alpha_G^{*-1} \in [25.7, 25.4]$$

Minimal SU(5)

Higgs and gage fields

$$\begin{aligned}
 H(\mathbf{5}) &= (H_C, H_u) & \Sigma(\mathbf{24}) &= (\Sigma_8, \Sigma_3, \Sigma_1, \Sigma_{(2,3)}, \Sigma_{(2,3^*)}) \\
 \bar{H}(\bar{\mathbf{5}}) &= (\bar{H}_C, H_d) & \mathcal{V}(\mathbf{24}) &= (G, W, B, (X, Y), (X, Y)^\dagger)
 \end{aligned}$$

$$W_H = \frac{1}{2} M \text{Tr} \Sigma^2 + \frac{1}{3} \lambda_\Sigma \text{Tr} \Sigma^3 + \lambda_H \bar{H} (\Sigma + 3V) \Sigma H,$$

$$\uparrow$$

$$\langle \Sigma \rangle = V \cdot \text{diag}(2, 2, 2, -3, -3)$$

doublet triplet splitting

$$\begin{aligned}
 M_{H_C} &= 5\lambda_H V \\
 M_{H_{u/d}} &= 0
 \end{aligned}$$

$$M_V = 5\sqrt{2}g_5 V$$

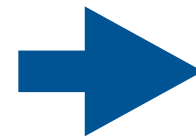
$$M_\Sigma = \frac{2}{5}\lambda_\Sigma V$$

mass	$(U(1) \times SU(2) \times SU(3))$	(b_1, b_2, b_3)
M_{H_C}	$(-\frac{1}{3}, \mathbf{1}, \mathbf{3}), (\frac{1}{3}, \mathbf{1}, \bar{\mathbf{3}})$	$(\frac{2}{5}, 0, 1)$
M_V	$(-\frac{5}{6}, \mathbf{2}, \mathbf{3}), (\frac{5}{6}, \mathbf{2}, \bar{\mathbf{3}})$	$(-10, -6, -4)$
M_Σ	$(0, \mathbf{3}, \mathbf{1}), (0, \mathbf{1}, \mathbf{8})$	$(0, 2, 3)$

formulae for T_G and Ω_G

$$\ln\left(\frac{T_G}{\Lambda}\right) = \sum_{\xi} \left(-\frac{5}{288}b_1^{\xi} - \frac{15}{76}b_2^{\xi} + \frac{25}{114}b_3^{\xi} \right) \ln\left(\frac{m_{\xi}}{\Lambda}\right)$$

$$\ln \Omega_G = \sum_{\xi} \left(\frac{10}{19}b_1^{\xi} - \frac{24}{19}b_2^{\xi} + \frac{14}{19}b_3^{\xi} \right) \ln\left(\frac{m_{\xi}}{\Lambda}\right)$$



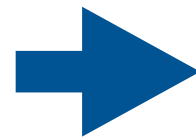
$$T_G = M_{H_C}^{\frac{4}{19}} (M_V^2 M_\Sigma)^{\frac{5}{19}}$$

$$\Omega_G = M_{H_C}^{\frac{18}{19}} (M_V^2 M_\Sigma)^{-\frac{6}{19}}$$

Condition for GCU

$$T_S = M_s^* \Omega_G$$

$$T_G = M_G^* \Omega_S$$



$$M_{H_C} = M_G^* \Omega_S \left(\frac{T_S}{M_s^*} \right)^{\frac{5}{6}}$$

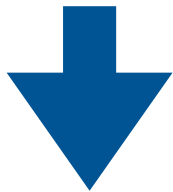
$$(M_V^2 M_\Sigma)^{\frac{1}{3}} = M_G^* \Omega_S \left(\frac{T_S}{M_s^*} \right)^{-\frac{2}{9}}$$

$$(3\alpha_2^{-1} - 2\alpha_3^{-1} - \alpha_1^{-1})(m_Z) = \frac{1}{2\pi} \left\{ \frac{12}{5} \ln \frac{M_{H_C}}{m_Z} - 2 \ln \frac{m_{SUSY}}{m_Z} \right\},$$

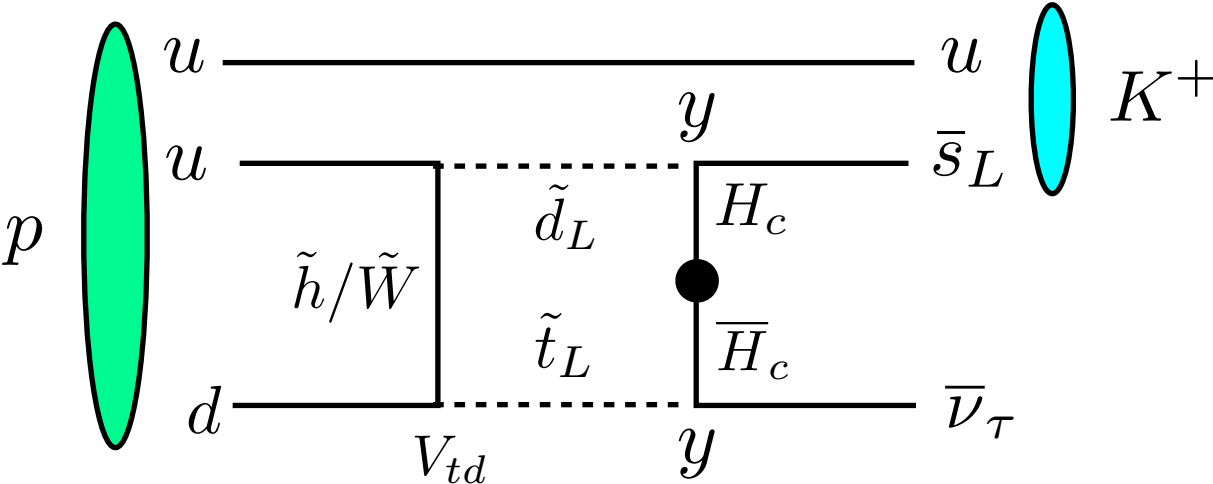
$$(5\alpha_1^{-1} - 3\alpha_2^{-1} - 2\alpha_3^{-1})(m_Z) = \frac{1}{2\pi} \left\{ 12 \ln \frac{M_V^2 M_\Sigma}{m_Z^3} + 8 \ln \frac{m_{SUSY}}{m_Z} \right\}.$$

Hisano,
Murayama,
Yanagida '92

$$M_{HC} = M_G^* \Omega_S \left(\frac{T_S}{M_s^*} \right)^{\frac{5}{6}}$$



D=5 proton decay



Hisano, Kobayashi,
Kuwahara, Nagata '13

$$T_S = \left[M_3^{-28} M_2^{32} \mu^{12} m_A^3 X_T \right]^{\frac{1}{19}}$$

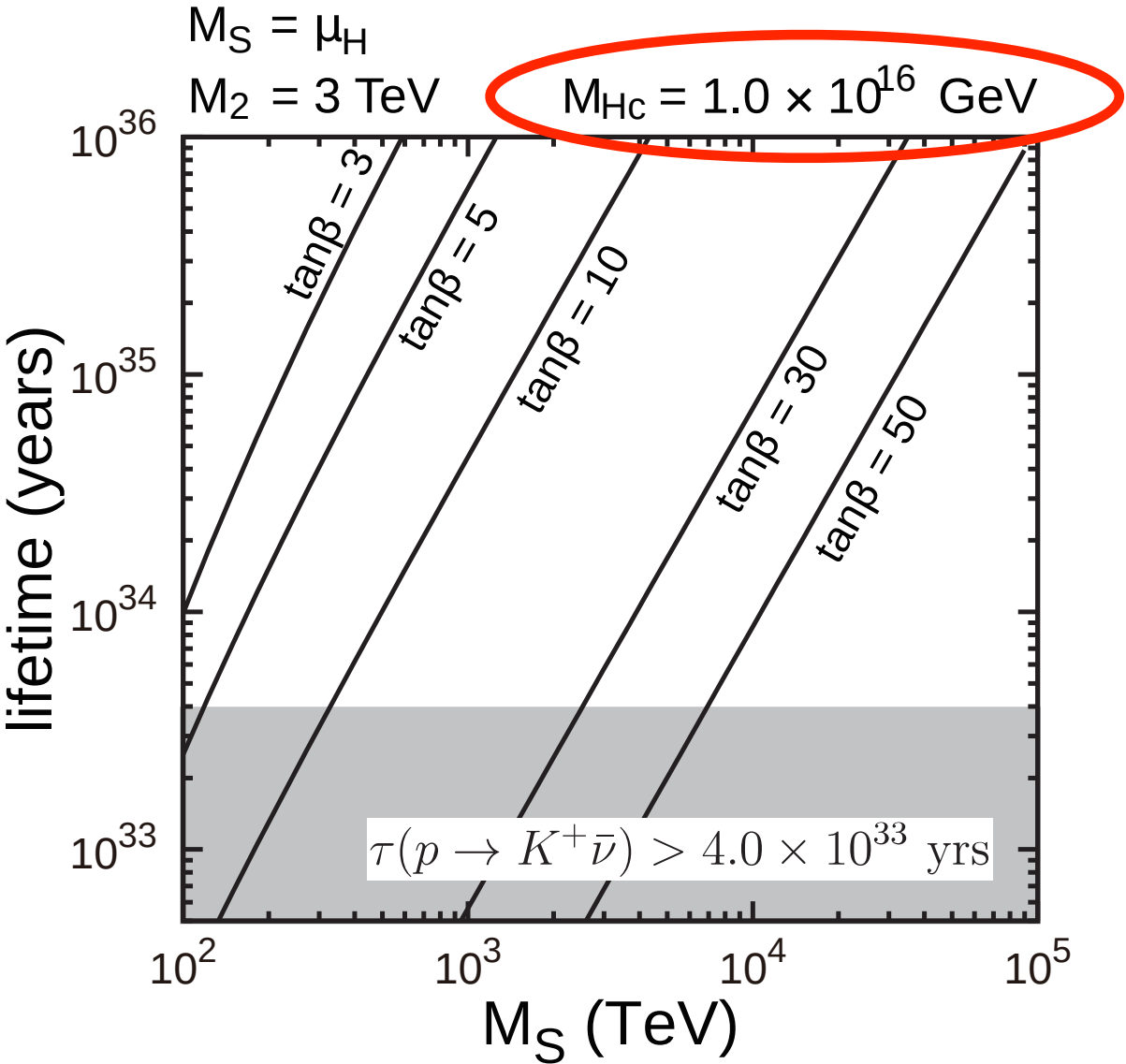
$$\Omega_S = \left[M_3^{-100} M_2^{60} \mu^{32} m_A^8 X_\Omega \right]^{\frac{1}{288}}$$

$$M_s^* = 2.08 \text{ TeV}$$

$$M_G^* = 1.27 \cdot 10^{16} \text{ GeV}$$

$$X_T \equiv \prod_{i=1 \dots 3} \left(\frac{m_{\tilde{l}_i}^3}{m_{\tilde{d}_{Ri}}^3} \right) \left(\frac{m_{\tilde{q}_i}^7}{m_{\tilde{e}_{Ri}}^2 m_{\tilde{u}_{Ri}}^5} \right)$$

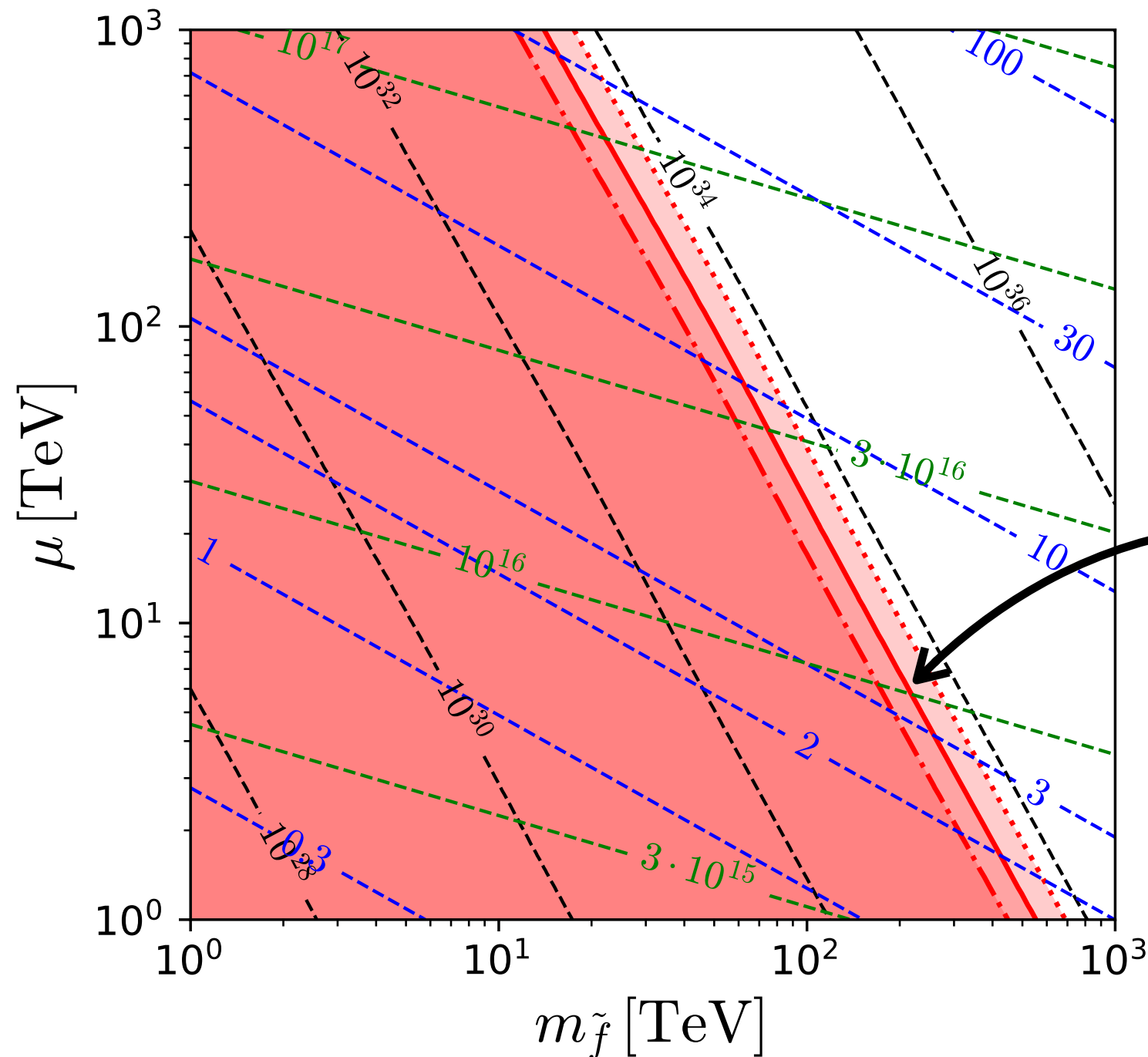
$$X_\Omega \equiv \prod_{i=1 \dots 3} \left(\frac{m_{\tilde{l}_i}^8}{m_{\tilde{d}_{Ri}}^8} \right) \left(\frac{m_{\tilde{q}_i}^6 m_{\tilde{e}_{Ri}}}{m_{\tilde{u}_{Ri}}^7} \right)$$



Vanilla SUSY

$$(M_V = M_\Sigma)$$

$$M_3 = m_{\tilde{f}} = m_A = 3M_2, \quad \tan\beta = 2$$



$$- - - \tau(p \rightarrow K^+ \bar{\nu})/\text{yrs}$$

$$- - - M_{H_C}/\text{GeV}$$

$$- - - T_S/\text{TeV}$$

$$\tau(p \rightarrow K^+ \bar{\nu}) > 4.0 \times 10^{33} \text{ yrs}$$

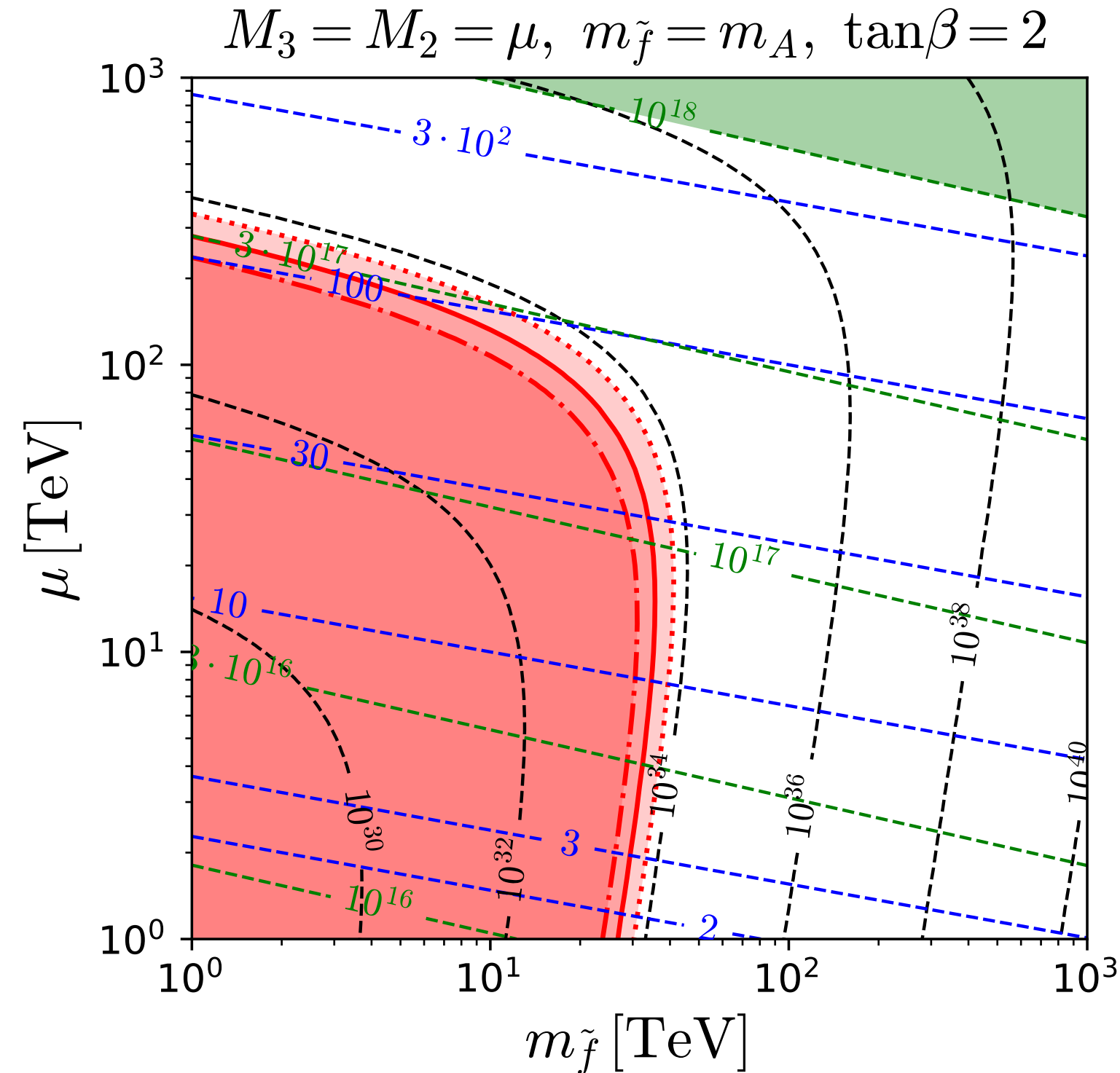
$$\alpha_s(m_Z) = 0.1183 \pm 0.0008$$

$$\Lambda \equiv \max\{M_{H_C}, M_V\}$$

$$24.7 < \alpha^{-1}(\Lambda) < 31.8 \quad (M_V = M_\Sigma)$$

$$3.77 \cdot 10^{15} \text{ GeV} < (M_V^2 M_\Sigma)^{\frac{1}{3}} < 1.83 \cdot 10^{16} \text{ GeV}$$

Universal Gaugino Mass



--- $\tau(p \rightarrow K^+ \bar{\nu})/\text{yrs}$

--- M_{H_C}/GeV

--- T_S/TeV

$\tau(p \rightarrow K^+ \bar{\nu}) > 4.0 \times 10^{33} \text{ yrs}$

$\alpha_s(m_Z) = 0.1183 \pm 0.0008$

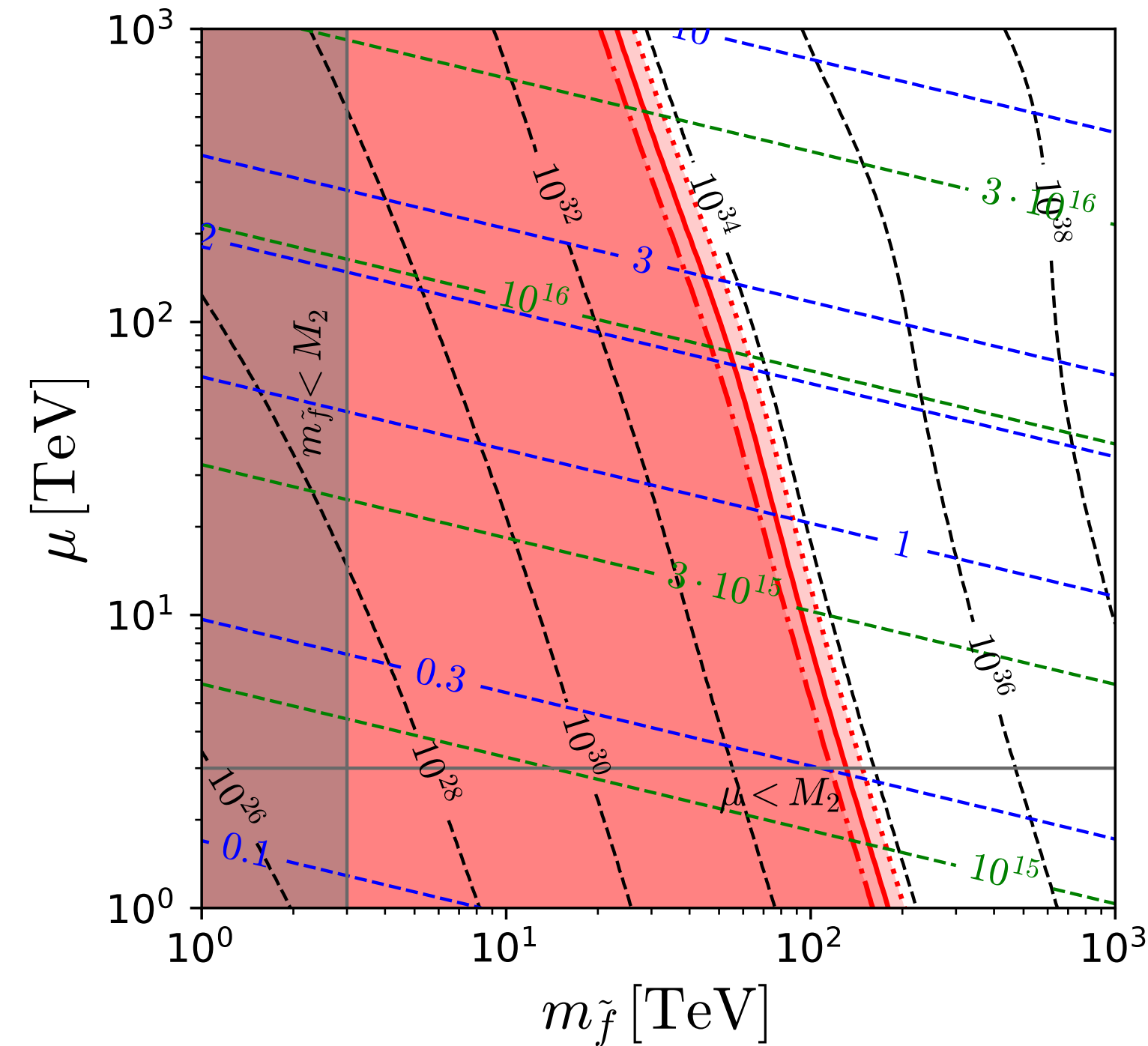
$\Lambda \equiv \max\{M_{H_C}, M_V\}$

$25.0 < \alpha^{-1}(\Lambda) < 33.0 \quad (M_V = M_\Sigma)$

$3.14 \cdot 10^{15} \text{ GeV} < (M_V^2 M_\Sigma)^{\frac{1}{3}} < 1.52 \cdot 10^{16} \text{ GeV}$

Pure Gravity Mediation

$$M_2 = 3 \text{ TeV}, M_3 = 7M_2, m_A = m_{\tilde{f}}, \tan\beta = 2$$



$$- - - \tau(p \rightarrow K^+ \bar{\nu})/\text{yrs}$$

$$- - - M_{H_C}/\text{GeV}$$

$$- - - T_S/\text{TeV}$$

$$\tau(p \rightarrow K^+ \bar{\nu}) > 4.0 \times 10^{33} \text{ yrs}$$

$$\alpha_s(m_Z) = 0.1183 \pm 0.0008$$

$$\Lambda \equiv \max\{M_{H_C}, M_V\}$$

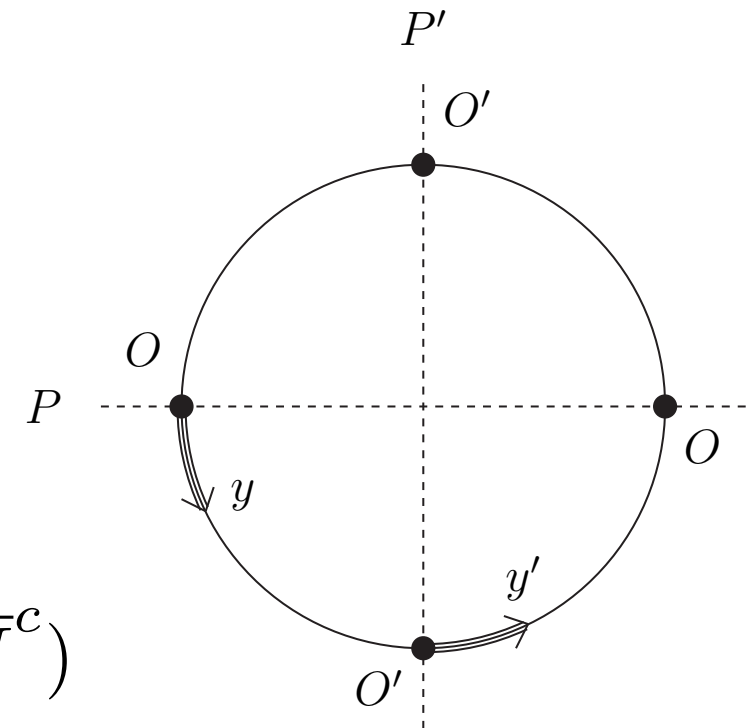
$$25.3 < \alpha^{-1}(\Lambda) < 29.5 \quad (M_V = M_\Sigma)$$

$$8.90 \cdot 10^{15} \text{ GeV} < (M_V^2 M_\Sigma)^{\frac{1}{3}} < 1.19 \cdot 10^{16} \text{ GeV}$$

Orbifold SU(5)

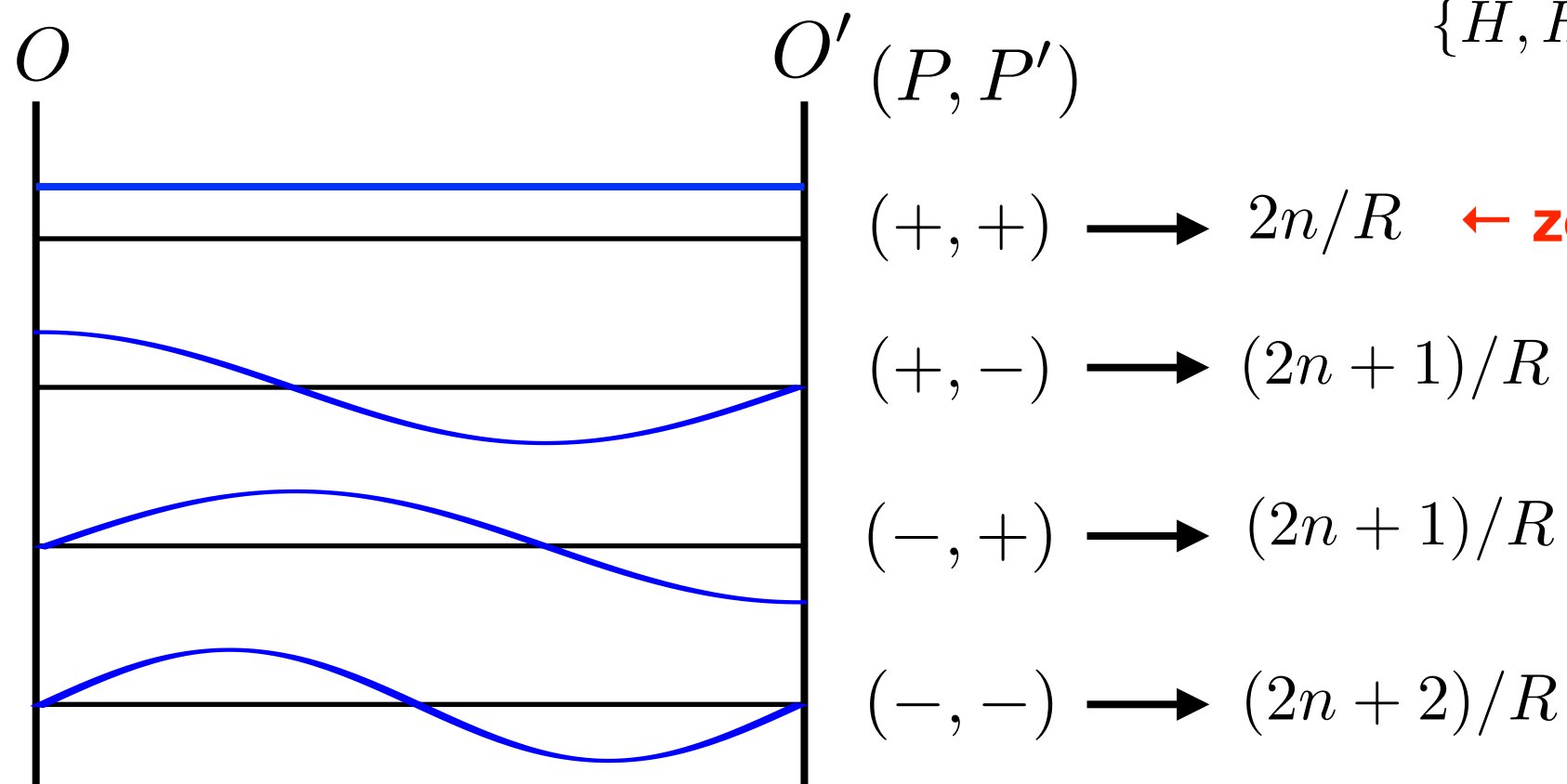
[Hall, Nomura '01]

- space-time = 4d Minkowski $\times S^1/(Z_2 \times Z_2')$
- two fixed points O, O' , & two orbifold parities P, P'
- N=2 4d SUSY in the bulk (8 spinor dof)
- 5d vector multiplets: $\mathcal{V}(\mathbf{24}) = (V, \Sigma)$
- 5d hipermultiplets: $\mathcal{H}(\mathbf{5}) = (H, H^c), \bar{\mathcal{H}}(\bar{\mathbf{5}}) = (\bar{H}, \bar{H}^c)$



$$\{H, H^c\} \supset \{(H_F, H_C), (H_F^c, H_C^c)\}$$

$$\{\bar{H}, \bar{H}^c\} \supset \{(H_{\bar{F}}, H_{\bar{C}}), (H_{\bar{F}}^c, H_{\bar{C}}^c)\}$$



$a =$ unbroken generators $\hat{a} =$ broken generators $n = 0, 1, 2, \dots$

KK mode	mass	(P, P')	4d fields	$\sum(b_1, b_2, b_3)$
zero	0	$(+, +)$	$V^a, H_F, H_{\bar{F}}$	
even	$(2n + 2)/R$	$(+, +)$	$V^a, H_F, H_{\bar{F}}$	$(\frac{6}{5}, -2, -6)$
		$(-, -)$	$\Sigma^a, H_F^c, H_{\bar{F}}^c$	
odd	$(2n + 1)/R$	$(+, -)$	$V^{\hat{a}}, H_C, H_{\bar{C}}$	$(-\frac{46}{5}, -6, -2)$
		$(-, +)$	$\Sigma^{\hat{a}}, H_C^c, H_{\bar{C}}^c$	

- doublet-triplet splitting can be achieved by the parity assignment
- X,Y fields are absent at the 3-brane at O' => SU(5) is explicitly broken at O'
- non-universal contribution to the gauge couplings from the 3-brane at O'

$$f^a \mathcal{W}^a \mathcal{W}^a \Big|_{O'}$$

the effect is negligible due to the dominant bulk contribution if the extra dimension is large, $R\Lambda > 4$ [Hall, Nomura '01]

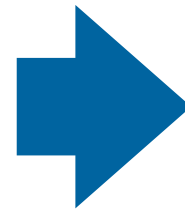
- The success of charge quantisation in 4d GUTs can be kept by placing the matter fields at the SU(5) symmetric brane at O .
- D=5 proton decay is absent due to an accidental $U(1)_R$ sym.

	Σ	H_5	$H_{\bar{5}}$	H_5^c	$H_{\bar{5}}^c$	T_{10}	$F_{\bar{5}}$	N_1
$U(1)_R$	0	0	0	2	2	1	1	1

- Above the KK mass scale, $M_c = 1/R$, KK modes appear and the three gauge couplings approach each other. They finally unify at the cut-off scale, Λ , where the 5d theory is incorporated into a more fundamental theory.

$$\ln\left(\frac{T_G}{\Lambda}\right) = \sum_{\xi} \left(-\frac{5}{288}b_1^{\xi} - \frac{15}{76}b_2^{\xi} + \frac{25}{114}b_3^{\xi} \right) \ln\left(\frac{m_{\xi}}{\Lambda}\right)$$

$$\ln \Omega_G = \sum_{\xi} \left(\frac{10}{19}b_1^{\xi} - \frac{24}{19}b_2^{\xi} + \frac{14}{19}b_3^{\xi} \right) \ln\left(\frac{m_{\xi}}{\Lambda}\right)$$



the total contributions from
(2k+1) and (2k+2) modes

$$c_o \ln \frac{2k+1}{R\Lambda} \quad c_e \ln \frac{2k+2}{R\Lambda}$$

$$c_o = -c_e = \frac{24}{19} \text{ for } \ln \Omega_G \quad c_o = -c_e = \frac{18}{19} \text{ for } \ln(T_G/\Lambda) \quad c_o = \frac{4}{19}, c_e = -\frac{156}{19} \text{ for } C_G$$

The index k runs from 0 to $k_{\max}^{o/e}$, where $\begin{cases} 2k_{\max}^o + 1 < R\Lambda & \text{for odd modes} \\ 2k_{\max}^e + 2 < R\Lambda & \text{for even modes} \end{cases}$

$$\Omega_G = \left[\frac{\prod_k^{k_{\max}^o} (2k+1)}{\prod_k^{k_{\max}^e} (2k+2)} \left(\frac{1}{r}\right)^{k_{\max}^o - k_{\max}^e} \right]^{\frac{24}{19}},$$

$$\frac{T_G}{\Lambda} = \left[\frac{\prod_k^{k_{\max}^o} (2k+1)}{\prod_k^{k_{\max}^e} (2k+2)} \left(\frac{1}{r}\right)^{k_{\max}^o - k_{\max}^e} \right]^{\frac{18}{19}},$$

$$(r \equiv R\Lambda)$$

$$C_G = \frac{4}{19} \ln \left[\prod_{k=0}^{k_{\max}^o} \frac{2k+1}{r} \right] - \frac{156}{19} \ln \left[\prod_{k=0}^{k_{\max}^e} \frac{2k+2}{r} \right]$$



$(r \equiv R\Lambda)$

$r \in$	$(1, 2]$	$(2, 3]$	$(3, 4]$	$(4, 5]$	$\dots$
Ω_G	$\left[\frac{1}{r}\right]^{\frac{24}{19}}$	$\left[\frac{1}{2}\right]^{\frac{24}{19}}$	$\left[\frac{1\cdot3}{2} \frac{1}{r}\right]^{\frac{24}{19}}$	$\left[\frac{1\cdot3}{2\cdot4}\right]^{\frac{24}{19}}$	$\dots$
T_G/Λ	$\left[\frac{1}{r}\right]^{\frac{18}{19}}$	$\left[\frac{1}{2}\right]^{\frac{18}{19}}$	$\left[\frac{1\cdot3}{2} \frac{1}{r}\right]^{\frac{18}{19}}$	$\left[\frac{1\cdot3}{2\cdot4}\right]^{\frac{18}{19}}$	$\dots$

$T_G/\Lambda = \Omega_G^{\frac{3}{4}}$

GCU condition

$$T_S = M_s^* \Omega_G$$

$$T_G = M_G^* \Omega_S$$

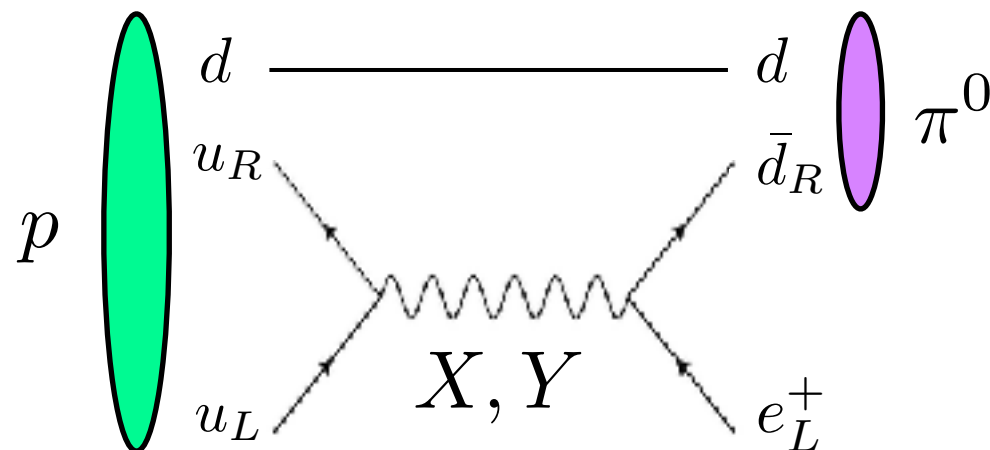
$$T_S = \left[M_3^{-28} M_2^{32} \mu^{12} m_A^3 X_T \right]^{\frac{1}{19}} = M_s^* \Omega_G \quad \text{non-trivial constraint on SUSY masses}$$

$$M_c = M_G^* \Omega_G^{-\frac{3}{4}} \Omega_S / r$$

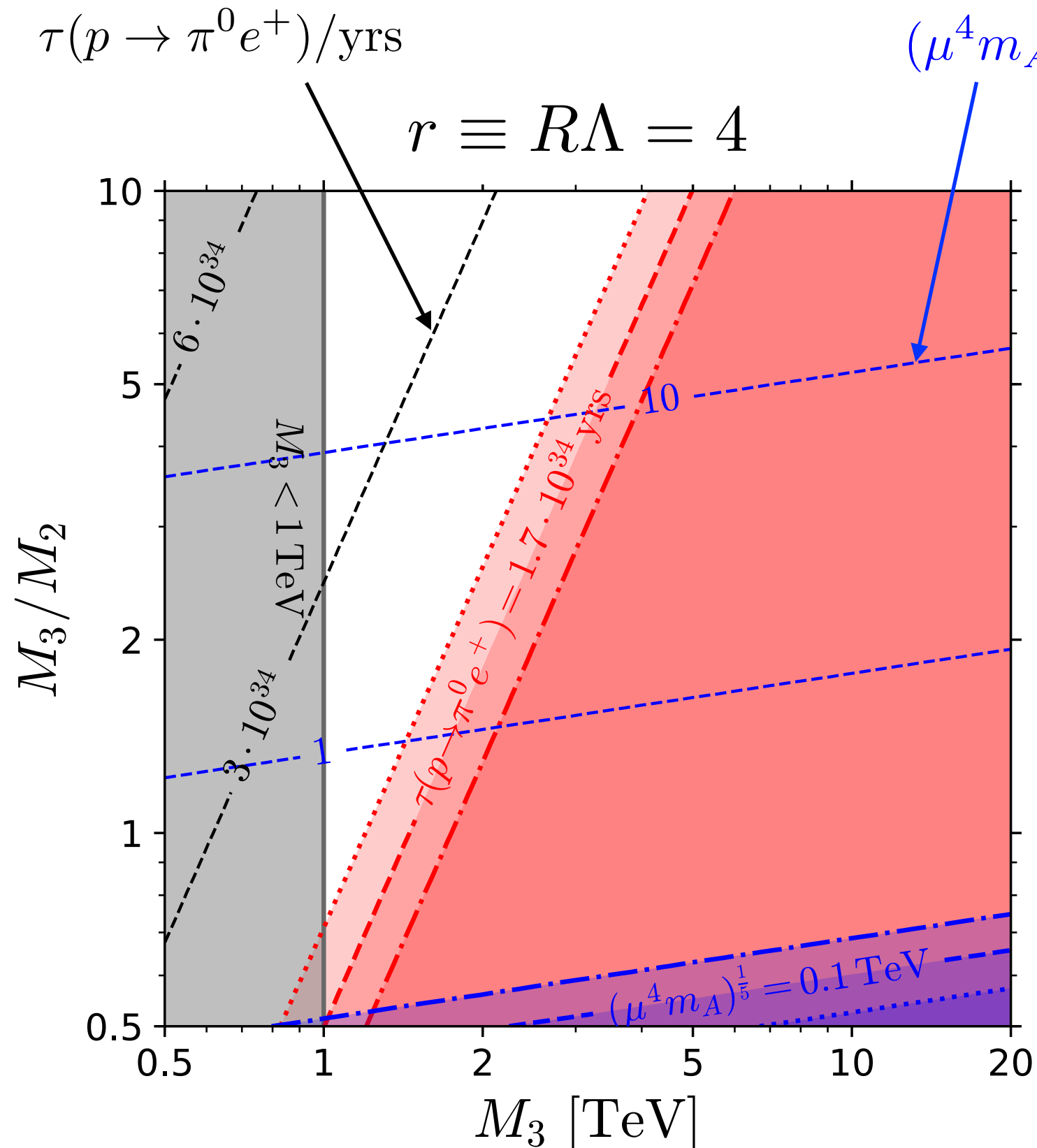
$$= M_G^* M_s^* \frac{19}{108} \Omega_G^{-\frac{31}{54}} M_3^{-\frac{19}{216}} M_2^{-\frac{19}{216}} X_T^{-\frac{1}{108}} X_\Omega^{-\frac{1}{288}} / r$$

compactification scale can be predicted from SUSY spectrum
and $r \Rightarrow$ allows to predict D=6 proton decay

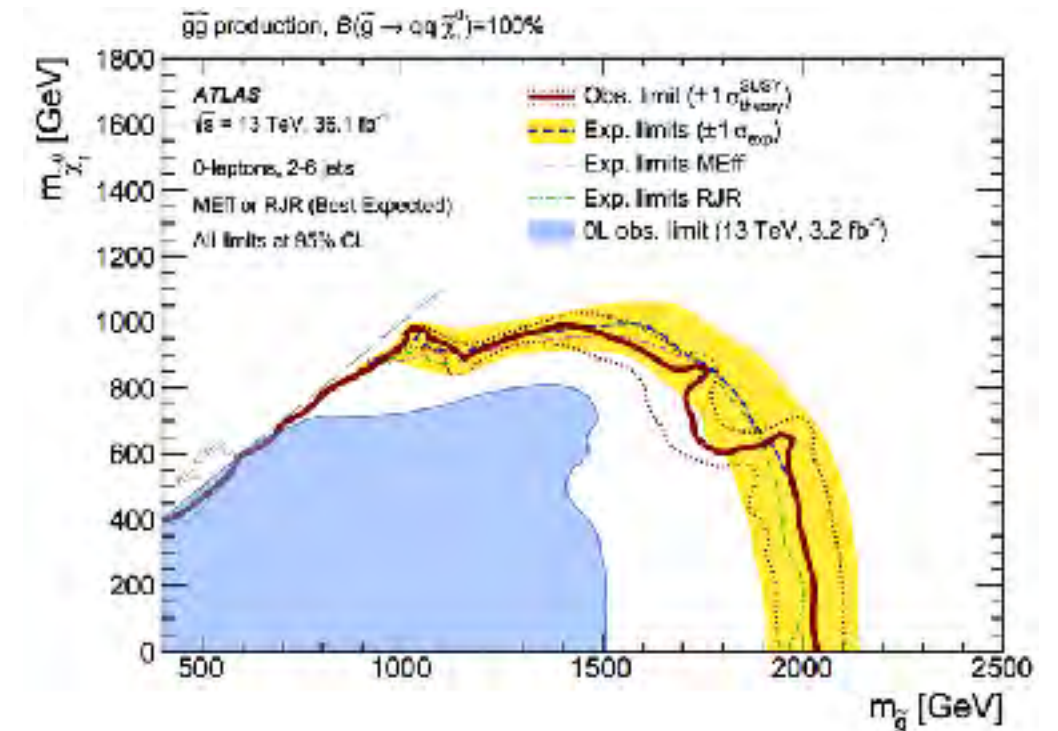
D=6 proton decay



A SUSY plane with GCU

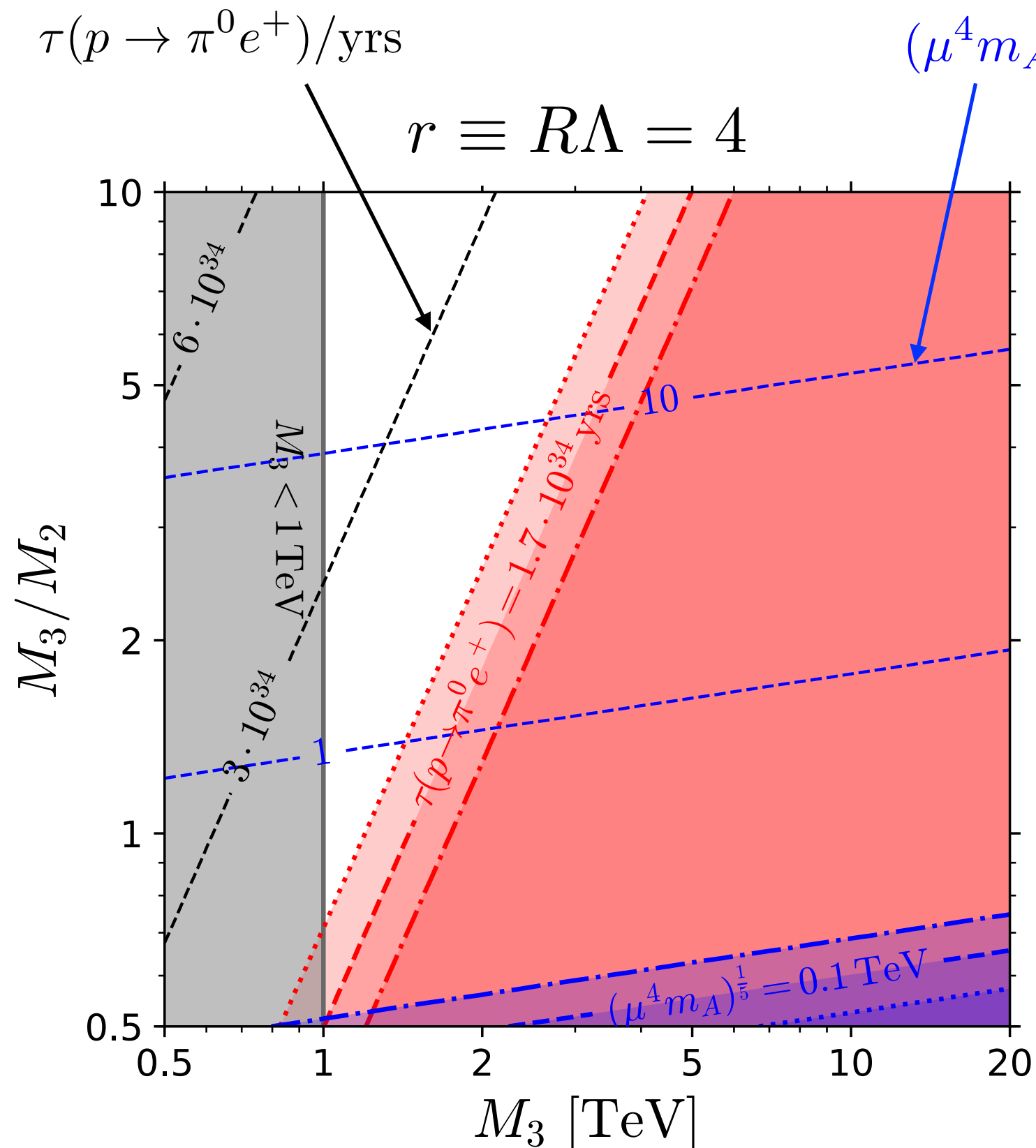


*) universal sfermion mass is assumed

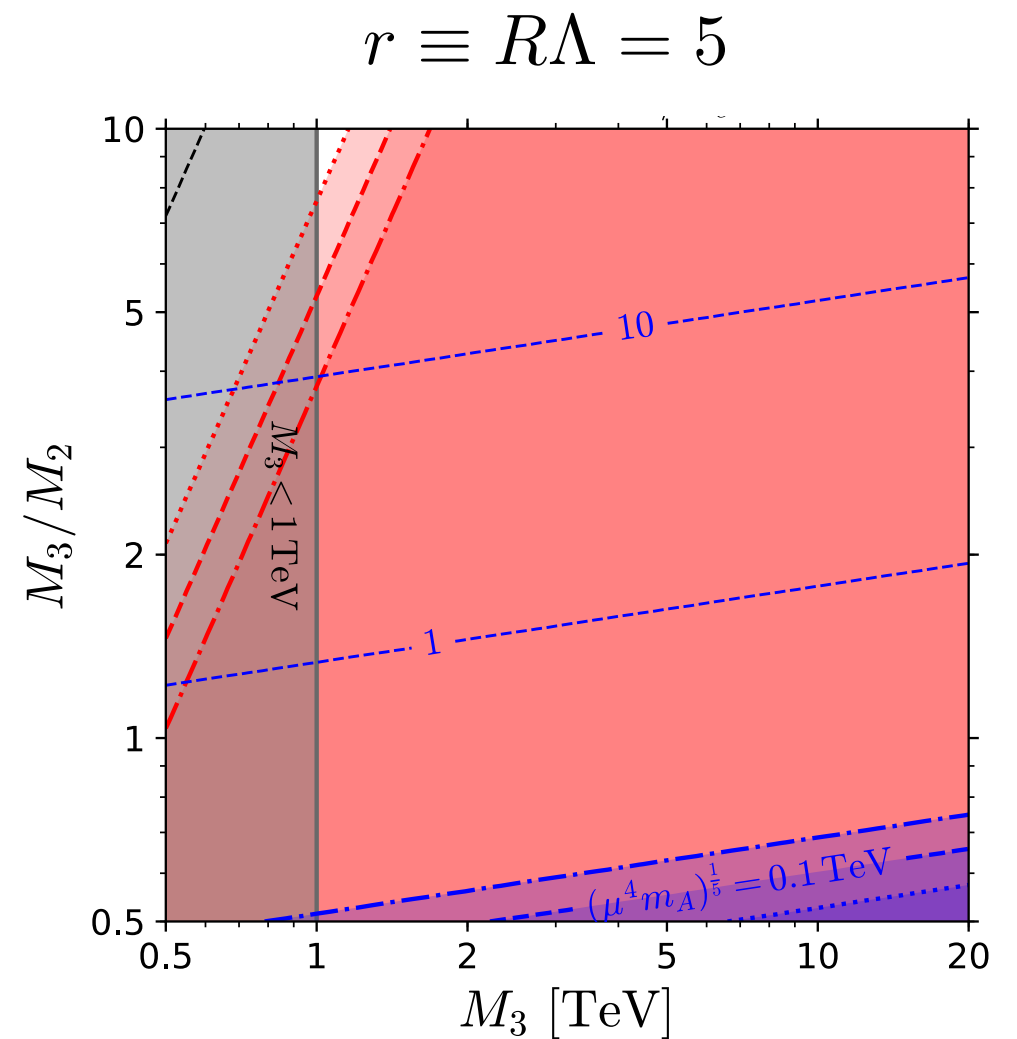


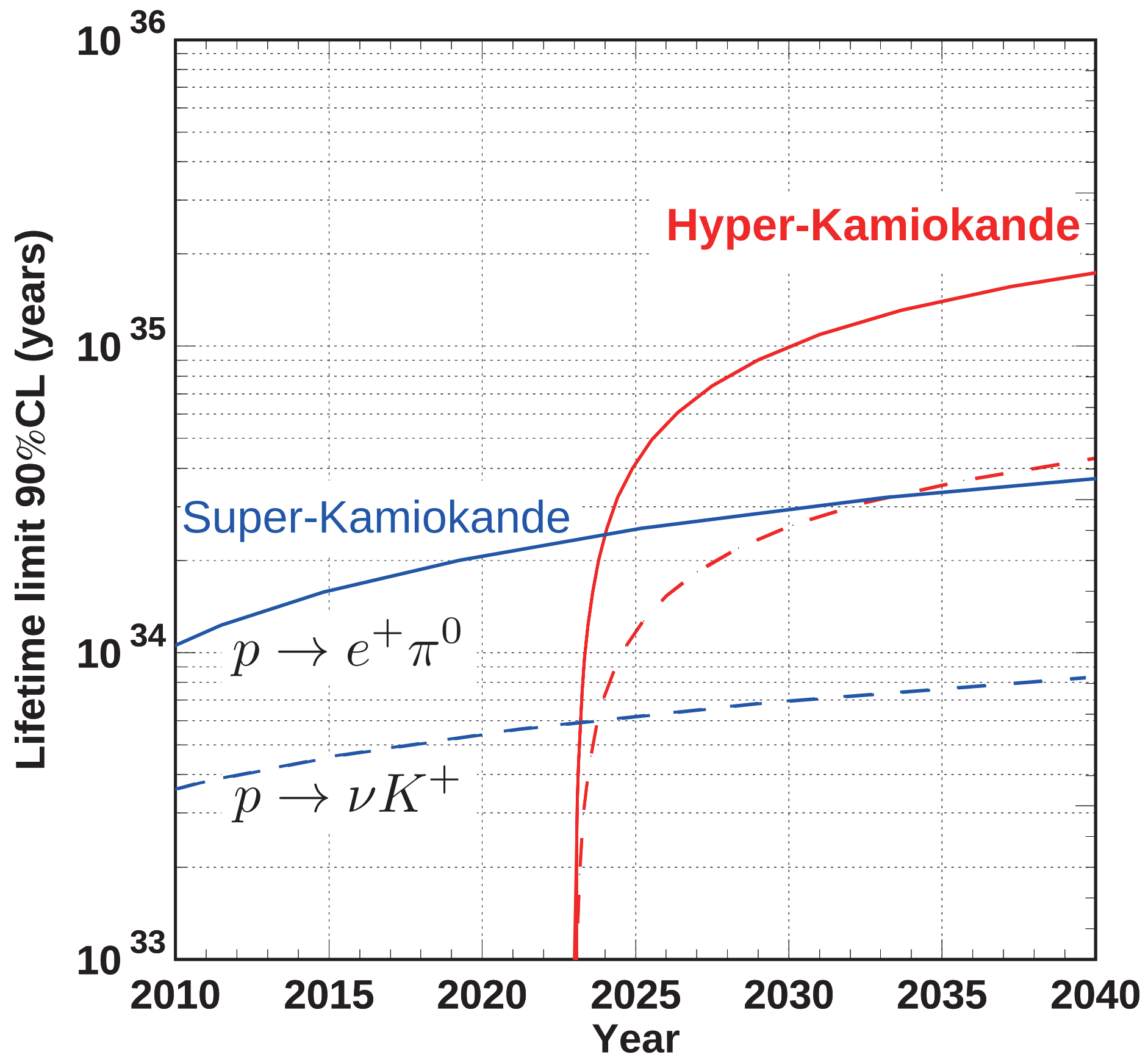
$M_3 > 1 - 2 \text{ TeV}$ from LHC

A SUSY plane with GCU



*) universal sfermion mass is assumed



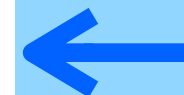


LHC

SK

**HL-LHC
FCC**

HK



$m_{\tilde{g}}$



Discussion

- The gauge coupling unification does not constrain the Bino mass, since it is singlet. The LHC signature varies depending on the Bino mass.

- The formula can easily be extended to non-minimal SUSY models.

η m_η	$\tilde{B}$ M_1	$\tilde{W}$ M_2	$\tilde{g}$ M_3	$\tilde{h}$ μ	A m_A	$\tilde{d}_R$ $m_{\tilde{d}_R}$	$\tilde{l}$ $m_{\tilde{l}}$	$\tilde{u}_R$ $m_{\tilde{u}_R}$	$\tilde{q}$ $m_{\tilde{q}}$	$\tilde{e}_R$ $m_{\tilde{e}_R}$
b_1^η	0	0	0	$\frac{2}{5}$	$\frac{1}{10}$	$\frac{1}{15}$	$\frac{1}{10}$	$\frac{4}{15}$	$\frac{1}{30}$	$\frac{1}{5}$
b_2^η	0	$\frac{4}{3}$	0	$\frac{2}{3}$	$\frac{1}{6}$	0	$\frac{1}{6}$	0	$\frac{1}{2}$	0
b_3^η	0	0	1	0	0	0	0	0	$\frac{1}{3}$	0

NMSSM => the same formula can be used

Table 1: Notation for the mass parameters and the values of β_a^η .

Non-singlet extension => straightforward to update the formulae for C_s, Ω_s, T_s

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A The condition of GCU

- The similar formula can be found for non-SUSY models.

In the MSSM, the gauge couplings at scale Q are given at 1-loop level by the RG equations:

$$\frac{2\pi}{\alpha_a(Q)} = \frac{2\pi}{\alpha_a(m_Z)} - b_a \ln\left(\frac{Q}{m_Z}\right) + s_a, \quad (\text{A.1})$$

Conclusions

- We have derived an analytic formula for the condition of GCU and the unified coupling at Λ in terms of the general SUSY and GUT spectra, including the dominant 2-loop effect and $\alpha_s(m_Z)$ uncertainty.

GCU condition

$$T_S = M_S^* \Omega_G \cap T_G = M_G^* \Omega_S$$

unified coupling

$$\alpha^{-1}(\Lambda) = \alpha_G^{*-1} + \frac{1}{2\pi} (C_S + C_G)$$

- Minimal SU(5):

The coloured Higgs mass is given as a function of low energy SUSY masses:
D=5 proton decay can be predicted by the SUSY spectrum.

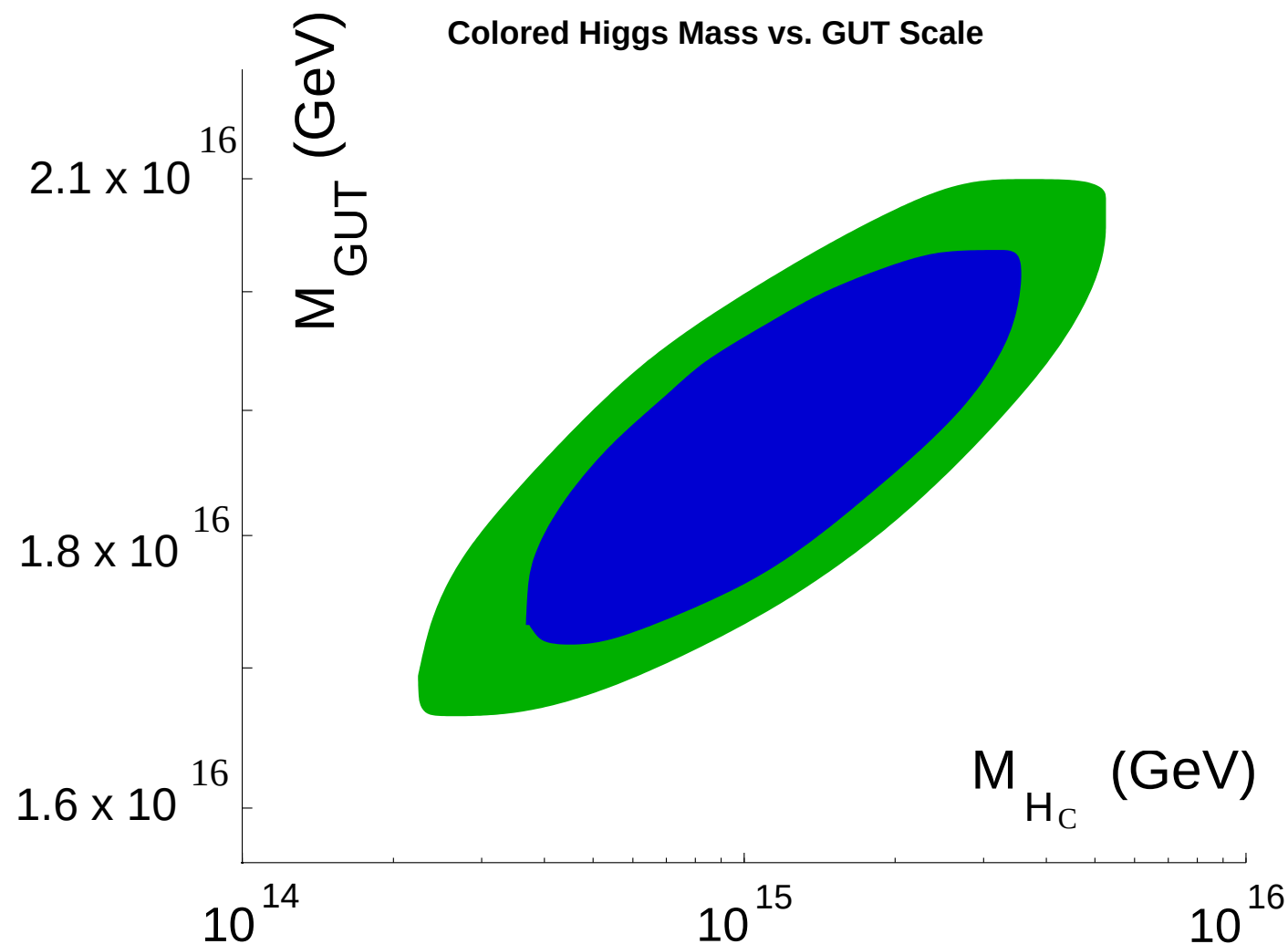
- Orbifold SUSY SU(5):

There is a non-trivial constraint on the SUSY spectrum. The X,Y boson mass is given as a function of low energy SUSY masses: D=6 proton decay can be predicted by the SUSY spectrum.

Backup

Is Minimal SU(5) excluded?

[Murayama, Pierce '01]



assumption:

$$M_3/M_2 = \alpha_3/\alpha_2$$

$$m_{\tilde{f}} = 1 \text{ TeV}$$

$$m_{\tilde{t}} \in (400, 800) \text{ GeV}$$

$$\tan \beta \in (1.8, 4)$$

$$M_2 \in (100, 400) \text{ GeV}$$

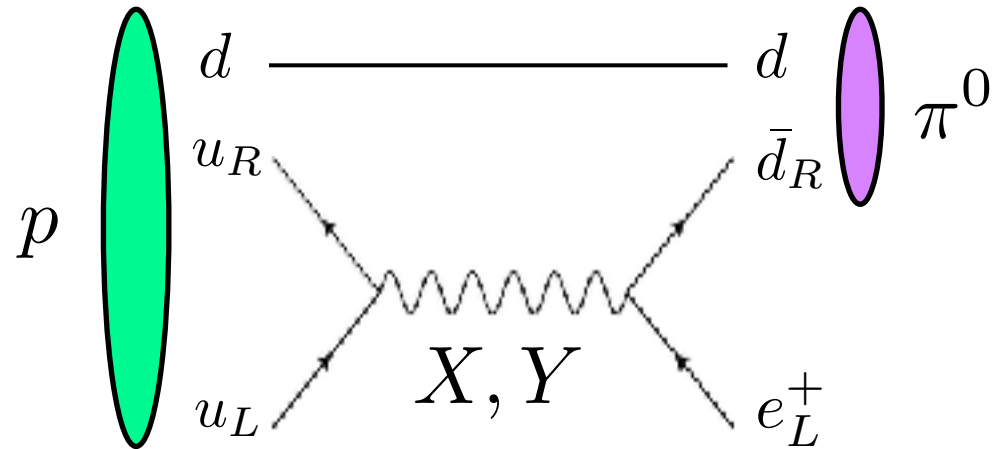
$$\mu \in (100, 1000) \text{ GeV}$$

FIG. 2. Plot showing 68% and 90% contours allowed by the renormalization group analysis for the color Higgs triplet mass, M_{HC} , and the GUT scale, $M_{GUT} \equiv (M_\Sigma M_V^2)^{1/3}$.

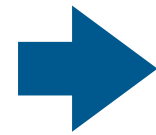
D=6

$$\frac{g^2}{\Lambda^2} (\mathbf{10}_i^* \mathbf{10}_i) (\mathbf{10}_j^* \mathbf{10}_j)$$

$$\frac{g^2}{\Lambda^2} (\mathbf{10}_i^* \mathbf{10}_i) (\bar{\mathbf{5}}_j^* \bar{\mathbf{5}}_j)$$



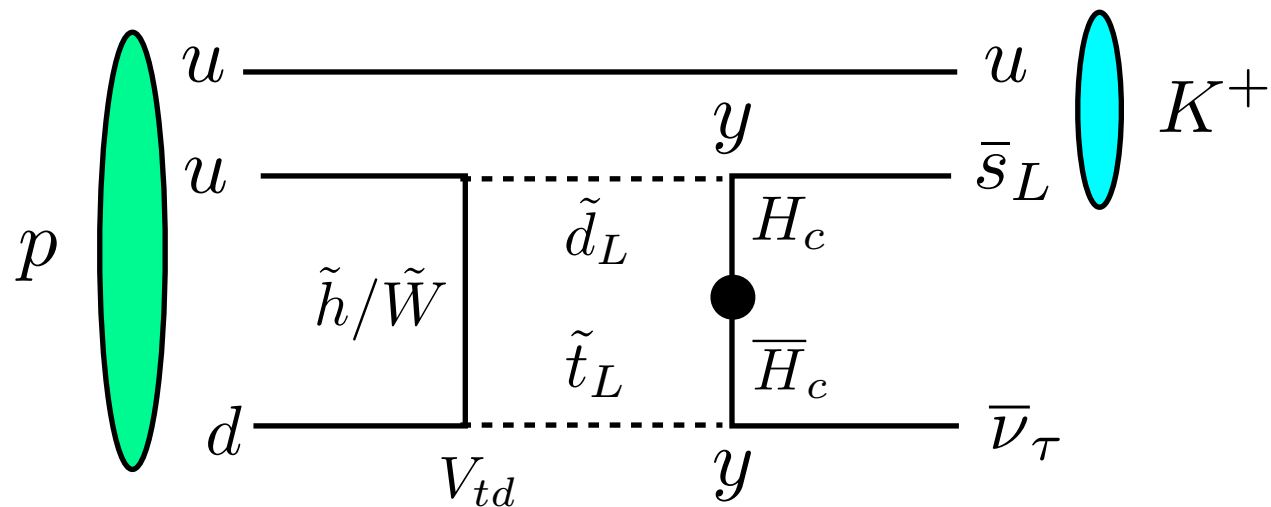
$$\tau_{p \rightarrow e^+ \pi^0} \sim \frac{1}{\alpha_G^2} \frac{M_{X,Y}^4}{m_p^5} > 1.7 \cdot 10^{34} \text{ years}$$



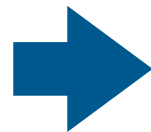
$$M_{X,Y} \gtrsim 6 \cdot 10^{15} \text{ GeV} \cdot \left(\frac{\alpha_G}{25.} \right)^2$$

D=5

$$\frac{y_i y_j}{\Lambda} \mathbf{10} \cdot \mathbf{10} \cdot \mathbf{10} \cdot \bar{\mathbf{5}}$$



$$\tau(p \rightarrow K^+ \bar{\nu}) > 5.9 \times 10^{33} \text{ yrs}$$



$$t_\beta^4 \left(\frac{10^3 \text{ GeV}}{M_{\text{SUSY}}} \right)^2 \left(\frac{10^{19} \text{ GeV}}{M_{H_c}} \right)^2 \lesssim 1$$

CMSSM

