Cosmology and Cluster Astrophysics with Cross-Correlations of HSC WL and Planck tSZ

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Contents

- **1. Introduction**
- 2. Method: model and covariance
- **3. Measurements and analysis of tSZ-WL cross-correlation with HSC and Planck**
- 4. Future prospects and summary

Cluster Cosmology and Astrophysics

<u>Cluster cosmology</u>

The abundance of clusters and mass profile of clusters are powerful probes into large-scale structures. However, there is a tension on cosmological parameters, e.g., **sigma8 tension**.

Non-thermal pressure

Some physical processes other than thermal pressure, e.g., **turbulence**, can also support the self-gravity of galaxy clusters. The mass under hydrostatic equilibrium (HSE) is parameterized as $M_{\rm HSE}$

$$\frac{M_{\rm HSE}}{M_{\rm true}} = 1 - b$$



Planck col. (2016)

Weak Gravitational Lensing

The **large-scale structures** induce weak gravitational lensing effect. We can probe into the matter distribution in an **unbiased** way.

The images of galaxies are distorted due to the foreground gravitational field, and the distortion can be detected by statistically analyzing many images.



Convergence field:

$$\begin{aligned} f(\theta) &= \frac{3}{2} \left(\frac{H_0}{c}\right)^2 \Omega_{\rm m} \\ &\times \int d\chi f(\chi_s, \chi) \delta(D_A(\chi)\theta, \chi) \end{aligned}$$

lensing kernel density



2D mass map from HSC



Cosmology with WL and tSZ



Auto 2pt Correlations



Power spectrum / 2pt correlation



Cosmology

Power spectrum / 2pt correlation







Auto 2pt Correlations



<u>tSZ</u>





Power spectrum / 2pt correlation



Cross 2pt Correlations

<u>WL</u>



-3.5

Cross-correlations!

Especially in the case of cross-correlation between high S/N and low S/N observables, the cross-correlation becomes more powerful!

X = WL, galaxies, CMB temp.

Y = **tSZ**, CMB pol., GW source



Cross-Correlations of tSZ-WL

♦ Measurements of tSZ-WL cross-correlations 3.5 × 10⁻⁹

The cross-correlation between thermal SZ effect and weak lensing has been measured from *Planck* and CFHTLenS (RCSLenS) data.

cf. van Waerbeke+ 2014, Hojjati+ 2016



separation [arcmin]



→We can constrain both of cosmological parameters (**cosmology**) and hydrostatic mass bias (**astrophysics**) with tSZ auto-power spectrum and tSZ-WL cross-correlations.

Cross-Correlations of tSZ-WL



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Theoretical prediction of cross spectra is based on <u>halo model</u>. All matter and gas is associated with halos.

$$\begin{split} C_{\ell}^{y\kappa} &= C_{\ell}^{y\kappa(1\mathrm{h})} + C_{\ell}^{y\kappa(2\mathrm{h})} \\ C_{\ell}^{y\kappa(1\mathrm{h})} &= \int dz \, \frac{d^2 V}{dz d\Omega} \int dM \frac{dn}{dM} y_{\ell}(M, z) \\ \kappa_{\ell}^{y\kappa(2\mathrm{h})} &= \int dz \, \frac{d^2 V}{dz d\Omega} P_{\mathrm{m}}(k = \ell / D_A, z) \\ &\times \int dM_1 dM_2 \frac{dn}{dM_1} b(M_1, z) y_{\ell}(M_1, z) \frac{dn}{dM_2} b(M_2, z) \\ \kappa_{\ell}(M_2, z) \end{split}$$

$$\xi^{y\kappa}(\theta) = \int \frac{\ell d\ell}{2\pi} C_{\ell}^{y\kappa} J_0(\ell\theta)$$

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 $C_{\ell}^{\gamma\kappa} = C_{\ell}^{\gamma\kappa(1h)} + C_{\ell}^{\gamma\kappa(2h)} \quad \text{well calibrated with} \\ \frac{N + body simulations}{N + body simulations}$ $C_{\ell}^{\gamma\kappa(1h)} = \int dz \frac{d^2V}{dzd\Omega} \int dM \frac{dn}{dM} y_{\ell}(M, z) \kappa_{\ell}(M, z)$ $C_{\ell}^{\gamma\kappa(2h)} = \int dz \frac{d^2V}{dzd\Omega} P_{\rm m}(k = \ell/D_A, z)$ $\times \int dM_1 dM_2 \frac{dn}{dM_1} b(M_1, z) y_{\ell}(M_1, z) \frac{dn}{dM_2} b(M_2, z) \kappa_{\ell}(M_2, z)$

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$$\begin{split} C_{\ell}^{y\kappa} &= C_{\ell}^{y\kappa(1\mathrm{h})} + C_{\ell}^{y\kappa(2\mathrm{h})} & \text{well calibrated with} \\ & \text{N-body simulations} \\ C_{\ell}^{y\kappa(1\mathrm{h})} &= \int dz \, \frac{d^2 V}{dz d\Omega} \int dM \frac{dn}{dM} y_{\ell}(M, z) \kappa_{\ell}(M, z) & \text{Projection of} \\ & \text{NFW profile} \\ C_{\ell}^{y\kappa(2\mathrm{h})} &= \int dz \, \frac{d^2 V}{dz d\Omega} P_{\mathrm{m}}(k = \ell/D_A, z) & \text{Compton-y} \\ & \times \int dM_1 dM_2 \frac{dn}{dM_1} b(M_1, z) y_{\ell}(M_1, z) \frac{dn}{dM_2} b(M_2, z) \kappa_{\ell}(M_2, z) \end{split}$$

$$\xi^{y\kappa}(\theta) = \int \frac{\ell d\ell}{2\pi} C_{\ell}^{y\kappa} J_0(\ell\theta)$$

Pressure Profile

- Universal pressure profile: We can fit the pressure profile with X-ray or SZ observation. The GNFW profile (Nagai+, 2007) is used.
- The fitted pressure profile using <u>62 nearby clusters</u> observed by *Planck*.
- We rescale the mass and radius with hydrostatic bias:

$$M_{500}^{\rm HSE} = M_{500}(1 - b_{\rm HSE})$$

$$P_{e}(r) = P_{500} \left(\frac{M_{500}^{\text{HSE}}}{3 \times 10^{14} h_{70}^{-1} \,\text{M}_{\odot}} \right)^{0.12} \times \frac{P_{0}}{(c_{500} x)^{\gamma} [1 + (c_{500} x)^{\alpha}]^{(\beta - \gamma)/\alpha}} \times \frac{R/R_{500}}{x = R/R_{500}}$$



Mock WL and tSZ Maps

<u>tSZ</u> + All-sky simulations GAMA15H GAMA15H Takahashi+ (2017) Shirasaki+ (2015) RA (J2000) RA (J2000) WIDE12H WIDE12H $\mathsf{Dec}_{\circ 0+} (\mathsf{J2000}) \\ ^{\circ 0+} _{\circ 0+}$ Dec (J2000 12^{h} $11^{h}52^{m}$ $11^{h}44^{n}$ $12^{h}08^{m}$ 12^{h} $11^{h}52^{m}$ $11^{h}44^{m}$ RA (J2000) RA (J2000) GAMA09H GAMA09H Dec (J2000) +0, +0, +0, (J2000) 09^{h} $08^{h}40^{n}$ $00^{h}20^{n}$ 08h40n 00^{h} RA (J2000) RA (J2000) VVDS VVDS Dec (J2000) (J2000 $22^{h}32^{m}$ $22^{h}24^{m}$ $22^{h}16^{m}$ $22^{h}08^{m}$ $22^{h}32^{m}$ $22^{h}24^{m}$ $22^{h}16^{m}$ $22^{h}08^{n}$ $22^{h}40^{m}$ $22^{h}40^{m}$ RA (J2000) RA (J2000) ХММ XMM 0000Z() -4 (J2000) x108 $02^{h}20^{m}$ $02^{h}20^{m}$ 02^{1} RA (J2000) RA (J2000) The mock measurements **HECTOMAP HECTOMAP** are used to evaluate the Dec $16^{h}24^{m}$ $16^{h}32^{m}$ $16^{h}16^{n}$ $16^{\rm h}$ covariance matrix. RA (RA (J2000) KO+ (2019)

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Subaru Hyper Suprime-Cam









Measurements of Cross-Correlation



tSZ-WL cross-correlatio

Measurements of Cross-Correlation



KO+ (2019)

Analysis

Data sets

tSZ auto-power spectrum only
tSZ-WL cross-correlations only
Joint of both

• For tSZ auto-power spectrum Bolliet+ (2018) CIB, IR point sources, radio sources, and correlated noise For tSZ-WL cross-correlations Shirasaki (2019) Radio sources (flat-spectrum radio quasars, BL Lac)

Priors on cosmological parameters

combination of low-z LSS probes
(HSC cosmic shear + JLA SN Ia + BOSS DR12 BAO/RSD)
Planck 2018 TT,TE,EE+lowE+lensing



Constraints on Parameters



KO+ (2019)

Constraints on Amplitude and Mass Bias

uctuation

Amplitude of matter



Constraints on Amplitude and Mass Bias



Constraints on Mass Bias



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Future Prospects

Cross-correlations with HSC WL and ACT tSZ

Atacama Cosmology Telescope (ACT) can measure the CMB temperature fluctuation with high resolution (beam size ~ 1 arcmin). By cross-correlating the data, we can probe into the fine structure of pressure profile.



Madhavacheril+ (2019)

Tomography technique

Similarly to cosmic shear analysis, the whole galaxy samples can be divided according to its photo-z. Then, we measure crosscorrelations for the subsamples. We can investigate the redshift evolution of hydrostatic bias.

In this aspect, HSC has advantages of high source number density.

Summary

- Weak lensing and the thermal Sunyaev-Zel'dovich effect are promising probes into the large-scale structure and thermodynamical properties of intra-cluster medium.
- Cross-correlation is a powerful statistic with high S/N significance provides additional information useful for breaking degeneracy.
- Halo model calculation and N-body simulations are used to predict the signal and estimate the covariance matrix. This study presents the first attempt to estimate covariance matrix from realistic mock simulations.
- HSC is the unique WL survey which can probe into the large-scale structures and cluster astrophysics at high redshifts, and the evolution of them by tomography.