

Progenitors of low-mass binary black hole mergers

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Very common!



Fractions estimated for all stars with initial masses above 15 M $_{\odot}$ (Sana et al. 2012)

Indeed very common!



from Sana et al. 2012

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What about this?



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Dynamical interaction channel

Key aspect: *dense stellar cluster*. Several interactions before formation



from Rodriguez et al. 2016

Isolated binary evolutionary channel (I)

Key aspect: stars need to evolve *chemical homogeneously*their entire life



Isolated binary evolutionary channel (II)

Key aspect: need a *common-envelope phase* (CE)



from Garcia et al. in prep



LIGO-Virgo | Frank Elavsky | Northwestern



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GW151226

GW170608

- Discovery date: December 26, 2015
- Hanford SNR of 10.5, Livingston SNR of 7.9
- Differences arising because of sensitivies



• Difficult signal visibility

- Discovery date: June 8, 2017
- Livingston SNR of 9, Hanford data not reliable at first



• Different sensitivies and difficult signal visibility again!

GW151226

GW170608

- $\bullet \,\, M_{
 m BH,1} = 14.2^{+8.3}_{-3.7}\,M_{\odot}$
- $\bullet \,\, M_{
 m BH,2} = 7.5^{+2.3}_{-2.3}\,M_{\odot}$
- $\bullet \,\, M_{
 m chirp} = 8.9^{+0.3}_{-0.3}\,M_{\odot}$
- $\bullet \,\, M_{
 m merged-BH} = 20.8^{+6.1}_{-1.7}\,M_{\odot}$
- redshift: $z = 0.09^{+0.03}_{-0.04}$



- $M_{
 m BH,1} = 12^{+7}_{-2}\,M_{\odot}$
- $\bullet \,\, M_{
 m BH,2} = 7^{+2}_{-2} \, M_{\odot}$
- $\bullet ~~M_{
 m chirp} = 7.9^{+0.2}_{-0.2} \, M_{\odot}$
- $M_{
 m merged-BH} = 18^{+4.8}_{-0.9}\,M_{\odot}$
- redshift: $z = 0.07^{+0.03}_{-0.03}$



Aims

1) Study the progenitor properties for these two GW events in the isolated binary evolutionary scenario going through a common-envelope phase



2) Obtain merger rates for them in 01, 02 and expected ones in 03

Methods:

First part

- Follow the complete evolution of the binary: from two nondegenerate stars up to the formation of two black holes
- Using a detailed stellar evolutionary code, publicly available, *MESA*, modified to include the common-envelope phase

Methods:

First part

Free parameters in simulations

Metallicity of the population, (Z)

Accretion efficiency during the stable mass-transfer phase, (ϵ)

Efficiency for the removal of a star envelope during a common-envelope phase, ($lpha_{ ext{CE}}$)

3D grids of initial masses and binary separations were created for each combination of the above parameters. Total number of simulations above 50 000!



Initial binary parameters for $lpha_{\rm CE}$ = 2

- Higher metallicities requires increasingly massive stars: directly related with winds
- In the low mass-accretion regime, only low metallicity binaries are progenitors. High metallicity ones produce low chirp masses.
- For $\epsilon > 0.2$ binaries with similar initial masses are also progenitors.



Map of initial parameters for $\alpha_{\rm CE}$ = 2 with BBH properties





Map of initial parameters for $\alpha_{\rm CE}$ = 2 with BBH properties





Initial binary parameters for



• Initial mass ratios are closer to unity.

- No solutions were found for mass-accretion efficiencies that are below ϵ < 0.2
- Also, binaries with separations lower than 60 merge during a CE phase as a consequence of having a lower efficiency for the CE ejection
 - Merger times are also lowered.

Map of initial parameters for $\alpha_{\rm CE}$ = 1 with BBH properties

Z = 0.004

Z = 0.001

1.0 0.2

 $q_{\rm BBH}$

0.8

 $\varepsilon = 0.4$

0.4

0.6

0.8

1.0



60

50

40

50

40

30

20

M_{i,2} [M_☉]

GW151226

GW170608

Map of initial parameters for $\alpha_{\rm CE}$ = 1 with BBH properties





Methods:

Second part

To estimate expected properties, detailed stellar models were rescaled by empirical **initial mass functions** (IMF) for the primary and secondary stars and by an **initial separation distribution** from the observed binary orbital period distribution

Weighted initial binary parameters

 $lpha_{
m CE}=2$



$lpha_{ m CE}=1$









Methods:

Second part

$$\mathcal{R}(Z, z(t)) = \mathcal{N}_{\text{corr}} \int_0^{t(z)} \int_{M_{i,1}} \int_{M_{i,2}} \int_{a_i} \int_0^{t(z)} \frac{\mathrm{d}N}{\mathrm{d}M_{i,1} \,\mathrm{d}M_{i,2} \,\mathrm{d}a_i \,\mathrm{d}t_m}$$
$$\widehat{\text{SFR}}(t'; Z)\delta\left[t(z) - (t_m + t')\right] \mathrm{d}t_m \mathrm{d}a_i \,\mathrm{d}M_{i,2} \,\mathrm{d}M_{i,1} \,\mathrm{d}t'$$



Merger rate density history

- The expected local merger rate densities are all larger for the highest value of CE removal efficiency. Related to 'size' of parameter space
- For high metallicties, rates decay with redshift because of the chemical evolution
- For the low CE efficiency, rates are largely dominated by low metallicities



$$R_{\rm D} = \frac{4\pi}{3} D_{\rm h}^3 \langle w^3 \rangle \langle (\mathcal{M}_{\rm c} \ / \ 1.2 \ M_\odot)^{15/6} \rangle \mathcal{R}(z=0)$$

Detectable merger rate at zero redshift

- The highest rate is obtained at ϵ = 0.4 for both CE efficiencies
- However, expected rates obtained are within a factor of 2, thus we are not able to distinguish a preferred value for the mass-accretion efficiency

Conclusions

With current and future campaings of observing GW, more we will now about their progenitors

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 Giving the rising power of computers, having a large grid of detailed binary calculations is possible. Even to calculate complete populations of progenitors

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Several uncertainties are still present in nowadays calculations, so estimates like rates can vary by order of magnitude when changing input physical parameters

THANK YOU!