

# Stochastic Background of Induced Gravitational Waves: Recent Developments

Shi PI 皮石

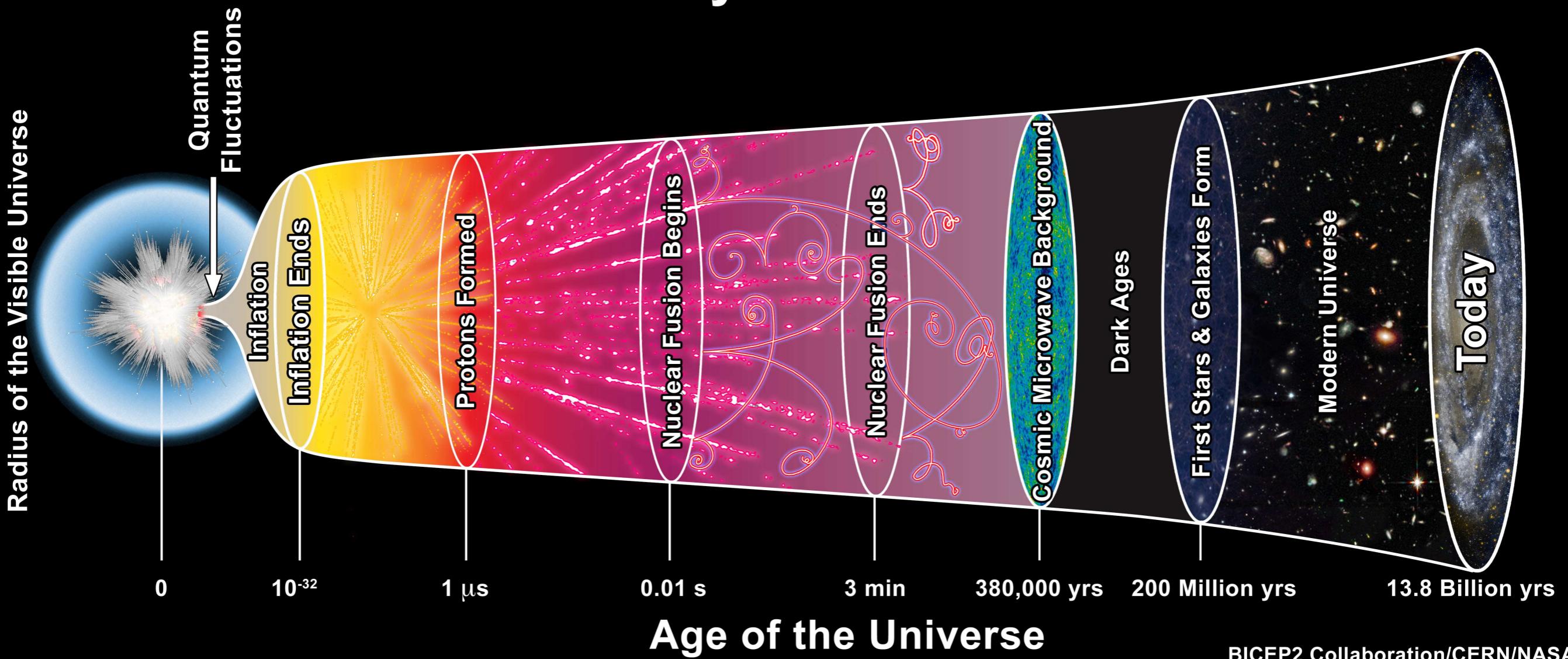
Kavli IPMU, the University of Tokyo → Institute for Theoretical Physics, CAS  
東京大学カブリ数物連携宇宙研究機構→中国科学院理论物理研究所

Kavli IPMU postdoc colloquium  
Zoom@Hubei China, 2020.9.29

# Content

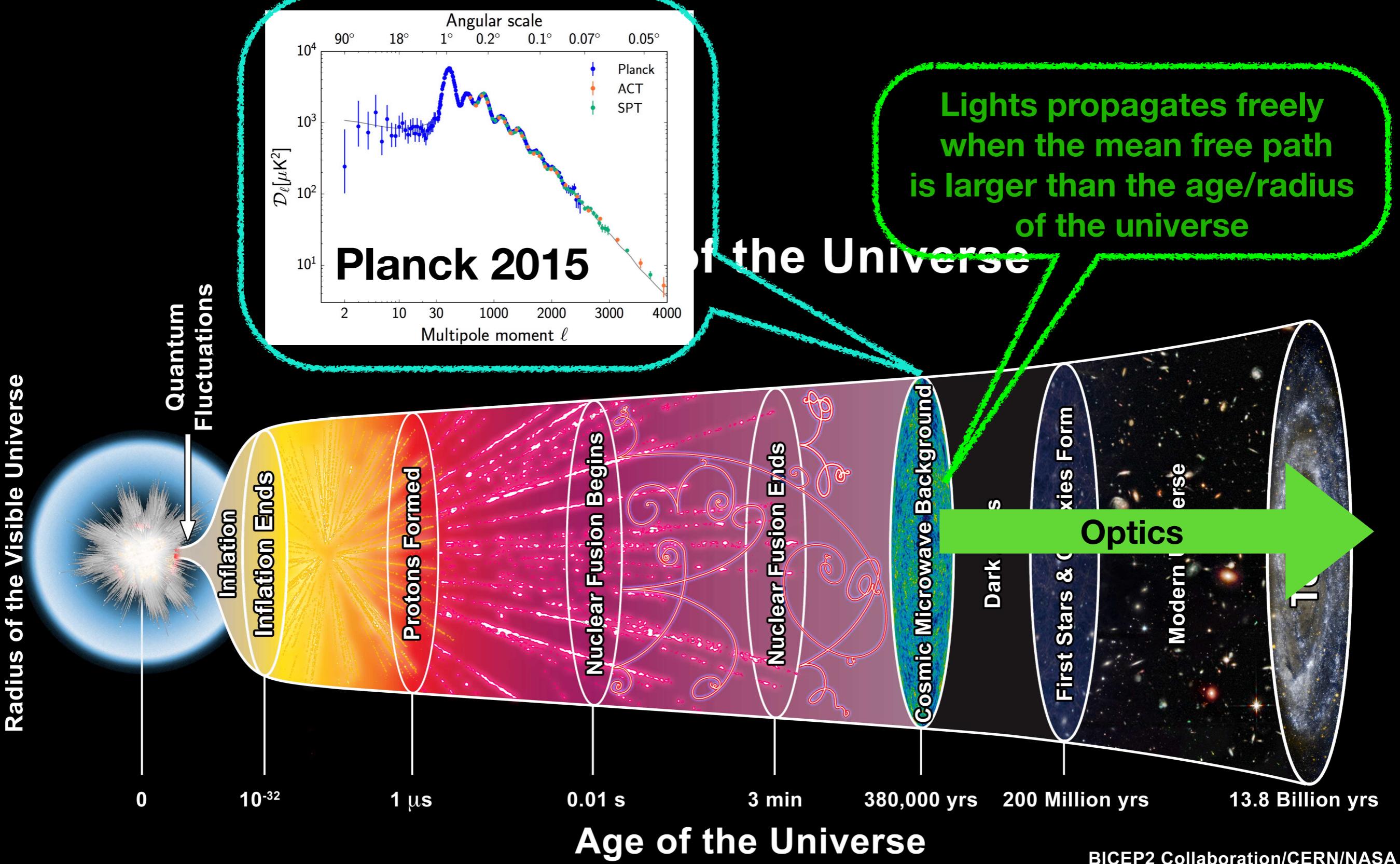
- Introduction
- Primordial perturbation, PBHs, and induced GWs
- NANOGrav hints on planet-mass PBHs
- Summary

# History of the Universe



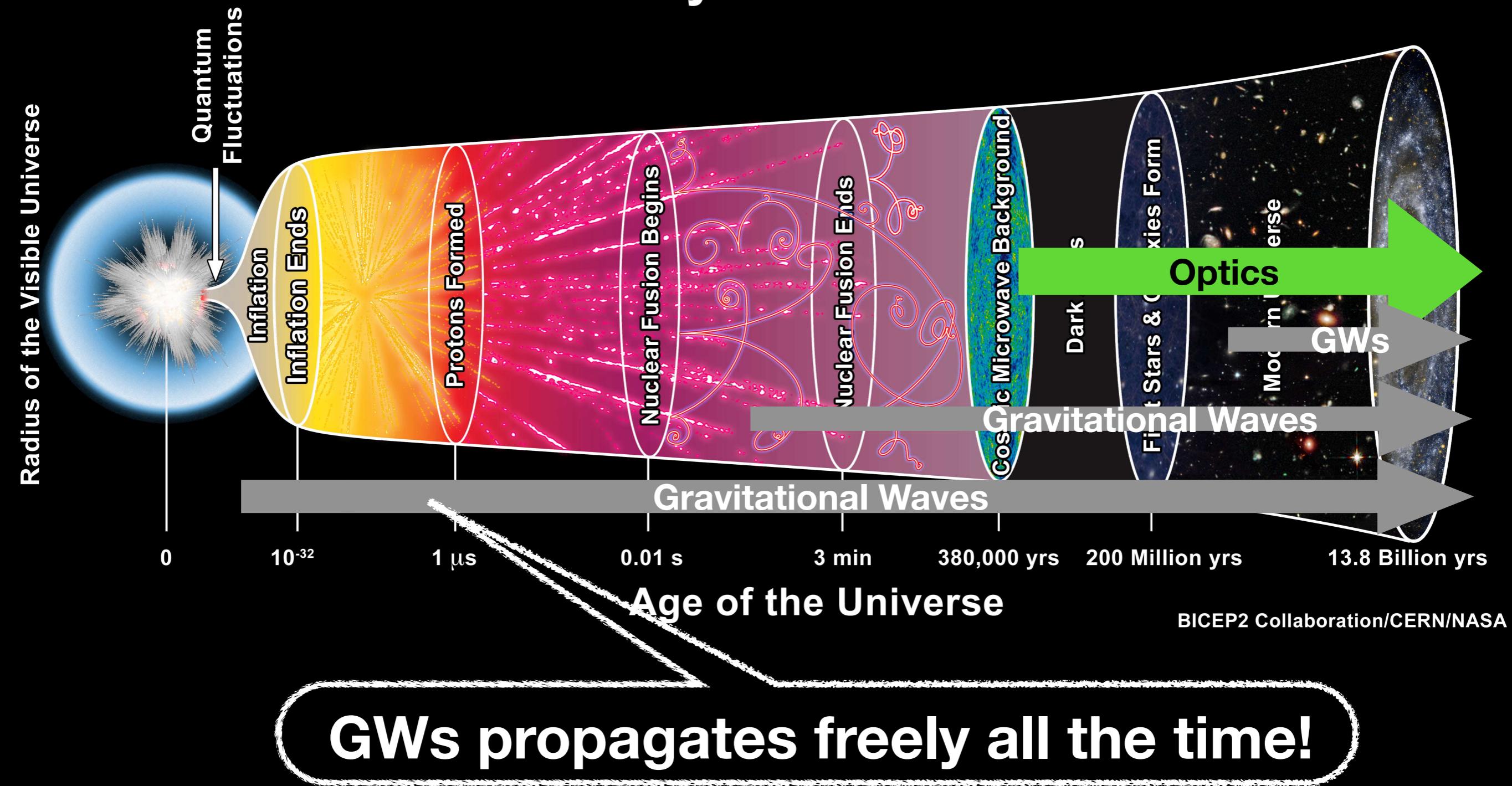
BICEP2 Collaboration/CERN/NASA

Radius of the Visible Universe

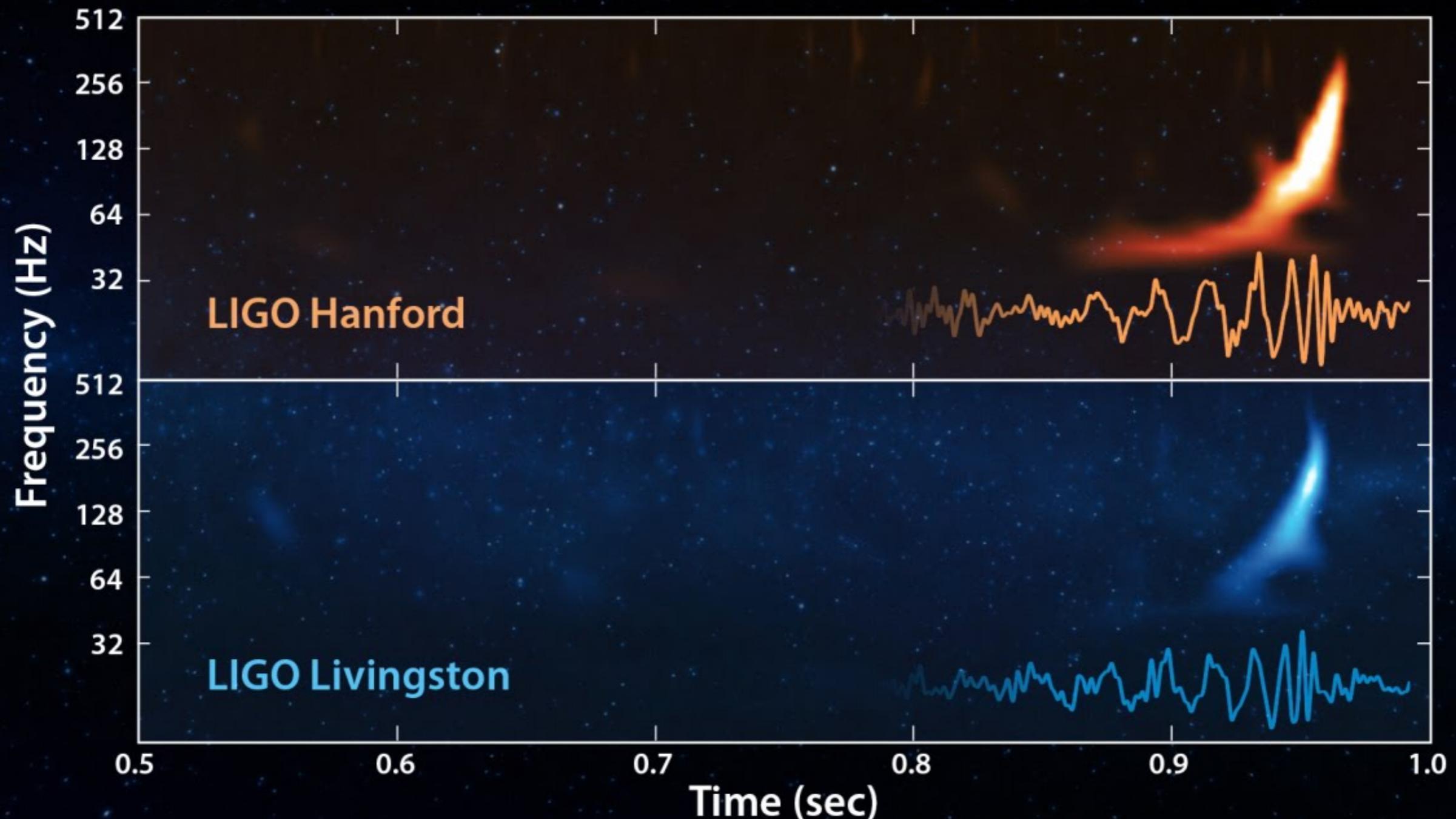


BICEP2 Collaboration/CERN/NASA

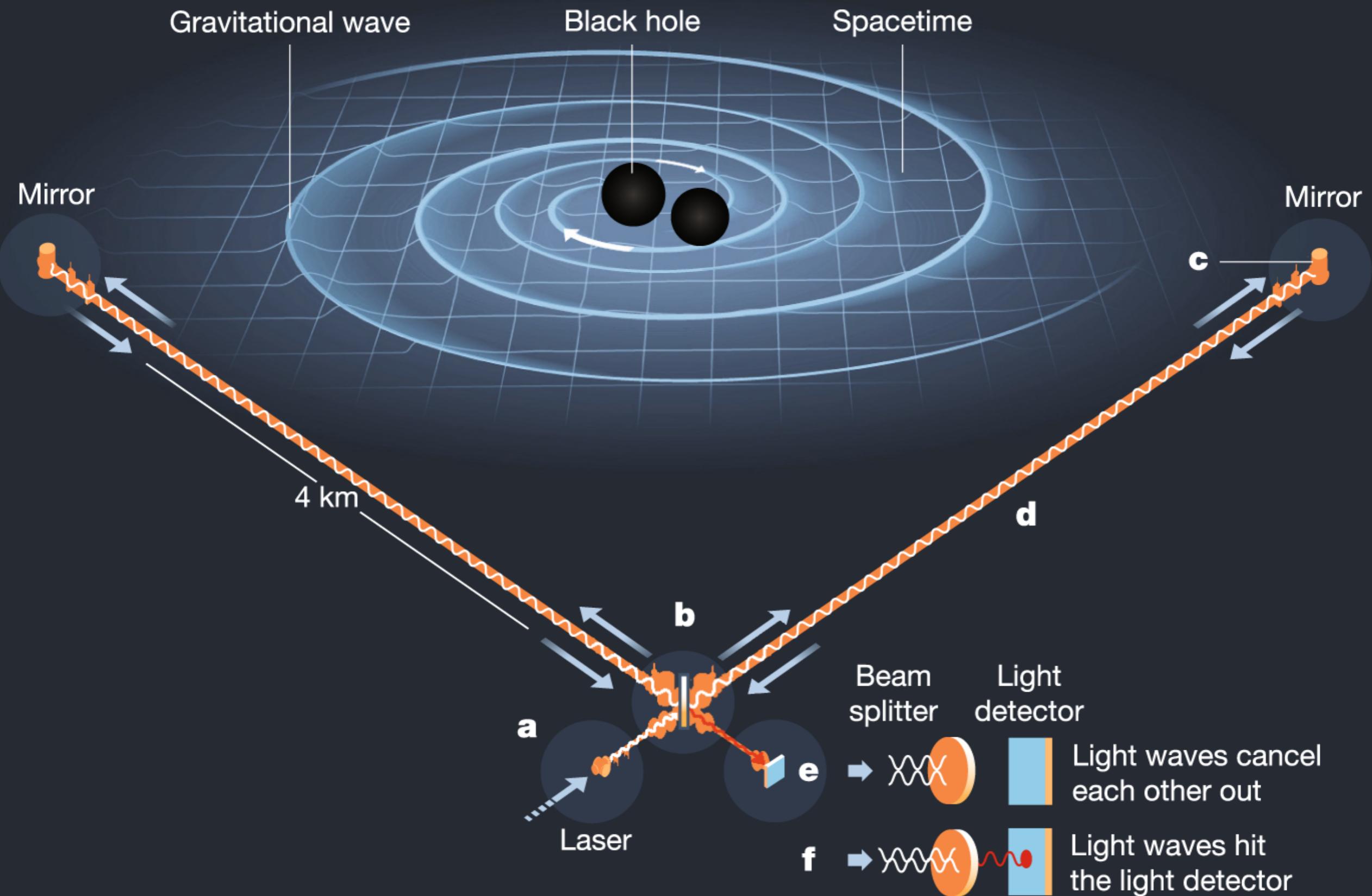
# History of the Universe



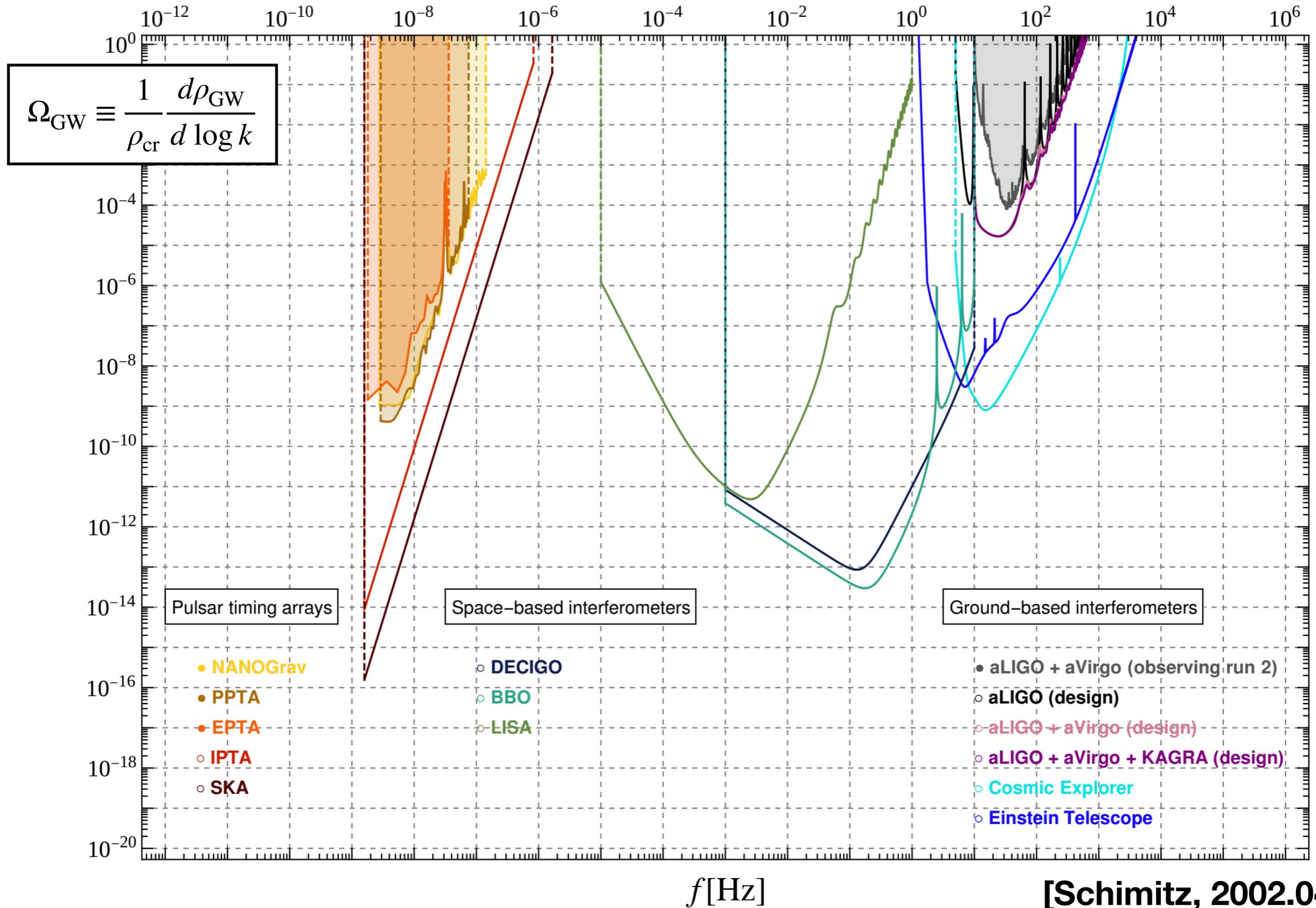
# LIGO detection of GWs



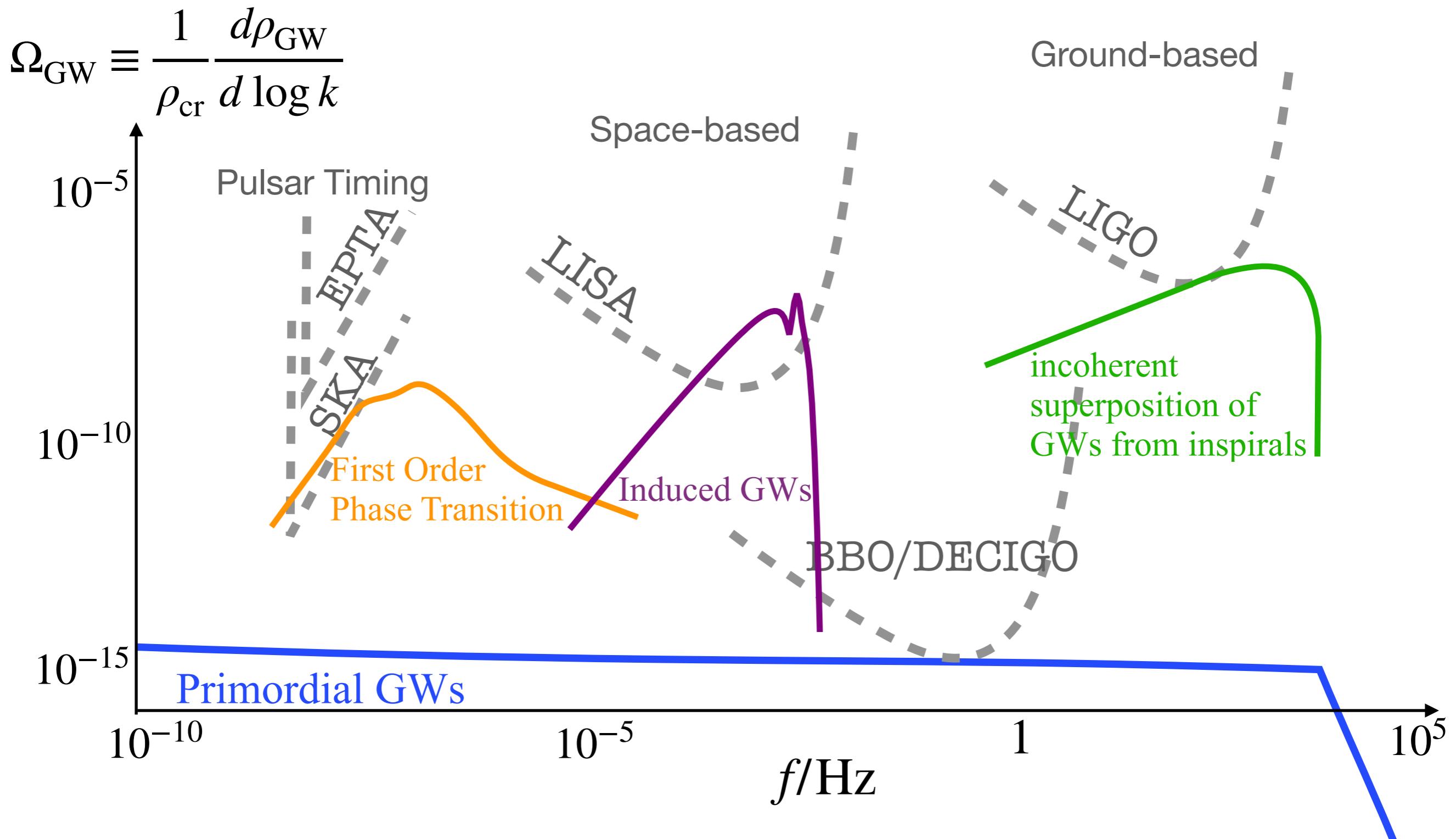
# Detection of GWs

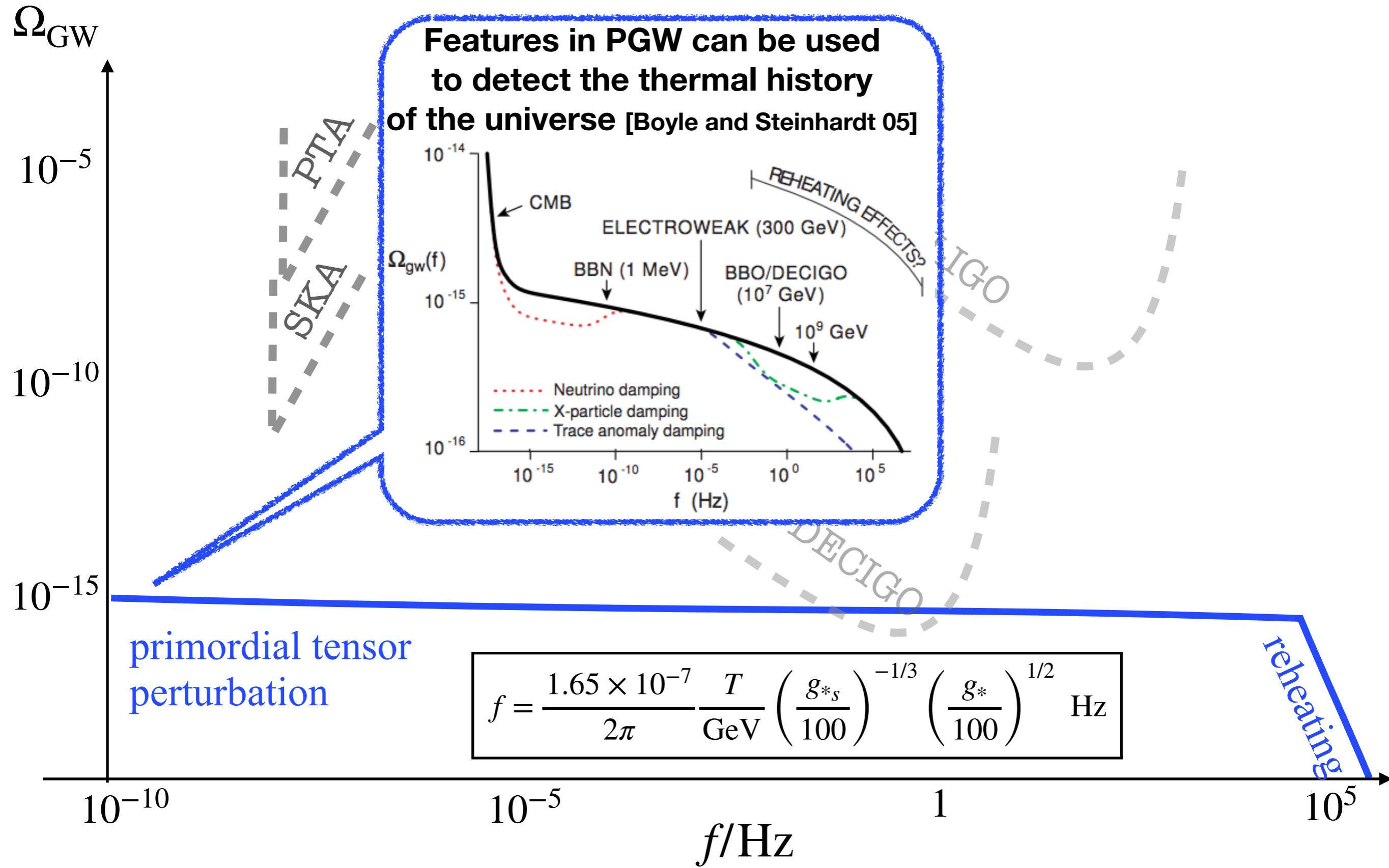


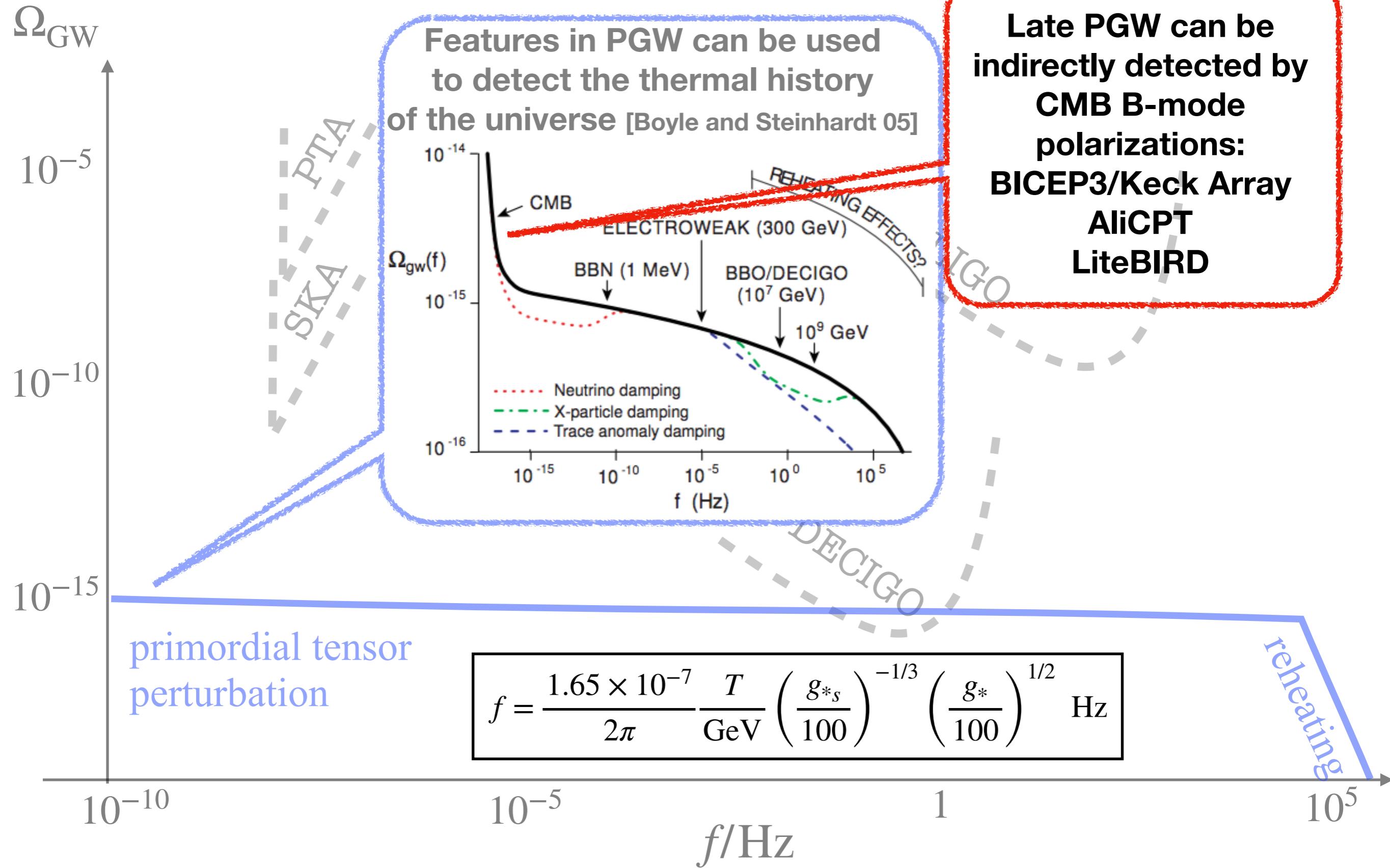
# Future GW detectors

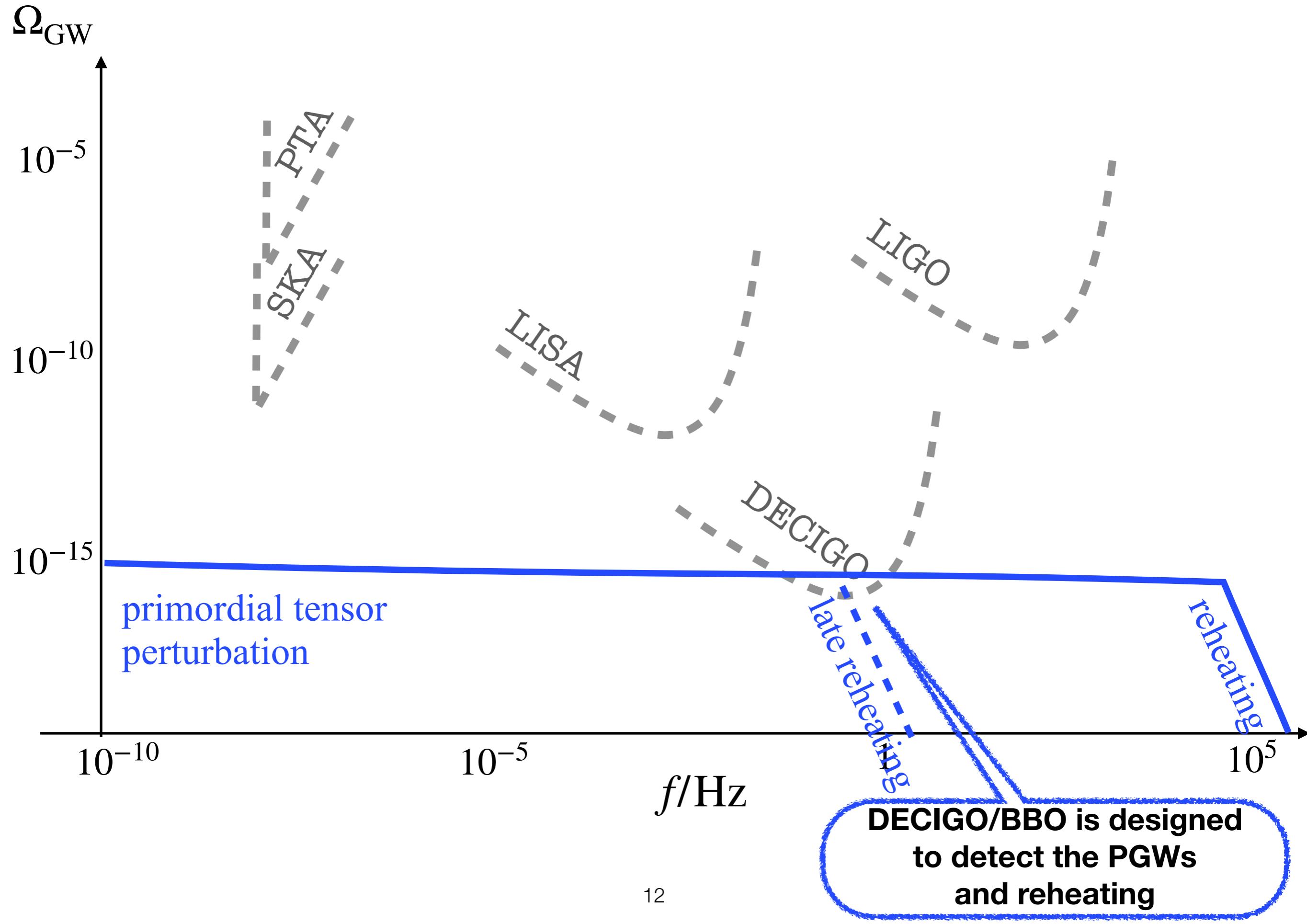


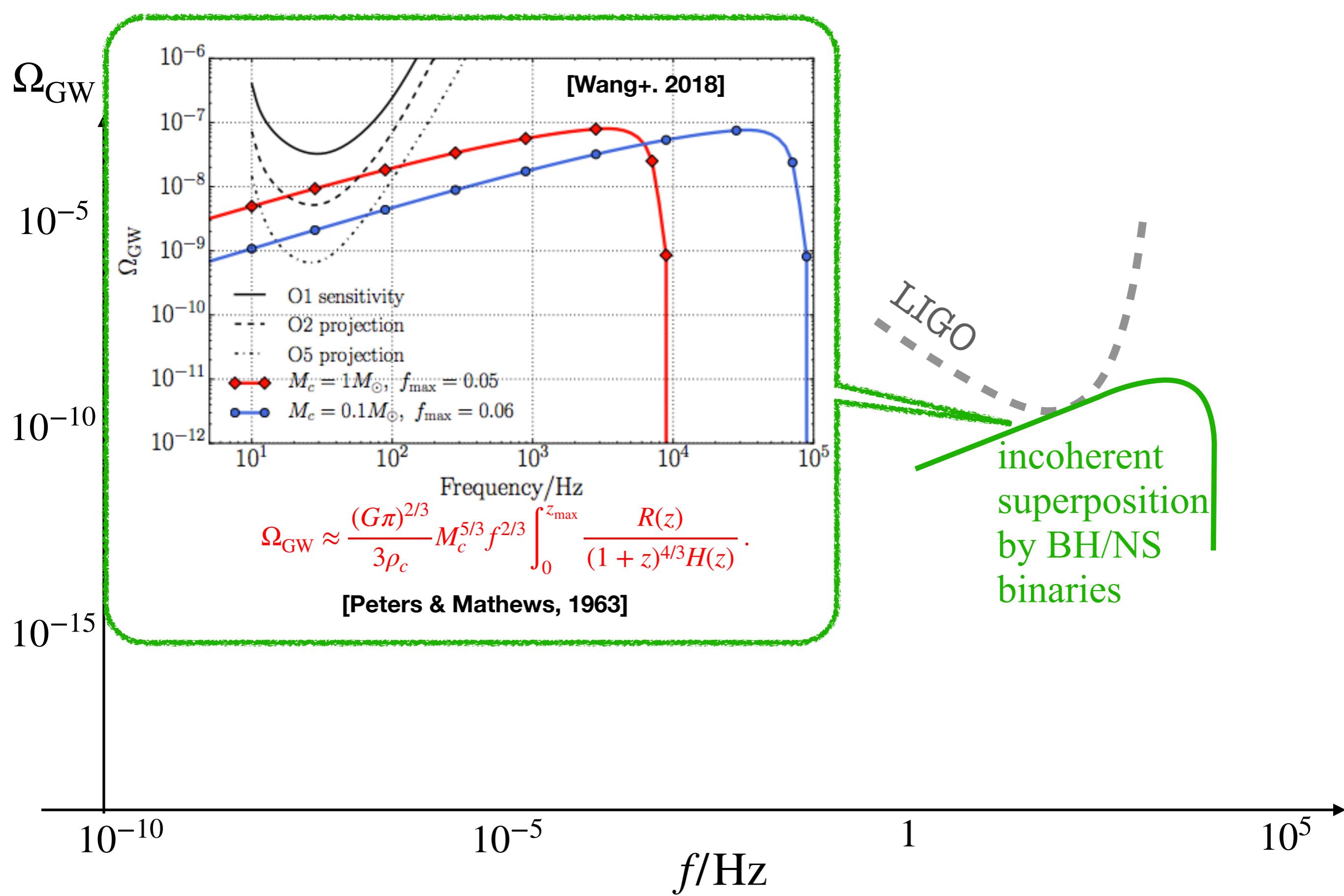
# Possible SGWB Sources

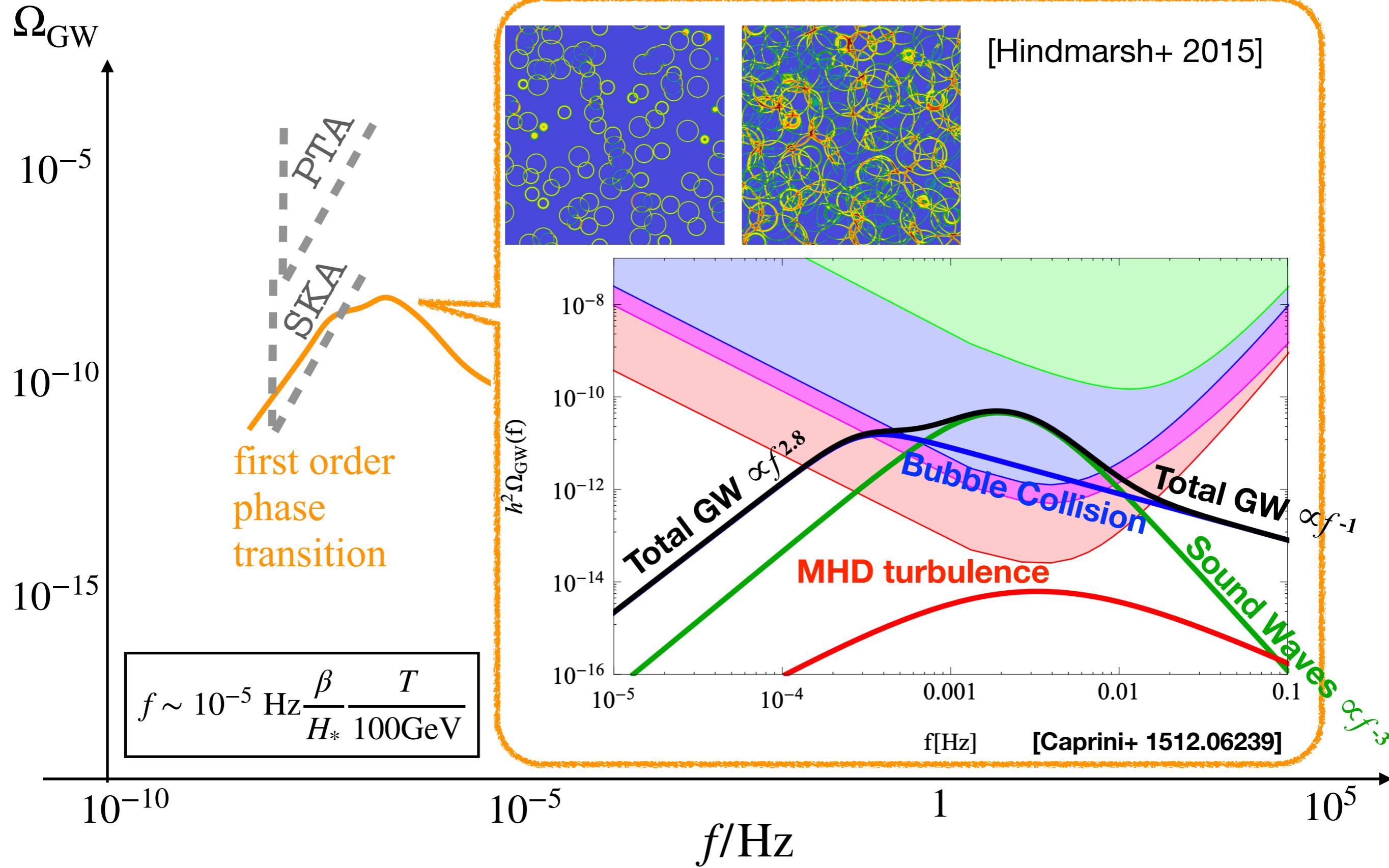


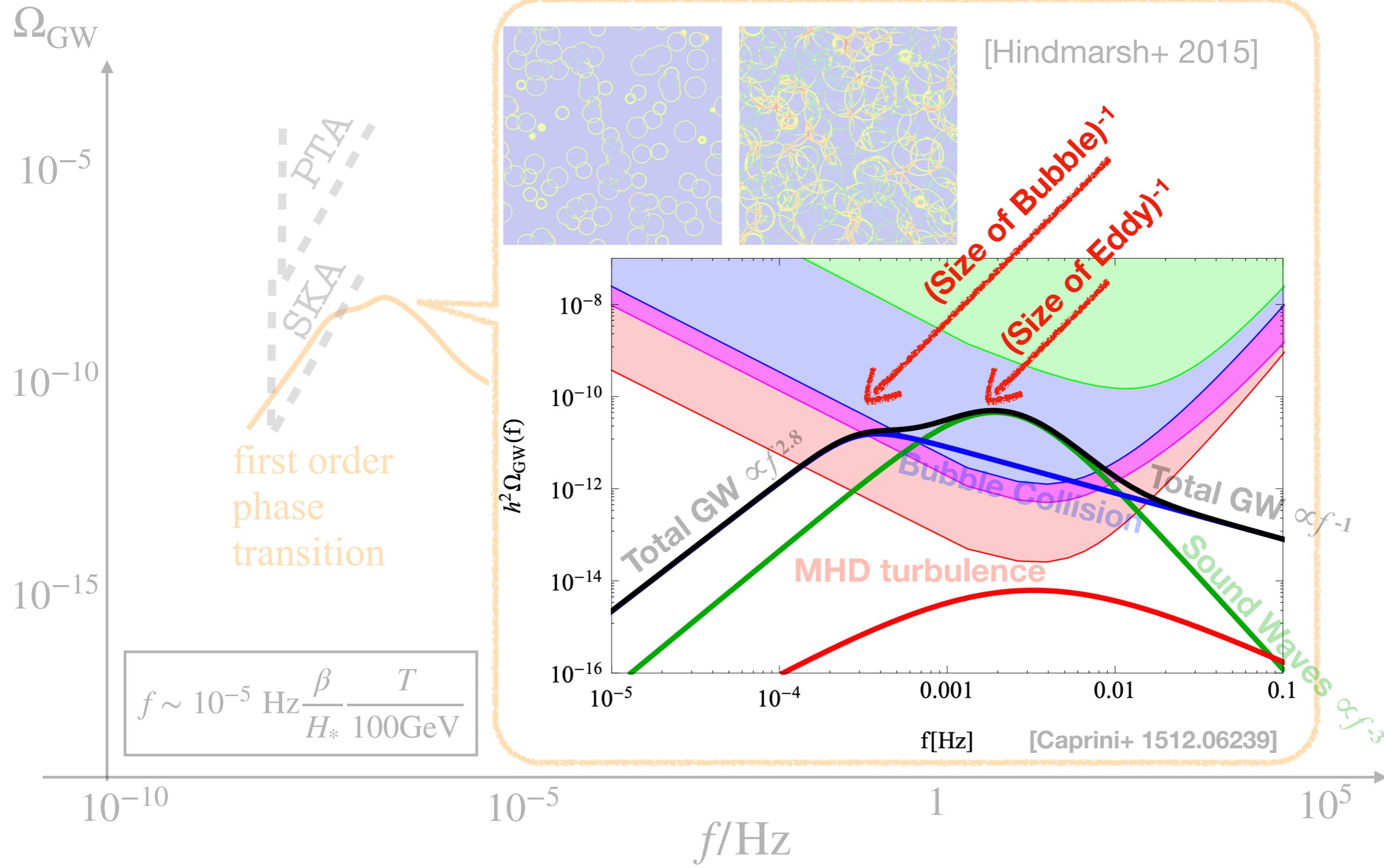


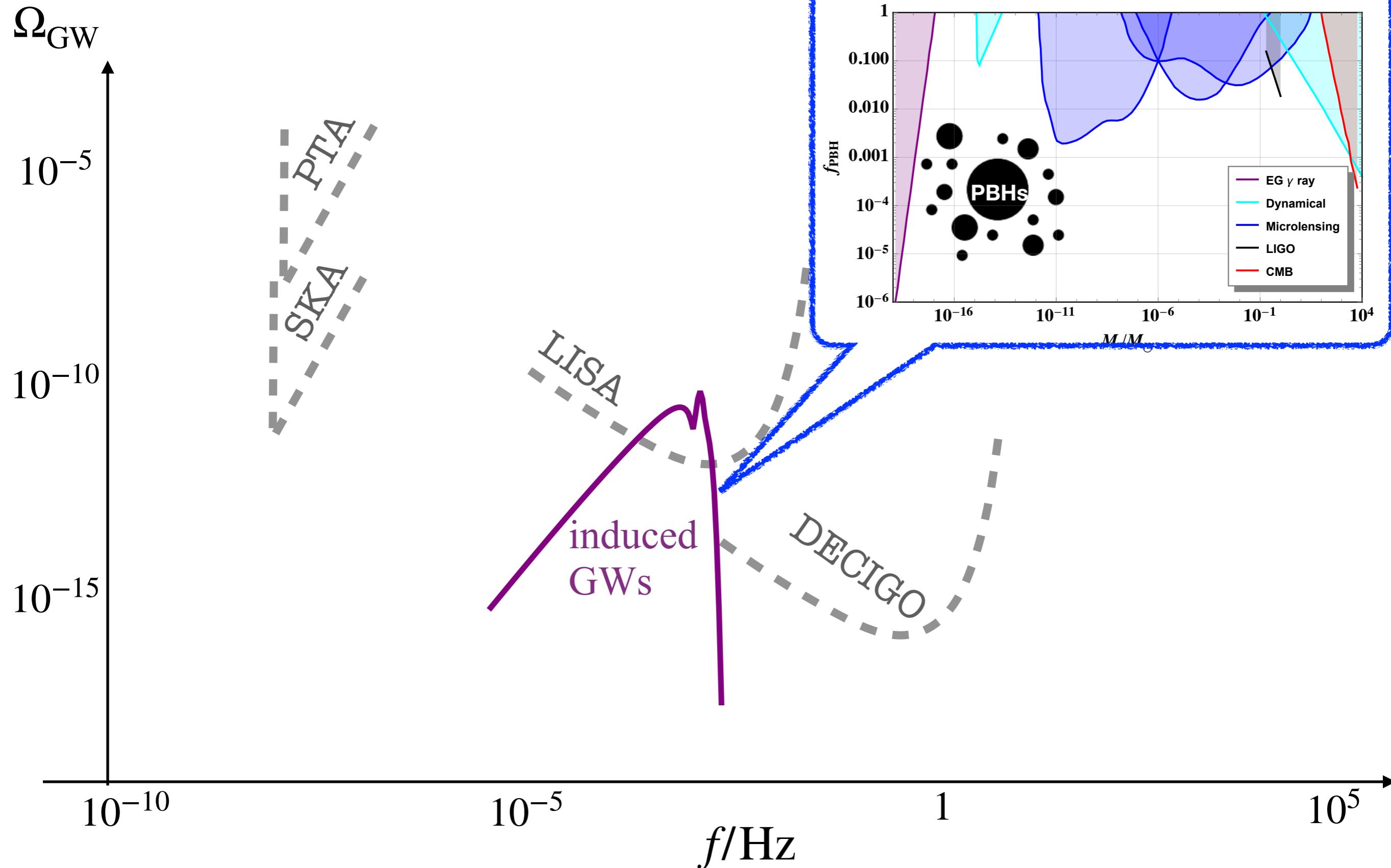












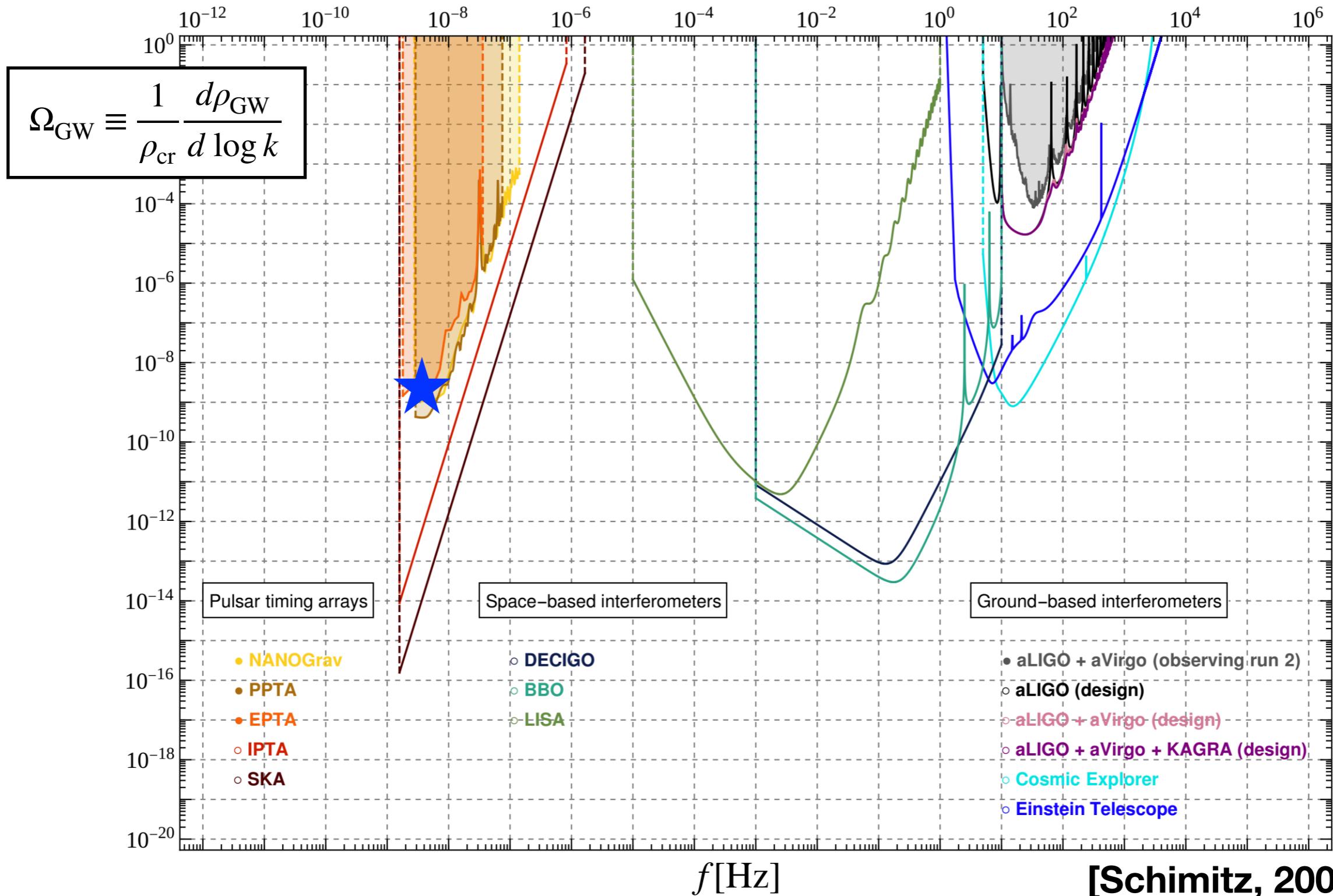
# PTA discovered SGWB?

## The NANOGrav 12.5-year Data Set: Search For An Isotropic Stochastic Gravitational-Wave Background

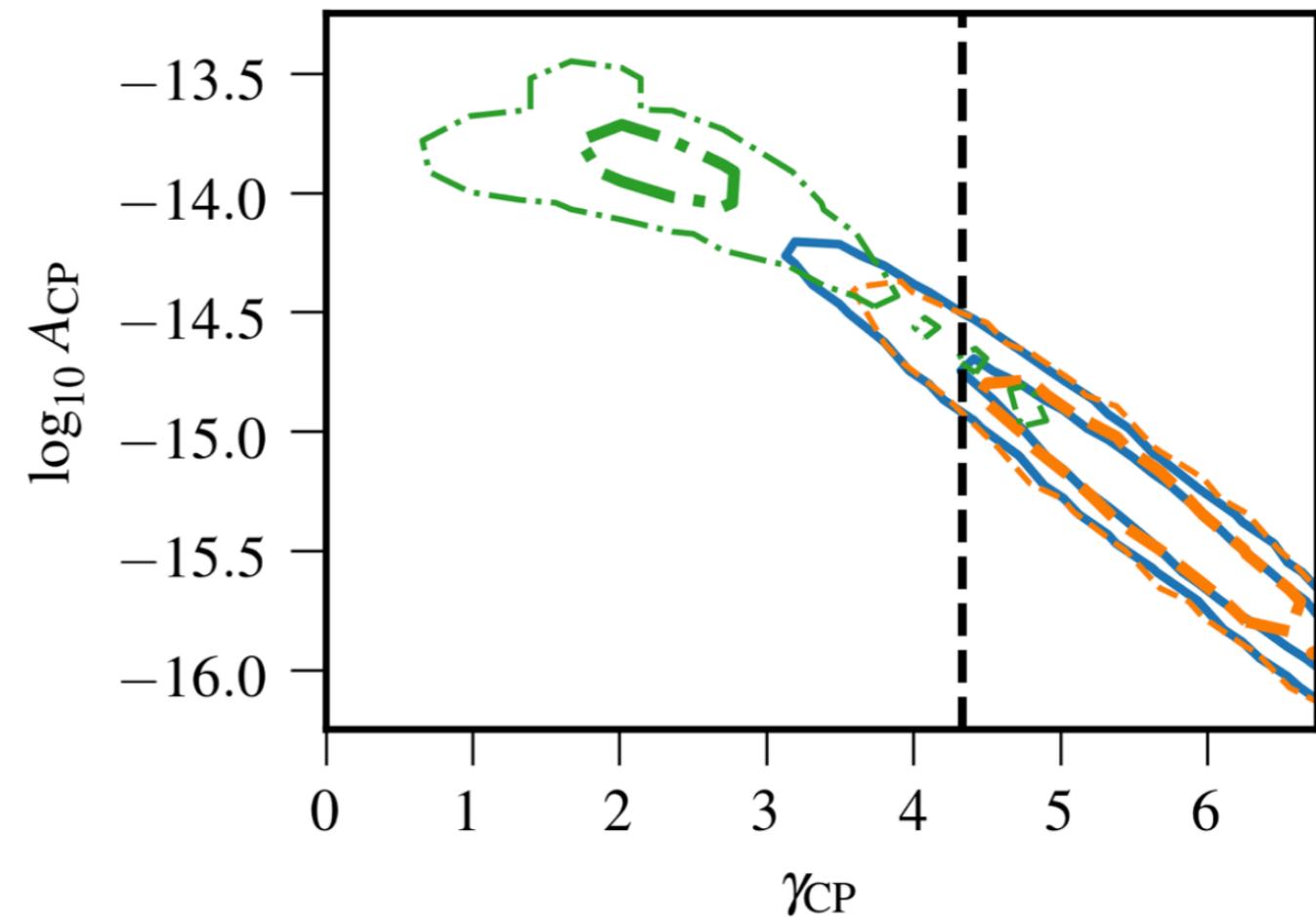
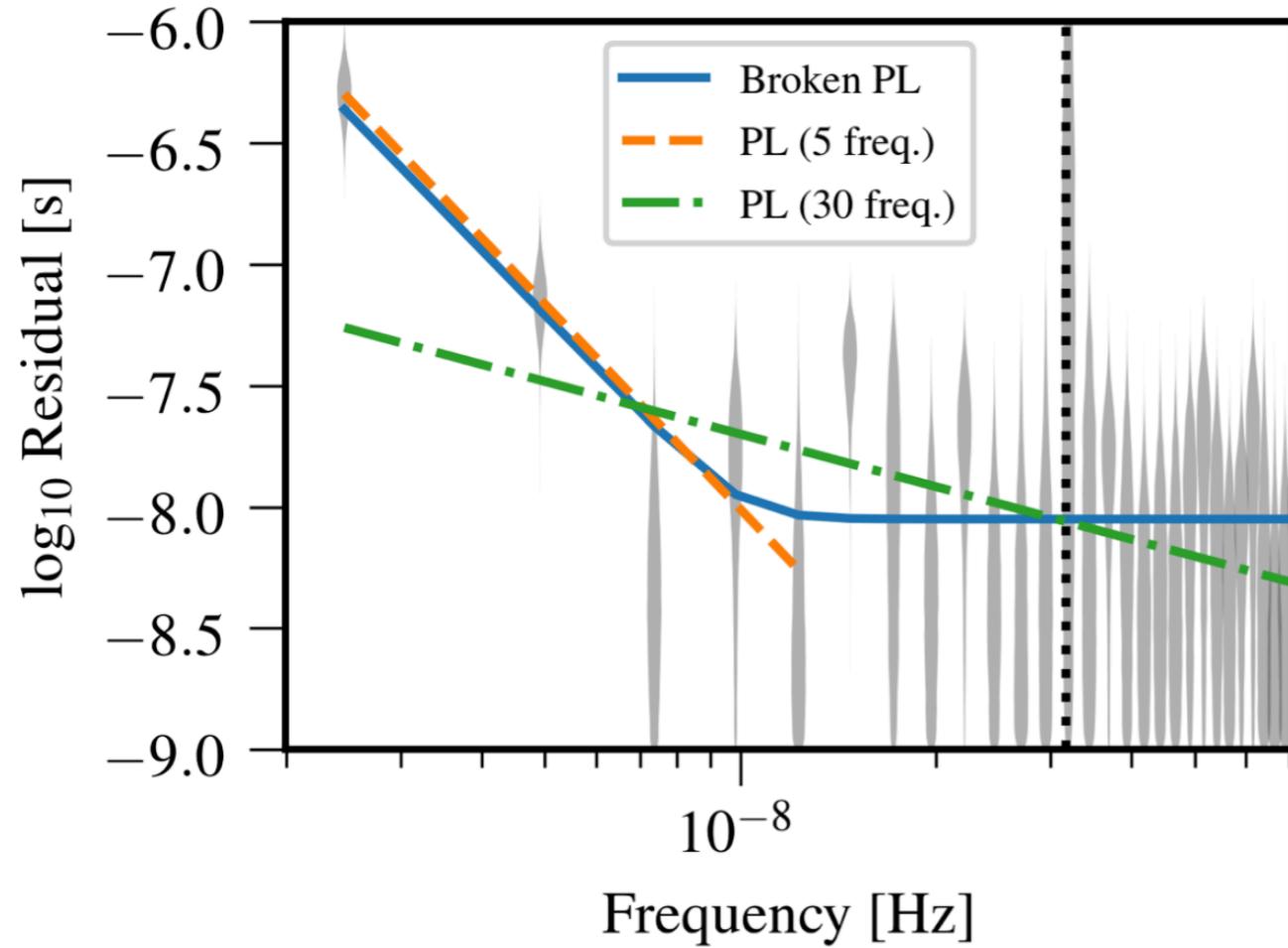
### ABSTRACT

We search for an isotropic stochastic gravitational-wave background (GWB) in the 12.5-year pulsar timing data set collected by the North American Nanohertz Observatory for Gravitational Waves (NANOGrav). Our analysis finds strong evidence of a stochastic process, modeled as a power-law, with common amplitude and spectral slope across pulsars. The Bayesian posterior of the amplitude for a  $f^{-2/3}$  power-law spectrum, expressed as characteristic GW strain, has median  $1.92 \times 10^{-15}$  and 5%–95% quantiles of  $1.37\text{--}2.67 \times 10^{-15}$  at a reference frequency of  $f_{\text{yr}} = 1 \text{ yr}^{-1}$ . The Bayes factor in favor of the common-spectrum process versus independent red-noise processes in each pulsar exceeds 10,000. However, we find no statistically significant evidence that this process has quadrupolar spatial correlations, which we would consider necessary to claim a GWB detection consistent with General Relativity. We find that the process has neither monopolar nor dipolar correlations, which may arise from, for example, reference clock or solar-system ephemeris systematics, respectively. The amplitude posterior has significant support above previously reported upper limits; we explain this in terms of the Bayesian priors assumed for intrinsic pulsar red noise. We examine potential implications for the supermassive black hole binary population under the hypothesis that the signal is indeed astrophysical in nature.

# PTA discovered SGWB?



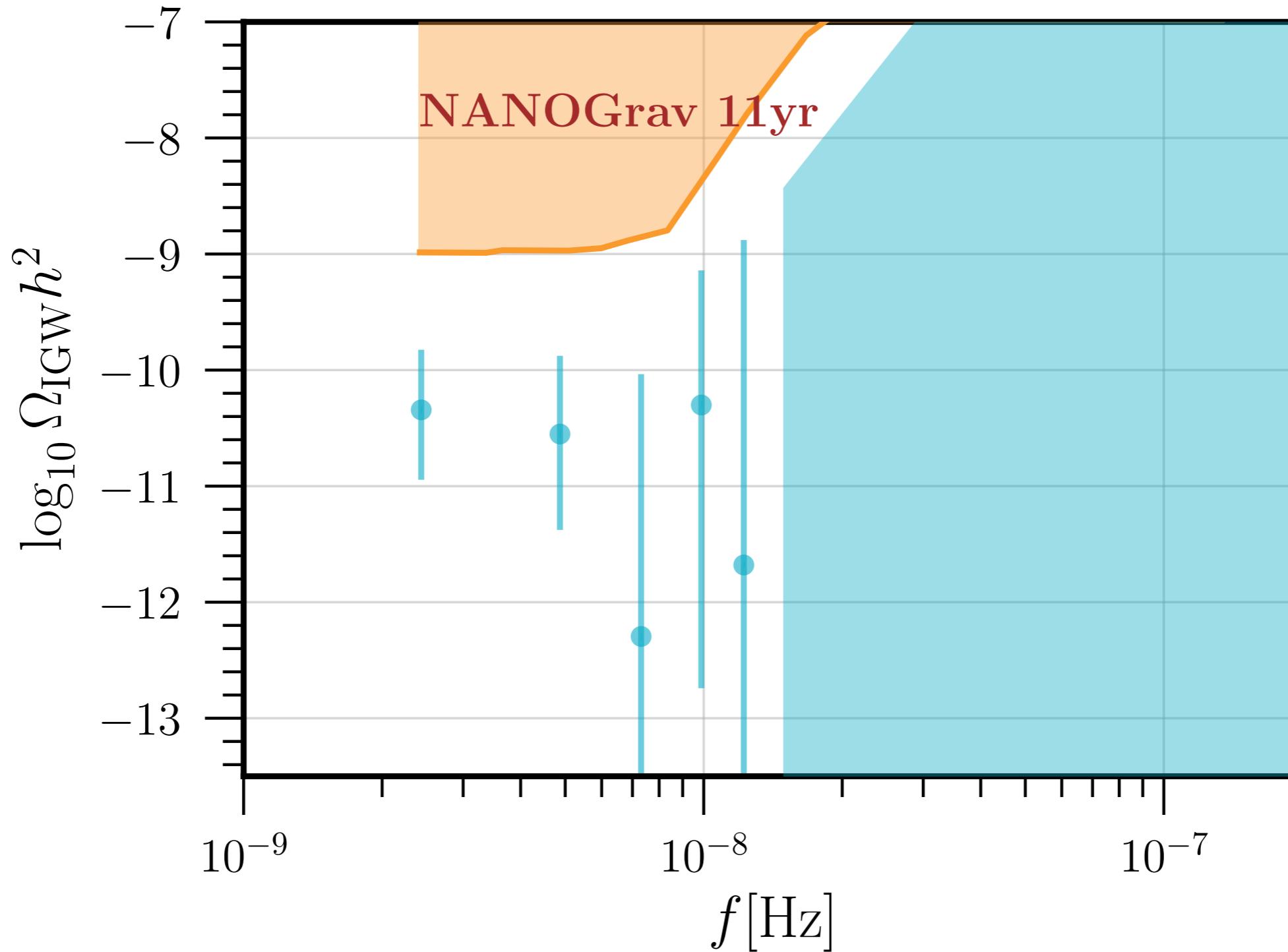
# PTA discovered SGWB?



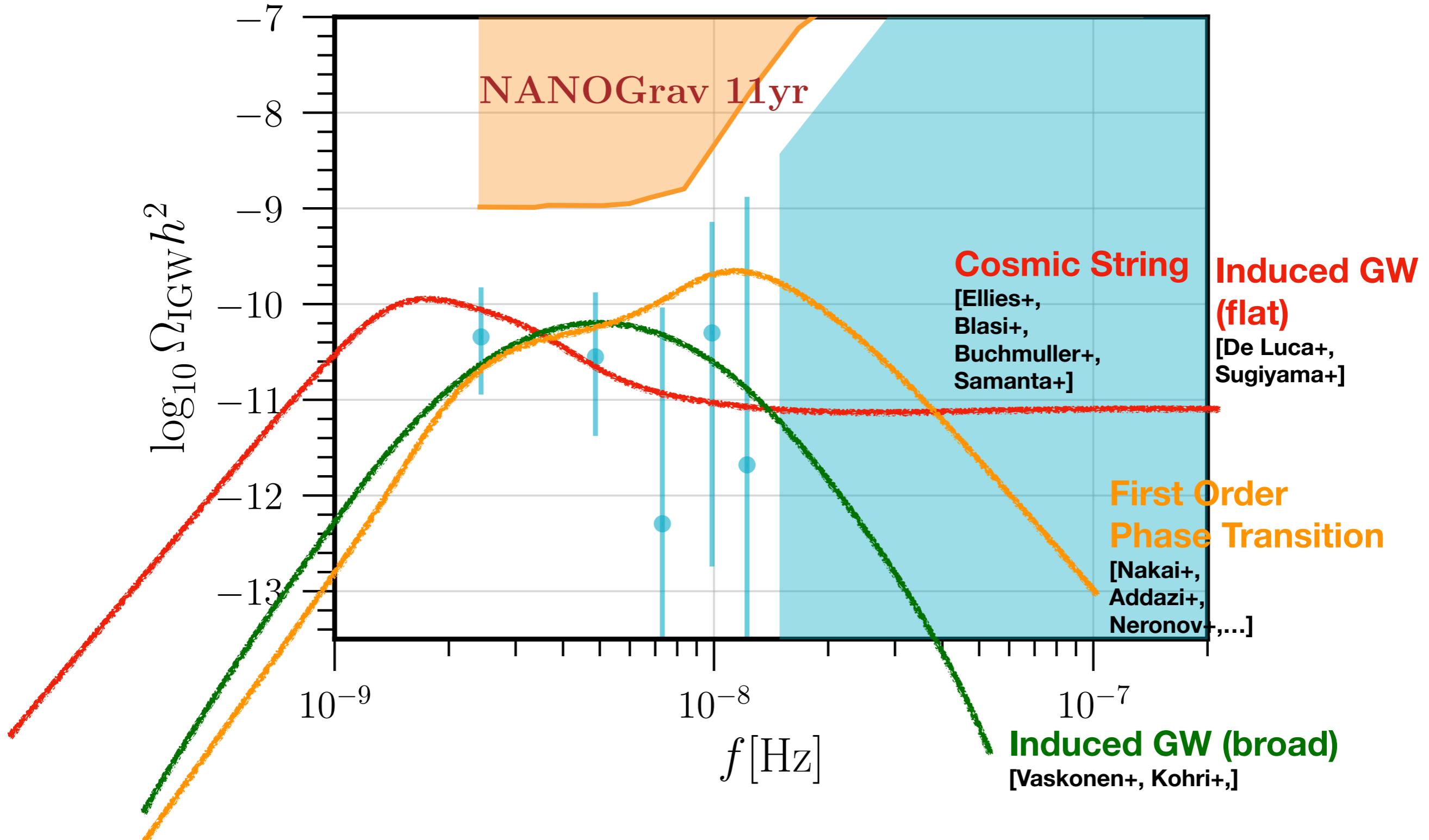
$$\Omega_{\text{GW}} = \frac{2\pi^2 f_{\text{yr}}^2}{3H_0^2} A_{\text{SGWB}}^2 \left( \frac{f}{f_{\text{yr}}} \right)^{5-\gamma}$$

$$f_{\text{yr}} = 3.17 \times 10^{-8} \text{ Hz}.$$

# PTA discovered SGWB?



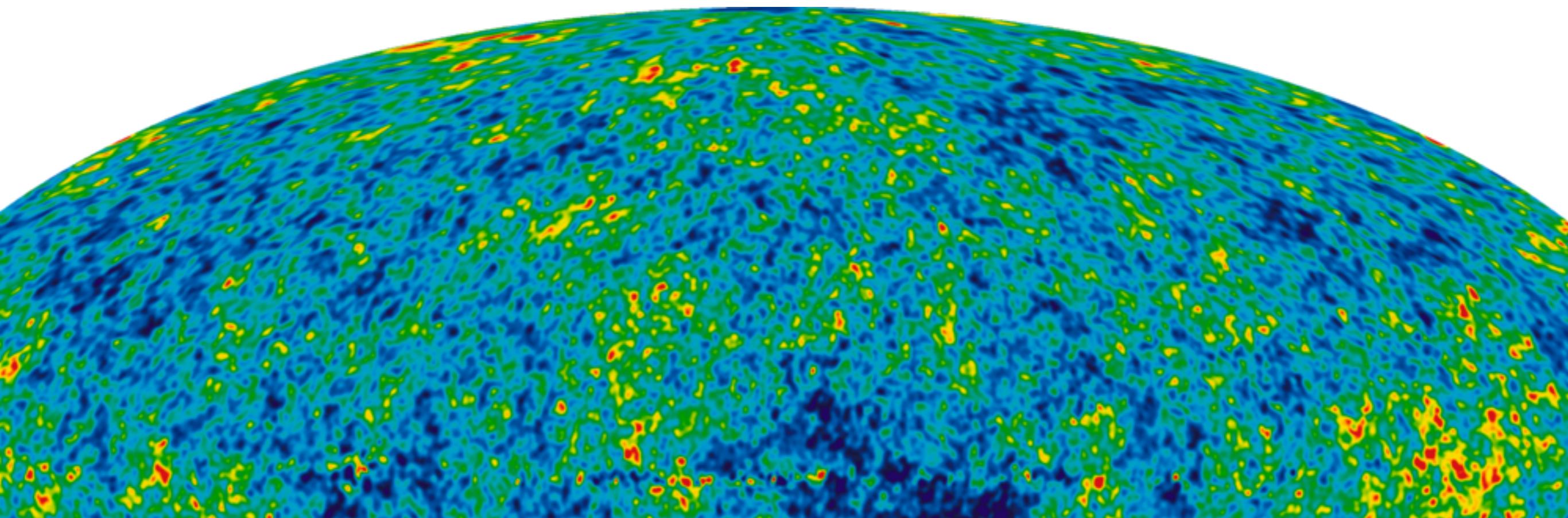
# PTA discovered SGWB?

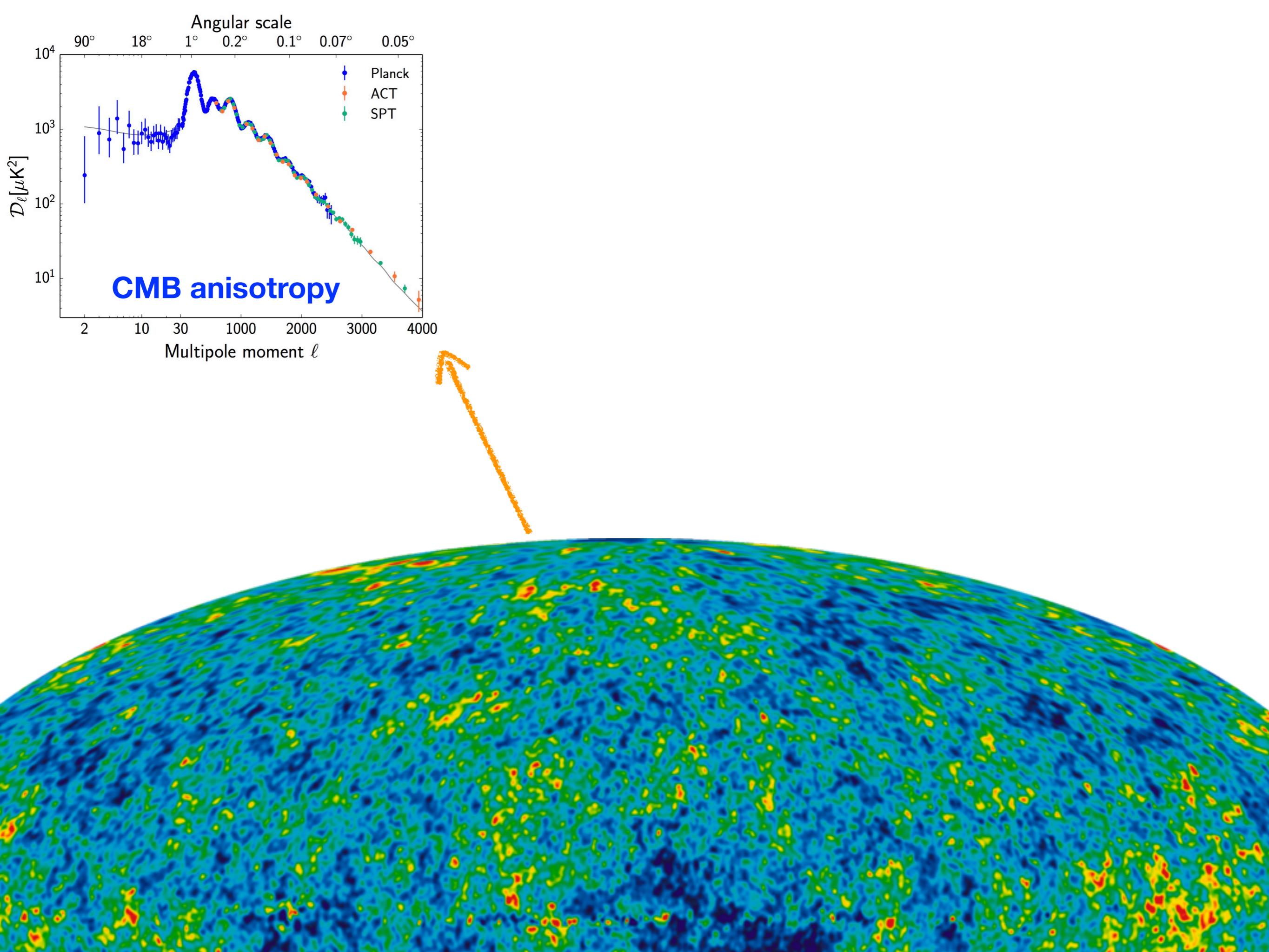


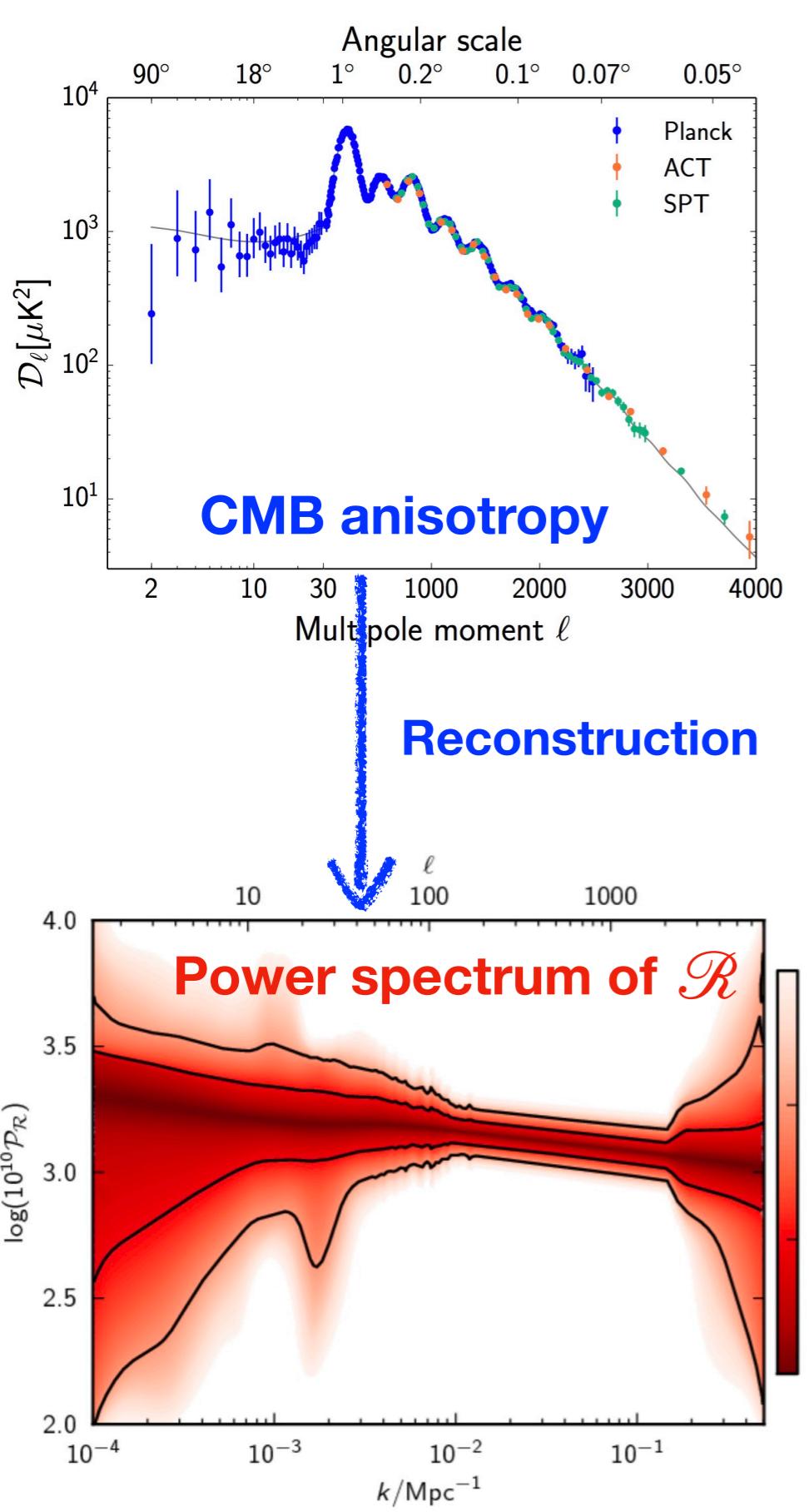
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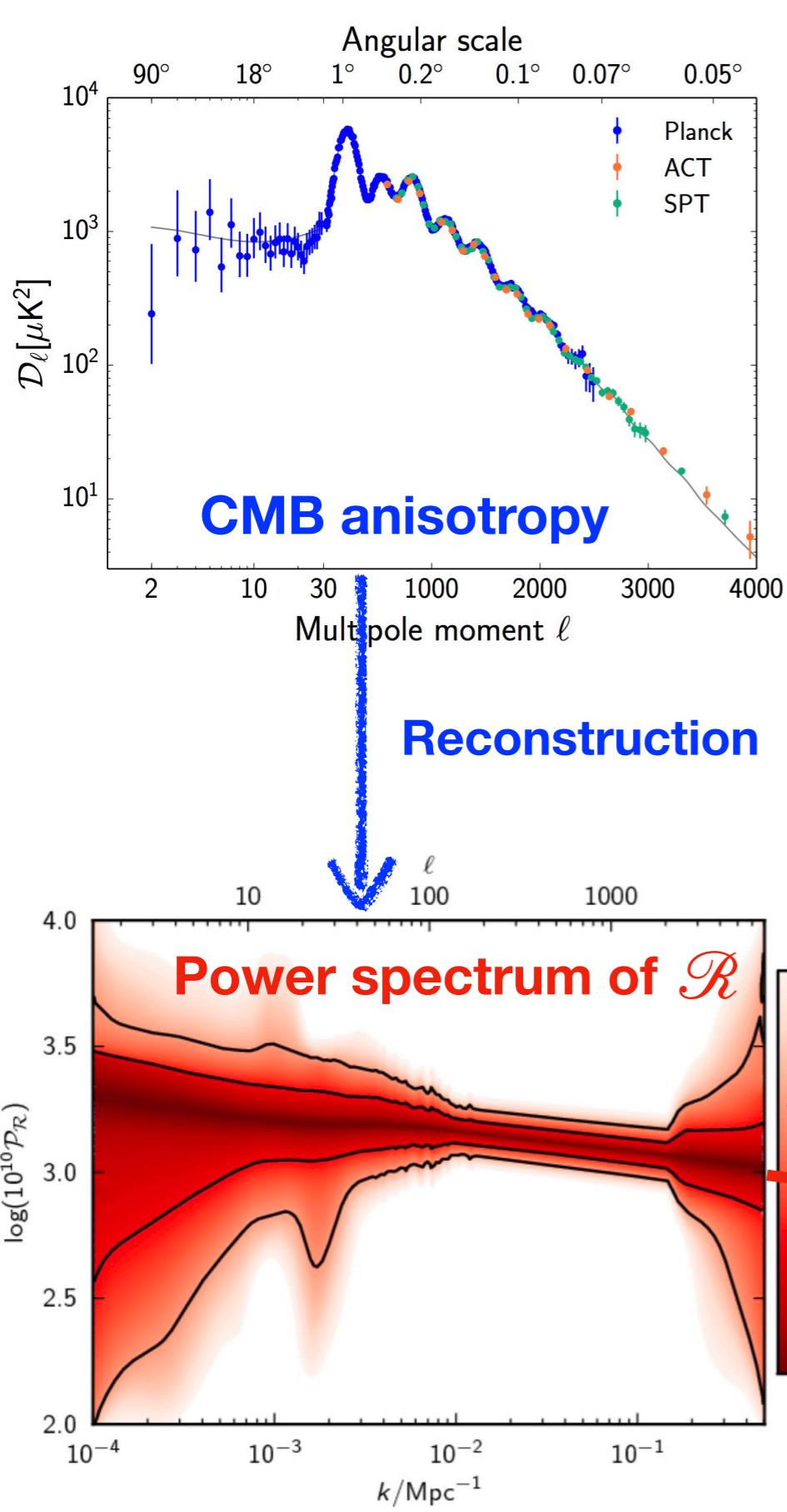
# Introduction to Induced GW (and PBH)

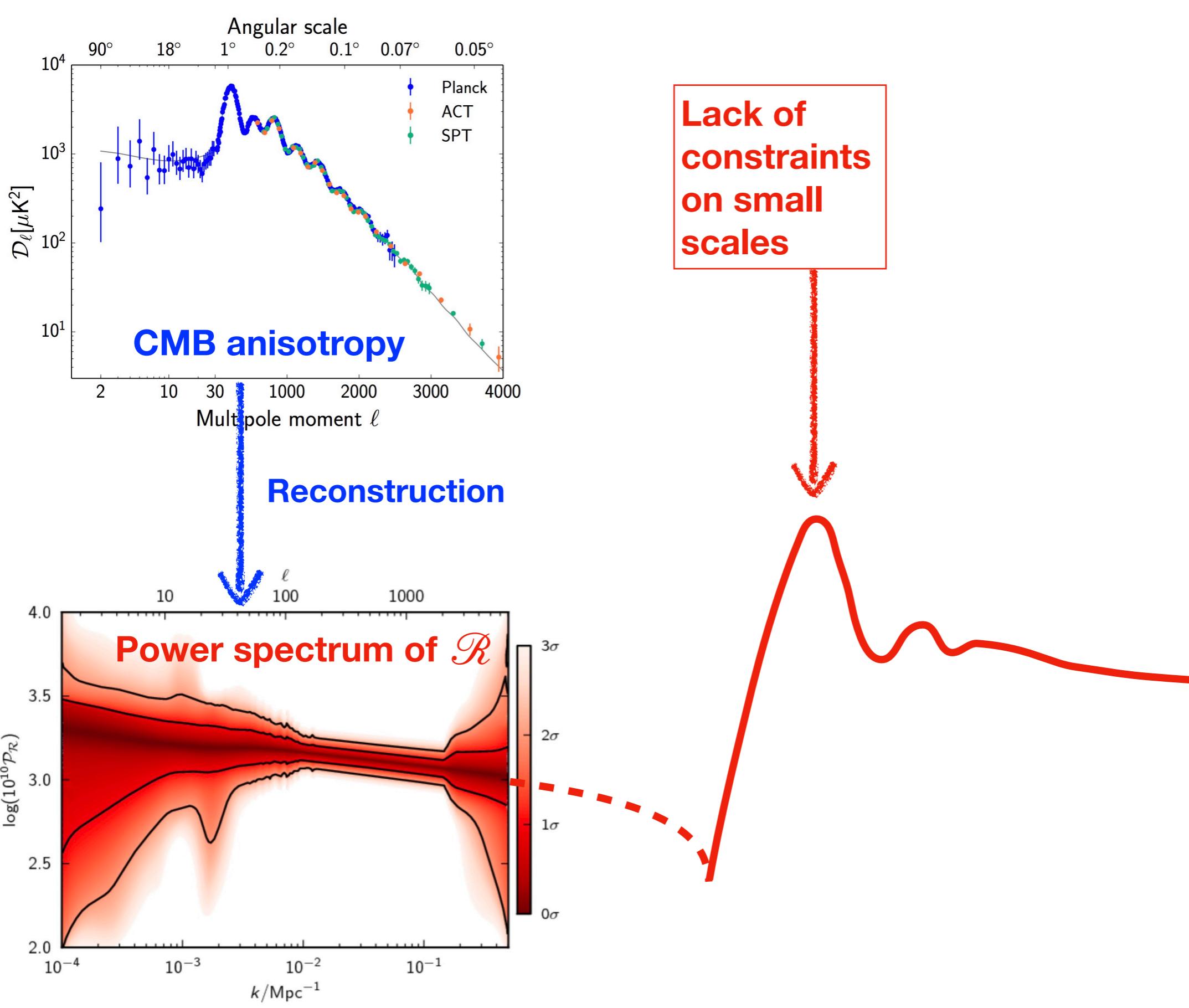


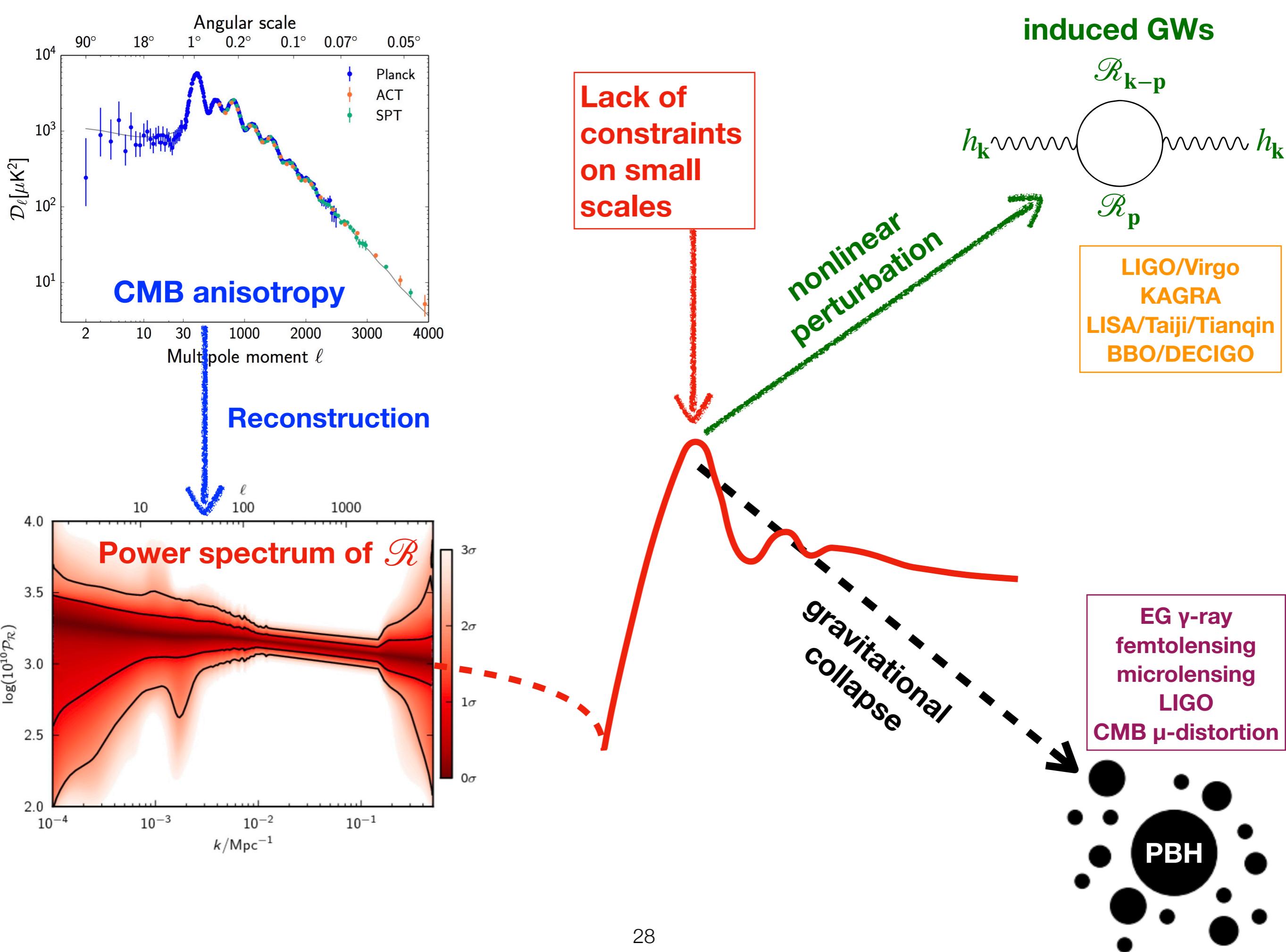


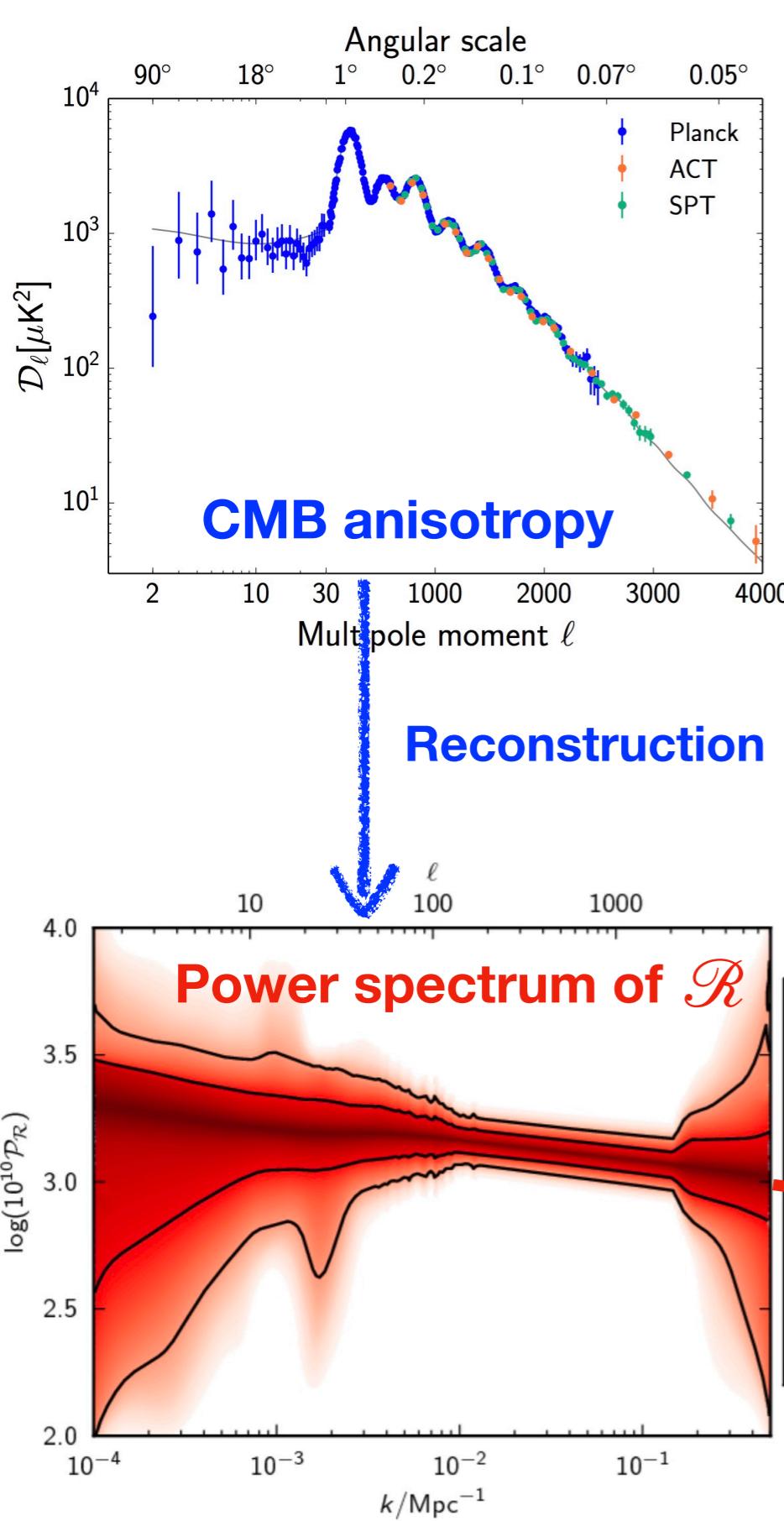


[Planck 2015, 4 knots]



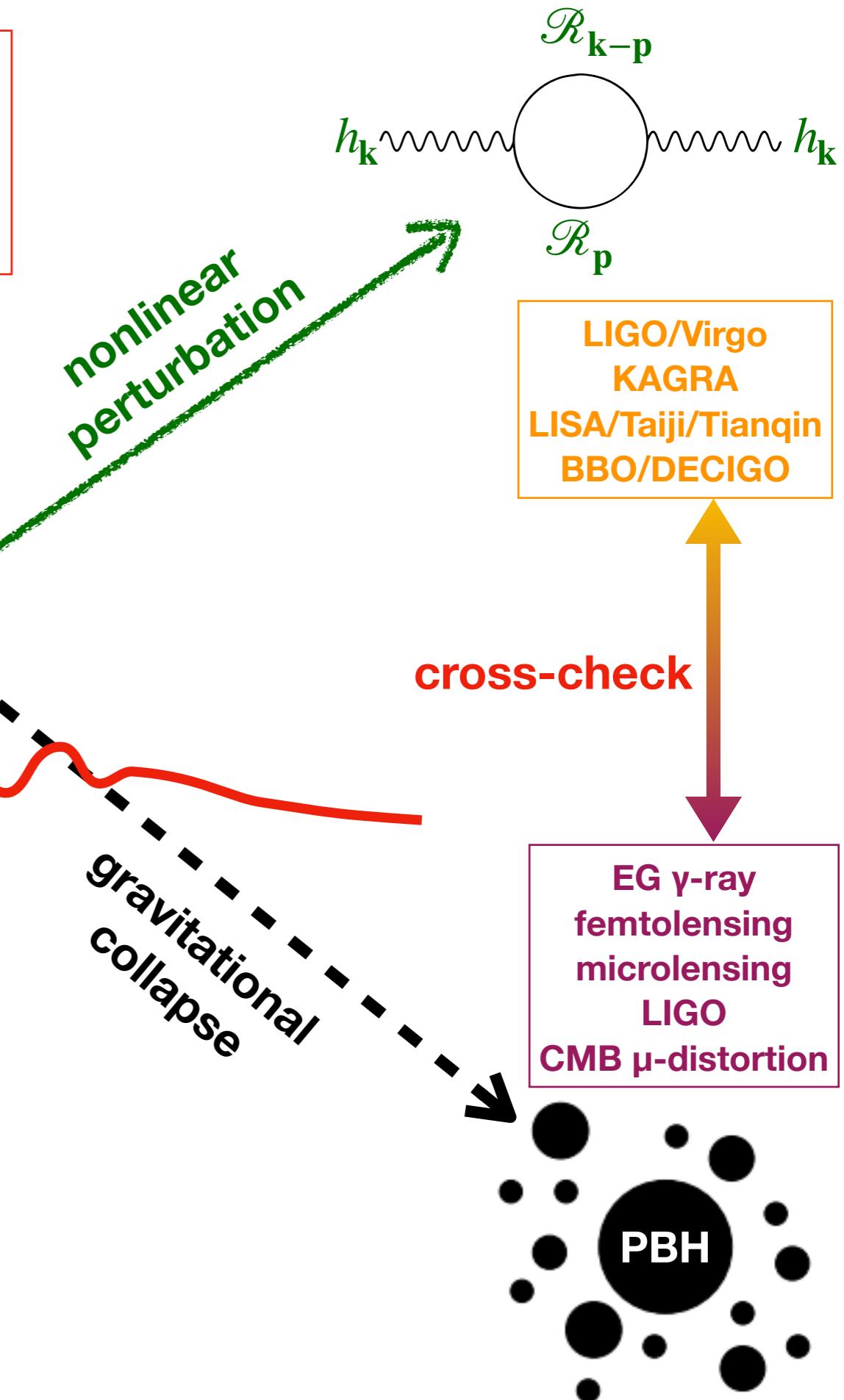


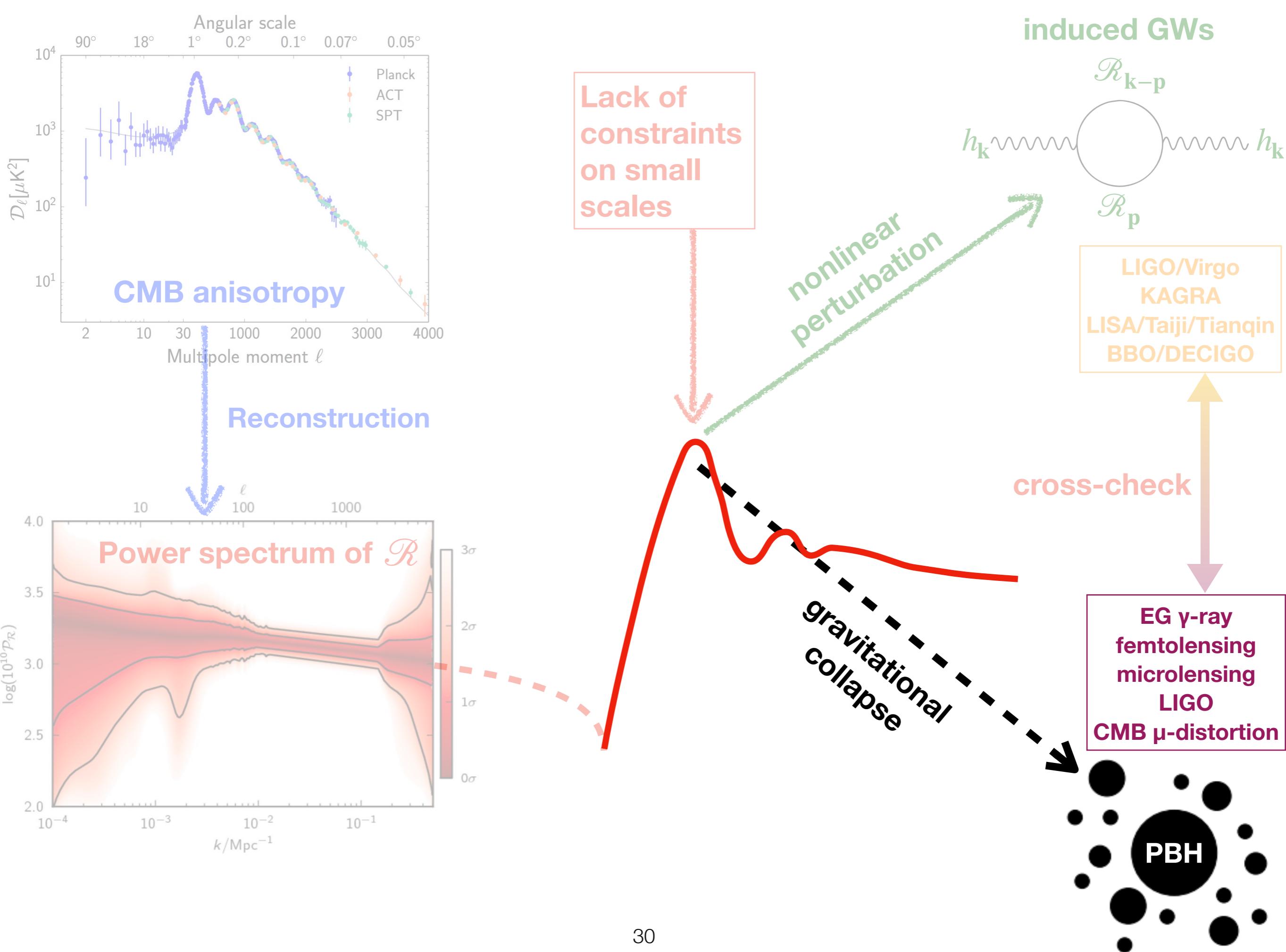




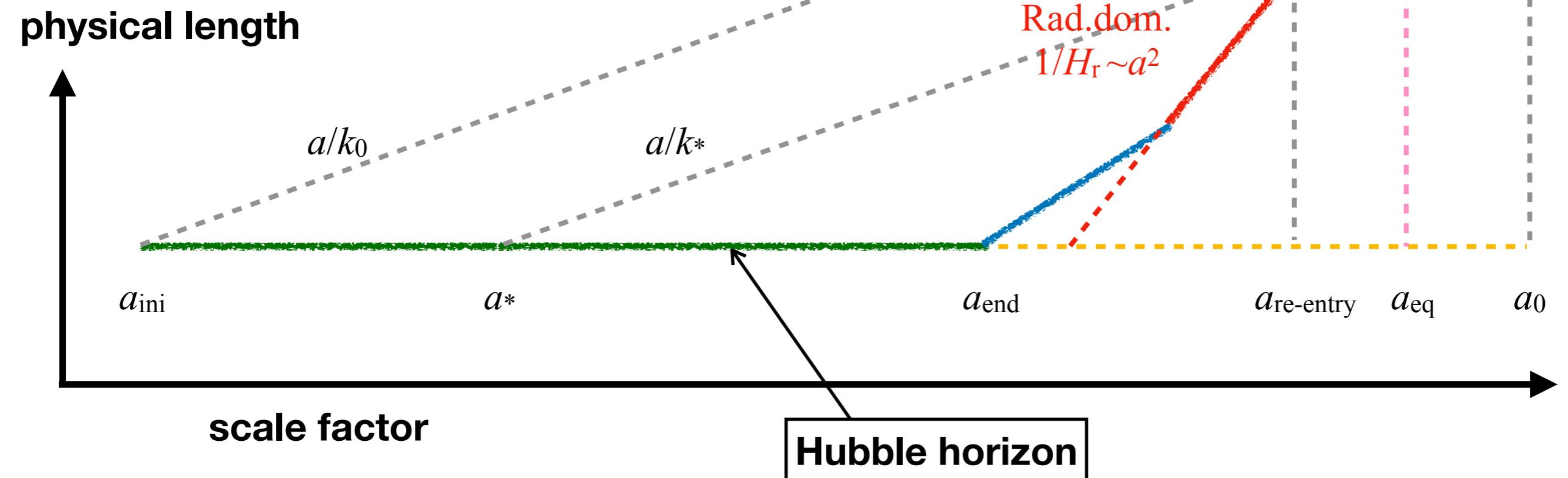
## Lack of constraints on small scales

# induced GWs

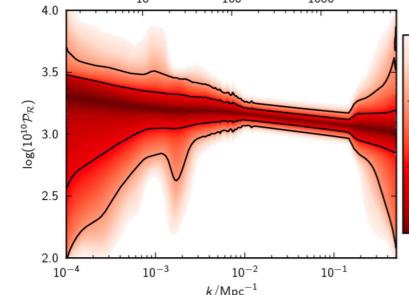
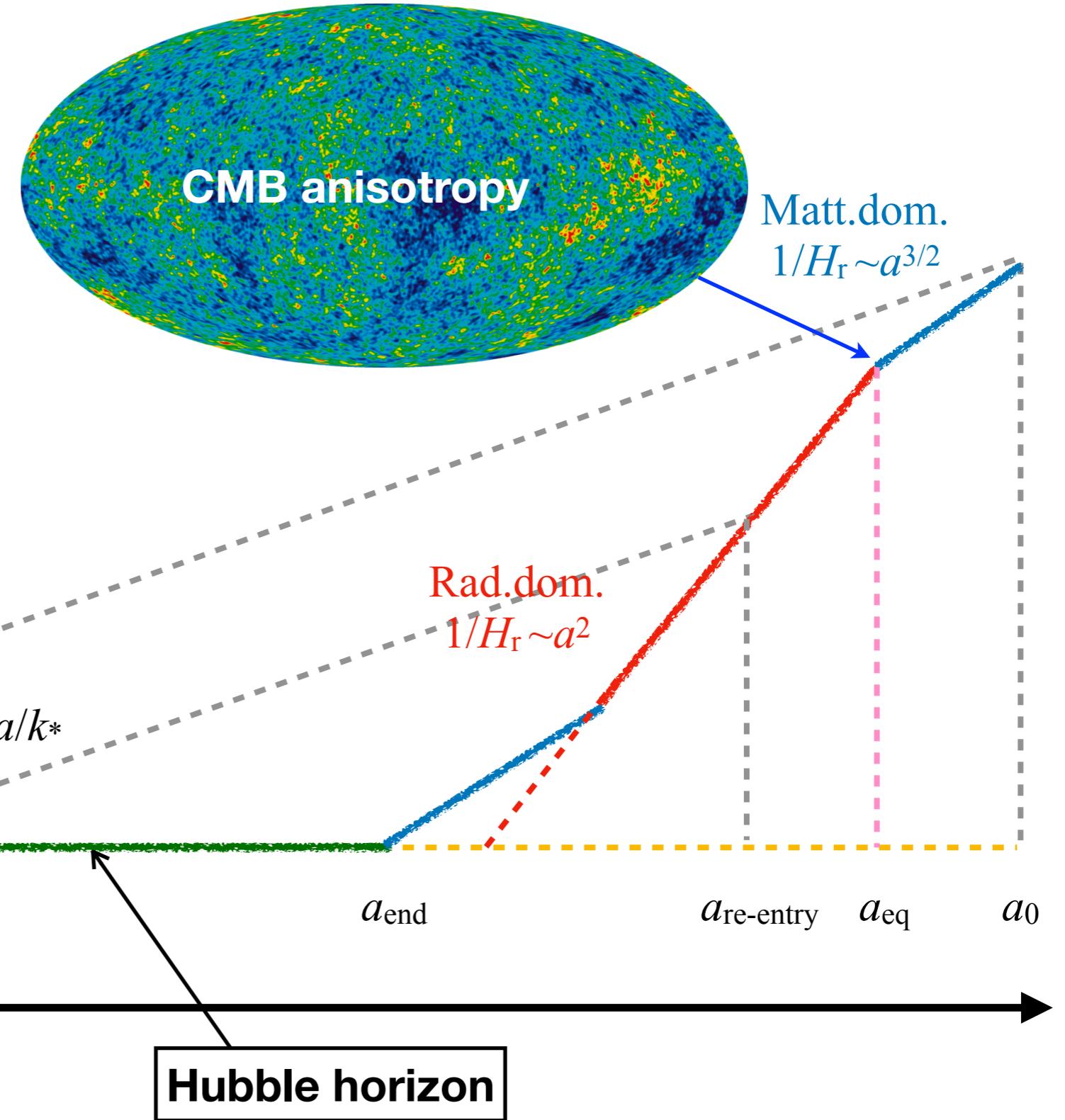




# The formation of Primordial Black Holes

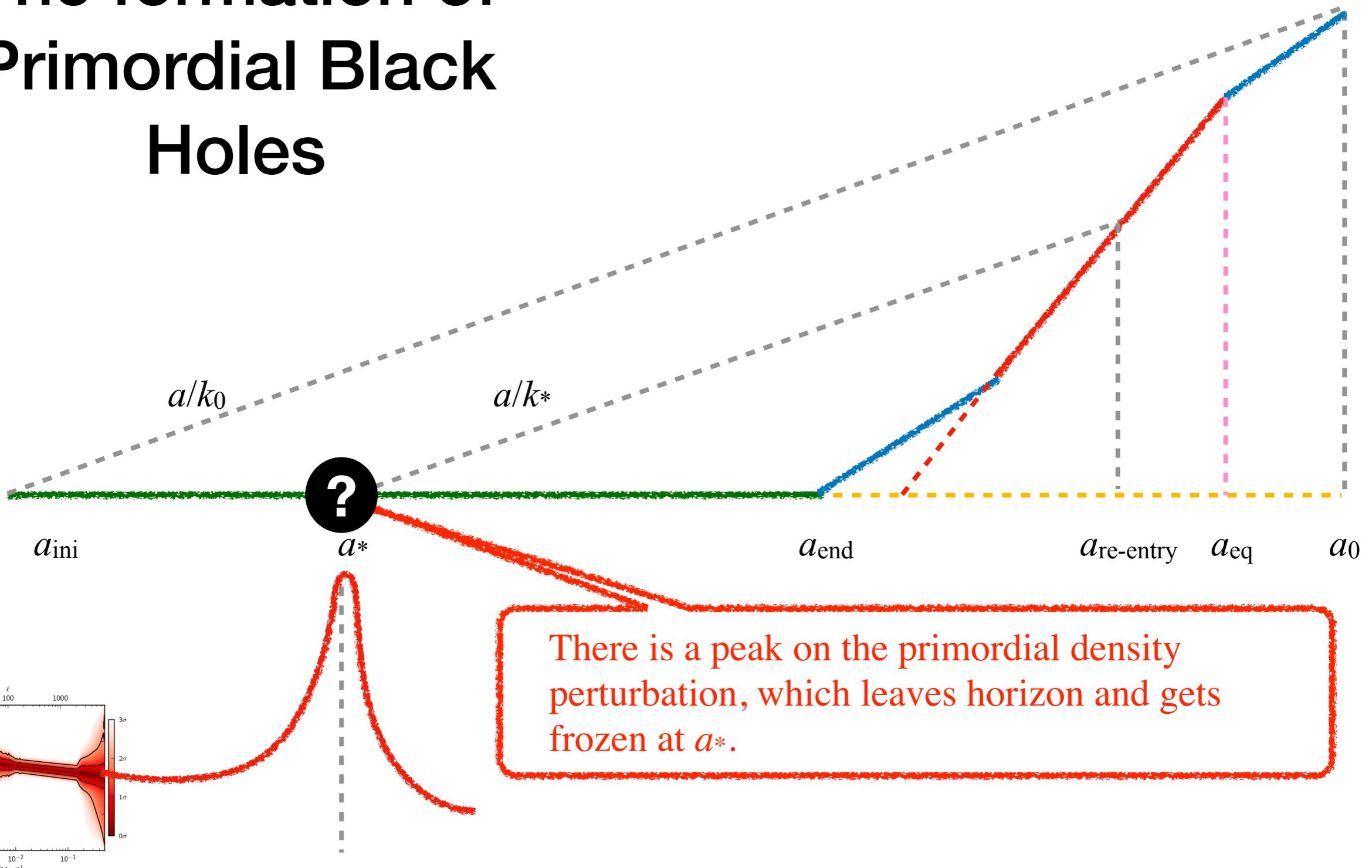


# The formation of Primordial Black Holes

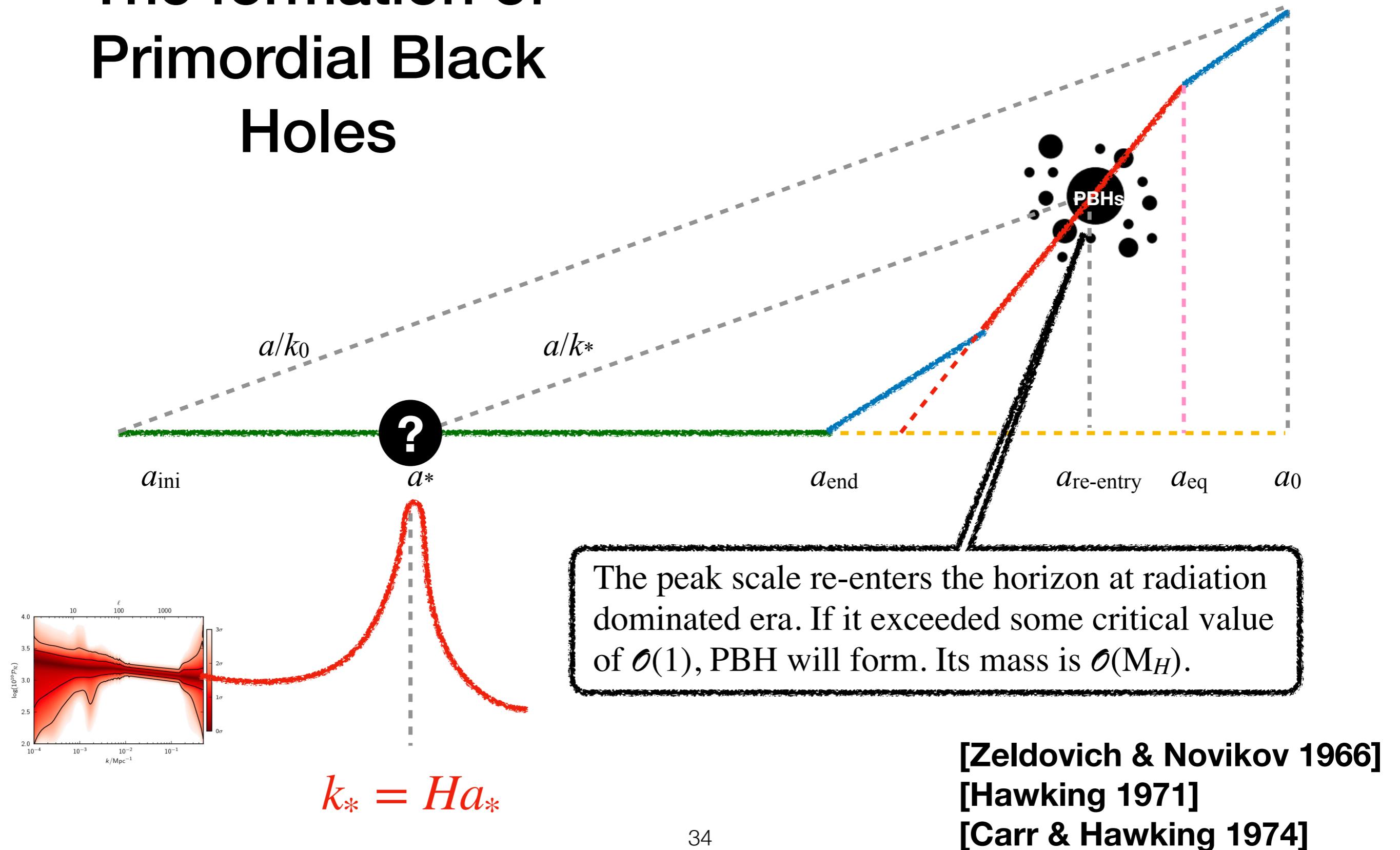


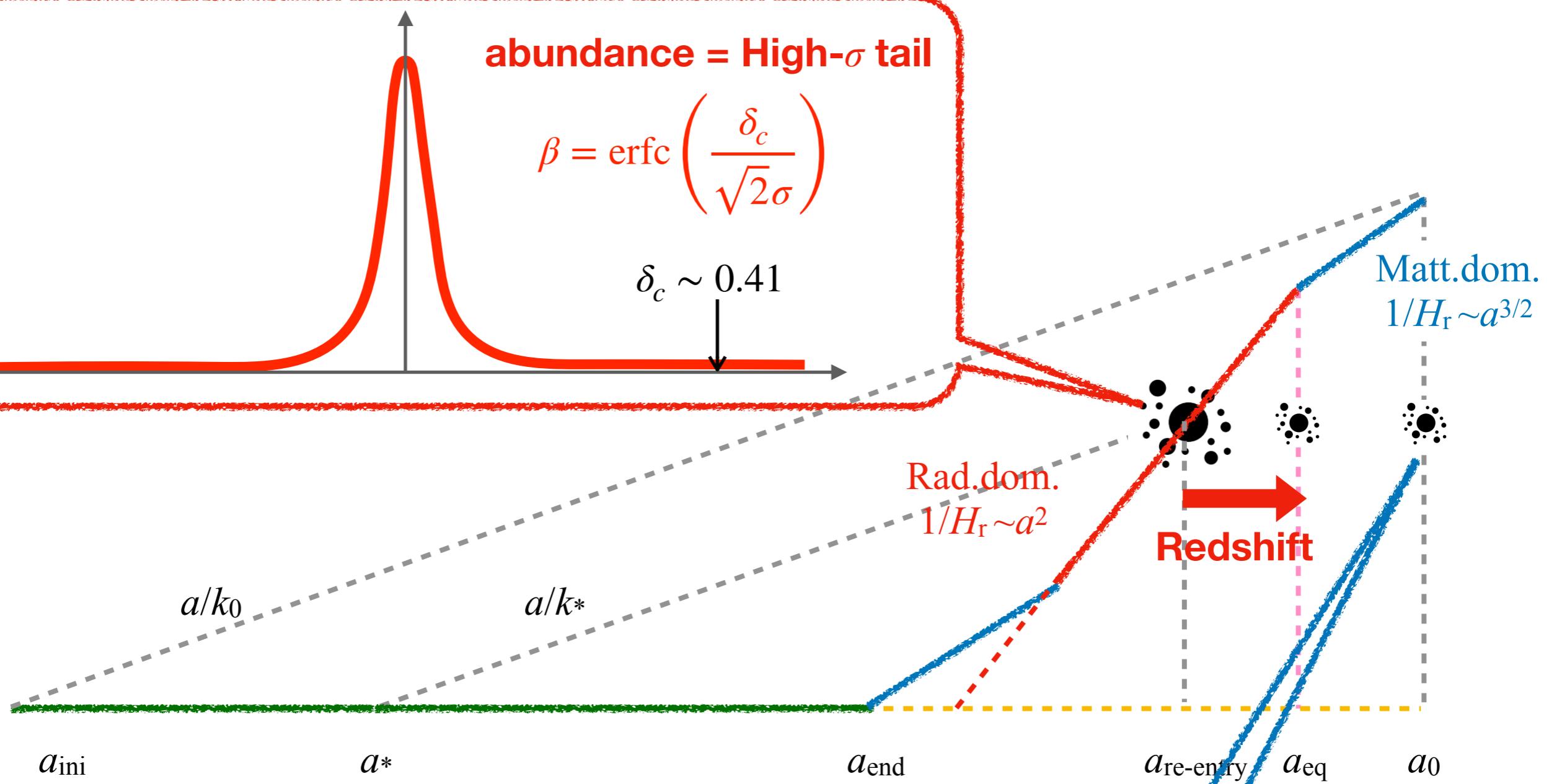
Power spectrum of  $\mathcal{R}$

# The formation of Primordial Black Holes



# The formation of Primordial Black Holes



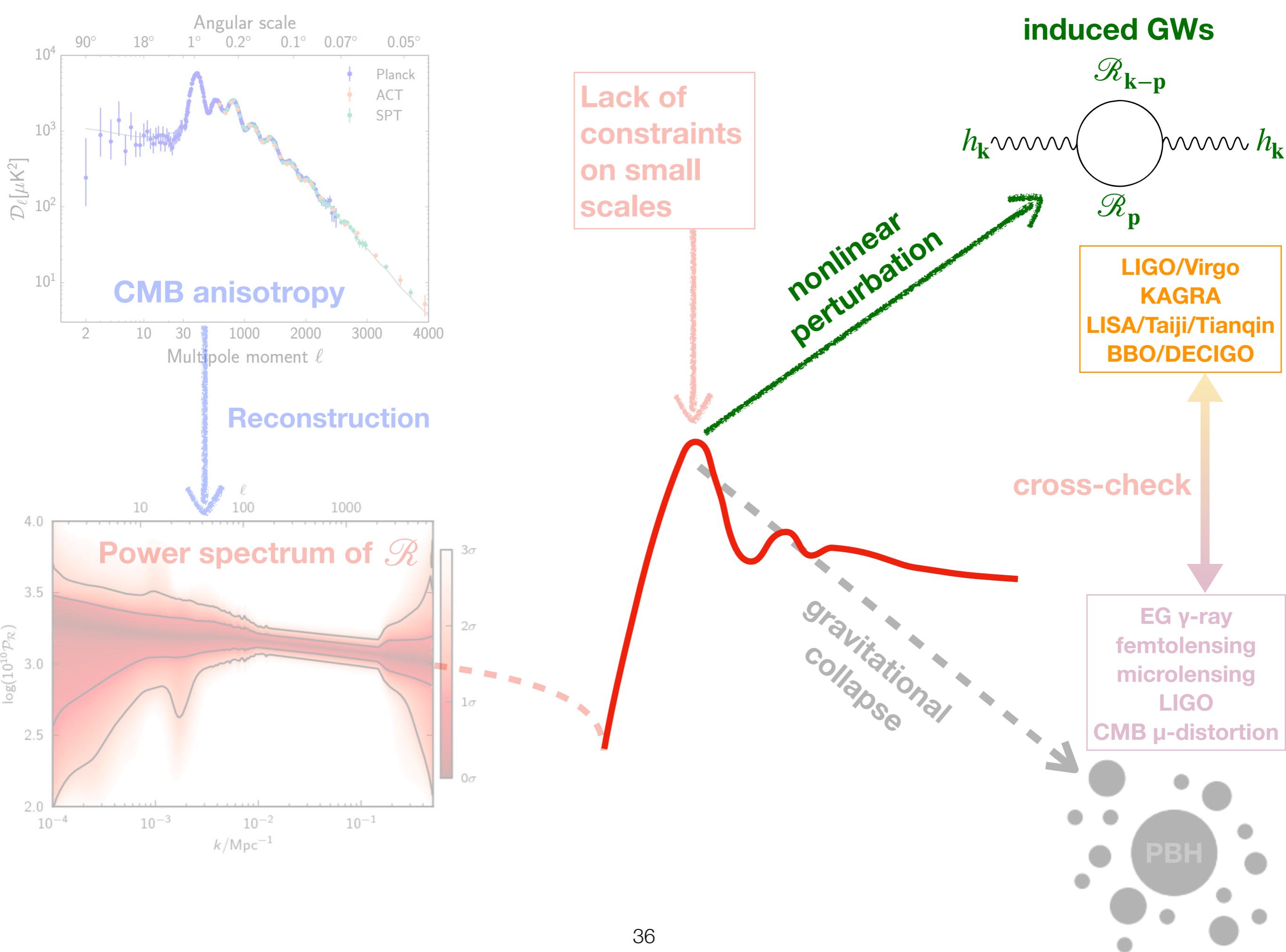


$$\Omega_{\text{PBH}} = \beta \frac{a_{\text{eq}}}{a_{\text{re}}} = \beta \frac{a_{\text{eq}}}{a_0} \frac{a_0}{a_{\text{re}}} \simeq \beta \Omega_r (1 + z_{\text{re}}(M))$$

$$M = \frac{c^3}{G H_{\text{re}}} = \frac{c^3}{G \Omega_r^{1/2} (1+z)^2 H_0}$$

**Redshifted after equality:**

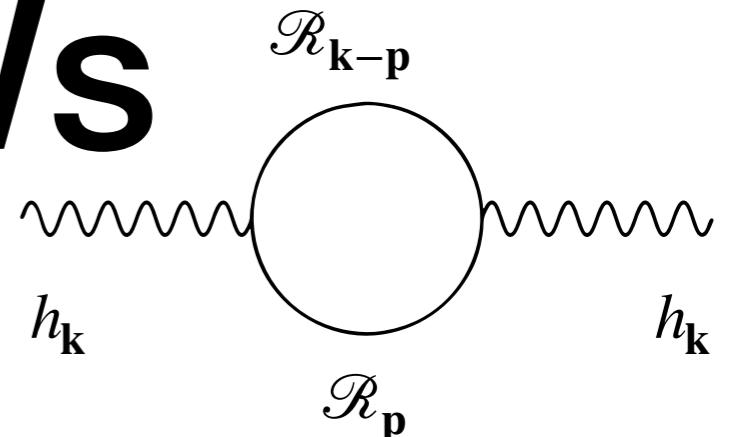
$$f_{\text{PBH}} \equiv \frac{\Omega_{\text{PBH}}}{\Omega_{\text{CDM}}} = 4.11 \times 10^8 \beta(M) \left( \frac{M}{M_\odot} \right)^{-1/2}$$



# Induced GWs

- EoM of the tensor perturbation:

$$h_{\mathbf{k}}'' + 2\mathcal{H}h_{\mathbf{k}}' + k^2 h_{\mathbf{k}} \sim \int d^3 p \Phi_{\mathbf{p}} \Phi_{\mathbf{k}-\mathbf{p}} \times (\text{transfer functions})$$



- The induced (secondary) GWs are

$$h_{\mathbf{k}} \sim \int d\eta \times (\text{Green function}) \int d^3 p \times (\text{Transfer function}) \times \Phi_{\mathbf{p}} \Phi_{\mathbf{k}-\mathbf{p}}$$

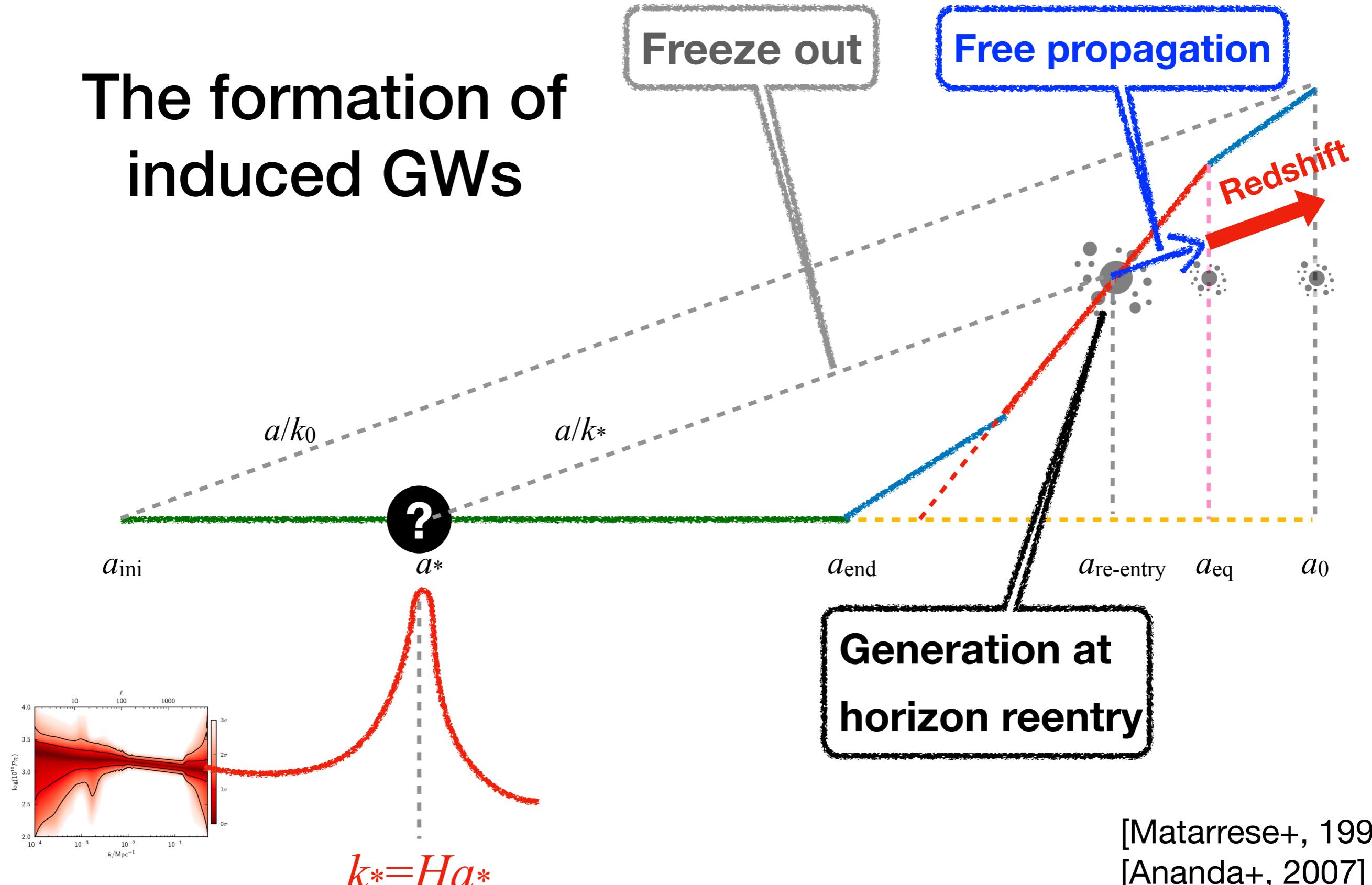
- The energy density parameter is then

$$\Omega_{\text{GW}} \sim \langle hh \rangle \sim \langle \Phi \Phi \Phi \Phi \rangle \sim \mathcal{P}_{\Phi}^2 \sim \mathcal{P}_{\mathcal{R}}^2$$

- In the radiation dominated universe we have  $\mathcal{R} = \frac{2}{3}\Phi$

[Baumann+ 2007, Ananda+ 2007, Saito and Yokoyama 2008]

# The formation of induced GWs



[Matarrese+, 1997]  
 [Ananda+, 2007]  
 [Baumann+, 2007]

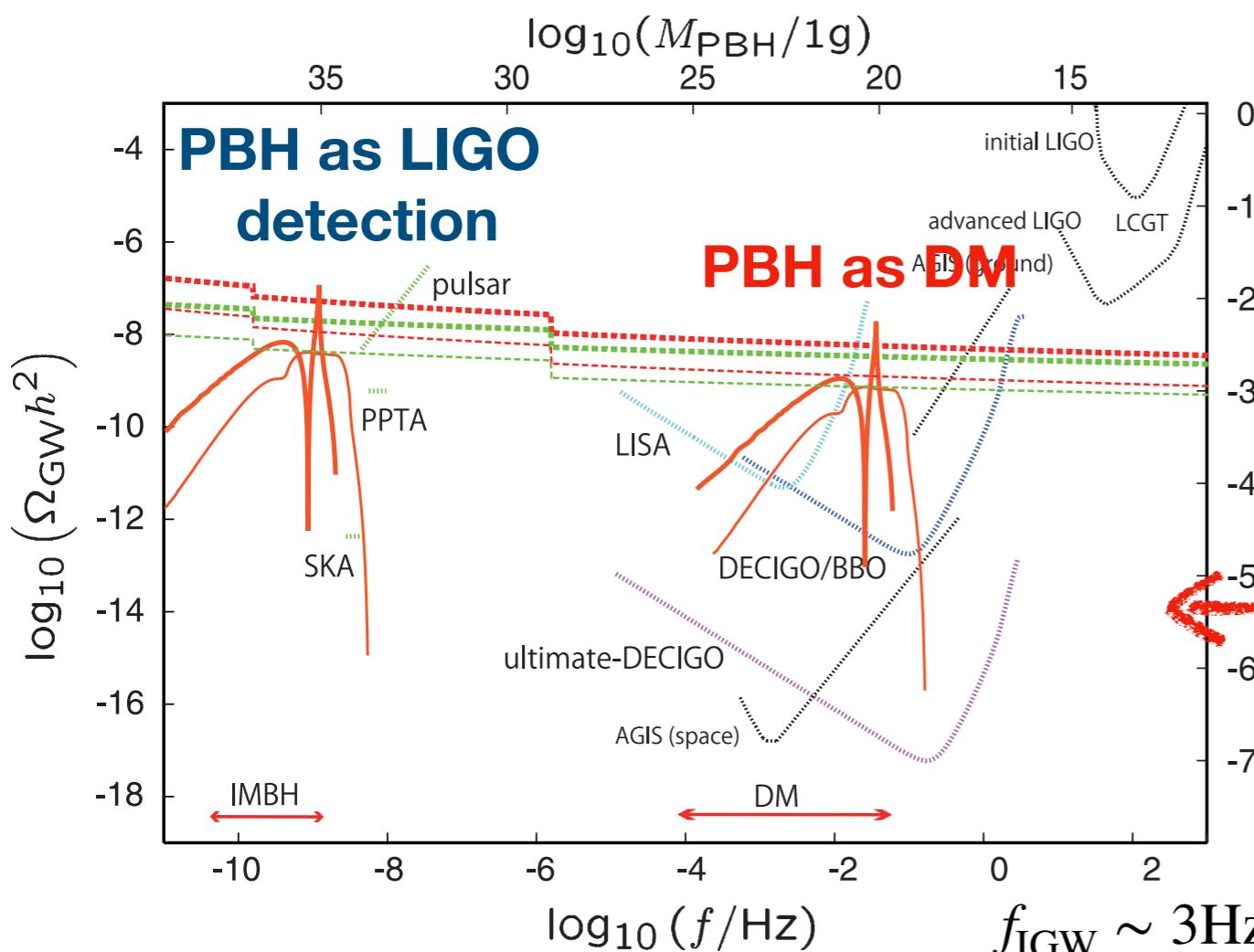
# Induced GWs

$$\Omega_{\text{GW}} \sim \langle hh \rangle \sim \langle \Phi \Phi \Phi \Phi \rangle \sim \mathcal{P}_\Phi^2 \sim \boxed{\mathcal{P}_{\mathcal{R}}^2}$$

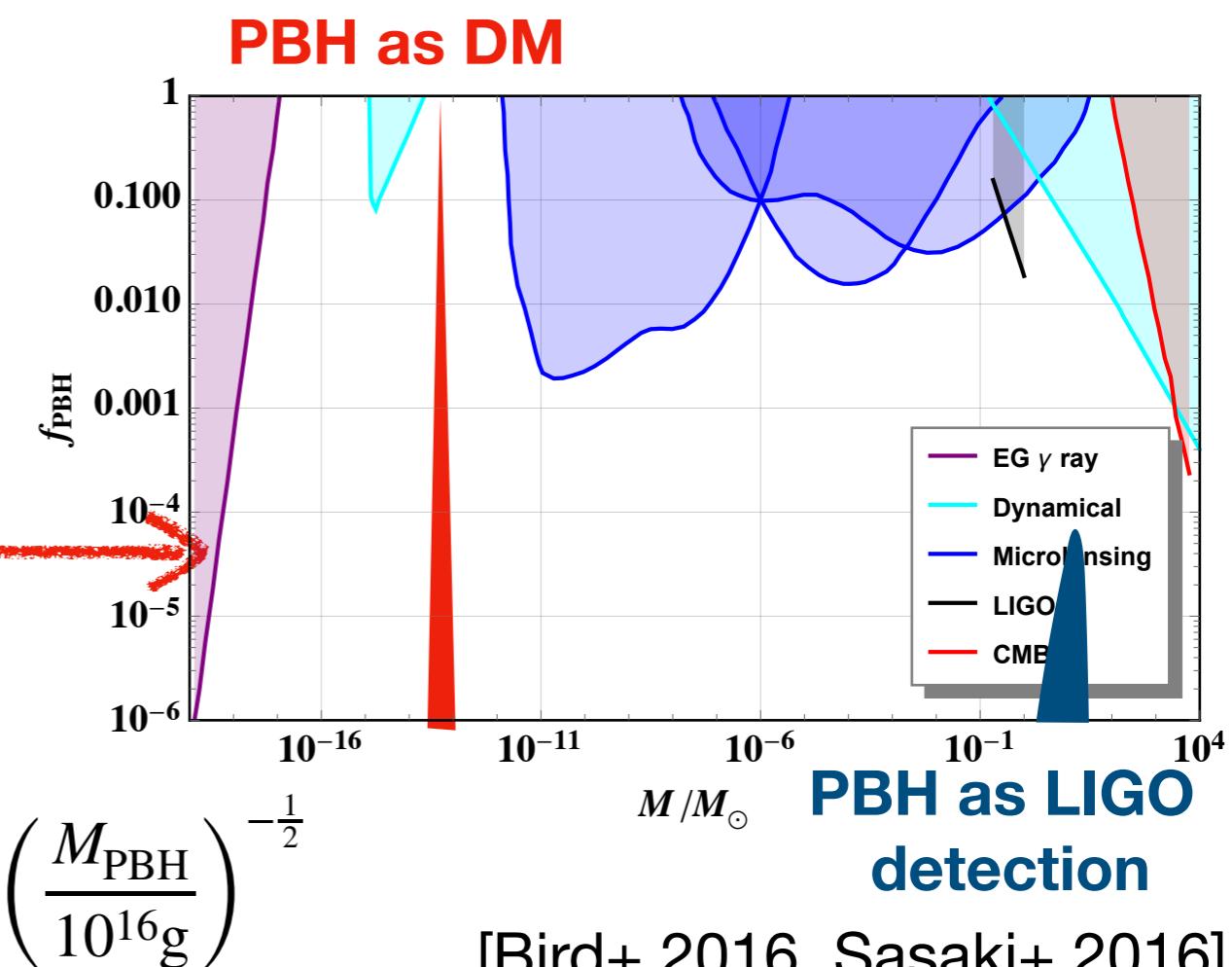
## PBH abundance

$$f_{\text{PBH}} \sim 4.11 \times 10^8 \beta(M) \left( \frac{M}{M_\odot} \right)^{-1/2}$$

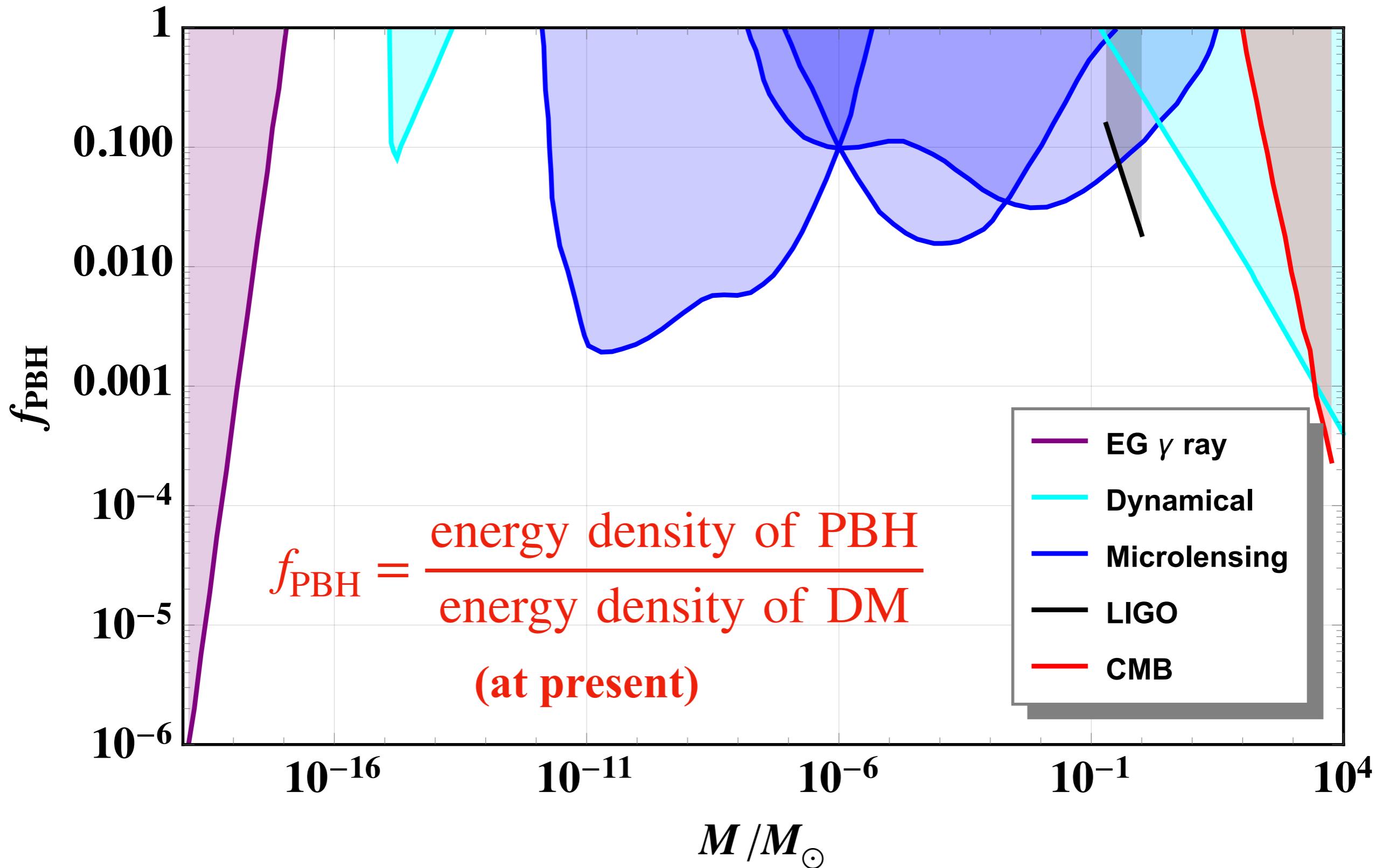
$$\beta \sim \text{erfc} \left( \frac{\mathcal{R}_c}{2 \boxed{\mathcal{P}^{1/2} \mathcal{R}}} \right)$$

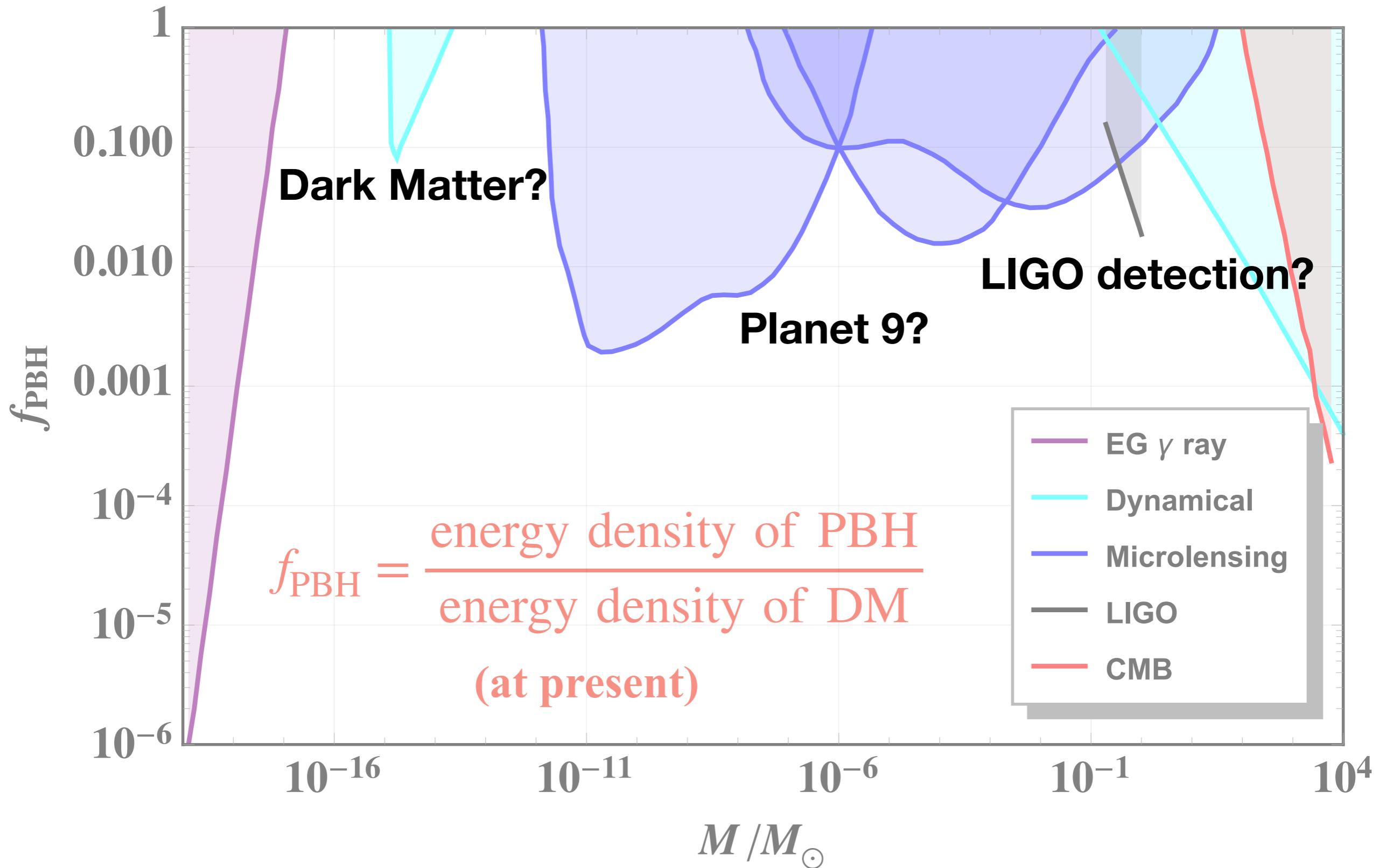


[Saito & Yokoyama 2008]

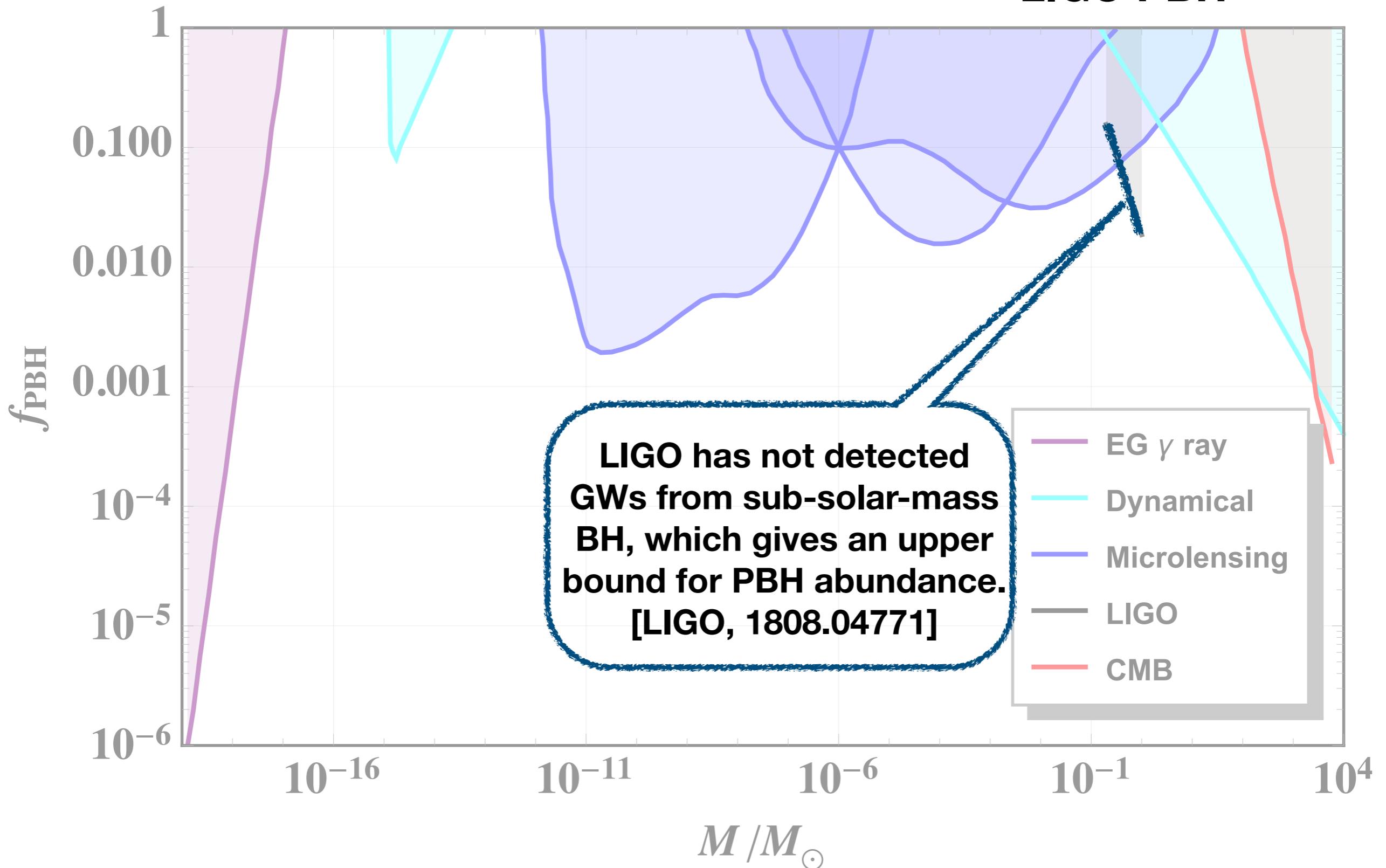


[Bird+ 2016, Sasaki+ 2016]

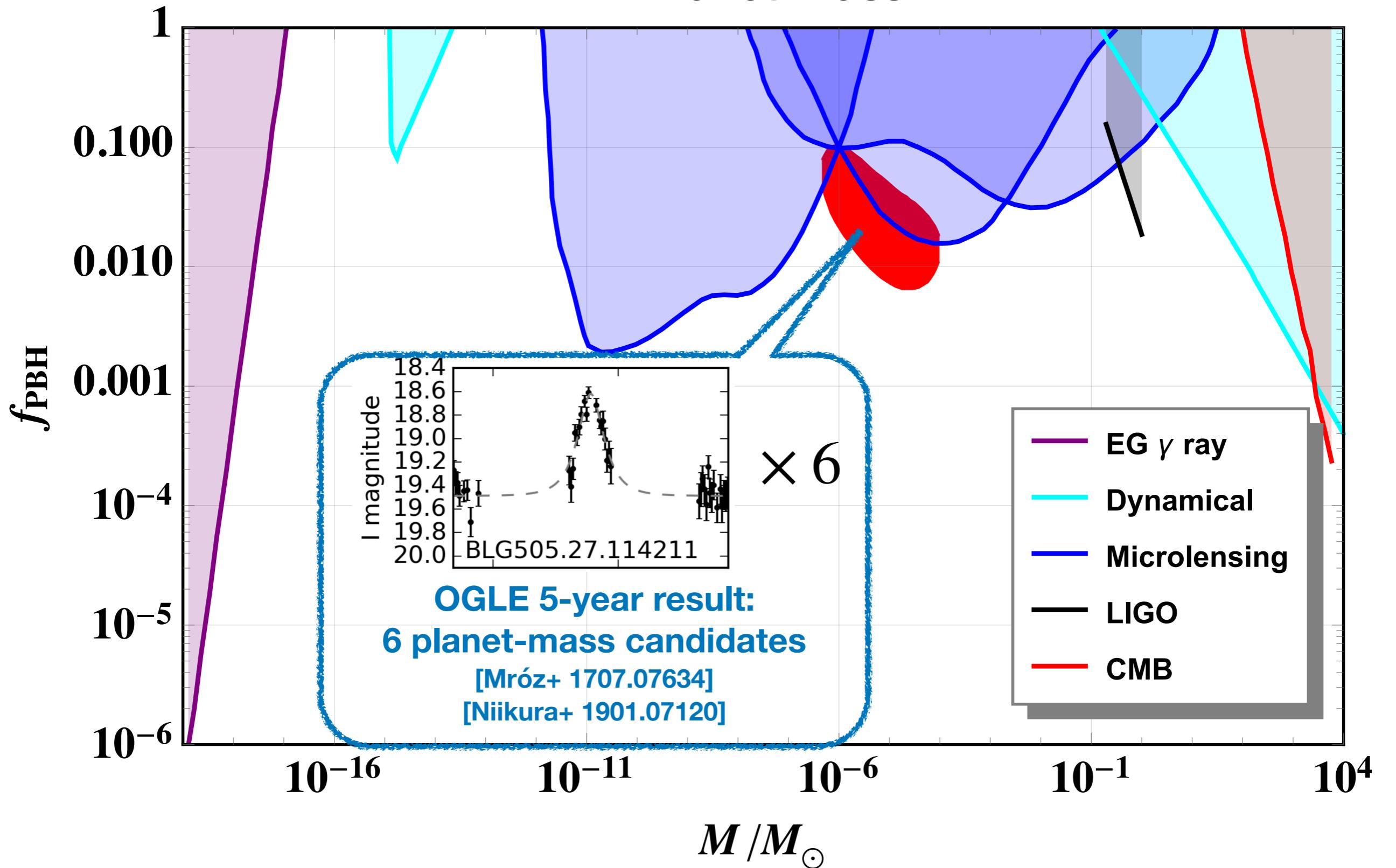




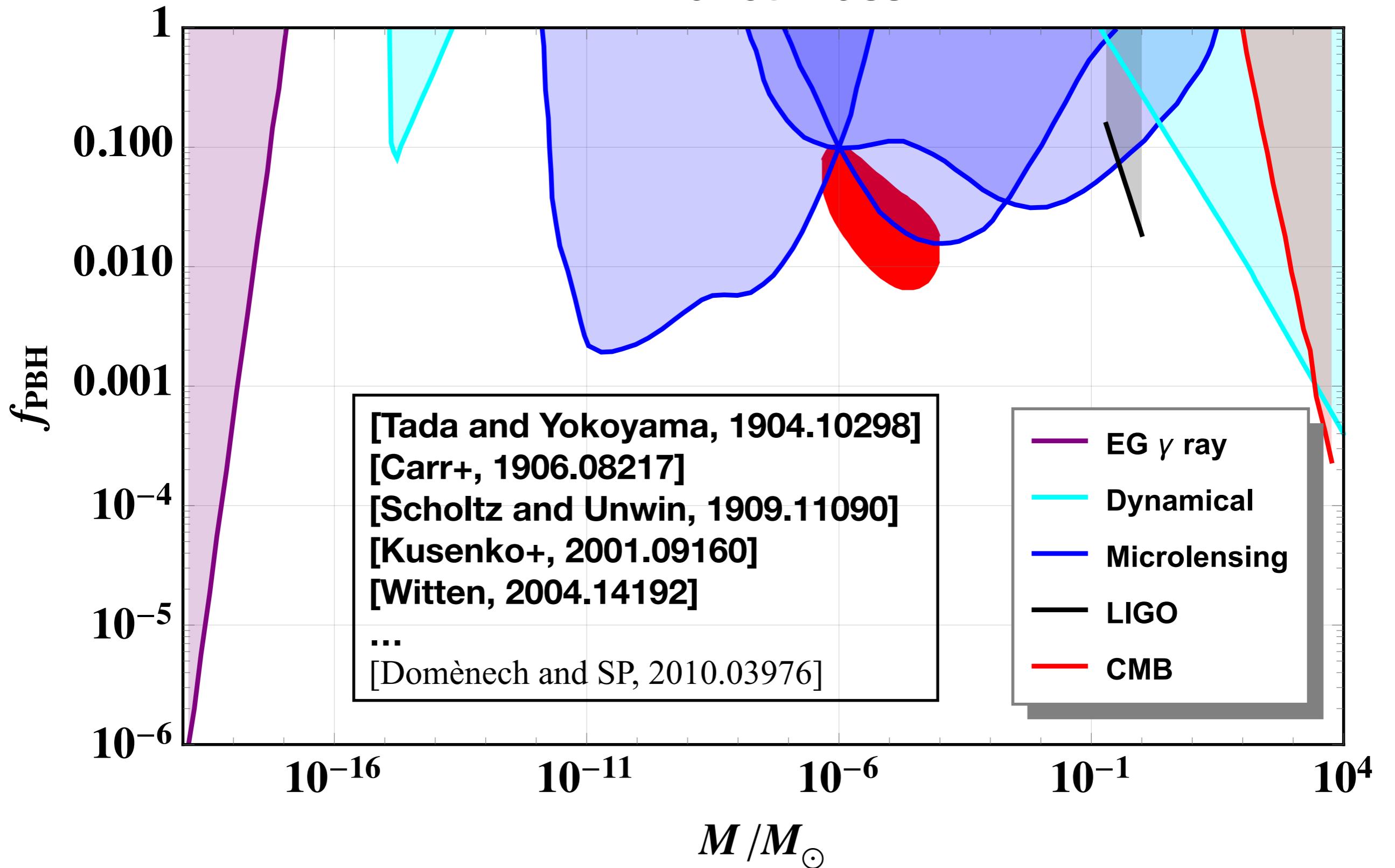
# LIGO PBH

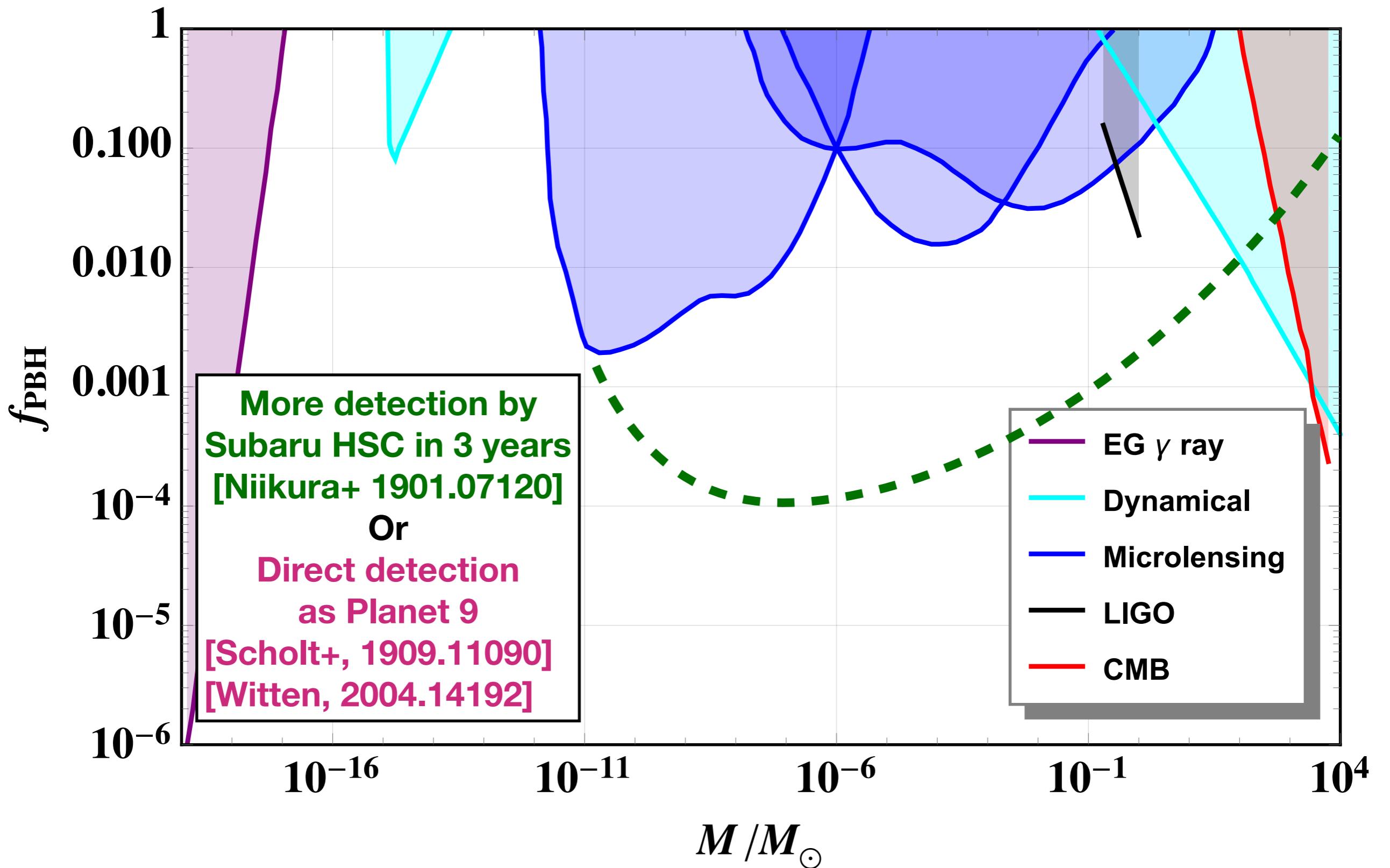


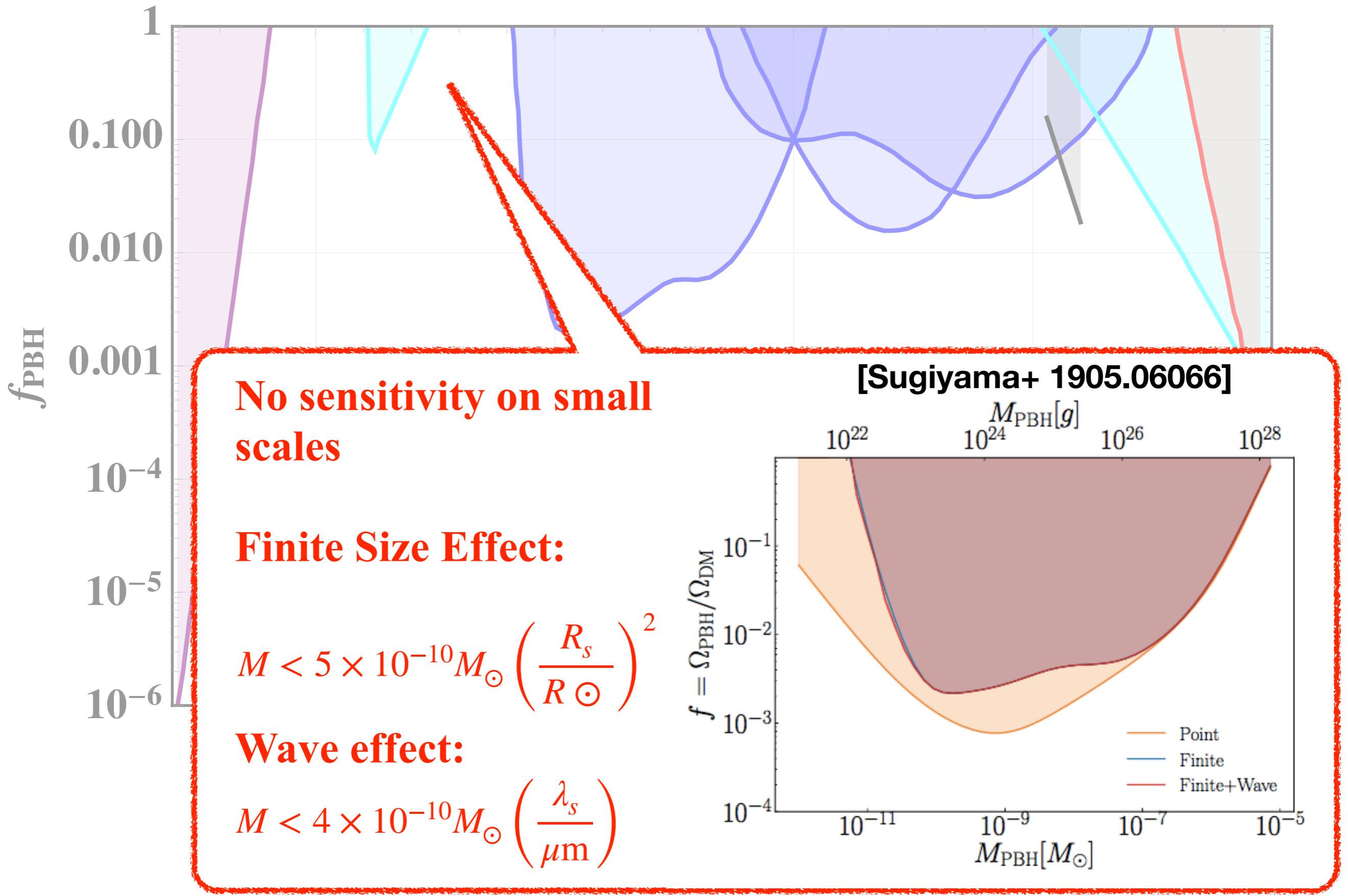
# Planet-mass PBH



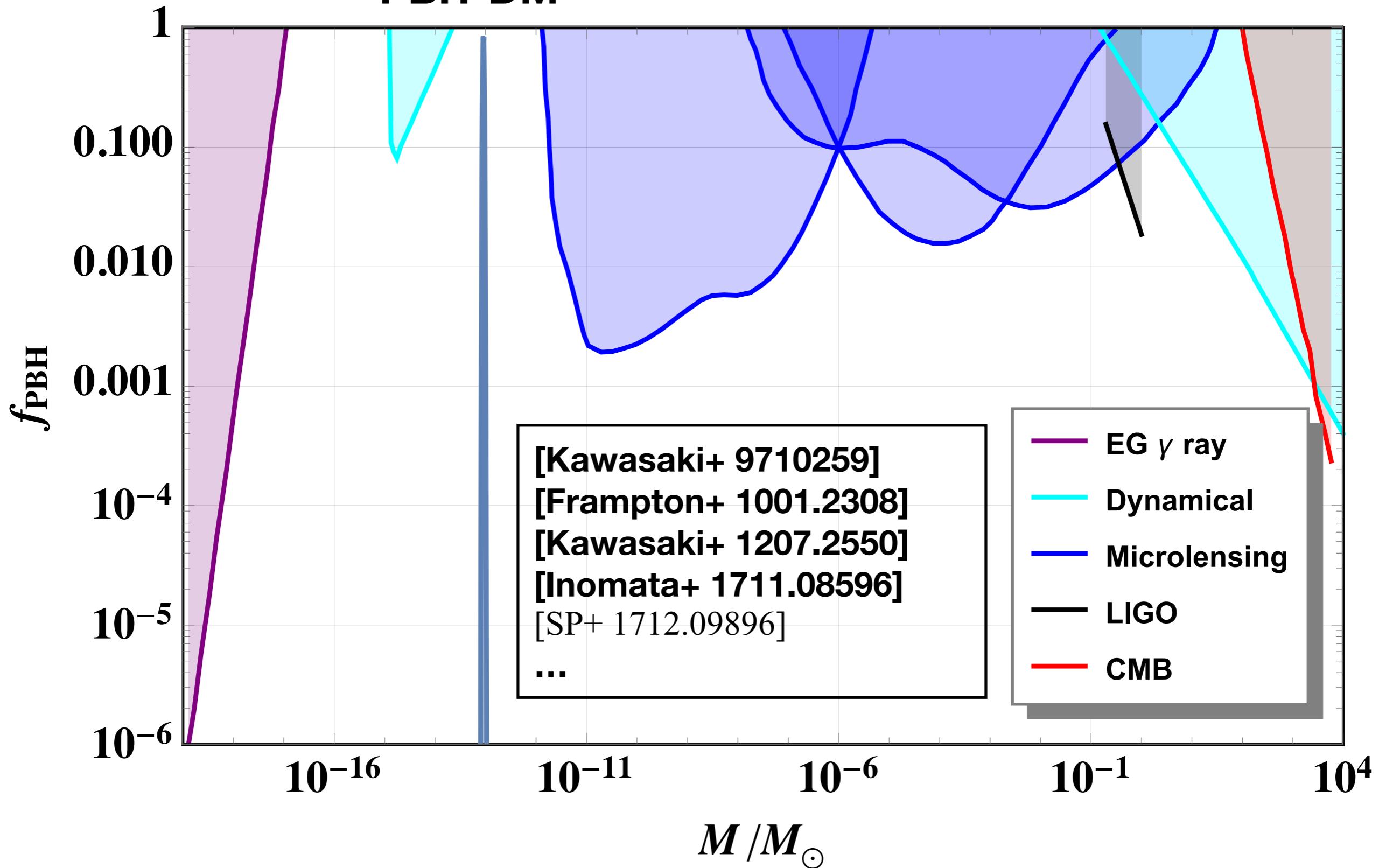
# Planet-mass PBH





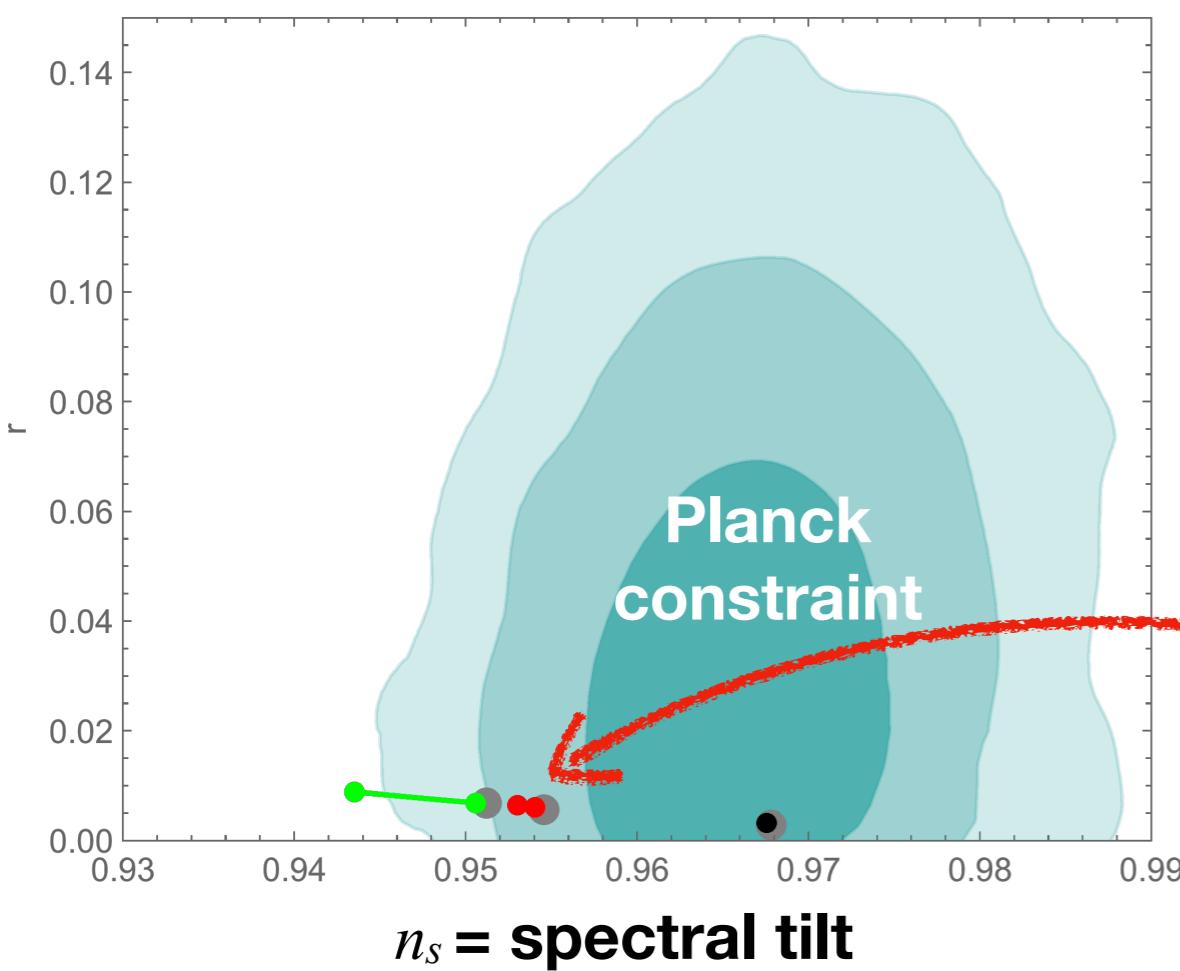


# PBH-DM

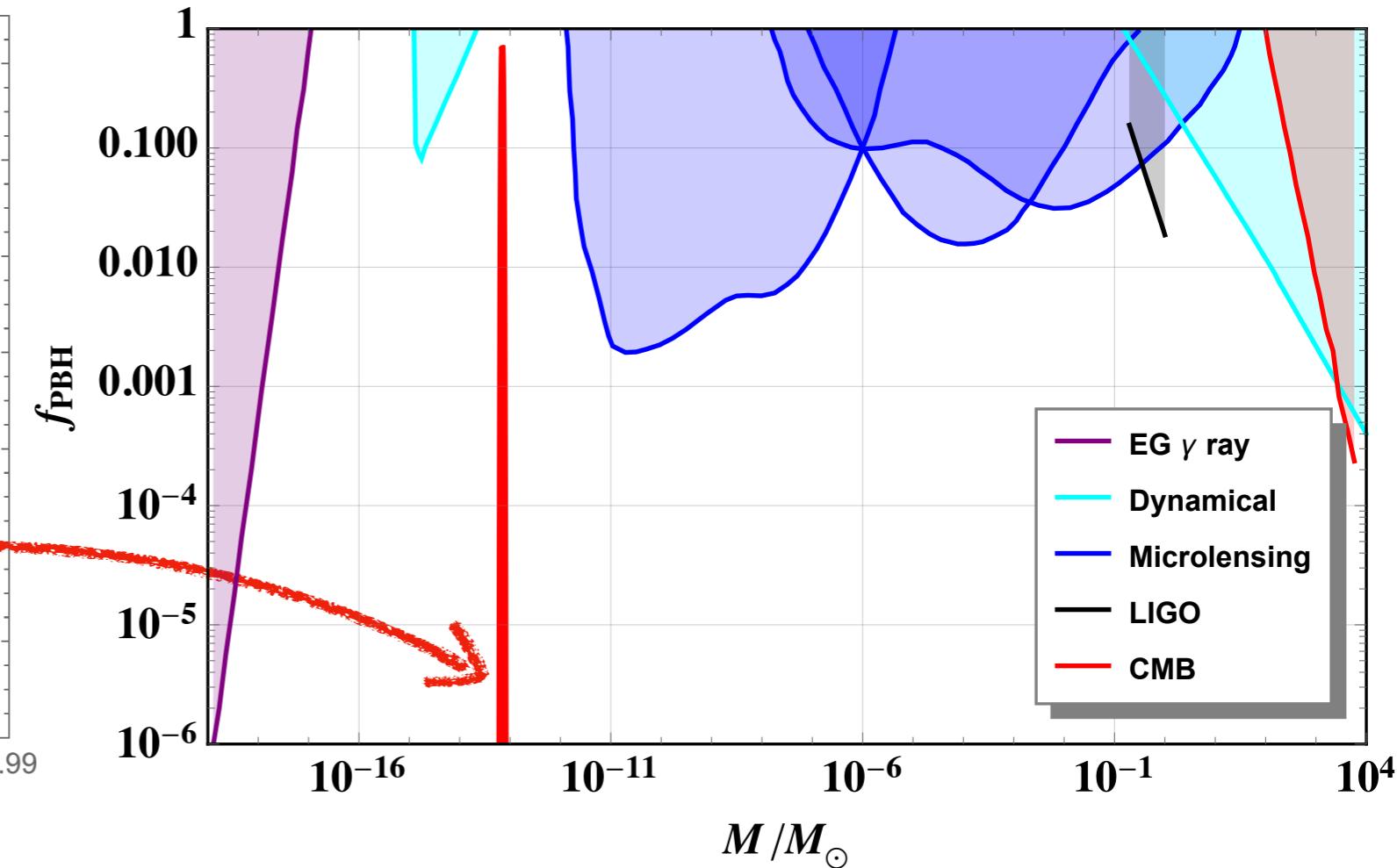


# PBH from $R^2$ -inflation

$$r = \frac{\text{tensor spectrum}}{\text{scalar spectrum}}$$



$$f_{\text{PBH}} = \frac{\text{PBH energy density}}{\text{CDM energy density}}$$

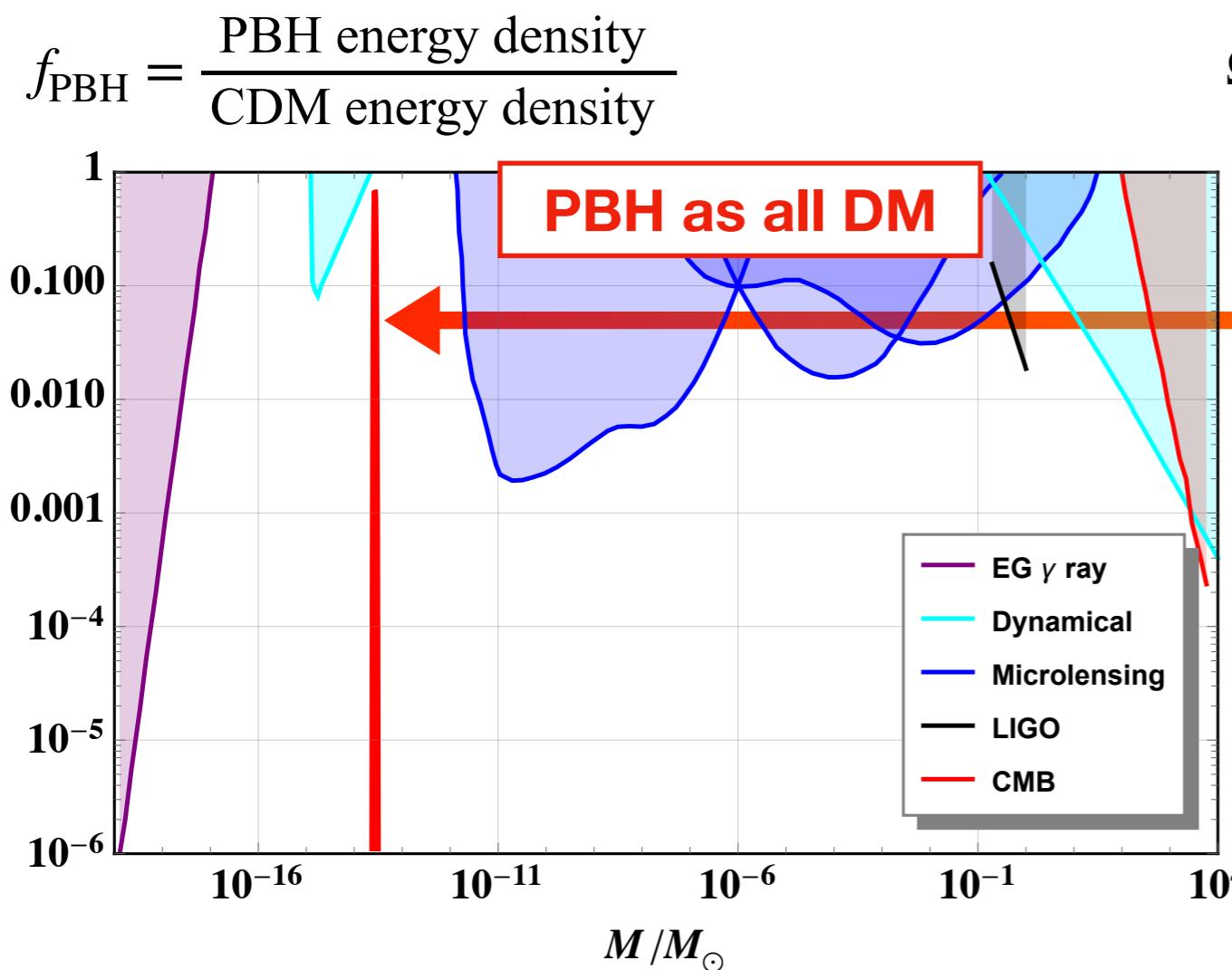


[SP, Zhang, Huang & Sasaki, JCAP1805, 042]

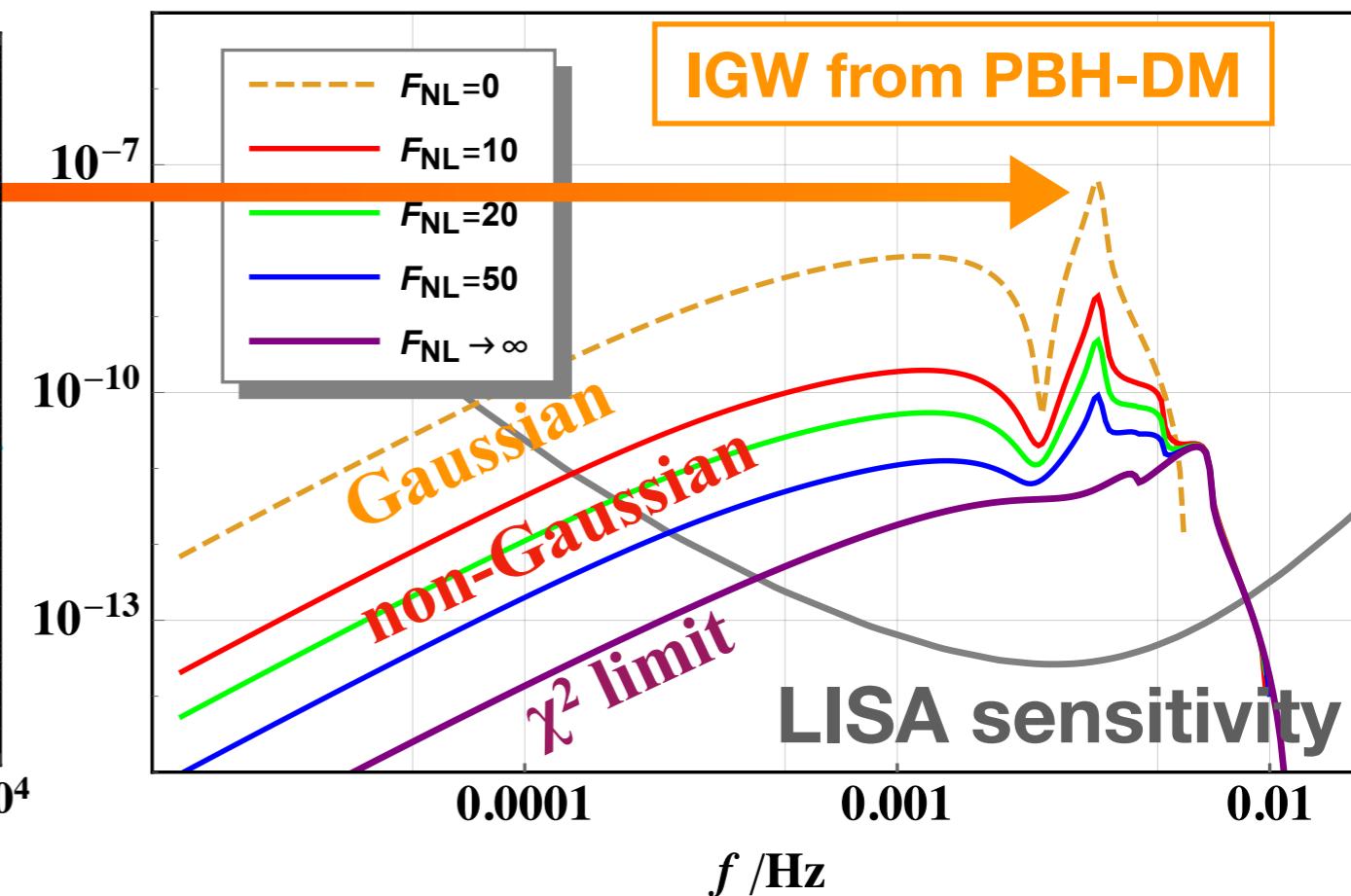
If PBH serves as all DM,  
the induced GWs must be  
detectable by LISA,  
independent of  $\mathcal{A}_{\mathcal{R}}$  or  $F_{NL}$ .

$$\mathcal{R} = \mathcal{R}_G + F_{NL}\mathcal{R}_G^2$$

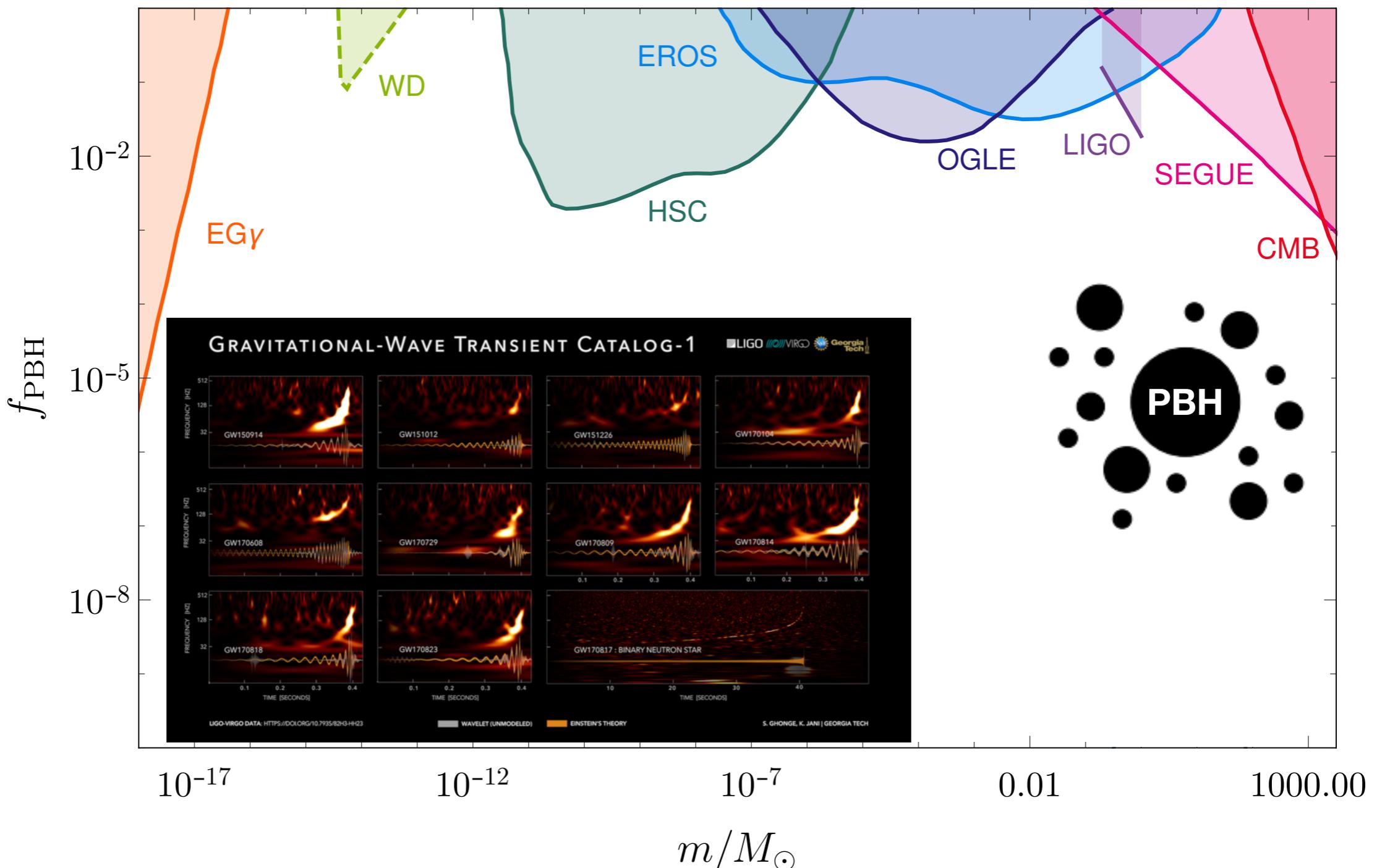
$$\Omega_{GW} \sim \mathcal{P}_{\mathcal{R},G}^2 + F_{NL}^2 \mathcal{P}_{\mathcal{R},G}^3 + F_{NL}^4 \mathcal{P}_{\mathcal{R},G}^4$$



$$\Omega_{GW} = \frac{\text{GW energy density}}{\text{critical energy density}}$$

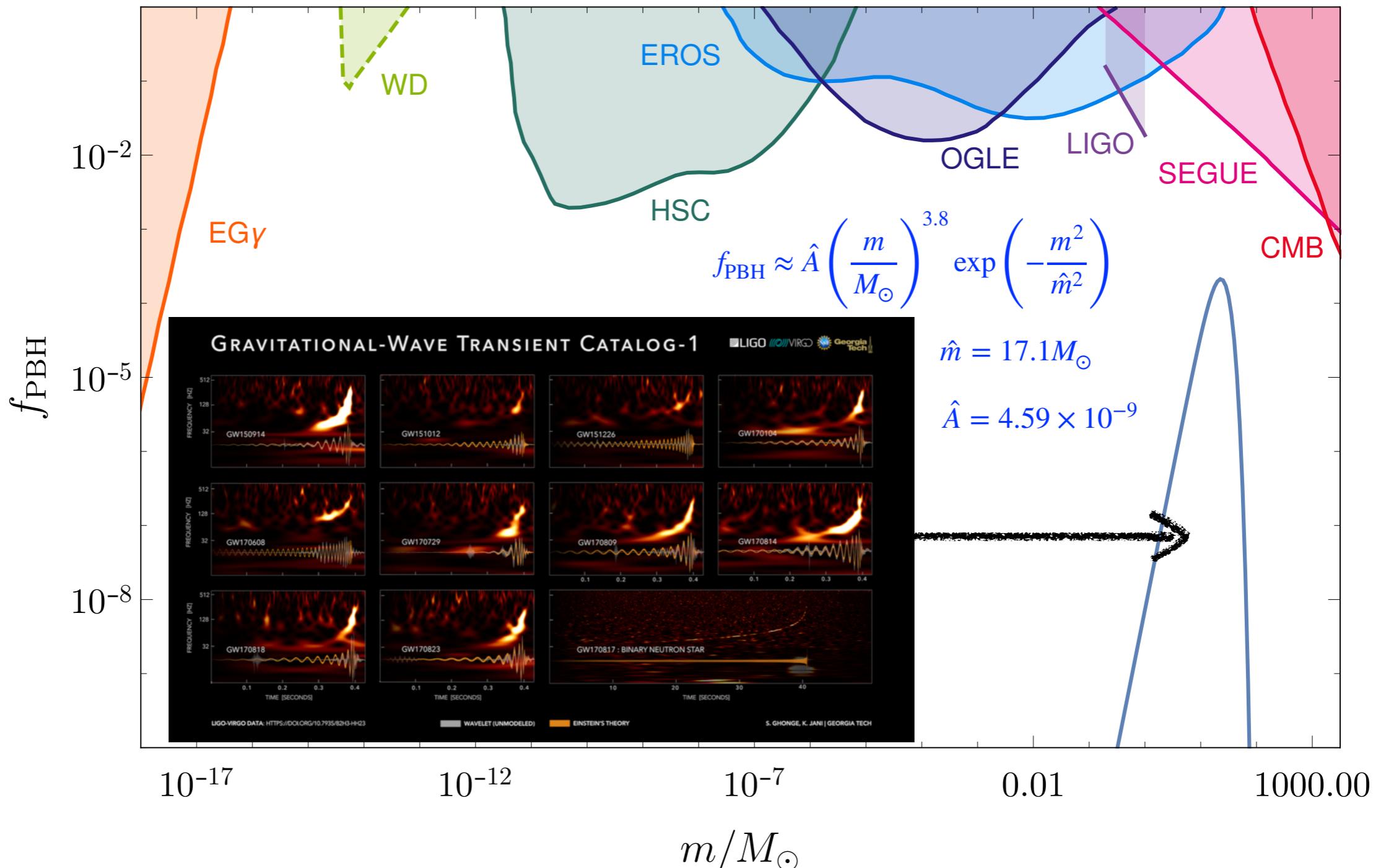


# PBHs as LIGO events



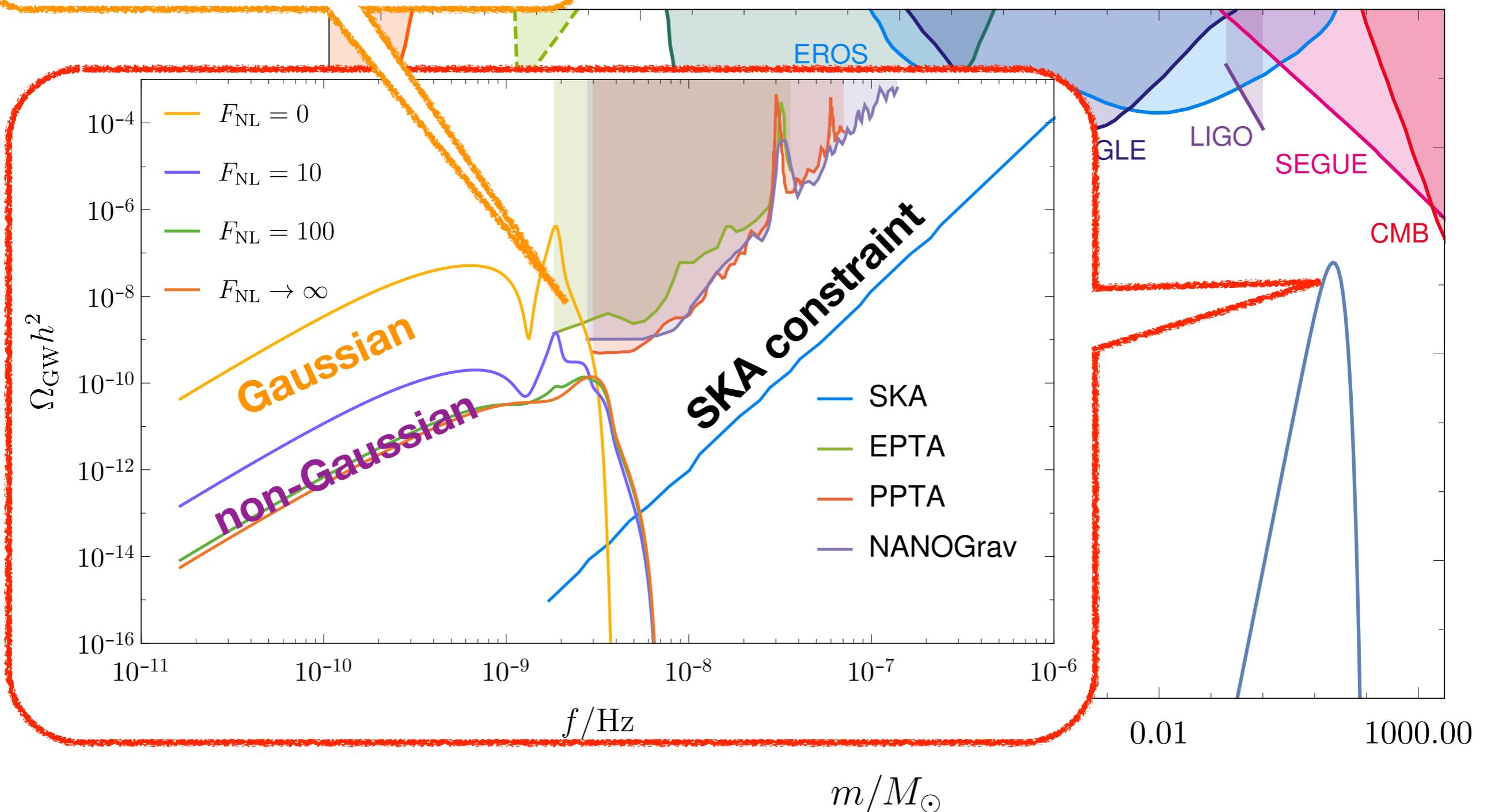
[Cai, SP, Wang, and Yang, JCAP1910, 059]

# PBHs as LIGO events



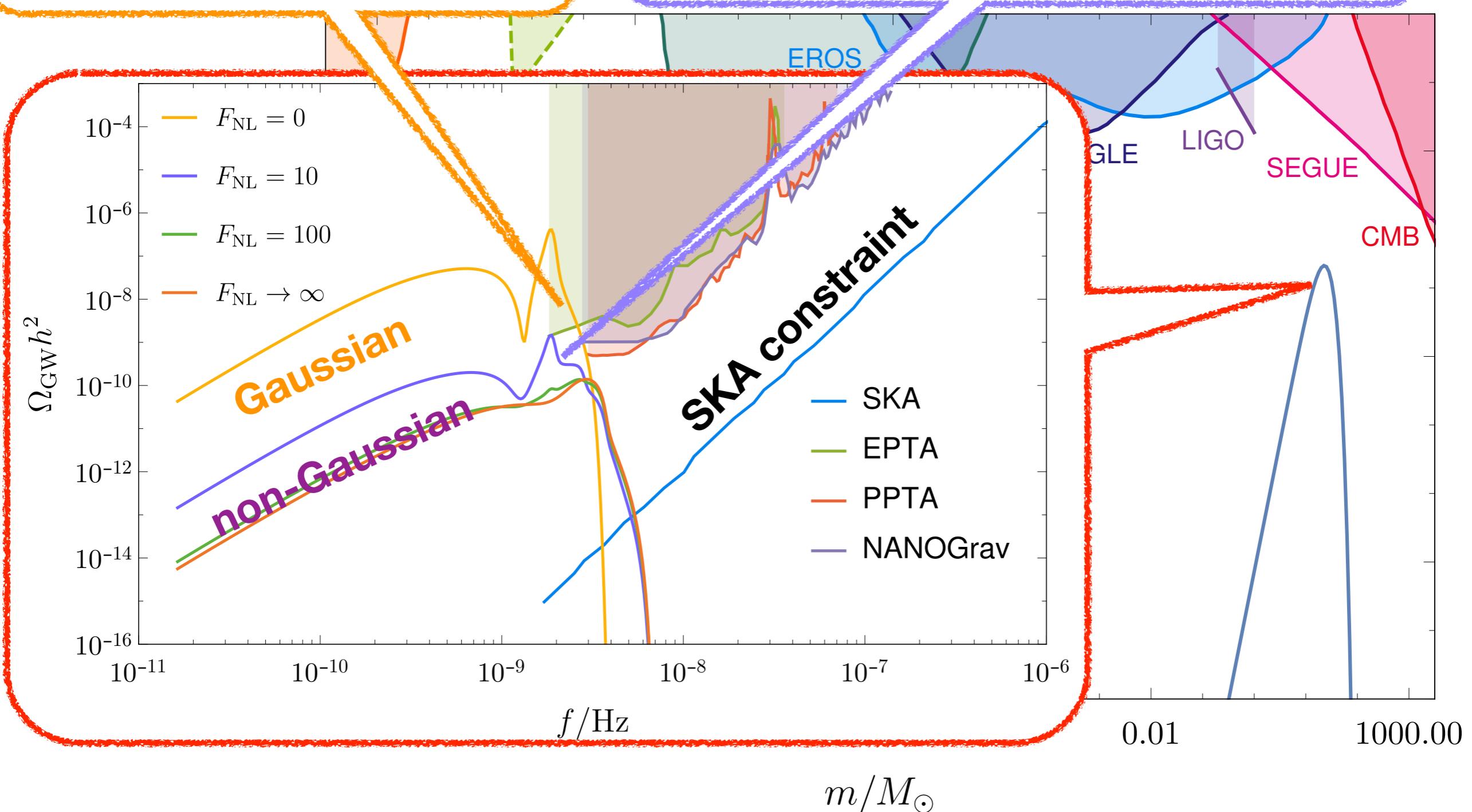
[Cai, SP, Wang, and Yang, JCAP1910, 059]

The IGW from LIGO-PBH  
is excluded by PTA...



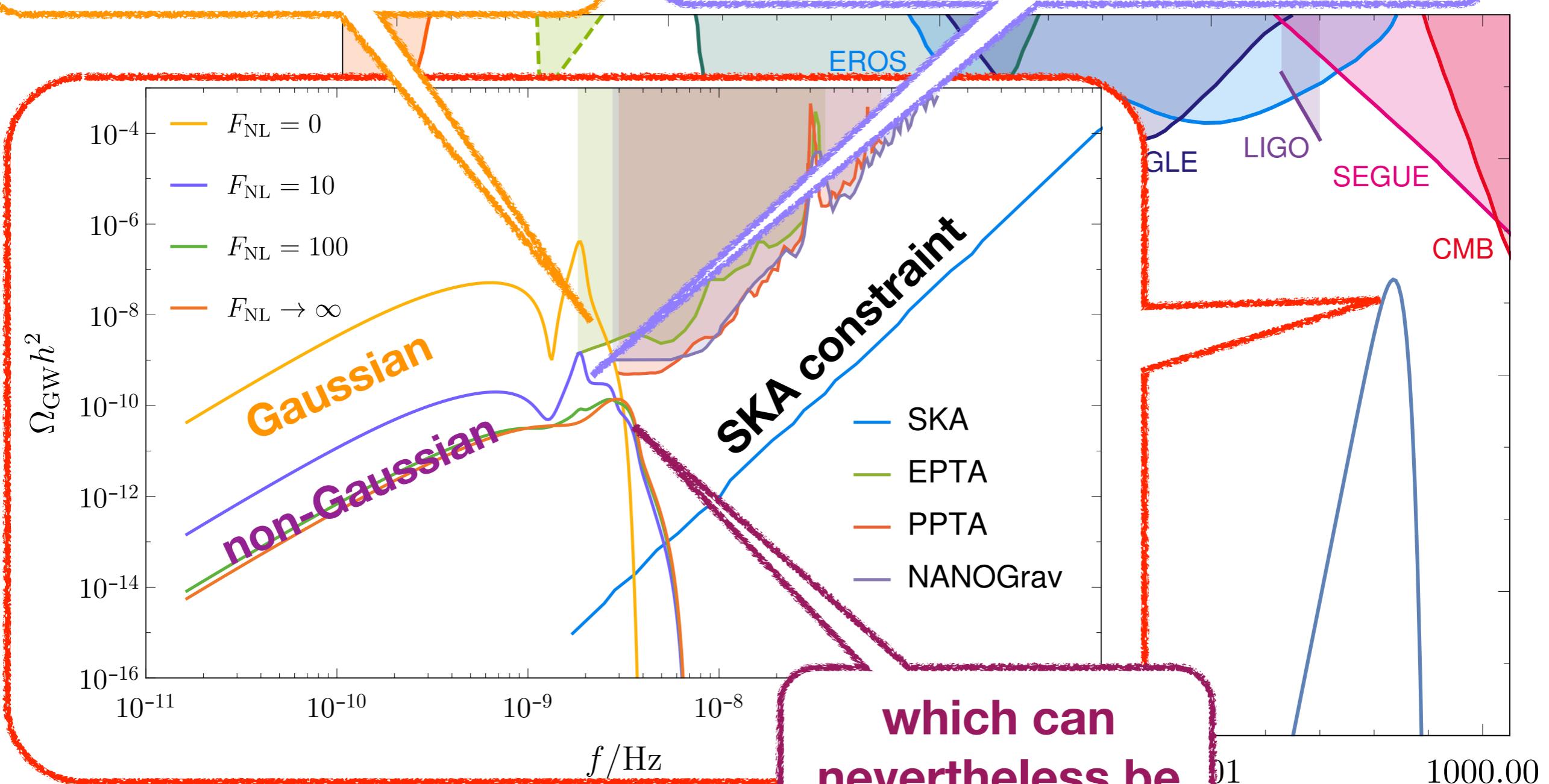
The IGW from LIGO-PBH  
is excluded by PTA...

...unless there is non-Gaussianity,



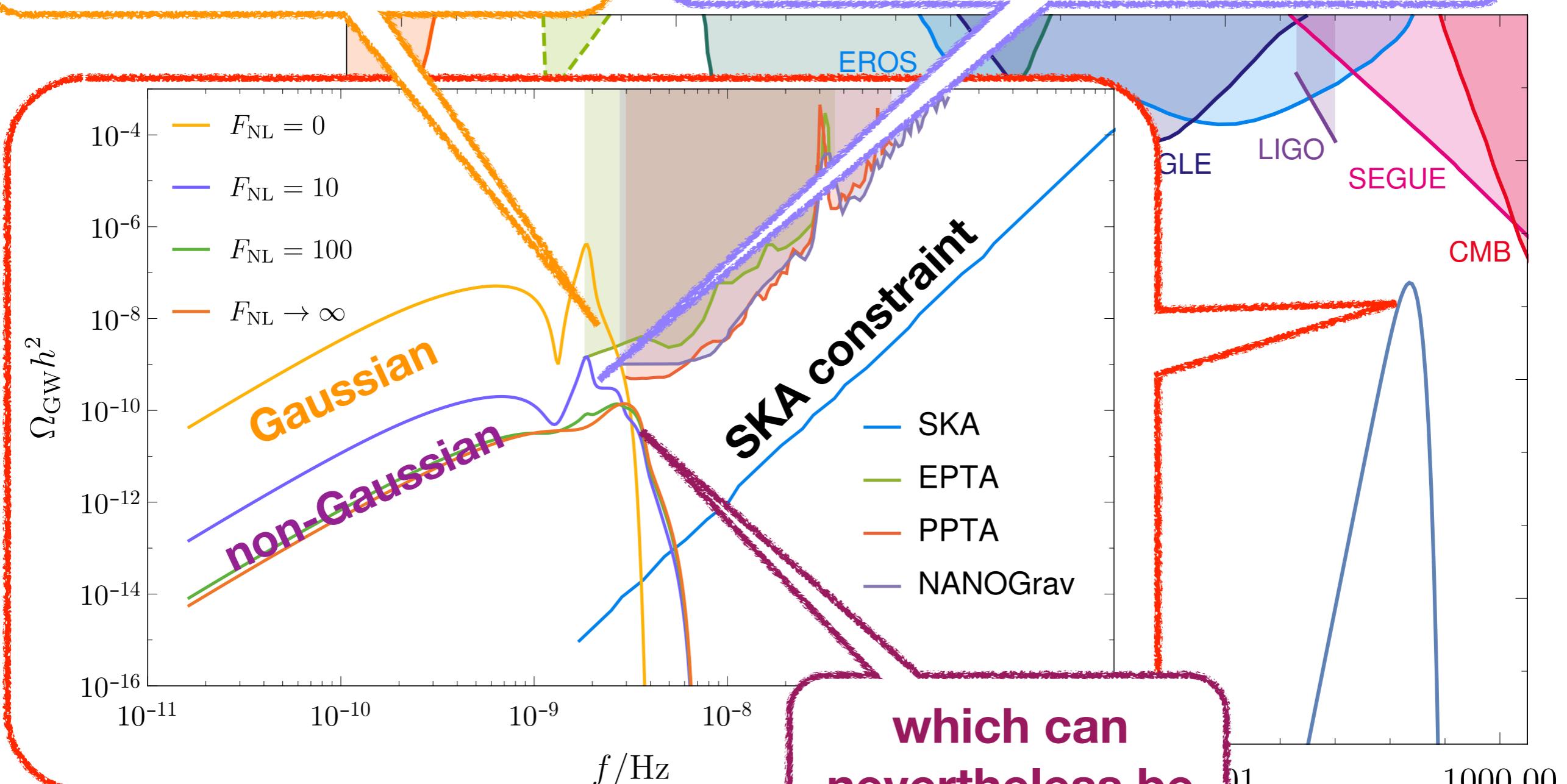
The IGW from LIGO-PBH  
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The IGW from LIGO-PBH  
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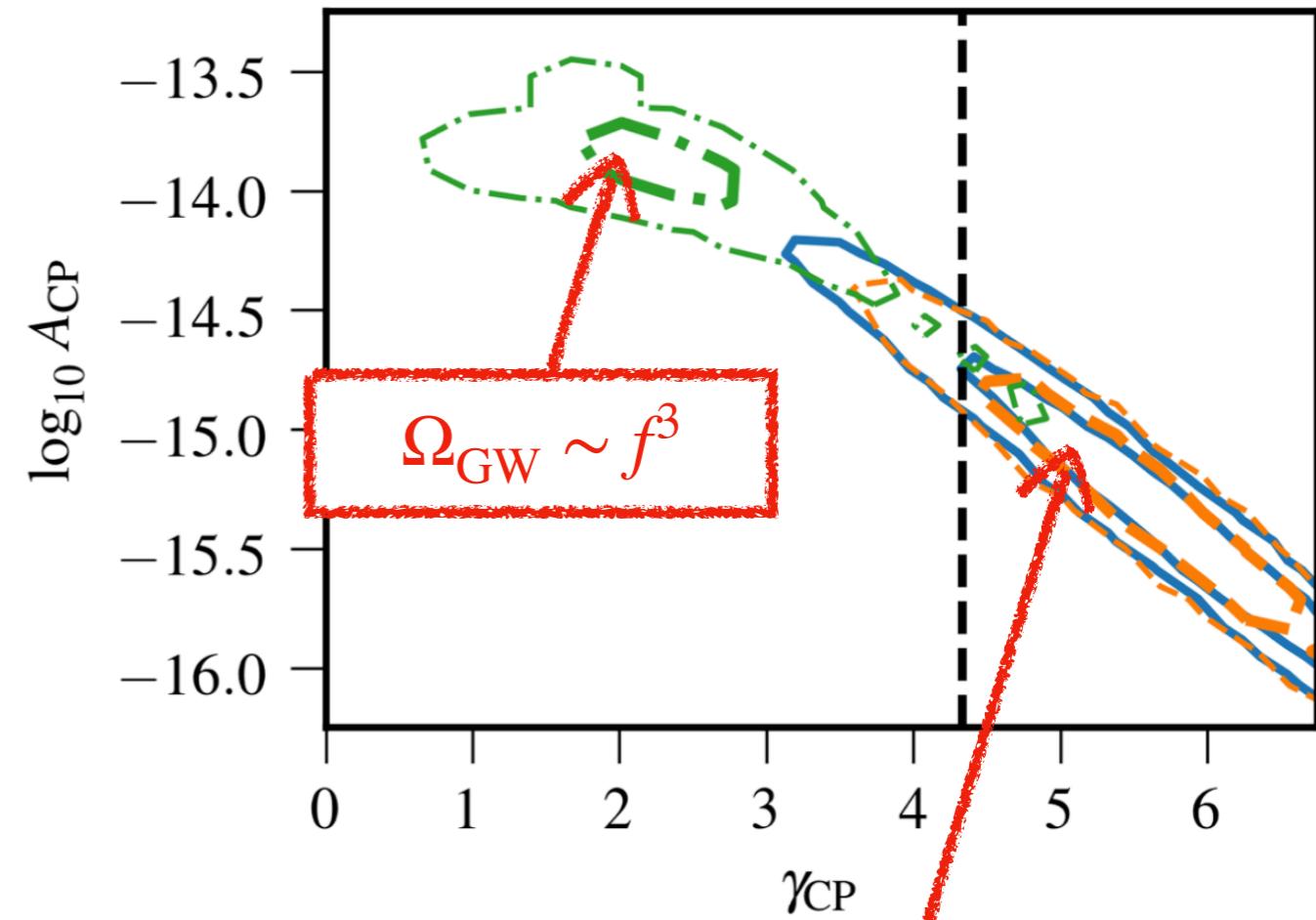
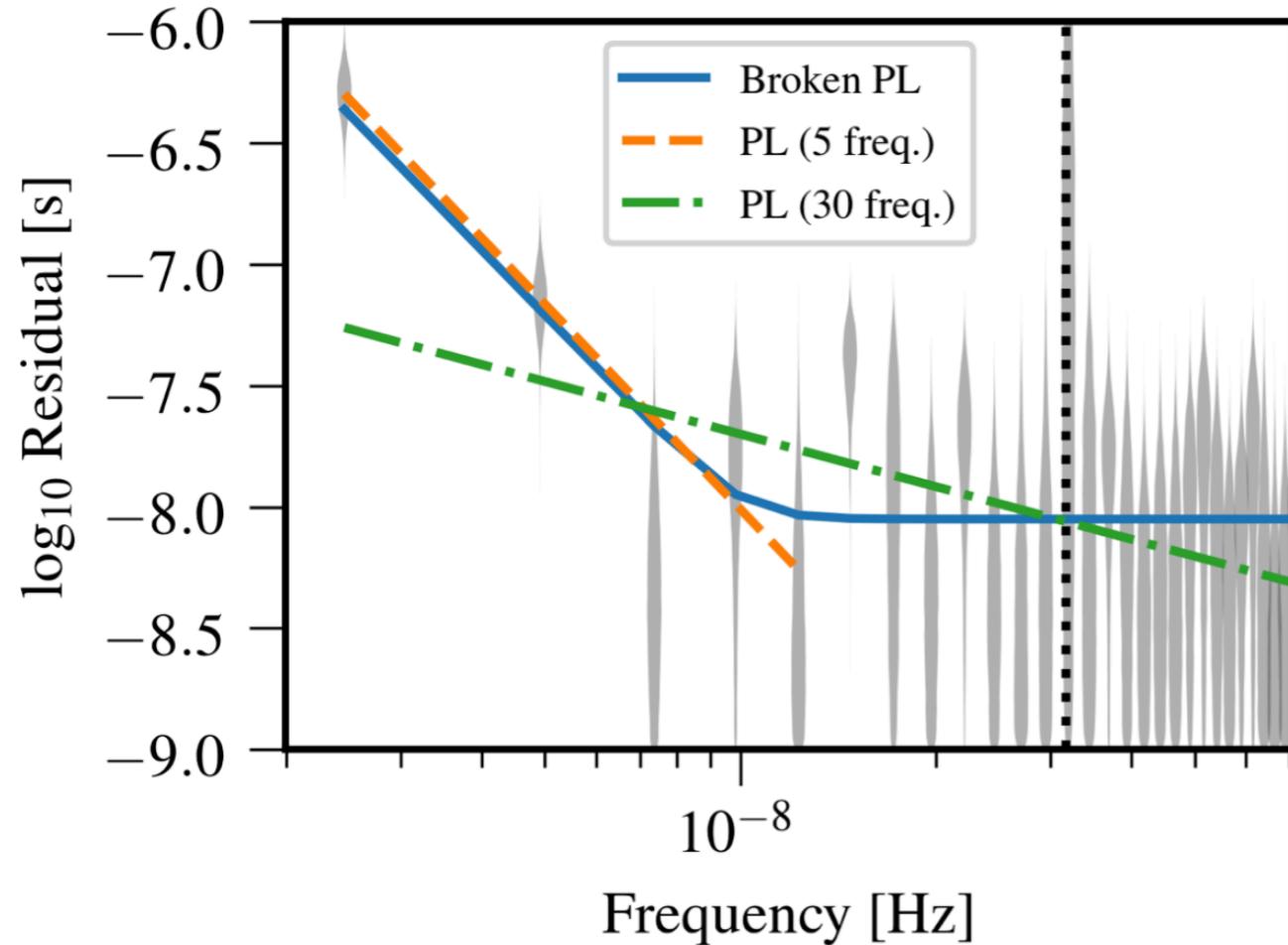
But wait...

[Cai, SP, Wang, and Yang, JCAP1910, 059]

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- Some properties for IGW: scaling and peak
- NANOGrav hints on planet-mass PBHs
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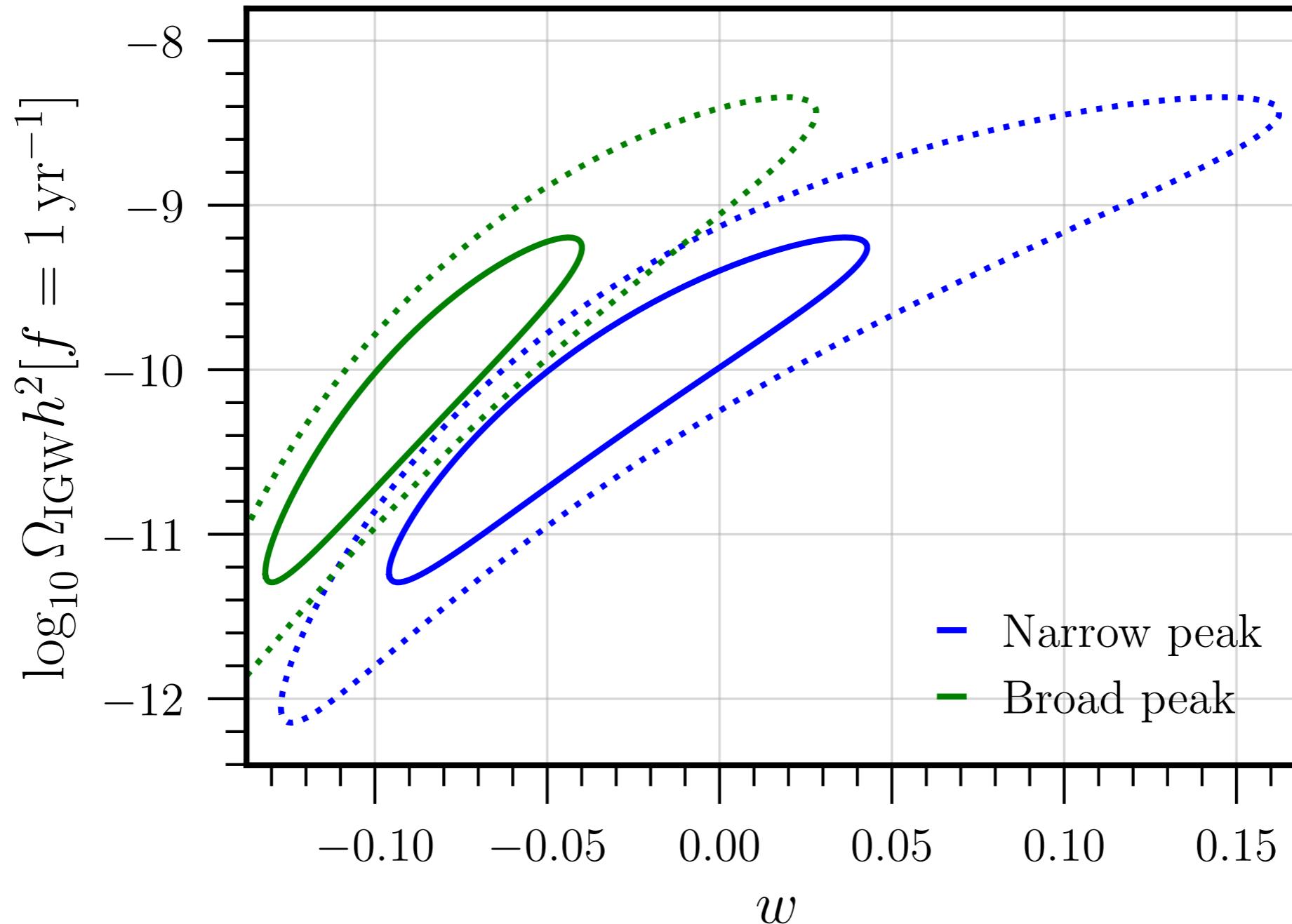
# PTA discovered SGWB?



$$\Omega_{\text{GW}} = \frac{2\pi^2 f_{\text{yr}}^2}{3H_0^2} A_{\text{SGWB}}^2 \left( \frac{f}{f_{\text{yr}}} \right)^{5-\gamma}$$

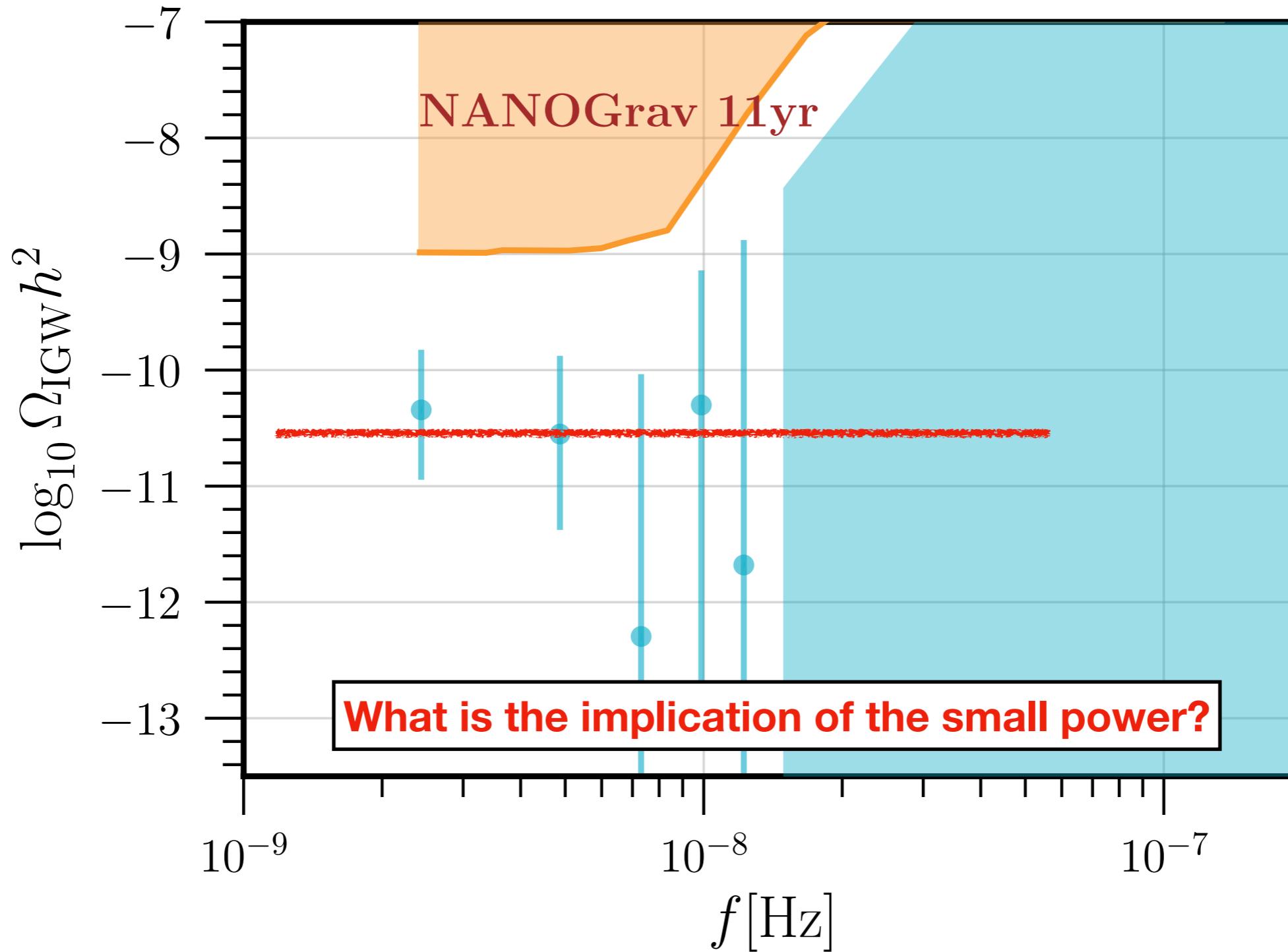
$$f_{\text{yr}} = 3.17 \times 10^{-8} \text{ Hz}.$$

# PTA discovered SGWB?

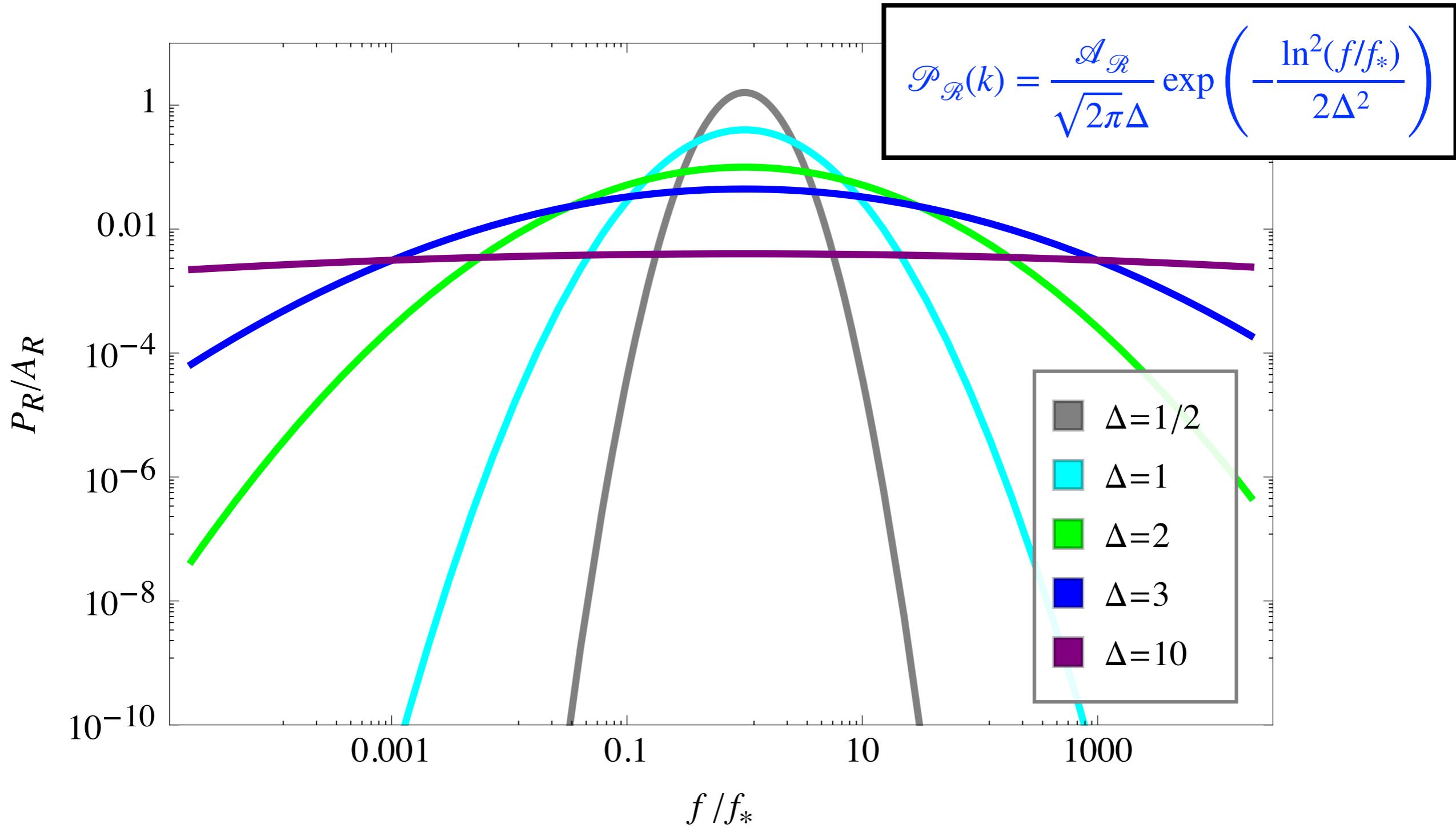


[Domènech and SP, 2010.03976]

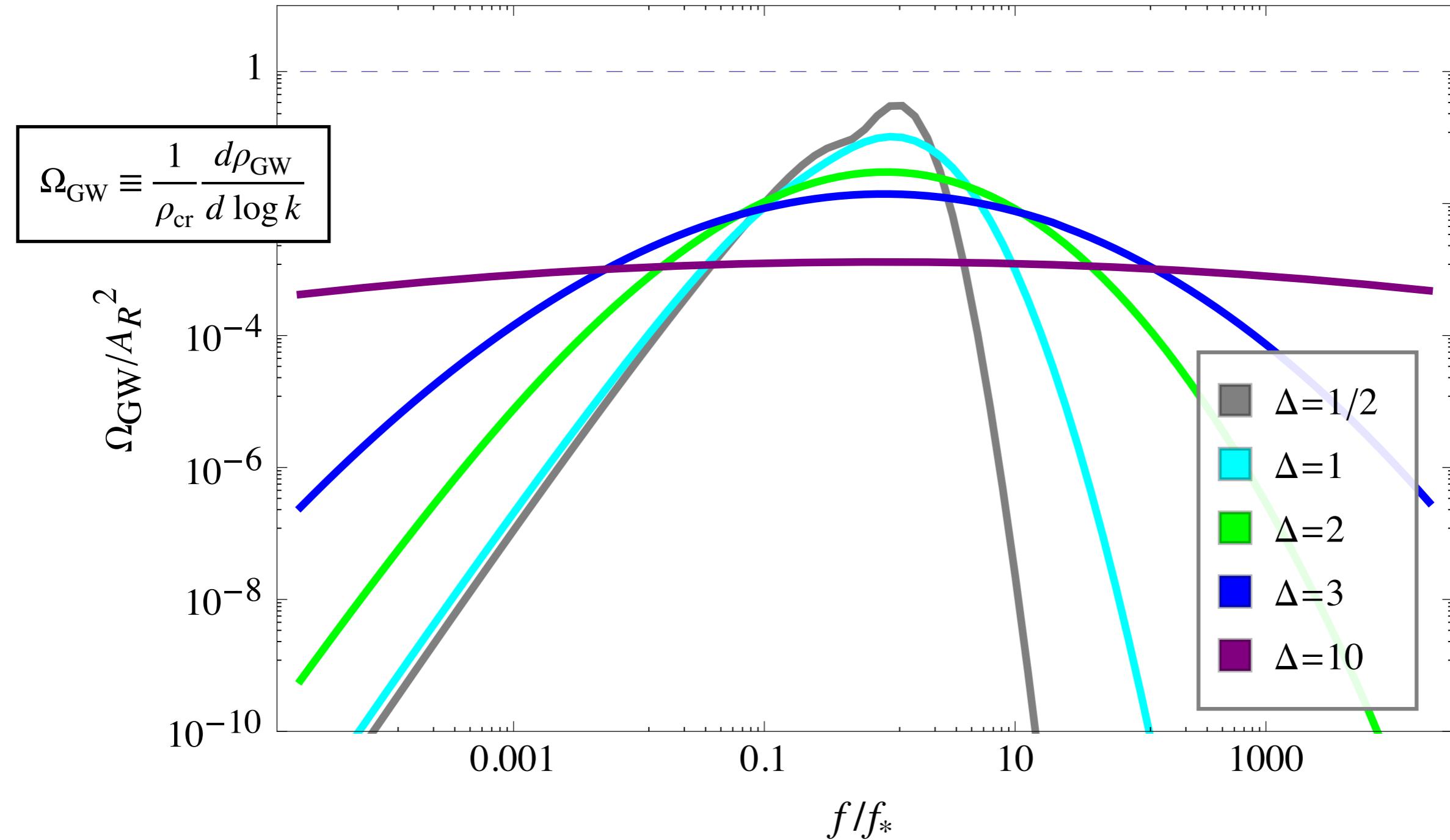
# PTA discovered SGWB?



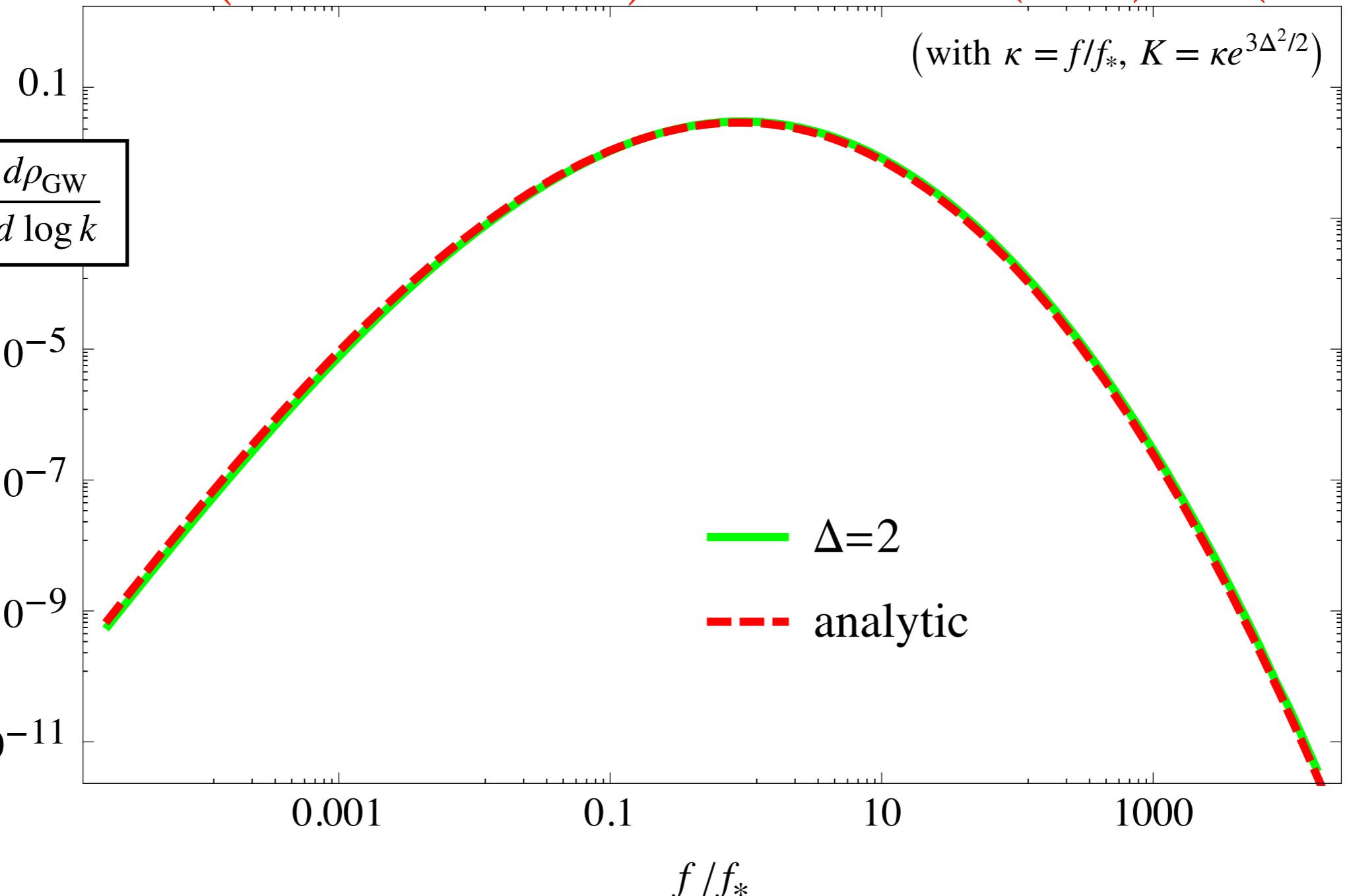
# Broad Peak of $\mathcal{P}_{\mathcal{R}}$



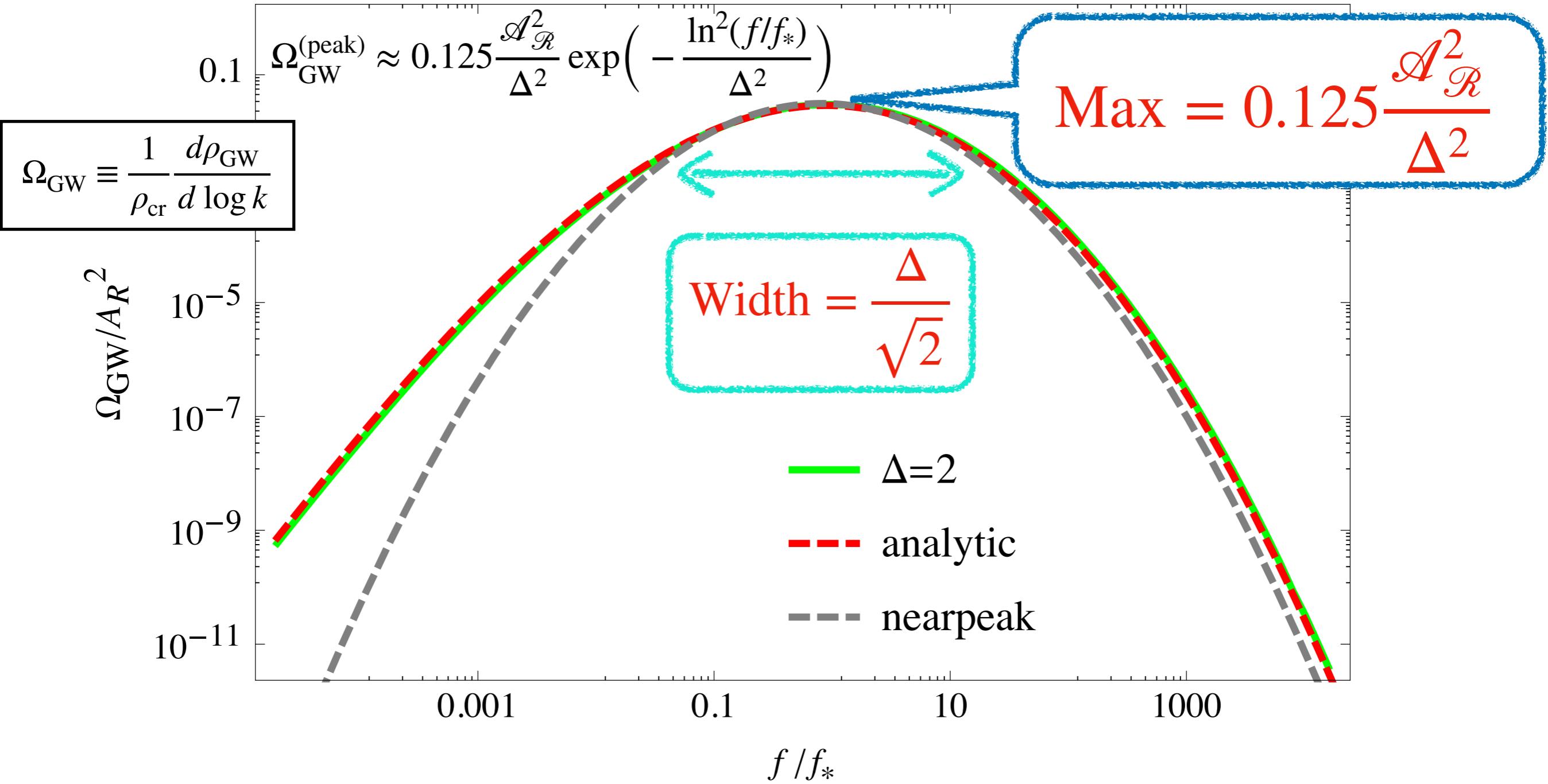
# IGW from Broad Peak



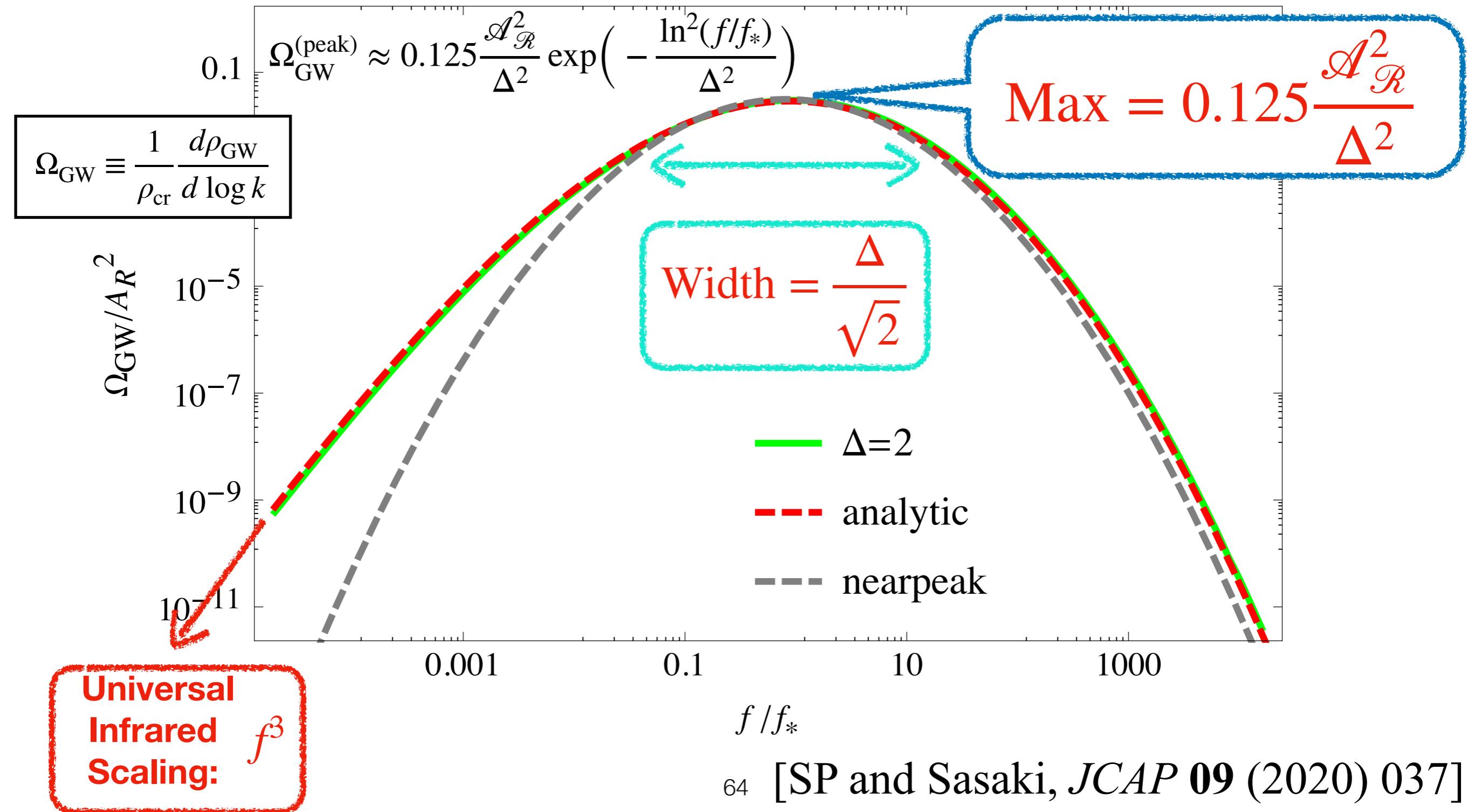
$$\frac{\Omega_{\text{GW}}}{\mathcal{A}_{\mathcal{R}}^2} \approx \frac{4}{5\sqrt{\pi}} \kappa^3 \frac{e^{\frac{9\Delta^2}{4}}}{\Delta} \left[ \left( \ln^2 K + \frac{\Delta^2}{2} \right) \text{erfc} \left( \frac{\ln K + \frac{1}{2} \ln \frac{3}{2}}{\Delta} \right) - \frac{\Delta}{\sqrt{\pi}} \exp \left( -\frac{\left( \ln K + \frac{1}{2} \ln \frac{3}{2} \right)^2}{\Delta^2} \right) \left( \ln K - \frac{1}{2} \ln \frac{3}{2} \right) \right] \\ + \frac{0.0659}{\Delta^2} \kappa^2 e^{\Delta^2} \exp \left( -\frac{\left( \ln \kappa + \Delta^2 - \frac{1}{2} \ln \frac{4}{3} \right)^2}{\Delta^2} \right) + \frac{1}{3} \sqrt{\frac{2}{\pi}} \kappa^{-4} \frac{e^{8\Delta^2}}{\Delta} \exp \left( -\frac{\ln^2 \kappa}{2\Delta^2} \right) \text{erfc} \left( \frac{4\Delta^2 - \ln(\kappa/4)}{\sqrt{2}\Delta} \right)$$



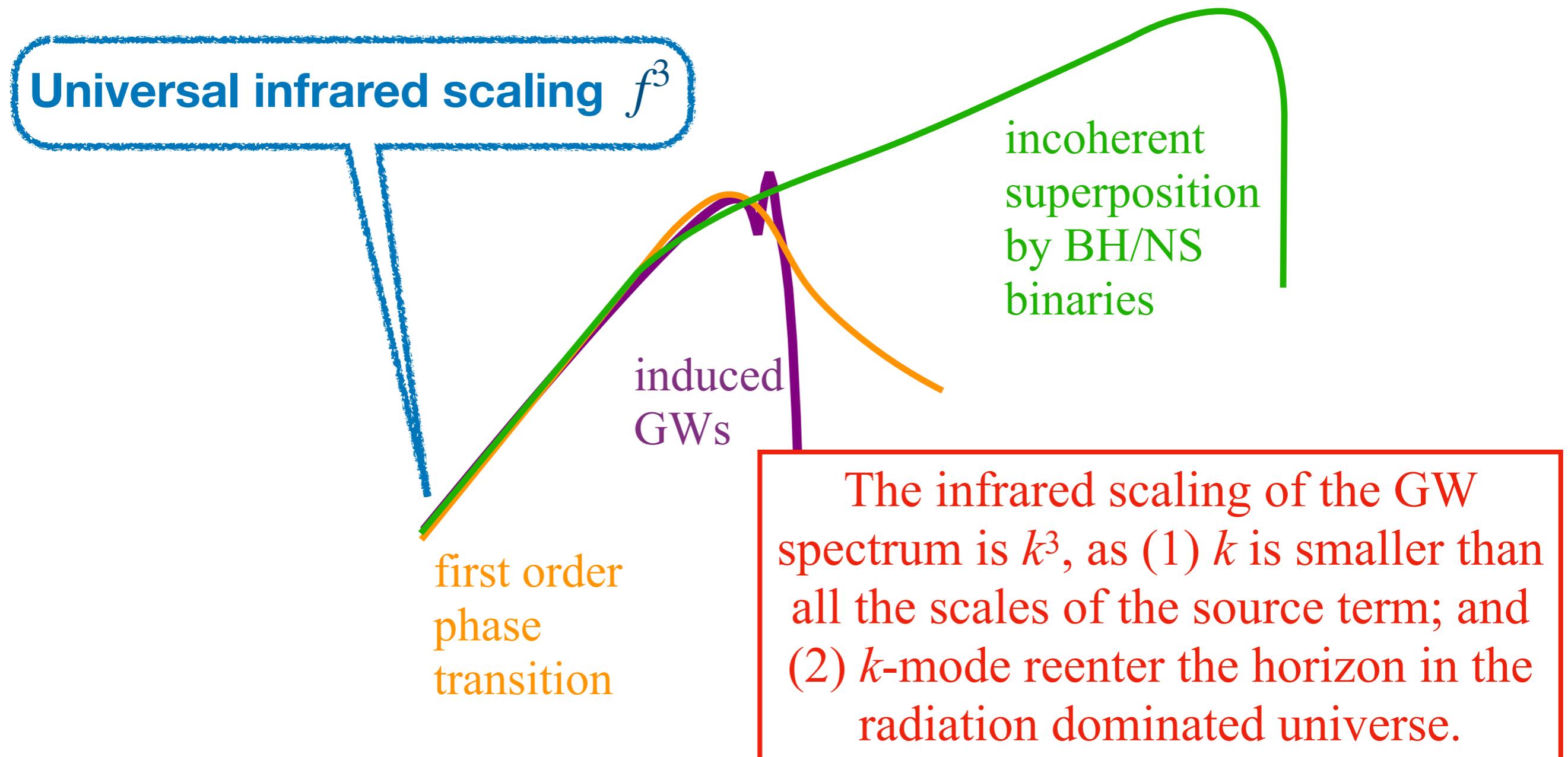
# IGW from Broad Peak



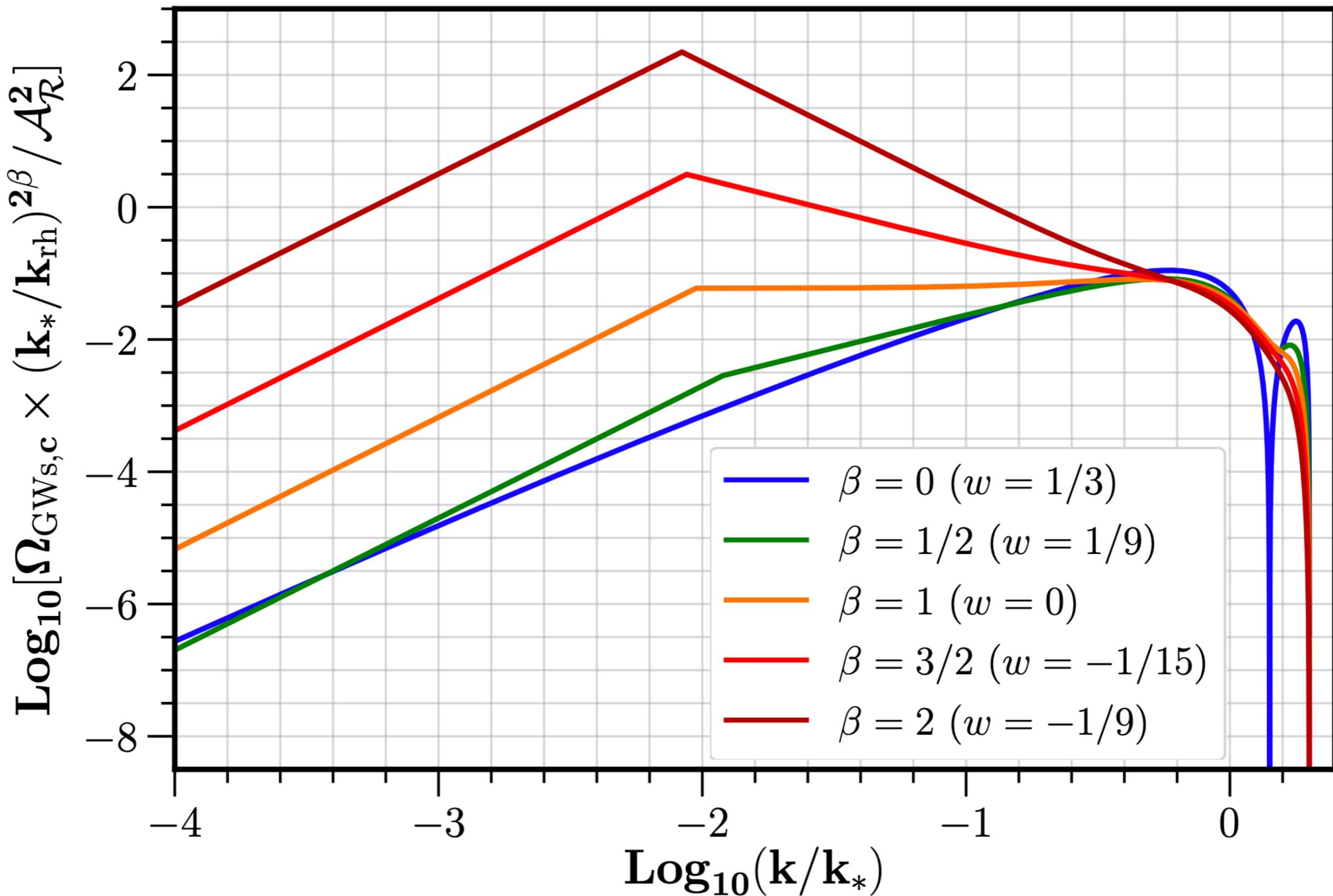
# Infrared Scaling



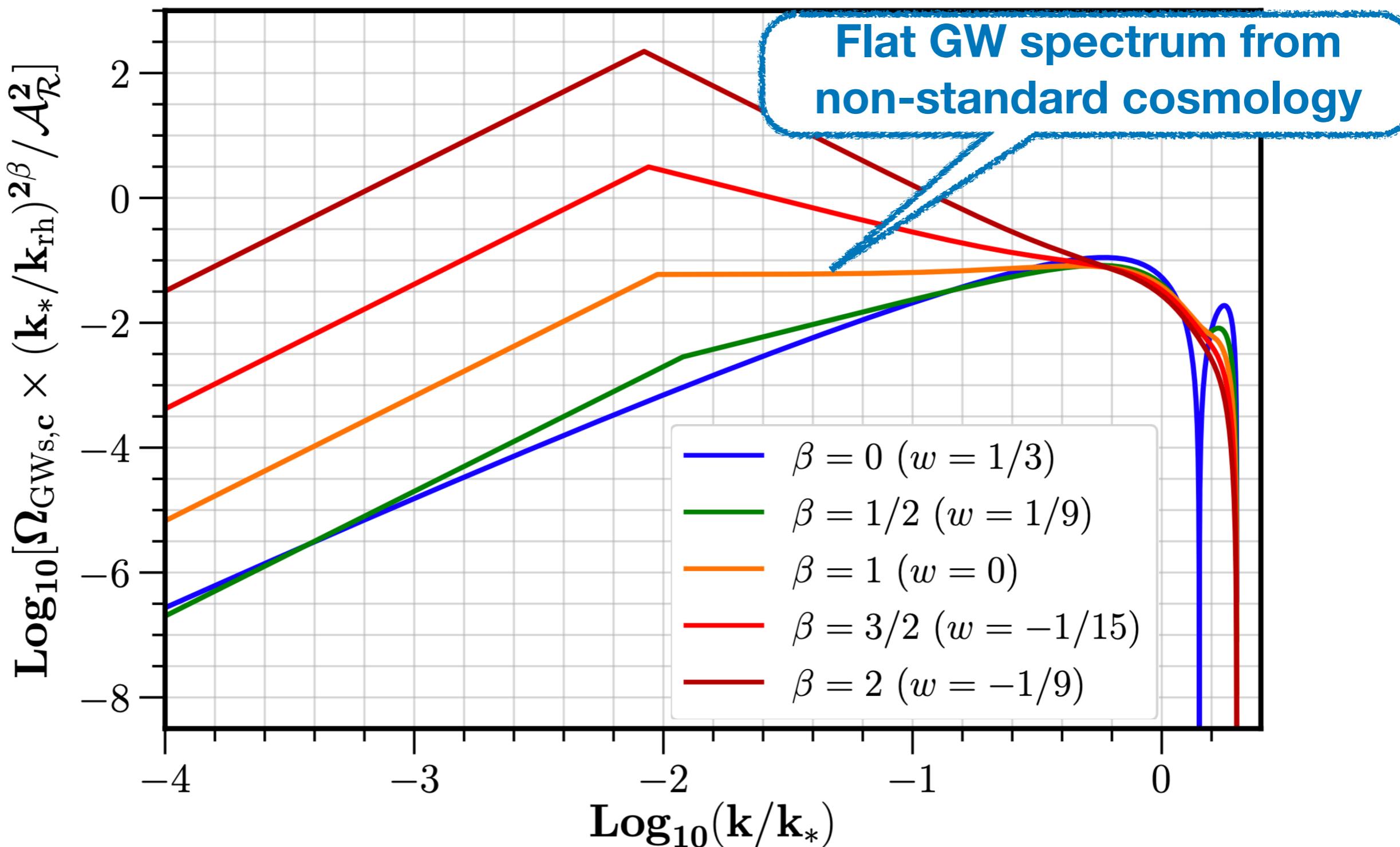
# Infrared Scaling of SGWB



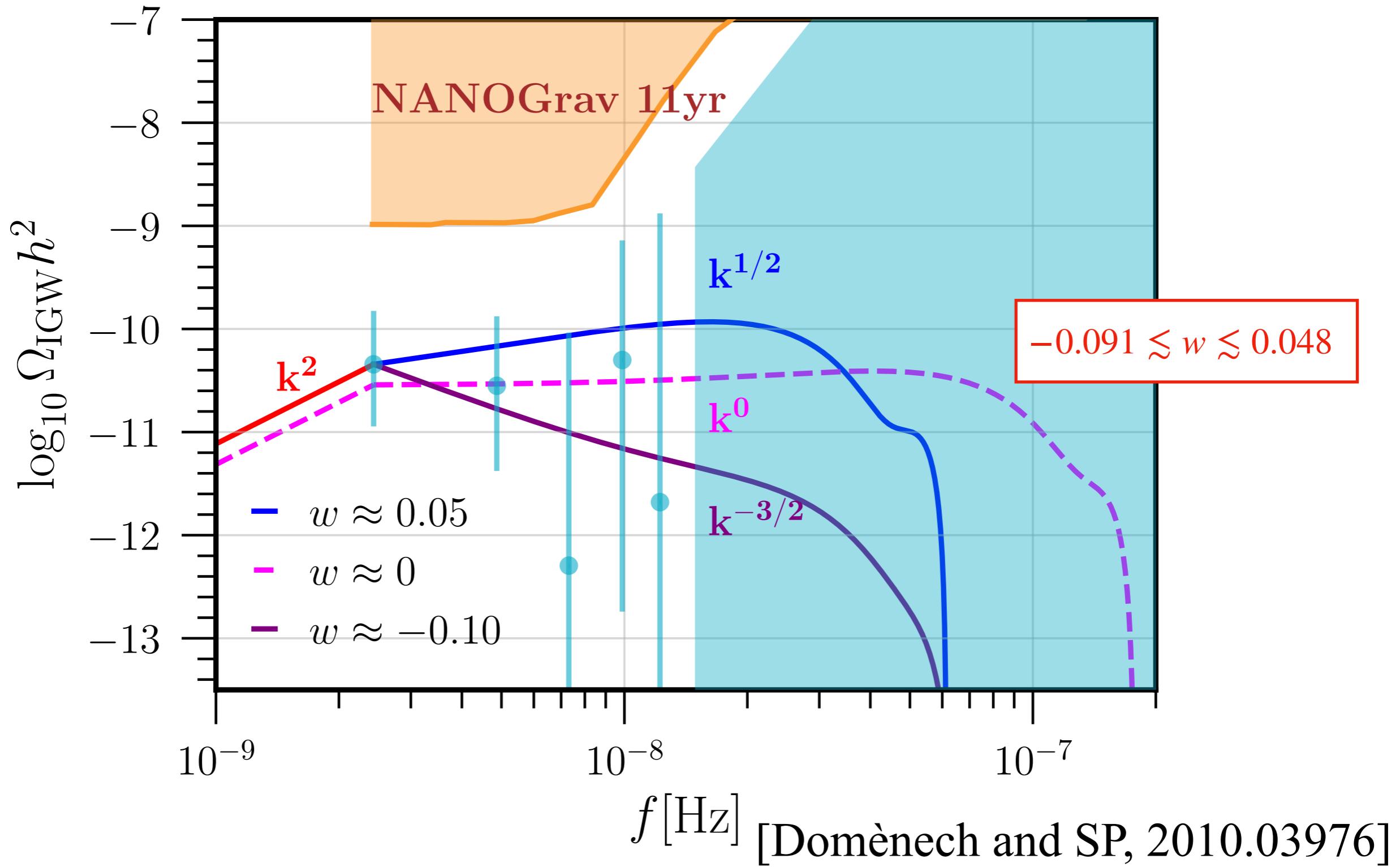
# Infrared Scaling of SGWB



# Infrared Scaling of SGWB

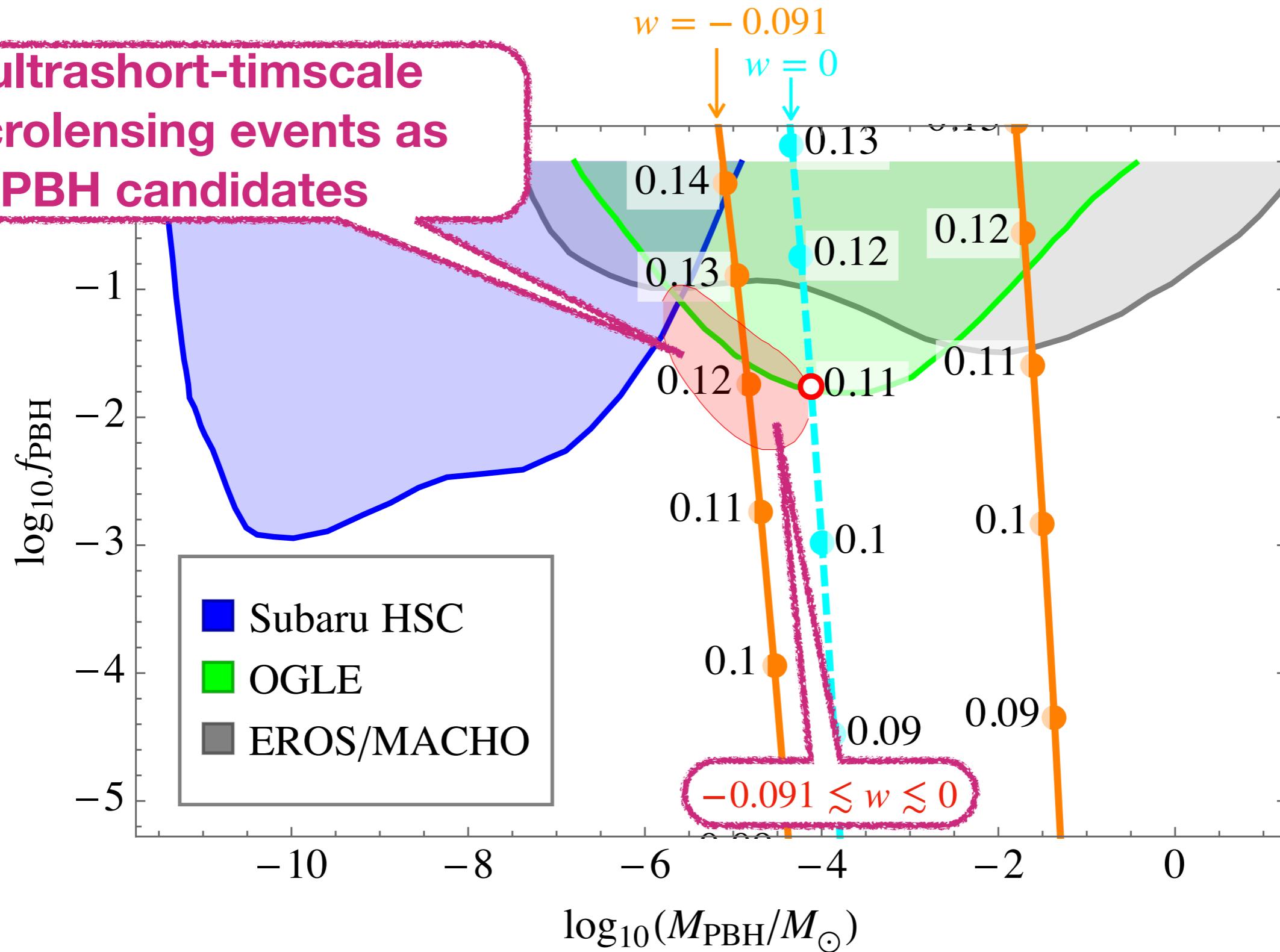


# NANOGrav SGWB



# NANOGrav PBHs

6 ultrashort-timescale  
microlensing events as  
PBH candidates



[Domènech and SP, 2010.03976]

# Summary

- Induced GWs is one of the most important scientific goal of the next generation GW detectors, which has fruitful phenomena together with PBH physics.
- IGW and PBH-DM: If PBHs can serve as all the DM, induced GWs must be detectable by space-based GW detectors.
- Shape of IGW: It depends crucially on the width of the scalar peak. We found analytical formulae for both cases in the radiation-dominated universe, which is useful for signal searching in the future.
- NANOGrav has detected common-spectrum PTA time residuals which might be the first detection of SGWB. We show that this can be connected to the recently reported planet-mass PBHs.

*Thank you  
for your attendance*