Kavli IPMU Seminar December 3 2020

Exploring the vector modes in ACDM with general-relativistic simulations

Cristian Barrera-Hinojosa Institute for Computational Cosmology Durham University





Based on 2010.08257 (C. B-H, B. Li, M. Bruni, J. He, 2020)

$$R_{\mu
u}-rac{1}{2}Rg_{\mu
u}=8\pi T_{\mu
u}$$

• Most general line element:

$$\mathrm{d}s^2 = -(1+2\Phi)\mathrm{d}t^2 + 2B_i\mathrm{d}x^i\mathrm{d}t + a^2[(1-2\Psi)\delta_{ij}+h_{ij}]\mathrm{d}x^i\mathrm{d}x^j$$

$$R_{\mu
u}-rac{1}{2}Rg_{\mu
u}=8\pi T_{\mu
u}$$

• Most general line element:

$$\mathrm{d}s^2 = -(1+2\mathbf{\Phi})\mathrm{d}t^2 + 2B_i\mathrm{d}x^i\mathrm{d}t + a^2[(1-2\mathbf{\Psi})\delta_{ij}+h_{ij}]\mathrm{d}x^i\mathrm{d}x^j$$

- Each mode encodes different physical effects.
 - Scalar = gravitational pull (Newtonian)

$$R_{\mu
u}-rac{1}{2}Rg_{\mu
u}=8\pi T_{\mu
u}$$

• Most general line element:

$$\mathrm{d}s^2 = -(1+2\mathbf{\Phi})\mathrm{d}t^2 + 2B_i\mathrm{d}x^i\mathrm{d}t + a^2[(1-2\mathbf{\Psi})\delta_{ij}+oldsymbol{h}_{ij}]\mathrm{d}x^i\mathrm{d}x^j$$

- Each mode encodes different physical effects.
 - Scalar = gravitational pull (Newtonian)
 - Tensor = ~GWs

$$R_{\mu
u}-rac{1}{2}Rg_{\mu
u}=8\pi T_{\mu
u}$$

• Most general line element:

 $\mathrm{d}s^2 = -(1+2\mathbf{\Phi})\mathrm{d}t^2 + 2\mathbf{B_i}\mathrm{d}x^i\mathrm{d}t + a^2[(1-2\mathbf{\Psi})\delta_{ij}+h_{ij}]\mathrm{d}x^i\mathrm{d}x^j$

- Each mode encodes different physical effects.
 - Scalar = gravitational pull (Newtonian)
 - Vector = frame-dragging
 - Tensor = ~GWs
- From the matter sector, we are interested in **vorticity**.



Detecting gravitomagnetic effects

• Solar System: V Gravity Probe B (Everitt+'11)



Detecting gravitomagnetic effects in cosmology

- Galactic scales: **?** Flat rotation curve of galaxies (Crosta+ '20).
- Cosmological scales: X observables currently dominated by scalar modes.
 - Gravitomagnetic effects need to be successfully isolated.
 - E.g., dipole feature in lensing convergence due to dark matter halo rotation (Tang+ '20).
- Cosmological vector modes may become non-negligible in future observations:
 - Ratio of vector potential and Newtonian potential can almost reach the percent level (Lu+ '09).
 - Vorticity can contaminate future observations (Bonvin+ '18).

Modelling the vector modes in ΛCDM



Modelling the vector modes in ΛCDM



The **GR**AMSES N-body simulation

GR extension of the particle-mesh RAMSES code (Teyssier '02).

 Gravity solved in a constrained formulation of Einstein Equations including scalar and vector modes (Bonazzola+ '04).

3+1 metric:
$$ds^2 = g_{\mu\nu}dx^{\mu}dx^{\nu} = -\alpha^2 dt^2 + \psi^4 \delta_{ij} (dx^i + \beta^i dt) (dx^j + \beta^j dt)$$

3+1 EOM: $\frac{du_i}{dt} = -\alpha u^0 \partial_i \alpha + u_j \partial_i \beta^j - \frac{u_j u_k}{2u^0} \partial_i \gamma^{jk}$
 $\frac{dx^i}{dt} = \frac{\gamma^{ij} u_j}{u^0} - \beta^i$ ~Gravitomagnetic force

- L = 256 Mpc/h, $10^9 \text{ dark matter particles}$. Planck cosmology.
- Initial conditions set by linear PT (CAMB) at z=49.
- Maximum resolution (AMR) ~ 2 kpc/h.



Results: Maps at *z*=0

Vector mode of the shift

$$egin{array}{c} B^i=eta^i_V \end{array}$$



Velocity power spectra

Velocity divergence $\theta = \nabla \cdot \mathbf{u}$

vorticity
$$\omega =
abla imes \mathbf{u}$$



• On small scales $heta \sim \omega$.

Vorticity power spectrum: shape

On large scales, the Effective Field Theory of the Large-scale Structure (EFTofLSS) predicts (Carrasco+ '14)

$$P_{\omega}(k)=A_{\omega}k^{n_{\omega}}$$

Literature:

Carrasco+ '14: n_ω≈3.6, 2.8 (theory)

```
Jelic-Cizmek+ '15 : n<sub>w</sub>≈2.5 (sim)
```

Hahn+ '15: n_{..}≈2.5 (sim)



Vorticity power spectrum: shape

On large scales, the Effective Field Theory of the Large-scale Structure (EFTofLSS) predicts (Carrasco+ '14)

$$P_{\omega}(k)=A_{\omega}k^{n_{\omega}}$$

Literature:

Carrasco+ '14: n_w≈3.6, 2.8 (theory)

Jelic-Cizmek+ '15 : n_w≈2.5 (sim)

Hahn+ '15: n_ω≈2.5 (sim)



Vorticity power spectrum: time evolution

We can try to parametrise the evolution in terms of a spectral index γ_{ω} as (Pueblas & Scoccimarro '09)

$$P_{\omega}(k;z) = \left(\frac{D_{+}(z)}{D_{+}(0)}\right)^{\gamma_{\omega}} P_{\omega}(k;z=0)$$

Literature:

Pueblas & Scoccimarro '09: γ₀≈7.0±0.3

Thomas+ 2015: $\gamma_{\omega} \gtrsim 7.0$

Jelic-Cizmek+ '15: γ_ω≳7.0



Vorticity power spectrum: time evolution (NL)

We can also parametrise the evolution on nonlinear scales with good approximation.

$$P_{\omega}(k;z) = \left(\frac{D_{+}(z)}{D_{+}(0)}\right)^{\gamma_{\omega}^{\mathsf{NL}}} P_{\omega}(k;z=0)$$

Best-fit value: $\gamma_{\omega} \approx 2.6$.

However, this seems to break down beyond z=1.



Gravitomagnetic potential: power spectrum



• **B**-spectrum consistent with Bruni+ '14, Thomas+ '15.

Gravitomagnetic potential: power spectrum



- **B**-spectrum consistent with Bruni+ '14, Thomas+ '15.
- **B**/ $|\Phi|$ remains close to PT (Lu+ '09), but the potentials can deviate ~1 order of magnitude.

Dark matter haloes

- Halo catalogue from the Friends-of-Friends halo finder *Rockstar*.
- From the catalogue, we select a sample of dark matter haloes split into different mass bins:

| Higher mass range | $M_h \ge 10^{14.5} \ h^{-1} M_{\odot}$ |
|-------------------------|--|
| Intermediate mass range | $\bar{M}_h = 10^{13.5} \ h^{-1} M_{\odot}$ |
| Lower mass range | $\bar{M}_h = 10^{12.5} \ h^{-1} M_{\odot}$ |

Dark matter haloes: potentials



Halo profiles: potentials

We calculate the **spherical average** of each field around halo centers. These are split into:

$$\Phi(r) = \Phi_{ ext{halo}} + \Phi_{ ext{env}}
onumber \ B(r) = B_{ ext{halo}} + B_{ ext{env}}$$

Then, we subtract their values at $r=2R_{200c}$.

We calculate the **ensemble average** and 1σ deviation in a given mass bin.

B/ $|\Phi|$ ratio has weak dependence on halo mass.



Dark matter haloes: forces



Halo profiles: forces

- Results for upper halo mass range are consistent with Adamek+ '16 at the 1σ level.
- Nonetheless, the gravitomagnetic acceleration can be 1 order of magnitude larger than in Adamek+ '16.
- Force ratio ~10⁻⁵, also has weak dependence on halo mass.



Conclusions

- I. We have studied the cosmological vector modes with a high-resolution GRAMSES simulation.
- II. Power-law fit for the vorticity power spectrum shape and evolution on both large and small scales.
- III. The gravitomagnetic potential and force peak towards the halo centers, correlating with halo mass.
- IV. Ratio **B**/ $|\Phi|$ remains flat around ~10⁻³ inside haloes, regardless of halo mass.
- V. Ratio between the gravitomagnetic and Newtonian force remains $\sim 10^{-5}$, regardless of halo mass.
- VI. Future work: study the impact of **B** on photons (e.g. lensing convergence) and cosmological observables.