

# Elliptic Quantum Toroidal Algebra $U_{q,t,p}(\mathfrak{gl}_{1,tor})$ and Affine Quiver Gauge Theories

Hitoshi Konno

Tokyo University of Marine Science & Technology

MS seminar at IPMU  
20 January 2022

- H.K and K.Oshima, “Elliptic Quantum Toroidal Algebra  $U_{q,t,p}(\mathfrak{gl}_{1,tor})$  and Affine Quiver Gauge Theories”, arXiv:2112.09885

# MSJ-SI 2023

## International Workshop and Mini School on Elliptic Integrable Systems, Representation Theory and Hypergeometric Functions

in

Tokyo

Past Workshops:

2004 Kyoto, 2008 Bonn, 2013 Leiden, 2017 Vienna, 2019 Stockholm

- 1 Introduction
  - $U_{q,p}(\widehat{\mathfrak{g}})$  and  $W_{p,p^*}(\mathfrak{g})$
  - $U_{q,t}(\mathfrak{gl}_{1,tor})$  and Linear Quiver Gauge Theories
- 2 Elliptic Quantum Toroidal Algebra  $U_{q,t,p}(\mathfrak{gl}_{1,tor})$
- 3 Typical  $U_{q,t,p}(\mathfrak{gl}_{1,tor})$ -modules
- 4 Intertwining Operators of  $U_{q,t,p}(\mathfrak{gl}_{1,tor})$ -modules
- 5  $U_{q,t,p}(\mathfrak{gl}_{1,tor})$  and Affine Quiver  $W$ -alg.  $W_{p,p^*}(\Gamma(\widehat{A}_0))$
- 6 Instanton Calculus in the 5d & 6d Lifts of the 4d  $\mathcal{N} = 2^*$  Theories

# Review of $U_{q,p}(\widehat{\mathfrak{g}})$ ( $\widehat{\mathfrak{g}}$ : affine Lie alg) and $W_{p,p^*}(\mathfrak{g})$

- elliptic (  $p$ : nome ) and dynamical analogue of  $U_q(\widehat{\mathfrak{g}})$  in the Drinfeld realization
- defined by gen. and rel.
  - elliptic Drinfeld currents  $E_j(z), F_j(z), \psi_j^\pm(z)$
  - $\Pi_j$  : dynamical parameters ( $\Leftrightarrow$  Kähler param. )
- Hopf algebroid str. as a co-alg. str.,
 
$$\Delta : U_{q,p} \rightarrow U_{q,p} \overset{\sim}{\otimes} U_{q,p} \quad \text{alg.hom.}$$
- $U_{q,p}(\widehat{\mathfrak{gl}}_N) \cong E_{q,p}(\widehat{\mathfrak{gl}}_N)$  (H.K '18)
 
$$E_{q,p}(\widehat{\mathfrak{gl}}_N) : \text{central extension of Felder's EQG } RLL = LLR^*$$

Level- $k$   $U_{q,p}(\widehat{\mathfrak{g}})$  realizes  $W_{p,p^*}(\mathfrak{g})$  ( $p^* = pq^{-2k}$ )

In general,

$W_{p,p^*}(\mathfrak{g})$  : the deformation of the coset type  $W$ -alg.

$$(\widehat{\mathfrak{g}})_{r-h^\vee-k} \oplus (\widehat{\mathfrak{g}})_k \supset (\widehat{\mathfrak{g}})_{r-h^\vee}$$

- For simply laced  $\widehat{\mathfrak{g}}$ :

the coset type  $W$ -alg.  $\cong$  the Hamiltonian reduction type  $W$ -alg.

Hence,  $W_{p,p^*}(\mathfrak{g}) \Leftrightarrow$  Frenkel-Reshetikhin's  $W_{q,t}(\mathfrak{g})$

by  $(p, p^*) \leftrightarrow (q, t)$  (at least for  $k = 1$ )

- For non-simply laced  $\widehat{\mathfrak{g}}$  : e.g.  $\widehat{\mathfrak{g}} = B_N^{(1)}$

$W_{p,p^*}(B_N)$  gives the deformation of Fateev-Lukyanov's  $WB_N$ -alg.

Note also  $WB_N \cong$  Hamiltonian reduction of  $B^{(1)}(0|N)$  ( Ito '90 )

# Example

$$\text{level-1 } U_{q,p}(\widehat{\mathfrak{sl}}_N) \quad W_{q,t}(\mathfrak{sl}_N) \quad \beta = \frac{r-1}{r}$$

$$p = q^{2r} \quad \leftrightarrow \quad q$$

$$p^* = pq^{-2} = q^{2(r-1)} \quad \leftrightarrow \quad t$$

$$F_j(z) \quad \leftrightarrow \quad S_j^+(z)$$

$$E_j(z) \quad \leftrightarrow \quad S_j^-(z)$$

$$\sum_{\mu} \Phi_{\mu}^D(p^{-1}z)^{-1} \Phi_{\mu}^D(z) \quad \leftrightarrow \quad T(z) = \sum_{\mu} \Lambda_{\mu}(z) \quad (\mu = 1, \dots, N)$$

$$\prod_{m=1}^a \Phi_m^D(zq^{A_m}) \quad \leftrightarrow \quad V_+^a(z) \quad (\text{Awata-Yamada '10})$$

$$\prod_{m=1}^a \Psi_m^{*D}(zq^{A'_m}) \quad \leftrightarrow \quad V_-^a(z) \quad a = 1, \dots, N-1$$

$$\prod_{m=1}^a \prod_{k=1}^{l-1} \Phi_m^D(zq^{A_m} p^{*B_k}) \quad \leftrightarrow \quad V_u^a(z) \quad u = p^l$$

# e.g. Intertwining Operators ( Vertex Operators ) of the $U_{q,p}(\widehat{\mathfrak{sl}}_N)$ -modules

$V_z = V[[z, z^{-1}]]$  : evaluation rep. ass.w. the vector rep.  $V = \bigoplus_{\mu=1}^N \mathbb{C} v_\mu$

$\mathcal{F}_\lambda$  : level-1 rep. with the h.w.  $\lambda$

$$\Phi(z) : \mathcal{F}_\lambda \rightarrow V_z \xrightarrow{\sim} \mathcal{F}_{\lambda'}$$

$$\text{s.t. } \Phi(z)x = \Delta(x)\Phi(z) \quad \forall x \in U_{q,p}$$

Components are defined by  $\Phi(z) = \sum_\mu v_\mu \tilde{\otimes} \Phi_\mu(z).$

Different coproducts  $\Delta$  (standard) and  $\Delta^D$  (the Drinfeld copro.) yield different intertwiners  $\Phi(z)$  and  $\Phi^D(z)$ , respectively.

- $\Phi^D(z) \rightsquigarrow W_{p,p^*}(\mathfrak{g})$
- $\Phi(z) \rightsquigarrow (\text{elliptic}) q\text{-KZ eq.}, \text{Stab}_{\mathfrak{C}}(F)$

# Integral Solution to the Elliptic $Q$ -KZ eq. via $U_{q,p}(\widehat{\mathfrak{sl}}_N)$

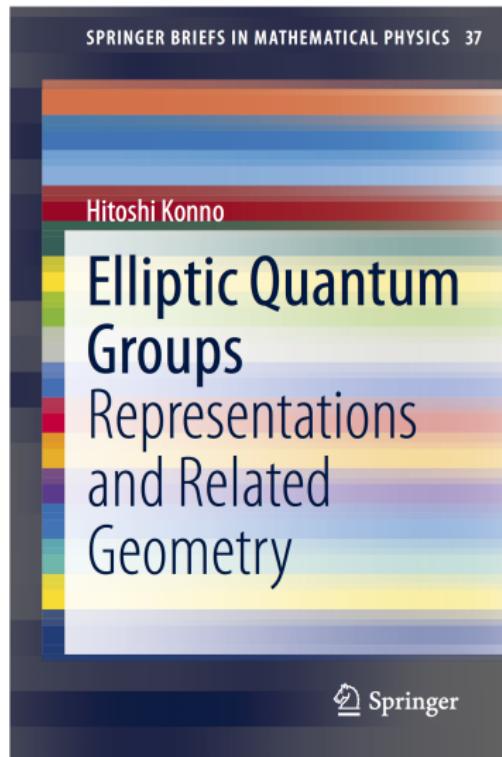
Theorem 1.1 (H.K '17, '18) *Under the zero-weight condition,*

$$\mathrm{tr}_{\mathcal{F}_\lambda} \left( Q^{-d} \Phi_{\mu_1}(z_1) \cdots \Phi_{\mu_n}(z_n) \right) = \oint_{\mathbb{T}^M} \underline{dt} \Phi(\underline{t}, \underline{z}) \mathrm{Stab}_{\mathfrak{C}}(F_I; p),$$

$$\begin{aligned} \Phi(\underline{t}, \underline{z}) &= \prod_{l=1}^{N-1} \exp \left\{ \frac{1}{\log p} \log(\Pi_l / \Pi_{l+1}) \log \left( \prod_{a=1}^{\lambda^{(l)}} t_a^{(l)} \right) \right\} \times \prod_{l=1}^{N-1} \left( \prod_{a=1}^{\lambda^{(l)}} t_a^{(l)} \right)^{\lambda_l - h_{\alpha_l}} \\ &\times \prod_{l=1}^{N-1} \left[ \prod_{a=1}^{\lambda^{(l)}} \prod_{b=1}^{\lambda^{(l+1)}} \frac{\Gamma(t_a^{(l)} / t_b^{(l+1)}; p, Q)}{\Gamma(p^* t_a^{(l)} / t_b^{(l+1)}; p, Q)} \prod_{1 \leq a < b \leq \lambda^{(l)}} \frac{\Gamma(p^* t_a^{(l)} / t_b^{(l)}, p^* t_b^{(l)} / t_a^{(l)}; p, Q)}{\Gamma(t_a^{(l)} / t_b^{(l)}, t_b^{(l)} / t_a^{(l)}; p, Q)} \right], \end{aligned}$$

$$t_a^{(N)} = z_a \quad (a = 1, \dots, n)$$

- This is an elliptic and dynamical analogue of Mimachi '96.
- Geometrically, this is an elliptic analogue of the vertex function with descendent in  $K_T(T^*Fl_n)$ . See also Aganagic-Frenkel-Okounkov "Quantum  $q$ -Langlands Corresp." '17



# Taking Residues at “ $\# = z_J p^{-d_J}$ ”

$$\begin{aligned}
 & \text{tr}_{\mathcal{F}_\lambda} \left( Q^{-d} \Phi_{\mu_1}(z_1) \cdots \Phi_{\mu_n}(z_n) \right) \\
 &= \sum_{d_{l,a} \in \mathfrak{D}} \prod_{l=1}^{N-1} \left[ \left( p^{*\lambda_{l+1}} \Pi_{l+1} / \Pi_l \right)^{\sum_{a=1}^{\lambda^{(l)}} d_{l,a}} \right. \\
 &\times \left. \prod_{a=1}^{\lambda^{(l)}} \prod_{b=1}^{\lambda^{(l+1)}} \frac{\theta_Q(pz_{j_b^{(l+1)}} / p^* z_{j_a^{(l)}}; p)_{d_{l,a} - d_{l+1,b}}}{\theta_Q(pz_{j_b^{(l+1)}} / z_{j_a^{(l)}}; p)_{d_{l,a} - d_{l+1,b}}} \prod_{1 \leq a \neq b \leq \lambda^{(l)}} \frac{\theta_Q(p^* z_{j_a^{(l)}} / z_{j_b^{(l)}}; p)_{d_{l,b} - d_{l,a}}}{\theta_Q(z_{j_a^{(l)}} / z_{j_b^{(l)}}; p)_{d_{l,b} - d_{l,a}}} \right]
 \end{aligned}$$

$$\times \text{Stab}_{\mathfrak{C}}(F_I; p)|_{F_J}$$

$$J : J^{(1)} \subset J^{(2)} \subset \cdots \subset J^{(N-1)} \subset [1, n]$$

$$J^{(l)} = \{j_1^{(l)} < \cdots < j_{\lambda^{(l)}}^{(l)}\}$$

For the full flag case, this gives an elliptic analogue of the 3d vortex partition function for the  $T[U(N)]$  theory.

# Trigonometric Limit $Q \rightarrow 0$

$$\mathrm{tr}_{\mathcal{F}_\lambda} \left( Q^{-d} \Phi_{\mu_1}(z_1) \cdots \Phi_{\mu_n}(z_n) \right)$$

$$\rightarrow \langle 0 | \Phi_{\mu_1}(z_1) \cdots \Phi_{\mu_n}(z_n) | 0 \rangle = \oint_{\mathbb{T}^M} \underline{dt} \Phi^{trig.}(\underline{t}, \underline{z}) \mathrm{Stab}_{\mathfrak{C}}(F_I; p),$$

$$\begin{aligned} & \Phi^{trig.}(\underline{t}, \underline{z}) \\ &= \prod_{l=1}^{N-1} \exp \left\{ \frac{1}{\log p} \log(\Pi_l / \Pi_{l+1}) \log \left( \prod_{a=1}^{\lambda^{(l)}} t_a^{(l)} \right) \right\} \times \prod_{l=1}^{N-1} \left( \prod_{a=1}^{\lambda^{(l)}} t_a^{(l)} \right)^{\lambda_l - h_{\alpha_l}} \\ & \times \prod_{l=1}^{N-1} \left[ \prod_{a=1}^{\lambda^{(l)}} \prod_{b=1}^{\lambda^{(l+1)}} \frac{(p^* t_a^{(l)} / t_b^{(l+1)}; p)_\infty}{(t_a^{(l)} / t_b^{(l+1)}; p)_\infty} \prod_{1 \leq a < b \leq \lambda^{(l)}} \frac{(t_a^{(l)} / t_b^{(l)}, t_b^{(l)} / t_a^{(l)}; p)_\infty}{(p^* t_a^{(l)} / t_b^{(l)}, p^* t_b^{(l)} / t_a^{(l)}; p)_\infty} \right] \end{aligned}$$

- This is the vertex function in  $K_T(T^*Fl_n)$
- The full flag case gives an integral formula for the Macdonald functions (Mimachi '96, Mimachi-Noumi '96, Noumi-Shiraishi '12)
- The sum of residues at “ $\underline{t} = \underline{z} J p^{-d_J}$ ” gives the 3d vortex partition function for the  $T[U(N)]$  theory.

# Rev. of $U_{q,t}(\mathfrak{gl}_{1,tor})$ and the Linear Quiver Gauge Theories

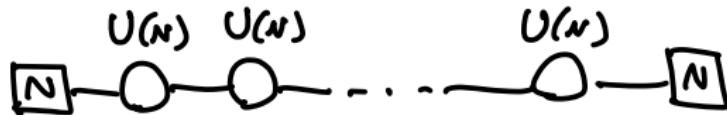
$U_{q,t}(\mathfrak{gl}_{1,tor})$  has been applied to the instanton calculus and the study of the AGT corresp. for the 5d & 6d lifts of the 4d  $\mathcal{N} = 2$  SUSY gauge theories. The essence is summarized as

- $W_{q,t}(\mathfrak{g})$  is realized on  $U_{q,t}(\mathfrak{gl}_{1,tor}) \otimes \cdots \otimes U_{q,t}(\mathfrak{gl}_{1,tor})$   
( Miki '07, Feigin-Hashizume-Hoshimo-Shiraishi-Yanagida '09,  
Berstein-Feigin-Merzon '15, Awata et.al. '16)
- Intertwiners of  $U_{q,t}(\mathfrak{gl}_{1,tor})$  w.r.t. the Drinfeld copro. realize the refined topological vertex (Awata-Feigin-Shiraishi '12)
- A certain block of composition of such intertwiners realizes the intertwiner of  $U_{q,p}(\widehat{\mathfrak{sl}}_N)$  (Zenkevich '18, Fukuda-Okubo-Shiraishi '19, '20)

# e.g. Triality in 5d Lift of $\mathcal{N} = 2$ SUSY Gauge Theories

Aganagic-Houzzi-Shakirov '14, Nedelin-Pasquetti-Zenkevich'19, ...

e.g.  $U(N)^{N-1}$  theory with  $N$  fundamental and anti-fundamental hypermultiplets  
( $A_{N-1}$  type linear quiver gauge theory)



## 5d Nekrasov Instanton PF

Higgsing ↘

.

↗ Higgsing (AGT corresp.?)

3d Vortex PF

↔

Corr. fnc. of the vertex ops.

↑  
taking residues

of  $W_{p,p^*}(\mathfrak{sl}_N)$  ("q-conf. block")

⇓

$\langle 0 | \Phi_{\mu_1}(z_1) \cdots \Phi_{\mu_N}(z_N) | 0 \rangle$   
in  $U_{q,p}(\widehat{\mathfrak{sl}}_N)$

# Triality in 6d Lift of $\mathcal{N} = 2$ SUSY Gauge Theories

Nieri '15, Iqbal-Kozcaz-Yau '15, ...

e.g.  $U(N)$  th. with  $N$  fundamental and anti-fundamental hypermultiplets  
( $A_1$  type linear quiver gauge theory)



6d Nekrasov Instanton PF

Higgsing ↕

↗ Higgsing (AGT corresp.?)

4d Vortex PF

$\iff$  Trace of the vertex ops. of  $W_{p,p^*}(\mathfrak{sl}_2)$   
taking residues ("elliptic conf. block")

⇓

$$\text{tr} \left( Q^{-d} \Phi_{\mu_1}(z_1) \cdots \Phi_{\mu_N}(z_N) \right)$$

in  $U_{q,p}(\widehat{\mathfrak{sl}}_2)$

This can be understood by the "trace" of  $U_{q,t}(\mathfrak{gl}_{1,tor})$  !

It should be distinguished from the elliptic quantum toroidal algebra !

# This talk : Elliptic Quantum Toroidal Algebra $U_{q,t,p}(\mathfrak{gl}_{1,tor})$

Main points:

- intertwiners  $\Phi(u)$ ,  $\Psi^*(u)$  w.r.t. the Drinfeld copro. and their “sifted inverse”  $\Phi^*(u)$ ,  $\Psi(u)$
- $U_{q,t,p}(\mathfrak{gl}_{1,tor})$  realizes the Jordan quiver  $W$ -alg.  $W_{p,p^*}(\Gamma(\widehat{A}_0))$   
 $SU(4)$   $\Omega$ -deformation parameters :  $q$ ,  $t$ ,  $p$ ,  $p^*$  s.t.  $q/t = p^*/p$
- Nekrasov instanton PFs of the 5d & 6d lifts of the 4d  $\mathcal{N} = 2^*$  gauge theories

# Elliptic Quantum Toroidal Algebra $U_{q,t,p}(\mathfrak{gl}_{1,tor})$

# Elliptic Quantum Toroidal Algebra $U_{q,t,p}(\mathfrak{gl}_{1,tor})$

Let  $q, t, p \in \mathbb{C}^*$ ,  $|q|, |t|, |p| < 1$ .

We define  $U_{q,t,p}(\mathfrak{gl}_{1,tor})$  to be a  $\mathbb{C}$ -algebra generated by

$$\alpha_m, \quad x_n^\pm, \quad K^{\pm 1}, \quad \gamma^{\pm 1/2} \quad (m \in \mathbb{Z} \setminus \{0\}, n \in \mathbb{Z}).$$

Let us set

$$\psi^+(z) = K \exp \left( - \sum_{m>0} \frac{p^m}{1-p^m} \alpha_{-m} (\gamma^{-1/2} z)^m \right) \exp \left( \sum_{m>0} \frac{1}{1-p^m} \alpha_{-m} (\gamma^{-1/2} z)^{-m} \right),$$

$$\psi^-(z) = K^{-1} \exp \left( - \sum_{m>0} \frac{1}{1-p^m} \alpha_{-m} (\gamma^{1/2} z)^m \right) \exp \left( \sum_{m>0} \frac{p^m}{1-p^m} \alpha_{-m} (\gamma^{1/2} z)^{-m} \right),$$

$$x^\pm(z) = \sum_{n \in \mathbb{Z}} x_n^\pm z^{-n}$$

: elliptic Drinfeld currents

$\gamma^{1/2}$ ,  $K$  : central,

$$[\alpha_m, \alpha_n] = -\frac{\kappa_m}{m} (\gamma^m - \gamma^{-m}) \gamma^{-m} \frac{1-p^m}{1-p^{*m}} \delta_{m+n,0}, \quad p^* = p\gamma^{-2}$$

$$[\alpha_n, x^+(z)] = -\frac{\kappa_n}{n} \frac{1-p^n}{1-p^{*n}} (\gamma^{-3n/2} z)^n x^+(z, p),$$

$$[\alpha_n, x^-(z)] = \frac{\kappa_n}{n} (\gamma^{-1/2} z)^n x^-(z, p),$$

$$[x^+(z), x^-(w)] = \frac{(1-q)(1-1/t)}{(1-q/t)} (\delta(\gamma^{-1} z/w) \psi^+(w) - \delta(\gamma z/w) \psi^-(\gamma^{-1} w)),$$

$$x^+(z)x^+(w) = g(z/w; p^*)^{-1} x^+(w)x^+(z),$$

$$x^-(z)x^-(w) = g(z/w; p)x^-(w)x^-(z),$$

$$\begin{aligned} g(p^* \frac{w}{z}, p^*)^{-1} g(p^* \frac{u}{w}, p^*)^{-1} g(p^* \frac{u}{z}, p^*)^{-1} & \left( \frac{w}{u} + \frac{w}{z} - \frac{z}{w} - \frac{u}{w} \right) x^+(z)x^+(w)x^+(u) \\ & + \text{permutations in } z, w, u = 0, \end{aligned}$$

$$\begin{aligned} g(p \frac{w}{z}, p) g(p \frac{u}{w}, p) g(p \frac{u}{z}, p) & \left( \frac{w}{u} + \frac{w}{z} - \frac{z}{w} - \frac{u}{w} \right) x^-(z)x^-(w)x^-(u) \\ & + \text{permutations in } z, w, u = 0, \end{aligned}$$

where

$$\kappa_m = (1-q^m)(1-t^{-m})(1-(t/q)^m),$$

$$g(z; p) = \frac{\theta_p(q^{-1}z)\theta_p((q/t)z)\theta_p(tz)}{\theta_p(qz)\theta_p((q/t)^{-1}z)\theta_p(t^{-1}z)}, \quad \theta_p(z) = (z; p)_\infty (p/z; p)_\infty$$

# Hopf Algebroid Structure via $\Delta^D$

Let  $\tilde{\otimes}$  denote a tensor product with extra condition

$$F(z; p^*) a \tilde{\otimes} b = a \tilde{\otimes} F(z; p) b, \quad p^* = p\gamma^{-2}$$

The following gives the Drinfeld coproduct for  $\mathcal{U}_{q,t,p} = U_{q,t,p}(\mathfrak{gl}_{1,tor})$ .

$$\Delta^D(\gamma^{\pm 1/2}) = \gamma^{\pm 1/2} \tilde{\otimes} \gamma^{\pm 1/2},$$

$$\Delta^D(\psi^\pm(z)) = \psi^\pm(\gamma_{(2)}^{\mp 1/2} z) \tilde{\otimes} \psi^\pm(\gamma_{(1)}^{\pm 1/2} z)$$

$$\Delta^D(x^+(z)) = 1 \tilde{\otimes} x^+(\gamma_{(1)}^{-1/2} z) + x^+(\gamma_{(2)}^{1/2} z) \tilde{\otimes} \psi^-(\gamma_{(1)}^{-1/2} z),$$

$$\Delta^D(x^-(z)) = x^-(\gamma_{(2)}^{-1/2} z) \tilde{\otimes} 1 + \psi^+(\gamma_{(2)}^{-1/2} z) \tilde{\otimes} x^-(\gamma_{(1)}^{1/2} z).$$

Here  $\gamma_{(1)} = \gamma \tilde{\otimes} 1$ ,  $\gamma_{(2)} = 1 \tilde{\otimes} \gamma$ .

# $U_{q,t,p}(\mathfrak{gl}_{1,tor})$ -modules

# Vector Representation $V(u)$ $u \in \mathbb{C}^*$

## Definition 3.1

We say that an  $\mathcal{U}_{q,t,p}$ -module has level  $(k, l) \in \mathbb{C}^2$  if  $\gamma^{1/2}$  acts by  $q^{k/2}$  and  $K$  acts by  $(q/t)^{l/2}$ .

For  $u \in \mathbb{C}^*$ ,  $V(u) := \text{Span}_{\mathbb{C}} \{ [u]_j \ (j \in \mathbb{Z}) \}$  has a level- $(0, 0)$   $\mathcal{U}_{q,t,p}$ -module structure by

$$x^+(z)[u]_j = a^+(p)\delta(q^j u/z)[u]_{j+1},$$

$$x^-(z)[u]_j = a^-(p)\delta(q^{j-1} u/z)[u]_{j-1},$$

$$\psi^\pm(z)[u]_j = \left. \frac{\theta_p(q^j t^{-1} u/z) \theta_p(q^{j-1} t u/z)}{\theta_p(q^j u/z) \theta_p(q^{j-1} u/z)} \right|_\pm [u]_j$$

where

$$a^\pm(p) = (1 - t^{\pm 1}) \frac{(p(t/q)^{\pm 1}; p)_\infty (pt^{\mp 1}; p)_\infty}{(p; p)_\infty (pq^{\mp 1}; p)_\infty}.$$

# Semi-Infinite Tensor Product Rep. $\mathcal{F}_u^{(0,1)}$ ( $q$ -Fock Rep.)

In the trig. case: Feigin<sup>2</sup>-Jimbo-Miwa-Mukhin '11

Applying  $\Delta^D$  repeatedly, one can extend the level  $(0,0)$   $\mathcal{U}_{q,t,p}$ -module structure inductively to the semi-infinite tensor product rep. of level  $(0,1)$

$$\mathcal{F}_u^{(0,1)} \subset V(u) \tilde{\otimes} V(u(t/q)^{-1}) \tilde{\otimes} V(u(t/q)^{-2}) \tilde{\otimes} \cdots$$

spanned by vectors

$$|\lambda\rangle_u = [u]_{\lambda_1-1} \tilde{\otimes} [u(t/q)^{-1}]_{\lambda_2-2} \tilde{\otimes} [u(t/q)^{-2}]_{\lambda_3-3} \tilde{\otimes} \cdots, \\ \lambda = (\lambda_1 \geq \lambda_2 \geq \cdots) \quad \lambda_l = 0 \text{ for } l \gg 1.$$

## Theorem 3.2

The following gives a level-(0,1) action of  $\mathcal{U}_{q,t,p}$  on  $\mathcal{F}_u^{(0,1)}$ .

$$x^+(z)|\lambda\rangle_u = a^+(p) \sum_{i=1}^{\ell(\lambda)+1} \delta(u_i/z) \prod_{j=1}^{i-1} \frac{\theta_p(tu_i/u_j)\theta_p(qu_i/tu_j)}{\theta_p(qu_i/u_j)\theta_p(u_i/u_j)} |\lambda + \mathbf{1}_i\rangle_u,$$

$$x^-(z)|\lambda\rangle_u = a^-(p)(q/t)^{1/2} \sum_{i=1}^{\ell(\lambda)} \delta(q^{-1}u_i/z) \prod_{j=i+1}^{\ell(\lambda)} \frac{\theta_p(qu_j/tu_i)}{\theta_p(u_j/tu_i)} \prod_{j=i+1}^{\ell(\lambda)+1} \frac{\theta_p(tu_j/u_i)}{\theta_p(qu_j/u_i)} |\lambda - \mathbf{1}_i\rangle_u,$$

$$\psi^+(z)|\lambda\rangle_u = (q/t)^{1/2} \prod_{j=1}^{\ell(\lambda)} \frac{\theta_p(t^{-1}u_j/z)}{\theta_p(q^{-1}u_j/z)} \prod_{j=1}^{\ell(\lambda)+1} \frac{\theta_p(tu_j/qz)}{\theta_p(u_j/z)} |\lambda\rangle_u,$$

$$\psi^-(z)|\lambda\rangle_u = (q/t)^{-1/2} \prod_{j=1}^{\ell(\lambda)} \frac{\theta_p(tz/u_j)}{\theta_p(qz/u_j)} \prod_{j=1}^{\ell(\lambda)+1} \frac{\theta_p(qz/tu_j)}{\theta_p(z/u_j)} |\lambda\rangle_u,$$

where we set  $u_i = q^{\lambda_i} t^{-i+1} u$ .

# Geometric Interpretation

## Conjecture 3.1

The action in Theorem 3.2 gives a level- $(0,1)$  action of  $\mathcal{U}_{q,t,p}$  on  $\bigoplus_{|\lambda|} \mathrm{E}_T(\mathrm{Hilb}_{|\lambda|}(\mathbb{C}^2))$  with  $T = \mathbb{C}^* \times \mathbb{C}^* \times \mathbb{C}^* \ni (q, t, u)$  via

$$|\lambda\rangle_u \text{'s } \Leftrightarrow \text{ fixed point classes } [\lambda] \text{'s}$$

Moreover,

$$[\lambda] = \sum_{\mu} \mathrm{Stab}_{\mathfrak{C}}^{-1}(\mu) \Big|_{\lambda} \mathrm{Stab}_{\mathfrak{C}}(\mu)$$

gives an *elliptic analogue of Macdonald symmetric function*.

Cf.  $U_{q,t}(\mathfrak{gl}_{1,tor}) \curvearrowright \bigoplus_n \mathrm{K}_T(\mathrm{Hilb}_n(\mathbb{C}^2))$  Feigin-Tsymbaliuk '11  
 $U_{q,p}(\widehat{\mathfrak{gl}}_N) \curvearrowright \bigoplus_{\lambda} \mathrm{E}_T(T^*Fl_{\lambda})$  H.K '18

Theorem 3.3 (*Level (0, 0) rep. in terms of the elliptic Ruijsenaars op.*)  
*(Cf. Miki '07 trig.case)*

On  $\mathbb{C}[x_1^{\pm 1}, \dots, x_N^{\pm 1}]$ ,

$$x^+(z) = a^+(p) \sum_{i=1}^N \prod_{j \neq i} \frac{\theta_p(tx_i/x_j)}{\theta_p(x_i/x_j)} \delta(x_i/z) T_{q,x_i},$$

$$x^-(z) = -a^-(p) \sum_{i=1}^N \prod_{j \neq i} \frac{\theta_p(t^{-1}x_i/x_j)}{\theta_p(x_i/x_j)} \delta(q^{-1}x_i/z) T_{q,x_i}^{-1},$$

$$\psi^\pm(z) = \prod_{j=1}^N \left. \frac{\theta_p(t^{-1}x_j/z)\theta_p(q^{-1}tx_j/z)}{\theta_p(x_j/z)\theta_p(q^{-1}x_j/z)} \right|_\pm,$$

or

$$\alpha_m = \frac{(1-t^{-m})(1-(q/t)^{-m})}{m} \sum_{j=1}^N x_j^m \quad (m \in \mathbb{Z} \setminus \{0\}).$$

In particular, the zero-mode  $x_0^+ = \oint_{|z|=0} \frac{dz}{2\pi iz} x^+(z)$  acts as the elliptic Ruijsenaars difference operator.

# Level $(1, N)$ Representation $\mathcal{F}_u^{(1,N)}$ of $U_{q,p}(\mathfrak{gl}_{1,tor})$

Proposition 3.4 ( Feigin-Hashizume-Hoshino-Shiraishi-Yanagida'09)

The following gives a level  $(1, N)$  representation on the Fock module  $\mathcal{F}_u^{(1,N)}$  of  $\alpha_m$  carrying a vacuum weight  $u \in \mathbb{C}^*$ .

$$\begin{aligned} x^+(z) &= uz^{-N} \left(\frac{t}{q}\right)^{\frac{3N}{4}} \exp \left\{ - \sum_{n>0} \frac{(t/q)^{n/4}}{1-(t/q)^n} \alpha_{-n} z^n \right\} \exp \left\{ \sum_{n>0} \frac{(t/q)^{3n/4}}{1-(t/q)^n} \alpha_n z^{-n} \right\}, \\ x^-(z) &= u^{-1} z^N \left(\frac{t}{q}\right)^{-\frac{3N}{4}} \exp \left\{ \sum_{n>0} \frac{(t/q)^{n/4}}{1-(t/q)^n} \alpha'_{-n} z^n \right\} \exp \left\{ - \sum_{n>0} \frac{(t/q)^{3n/4}}{1-(t/q)^n} \alpha'_n z^{-n} \right\}, \\ \psi^+(z) &= \left(\frac{t}{q}\right)^{-\frac{N}{2}} \exp \left\{ - \sum_{n>0} \frac{p^n(t/q)^{-n/4}}{1-p^n} \alpha_{-n} z^n \right\} \exp \left\{ \sum_{n>0} \frac{(t/q)^{n/4}}{1-p^n} \alpha_n z^{-n} \right\}, \\ \psi^-(z) &= \left(\frac{t}{q}\right)^{\frac{N}{2}} \exp \left\{ - \sum_{n>0} \frac{(t/q)^{n/4}}{1-p^n} \alpha_{-n} z^n \right\} \exp \left\{ \sum_{n>0} \frac{p^n(t/q)^{-n/4}}{1-p^n} \alpha_n z^{-n} \right\}, \end{aligned}$$

where

$$\alpha'_m = \frac{1-p^{*m}}{1-p^m} \gamma^m \alpha_m \quad (m \in \mathbb{Z}_{\neq 0}), \quad \gamma^{1/2} = (t/q)^{1/4}, \text{ hence } p^* = pq/t.$$

# Intertwining Operators of $U_{q,t,p}(\mathfrak{gl}_{1,tor})$ -modules

$$\text{Type I} \quad \Phi(u) : \mathcal{F}_{-uv}^{(1,N+1)} \rightarrow \mathcal{F}_u^{(0,1)} \tilde{\otimes} \mathcal{F}_v^{(1,N)}$$

$$\Delta^D(x)\Phi(u) = \Phi(u)x \quad (\forall x \in \mathcal{U}_{q,t,p})$$

### Theorem 4.1

$$\Phi(u) = \sum_{\lambda} |\lambda\rangle_u \tilde{\otimes} \Phi_{\lambda}(u),$$

$$\Phi_{\lambda}(u) = \frac{q^{n(\lambda')} t^*(\lambda, u, v, N) N_{\lambda}(p)}{c_{\lambda}} \Phi_{\emptyset}(u) \prod_{(i,j) \in \lambda} \tilde{x}^{-}((t/q)^{1/4} t^{-i+1} q^{j-1} u),$$

$$\Phi_{\emptyset}(u) = \exp \left\{ - \sum_{m>0} \frac{\alpha'_{-m}}{\kappa_m} ((t/q)^{1/2} u)^m \right\} \exp \left\{ \sum_{m>0} \frac{\alpha'_m}{\kappa_m} ((t/q)^{1/2} u)^{-m} \right\}.$$

where  $\tilde{x}^{-}(z) = uz^{-N}(t/q)^{3N/4}x^{-}(z)$ ,  $N_{\lambda}(0) = N'_{\lambda}(0) = 1$ ,

$$\langle P_{\lambda}, P_{\lambda} \rangle_{q,t} = \frac{c'_{\lambda}}{c_{\lambda}} \quad \sim \quad \frac{c'_{\lambda} N_{\lambda}(p)}{c_{\lambda} N'_{\lambda}(p)} = \frac{\prod_{\square \in \lambda} \theta_p(q^{a(\square)+1} t^{\ell(\square)})}{\prod_{\square \in \lambda} \theta_p(q^{a(\square)} t^{\ell(\square)+1})} =: \frac{c'_{\lambda}(p)}{c_{\lambda}(p)},$$

$$t^*(\lambda, u, v, N) = (q^{-1}v)^{-|\lambda|} (-u)^{N|\lambda|} f_{\lambda}(q, t)^N.$$

Cf. Trig. case : Awata-Feigin-Shiraishi '12

$$\text{Type II dual} \quad \Psi^*(v) : \mathcal{F}_u^{(1,N)} \tilde{\otimes} \mathcal{F}_v^{(0,1)} \rightarrow \mathcal{F}_{-uv}^{(1,N+1)}$$

$$x\Psi^*(v) = \Psi^*(v)\Delta^D(x) \quad (\forall x \in \mathcal{U}_{q,t,p})$$

### Theorem 4.2

$$\begin{aligned} \Psi_\lambda^*(v)\xi &= \Psi^*(v)(\xi \tilde{\otimes} |\lambda\rangle'_v) \quad \xi \in \mathcal{F}_u^{(1,N)}, \quad |\lambda\rangle'_v = \frac{c'_\lambda(p)}{c_\lambda(p)} |\lambda\rangle_v, \\ \Psi_\lambda^*(v) &= \frac{q^{n(\lambda')}}{c_\lambda} \frac{t(\lambda, u, v, N)}{\textcolor{red}{N'_\lambda}(p)} \Psi_\emptyset^*(v) \prod_{(i,j) \in \lambda} \tilde{x}^+((t/q)^{1/4} t^{-i+1} q^{j-1} v), \\ \Psi_\emptyset^*(u) &= \exp \left\{ \sum_{m>0} \frac{\alpha_{-m}}{\kappa_m} ((t/q)^{1/2} u)^m \right\} \exp \left\{ - \sum_{m>0} \frac{\alpha_m}{\kappa_m} ((t/q)^{1/2} u)^{-m} \right\}. \end{aligned}$$

where  $\tilde{x}^+(z) = u^{-1} z^N (t/q)^{-3N/4} x^+(z)$ ,

$$t(\lambda, u, v, N) = (-uv)^{|\lambda|} (-v)^{-(N+1)|\lambda|} f_\lambda(q, t)^{-N-1}.$$

# Type I “shifted inverse” $\Phi^*(v) : \mathcal{F}_u^{(0,1)} \tilde{\otimes} \mathcal{F}_v^{(1,N)} \rightarrow \mathcal{F}_{-uv}^{(1,N+1)}$

We define  $\Phi^*(u)$  by

$$\begin{aligned}\Phi_\lambda^*(v)\xi &= \Phi^*(v)(|\lambda\rangle_u \tilde{\otimes} \xi) \quad \xi \in \mathcal{F}_v^{(1,N)} \\ \Phi_\lambda^*(u) &= \frac{q^{n(\lambda')} t(\lambda, v, p^{-1}u, N) N'_\lambda(p)}{c'_\lambda} : \tilde{\Phi}_\lambda(p^{-1}u)^{-1} : \\ \tilde{\Phi}_\lambda(u) &=: \Phi_\emptyset(u) \prod_{(i,j) \in \lambda} \tilde{x}^-((t/q)^{1/4} t^{-i+1} q^{j-1} u) :\end{aligned}$$

Then  $\Phi^*(u)$  satisfies similar relations to the type I dual intertwining relations “ $x\Phi^*(v) = \Phi^*(v)\Delta^D(x)$ ”, but not exactly the same !

Hopefully, there is a suitable dual representation to  $\mathcal{F}_u^{(0,1)}$ , and  $\Phi^*(u)$  can be interpreted as the dual intertwining operator to  $\Phi(u)$  w.r.t. such dual representation.

The same happens for the shifted inverse of the type II dual VO.

# $U_{q,t,p}(\mathfrak{gl}_{1,tor})$ and Affine Quiver $W$ -alg. $W_{p,p^*}(\Gamma(\widehat{A}_0))$

Remember :  $U_{q,p}(\widehat{\mathfrak{g}})$  realizes  $W_{p,p^*}(\mathfrak{g})$

e.g. the level-1 elliptic currents  $E_i(z), F_i(z)$  of  $U_{q,p}(\widehat{\mathfrak{sl}}_N)$   
 $\iff$  screening currents  $S_i^+(z), S_i^-(z)$  of  $W_{p,p^*}(\mathfrak{gl}_N)$

Theorem 5.1 (H.K '14)

Let  $\Phi^D(z) : \mathcal{V} \rightarrow V_z \widetilde{\otimes} \mathcal{V}'$  s.t.  $\Phi^D(z)x = \Delta^D(x)\Phi^D(z) \quad \forall x \in U_{q,p}$

Then the generating functions  $T(z)$  of  $W_{p,p^*}(\mathfrak{gl}_N)$  is realized as

$$T(z) = \sum_{\mu=1}^N \Phi_\mu^D(p^{-1}z)^{-1} \Phi_\mu^D(z)$$

We have the same statement for  $D_N^{(1)}$  and  $B_N^{(1)}$ .

# $U_{q,t,p}(\mathfrak{gl}_{1,tor})$ and Affine Quiver $W$ -algebra $W_{p,p^*}(\Gamma(\widehat{A}_0))$

Screening currents :

The level  $(1, N)$  rep. of  $\mathcal{U}_{q,t,p}$  :  $\gamma^{1/2} = (t/q)^{1/4}$ ,  $p^* = p q/t$

Setting  $s_m^+ = \frac{(t/q)^{m/2}}{1 - (t/q)^m} \alpha_m$ ,  $s_m^- = \frac{(t/q)^{m/2}}{1 - (t/q)^m} \alpha'_m$ , we have

$$x^\pm((t/q)^{1/4}z) = \left( \frac{u}{z^N} \left( \frac{t}{q} \right)^{N/2} \right)^{\pm 1} : \exp \left\{ \pm \sum_{m \neq 0} s_m^\pm z^{-m} \right\} :,$$

and

$$[s_m^+, s_n^+] = -\frac{1}{m} \frac{1-p^m}{1-p^{*m}} (1-q^{-m})(1-t^m) \delta_{m+n,0},$$

$$[s_m^-, s_n^-] = -\frac{1}{m} \frac{1-p^{*-m}}{1-p^{-m}} (1-q^{-m})(1-t^m) \delta_{m+n,0}.$$

One of  $x^\pm((t/q)^{1/4}z)$  coincides with the screening current of the affine quiver  $W$ -algebra  $W_{p,p^*}(\Gamma(\widehat{A}_0))$  in **Kimura-Pestun'15** with the  $SU(4)$   $\Omega$ -deformation parameters  $p, p^*, q, t$  s.t.  $p/p^* = t/q$  (**Nekrasov '16**).

# Generating Function

$$T(u) = \Phi^*(u)\Phi(u) = \sum_{\lambda} \Phi_{\lambda}^*(u)\Phi_{\lambda}(u) = \sum_{\lambda} \mathcal{C}_{\lambda}(q, t, p) : \widetilde{\Phi^*}_{\lambda}(u)\widetilde{\Phi}_{\lambda}(u) :$$

One finds

$$: \widetilde{\Phi^*}_{\lambda}(u)\widetilde{\Phi}_{\lambda}(u) := \prod_{\square \in A(\lambda)} Y(u/q^{\square}) \prod_{\blacksquare \in R(\lambda)} Y((q/t)u/q^{\blacksquare})^{-1} :$$

where  $q^{\square} = t^{i-1}q^{-j+1}$  for  $\square = (i, j) \in \lambda$  etc. ,

$$Y(u) =: \exp \left\{ \sum_{m \neq 0} y_m u^{-m} \right\} :$$

with  $y_m = \frac{1 - p^m}{\kappa_m} (q/t)^{m/2} \alpha'_m$  satisfying

$$[y_m, y_n] = -\frac{1}{m} \frac{(1 - p^{*m})(1 - p^{-m})}{(1 - q^m)(1 - t^{-m})} \delta_{m+n,0}.$$

Hence this operator part coincides with the one of  $W_{p,p^*}(\Gamma(\widehat{A}_0))$  in  
**Kimura-Pestun'15**.

OPE coefficient  $\times$  normalization factors of  $\Phi_\lambda^*(u), \Phi_\lambda(u)$  yields

$$\mathcal{C}_\lambda(q, t, p) = \mathcal{C} \, \mathfrak{q}^{|\lambda|} Z_\lambda^{\widehat{A}_0}(t, q^{-1}, p)$$

where

$$\mathfrak{q} = p^{*-1} p^{N-1} (t/q)^{1/2}, \quad \mathcal{C} = \frac{(p^{-1}t; q, t, p)_\infty}{(q; q, t, p)_\infty},$$

$$Z_\lambda^{\widehat{A}_0}(t, q^{-1}, p^*) = \prod_{\square \in \lambda} \frac{(1 - pq^{a(\square)+1} t^{\ell(\square)})(1 - pq^{-a(\square)} t^{-\ell(\square)-1})}{(1 - q^{a(\square)+1} t^{\ell(\square)})(1 - q^{-a(\square)} t^{-\ell(\square)-1})}.$$

Note that  $\sum_{\lambda, |\lambda|=k} Z_\lambda^{\widehat{A}_0}(t, q^{-1}, p^*)$  is the equivariant Hirzebruch  $\chi_y$ -genus with  $y = p$  of  $\text{Hilb}_k(\mathbb{C}^2)$  (Li-Liu-Zhou'04). Hence

$$\begin{aligned} T(u) &= \sum_{\lambda} \Phi_\lambda^*(u) \Phi_\lambda(u) \\ &= \mathcal{C} \sum_{\lambda} \mathfrak{q}^{|\lambda|} Z_\lambda^{\widehat{A}_0}(t, q^{-1}, p) : \prod_{\square \in A(\lambda)} Y(u/q^\square) \prod_{\blacksquare \in R(\lambda)} Y((q/t)u/q^\blacksquare)^{-1} : \end{aligned}$$

coincides with the gen. fnc. of  $W_{p,p^*}(\Gamma(\widehat{A}_0))$  in Kimura-Pestun'15.

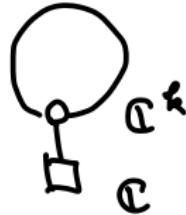
# Instanton Calculus in the 5d & 6d Lifts of the 4d $\mathcal{N} = 2^*$ Theories

# The 5d & 6d lifts of the $\mathcal{N} = 2^*$ $U(1)$ Theory

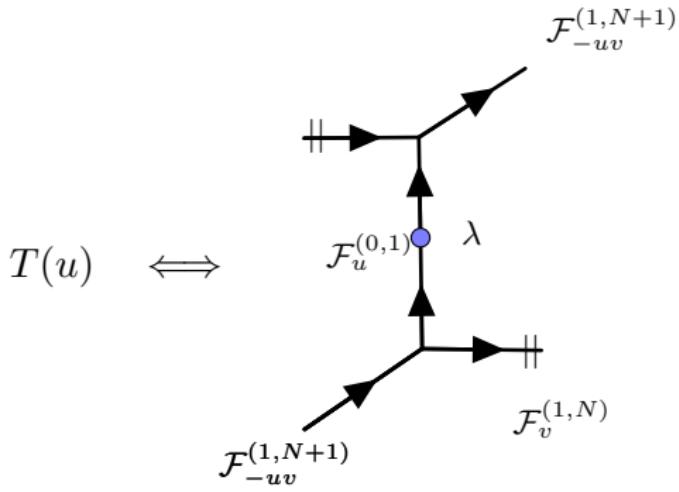
$$T(u) = \mathcal{C} \sum_{\lambda} q^{|\lambda|} Z_{\lambda}^{\widehat{A}_0}(t, q^{-1}, p) : \prod_{\square \in A(\lambda)} Y(u/q^{\square}) \prod_{\blacksquare \in R(\lambda)} Y((q/t)u/q^{\blacksquare})^{-1} :$$

From this we immediately obtain the instanton PF of the 5d lift of the  $\mathcal{N} = 2^*$   $U(1)$  Theory

$$\langle 0 | T(u) | 0 \rangle = \mathcal{C} \sum_{\lambda} q^{|\lambda|} Z_{\lambda}^{\widehat{A}_0}(t, q^{-1}, p) = \mathcal{C} \sum_{k \geq 0} q^k \underbrace{\sum_{\substack{\lambda \\ |\lambda|=k}} Z_{\lambda}^{\widehat{A}_0}(t, q^{-1}, p)}_{\chi_p\text{-genus of } \mathrm{Hilb}_k(\mathbb{C}^2)}.$$



This result and  $T(u) = \sum_{\lambda} \Phi_{\lambda}^*(u) \Phi_{\lambda}(u)$  indicate



(Hollowood-Iqbal-Vafa'08)

# The 5d & 6d lifts of the $\mathcal{N} = 2^*$ $U(1)$ Theory

The trace gives the instanton PF of the 6d lift of the  $\mathcal{N} = 2^*$   $U(1)$  theory

$$\text{tr}_{\mathcal{F}_{-uv}^{(1,N+1)}} Q^{-d} T(u) = \mathcal{C}_Q \sum_{\lambda} \mathfrak{q}^{|\lambda|} Z_{\lambda}^{\widehat{A}_0}(t, q^{-1}, p; Q), \quad v \in \mathbb{C}^*$$

where

$$Z_{\lambda}^{\widehat{A}_0}(t, q^{-1}, p; Q) = \prod_{\square \in \lambda} \frac{\theta_Q(p q^{a(\square)+1} t^{\ell(\square)}) \theta_Q(p q^{-a(\square)} t^{-\ell(\square)-1})}{\theta_Q(q^{a(\square)+1} t^{\ell(\square)}) \theta_Q(q^{-a(\square)} t^{-\ell(\square)-1})}.$$

Note :

$\sum_{\lambda, |\lambda|=k} Z_{\lambda}^{\widehat{A}_0}(t, q^{-1}, p; Q)$  is the equivariant elliptic genus of  $\text{Hilb}_k(\mathbb{C}^2)$ .

(Haghighat-Iqbal-Kozcaz-Lockhart-Vafa'15)

# The 5d & 6d Lifts of the $\mathcal{N} = 2^*$ $U(M)$ Theory

$$\begin{aligned}
 T(u_1) \cdots T(u_M) &= \mathcal{C}_M \sum_{k=0}^{\infty} \mathfrak{q}_M^k \sum_{\substack{\lambda^{(1)}, \dots, \lambda^{(M)} \\ \sum |\lambda^{(j)}| = k}} \prod_{i,j}^M Z_{\lambda^{(i)}, \lambda^{(j)}}(u_{j,i}; t, q^{-1}, p) \\
 &\times : \prod_{l=1}^M \prod_{\square \in A(\lambda^{(l)})} Y(u_l/q^\square) \prod_{\blacksquare \in R(\lambda^{(l)})} Y((q/t)u_l/q^\blacksquare)^{-1} :,
 \end{aligned}$$

where  $u_{j,i} = u_j/u_i$ ,

$$\begin{aligned}
 \mathfrak{q}_M &= \mathfrak{q} p^{-(M-1)} = p^{*-1} p^{M+N} (t/q)^{1/2}, \\
 \mathcal{Z}_{\lambda,\mu}(u; t, q^{-1}, p) &= \prod_{\square \in \lambda} \frac{(1 - puq^{a_\mu(\square)+1}t^{\ell_\lambda(\square)})}{(1 - uq^{a_\mu(\square)+1}t^{\ell_\lambda(\square)})} \prod_{\blacksquare \in \mu} \frac{(1 - puq^{-a_\lambda(\blacksquare)}t^{-\ell_\mu(\blacksquare)-1})}{(1 - uq^{a_\lambda(\blacksquare)}t^{-\ell_\mu(\blacksquare)-1})}
 \end{aligned}$$

$T(u_1) \cdots T(u_M)$  gives the higher dim. extension of  $W_{p,p^*}(\Gamma(\widehat{A}_0))$ .

We obtain the instanton PF of the 5d lift of the  $\mathcal{N} = 2^*$   $U(M)$  theory

$$\langle 0 | T(u_1) \cdots T(u_M) | 0 \rangle = \mathcal{C}_M \sum_{k=0}^{\infty} \mathfrak{q}_M^k \chi_p(\mathfrak{M}_{k,M}),$$

where

$$\chi_p(\mathfrak{M}_{k,M}) = \sum_{\substack{\lambda^{(1)}, \dots, \lambda^{(M)} \\ \sum |\lambda^{(j)}| = k}} \prod_{i,j}^M Z_{\lambda^{(i)}, \lambda^{(j)}}(u_{j,i}; t, q^{-1}, p)$$

is the  $\chi_y$  ( $y = p$ ) genus of the moduli space of rank  $M$  instantons with charge  $k$ .



(Haghighat-Iqbal-Kozcaz-Lockhart-Vafa '15)

The trace gives the instanton PF of the 6d lift of the  $\mathcal{N} = 2^*$   $U(M)$  theory

$$\mathrm{tr}_{\mathcal{F}_{-u_1 v_1}^{(1, N+1)}} Q^{-d} T(u_1) \cdots T(u_M) = \mathcal{C}_{Q, M} \sum_{k=0}^{\infty} \mathfrak{q}_M^k \mathcal{E}_{p, Q}(\mathfrak{M}_{k, M}),$$

where

$$\mathcal{E}_{p, Q}(\mathfrak{M}_{k, M}) = \sum_{\substack{\lambda^{(1)}, \dots, \lambda^{(M)} \\ \sum_j |\lambda^{(j)}| = k}} \prod_{i, j=1}^M \mathcal{Z}_{\lambda^{(i)}, \lambda^{(j)}}(u_{j, i}; t, q^{-1}, p; Q),$$

$$\begin{aligned} & \mathcal{Z}_{\lambda^{(i)}, \lambda^{(j)}}(u; t, q^{-1}, p; Q) \\ &= \prod_{\square \in \lambda^{(i)}} \frac{\theta_Q(p u q^{a_{\lambda^{(j)}}(\square)+1} t^{\ell_{\lambda^{(i)}}(\square)})}{\theta_Q(u q^{a_{\lambda^{(j)}}(\square)+1} t^{\ell_{\lambda^{(i)}}(\square)})} \prod_{\blacksquare \in \lambda^{(j)}} \frac{\theta_Q(p u q^{-a_{\lambda^{(i)}}(\blacksquare)} t^{-\ell_{\lambda^{(j)}}(\blacksquare)-1})}{\theta_Q(u q^{a_{\lambda^{(i)}}(\blacksquare)} t^{-\ell_{\lambda^{(j)}}(\blacksquare)-1})}. \end{aligned}$$

Note :

$\mathcal{E}_{p, Q}(\mathfrak{M}_{k, M})$  gives the equivariant elliptic genus of the moduli space of rank  $M$  instantons with charge  $k$ .

Hence we have shown a new AGT correspondence :

Instanton PF of the 5d & 6d lifts of the 4d  $\mathcal{N} = 2^*$  th.

$$\iff \text{corr. fnc. of } W_{p,p^*}(\Gamma(\widehat{A}_0))$$

via a realization of  $W_{p,p^*}(\Gamma(\widehat{A}_0))$  by  $U_{q,t,p}(\mathfrak{gl}_{1,tor})$