Was There an Electroweak Phase Transition ?

M.J. Ramsey-Musolf

- T.D. Lee Institute/Shanghai Jiao Tong Univ.
- UMass Amherst
- Caltech

About MJRM:



Science



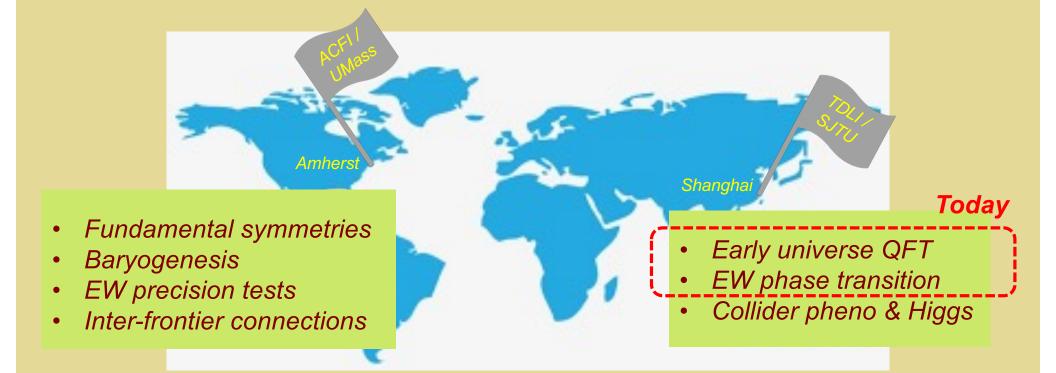
Family



Friends

My pronouns: he/him/his # MeToo IPMU Seminar March 3-4, 2022

MJRM: Scientist & "Ambassador"

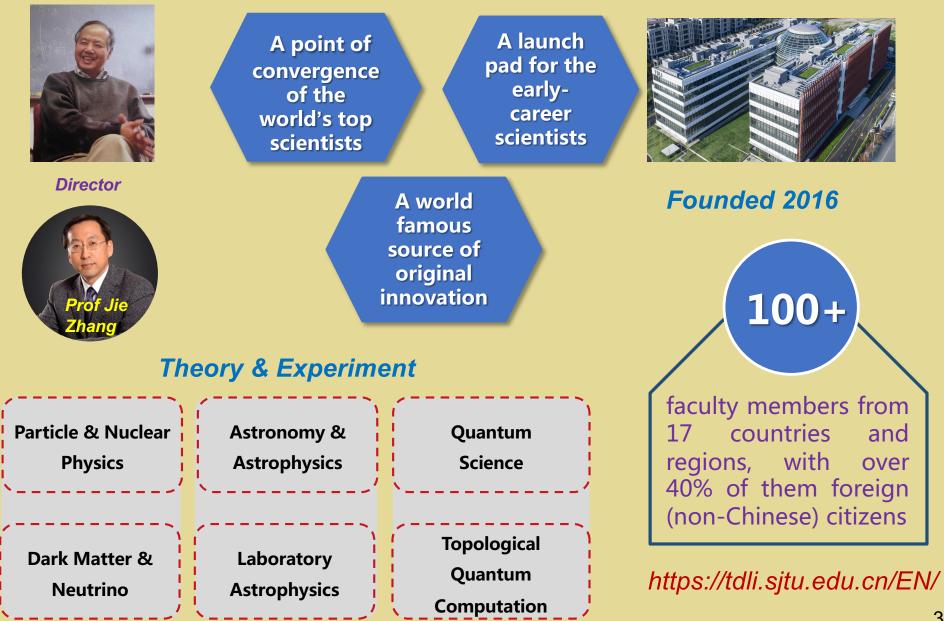




- Global effort: 18 researchers
- Foster scientific connections
- Science First ! 科学 第一 !



T. D. Lee Institute / Shanghai Jiao Tong U.



TDLI/SJTU: Particle & Nuclear Physics















Underground Experimental Group

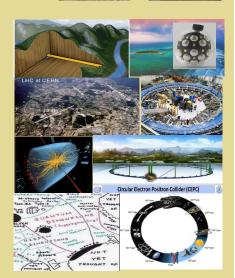
- Dark Matter and Axion (PandaX). 1.
- Neutrinoless Double Beta Decay (PandaX). 2.
- 3 Neutrino mass, Reactor and Cosmic Experiments (JUNO, ICECUBE, Hai-Ling Neutrino Telescope).

Collider Experiment Group

LHC Physics. 2. CEPC R&D. 3. Muon g-2. 4. Dark photon. 1.

Theory Group

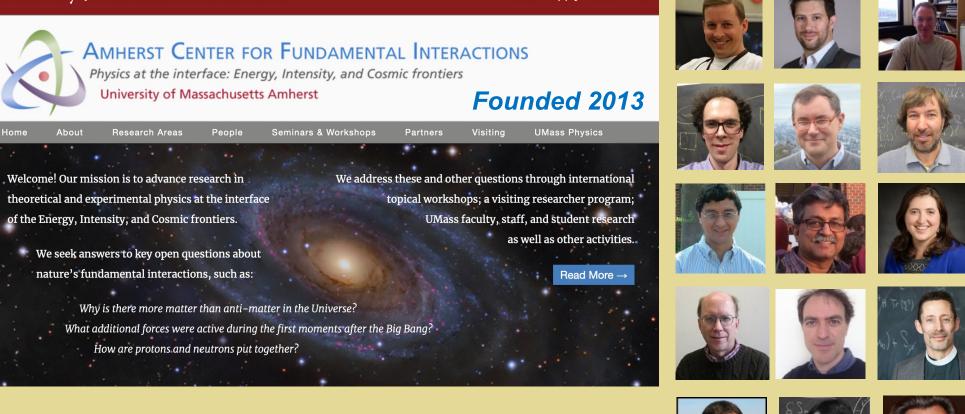
- Dark Matter, Dark Energy, Inflation, Phase Transition 1. In the Early Universe, Gravitational Waves, and Unification of Different Interactions.
- Lattice QCD Calculations, Higgs, Neutrino and Flavor 2. Physics, New Physics and Collider Phenomenology.



Amherst Center for Fundamental Interactions

University of Massachusetts Amherst

Visit Apply Give **Q**



- Experimental & theoretical research at the energy, intensity, and cosmic frontiers
- Targeted topical workshop program

https://www.physics.umass.edu/acfi/











Goals for this Talk

- Motivate the scientific opportunity associated with a possible EWPT in BSM scenarios
- Introduce the EFTs used in studying the thermal history of EW symmetry breaking
- Illustrate the interplay with non-perturbative (lattice) computations
- Highlight recent results that draw on this interplay

Key Ideas for this Talk

- The "electroweak temperature" → a scale provided by nature that gives us a clear BSM target for colliders & GW probes
- Simple arguments → BSM physics that changes the thermal history of EWSB cannot be too heavy or too feebly coupled to the SM
- Robust test of theory requires a new era of EFT & non-perturbative computations -> new results highlight this theoretical frontier

Key Ideas for this Talk

- MJRM: 1912.07189
- Recent EFT + Non-perturbative:
 - L. Niemi, H.H. Patel, MJRM, T.V.I. Tenkanen, D. J. Weir: 1802.10500
 - O. Gould, J. Kozaczuk, L. Niemi, MJRM, T.V.I. Tenkanen, D.J. Weir: 1903.11604
 - L. Niemi, MJRM, T.V.I. Tenkanen, D.J. Weir: 2005.11332

Acknowledgments

- Apologies for omissions of references to other important work
- Collaborators (this talk):
 - T. V. I. Tenkanen *
 - L. Niemi *
 - D. J. Weir
 - O. Gould
 - J. Kozaczuk
 - P. Schicho
 - J. Hirvonen
 - J. Lofgren
 - H. Patel
 - S. Arunasalam *

* TDLI / SJTU

Outline

- I. Context & Questions
- II. EWPT: A Collider & GW Target
- III. Theoretical Robustness
- IV. Outlook

I. Context & Questions

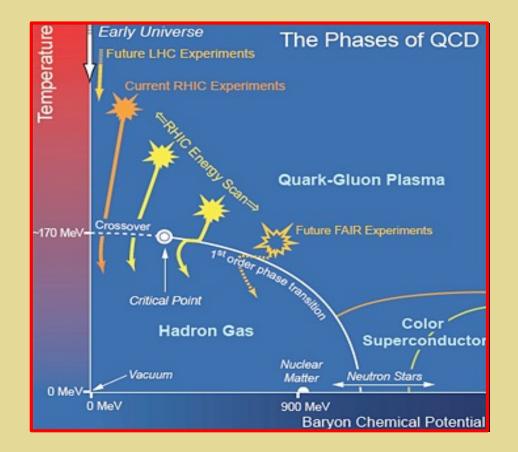
Electroweak Phase Transition

- Higgs discovery → What was the thermal history of EWSB ?
- Baryogenesis → Was the matter-antimatter asymmetry generated in conjunction with EWSB (EW baryogenesis) ?
- Gravitational waves → If a signal observed in LISA, could a cosmological phase transition be responsible ?

Electroweak Phase Transition

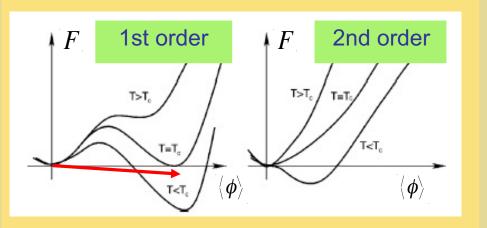
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Thermal History of Symmetry Breaking



QCD Phase Diagram \rightarrow EW Theory Analog?

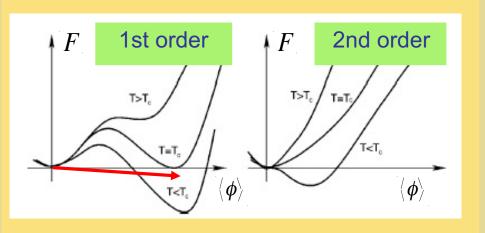
EWSB Transition: St'd Model



Increasing m_h

15

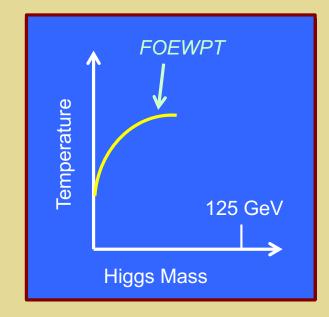
EWSB Transition: St'd Model



Increasing m_h

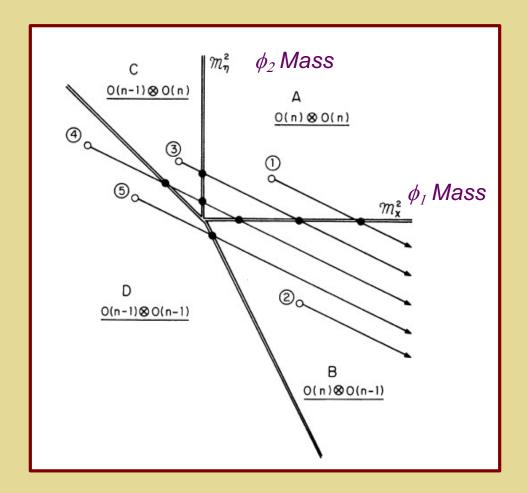
Lattice	Authors	$M_{\rm h}^C$ (GeV)
4D Isotropic	[76]	80 ± 7
4D Anisotropic	[74]	72.4 ± 1.7
3D Isotropic	[72]	72.3 ± 0.7
3D Isotropic	[70]	72.4 ± 0.9

SM EW: Cross over transition

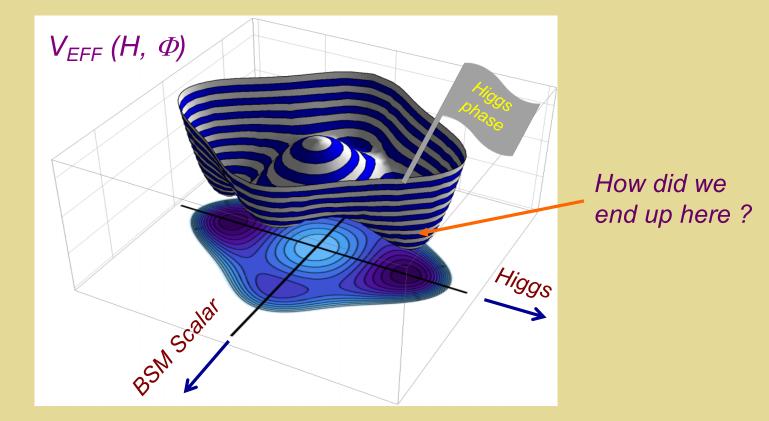


EW Phase Diagram

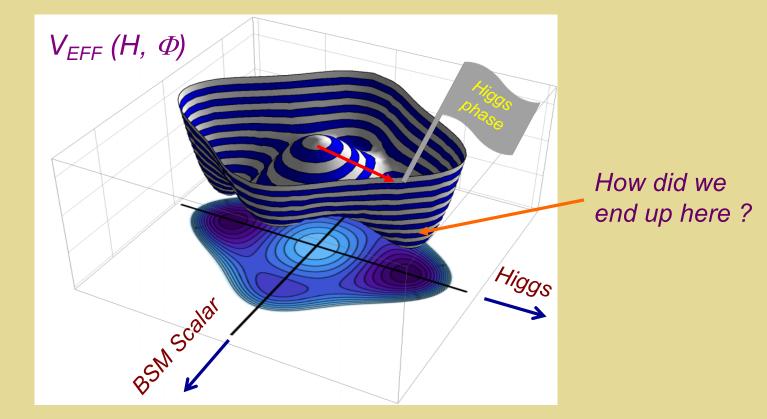
How does this picture change in presence of new TeV scale physics ? What is the phase diagram ? SFOEWPT ?



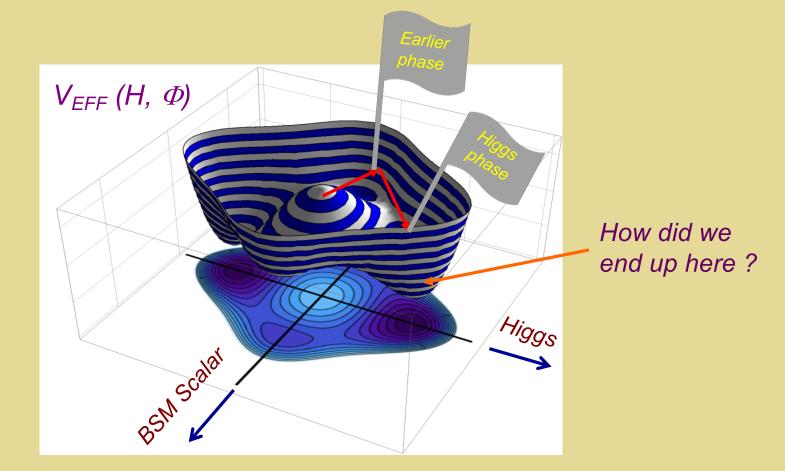
S. Weinberg, PRD 9 (1974) 3357



Extrema can evolve differently as T evolves → rich possibilities for symmetry breaking

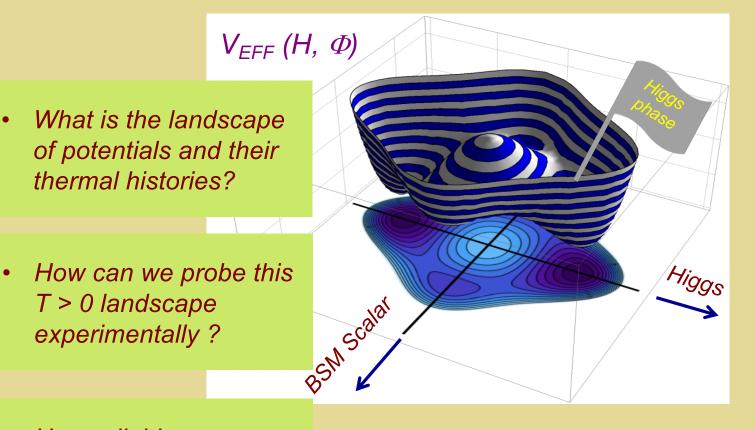


Extrema can evolve differently as T evolves → rich possibilities for symmetry breaking



Extrema can evolve differently as T evolves → rich possibilities for symmetry breaking

Thermal History of EWSB



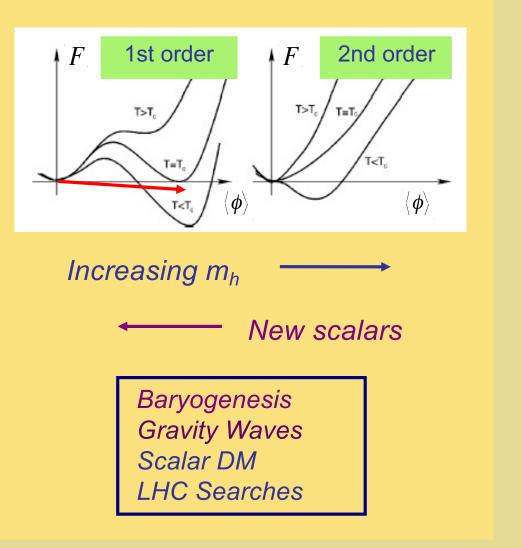
 How reliably can we compute the thermodynamics ?

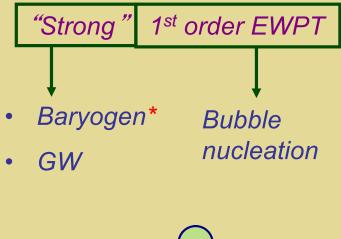
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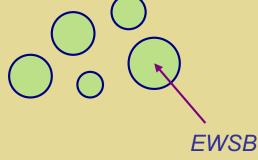
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EW Phase Transition: Baryogen & GW







* Need BSM CPV

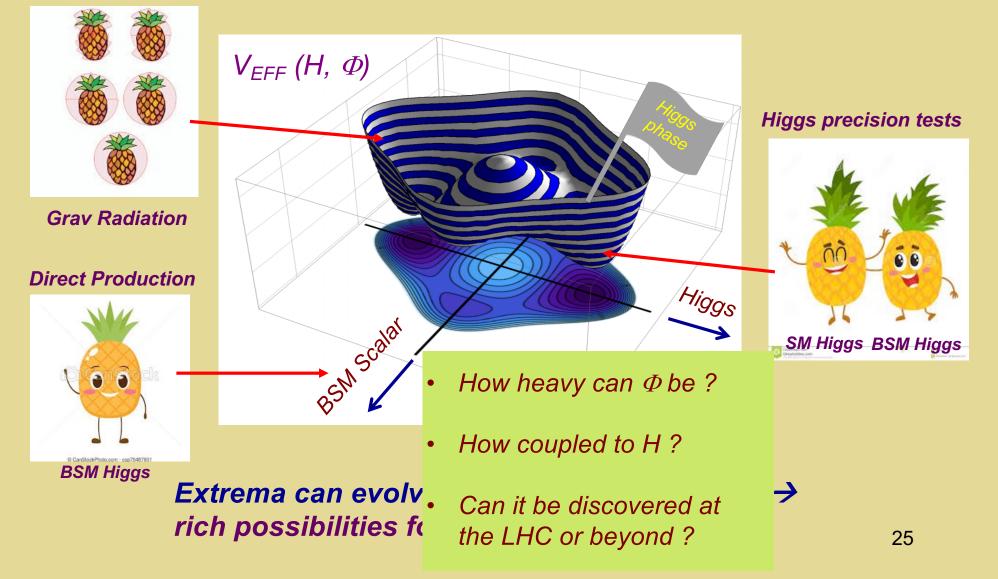
II. EWPT: A Collider Target

MJRM 1912.07189

Mass scale Precision

Experimental Probes

Bubble Collisions



T_{EW} Sets a Scale for Colliders

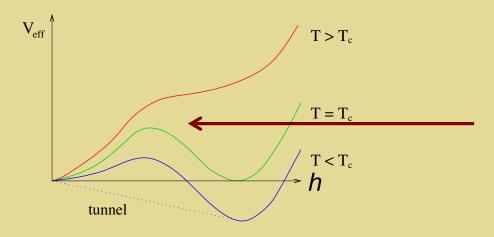
High-T SM Effective Potential

$$V(h,T)_{\rm SM} = D(T^2 - T_0^2) \, h^2 + \lambda \, h^4 \ \ \textbf{+} \ \dots$$

$$T_0^2 = (8\lambda + \text{ loops}) \left(4\lambda + \frac{3}{2}g^2 + \frac{1}{2}g'^2 + 2y_t^2 + \cdots \right)^{-1} v^2$$

$$T_0 \sim 140 \; \text{GeV} \equiv T_{EW}$$

First Order EWPT from BSM Physics

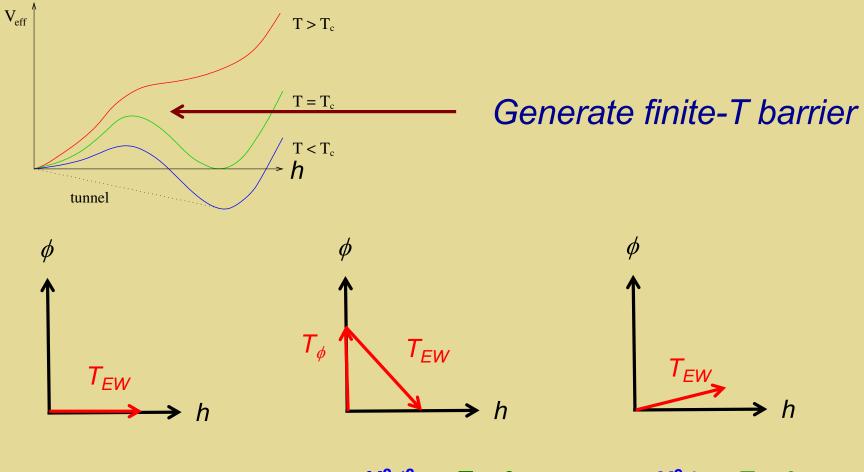


Generate finite-T barrier

Introduce new scalar ϕ interaction with h via the Higgs Portal



First Order EWPT from BSM Physics



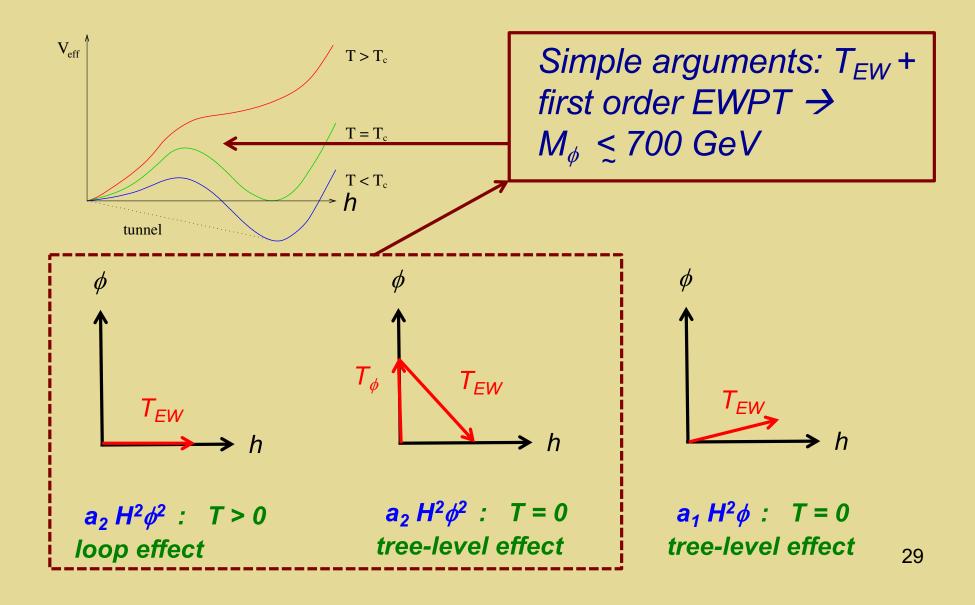
 $a_2 H^2 \phi^2$: T > 0loop effect

 $a_2 H^2 \phi^2$: T = 0tree-level effect

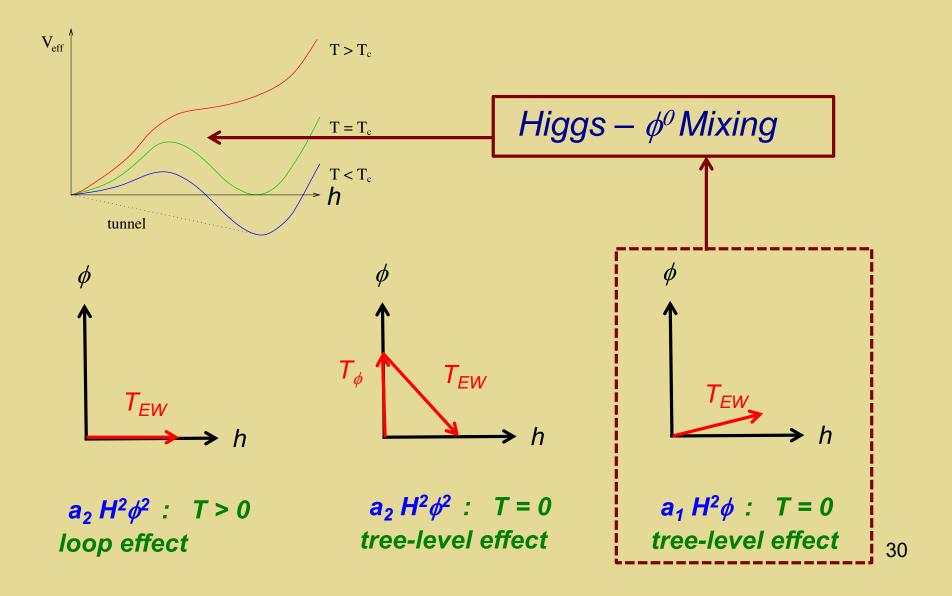
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MJRM: 1912.07189

First Order EWPT from BSM Physics

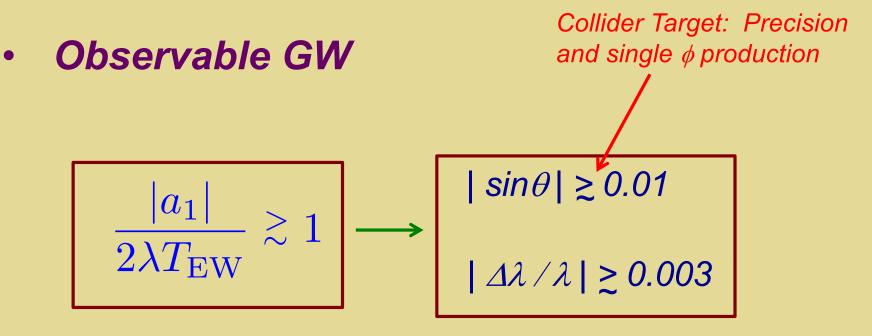


First Order EWPT from BSM Physics



Strong First Order EWPT

Prevent baryon number washout



Models & Phenomenology

What BSM Scenarios?

SM + Scalar Singlet	Espinosa, Quiros 93, Benson 93, Choi, Volkas 93, Vergara 96, Branco, Delepine, Emmanuel- Costa, Gonzalez 98, Ham, Jeong, Oh 04, Ahriche 07, Espinosa, Quiros 07, Profumo, Ramsey-Musolf, Shaughnessy 07, Noble, Perelstein 07, Espinosa, Konstandin, No, Quiros 08, Barger, Langacker, McCaskey, Ramsey-Musolf, Shaughnessy 09, Ashoorioon, Konstandin 09, Das, Fox, Kumar, Weiner 09, Espinosa, Konstandin, Riva 11, Chung, Long 11, Barger, Chung, Long, Wang 12, Huang, Shu, Zhang 12, Fairbairn, Hogan 13, Katz, Perelstein 14, Profumo, Ramsey-Musolf, Wainwright, Winslow 14, Jiang, Bian, Huang, Shu 15, Kozaczuk 15, Cline, Kainulainen, Tucker-Smith 17, Kurup, Perelstein 17, Chen, Kozaczuk, Lewis 17, Gould, Kozaczuk, Niemi, Ramsey-Musolf, Tenkanen, Weir 19
SM + Scalar Doublet (2HDM)	Turok, Zadrozny 92, Davies, Froggatt, Jenkins, Moorhouse 94, Cline, Lemieux 97, Huber 06, Froome, Huber, Seniuch 06, Cline, Kainulainen, Trott 11, Dorsch, Huber, No 13, Dorsch, Huber, Mimasu, No 14, Basler, Krause, Muhlleitner, Wittbrodt, Wlotzka 16, Dorsch, Huber, Mimasu, No 17, Bernon, Bian, Jiang 17, Andersen, Gorda, Helset, Niemi, Tenkanen, Tranberg, Vuorinen, Weir 18
SM + Scalar Triplet	Patel, Ramsey-Musolf 12, Niemi, Patel, Ramsey-Musolf, Tenkanen, Weir 18
MSSM	Carena, Quiros, Wagner 96, Delepine, Gerard, Gonzalez Felipe, Weyers 96, Cline, Kainulainen 96, Laine, Rummukainen 98, Carena, Nardini, Quiros, Wagner 09, Cohen, Morrissey, Pierce 12, Curtin, Jaiswal, Meade 12, Carena, Nardini, Quiros, Wagner 13, Katz, Perelstein, Ramsey-Musolf, Winslow 14
NMSSM	Pietroni 93, Davies, Froggatt, Moorhouse 95, Huber, Schmidt 01, Ham, Oh, Kim, Yoo, Son 04, Menon, Morrissey, Wagner 04, Funakubo, Tao, Yokoda 05, Huber, Konstandin, Prokopec, Schmidt 07, Chung, Long 10, Kozaczuk, Profumo, Stephenson Haskins, Wainwright 15

Thanks: J. M. No

Extensive references in MJRM: 1912.07189

III. Theoretical Robustness

- L. Niemi, H. Patel, MRM, T. Tenkanen, D. Weir 1802.10500
- O. Gould, J. Kozaczuk, L. Niemi, MJRM, T.V.I. Tenkanen, D.J. Weir: 1903.11604
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Models & Phenomenology

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34

Inputs from Thermal QFT

Thermodynamics

- Phase diagram: first order EWPT?
- Latent heat: GW

Dynamics

- Nucleation rate: transition occurs? T_N ? Transition duration (GW) ?
- EW sphaleron rate: baryon number preserved?

How reliable is the theory ?

EWPT & Perturbation Theory: IR Problem

Bosonic loop at T>0

$$I(T) = g^{2} \int \frac{d^{3}p}{(2\pi)^{3}} f_{B}(E,T) \frac{1}{(p^{2}+m^{2})^{n}} \longrightarrow \boxed{\frac{g^{2}T}{m}} \int_{I.R.} \frac{d^{3}p}{(2\pi)^{3}} \frac{1}{(p^{2}+m^{2})^{n}}$$
Small p regime
Effective expansion parameter

$$f_B(E,T) \longrightarrow \frac{T}{m}$$

Field-dependent thermal mass

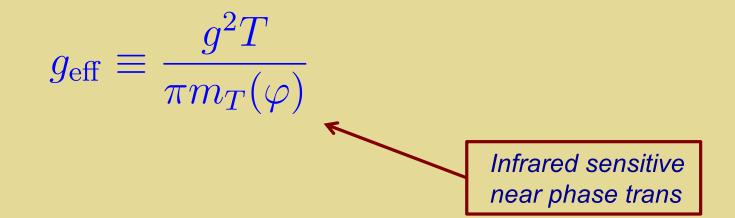
$$m^2(\varphi,T) \sim C_1 g^2 \varphi^2 + C_2 g^2 T^2 \equiv m_T^2(\varphi)$$

Near phase transition: $\varphi \sim 0$

•
$$m_T(\varphi) < g T$$

EWPT & Perturbation Theory

Expansion parameter

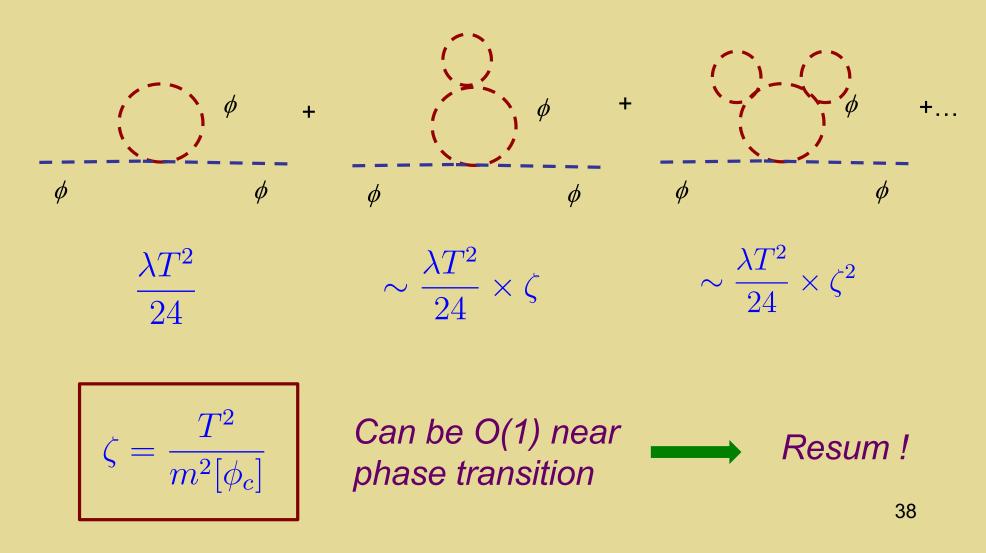


SM lattice studies: $g_{eff} \sim 0.8$ in vicinity of EWPT for $m_H \sim 70 \text{ GeV}^*$

* Kajantie et al, NPB 466 (1996) 189; hep/lat 9510020 [see sec 10.1]

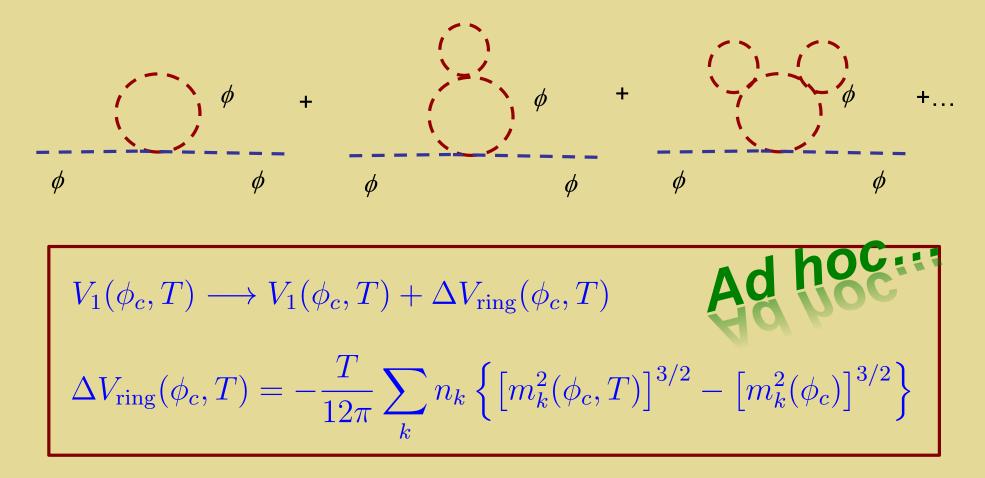
Power Counting & Resummations

Thermal masses *∆m*² (*T*):



Power Counting & Resummations

"Daisy" or "ring" resummation



Systematic resummation: dimensional reduced 3D EFT

Challenges for Theory

Perturbation theory

- I.R. problem: poor convergence
- Thermal resummations
- Gauge Invariance
 (radiative barriers)
- RG invariance at T>0

Non-perturbative (I.R.)

 Computationally and labor intensive

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BSM proposals

Theory Meets Phenomenology

A. Non-perturbative

- Most reliable determination of character
 of EWPT & dependence on parameters
- Broad survey of scenarios & parameter space not viable
- A. Perturbative
 - Most feasible approach to survey broad ranges of models, analyze parameter space, & predict experimental signatures
 - Quantitative reliability needs to be verified

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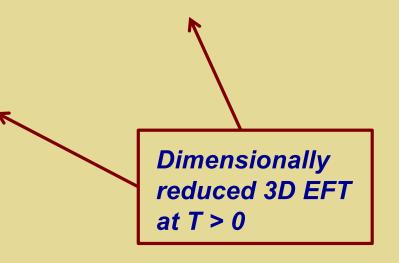
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BSM proposals

Non-perturbative (I.R.)

 Computationally and labor intensive



44

Strategy

- Employ dimensionally-reduced 3D EFT in two regimes:
 - Heavy BSM scalars → integrate out and "repurpose" existing lattice computations
 - Light BSM scalars → perform new lattice simulations
 - Compare with perturbative computations at benchmark parameter points in selected models

Inputs from Thermal QFT: EFTs

Thermodynamics

- Phase diagram: first order EWPT?
- Latent heat: GW

EFT 1

Dynamics

EFT 2

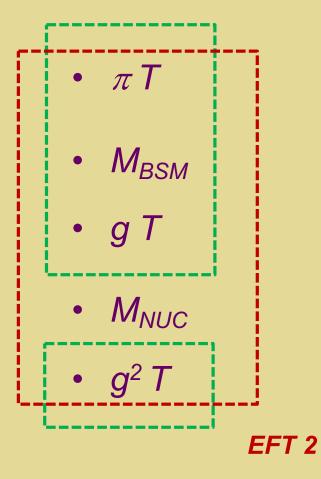
- Nucleation rate: transition occurs? T_N ? Transition duration (GW) ?
- EW sphaleron rate: baryon number preserved?

EFT 3



High-T EFT: Dimensional Reduction

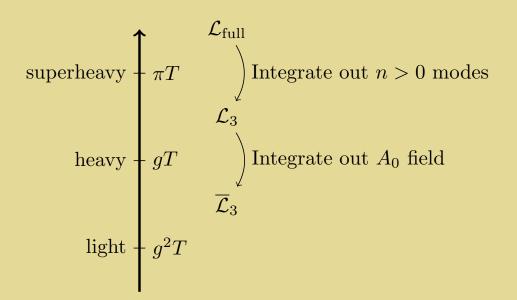
DR 3dEFT: Scales



Non-zero Matsubara modes BSM mass scale: can be > or < π T Thermal masses Nucleation scale ~ 1/r_{bubble} Light scale

EFT 1

Meeting ground: 3-D high-T effective theory



Matching: Two Elements

Dimensional Reduction

All integrals are 3D with prefactor T \rightarrow Rescale fields, couplings...

$$\int \frac{d^4k}{(2\pi)^4} \longrightarrow \frac{1}{\beta} \sum_n \int \frac{d^3k}{(2\pi)^3}$$

•
$$\varphi^2_{4d} = T \varphi^2_{3d}$$

• $T \lambda_{4d} = \lambda_{3d}$

Thermal Loops

Equate Greens functions

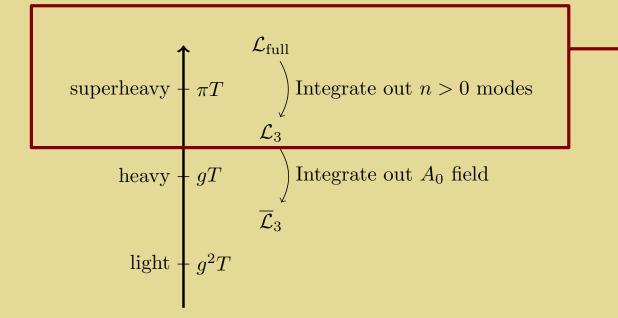
$$\phi_{\rm 3d}^2 = \frac{1}{T} \left[1 + \hat{\Pi}'_{\phi}(0,0) \right] \phi^2$$

$$a_{2,3} = T \left[a_2 - a_2 (\hat{\Pi}'_H(0) + \hat{\Pi}'_{\Sigma}(0)) + \hat{\Gamma}(0) \right]$$

Field

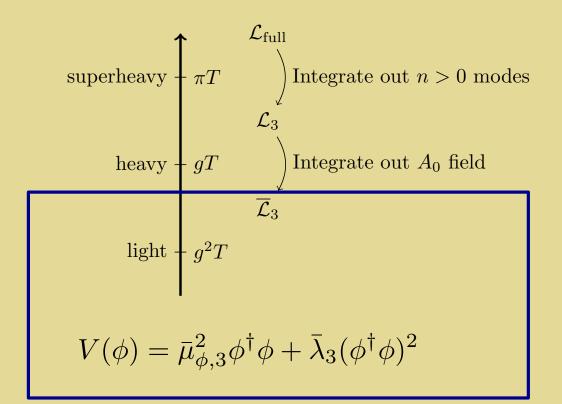
Quartic coupling

Meeting ground: 3-D high-T effective theory



Thermal resummations: systematically implemented

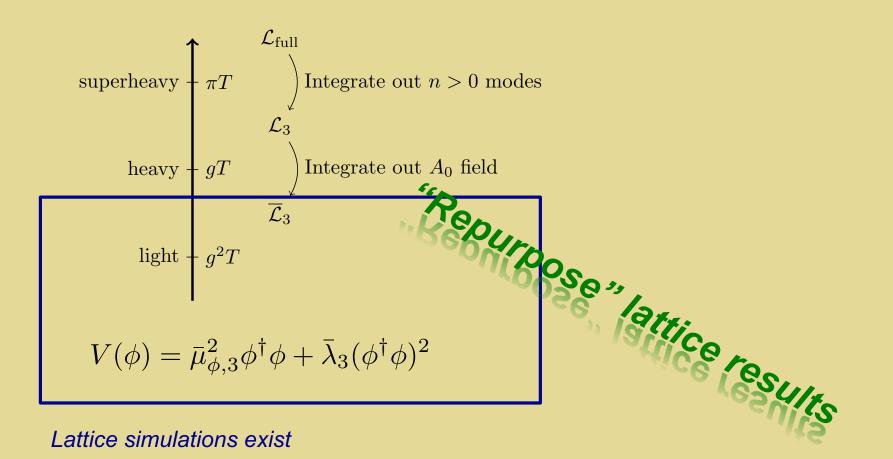
Meeting ground: 3-D high-T effective theory



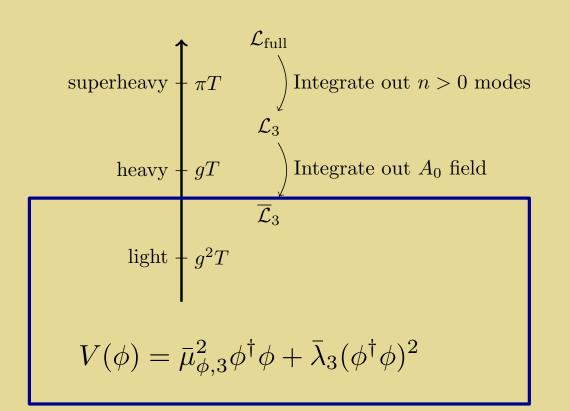
When \mathcal{L}_{full} contains BSM interactions, λ_3 and $\mu_{\phi,3}$ can accommodate first order EWPT and $m_h = 125$ GeV

Lattice simulations exist

Meeting ground: 3-D high-T effective theory

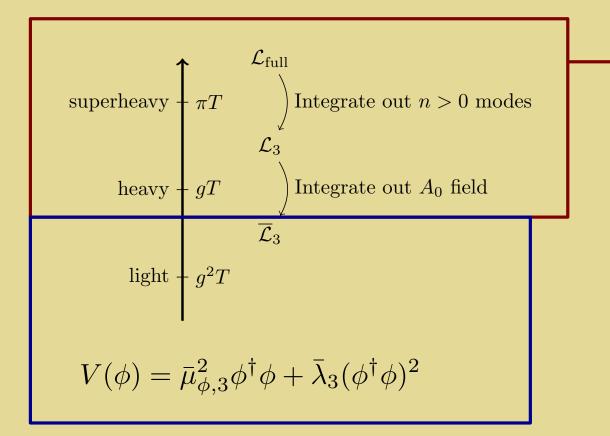


Meeting ground: 3-D high-T effective theory



- Assume BSM fields are "heavy" or "supeheavy" : integrate out
- Effective "SM-like" theory parameters are functions of BSM parameters
- Use existing lattice computations for SM-like effective theory & matching onto full theory to determine FOEWPT-viable parameter space regions

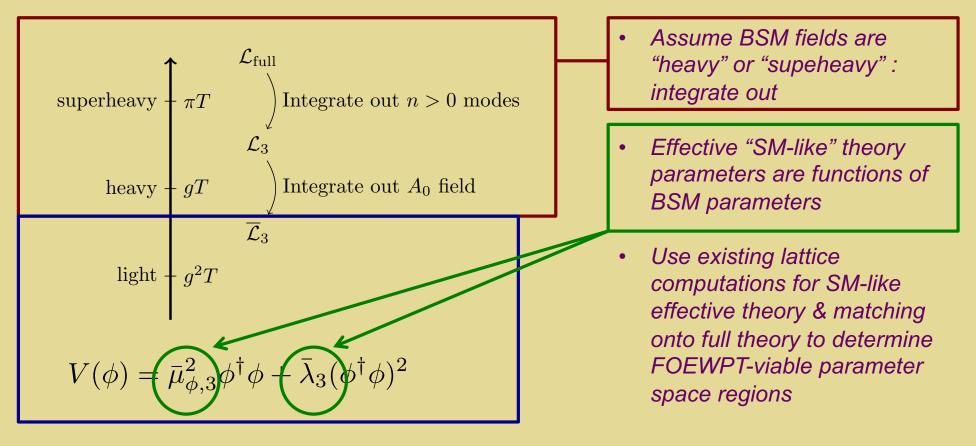
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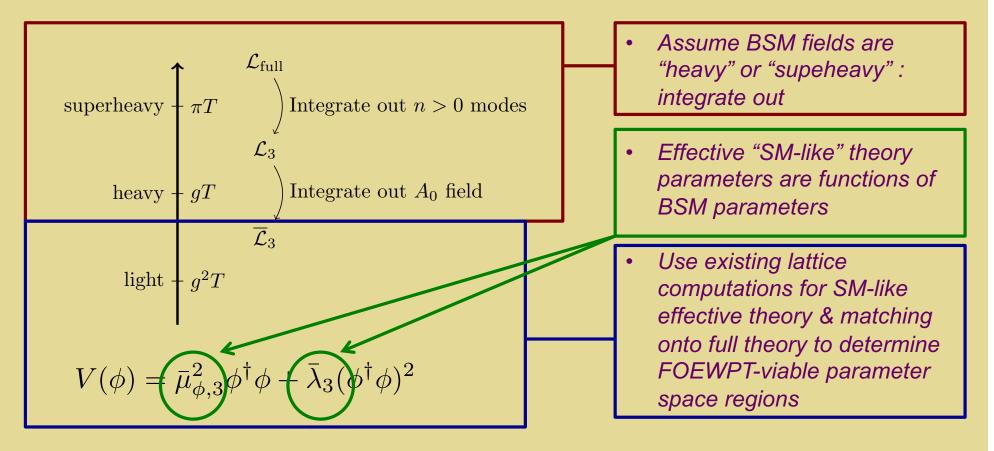
Benchmarking PT: Recent Progress

Meeting ground: 3-D high-T effective theory



Benchmarking PT: Recent Progress

Meeting ground: 3-D high-T effective theory



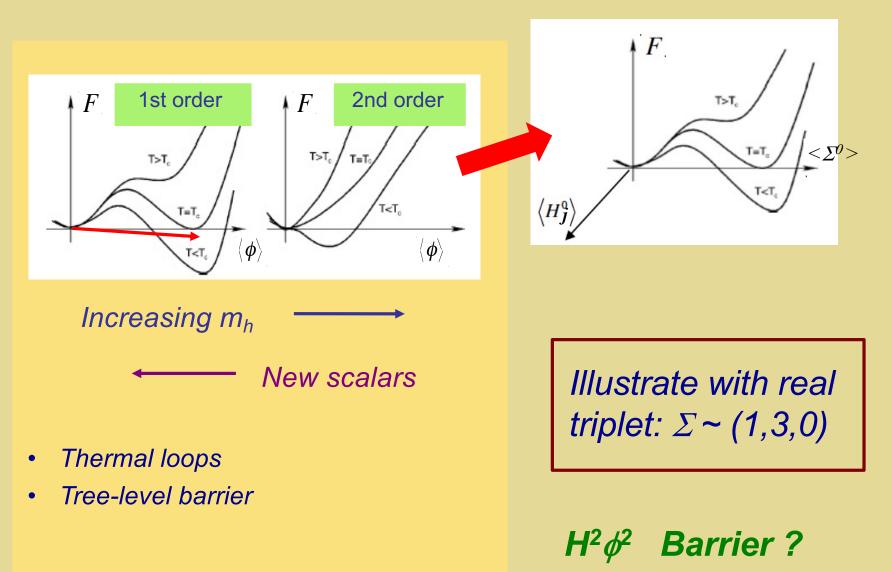
Model Illustrations



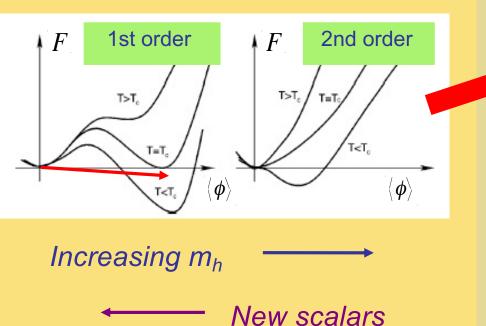
Simple Higgs portal models:

- Real gauge singlet (SM + 1)
- Real EW triplet (SM + 3)

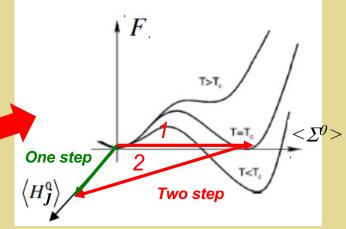
EW Multiplets: EWPT

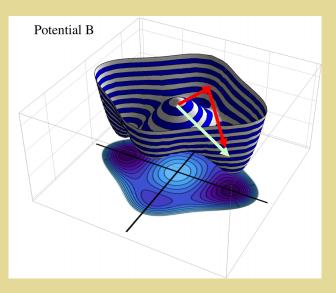


EW Multiplets: Two-Step EWPT

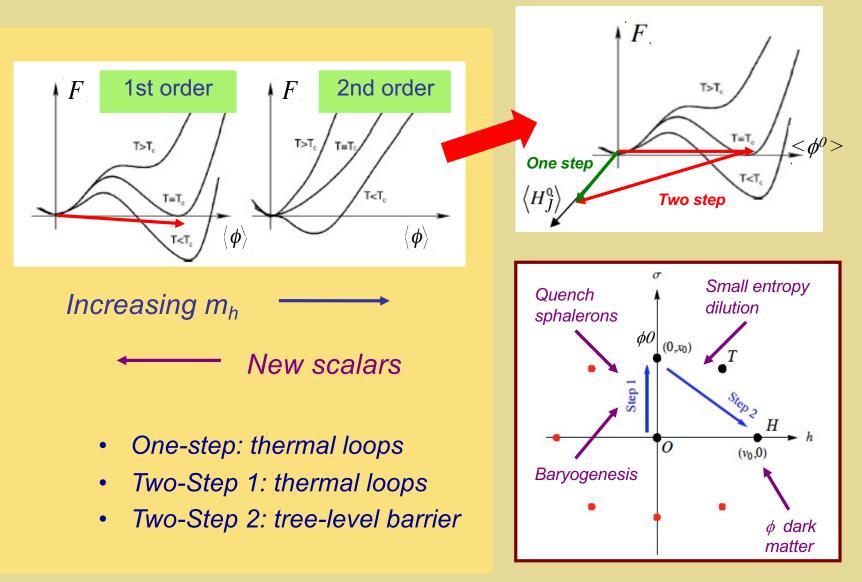


- One-step: Sym phase → Higgs phase
- Two-step: successive EW broken
 phases



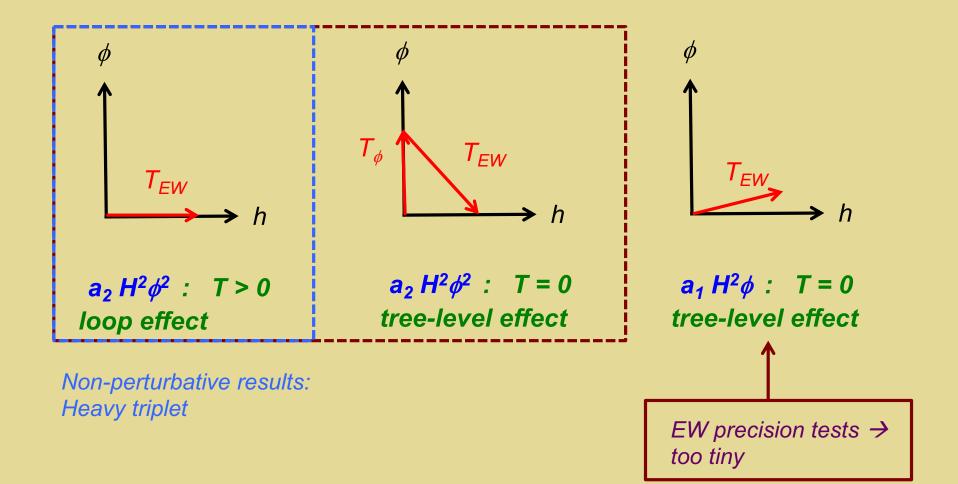


EW Multiplets: Two-Step EWPT



Patel, R-M: arXiv 1212.5652 ; Blinov et al: 1504.05195

Real Triplet



Super Heavy Triplet: EFT 1

$$\begin{split} V(\phi,\Sigma) &= -\,\mu_{\phi}^2 \phi^{\dagger} \phi - \frac{1}{2} \mu_{\Sigma}^2 \Sigma^a \Sigma^a + \lambda (\phi^{\dagger} \phi)^2 \\ &+ \frac{b_4}{4} (\Sigma^a \Sigma^a)^2 + \frac{a_2}{2} \phi^{\dagger} \phi \Sigma^a \Sigma^a, \end{split}$$

Full theory potential

$$\begin{split} \lambda_3 &= T \Biggl\{ \lambda(\Lambda) + \frac{1}{(4\pi)^2} \Biggl[\frac{1}{8} \Bigl(3g^4 + {g'}^4 + 2g^2 {g'}^2 \Bigr) + 3L_f \Bigl(y_t^4 - 2\lambda y_t^2 \Bigr) - L_b \Bigl(\frac{3}{16} \Bigl(3g^4 + {g'}^4 + 2g^2 {g'}^2 \Bigr) \\ &- \frac{3}{2} \Bigl(3g^2 + {g'}^2 - 8\lambda \Bigr) \lambda + \boxed{\frac{3}{4} a_2^2} \Biggr] \Biggr\}, \end{split}$$

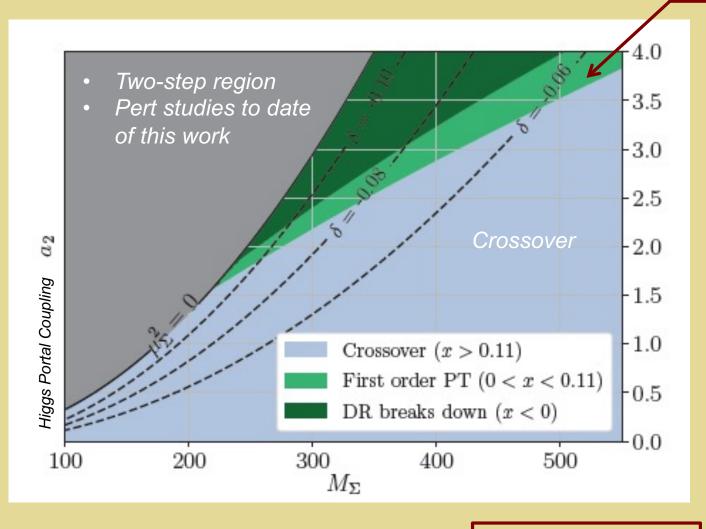
Matching: heavy theory

Allows for first order EWPT with λ (Λ) consistent with m_h = 125 GeV

Niemi, R-M, Tenkanen, Weir 2005.11332

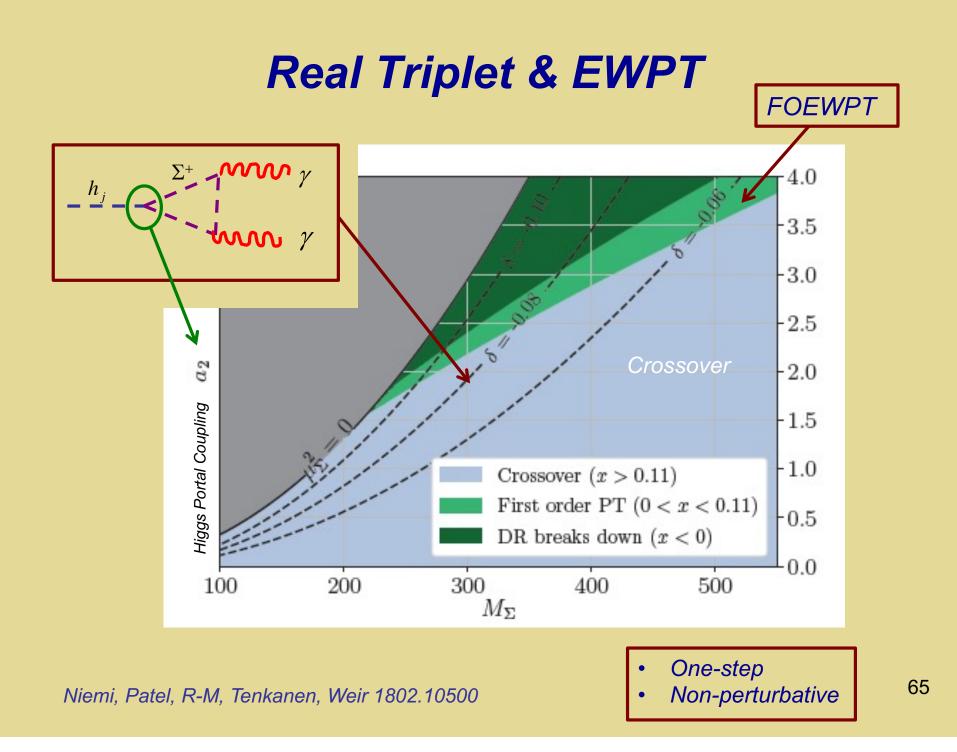
Real Triplet: One-Step EWPT

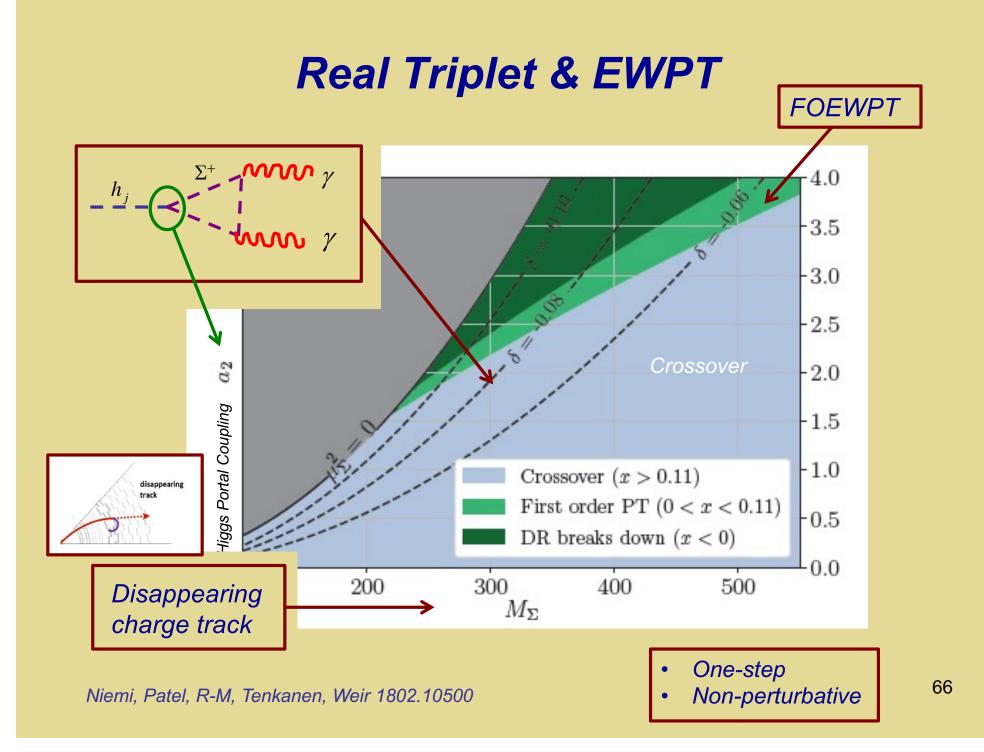
FOEWPT

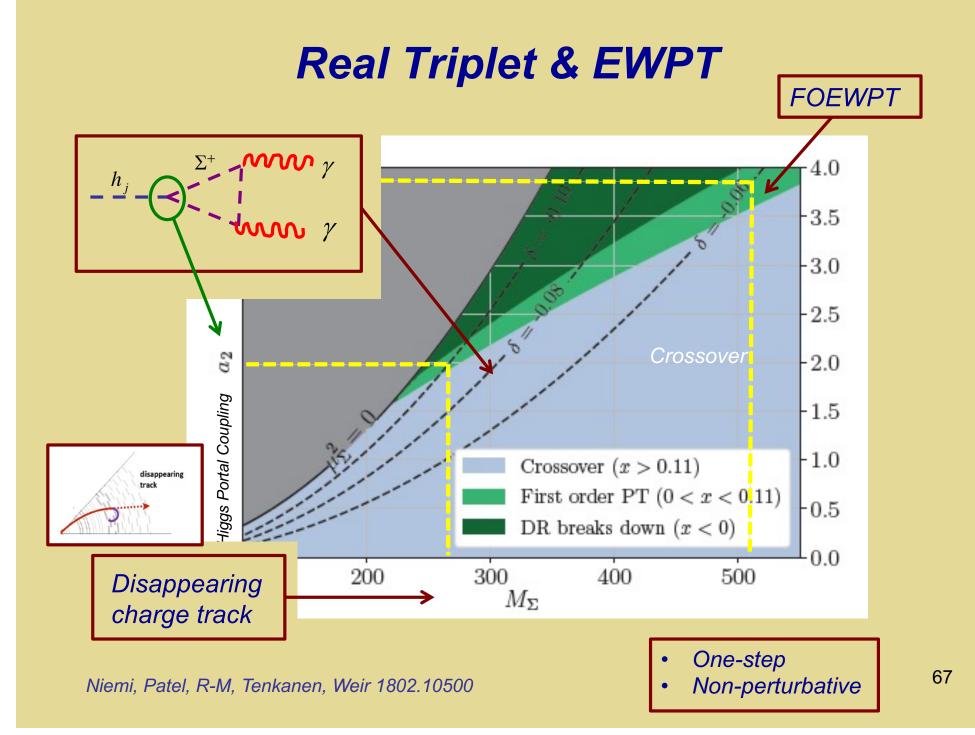


Niemi, Patel, R-M, Tenkanen, Weir 1802.10500

- One-step
- Non-perturbative

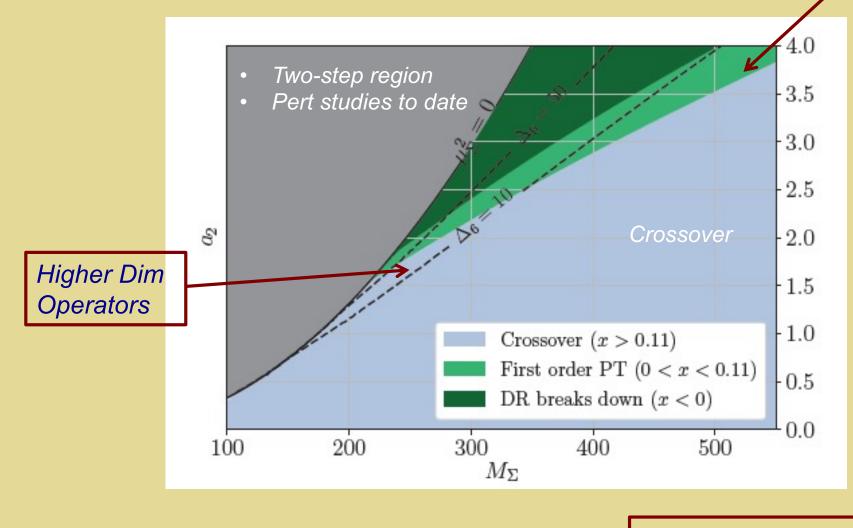






Real Triplet: One-Step EWPT

FOEWPT



Niemi, Patel, R-M, Tenkanen, Weir 1802.10500

- One-step
- Non-perturbative

68

Dynamical Real Triplet: EFT 1

$$\begin{split} V(\phi,\Sigma) &= -\,\mu_{\phi}^2 \phi^{\dagger} \phi - \frac{1}{2} \mu_{\Sigma}^2 \Sigma^a \Sigma^a + \lambda (\phi^{\dagger} \phi)^2 \\ &+ \frac{b_4}{4} (\Sigma^a \Sigma^a)^2 + \frac{a_2}{2} \phi^{\dagger} \phi \Sigma^a \Sigma^a, \end{split}$$

Full theory potential

$$\lambda_{3} = T \Biggl\{ \lambda(\Lambda) + \frac{1}{(4\pi)^{2}} \Biggl[\frac{1}{8} \Bigl(3g^{4} + {g'}^{4} + 2g^{2}{g'}^{2} \Bigr) + 3L_{f} \Bigl(y_{t}^{4} - 2\lambda y_{t}^{2} \Bigr) - L_{b} \Bigl(\frac{3}{16} \Bigl(3g^{4} + {g'}^{4} + 2g^{2}{g'}^{2} \Bigr) - \frac{3}{2} \Bigl(3g^{2} + {g'}^{2} - 8\lambda \Bigr) \lambda \Bigl(\frac{3}{4}a_{2}^{2} \Bigr) \Biggr] \Biggr\}$$

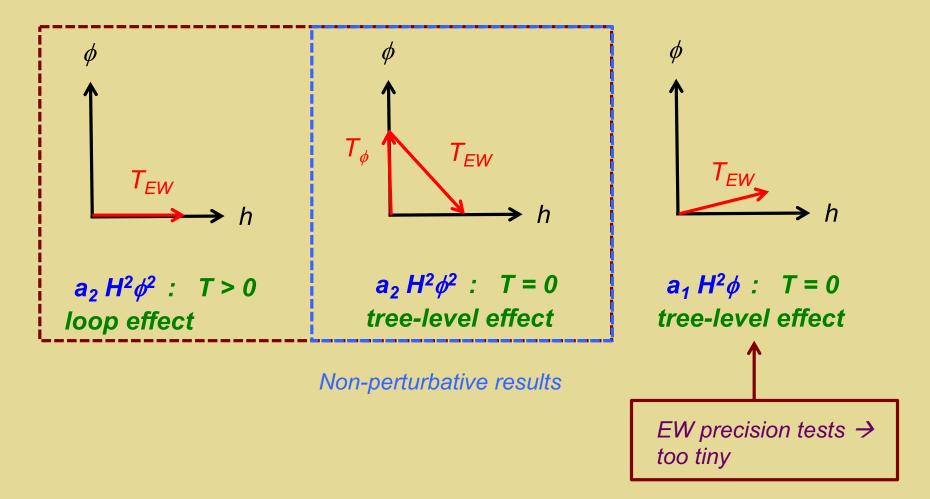
$$\mathcal{L}_{3d} = \frac{1}{4} (F_{ij}^{a})^{2} + |D_{i}\phi|^{2} + \frac{1}{2} (D_{i}\Sigma^{a})^{2} + \bar{\mu}_{\phi}^{2}\phi^{\dagger}\phi + \bar{\lambda}(\phi^{\dagger}\phi)^{2} + \frac{\bar{\mu}_{\Sigma}^{2}}{2} \Sigma^{a}\Sigma^{a} + \frac{\bar{b}_{4}}{4} (\Sigma^{a}\Sigma^{a})^{2} + \frac{\bar{a}_{2}}{2} \phi^{\dagger}\phi\Sigma^{a}\Sigma^{a}.$$
(2)

Matching: heavy theory

Light theory

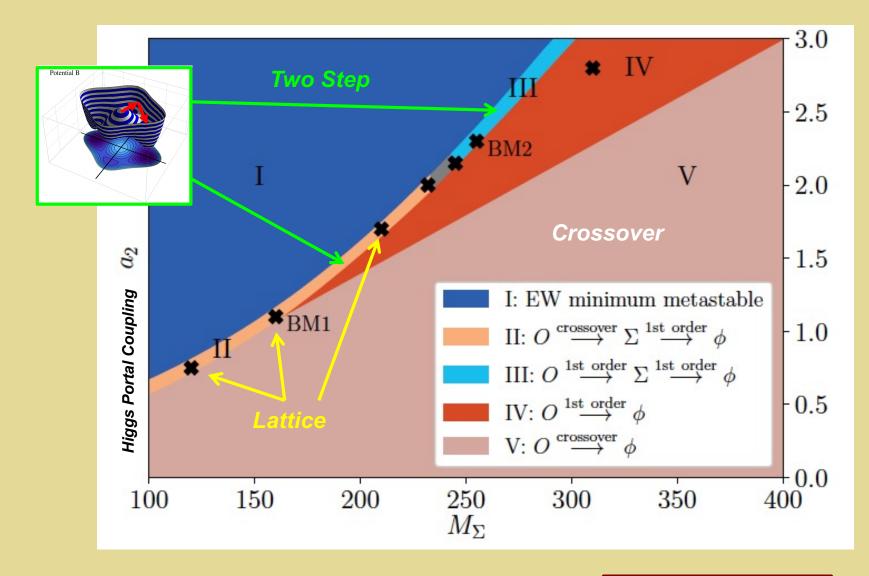
(2)

Real Triplet



70

Real Triplet & EWPT: Novel EWSB

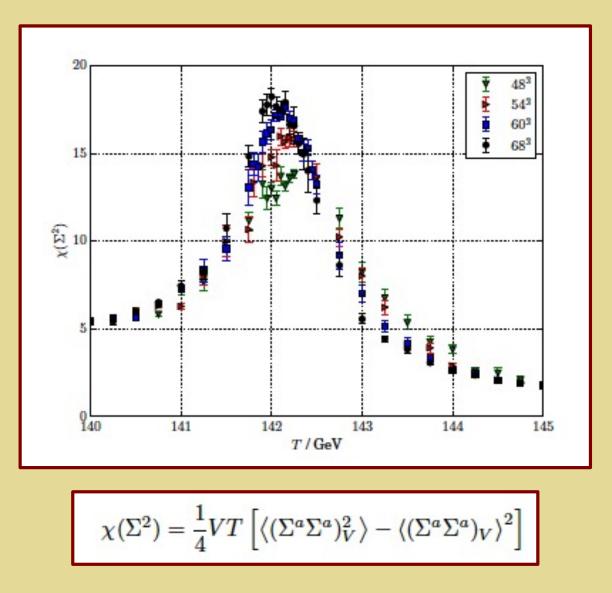


Niemi, R-M, Tenkanen, Weir 2005.11332

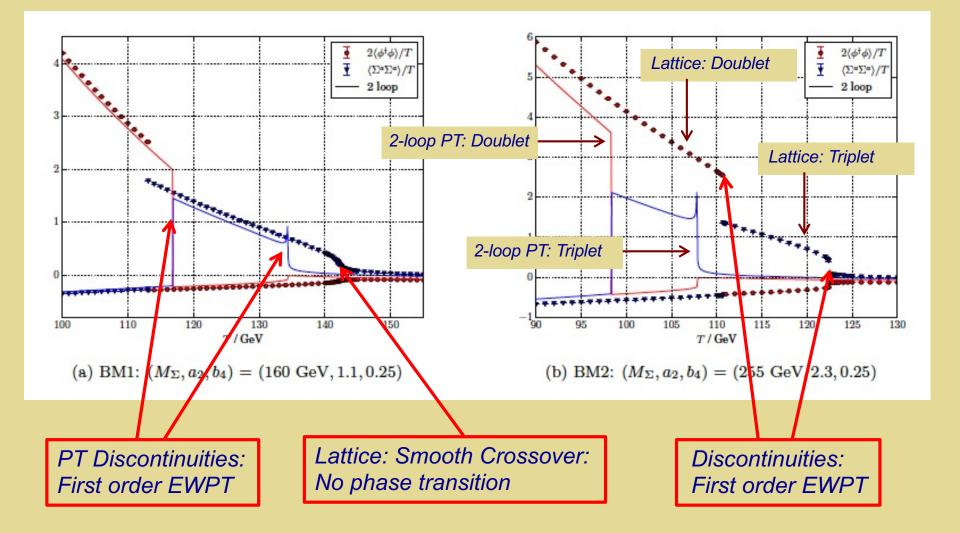
- 1 or 2 step
- Non-perturbative

71

Real Triplet: Crossover vs 2nd Order



Real Triplet & EWPT: Benchmark PT



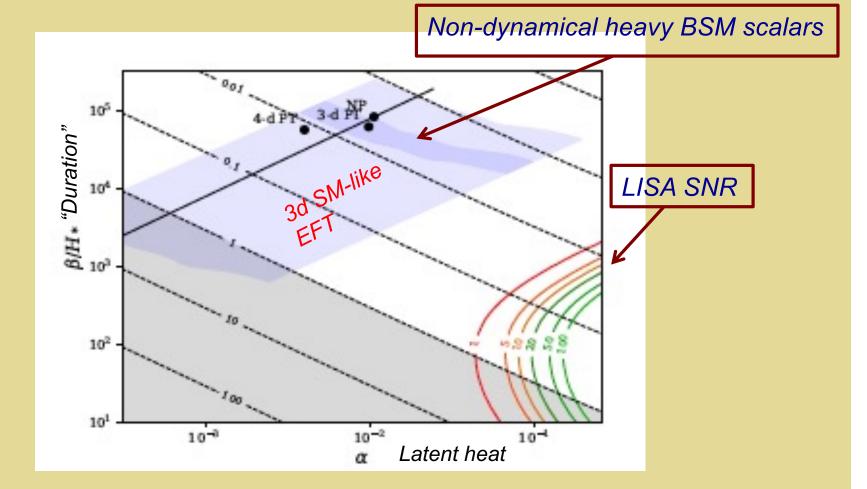
Model Illustrations



Simple Higgs portal models:

- Real gauge singlet (SM + 1)
- Real EW triplet (SM + 3)

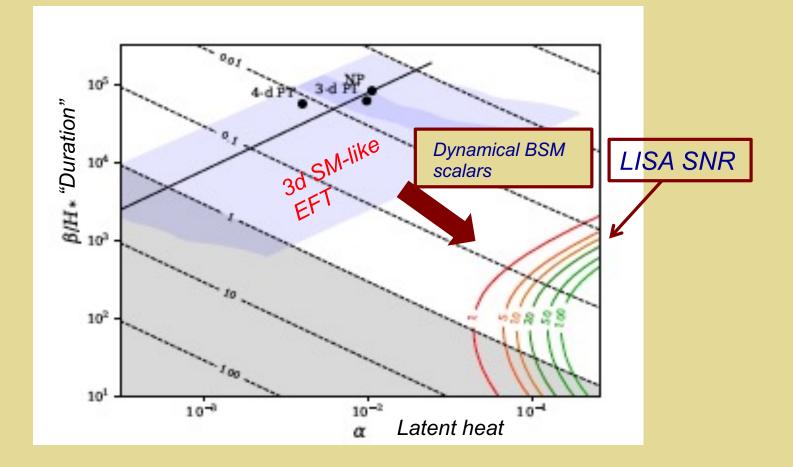
Heavy Real Singlet: EWPT & GW



Gould, Kozaczuk, Niemi, R-M, Tenkanen, Weir 1903.11604

- One-step
- Non-perturbative

Heavy Real Singlet: EWPT & GW



Gould, Kozaczuk, Niemi, R-M, Tenkanen, Weir 1903.11604

- One-step
- Non-perturbative

Other Selected Recent Work

- Gauge invariance in Γ_N (EFT2) : radiative barriers
 - Arunasalam, MJRM: 2105.07588
 - Lofgren, MJRM, Schicho, Tenkanen 2112.05472
 - Hirvonen, Lofgren, MJRM, Schicho, Tenkanen 2112.08912
- RGE at T> 0
 - Gould, Tenkanen 2104.04399
- Nucleation: EFT2
 - Gould, Hirvonen 2108.04377

IV. Outlook - 1

- Determining the thermal history of EWSB is field theoretically interesting in its own right and of practical importance for baryogenesis and GW
- The scale T_{EW} → any new physics that modifies the SM crossover transition to a first order transition must live at M < 1 TeV and couple with sufficient strength to yield (in principle) observable shifts in Higgs boson properties
- Searches for new scalars and precision Higgs measurements at the LHC and prospective next generation colliders could conclusively determine the nature of the EWSB transition

IV. Outlook - 2

- Realizing this opportunity requires a new generation of robust theoretical computations, using EFT & nonperturbative methods, to benchmark perturbative calculations
- There are exciting opportunities for talented and ambitious theorists to make significant contributions to this growing frontier





Temperature Dependence of V(ϕ **)**

"Imaginary time":

$$t \longrightarrow i\tau$$

Im time Greens functions:

$$\mathcal{G}_{\beta}(\tau; \vec{x} - \vec{x'}) \equiv G(0, i\tau; \vec{x} - \vec{x'})_{\beta}$$

KMS:

$$\mathcal{G}_{\beta}(\tau;\vec{r}) = \mathcal{G}_{\beta}(\tau+\beta;\vec{r})$$

Temperature Dependence of V(ϕ **)**

$$\left(\frac{\partial^2}{\partial \tau^2} + \nabla^2 - m^2\right) \mathcal{G}_{\beta}(\tau; \vec{r}) = -\delta^3(r)\delta(\tau)$$
$$\mathcal{G}_{\beta}(\tau; \vec{r}) = \mathcal{G}_{\beta}(\tau + \beta; \vec{r})$$

Solutions:

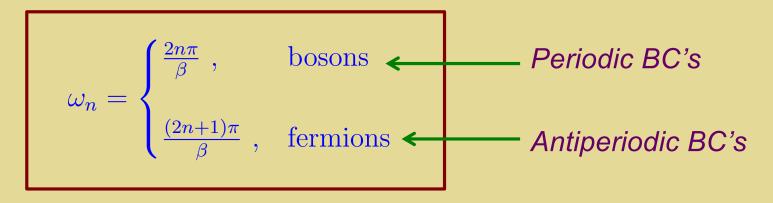
$$\mathcal{G}_{\beta}(\tau;\vec{r}) = \frac{1}{\beta} \sum_{n} \int \frac{d^{3}k}{(2\pi)^{3}} e^{-i(\omega_{n}\tau - \vec{k}\cdot\vec{r})} \mathcal{G}_{\beta}(\vec{k},\omega_{n})$$
$$\mathcal{G}_{\beta}(\vec{k},\omega_{n}) = \frac{1}{\omega_{n}^{2} + \vec{k}^{2} + m^{2}}$$

82

Temperature Dependence of V(\phi)

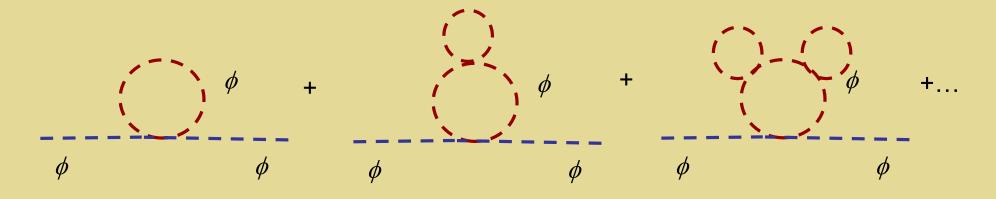
$$\mathcal{G}_{\beta}(\tau;\vec{r}) = \frac{1}{\beta} \sum_{n} \int \frac{d^{3}k}{(2\pi)^{3}} e^{-i(\omega_{n}\tau - \vec{k}\cdot\vec{r})} \mathcal{G}_{\beta}(\vec{k},\omega_{n})$$
$$\mathcal{G}_{\beta}(\vec{k},\omega_{n}) = \frac{1}{\omega_{n}^{2} + \vec{k}^{2} + m^{2}}$$

"Matsubara frequencies"



Power Counting & Resummations

"Daisy" or "ring" resummation



$$egin{aligned} J_B(y) &pprox -rac{\pi^4}{45} + rac{\pi^2}{12} y^2 - rac{\pi}{6} y^3 - rac{y^4}{32} \ln\left(rac{y^2}{a_B}
ight) \ J_F(y) &pprox rac{7\pi^4}{360} - rac{\pi^2}{24} y^2 - rac{y^4}{32} \ln\left(rac{y^2}{a_F}
ight) \ , \end{aligned}$$

- Nonanalytic in ϕ_c
- Origin of barrier in V_{EFF}
- Susceptible to higher order ring diagrams: "screening"