Interpretation of muon (g-2) and W-mass in MSSM

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 - Based on :
 - hep-ph/2203.15710
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THE MUON ANOMALOUS MAGNETIC MOMENT











Muon coupling to magnetic field:

$$\bar{u}(p') \left[\gamma^{\mu} F_1(q^2) + \frac{i}{2m_{\mu}} \sigma^{\mu} \right]$$

Effective Operator:

 $\mu^{\mu\nu}q_{\nu}F_2(q^2) | u(p)A_{\mu}$ $F_2(0) = a_{\mu}$

 $\mathscr{L}_{\text{eff}} = \frac{ea_{\mu}}{2m_{\mu}}\bar{\mu}\sigma_{\mu\nu}\mu F^{\mu\nu}$

<u>New physics</u> contribution

Chiral Enhancement, Large Coupling, Light non-standard particles

 αm , πm_W^2

Muon (g-2) in SUSY

SUSY contributions from Chargino, Sleptons and Neutralino

T. Moroi '96, M.Endo, K. Hamaguchi, S. Iwamoto, Yoshinaga '13

Mass insertion approximation

$$a_{\mu}(\tilde{B}, \tilde{H}, \tilde{\mu}_{L}) = \frac{\alpha_{Y}}{8\pi} \frac{m_{\mu}^{2}}{M_{1}\mu} \tan\beta \cdot f_{N} \left(\frac{M_{1}^{2}}{m_{\tilde{\mu}_{L}}}, \frac{\mu^{2}}{m_{\tilde{\mu}_{R}}}\right)$$
$$a_{\mu}(\tilde{\mu}_{L}, \tilde{\mu}_{R}, \tilde{B}) = \frac{\alpha_{Y}}{4\pi} \frac{m_{\mu}^{2}M_{1}\mu}{m_{\tilde{\mu}_{L}}^{2}m_{\tilde{\mu}_{R}}^{2}} \tan\beta \cdot f_{N} \left(\frac{m_{\tilde{\mu}_{L}}^{2}}{M_{1}^{2}}, \frac{m_{\mu}^{2}}{M_{1}^{2}}\right)$$

SUSY contributions from Chargino-Sneutrino and Smuon-Neutralino loop

SM EW 1 loop :
$$\frac{\alpha}{\pi} \frac{m_{\mu}^2}{M_W^2}$$
.

• SUSY can easily explain anomaly : upper limits on EW super partner masses

K.Hagiwara, K.Ma, S. Mukhopadhyaya '17, T. Yanagida, W.Yu, N.Yokozaki '16, '20, '21 S. AbdusSalam et. al '11, E. Bagnashi et. al '15, P. Cox, C. Han, T. Yanagida '18, '19 '21, M.Endo, K. Hamaguchi, S. Iwamoto, T. Kitahara '20 '21

MSSM , 1 loop :
$$\frac{\alpha}{\pi} \frac{m_{\mu}^2}{M_{SUSY}^2} \times tan\beta$$

THE W-BOSON MASS

Theoretical Prediction of W-mass

 M_W^2 1

One loop correction to the SM:

 $\Delta r = \Delta$

 $M_W^{\rm SM} = 80.357 \pm 0.006 \,\,{\rm GeV}$ PDG 2020 $M_W^{\rm exp} = 80.379 \pm 0.012 \,\,{\rm GeV}$

$$-\frac{M_W^2}{M_Z^2}\right) = \frac{\pi\alpha}{\sqrt{2}G_\mu}(1+\Delta r)$$

$$\alpha \left(\ln \frac{M_f}{M_Z} \right) + \Delta \rho(m_t^2) + \Delta r_{\rm rem} \left(\ln \frac{M_H}{M_Z} \right)$$

Experimental average w/o CDF-II result

Radiative corrections

W-Mass at CDF-II

Tevatron at 1.96 TeV with 8.8 fb⁻¹ $M_W|_{\text{CDF-II}} = 80433 \pm 9.4 \text{ (MeV)}$

$M_W|_{\text{CDF+D0}} = 80427 \pm 8.9 \text{ (MeV)}$

 $M_W|_{\text{CDF+D0+LEP}} = 80424 \pm 8.7 \text{ (MeV)}$

 $M_W^{\rm SM} = 80.357 \pm 0.006 \,\,{\rm GeV}$

One-loop contribution

One loop contribution to Δr can be divided into four classes.

- SUSY contributions of squarks and sleptons to gauge boson self-energies.
- SM contributions from the gauge and Higgs sector that contains the vertex and box diagrams.
- SUSY contributions from neutralinos and charginos in self-energies, vertex diagram and box diagram.

$$\Delta r = \Delta r(M_W, M_Z, m_t, \alpha, \alpha_s, \dots, X)$$

 $X = M_{H^{\rm SM}} \quad (SM),$ $X = M_h, M_H, M_A, M_{H^{\pm}}, \tan\beta, M_{\tilde{f}}, A_f, m_{\tilde{\chi}^{0,\pm}}, \dots$

(MSSM)

 Δr is evaluated in an iterative procedure.

One-loop contribution

Isospin splitting :

Leading contribution induced by the mass splitting

SM CONTRIBUTION

 $\nabla WW(0)$ $\Sigma_T^{ZZ}(0)$ M_W^2 M_Z^2

In MSSM, particularly important for third generation squarks.

Previous study

Parameter	Minimum	Maximu
μ	-2000	2000
$M_{\tilde{E}_{1,2,3}} = M_{\tilde{L}_{1,2,3}}$	100	2000
$M_{\tilde{Q}_{1,2}} = M_{\tilde{U}_{1,2}} = M_{\tilde{D}_{1,2}}$	500	2000
$M_{ ilde{Q}_3}$	100	2000
$M_{ ilde{U}_3}$	100	2000
$M_{ ilde{D}_3}$	100	2000
$A_e = A_{\mu} = A_{\tau}$	-3 $M_{ ilde{E}}$	$3M_{ ilde{E}}$
$A_u = A_d = A_c = A_s$	$-3 M_{ ilde{Q}_{12}}$	$3 M_{ ilde{Q}_{12}}$
A_b	$-3 \max(M_{\tilde{Q}_3}, M_{\tilde{D}_3})$	$3\max(M_{\tilde{Q}_3},$
A_t	$-3\max(M_{\tilde{Q}_3}, M_{\tilde{U}_3})$	$3\max(M_{\tilde{Q}_3},$
aneta	1	60
M_3	500	2000
M_A	90	1000
M_2	100	1000

 $M_1 = 5/3 \, s_{\rm w}^2 / c_{\rm w}^2 \, M_2$

S. Heinemeyer, W. Hollik, G. Weiglein, L. Zeuge '13

Motivation

- \checkmark The lightest electroweakinos and sleptons are important for both observables.
- ✓ Assuming that the neutralino contributes to the observed DM relic density, is there a correlation with specific DM mechanisms?

 \checkmark Is there a correlation between the values assumed by muon (g-2) and MW in the MSSM?

ELECTROWEAK-MSSM

The MSSM

SUSY particles

<u>Slepton Mass Matrix</u>

$$M_{\tilde{L}}^{2} = \begin{pmatrix} m_{l}^{2} + m_{LL}^{2} & m_{l}X_{l} \\ m_{l}X_{l} & m_{l}^{2} + m_{RR}^{2} \end{pmatrix}$$

PARAMETERS

First two gens. $m_{\tilde{l}_1} \sim m_{LL}$ $m_{\tilde{l}_2} \sim m_{RR}$

$$m_{LL}^{2} = m_{\tilde{L}}^{2} + (I_{l}^{3L} - Q_{f}s_{w}^{2})M_{z}^{2}c_{2\beta}$$
$$m_{RR}^{2} = m_{\tilde{R}}^{2} + Q_{f}s_{w}^{2}M_{z}^{2}c_{2\beta}$$
$$X_{l} = A_{l} - \mu(\tan\beta)^{2I_{l}^{3L}}$$

$M_1, M_2, \mu, \tan\beta, m_{\tilde{L}}, m_{\tilde{R}}$

CONSTRAINTS FROM DM AND LHC SEARCHES ON EW-MSSM

Electroweak MSSM at LHC

- **★** EW sector may be hiding the key to new physics.
- Modest production cross section, mass bounds from the LHC comparably weak. \star
- ★ May show up elsewhere : DM experiments, $(g 2)_{\mu}$..

Relevant searches at the LHC

ATLAS [1803.02762] 13 TeV, 36 fb^{-1}

Proper recasting is important **—>** checkMATE

Relevant searches at the LHC

• <u>Slepton pair production</u>

ATLAS [1908.08215]

13 TeV, 139 fb^{-1}

ATLAS 1911.12606

Proper recasting is important **—>** checkMATE

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DM Constraints

A well-tempered bino-wino or bino-higgsino LSP is favorable for chargino co-annihilation while a bino dominated LSP will work for slepton co-annihilation.

<u>Classification based on DM nature</u>

Correct abundance $\Omega_{\rm DM} h^2 = 0.120 \pm 0.001$

Bino (Slepton Co-ann Case-L)

Bino-Wino

Bino (Slepton Coann Case-R)

under-abundant DM requirement follows the (g-2) preferred mass region.

Higgsino

Wino

<u>Muon (g-2)</u>

$\Delta a_{\mu} = (25.1 \pm 5.9) \times 10^{-10}$

Dark Matter Results

Correct (low) Relic abundance.

 $\Omega_{CDM} h^2 = (\leq) 0.120 \pm 0.001$

Direct detection SI bounds from XENON1T

J1T

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New analysis implementation

And, Compressed spectra searches.

Bino-Wino Co-annihilation (Correct abundance) 1000Bino-wino co-annihilation 800 $m_{\chi_1}^{\sim_{\pm}}$ (GeV) 600 $\tilde{\chi}_{1}^{\pm}$ -coannihilation 400 (g-2)_μ spectrum) respectively. (g-2)_{μ}+ Ω h² 200 $(g-2)_{\mu}+\Omega h^{2}+DD$ ***** $(g-2)_{\mu}+\Omega h^{2}+DD+LHC$ 800 1000 200 400 600 $m_{\widetilde{\chi}_1^0}$ (GeV)

100 GeV $\leq M_1 \leq 1$ TeV, $M_1 \leq M_2 \leq 1.1M_1$, $1.1M_1 \le \mu \le 10M_1, \quad 5 \le \tan \beta \le 60,$ 100 GeV $\leq m_{\tilde{l}_L} \leq 1$ TeV, $m_{\tilde{l}_R} = m_{\tilde{l}_L}$.

Upper and lower bounds from $(g - 2)_{\mu}$ and LHC searches (including compressed

NLSP mass upper bound around 750 GeV.

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Bino-Wino Co-annihilation

Additional LHC bounds come from slepton searches.

- Slepton-pair production \rightarrow (2*l* + missing E_T) provides important search channel
- Considerable BR for $\tilde{e}_L(\tilde{\mu}_L) \rightarrow \tilde{\chi}_1^{\pm} \nu_e(\nu_\mu)$ Less no. of signal leptons.

Slepton Co-annihilation: Case-L (Correct abundance) Case-L: SU(2) doublet

 $100 \text{ GeV} \le M_1 \le 1 \text{ TeV}, \quad M_1 \le M_2 \le 10M_1,$ $1.1M_1 \le \mu \le 10M_1, \quad 5 \le \tan \beta \le 60,$ $M_1 \text{ GeV} \le m_{\tilde{l}_L} \le 1.2M_1, \quad M_1 \le m_{\tilde{l}_R} \le 10M_1.$

The left-sleptons and sneutrinos are close in mass to the LSP. NLSP mass upper bound around 750 GeV.

Slepton Co-annihilation: Case-L (Correct abundance)

Additional LHC bounds come from chargino plus heavier neutralino searches.

ATLAS 13 TeV limit

 $(3l + missing E_T)$ exclusion limit weakens

 $BR(\tilde{\chi}_1^{\pm} \to \tilde{\tau}_1 \nu_{\tau}) \text{ and } BR(\tilde{\chi}_2^0 \to \tilde{\tau}_1 \tau), BR(\tilde{\chi}_2^0 \to \tilde{\nu}\nu)$

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Higgsino:

FUTURE DIRECT DETECTION AND LHC CONSTRAINTS WILL BE IMPORTANT FOR THESE SCENARIOS.

(Upper limit from relic abundance)

Higgsino

100 GeV $\leq \mu \leq 1.2$ TeV, $1.1\mu \leq M_1 \leq 10\mu$, $1.1\mu \le M_2 \le 10\mu, \quad 5 \le \tan\beta \le 60,$ 100 GeV $\leq m_{\tilde{l}_L}, m_{\tilde{l}_R} \leq 2$ TeV .

> Chargino-neutralino compressed spectrum searches are important in addition to slepton searches.

> > EUR.PHYS.J.C 81 (2021) 12, 1069 ³¹

FUTURE DIRECT DETECTION AND LHC CONSTRAINTS WILL BE IMPORTANT FOR THESE SCENARIOS.

Wino:

(Upper limit from relic abundance)

Wino

 $100 \text{ GeV} \le M_2 \le 1.5 \text{ TeV}, \quad 1.1M_2 \le M_1 \le 10M_2,$ $1.1M_2 \le \mu \le 10M_2, \quad 5 \le \tan \beta \le 60,$ 100 GeV $\leq m_{\tilde{l}_L}, m_{\tilde{l}_R} \leq 2$ TeV.

> Disappearing track searches are relevant in addition to slepton searches.

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Future prospects (under abundant DM)

Compressed Chargino-Neutralino spectrum at future lepton colliders has high hope. 'Wino and Higgsino Factory'

$$\Omega_{CDM} h^2 \le 0.120 \pm 0.001$$

 $\Delta a_{\mu} = (25.1 \pm 5.9) \times 10^{-10}$

Direct detection SI bounds from XENON1T

ILC-1 TeV reach

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Interpretation of muon (g-2) and MW

FeynHiggsv2.18.1

CDF-II = 80.433 + 0.009

• Heavy EW masses corresponding to low $\Delta a_{\mu}^{\text{MSSM}}$ recovers the SM prediction, the decoupling limit.

• The scenario of Wino DM and slepton co-annihilation (case-L) can give rise to sizable contribution to MW upto 25 and 20 MeV respectively. 80.41 čoann. case-R $\tilde{\chi}^{\pm}$ coann. $(\tilde{B}-\tilde{W})$

- LHC searches.
- contributions are relevant.

• Lightest chargino of mass $m_{\tilde{\chi}_1^{\pm}} \leq 200$ GeV in Wino DM case give rise to the largest value for M_W^{MSSM} while lefthanded smuon less than 250 GeV is favored for slepton co-annihilation in consistent with muon (g-2), DM and

• For wino DM, the vertex and box diagram is important 0 while $f_{0}r_{o}$ the self-energy diagrams $\tilde{\chi}^{\pm}$ coann. $(\tilde{B}-\tilde{W})$

Conclusions

- along with the direct collider limits.
- DM and muon (g-2) constraint put effective upper limit on EW SUSY NLSP masses while LHC limits restrict the mass ranges from below.
- LHC exclusion bound strongly depends on EW gaugino composition. Proper recasting of ATLAS/CMS analysis relaxes the existing bound.
- Searches at future lepton colliders i.e. ILC (1 TeV) will be conclusive.
- The light EW sector consistent with muon (g-2), DM and LHC can contribute to the W-mass up to maximum 25 MeV. The slepton coannihilation (case-L) and Wino DM scenario give the largest contribution to MW.
- If the W-mass anomaly persists in future with a substantially large deviation from the current PDG value, an analysis including the light stop/sbottom sector will be necessary.

* It is possible to constrain the EW MSSM with the help of indirect constraints

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W-Mass and CDF-II Anomaly

Science 376, 170–176 (2022)

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SUSY contributions to $(g-2)_{\mu}$

$$\Delta a_{\mu}(\tilde{W}, \tilde{H}, \tilde{\nu}_{\mu}) \simeq 15 \times 10^{-9} \left(\frac{\tan\beta}{10}\right) \left(\frac{(100 \,\mathrm{GeV})^2}{M_2\mu} \Delta a_{\mu}(\tilde{W}, \tilde{H}, \tilde{\mu}_L) \simeq -2.5 \times 10^{-9} \left(\frac{\tan\beta}{10}\right) \left(\frac{(100 \,\mathrm{GeV})^2}{M_2\mu} \Delta a_{\mu}(\tilde{B}, \tilde{H}, \tilde{\mu}_L) \simeq 0.76 \times 10^{-9} \left(\frac{\tan\beta}{10}\right) \left(\frac{(100 \,\mathrm{GeV})^2}{M_1\mu} \Delta a_{\mu}(\tilde{B}, \tilde{H}, \tilde{\mu}_R) \simeq -1.5 \times 10^{-9} \left(\frac{\tan\beta}{10}\right) \left(\frac{(100 \,\mathrm{GeV})^2}{M_1\mu} \Delta a_{\mu}(\tilde{\mu}_L, \tilde{\mu}_R, \tilde{B}) \simeq 1.5 \times 10^{-9} \left(\frac{\tan\beta}{10}\right) \left(\frac{(100 \,\mathrm{GeV})^2}{m_{\tilde{\mu}_L}^2 m_{\tilde{\mu}_R}^2/2} \right)$$

Endo, Hamaguchi, Iwamoto, Yoshinaga'13

Lattice result

