The central dogma and cosmological horizons

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[ES, Susskind 2201.03603]







Classical black hole [Schwarzschild 1916]



"The war treated me kindly enough, in spite of the heavy gunfire, to allow me to get away from it all and take this walk in the land of your ideas."











Like a gas in a box, black holes have a **temperature** and an **entropy**:

$$S = \frac{k_B c^3}{\hbar} \frac{\text{Area}}{4G}$$

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Quantum black hole

Boltzmann provided atomic description for gas:



Black hole central dogma: from the outside, a black hole can be described in terms of a quantum system with $\dim(\mathcal{H}_{BH}) = \frac{\text{Area}}{4G_N}$, which evolves unitarily. [Bekenstein, Hawking, 't Hooft, Susskind,...]

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Computations of the Page curve are evidence for central dogma. [Penington] [Almheiri, Engelhardt, Marolf, Maxfield] [Penington, Shenker, Stanford, Yang] [Almheiri, Hartman, Maldacena, ES, Tajdini]

JT gravity coupled to CFT with wings



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$$I = I_{grav}[\phi, g] + I_{CFT}[g, \psi_i]$$

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CFT has transparent boundary conditions at AdS_2 boundary.

$$\mathcal{H}=\mathcal{H}_{\rm BH}\otimes\mathcal{H}_{\rm flat}$$

Page curve and island rule

Entropy of Hawking radiation unbounded for eternal black hole:



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Island rule is extension of Ryu-Takayanagi formula in AdS/CFT:

$$S_R = \min_{I} \exp_{I} \left[S_{\text{matter}}(R \cup I) + S_{\text{grav}}(I) \right]$$



Island "encoded" in exterior.

Black hole event horizon





Cosmological event horizon



Cosmic central dogma

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Will focus on de Sitter spacetime in this talk:

$$ds^{2} = -\left(1 - r^{2}/\ell^{2}\right) dt^{2} + \frac{dr^{2}}{1 - r^{2}/\ell^{2}} + r^{2}d\Omega_{d-1}^{2}$$
$$T = \frac{1}{2\pi\ell}, \qquad S = \frac{\ell^{d-1}\operatorname{Area}(S^{d-1})}{4G_{N}}$$

CAUTION!

Little microscopic support for Gibbons-Hawking entropy.

Zero-point entropy difficult to interpret.

Black hole horizon encodes interior; which side does cosmological horizon encode?

Cosmological horizon both more universal and more observer-dependent than BH horizon.

No asymptotic region (at fixed time) with weak gravity, backreaction on system important.

Overlapping islands.

. . .

BREAK



Can gray region in dS_d probe beyond horizon?



Entropy of region in microscopic description given by extremizing generalized entropy $S_{\text{matter}} + S_{\text{grav}}$ with respect to left endpoint.

Need entropy for region at t = 0 in dS_d in Hartle-Hawking state.



Region at t = 0 maps to annulus in plane, use strong subadditivity [Hirata, Takayanagi]:

$$S_0'(R) \ge 0, \qquad S_0''(R) \le 0$$

 $S_0(R \to 1) \sim -\frac{1}{(R-1)^{d-2}}, \qquad S_0(R \to \infty) \approx -2F$



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As r_2 moves rightward, r_1 also moves rightward; violation of entanglement wedge nesting!

Use of dS bifurcate horizon to compute entropy seems prohibited; very different than black hole bifurcate horizon.

(True QES lives right at cutoff r_c .)

Anchor to horizon – AdS analogy

Let's anchor to horizon [Sanches, Weinberg; Nomura, Rath, Salzetta; Susskind]. Which side is encoded?

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Prescription is to find extremal surface on *both* sides of AdS boundary.

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Entanglement wedge in bilayer proposal \implies QES extension.

Anchor to horizon: pure de Sitter



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 $S(H_L) = A/4G$; EW = interior. Horizon is now maximin!

Monolayer theory gives same answer for entropies.

Schwarzschild black hole in de Sitter:

$$ds^{2} = -(1 - 2m/r^{d-2} - r^{2}/\ell^{2})dt^{2} + \frac{dr^{2}}{1 - 2m/r^{d-2} - r^{2}/\ell^{2}} + r^{2}d\Omega_{d-1}^{2}$$





 $S(H_L) = A_{\rm CH}/4G + A_{\rm BH}/4G$; EW = region between BH and CH.



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Monolayer theory gives same answer upon extremizing between BH horizons as well.

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Large-N thermodynamics

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Pattern of higher form symmetry breaking in holographic CFTs, through Eguchi-Kawai mechanism, makes them very similar sometimes $_{\rm [ES \ 16, \ 20]}$

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but not all the time



Model

Model with finite dim \mathcal{H} : Heisenberg antiferromagnet for two qubits

$$H = J\sigma \cdot \tau$$

Entropy for thermal state:



Pick same interval on both horizons in dS_3 . Four possible saddles:



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Geodesics are pieces of infinite family of degenerate length- π geodesics connecting static patch origin r = 0 at t = 0 to its antipodal point



Connected surface dominates at early times, grows to maximum length 2π and disappears; transition to disconnected surface occurs before it disappears. Analog of Hartman-Maldacena transition for dS.

Area of connected surface grows without bound for $dS_{d>3}$.

Summary + Future

Cosmological horizon very different than black hole horizon (minimax vs maximin), does not naively work as a quantum extremal surface.

Anchoring to horizon allows you to use the horizon and its associated entropy without violating entanglement wedge nesting.

Does anchoring to a black hole event horizon make sense? Probe in AdS/CFT.

More general cosmologies cannot be encoded by pair of horizons; how should encoding work?

Microscopic model matching these features.