

A journey through the non supersymmetric landscape

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Motivation

String Theory provides the best candidate for an unified theory of quantum gravity.

- In this framework **supersymmetry** occupies a prominent place, both theoretically and phenomenologically.
- Despite this, we do not observe supersymmetry within the energy ranges of current experiments.
- Therefore it is crucially important to investigate the consistency and the structure of String Theories where supersymmetry is absent, both in spacetime and branes.

Goals of this talk



~~SUSY~~ theories are (perturbative) anomaly free $\xrightarrow{\triangle!}$ Sugimoto: possible brane anomaly

[Alvarez-Gaume, Ginsparg, Moore, Vafa '86/ Dixon, Harvey '86/ Sagnotti '95,'97/ Sugimoto '99/ Dudas, Mourad '01, Sagnotti '02]

1. Verify **anomaly cancellation** via inflow on the branes' worldvolume.
2. Study the **structure of the gauge group** and hints of new dualities to heterotic strings.
3. **Constrain 10d $Sp(16)/\mathbb{Z}_2$ gravitational QFTs** by means of anomaly cancellation.

The $SO(16)^2$ heterotic theory endows a **discrete topological term** in the spacetime action

[Tachikawa, Zhang '24]

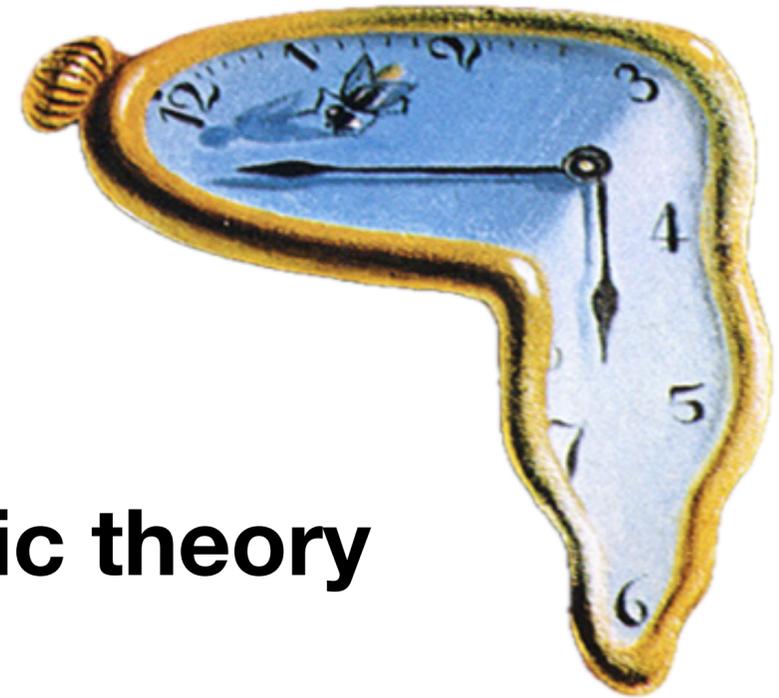
4. Prove the consistency of a **tentative NS5-brane spectrum**.

[Basile, Debray, Delgado, Montero '23]

Present my latest study on **Type II ~~SUSY~~ orbifolds** with zero 1-loop cosmological constant.

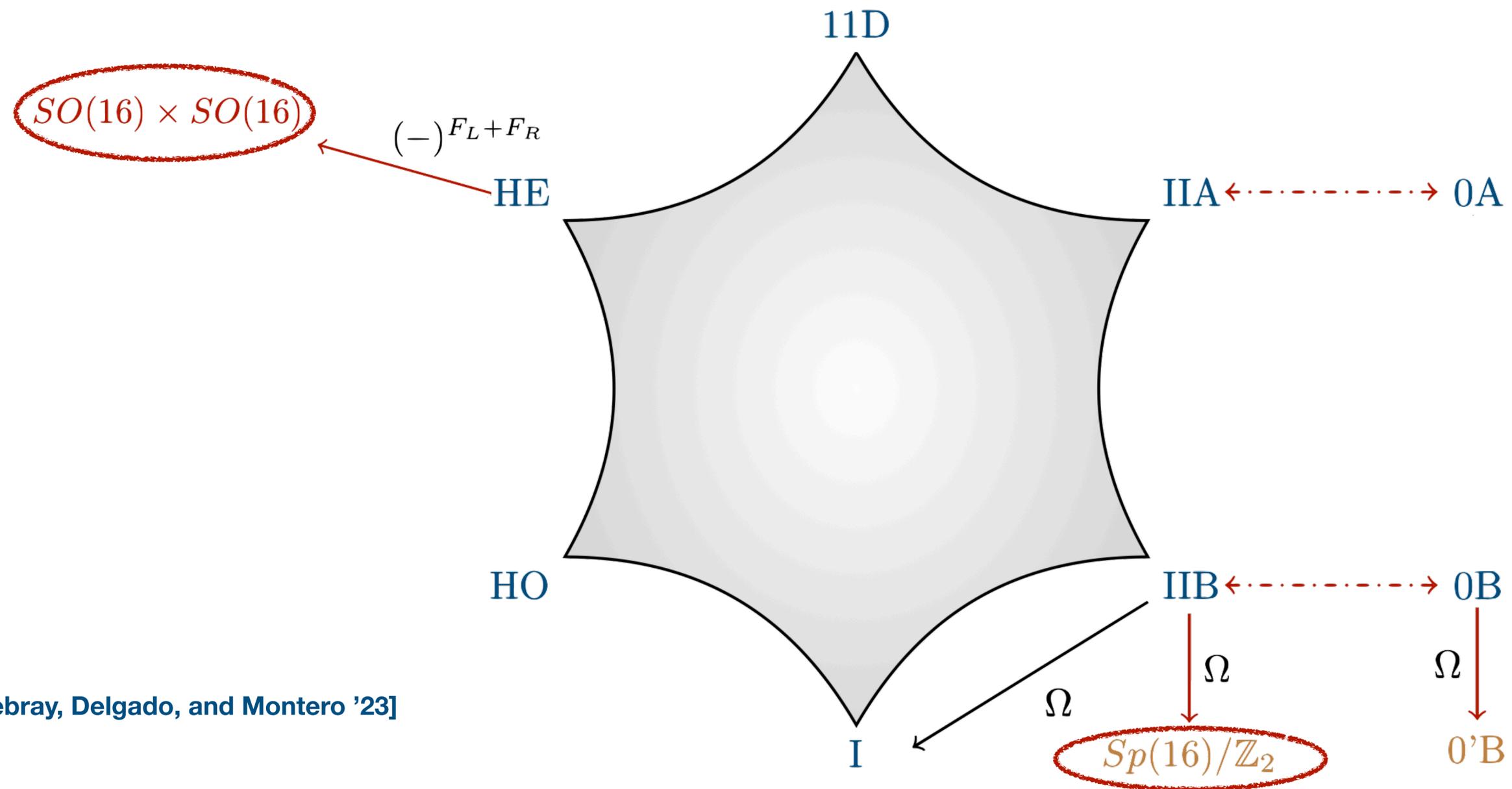
Outline

1. ~~SUSY~~ Strings & Anomalies
2. The structure of the Sugimoto model
3. Discrete topological term in the ~~SUSY~~ heterotic theory
4. ~~SUSY~~ asymmetric orbifolds



1. ~~SUSY~~ Strings & Anomalies

Non-supersymmetric String Theories



[Basile, Debray, Delgado, and Montero '23]

There are 3 non-supersymmetric 10d models without tachyons: Sagnotti, Sugimoto and the heterotic model.

[Sagnotti '95/ Sagnotti '97]

[Sugimoto '99]

[Alvarez-Gaume, Ginsparg, Moore, Vafa '86/ Dixon, Harvey '86]

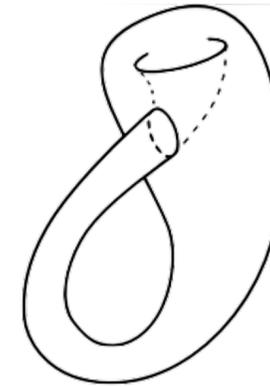
Sugimoto model

The (other) orientifold of Type IIB

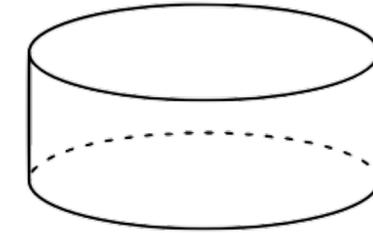
- Anomaly cancellation: $O9 + 16\bar{D}9$

\implies Brane supersymmetry breaking

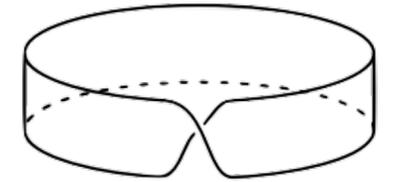
[Antoniadis, Dudas, Sagnotti '99/ Angelantonj '00/ Aldazabal, Uranga '99/
Angelantonj, Antoniadis, D'Appollonio, Dudas, A. Sagnotti '00]



Klein bottle



Annulus



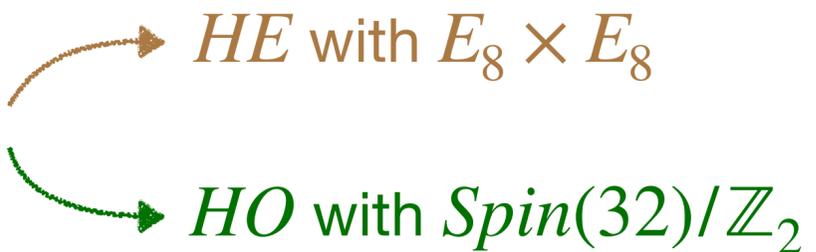
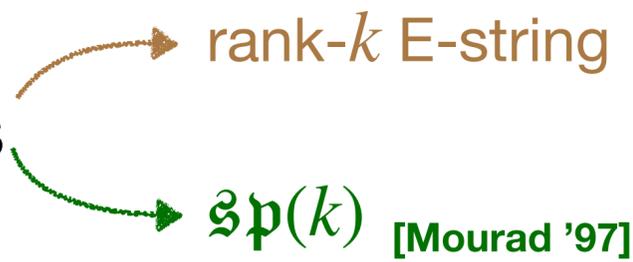
Möbius strip

- Open string sector: $\mathfrak{sp}(16)$ gauge algebra.
- Chiral massless spectrum: Gravitino, Dilatino and $495 + 1$ (anti-symmetric).
- Two dual RR gauge fields B_2 and $B_6 \implies$ charged $D1$ and $D5$ - branes.

~~SUSY~~ Heterotic Theory

The $SO(16) \times SO(16)$ is the only 10d ~~SUSY~~ heterotic theory **without the tachyon**.

The **supersymmetric** heterotic theories are both endowed of

- Spacetime gauge group 
 - HE with $E_8 \times E_8$
 - HO with $Spin(32)/\mathbb{Z}_2$
- NS5-branes 
 - rank- k E-string
 - $\mathfrak{sp}(k)$ [Mourad '97]
- Green-Schwarz mechanism in spacetime

The ~~SUSY~~ chiral spectrum can be understood as a **subtraction** of the SUSY ones:

$$I_{12}^{SO(16)^2} = I_{12}^{HO} - I_{12}^{HE}$$

Green-Schwarz mechanism & anomaly inflow

A general d-dimensional **anomaly polynomial** can be written as

$$I_{d+2} = X_{d+2} + X_2 \wedge X_d + X_4 \wedge X_{d-2} + \dots$$

Terms of the form $\int B_p \wedge X_{d-p}$ may **cancel** the anomaly if the polynomial **factorizes** and

[Green, Schwarz '84/
Sagnotti '92]

$$\delta_\varepsilon B_p = X_p^{(1)} \Leftrightarrow dH_{p+1} = X_{p+2}$$

String structure for p=2

Lower dimensional defects host chiral QFT that may be anomalous.

The **coupling term** contributes to the anomaly with an **inflow**

$$\int_{W_p} B_p \cdot$$



Consistency \Leftrightarrow **matching** between worldvolume anomaly and inflow.

[Callan, Harvey '85/ Blum, Harvey '94/ Freed, Harvey, Minasian, Moore '98/ Lawrie, Schäfer-Nameki, Weigand '17/ Kim, Shiu, Vafa '19/
Angelantonj, Bonnetoy, Condeescu, Dudas '20/ Martucci, Risco, Weigand '23/ Hamada, Vafa '21/ Bedroya, Hamada, Montero, Vafa '21...]

Global anomalies

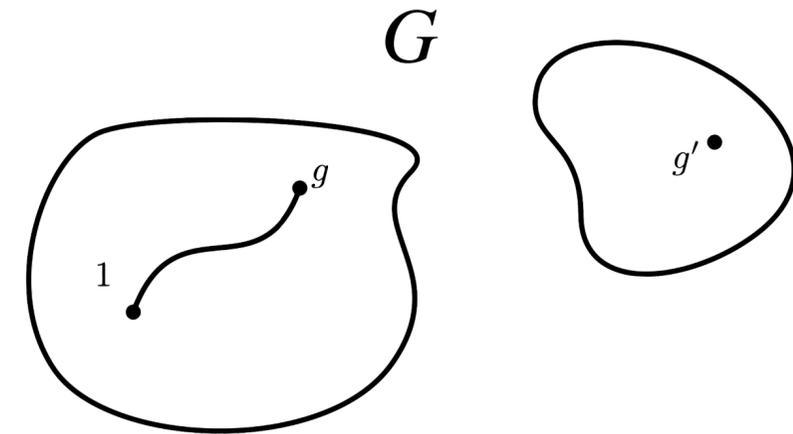
We have discussed perturbative anomalies, related to “small” gauge transformations.

The treatment of “large” gauge transformations is more involved.

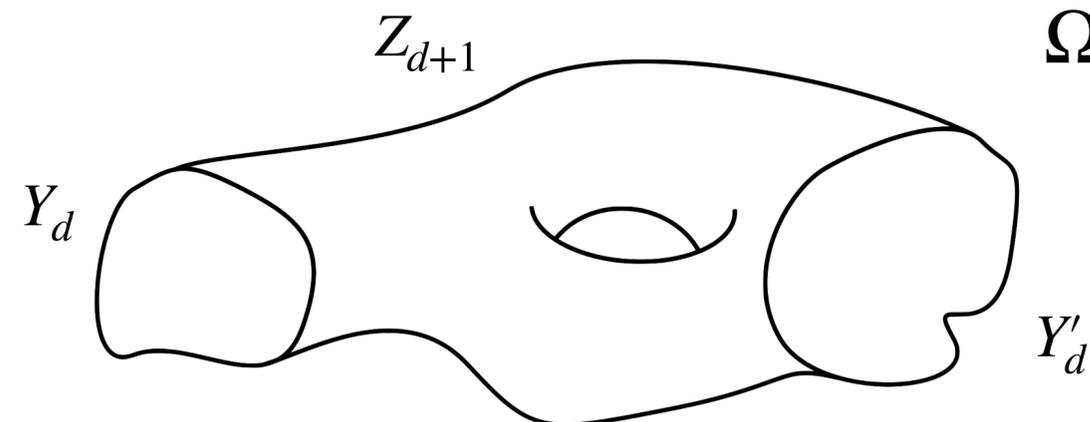
These led to global anomalies, encoded in the homotopy groups

[Witten '82]

$$\pi_d(G)$$

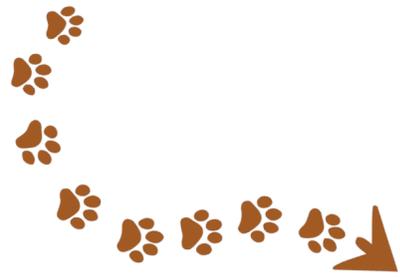


Modern understanding: Dai-Freed anomalies and bordism groups. [García-Etxebarria, Montero '18/ Hsieh '18]



$$\Omega_{d+1}^{spin}(BG) = \frac{\{\text{Mfld w/ properties}\}}{\sim}$$

Anomaly inflow



**Gauge group topology
&
Dualities**



**Bottom-up anomaly
constraints**

2. The structure of the Sugimoto model

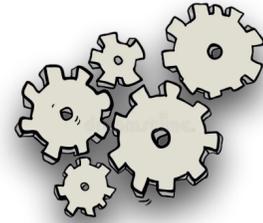
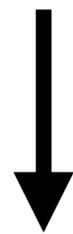
[V.L., L. Lin '24]

Brane anomalies in the Sugimoto model

From **open string amplitudes** we can read the brane's spectrum:

D1

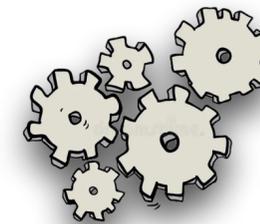
Chirality	$so(8)_T$	$sp(16)_B$	$sp(n)_W$
(+)	8_c	1	
(-)	8_s	1	
(+)	1	32	



$$I_4 = nX_4 \xrightarrow{\text{descent}} -\delta S_{D1}.$$

D5

Chirality	$so(4)_T$	$sp(16)_B$	$so(n)_W$
(+)	2_c	1	
(-)	2_s	1	
(+)	1	32	



$$I_8 = nX_8 + (X_4 + n\chi(N)) \wedge Y_4 \xrightarrow{\text{descent}} -\delta S_{D5}.$$

Anomaly inflow **cancels** $D1$ and $D5$ brane anomaly \implies **Sugimoto model is (perturbative) anomaly free!**



Furthermore, the Sugimoto model has no **global anomalies**.

Investigating the Gauge Group

The 10d spectrum of Sugimoto is invariant under the \mathbb{Z}_2 -center of $Sp(16)$.

Problem:

Do the (bifundamental) representations on the branes change this situation?

Workflow:

- Canonically quantize these modes on the brane.
- Consider the branching products invariant under the worldvolume gauge algebra.
- Check if the corresponding spacetime representations break the \mathbb{Z}_2 center.

D1- and D5-branes' mode quantization

As for the **D0 in Type I**, we quantize the bifundamental on the D1-brane:

$$\begin{array}{ccc}
 \text{[Witten '23]} & & \\
 (32, 2n) \text{ of } \mathfrak{sp}(16)_B \oplus \mathfrak{sp}(n)_W & \xrightarrow{\text{Quantization}} & \{\chi_{i,\alpha}, \chi_{j,\beta}\} = \omega_{ij}\omega_{\alpha\beta} & \xrightarrow{\text{Indices rearrangement}} & \{\chi_I, \chi_J\} = \delta_{IJ}
 \end{array}$$

This is a **Clifford algebra** over the vector **64n** of $\mathfrak{so}(64n)$, associated with the branching:

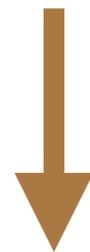
$$\mathfrak{so}(4mn) \supset \mathfrak{sp}(m) \oplus \mathfrak{sp}(n) \quad 4mn \rightarrow (2m, 2n) \quad \text{Physical states} \equiv \text{singlets}$$

Then, the Hilbert space of the D1-brane accounts for the spinor representations.

The computation for the D5-brane modes is analogous.

We decompose the equivalence classes of the spinors using the **projection matrix**.

[Yamatsu '20]



D1 and D5 worldvolume gauge-invariant irreps **must** be also invariant under \mathbb{Z}_2 .

Duality to Heterotic Strings?

We have found strong evidence that the **10d gauge group** of Sugimoto is $Sp(16)/\mathbb{Z}_2$.

In SUSY strings Type I $\xleftrightarrow{\text{S-duality}}$ $Spin(32)/\mathbb{Z}_2$ heterotic.

$$\Lambda_r \longleftrightarrow \Lambda_{cr} \implies \mathfrak{sp}(16) \text{ S-dual algebra } \mathfrak{g} = \mathfrak{so}(33)$$

S duality:

$$\Lambda_c \longleftrightarrow \Lambda_{cc} \implies Sp(16)/\mathbb{Z}_2 \text{ S-dual group } \hat{G} = Spin(33)$$

This algebra/group does not appear as the gauge symmetry in ~~SUSY~~ heterotic theories.

BUT through the embedding $\mathfrak{so}(33) \supset \mathfrak{so}(32)$, $528 \rightarrow 496 \oplus 32$

$Spin(32)$ ~~SUSY~~ heterotic string is a candidate to be the **possible S-dual** of Sugimoto.

[Fraiman, Grana, Parra De Freitas, Sethi '24]

[See also Basile '23]

Bottom-up constraints for $Sp(16)/\mathbb{Z}_2$ theories

Consider the **set** of chiral 10d ~~SUSY~~ $Sp(16)/\mathbb{Z}_2$ theories coupled to gravity.

No 1-form symmetry anomaly. [Apruzzi, Dierigl, Lin '20/
Cvetič, Dierigl, Lin, Zhang '20]

$$I_{12} = n_1 I_{12\mathbf{1}} + n_g I_{12\mathbf{g}} + n_{\begin{smallmatrix} \square \\ \square \end{smallmatrix}} I_{12\begin{smallmatrix} \square \\ \square \end{smallmatrix}} + n_{\begin{smallmatrix} \square & \square \\ \square & \square \end{smallmatrix}} I_{12\begin{smallmatrix} \square & \square \\ \square & \square \end{smallmatrix}}.$$

Green-Schwarz: • $X_{12} = 0 \implies n_{\begin{smallmatrix} \square & \square \\ \square & \square \end{smallmatrix}} = 0$ and $n_1 = 495n_g - 496n_{\begin{smallmatrix} \square \\ \square \end{smallmatrix}}$

$\tilde{X}_4 = \frac{1}{2} [trR^2 - \beta trF^2]$ • Imposing $I_{12} = \tilde{X}_4 \wedge \tilde{X}_8 \implies n_g = n_{\begin{smallmatrix} \square \\ \square \end{smallmatrix}}$ and $\beta = 1$.

⚠ For the structure of $Sp(16)/\mathbb{Z}_2$ this is the only possible factorization $\implies \exists B_2$ and charged 1 and 5 branes

Anomaly inflow: anomaly cancellation requires $n_{\begin{smallmatrix} \square \\ \square \end{smallmatrix}} = 1$

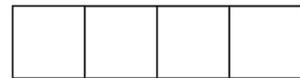
The **only possible** 10d realization of $Sp(16)/\mathbb{Z}_2$ is the Sugimoto model! 🏆

Consistency of tentative $Sp(16)/\mathbb{Z}_2$ models

Extending our search of $Sp(16)/\mathbb{Z}_2$ models by **enlarging** the set of matter representations.

Allow for **higher** rank representations of $Sp(16)/\mathbb{Z}_2$, both in **space-time** and **world-volume**.

$$I_{12} = n_1 I_{12\mathbf{1}} + n_g I_{12\mathbf{g}} + n_{\begin{smallmatrix} \square \\ \square \end{smallmatrix}} I_{12\begin{smallmatrix} \square \\ \square \end{smallmatrix}} + n_{\begin{smallmatrix} \square & \square \\ \square & \square \end{smallmatrix}} I_{12\begin{smallmatrix} \square & \square \\ \square & \square \end{smallmatrix}} + \dots$$



✗ Obstruction to GS mechanism

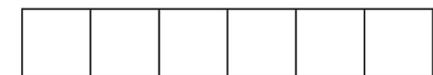


$$\beta = \frac{4095}{3071}$$

✓ Anomaly cancellation in space-time

✗ Obstruction to anomaly inflow onto 1-brane

$$\beta = -2, -4$$



✓ Anomaly cancellation in space-time

✓ Anomaly inflow onto 1-brane

✗ Obstruction to anomaly inflow onto 5-brane

3. Discrete topological term in the ~~SUSY~~ Heterotic theory

[I. Basile, V.L. '25]

Discrete topological term

[Tachikawa, Zhang '24]

Global counterpart of the Green-Schwarz mechanism, classified by

$$\Omega_{10}^{String} = \mathbb{Z}_2 \times \mathbb{Z}_3$$

We consider a ten-dimensional theory with a Green-Schwarz mechanism

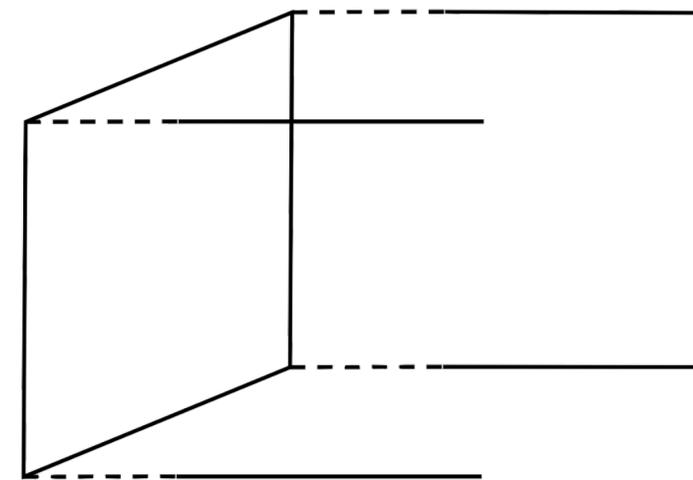
$$\int_{M_{10}} B_2 \wedge Y = \int_{M_{10}} H_3 \wedge CS[Y]$$

If we choose M_{10} as

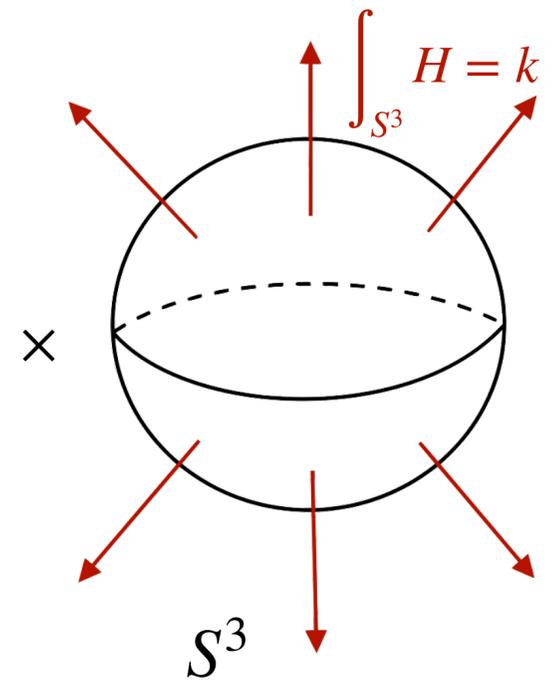
$$S^3 \hookrightarrow M_{10} \longrightarrow N_7 \cong \mathbb{R}_{>0} \times N'_6 \text{ with } H\text{-flux}$$

we recover the inflow onto the NS5-brane

$$k \int_{N_7} CS[Y]$$



$$N_7 \cong \mathbb{R}_{>0} \times N'_6$$



Strategy

Spacetime

Worldvolume

Green-Schwarz term

Anomaly inflow

Discrete

Global anomaly

Matching

Continuous

NS5 anomaly

Local anomaly

The global anomaly

Our aim is to **compute the global anomaly** of the NS5 worldvolume theory.

These objects break the spacetime Lorentz symmetry to

$$SO(1,9) \longrightarrow SO(1,5) \times SO(4) \cong SO(1,5) \times \underbrace{SU(2)_R}_{\text{“R-symmetry”}} \times SU(2)_L$$

The $SU(2)_R$ bundle is **enough** to detect the discrete topological term.

The **global anomaly of $SU(2)$** is encoded by

$$\pi_6(SU(2)) = \mathbb{Z}_{12} \quad \text{but with the spin structure} \quad \Omega_7^{spin}(BSU(2)) = 0$$

Global anomaly on the NS5 \Leftrightarrow Introduction of H_3 with $dH_3 = Z_4$

Evaluation

[Tachikawa, Zhang '24]

We have an $SU(2)$ theory with a 6d Green-Schwarz mechanism:

$$I_8 = Z_4 \wedge W_4 \text{ with } - \int_{N'_6} B_2 \wedge W_4 \quad \text{where} \quad I_8 = I_8^{Spin(32)/\mathbb{Z}_2} - I_8^{E_8 \times E_8}$$

[Ohomori, Shimizu, Tachikawa '14]

The perturbative anomaly is cancelled but the global one remains:

$$X = \eta[\mathcal{M}_7, Z_4 W_4] - \int_{\mathcal{M}_7} H_3 \wedge W_4$$

Using a standard method to compute it, we find

$$X = \frac{t}{12} \quad \text{and} \quad t = 8 \pmod{12}$$



Our approach

We consider a **general** $SU(2)$ spectrum providing the anomaly

$$I_8 = n_1 I_1 + n_2 I_2 + n_3 I_3 + n_{SD} I_{SD}$$

Factorization

+

Matching with the proposed spectrum

[Basile, Debray, Delgado, and Montero '23]



$$\begin{cases} n_{SD} = -1, \\ n_1 = 32 + 13n_3, \\ n_2 = -2 - 8n_3. \end{cases}$$

We compute the global anomaly and find

$$t = 124 + 36n_3 \cong 4 \pmod{12}$$

The analysis with the $SO(16) \times SO(16)$ bundle gives **no further restrictions** and suggests

$$n_3 = 0.$$

3. ~~SUSY~~ asymmetric orbifolds

[M. Montero, V.L., M. Tartaglia '26]

~~SUSY~~ asymmetric orbifolds



A **toroidal orbifold** T^n/Γ is the quotient space of an n -torus by a finite isometry group Γ :

$$g \in \Gamma, \quad g|p_L, p_R\rangle = e^{2\pi i(-v_L \cdot p_L + v_R \cdot p_R)} |\theta_L p_L, \theta_R p_R\rangle \quad (v_L, v_R) \text{ shifts, } (\theta_L, \theta_R) \text{ twists}$$

$$L = R$$

Symmetric orbifolds

[Dixon, Harvey, Vafa, Witten '85/86]

$$L \neq R$$

Asymmetric orbifolds

[Narain, Sarmadi, Vafa '87]

Considering Type II compactifications with $n = 6$, the one-loop cosmological constant is

$$\Lambda_1 \sim \int_{\mathcal{F}} \frac{d^2\tau}{\text{Im}\tau} \mathcal{Z}(\tau, \bar{\tau}) \quad \mathcal{Z}_{\text{orbifold}}(\tau, \bar{\tau}) = \frac{1}{N} \sum_{(f,g)} \mathcal{Z}[f, g] \quad \text{with} \quad fgf^{-1}g^{-1} = 1$$

If each commuting pair (f, g) **preserves some supercharge**, but **not the same** in all sectors

4d ~~SUSY~~ model with $\Lambda_1 = 0$.

~~SUSY~~ asymmetric orbifolds II

Besides some $\mathbb{Z}_2 \times \mathbb{Z}_2$ models, there is a lack of knowledge about these theories.

[Kachru, Kumar, Silverstein '98/ Harvey '98/ Shiu, Tye '98/ Angelantonj, Antoniadis, Forger '99/
Angelantonj, Cardella '04/ Satoh, Sugawara, Wada '15/ Ayoama, Sugawara '20]

We search for orbifold groups Γ with an action on the supercharges such that

1. $\forall g_i \in \Gamma$ has an invariant subspace \implies each sector preserves a supercharge

2. No global invariant subspace + invariant base torus \implies no SUSY

3. $(1, g_i)$ are the only commuting pairs \implies avoid contributions like $\mathcal{L}[f, g] \neq 0$

1. and 2. $\implies \mathbb{Z}_3 \times \mathbb{Z}_3, \mathbb{Z}_4 \times \mathbb{Z}_4$ 3. \implies add shifts \implies anomalies! 

No-go for these models

New models

To find new models we have to drop the requirement of the **invariant base torus**

⇒ check explicitly the **absence of gravitini** in twisted sectors, one by one!



We find **240** $\mathbb{Z}_3 \times \mathbb{Z}_3$ and **1125** $\mathbb{Z}_4 \times \mathbb{Z}_4$ actions.

Explicit examples: have a look at the paper!



You will also find two **non-Abelian** models, $S_3 \times \mathbb{Z}_3$ and D_6 .

We are currently classifying **non-abelian actions**, stay tuned!



Conclusions

- The **Sugimoto model** is consistent and $Sp(16)/\mathbb{Z}_2$ naturally emerges as gauge group.
- The **tentative spectrum** of the NS5-brane gives a **correct topological term**.
- We observed a **different** \mathbb{Z}_3 value computing the DTT from the worldvolume.
- We found **new** ~~SUSY~~ asymmetric orbifolds with $\Lambda_1 = 0$.

Outlook

- Net of dualities and string universality also in the ~~SUSY~~ landscape?
- The subtraction of the spectra is **structural feature** of the ~~SUSY~~ theory?
- **M-theoretical embedding** of the ~~SUSY~~ heterotic theory?
- New models with non-abelian groups? What about Λ_n ?



Thank you!
ありがとう!