

Super-Planckian scattering at the LHC

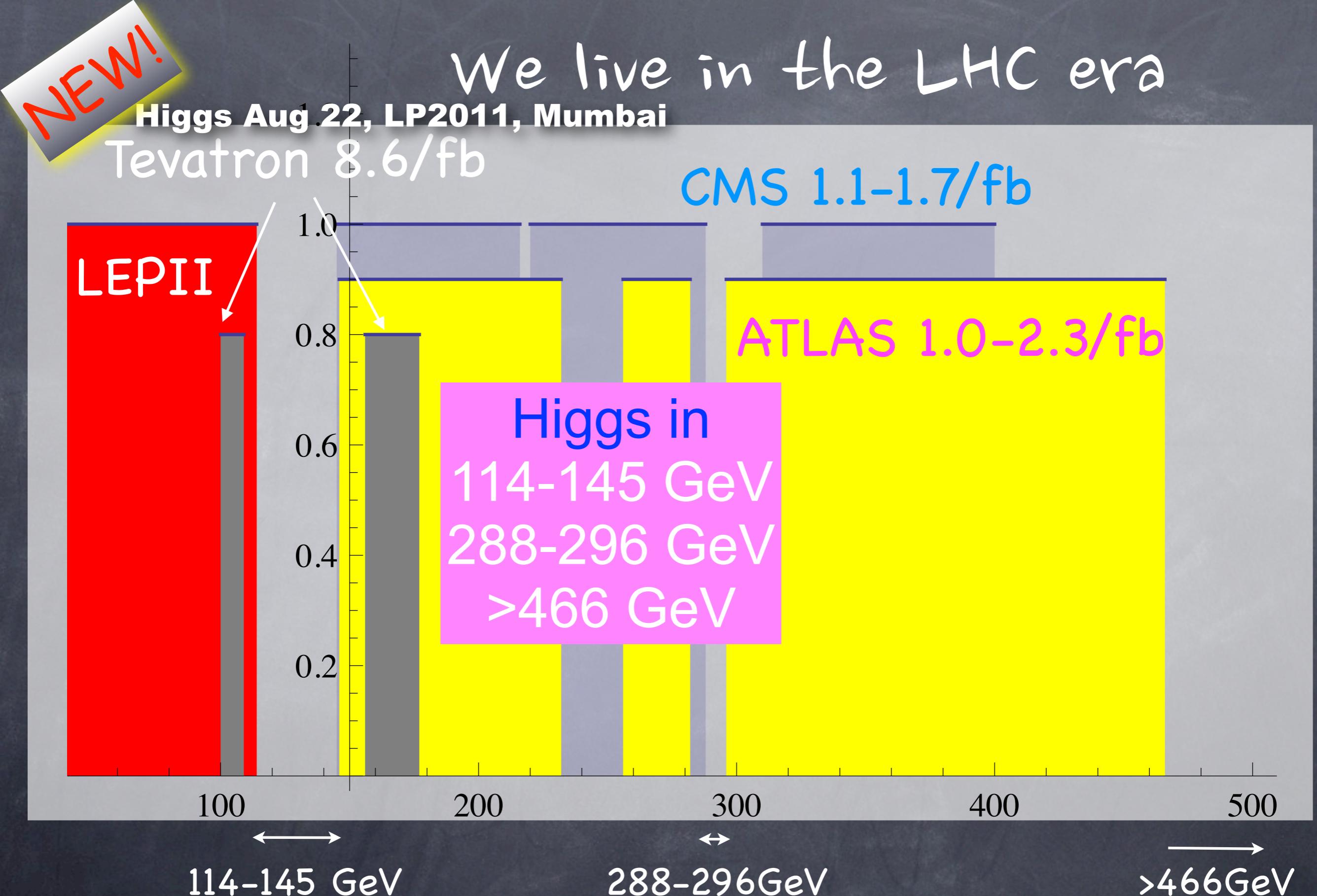
Seongchan Park, CNU

ACP seminar
Aug. 24. IPMU

S. Park, Phys.Lett. B701 (2011) 587-590

Outline

- Large volume compactification,
RS, ADD, CHS
- Trans-Planckian scattering, Black hole
production
- BH Signature
- The recent CMS result on BH, critical
comment/ Future



$m(\text{Higgs}) \sim 100 \text{ GeV}$

experiment

theory

Higgs in
114-145 GeV
288-296 GeV
 $>466 \text{ GeV}$

$$m_H^2 = m_{\text{tree}}^2 + \frac{\lambda_t^2}{16\pi^2} \Lambda_{\text{NP}}^2$$

Hierarchy Problem

A delicate Fine tuning is required if

$$\Lambda \gg 100 \text{ GeV}$$

Some well known (Stringy) solutions

SUSY

$$m_H^2 \sim m_{\text{soft}}^2 \ll \Lambda^2$$

XD

$$\Lambda_{\text{NP}} \sim \text{TeV}$$

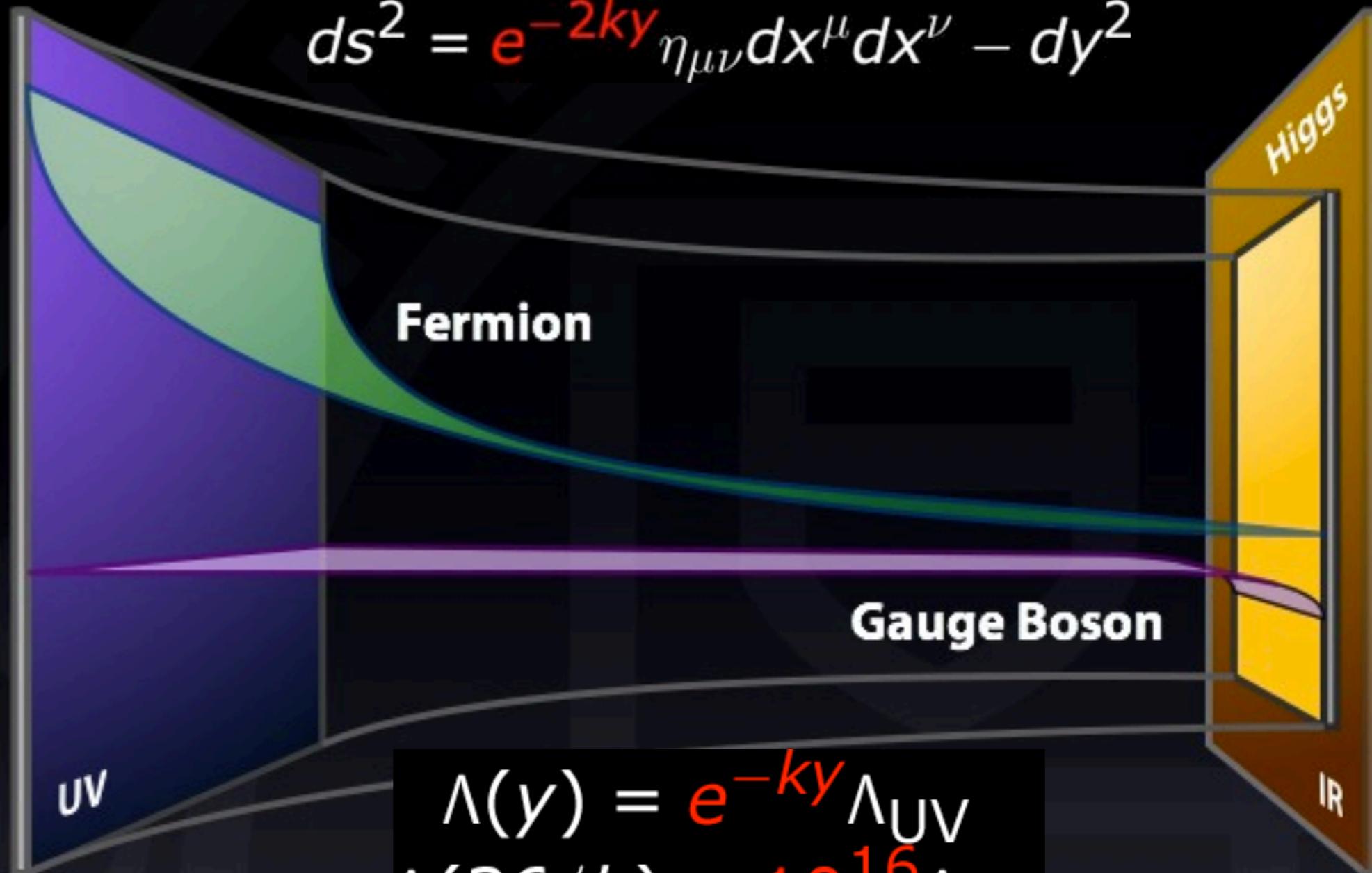
LHC: squarks, gluinos

KK states
strong dynamics

this talk

Randall-Sundrum

$$ds^2 = e^{-2ky} \eta_{\mu\nu} dx^\mu dx^\nu - dy^2$$

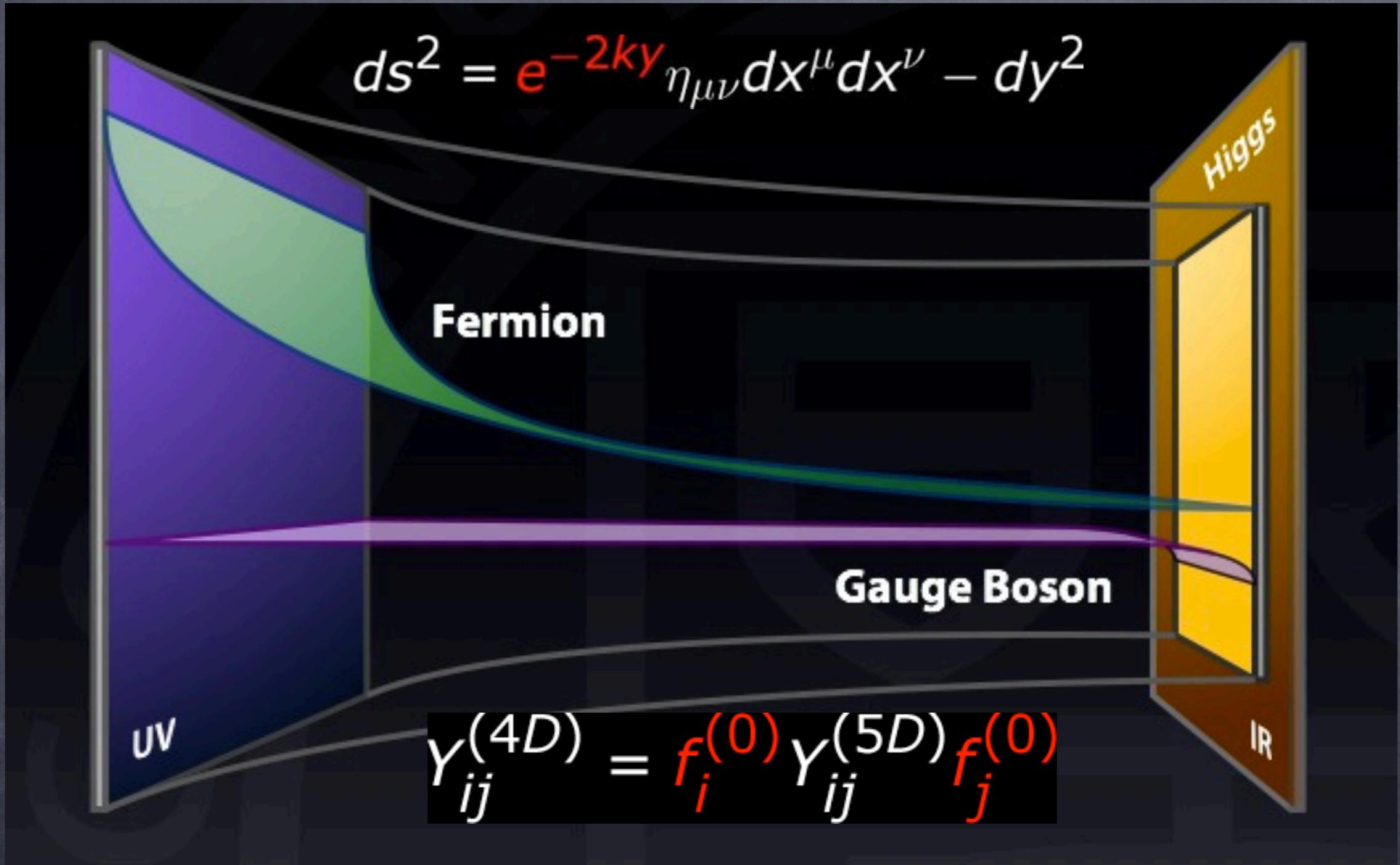


$$\Lambda(y) = e^{-ky} \Lambda_{\text{UV}}$$

$$\Lambda(36/k) \sim 10^{16} \Lambda_{\text{UV}}$$

Randall, Sundrum (99); Davoudiasl, Hewett, Rizzo (99); Grossman, Neubert (00); Gherghetta, Pomarol (00); **Bulk Higgs:** Agashe, Contino, Pomarol (04); Davoudiasl, Lille, Rizzo (05)

Randall-Sundrum



Flavor: Huber, Shafi (03); Burdman (03); Kalil, Mohapatra (04); Agashe, Perez, Soni (04); Chen (05); Agashe, Blechman, Petriello (06); Davidson, Isidori, Uhlig (07); Csáki, Falkowski, Weiler (08) Chen, H.B. Yu (08); Chen, Mahanthappa, F. Yu (09) , ...

Large Torus compactification: ADD

$$S = \frac{M_{\text{Pl}}^2}{2} \int d^4x \sqrt{g} R + \int d^4x \sqrt{g} \mathcal{L}_M$$

$$1/8\pi G = M_{\text{Pl}}^2 \quad [R] = L^{-2} = M^2$$

$$S = \frac{M_D^{2+n}}{2} \int dx^{4+n} \sqrt{g_D} R_D$$

factorizable
geometry

$$\frac{M_D^{2+n}}{2} \int d^n y \sqrt{g} \int d^4x \sqrt{g_4} R + \dots$$



$$\frac{M_{\text{Pl}}^2}{2} = \frac{M_D^{2+n}}{2} [\text{Vol}]$$

High Planck Energy is
due to “large volume”
of extra dimension

What we expect to see-(I)

“small mass gap”

$$\Delta m \sim m_{KK} = k e^{-kL}$$

$$\Delta m \sim [Vol]^{-1/n}$$

Light KK-excitations (\sim Tev):

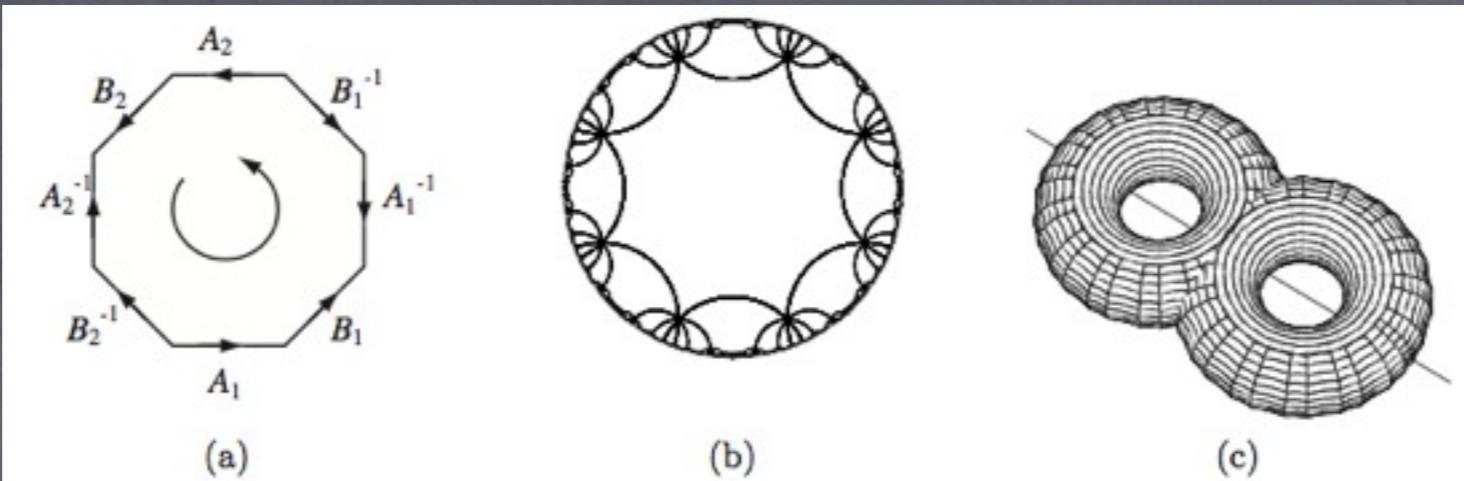
KK-graviton, KK-Z, KK-W, etc..

G+gamma, jj, ll .. @LHC

$$m_{KK} \geq 1.1 \text{ TeV}$$

(NOTE) not always the case e.g. CHS

[Orlando, Park 2010]



Compact Hyperbolic Space

{ Mostow Rigidity theorem ($d > 2$)

[G. Mostow, 1968]

Selberg conjecture

$$k_1 \ell \geq 1/2$$

[A. Selberg, 1965]

proved up to $k_1 \ell \geq \sqrt{171}/784 \approx 0.22$

[W. Luo, Z. Rudnick, and P. Sarnak, 1995]

(by number theory)

Maximally Symmetric Spaces

n-maximally symmetric Einstein space in \mathbb{R}^{n+1}

$$\epsilon_0 X_0^2 + X_1^2 + \cdots + X_{n-1}^2 + \epsilon_n X_n^2 = \epsilon L^2$$

$(\epsilon_0, \epsilon_n, \epsilon)$	curvature
$AdS_n = (-1, -1, -1)$	<0
$dS_n = (-, +, +)$	>0
$S^n = (1, 1, 1)$	>0
$H_n = (-1, 1, -1)$	<0

$$\left. \begin{aligned} R &= \frac{\epsilon}{L^2} n(n-1) \\ R_{\mu\nu} &= R g_{\mu\nu}/n \\ R_{\mu\nu\rho\sigma} &= R(g_{\mu\sigma}g_{\nu\rho} - g_{\mu\rho}g_{\nu\sigma})/n(n-1) \end{aligned} \right\} \propto \epsilon$$

Compact Hyperbolic space

$$H_n = \frac{SO(1, n)}{SO(n)} \xrightarrow{\Gamma \in SO(1, n)} \text{“Lattice”} \quad H_n/\Gamma$$

CHS

H_n/Γ

$$\left\{ \begin{array}{ll} \text{Mostow Rigidity theorem (d>2)} & \\ & [\text{G. Mostow, 1968}] \\ \text{Selberg conjecture} & k_1 \ell \geq 1/2 \\ & [\text{A. Selberg, 1965}] \end{array} \right.$$

proved up to $k_1 \ell \geq \sqrt{171}/784 \approx 0.22$

[W. Luo, Z. Rudnick, and P. Sarnak, 1995]
(by number theory)

Mostow theorem

1. The geometry of a finite volume hyperbolic manifold of dimension greater than two is determined by its fundamental group
2. once we have fixed the curvature and the volume, there are no more moduli.

**The theorem is not valid in $d = 2$ dimensions. In fact, a Riemann surface of genus $g > 1$ (which can always be represented as quotients H_2/Γ , $\Gamma \subset SO(1, 2)$), has a $6(g - 1)$ -dimensional moduli space.

two independent length scales of H_d/Γ

$$\ell_c$$

:local length scale, curvature
appears in EOM
 $\sim 1/[\text{Mass Gap}]$

$$k_1 \ell \geq 1/2$$

$$\ell_\Gamma$$

:global length scale, volume
appears in the quotient group
of the compactification

$$V[H_d/\Gamma] = \ell_c^d e^{\alpha \ell_\Gamma / \ell_c}$$

CHS in an M-theory solution

[Orlando, SCP 006]

$$AdS_{7-d} \times H^d/\Gamma \times S^4$$

$$\left\{ \begin{array}{l} V[H_d/\Gamma] = \ell_c^d e^{\alpha \ell_\Gamma / \ell_c} \\ \text{Large volume} = \text{Hierarchy solved} \\ \text{No light KK-state expected} \\ k_1 \ell \geq 1/2 \end{array} \right.$$

Compact hyperbolic extra dimensions: a M-theory solution and its implications for the LHC.
Domenico Orlando, Seong Chan Park (Tokyo U., IPMU). IPMU10-0096. Jun 2010. 17 pp.
Published in JHEP 1008 (2010) 006
e-Print: arXiv:1006.1901 [hep-th]

What we expect to see (2)

“Low scale gravity”

$$M_D < \sqrt{s}$$

Trans-Planckian scattering

*Gravity dominance [^tHooft 1986]

*Black hole production [T. Banks, Fischler 1999],
[Dimopoulos, Landsberg 2001], [Giddings Thomas 2002]

LHC

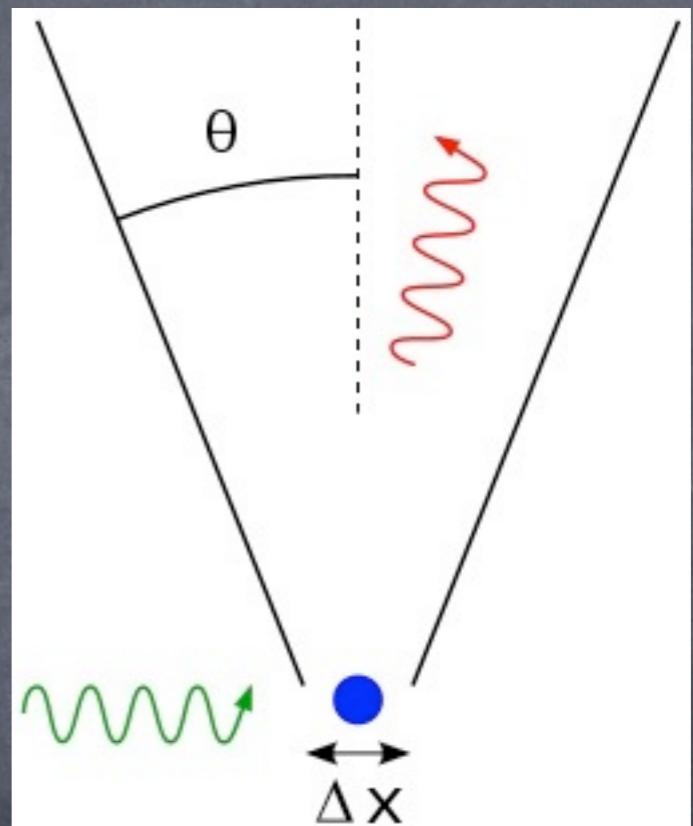
A super-Planckian machine?

- $E = 7 \text{ TeV}$ in pp frame (will be 14 TeV in 2 yrs)
- $P = \{u, d, s, c, b, g\}$
- Partons carry a fraction of 7 TeV (pdf suppression): $E_{ij} = E F_{i/p} F_{j/p}$. Say, $E/5$ of the energy is available.
- Possible to have $E > \Lambda \sim \text{TeV}$

Heisenberg's Microscope

If one wants to be clear about what is meant by "spattering energy," for example finding greater than one has to specify definite experiments by which the "position of an electron" can be measured; otherwise this term has no form.

G. 't Hooft 1987
--Heisenberg, in uncertainty paper, 1927



a baby black hole forms

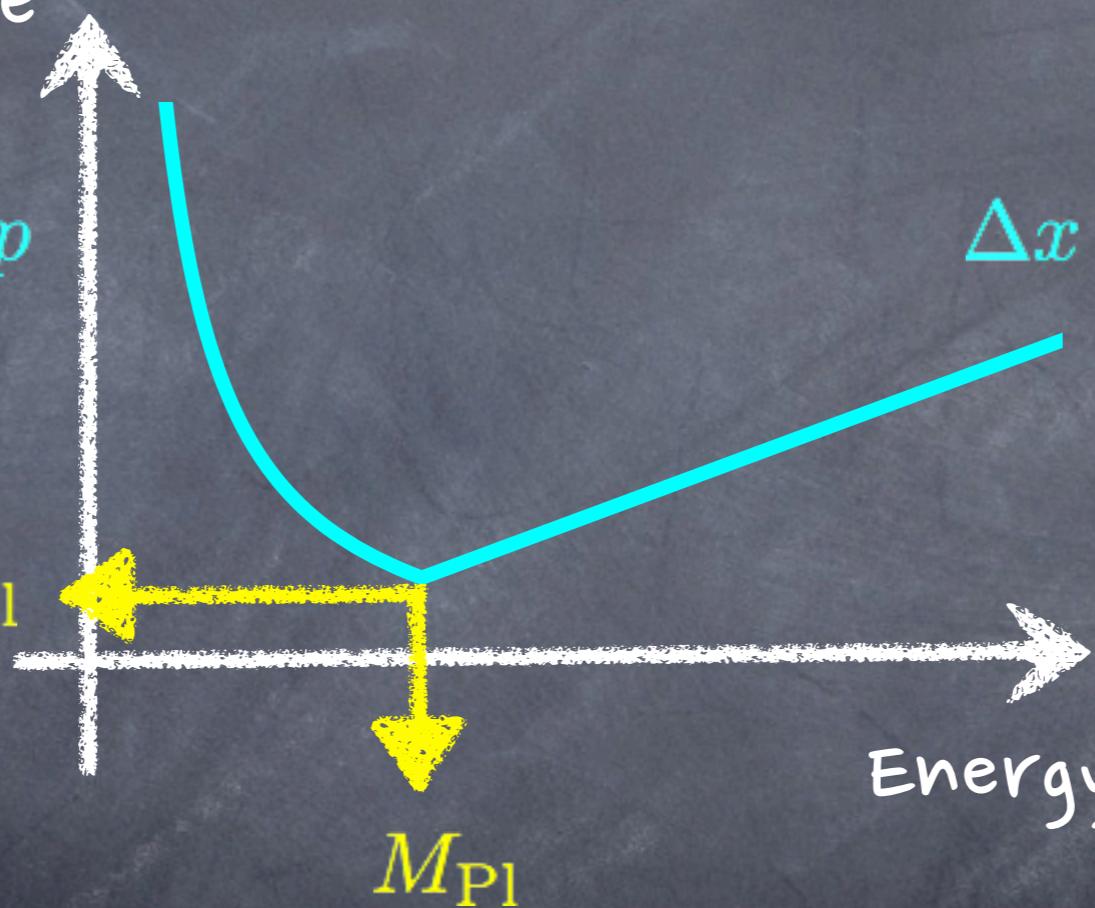
$$\Delta x \sim \lambda \sim \hbar / \Delta p$$

't Hooft-Heisenberg's uncertainty principle

Probing distance

$$\Delta x \sim \hbar / \Delta p$$

$$\ell_{\text{Pl}} \sim 1/M_{\text{Pl}}$$



$$\Delta x \sim r_{\text{BH}}(E)$$

Energy/Momentum

$$[x, p] = i(1 - \alpha p)$$

T. Yoneya (1989)

Trans-Planckian Scattering

initial state



$$E_{\text{CM}} > M_P$$

what's the final state?

Trans-Planckian Scattering

A naive effective field theory
approach fails

$$\mathcal{L}_{\text{eff}} \ni \psi\psi\psi\psi/M_P^2$$

becomes non-unitary at $E > M_P$

Relevant scales

$$\ell_{\text{de-Broglie}} \sim 1/E$$

$$\ell_{\text{Planck}}$$

$$\ell_{\text{EventHorizon}} \sim (GE)^{1/(D-3)}$$

$$b \sim J/E$$

$$\ell_c$$

Black hole forms

$$b < \ell_H$$

$$\ell_H \gg \ell_P$$

$$b > \ell_H \gg \ell_P$$

$$\ell_c > b > \ell_H >> \ell_p$$

Trans-Planckian scattering

[D. Gross et.al. 1987, 1988]

[G. Veneziano 1992, 1993, 2010]

[R. Rattazzi 2004]

$$\sqrt{s} \gg M_D$$

$$b \gg 1/M_P$$

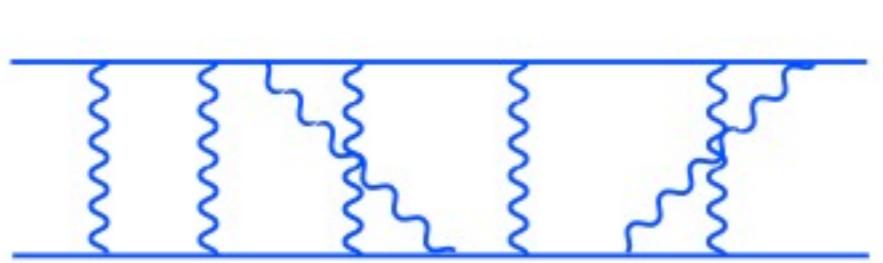
: small QG effect (R/M_P^2)

Born approxima-

Eikonal approxim-

black hole formation

cal, quasi-quantum description holds!

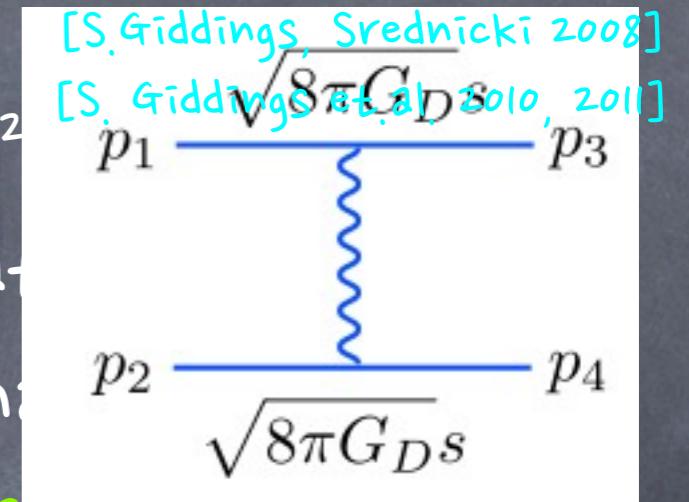


$$b \sim 1/M_P$$

: QG [string] domain
non-perturbative

Not much things we know

2->2 dominance (?)



Small curvature perturbation is valid for trans-Planckian BH

Einstein-Hilbert

$$S = \int d^4x \sqrt{g} \frac{M^2}{2} \left(R + O\left(\frac{R^2}{M^2}\right) \right)$$

↑
↓
includes higher derivative terms (>2)

$$O(R^2) = aR \frac{2R}{M^2} + bR \frac{\mu\nu\rho\sigma}{r_H^2 M^2} R_{\mu\nu\rho\sigma} + cR \frac{\mu\nu}{M^2} R_{\mu\nu}$$

$(a, b, c) = (1, 1, -4)$ Gauss-Bonnet
 $r_H \sim E/M^2$
Ghost-Free, D>4 [B. Zwiebach 1985]

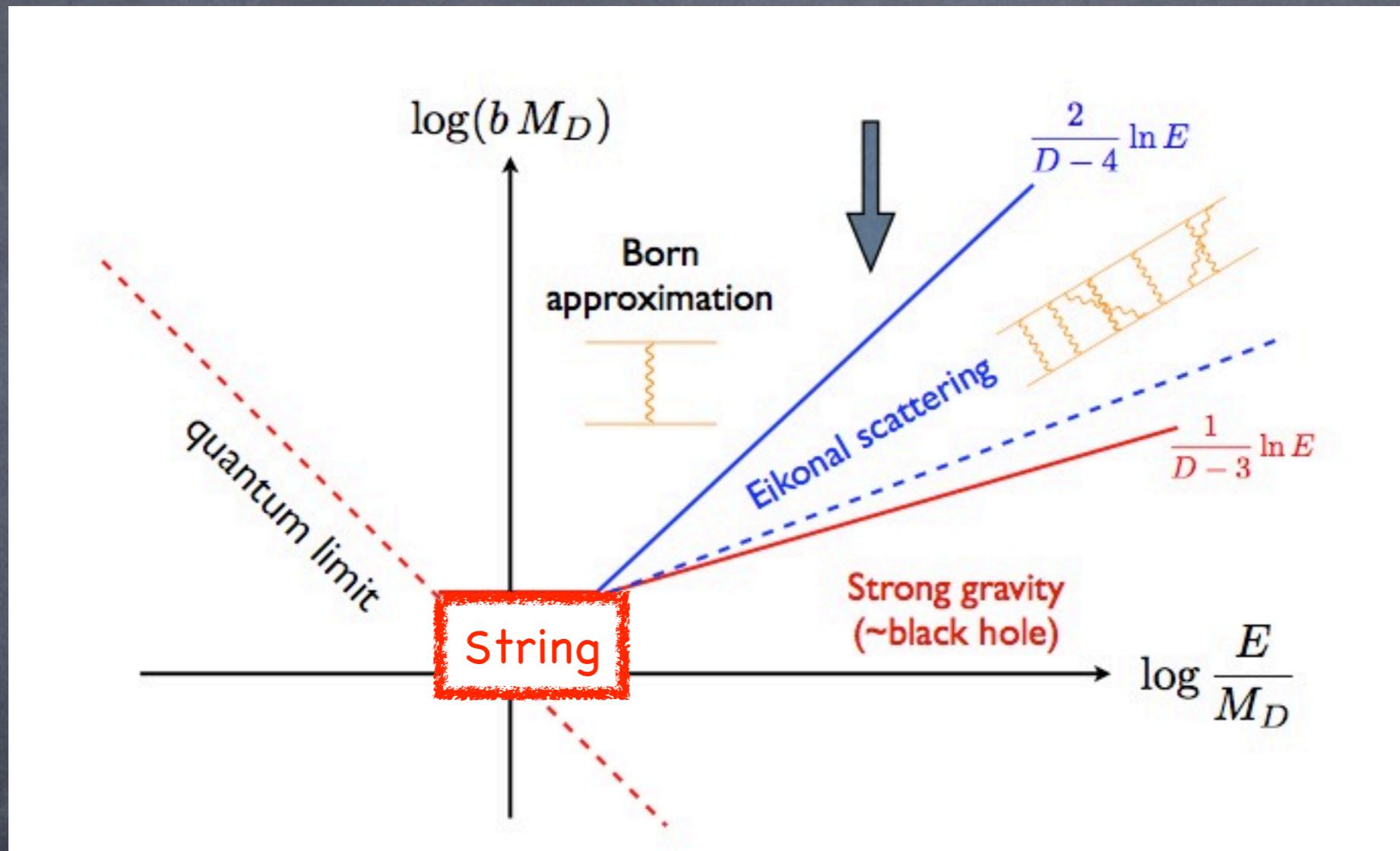
Heterotic string
6D Calabi-Yau compactification of M-theory

for BH in $D > 4$

$$S_G = \int d^D x \sqrt{-g} \frac{M_D^{D-2}}{2} R [1 + \Delta],$$

$$\Delta \sim \frac{R}{M_D^2} \sim \frac{1}{r_{BH}^2 M_D^2} \sim \left(\frac{M_D}{M} \right)^{\frac{2}{D-3}}$$

if $M \gg M_D$, QG correction is small



validity domain of semiclassical approximation

Higher dimensional

$$\ell_D \ll r_{\text{BH}} \ll r_c$$

QG irrelevant

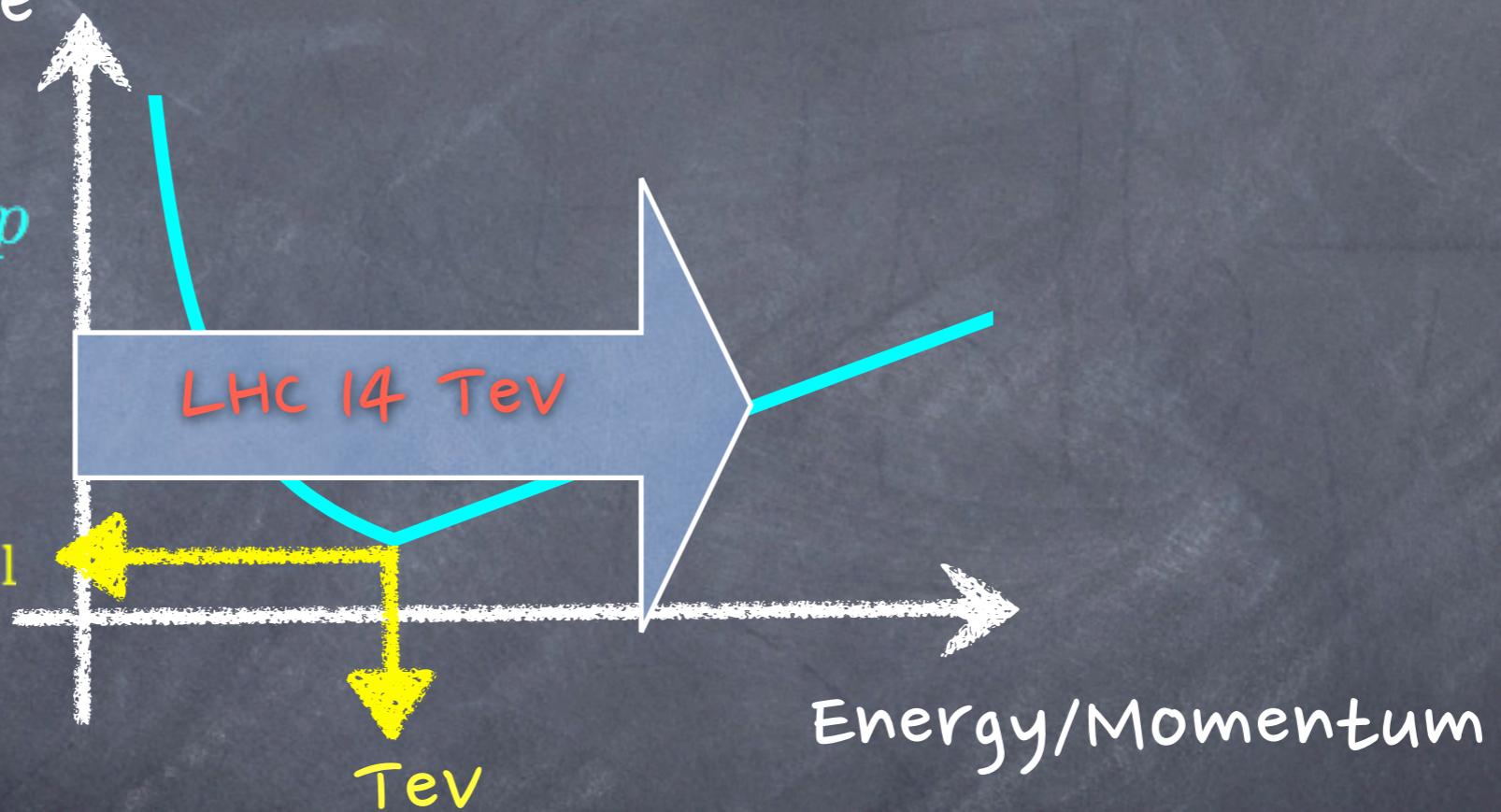
- Reliable calculation available
- Large entropy (= Large surface area)

BH @ LHC

Probing Distance

$$\Delta x \sim \hbar / \Delta p$$

$$\ell_{\text{Pl}} \sim 1/M_{\text{Pl}}$$



BH(D>4) Signal

[Ida-Oda-SCP 2003,2005,2006]

- BH decays through Hawking radiation: Thermal radiation with Large multiplicity (large entropy >10).
- $T > \text{TeV}$, all kinds of SM particles will come out.
- Doesn't depend on the flavor but spin of radiated particle; details are determined by Greybody factor. (vector>spinor>scalar)
- gluon > quark (>photon) > lepton, neutrino> Higgs

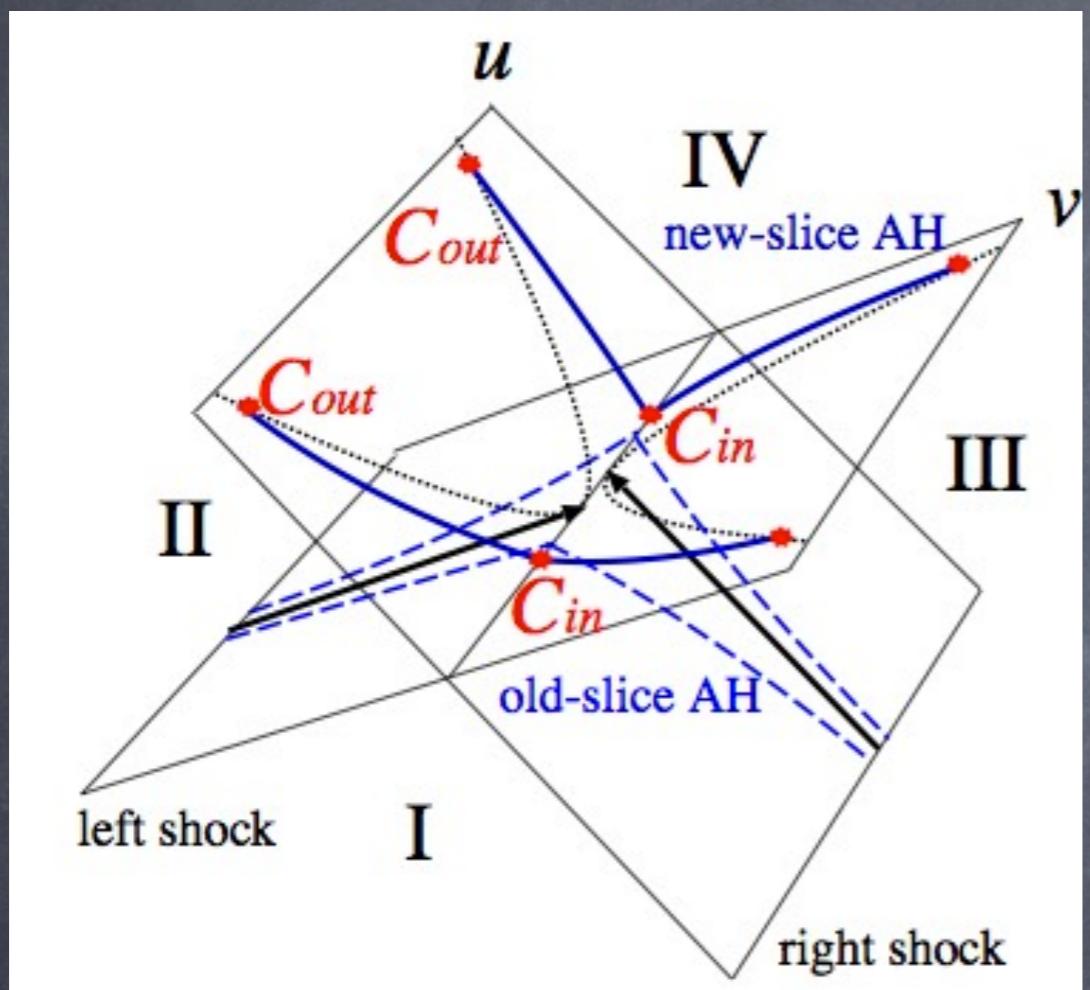
MC for BH events

- ⦿ Two MC simulation codes are used by the LHC
 - **BlackMax2.0** (case western, SUNY buffalo, oxford, London): CMS
 - **CHARYBDIS2.0** (cambridge): ATLAS
- ⦿ Both are based on [Yoshino et.al. 2003, 2005] (for production), [Ida-oda-SCP 2003, 2005, 2006] (for decay).
- ⦿ valid in **semi-classical** domain.

Yoshino-Rychkov

[Yoshino-Nambu 2003]

[Yoshino-Rychkov 2005]



*initial:
infinitely boosted BH
(Aichelberg-Sexl sol.)

*final:
formation of AH

*Theorem: BH forms when AH forms

*An active topic in NR

[Pretorius et.al. 2008, 2010]
[Shibata et.al. 2008, 2010, 2011]

Ida-Oda-SCP

emission rate

$$-\frac{dE_{s,l,m}}{dt} = \frac{g_s}{2\pi} \int d\omega \frac{s \Gamma_{lm}\omega}{e^{\omega - m\Omega/T} + 1}$$

[Ida*-Oda-Park* 2003,2005,2006]



Observer a
distance
"black hole"
looks grey"



Search for microscopic black hole signatures at the Large Hadron Collider[☆]

CMS Collaboration*

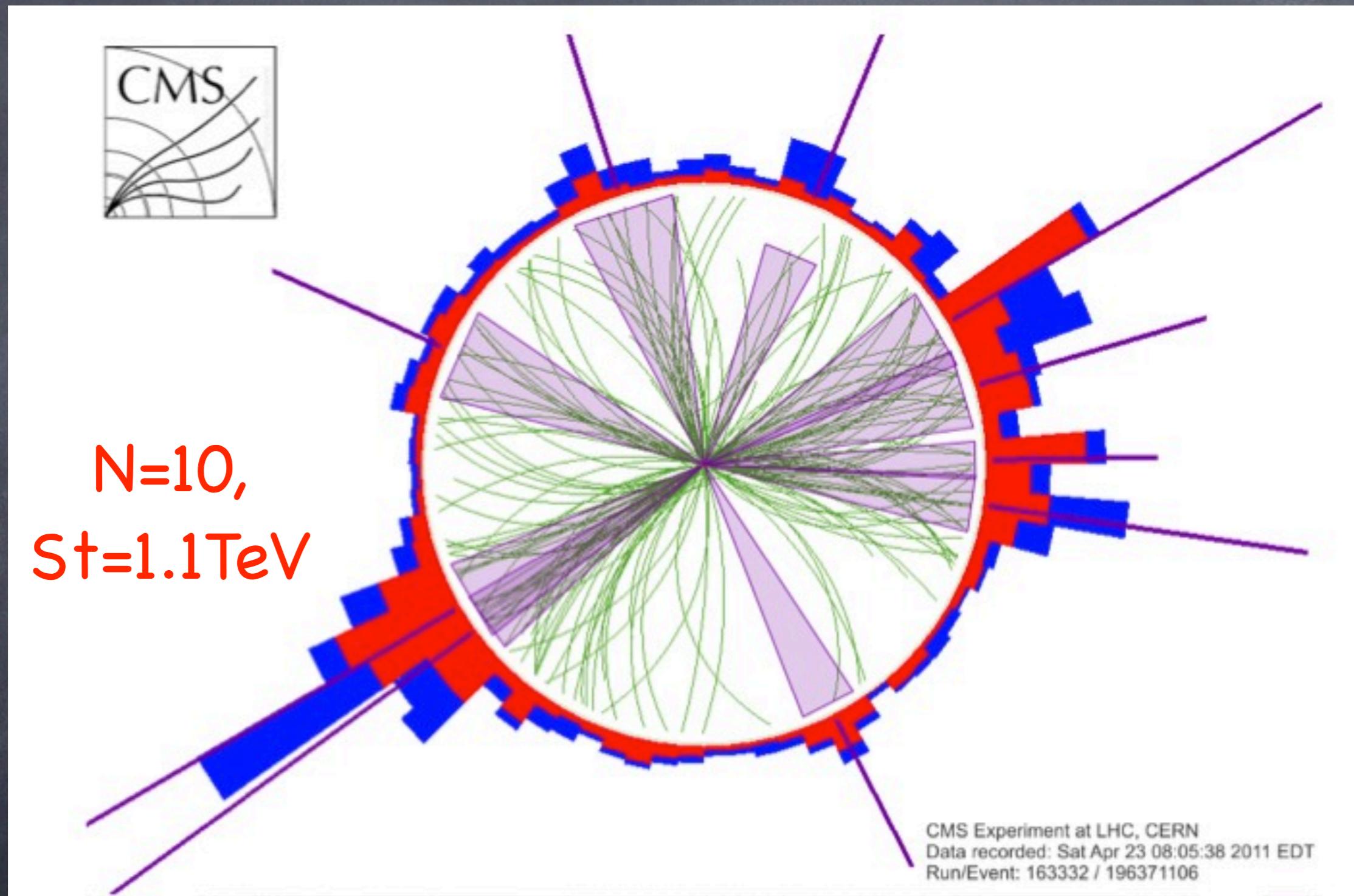
CERN, Switzerland

A B S T R A C T

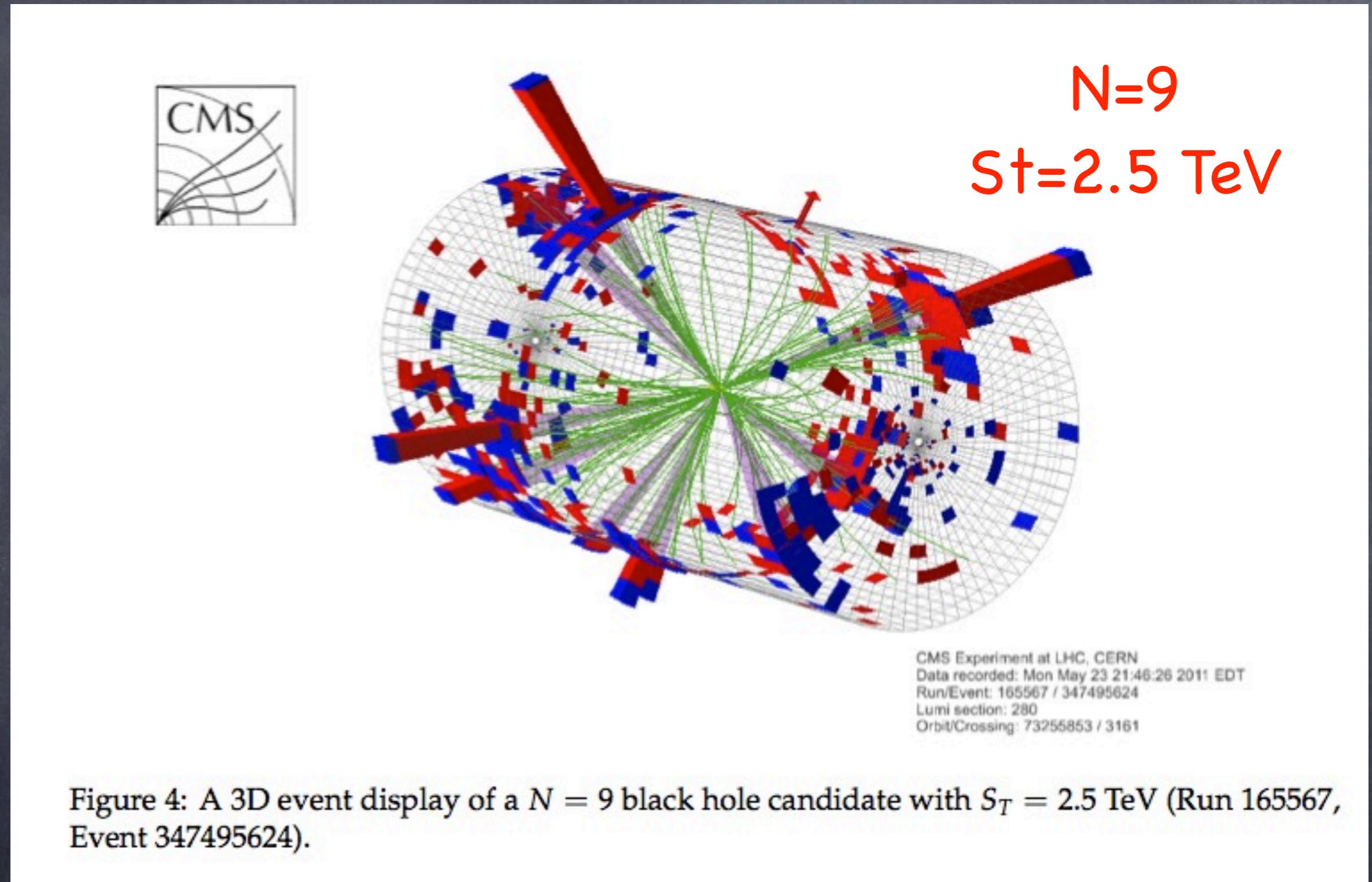
A search for microscopic black hole production and decay in pp collisions at a center-of-mass energy of 7 TeV has been conducted by the CMS Collaboration at the LHC, using a data sample corresponding to an integrated luminosity of 35 pb^{-1} . Events with large total transverse energy are analyzed for the presence of multiple high-energy jets, leptons, and photons, typical of a signal expected from a microscopic black hole. Good agreement with the standard model backgrounds, dominated by QCD multijet production, is observed for various final-state multiplicities and model independent limits on new physics in these final states are set. Using simple semi-classical approximation, limits on the minimum black hole mass are derived as well, in the range 3.5–4.5 TeV. These are the first direct limits on black hole production at a particle accelerator.

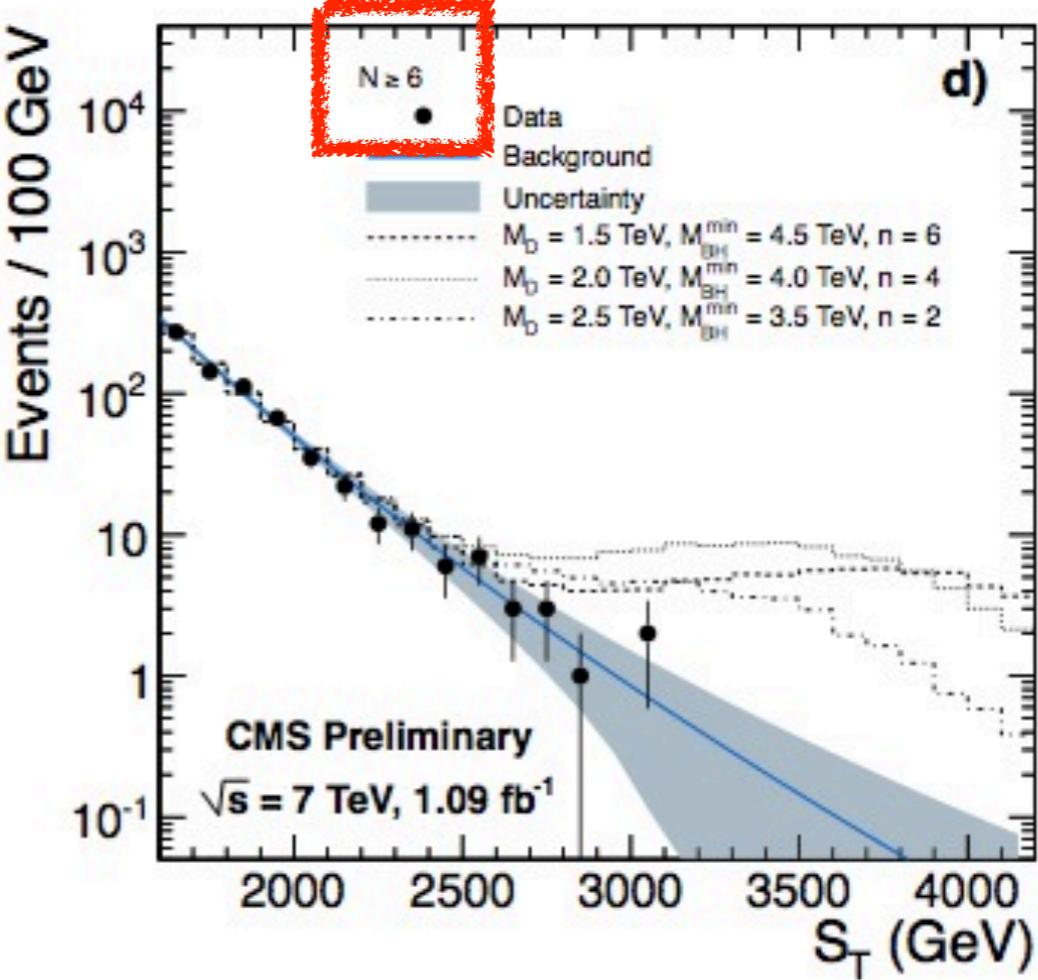
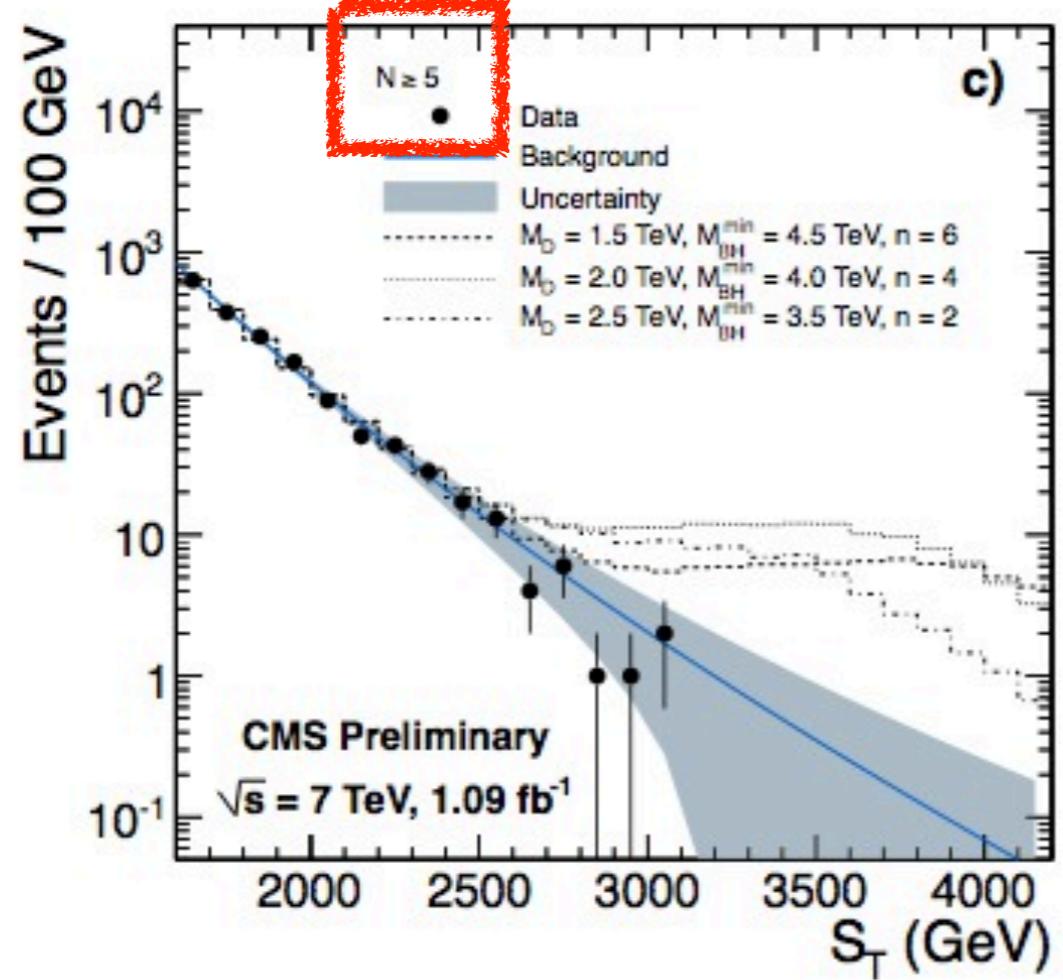
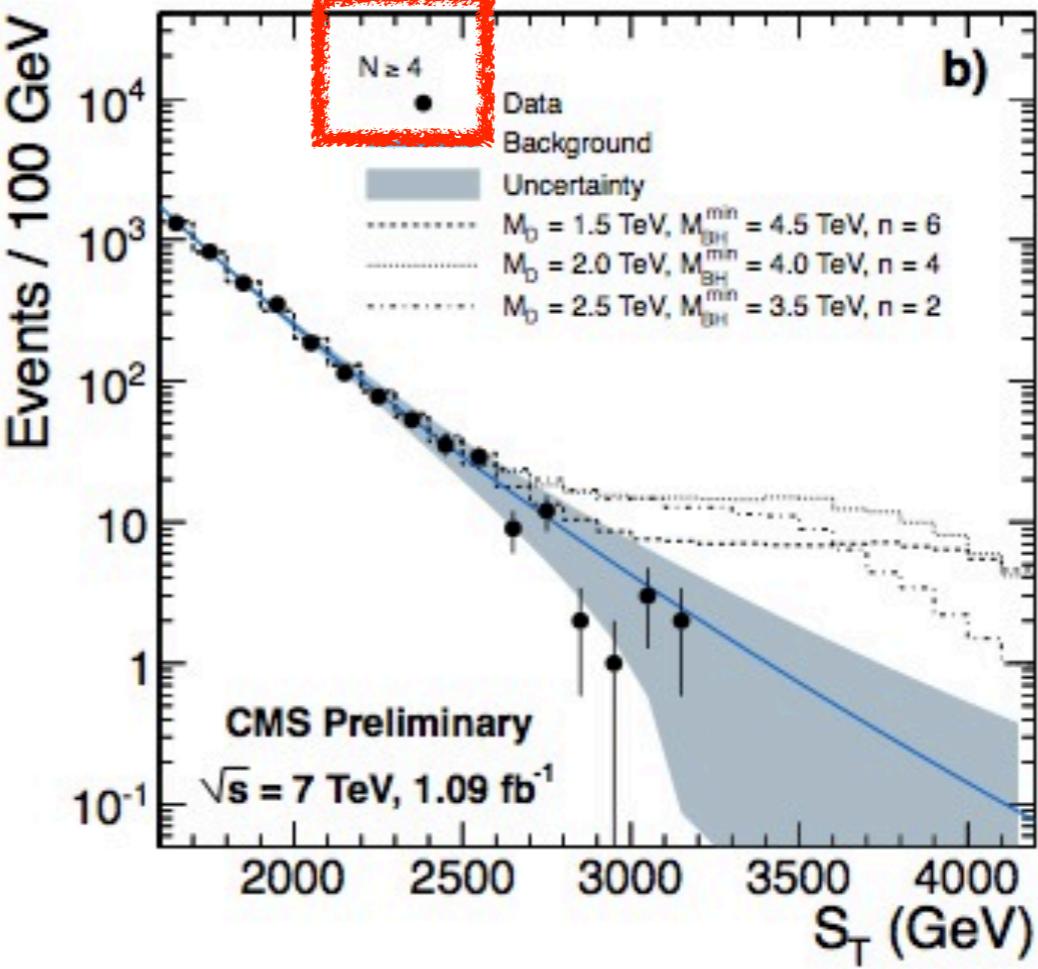
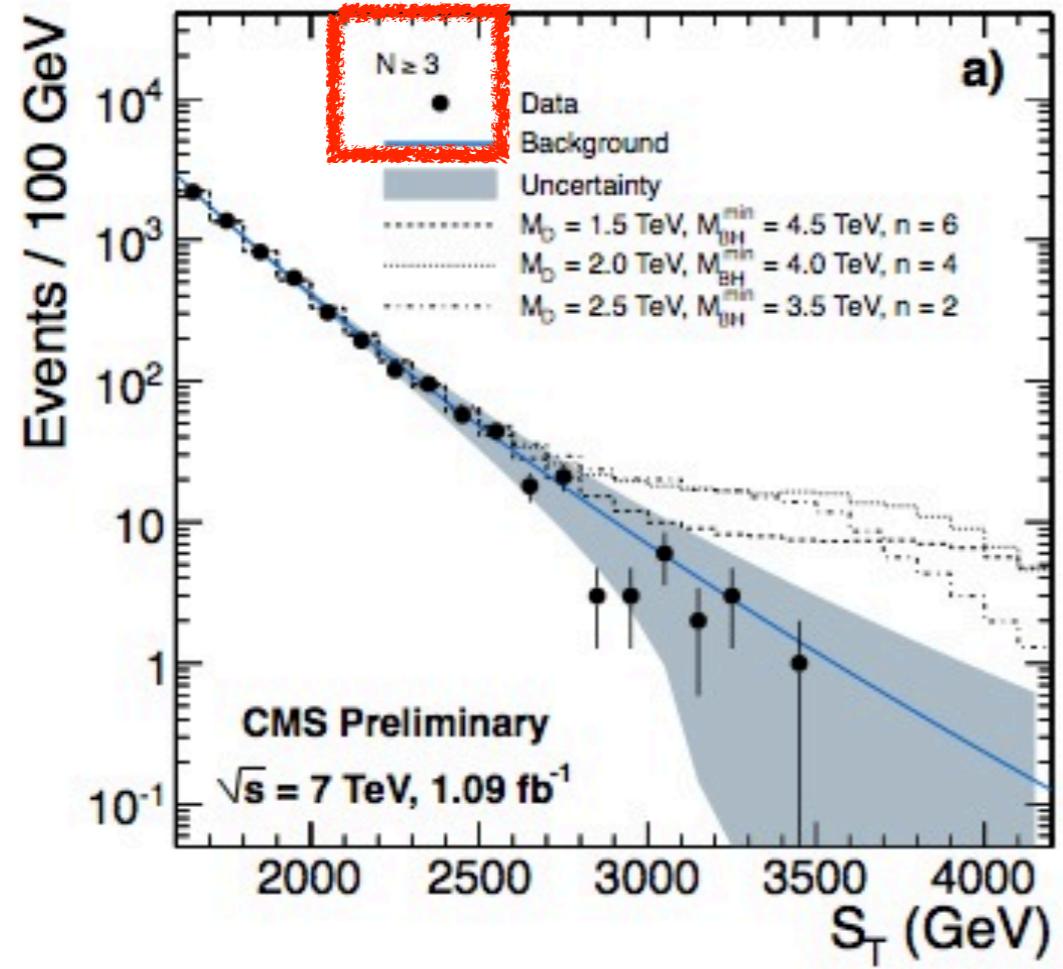
**also see LHC analysis
(with 1.1/fb) 2011 July 21 @EPS
(with 2/fb) 2011 Aug @LP11??

BH Candidate event (1)



BH Candidate event (2)





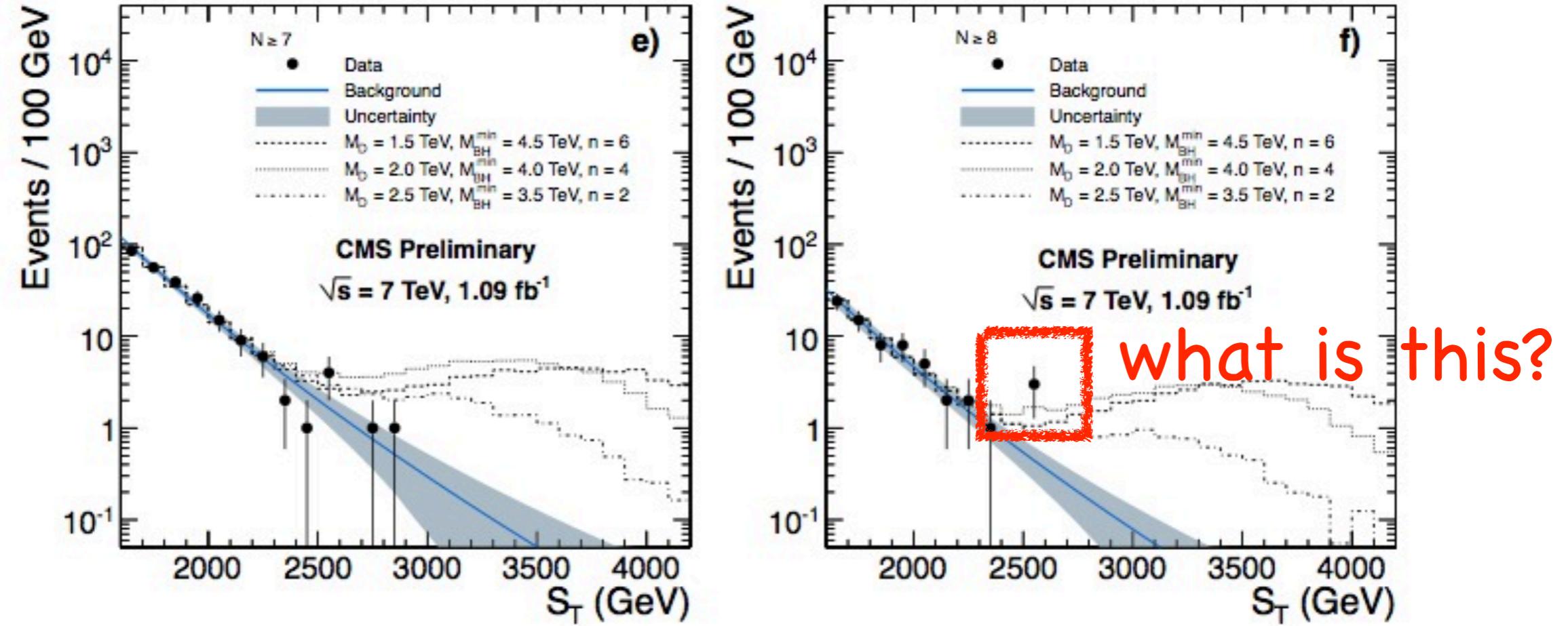


Figure 2: Total transverse energy, S_T , for events with: a) $N \geq 3$, b) $N \geq 4$, c) $N \geq 5$, d) $N \geq 6$, e) $N \geq 7$, and f) $N \geq 8$ photons, electrons, muons, or jets in the final state. Data are depicted as points with error bars; shaded band is the background prediction (solid line) with its uncertainty. Also shown are black hole signals for three different parameter sets.

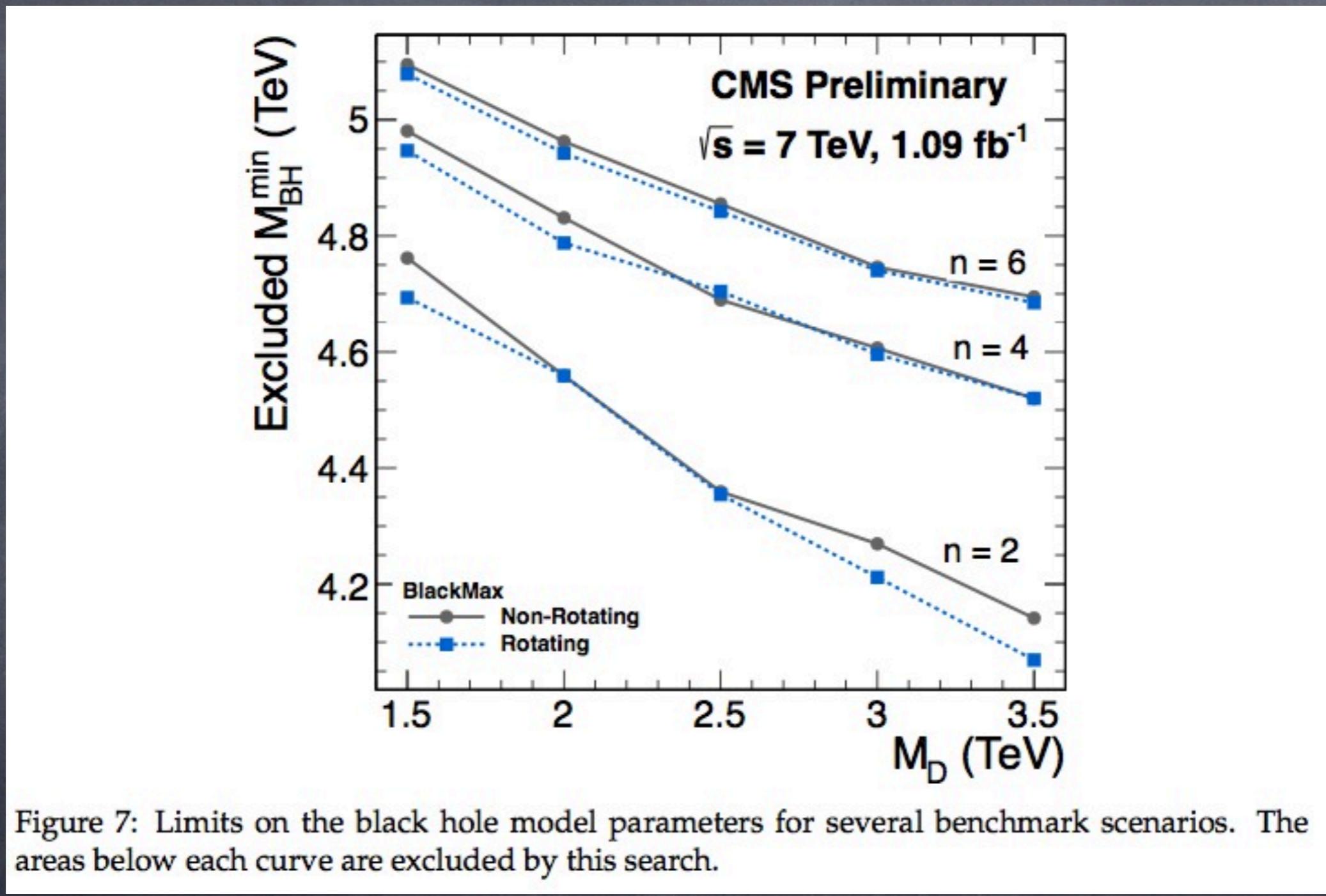


Figure 7: Limits on the black hole model parameters for several benchmark scenarios. The areas below each curve are excluded by this search.

CMS summary

Lower bound of BH mass

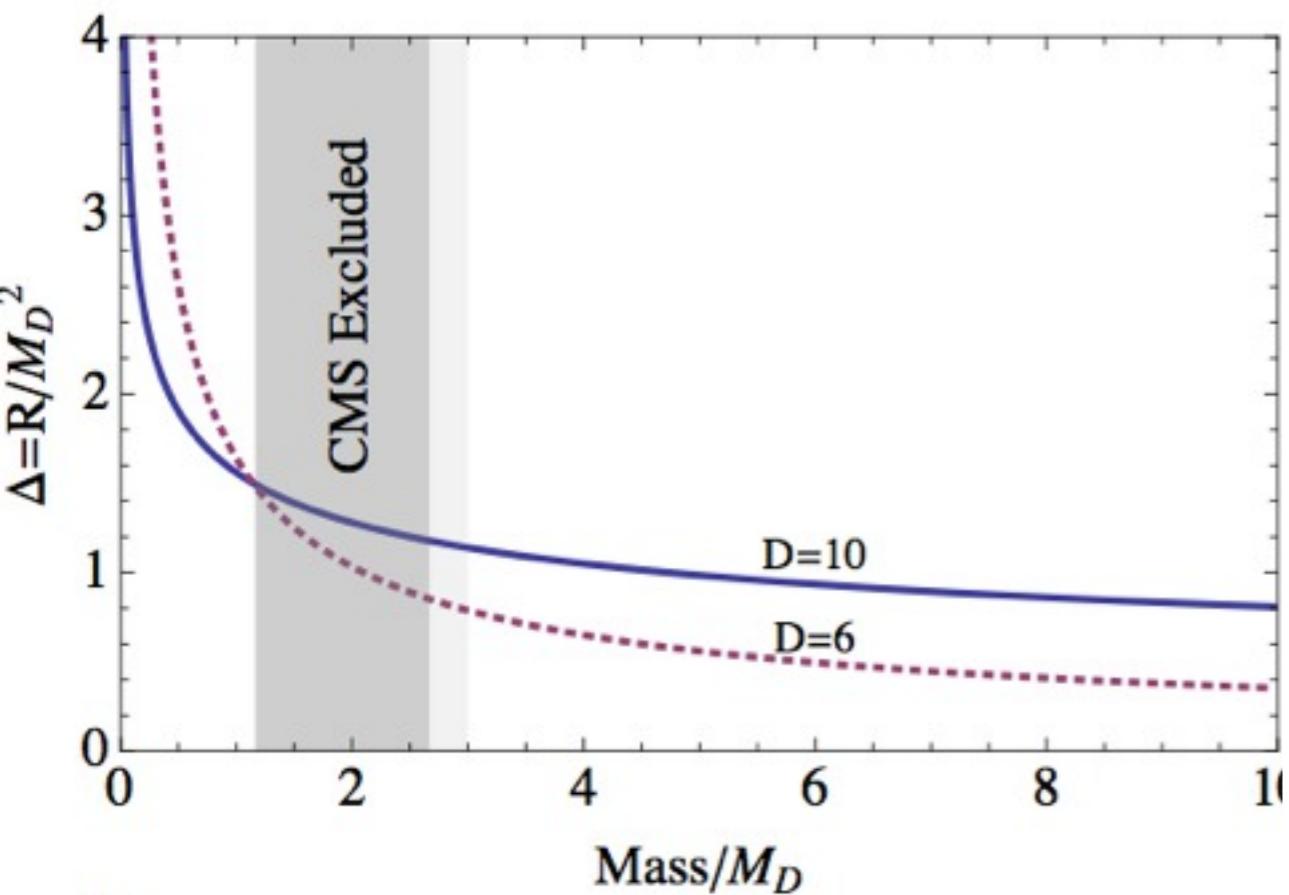
$$\left(\frac{M}{M_D}\right)_{\text{CMS}} \in \left[\frac{3.5}{3.0}, \frac{4.0}{1.5}\right]_{D=6},$$
$$\in \left[\frac{4.0}{3.5}, \frac{4.5}{1.5}\right]_{D=10},$$

or

$$(r_{BH})_{\text{CMS}} \in [0.9, 2.1] \ell_6,$$
$$\in [0.9, 2.4] \ell_{10}.$$

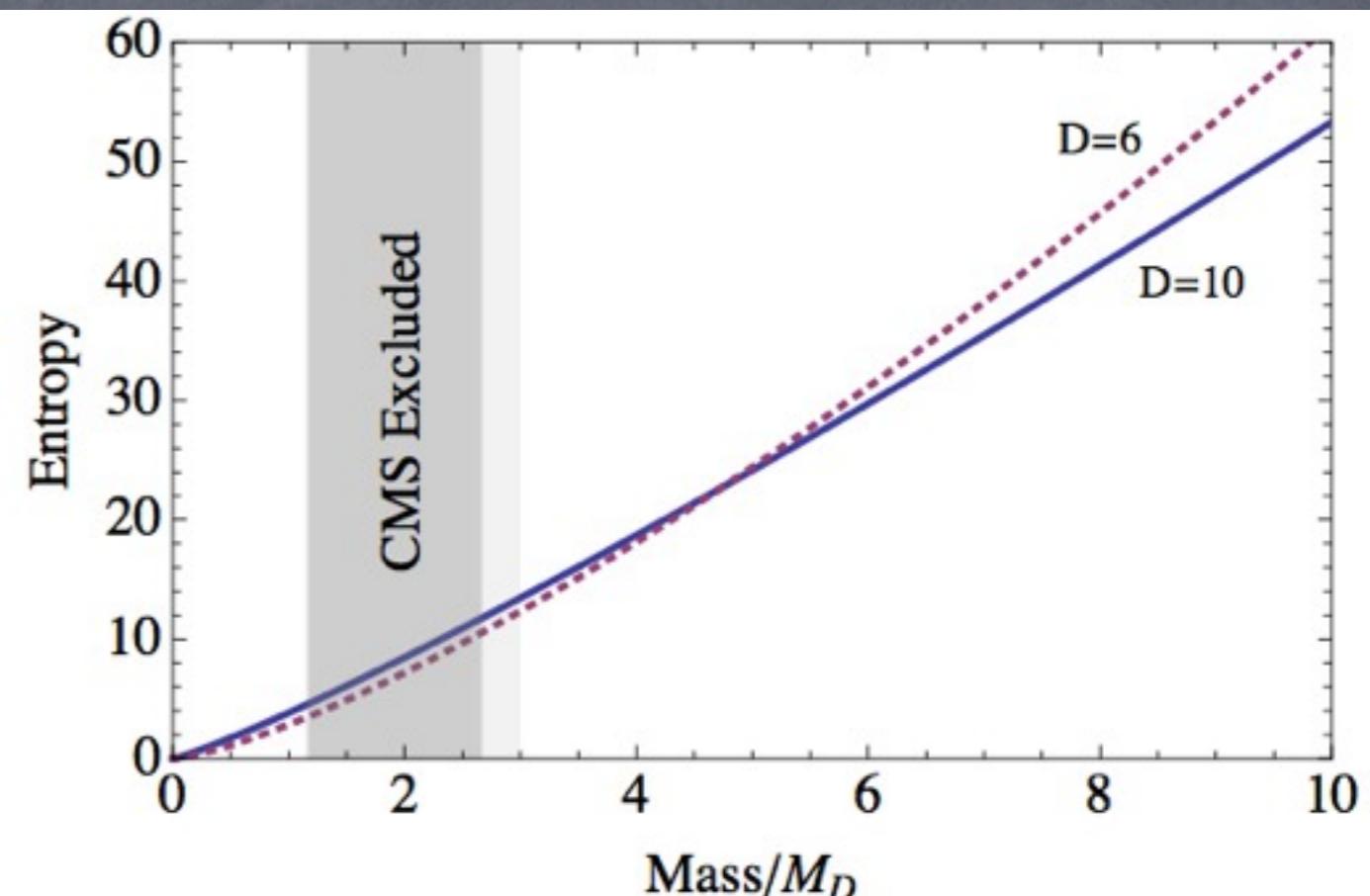
Do not belong to semi-classical domain! :-(

QG correction



$\Delta > 100\%$

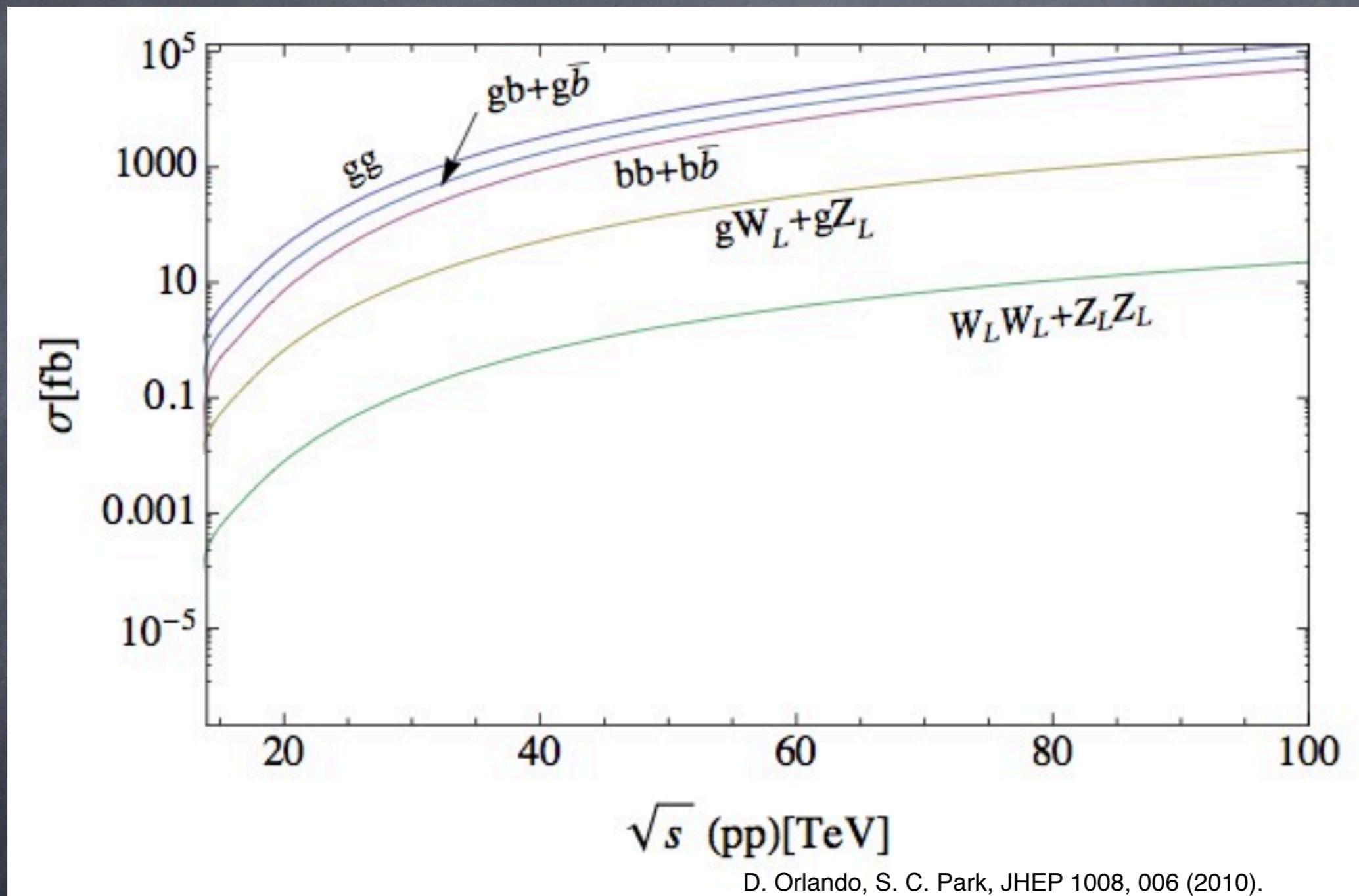
Entropy



Entropy < 10

Model dependence:

RS with bulk SM



conclusion & future prospects (1/3)

- TeV gravity realized with large volume compactification (ADD, RS, CHS ..): hierarchy problem solved.
- Two ways to see LED: Low mass gap ($\sim 1/\text{vol}$) or strong gravity
- The LHC can be the first BH factory (still valid!)

conclusion & future prospects (2/3)

- caution! The first result(7 Tev run) is out of the validity range of semi-classical approximation.
- Model independent “bound” can be useful in the future.
- Future runs with 14 Tev or higher (100 Tev) are more interesting for BH search.

conclusion & future prospects (3/3)

- Scattering at $E \sim M_P$ with a definite QG theory.
- AdS/CFT ... BH \sim Thermal state in CFT.
Glueball? string ball?
- The final state of BH at $M \sim M_P$.