## GRAVITATIONAL LENSING IN A DARK MATTER FREE BRANEWORLD MODEL

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#### All audience of this talk

🗆 Dr. Lam Tsz Yan;

- Prof. Shinji Mukohyama; and
- Prof. Masamune Oguri

## The University of Hong Kong

#### Main building and physics building



# Some Views from Physics building

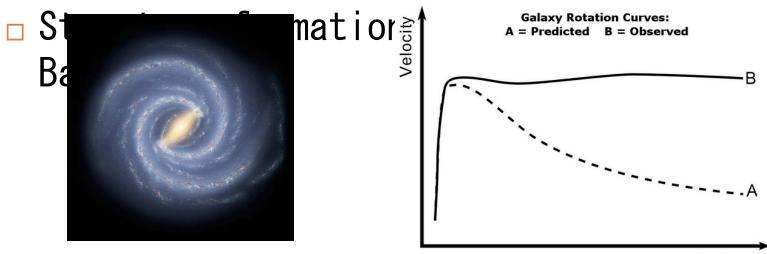


## Outline of presentation

- Background of this project
- Introduction to Randall Sundrum braneworld
- Motivations of this study
- Rotation Curves in braneworlds
- Gravitational Lensing in braneworlds

## Dark Matter problem

- Missing mass at galactic scales and galactic cluster observations
- e.g. Rotation Curves of galaxies, X-ray clusters images, gravitational lensing; and



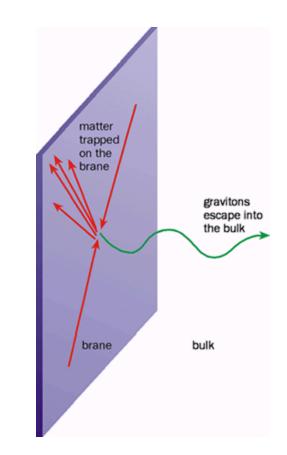
Distance

## Braneworld as a solution

- 1. Evidences of Dark matter always inferred from gravitational motion
- 2. No simple theoretical framework for dark matter
- Supersymmetry offers a solution
- Detection of SUSY particle has never succeeded (yet?)
- Behavior of SUSY still not clear, alternative approach can not be ignored.

## Braneworld

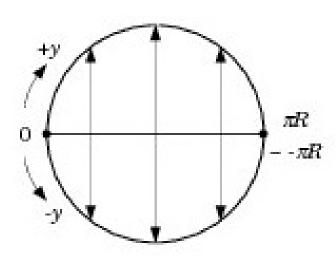
- Offers an explanation on why gravity is weak
- The idea is not new e.g. Kaluza-Klein theory in 1921
- RS models (Randall, Sundrum, 1999)
- Motivated by String theory/M theory



## Randall-Sundrum models

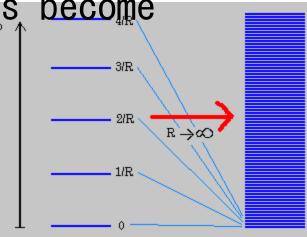
#### □ 2-brane model

- Suggested by Horava-Witten solution
- 1-brane model one of the brane located at infinity
- Showing the possibility of infinite extra dimension



## Major prediction of Braneworld

- 🗆 Kaluza Klein (KK) Modes
- In 2-branes model
- It is like standing waves between two branes
- Predictions in TeV physics
- In 1-branes model KK modes become spectrum
- No predictions in TeV
- But also allows more study about it



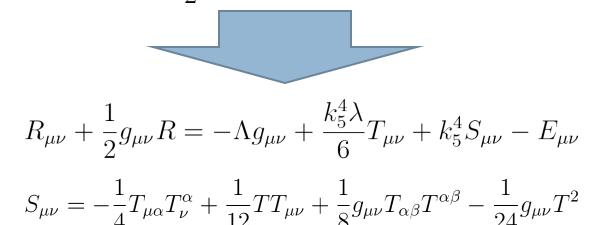
## Original RS 1-brane construction

- □ Gravity is a 5D interaction
- 5D spacetime (bulk) have large negative curvature, it is done by introducing a negative cosmological constant
- Standard Model fields confined on 4D brane, and there is a brane tension fine tuned with 5D cosmological constant
- Brane is mirror symmetric
- Construct action of 2-brane model and limit another brane to infinity

## Braneworld as modified 4D Einstein theory

#### The matter on membrane infer discontinuity of extrinsic curvature (Israel, 1966)

□ Project 5D Einstein equation to 4D (Shiromizu, Maeda<sub>R</sub> = Sasak in 5D 2000)



## What is $E_{\mu\nu}$

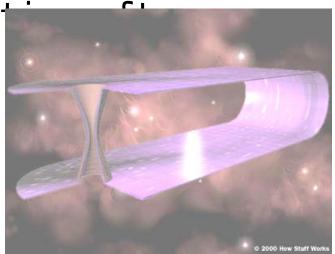
- A traceless tensor
- □ Projection of 5D Weyl tensor to the brane  $E_{\mu\nu} = {}^{(5)} C_{ABCD} n^C n^D g^A_\mu g^B_\nu$
- Weyl tensor = component of Curvature that is not governed by Einstein's equations
- It contains freedom for 5D Gravitational wave, therefore KK spectrum
- KK spectrum could "source" by the brane
   via
    $abla^{\mu}S_{\mu\nu} = k_5^4 \nabla^{\mu}E_{\mu\nu}$
- Is it possible to be dark matter?

## Motivations from wormhole

- Origins of Dark Matter
- □ Galactic scale evidence of extra dimension
- Inter Galactic wormhole may be possible
- Wormholes violate energy conditions in GRDauk Matter"braneworddseeobo, 2007) everywhere is the key to sustain it.

# A little advertisement on wormhole

- □ Wong, Harko, Cheng, 2011
- Braneworld wormhole could evolve with the Universe, and expand with the Universe
- Wormholes that exist before inflation collapse
- Expanding wormholes existing inflation could be still size



## Testify or Falsify

- Explaining dark matter phenomenology
- The braneworld corrections depend on geometry
- More empirical results are required to determine the true structure of spacetime
- Predict different observations from the same object
- Galactic rotation curves (Gergely et al, 2011)

## HSB and LSB

- □ Low Surface Brightness galaxies - low visible mass content, diffuse gas riph ≥ 23mag/arsec<sup>2</sup>
- LSB galaxies are conventionally dark matter rich
   HSB very different from LSB

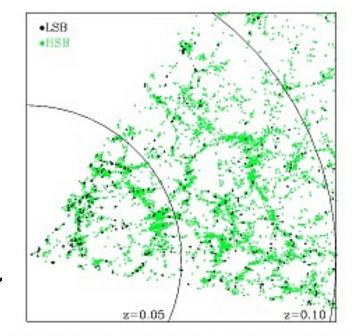


Fig. 1. One of the two analysed pie slices. The right ascension range  $354^\circ \le \alpha \le 53^\circ$  containing ~400 LSB galaxies and ~8000 HSB galaxies is plotted.



## Bulk and brane

The geometry of bulk leaves an imprint on the brane.

□ The brane spherical geometry  $ds^2 = g_{tt}(r)dt^2 + g_{rr}(r)dr^2 + r^2d\Omega^2$ 

□ Possible bulk configuration  $^{(5)}ds^2 = -M(r,y)^2 dt^2 + N(r,y)^2 dr^2 + Q(r,y)^2 d\Omega^2 + dy^2$ 

Calculate 5D Weyl tensor

- Project it to the brane
- 🗆 Simplify with 5D Einstein equati



## The form of $E_{\mu\nu}$

$$\begin{split} E_t^t &= \frac{N_{,y,y}}{N} + \frac{2Q_{,y,y}}{Q} - \frac{\Lambda_5}{2} \\ E_r^r &= -\frac{N_{,y,y}}{N} + \frac{\Lambda_5}{6} \\ E_\theta^\theta &= E_\phi^\phi = -\frac{Q_{,y,y}}{Q} + \frac{\Lambda_5}{6} \end{split}$$

- 2<sup>nd</sup> derivative of metric along extra dimension
- Additional constraint from the conservation equation

Only the r component is non-trivial

## Energy and Pressure of $E_{\mu\nu}$

- Ist approach: Hypothesize high energy events that could source E in the formation history. Based on the physics we can model the remnant form of E, and test it with observations
- Alternative approach in Gergely, et al., 2011
- Components of E view as energy U and pressure P of some fluid
- □ Guess equation of state based on Schwarzschild case in 2+1+1 decomposition

## LSB metric and Visible mass matching

$$-g_{tt} \approx 1 - \frac{2GM}{c^2 r} + \frac{2\gamma}{1 - \alpha} \left(\frac{r}{r_c}\right)^{\alpha - 1}$$
$$g_{rr} \approx \left\{1 - \frac{2G(M + M_b)}{c^2 r} + \gamma \left[\left(\frac{r}{r_c}\right)^{\alpha - 1} - 1\right]\right\}^{-1}$$

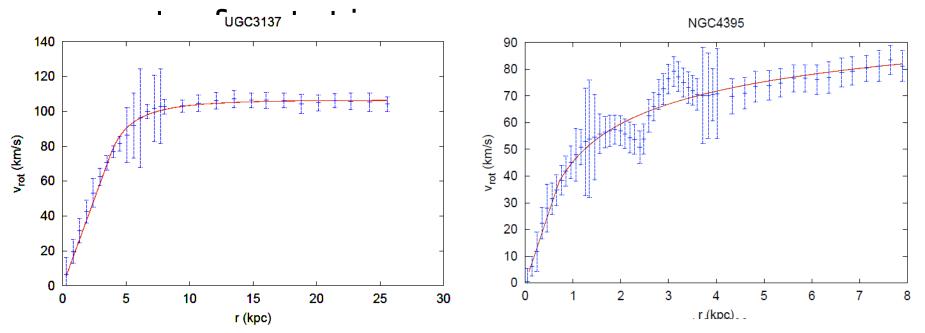
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- $\square \alpha$  and  $\vee$  depend on a and B
- Mb is a degenerate parameter in rotation curve
- rc is obtained from matching Baryonic
  mass

## Rotation Curve studies

Data from Blok, 2001

Gergely, et al., 2011 – approximate Newtonian in linear part of rotation curve and braneworld dominant in outer

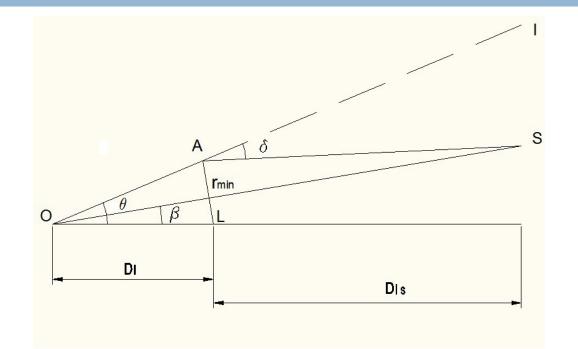


## List of Candidate galaxies

#### k is introduced in step function to switch from baryonic mass dominant region

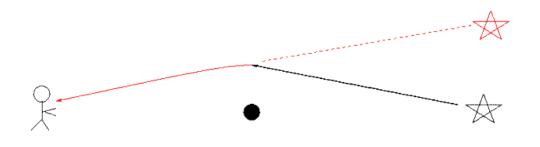
Galaxy	k	$M_0$	$r_c$	$\alpha$	$\gamma$	$\chi^2_{\rm min}$
	${\rm kpc}^{-1}$	$\odot$	kpc			
DDO 189	57.5	$4.05 \times 10^{8}$	1.25	0.3	$6.43 \times 10^{-8}$	0.742
NGC 2366	46.0	$1.05 \times 10^{9}$	1.47	0.8	$1.12 \times 10^{-7}$	2.538
NGC 3274	138.1	$4.38 \times 10^{8}$	0.69	-0.4	$6.73 \times 10^{-8}$	18.099
NGC $4395$	30.0	$2.37 \times 10^{8}$	0.71	0.9	$3.43 \times 10^{-7}$	27.98
NGC $4455$	99.7	$2.26 \times 10^{8}$	1.03	0.9	$2.72 \times 10^{-7}$	7.129
NGC $5023$	86.3	$2.69 \times 10^{8}$	0.74	0.9	$4.53 \times 10^{-7}$	10.614
$UGC \ 10310$	36.4	$1.28 \times 10^{9}$	2.6	0.4	$1.12 \times 10^{-7}$	0.729
UGC 1230	15.3	$3.87 \times 10^{9}$	3.22	-1.7	$1.12 \times 10^{-7}$	0.539
UGC $3137$	34.5	$5.32 \times 10^{9}$	3.87	-0.5	$1.23 \times 10^{-7}$	4.877

## Simple 1D lensing



Lens equation  $\tan |\theta| - \tan(s\beta) - \frac{D_{ls}}{D_s} [\tan |\theta| + \tan(\delta - |\theta|)] = 0 \quad (1)$ 

## Minimal approach radius



### □ Null geodesic

$$g_{\mu\nu}\frac{dx^{\mu}}{d\lambda}\frac{dx^{\nu}}{d\lambda} = g_{tt}(r)t'^{2} + g_{r}r(r)r'^{2} + r^{2}\phi'^{2} = 0$$

- □ Minimum radius on trajectoŕy= 0
- Identify the constants of motion with lensing geometry

$$g_{tt}(r_{\min})D_l^2\sin(\theta) + r_{\min}^2 = 0$$
 (2)

## Deflection angle expanded

Compare angle between asymptotes

$$\delta(r_{\min}) = 2 \int_{r_{\min}}^{\infty} \mathcal{I} - \pi$$
$$\mathcal{I}(r) = \frac{1}{r} \left\{ \frac{g_{rr}(r)}{[g_{tt}(r_{\min})/g_{tt}(r)] (r/r_{\min})^2 - 1} \right\}^{1/2}$$
(3)

Ъ∫

θ

Deflection angle can be studied by treating minimal approach radius like impact parameter

### Image angle and remarks

Lens equation approximated for small angle  $\theta^+ = \beta + \frac{D_{ls}}{D_s} \delta(r_{\min}^+, M_b)$  $\theta^- = \beta - \frac{D_{ls}}{D_s} \delta(r_{\min}^-, M_b)$ 

• With equation 
$$\left(\frac{2}{D_l}\right)^2 \frac{1}{g_{tt}(r_{\min}^{\pm})}$$

A single lensing observation from a known source should able to rule out the model

## Post Newtonian Expansion

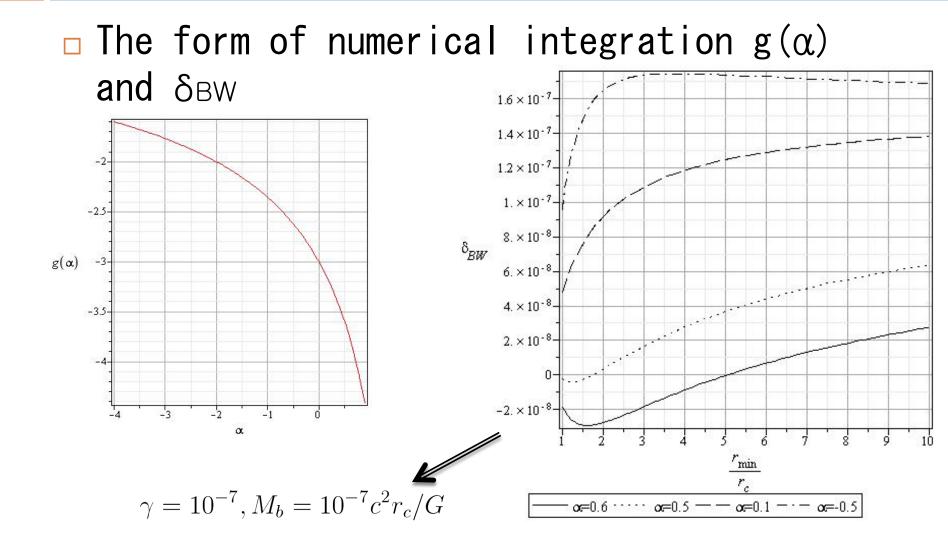
Focus on the deflection angle
 Expand the integral as series of small parameter δ = δ<sub>GR</sub> + δ<sub>BW</sub>

$$\delta_{\rm GR} = \frac{4GM}{c^2 r_{\rm min}}$$

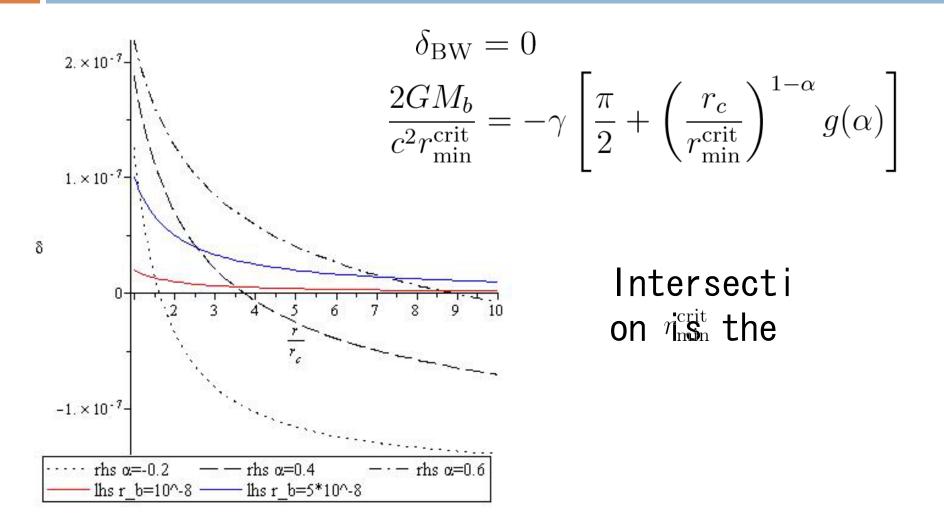
$$\delta_{\rm BW} = \frac{2GM_b}{c^2 r_{\rm min}} + \gamma \left[\frac{\pi}{2} + \left(\frac{r_c}{r_{\rm min}}\right)^{1-\alpha} g(\alpha)\right]$$

$$g(\alpha) = \int_{1}^{\infty} \frac{du}{u} \sqrt{\frac{1}{u^2 - 1}} \left[ \frac{(1 - \alpha)u^{\alpha - 1} + (1 + \alpha)u^{\alpha + 1} - 2u^2}{(1 - \alpha)(u^2 - 1)} \right]$$

## Braneworld contribution



### Critical approach radius



## Lensing profile comparison

- Use the parameters fixed by rotation curves
- Compare LSB braneworld lensing with dark matter lensing
- To see if there could be discriminative effect on the lensing profile
- View dark matter as a correction to baryonic mass lensing

## The lensing of Dark Matter

Pseudo isothermal halo vs NFW model
 Both can explain rotation curve of LSB
 Pseudo isothermal

$$\rho_{\rm iso}(r) = \frac{\rho_0}{1 + \left(\frac{r}{r_s}\right)^2} \longrightarrow M_{\rm iso}(r) = 4\pi r_s^2 \rho_0 \left(r - r_s \arctan(\frac{r}{r_s})\right)$$

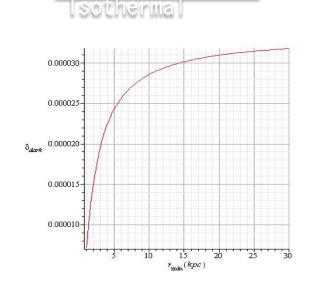
NFW motivated by N body simulations

$$\rho_{\rm NFW}(r) = \frac{\rho_i}{\frac{r}{r_s} \left(1 + \frac{r}{r_s}\right)^2} \qquad \longrightarrow \qquad M_{\rm NFW}(r) = 4\pi r_s^3 \rho_i \left[\frac{\left(1 + \frac{r}{r_s}\right) \ln\left(1 + \frac{r}{r_s}\right) - \frac{r}{r_s}}{1 + \frac{r}{r_s}}\right]$$

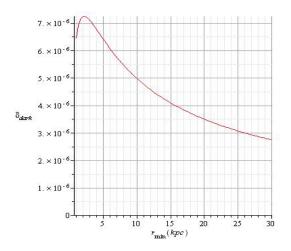
## Isothermal halo lensing

- Assume only mass inside radius of approach cause deflection
- Recall that braneworld deflection angle converge to a constant, it is more

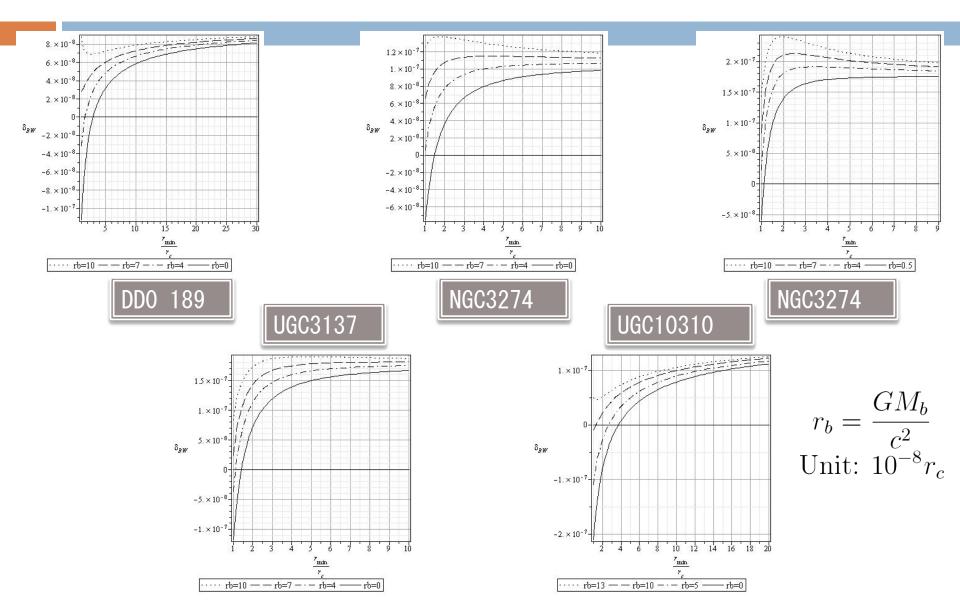
isothermal hal NEW



Simi Pseudo

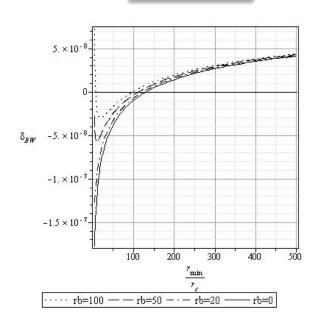


## Isothermal Halo like cases



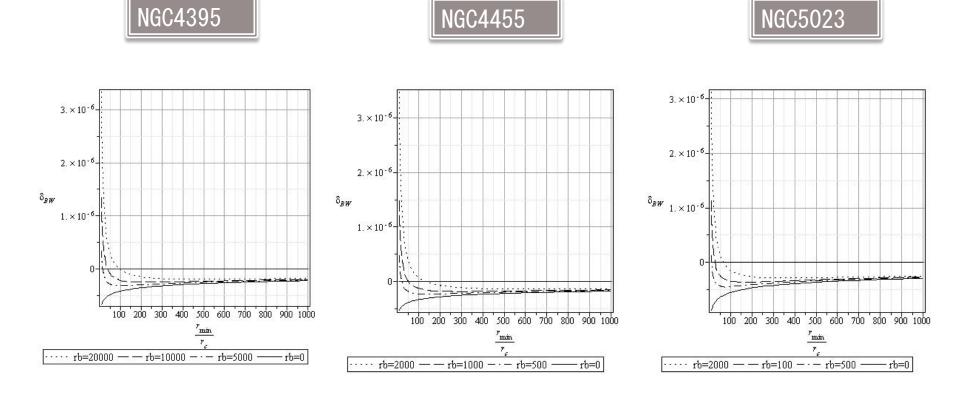
## Unusual Cases

### Negative contribution occurs at large radius

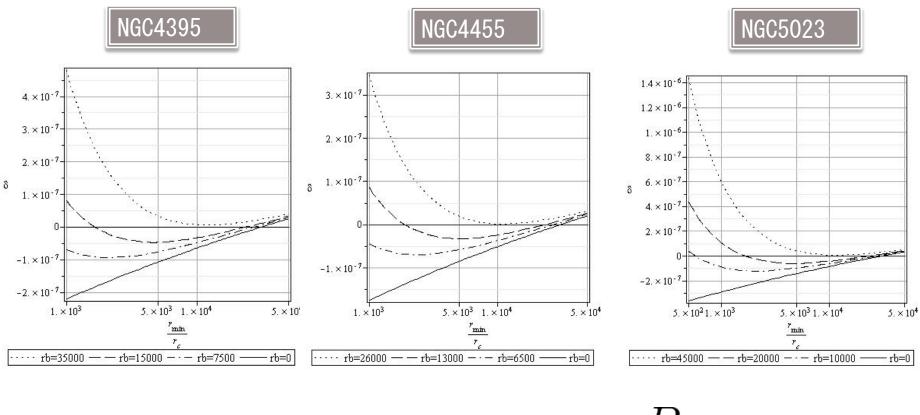


## Unusual Cases

#### Braneworld contribution remain negative at large impact



# Scattering at Imaginative range



 $\Box \alpha \sim 1 \qquad \longrightarrow \qquad P = U + \frac{B}{r^2}$ 

## Another way of comparison

- Assume a well description of rotation curves by braneworld and dark matter model
- The metric will be determined up to a integration constant
- Compare the lensing effect predicted by braneworld versus dark matter

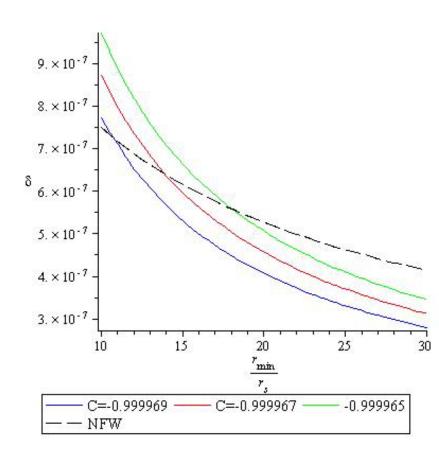
## Formulation

- □ Assume a correction on Schwarzschild metric  $\left[1 - \frac{r_b}{r} + \nu(r)\right] dt^2 + \left[1 - \frac{r_b}{r} + \mu(r)\right]^{-1} dr^2 + r^2 d\Omega$
- □ Assume ∨ fixed by tangential velocity profile in rotat<sup>v</sup>ion cyft
- Dobtain the differential equation of  $\mu$  by traceless  $\frac{d\mu(r)}{dr} = -\frac{\mu(r)\left[1 + M'(r)\right]}{r} \frac{M'(r)}{r}$ ke lowest order, e. g.
- Calculate the braneworld contribution to

## Results

 If the rotation curves is well described by a Newtonian motion in NFW density profile

It is general for
 Braneworld predict
 different deflection
 against dark matter



## Summary

- Braneworld contributions to the brane come from bulk geometry
- Galactic rotation curves as a probe to bulk geometry
- Braneworld lensing indicate unique features as compared to the dark matter case
- Future Studies
  - Determining more precisely the equation of state
  - "Corresponding" mass profile in braneworld, so to obtain 2D lensing image by programming code like glafic (Oguri, 2010)
  - HSB lensing study would allow comparison with more observations
  - What kind of high energy processes in the early