THE ACCELERATING UNIVERSE AND THE COSMIC MYSTERY OF DARK ENERGY



Grigoris Panotopoulos Okinawa Institute of Science and Technology Graduate University

IPMU, September 2012

Outline

- Introduction Motivation
- Dynamical dark energy
- Geometrical dark energy
- Statefinder diagnostics
- Conclusions

Evolution of the Universe



1998: The accelerating universe breakthrough of the year





Theory: Friedmann equations

 Derived from Einstein's theory of General Relativity

First Friedmann equation

Second Friedmann equation

$$H^{2} = \left(\frac{\dot{a}}{a}\right)^{2} = \frac{8\pi G}{3}\rho - \frac{kc^{2}}{a^{2}} + \frac{\Lambda c^{2}}{3}$$

Curvature Cosmological
Constant



APpuquean

A. Friedmann (1888-1925)

$$\begin{aligned} \dot{H} + H^2 = \frac{\ddot{a}}{a} = -\frac{4\pi G}{3}\left(\rho + \frac{3p}{c^2}\right) + \frac{\Lambda c^2}{3} \end{aligned}$$
 Energy density Pressure

Acceleration (of Universe's expansion)

Image: 2nd Friedmann equation gives the acceleration / deceleration with which the universe expands

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}\rho\left(1+3w\right)$$

w < -1/3: acceleration w > -1/3: deceleration

 $w\rho$

Energy conservation (continuity)

$$\dot{\rho} = -3H\left(\rho + \frac{p}{c^2}\right)$$

Useful Quantities

 Critical energy density (the energy density for a flat universe)

$$\rho_c = \frac{3H^2}{8\pi G}$$

Normalized densities

 \mathcal{S}

Matter

$$\Omega = rac{
ho}{
ho_c}$$
 $\Omega_k = -rac{k}{a^2 H^2}$ Curvature

This way the 1st Friedmann equation can be written as

$$\Omega_k + \Omega = 1 \bigg]$$





Basic prediction of inflation: The Universe is flat

Age of the Universe and Hubble constant

Universe older than oldest objects: t₀ > 12 Gy
 Hubble Space Telescope: H₀ = 72 Km/(Mpc sec)

With matter only, $H_0 t_0 = \frac{2}{3}$ \rightarrow Universe is too young

• With a cosmological constant $(\Omega_{\Lambda} = 0.7)$ $H_0 t_0 \cong 0.96, \quad t_0 \cong 13.7 \text{ Gy}$

Primordial Nucleosynthesis



Today's picture of the Universe

- 3 independent data sets coincide
 - Supernovae
 - CMB
 - Galaxy surveys

Concordance Cosmological Model



Today's picture of the Universe

Dark energy dominates in the (flat) Universe



Matter 27% (baryons + cold dark matter)

Dark Energy 73%

-1.1 < w < -0.9



What is Dark Energy?

- In the simplest case:
 Cosmological constant
 Introduced by Einstein for a
 - static universe
 - Pros
 - Allowed by all symmetries
 - ACDM agrees with data

Cons

- The cosmological and
- the coincidence problems



Cosmological Constant

$$G_{\mu\nu} = -\Lambda g_{\mu\nu}$$

• Fluid with w = -1

Very different evolution

 Value much lower than expected



Field equation for gravity

Einstein's General Relativity with dust or radiation only

 $G_{\mu\nu}=8\pi G T_{\mu\nu}$

TheoryDecelerating expansion

Observationaccelerating expansion

Disagreement between theory and observation!

Two choices

Geometrical Dark Energy
 Modify left hand side

Dynamical Dark Energy
 Modify right hand side

$$G_{\mu\nu} + \frac{G_{\mu\nu}^{dark}}{G_{\mu\nu}} = T_{\mu\nu}$$

 $G_{\mu\nu} = T_{\mu\nu} + T_{\mu\nu}^{dark}$

 \rightarrow New gravitational theory

→ New dynamical component

A very active field

- S. Nojiri, S. D. Odintsov and M. Sami, arXiv:hepth/0605039
- V. Sahni and Y. Shtanov, arXiv:astro-ph/0202346
- R. A. Brown, R. Maartens, E. Papantonopoulos and V. Zamarias, arXiv:gr-qc/0508116
- P. S. Apostolopoulos and N. Tetradis, arXiv:hepth/0604014; arXiv:astro-ph/0605450;
- C. Wetterich, L. P. Chimento, R. Lazkoz, R. Maartens and I. Quiros, Nucl. Phys. B 302 (1988) 668;
- B.Ratra and P.J.E.Peebles, Phys. Rev. D 37 (1988) 3406;
- R. R. Caldwell, R. Dave and P. J. Steinhardt, arXiv:astroph/9708069

The Coincidence Problem

Q: Why are the Omegas of matter and dark energy so similar in magnitude?

A: There are two possibilities
1. Due to special initial conditions:
if current universe is a finite point in phase-space

2. Due to the specific values of parameters: if current universe is close to a fixed point

Not so simple to realize!

Cosmology of type
$$H^2 = 2\gamma(\rho + \rho_{DE})$$

Without energy exchange

$$\dot{\rho} + 3H(\rho + p) = 0$$

With energy exchange

Fixed point \rightarrow deceleration

$$\dot{\rho} + 3(1+w)H\rho = -T$$
$$\dot{\rho}_{DE} + 3(1+w_{DE})H\rho_{DE} = T$$

Fixed point \rightarrow acceleration

Superstring theory: basic idea

- Really fundamental objects are one-dimensional (strings)
- In low energies strings look
 like point-like particles
- All known particles are different oscillatory modes of the string



Extended Objects: Branes

- String theory does not contain strings only
- Normally, open strings satisfy
 Neumann boundary conditions
- Dirichlet boundary conditions also make sense



- End points are stuck on a hypersurface, interpreted as a heavy solitonic object, a D-brane
- Brane –world idea: we are confined on such an object.

A simple brane model (DGP)

Action

$$S = \int d^5x \sqrt{-g} M^3 R + \int d^4x \sqrt{-h} \left(m^2 \hat{R} - \mathcal{L}_{SM} \right)$$

- One extra dimension
- Gravity in 5D, our world in 4D
- Reduced to known gravity and cosmology in the early universe
- New gravity and cosmology at late times

Dvali, Gabadadze, Porrati, hep-th/0005016; hep-th/0008054

Cosmology for DGP

Friedmann equation

$$\left(\frac{H}{r_c} = H^2 - \frac{8\pi G}{3}\rho\right) \qquad \left(r_c = \frac{m^2}{M^3}\right)$$

• For early times, we recover standard cosmology $H r_c \to \infty$

- At recent times, $\rho \rightarrow 0, H \rightarrow 1/r_c$
- □ Same number of parameters as ΛCDM , $r_c \cong H_0^{-1}$

$$\Omega_{r_c} = \frac{1}{4r_c^2 H_0^2} = \left(\frac{1-\Omega_M}{2}\right)^2$$

Deffayet, Dvali, Gabadadze, astro-ph/0105068

A more realistic model

$$S = \int d^5x \sqrt{-g} \left(M^3 R - \Lambda \right) + \int d^4x \sqrt{-h} \left(m^2 \hat{R} - V \right)$$

$$G_{AC} = \frac{1}{2M^3} T_{AC}|_{tot}$$

+ Matter in 5 dimensions (unspecified)+ Fluid on the brane

Kofinas, Panotopoulos, Tomaras, hep-th/0510207

Cosmological Solution

$$ds^{2} = -n(t, y)^{2} dt^{2} + a(t, y)^{2} \gamma_{ij} dx^{i} dx^{j} + b(t, y)^{2} dy^{2}$$

$$G_{00} = 3 \left\{ \frac{\dot{a}}{a} \left(\frac{\dot{a}}{a} + \frac{\dot{b}}{b} \right) - \frac{n^{2}}{b^{2}} \left[\frac{a''}{a} + \frac{a'}{a} \left(\frac{a'}{a} - \frac{b'}{b} \right) \right] + \frac{kn^{2}}{a^{2}} \right\}$$

$$G_{05} = 3 \left\{ \frac{n'}{a} \frac{\dot{a}}{a} + \frac{a'}{a} \frac{\dot{b}}{b} - \frac{\dot{a}'}{a} \right\}$$

$$G_{55} = 3 \left\{ \frac{a'}{a} \left(\frac{a'}{a} + \frac{n'}{n} \right) - \frac{b^{2}}{n^{2}} \left[\frac{\ddot{a}}{a} + \frac{\dot{a}}{a} \left(\frac{\dot{a}}{a} - \frac{\dot{n}}{n} \right) \right] - \frac{kb^{2}}{a^{2}} \right\}$$

$$G_{ij} = \frac{a^{2}}{b^{2}} \gamma_{ij} \left\{ \frac{a'}{a} \left(\frac{a'}{a} + \frac{2n'}{n} \right) - \frac{b'}{b} \left(\frac{n'}{n} + \frac{2a'}{a} \right) + \frac{2a''}{a} + \frac{n''}{n} \right\}$$

$$+ \frac{a^{2}}{n^{2}} \gamma_{ij} \left\{ \frac{\dot{a}}{a} \left(\frac{2\dot{n}}{n} - \frac{\dot{a}}{a} \right) - \frac{2\ddot{a}}{a} + \frac{\dot{b}}{b} \left(\frac{\dot{n}}{n} - \frac{2\dot{a}}{a} \right) - \frac{\ddot{b}}{b} \right\} - k\gamma_{ij}$$

Cosmological equations

$$\lambda = \frac{2r}{m^2} + \frac{12}{r_c^2} - \frac{\pi}{M^3} \quad \mu = \frac{r}{6m^2} + \frac{2}{r_c^2} \quad \gamma = \frac{1}{12m^2} \quad \beta = \frac{1}{\sqrt{3}r_c}$$

Phase-space analysis

Л

1

$$\omega_{m} + \omega_{\psi} = 1 \qquad ' = \frac{1}{D} \frac{a}{dt}$$

$$\omega_{m}' = \omega_{m} \Big[(1+3w)(\omega_{m}-1)Z - \frac{A}{\sqrt{|\mu|}} \Big(\frac{|\mu|\omega_{m}}{2\gamma} \Big)^{\nu-1} (1-Z^{2})^{\frac{3}{2}-\nu} -2Z(1-Z^{2}) \frac{1-Z^{2}-3(1-3w)\beta^{2}\mu^{-1}\omega_{m}}{1-\omega_{m}} \Big]$$

$$Z' = (1 - Z^2) \left[(1 - Z^2) \frac{1 - Z^2 - 3(1 - 3w)\beta^2 \mu^{-1} \omega_m}{1 - \omega_m} - 1 - \frac{1 + 3w}{2} \omega_m \right]$$

New quantities for dynamical study

$$D = \sqrt{H^2 - \mu} \qquad Z = \frac{H}{D} \qquad \omega_m = \frac{2\gamma\rho}{D^2} \qquad \omega_\psi = \frac{\beta\psi}{D^2}$$

Critical points and their stability

	$\nu < 3/2$	$\nu = 3/2$	$\nu > 3/2$			
No. of F.P.	1	0 or 1	1			
Nature	А	A	S			
Table 1: The fixed points for $w=0$, influx						

	$\nu < 3/2$	$\nu = 3/2$	$3/2 < \nu < 2$	$\nu = 2$	$\nu > 2$		
No. of F.P.	1	0 or 1	0 or 2	0 or 1	1		
Nature	A	А	A,S	S	S		
Table 2: The fixed points for $w=1/3$, influx							

Kofinas, Panotopoulos, Tomaras, hep-th/0510207

Numerical results for brane model

• Evolution in the $\omega_m - Z$ plane for k = 0, w = 0, A < 0



Coincidence problem with dynamical dark energy

- Cosmic acceleration → scalar field (idea from inflation)
- Exponential potential \rightarrow cosmological scaling solutions ($\frac{\Omega_{DE}}{\Omega_M} = const.$)
- Interaction between DE and DM
 - Possibility that cannot be excluded
 - Can be realized in brane-models with brane-bulk energy exchange and in 4D models with -varying mass
 - We can have $w_{eff} < -1$ (allowed by data but problematic in theory) even though $w_{\phi} > -1$

Interacting dynamical dark energy: Quintessence

 Cosmological constant problem OK (not vacuum energy anymore)

- Coincidence problem \rightarrow Interaction between DE & DM
- Usually assume source $Q \propto \rho_{dm}$ (linear)
- Model with $Q \propto \rho_{dm} \rho_{\phi}$ (0911.3089, quadratic)
- Our idea: Lagrangian description & comparison to data



Dark energy \rightarrow Canonical scalar field ϕ (Quintessence) Dark matter \rightarrow Fermion Ψ • Self-interaction potential $V(\phi)$ Interaction \rightarrow Lagrangian mass term for dark matter $m_d \overline{m(\phi)} \overline{\Psi} \overline{\Psi}$

Mena, Honorez, Panotopoulos, 1009.5263

Equations of motion

For dark matter

For scalar field

Source:

$$\dot{\rho}_{dm} + 3H\rho_{dm} = Q$$
$$\ddot{\phi} + 3H\dot{\phi} + V_{,\phi} = -Q/\dot{\phi}$$
$$Q = \frac{\partial \ln m_{dm}(\phi)}{\partial \phi}\rho_{dm}\dot{\phi}$$

• Requirement $Q \propto \rho_{dm} \rho_{\phi} \approx \rho_{dm} V(\phi)$ satisfied by

$$V(\phi) = M^4 \exp[-\alpha \phi / M_{pl}]$$
$$m_{dm}(\phi) = \exp\left[\left(V(\phi) / \rho_{cr}^0\right)^n\right]$$

Phase-space analysis

Define new dimensionless variables

$$y^2 = \frac{\kappa^2 V}{3H^2}$$
 $x^2 = \frac{\kappa^2 \dot{\varphi}^2}{6H^2}$ $z = \frac{H_0}{H + H_0}$

 $\kappa^2 = 8\pi G$

1st Friedmann eq. (constraint)

$$\Omega_{dm} = 1 - x^2 - y^2$$

2nd Friedmann eq. (dynamical)

$$\frac{\dot{H}}{H^2} = -\frac{3}{2}(1+x^2-y^2)$$

New dynamical equations

$$\begin{aligned} x' &= -3x + \frac{3}{4} \frac{\alpha}{\sqrt{3\pi}} y^2 + \frac{3}{4} \frac{\alpha}{\sqrt{3\pi}} y^{2n} \frac{(1-z)^{2n}}{z^{2n}} (1-x^2-y^2) \\ &+ \frac{3}{2} x (1+x^2-y^2) \\ y' &= -\alpha \frac{\sqrt{3}}{4\sqrt{\pi}} xy + \frac{3}{2} y (1+x^2-y^2) \\ z' &= \frac{3}{2} z (1-z) (1+x^2-y^2) \end{aligned}$$

$$\begin{aligned} N &= \log(a) \end{aligned}$$

Stable fixed point \rightarrow acceleration

$$x_* = \frac{\alpha}{4\sqrt{3\pi}} \qquad y_* = \frac{1}{4}\sqrt{16 - \frac{\alpha^2}{3\pi}} \qquad z_* = 1$$

 Existence: $\Omega_{\phi*} = x_*^2 + y_*^2 \le 1$ Stability: $x = x_* + u$ $y = y_* + v$ $\begin{pmatrix} \dot{u} \\ \dot{v} \end{pmatrix} = \mathcal{M} \begin{pmatrix} u \\ v \end{pmatrix}$ $\operatorname{Re} \lambda_1, \lambda_2 < 0$

• Acceleration:
$$w_{T*} = x_*^2 - y_*^2 < -\frac{1}{3} \rightarrow \alpha < 4\sqrt{\pi}$$

Comparison with data

Supernovae

$$d_L(z) = c(1+z) \int_0^z H(z)^{-1} dz \quad \mu = 5 \log\left(\frac{d_L}{Mpc}\right) + 25$$

CMB

$$R = 1.7 \pm 0.03 \qquad \qquad R = (\Omega_m H_0^2)^{1/2} \int_0^{1089} dz / H(z)$$

BAO

$$A = 0.469 \pm 0.017$$
$$A = \sqrt{\Omega_m H_0^2} \left(\frac{d_L (z = 0.35)^2}{H(z = 0.35)(1 + 0.35)^2 0.35^2)} \right)^{1/3}$$

Global χ^2 analysis

Supernovae

$$\chi^2_{SNIa}(c_i) = \sum_{z,z'} \left(\mu(c_i, z) - \mu_{obs}(z) \right) C_{z;z'}^{-1} \left(\mu(c_i, z') - \mu_{obs}(z') \right)$$

CMB

$$\chi^2_{CMB}(c_i) = [(R(c_i) - R)/\sigma_R]^2$$

BAO

$$\chi^2_{BAO}(c_i) = \left[(A(c_i, z = 0.35) - A) / \sigma_{A(z=0.35)} \right]^2$$

$$\chi^{2}_{tot}(c_{i}) = \chi^{2}_{SNIa}(c_{i}) + \chi^{2}_{BAO}(c_{i}) + \chi^{2}_{CMB}(c_{i})$$

Numerical Results for 4D model



Diagnostic for different cosmological models

- Many models with similar expansion history that cannot be excluded by data
- Quantities that can be measured and computed within a model

Appropriate quantities

$$r = \frac{\ddot{a}}{aH^3} \qquad s = \frac{r-1}{3(q-\frac{1}{2})}$$
$$H = \frac{\dot{a}}{a} \qquad q = -\frac{\ddot{a}}{aH^2}$$

Sahni, Saini, Starobinsky, Alam, astro-ph/0201498; Alam, Sahni, Saini, Starobinsky, astro-ph/0303009

4D model with interacting dark energy

$$H^2 = \frac{\kappa^2}{3} \rho$$

$$\dot{H} = -\frac{\kappa^2}{2} \left(\rho + p\right)$$

$$Q = \dot{\rho}_m + 3H\rho_m$$

$$-Q = \dot{\rho}_X + 3H(\rho_X + p_X)$$

$$\rho = \rho_m + \rho_X \qquad p = p_m + p_X = p_X = w\rho_X$$
$$\kappa^2 = 8\pi G \qquad Q = \delta H \rho_m$$

Use of dimensionless quantities

$$\Omega_m = \frac{\kappa^2 \rho_m}{3H^2} \quad \Omega_X = \frac{\kappa^2 \rho_X}{3H^2} \quad \Omega_m + \Omega_X = 1$$

Upon comparison to observational data

 $\delta = -0.03 \qquad w = -1.02 \qquad \Omega_{X,0} = 0.73$

astro-ph/0702015

Statefinders

$$s = 1 + w + \frac{\delta}{3}(\frac{1}{\Omega_X} - 1) \quad r = 1 + \frac{9}{2} w \Omega_X s$$

$$\Omega'_X = -(1 - \Omega_X)(\delta + 3w \Omega_X) \quad N = \ln a \quad \Omega'_m = -\Omega'_X$$

• Critical points: two fixed points
Stable $\Omega_{*,1} = 1$ $\Omega_{*,2} = -\frac{\delta}{3w}$ Unstable
• At the stable critical point
 $s = 1 + w \qquad r = 1 + \frac{9}{2}w(1 + w)$
• Special case: cosmological constant
 $w = -1 \qquad \delta = 0 \qquad s = 0 \qquad r = 1$

Numerical results for the 4D model

Parameter s vs. red-shift
Parameter r vs. red-shift



G. Panotopoulos, 0712.1177

Numerical results for the brane model

Parameter s vs. red-shift
Parameter r vs. red-shift



G. Panotopoulos, 0712.1177

(s-r) plane for both models

■ 5D brane model

■ 4D model in GR



G. Panotopoulos, 0712.1177

Conclusions

- Cosmic acceleration \rightarrow Dark energy
- Cosmological constant & ACDM: most economical model
- Other possibilities
 - Dynamical dark energy (quintessence)
 - Geometrical dark energy (brane models)

Statefinders: can discriminate between different dark energy models with same expansion history