

# UNDERSTANDING THE SHAPE OF NON-GAUSSIANITY



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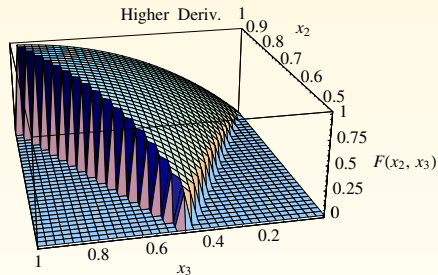
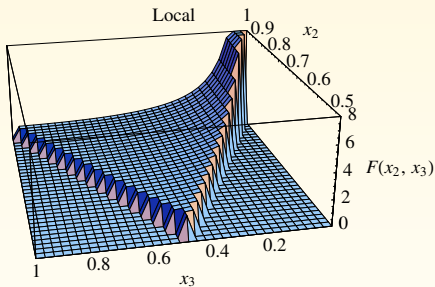
*Institute for the Physics and Mathematics of the Universe*

*University of Tokyo, Kashiwa, Japan*

*19 November 2012*

# THE MANY SHAPES OF NON-GAUSSIANITY

Babich, Creminelli and Zaldarriaga [astro-ph/0405356]

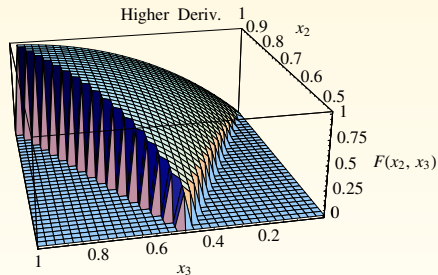
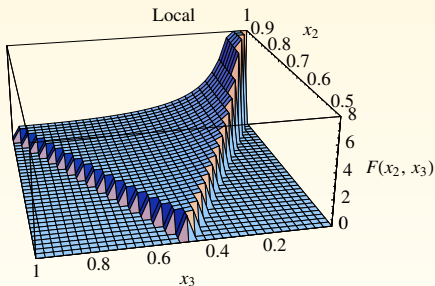


But wait, how exactly  $k$ 's got in there? Wave vector can enter through:

- gradient terms:  $\vec{\partial}\phi \rightarrow \vec{k}\phi_k$
- time derivatives:  $\dot{\phi} \rightarrow |k|\phi_k$
- spatial correlations:  $P(\vec{k}) = 1/|k|^3$

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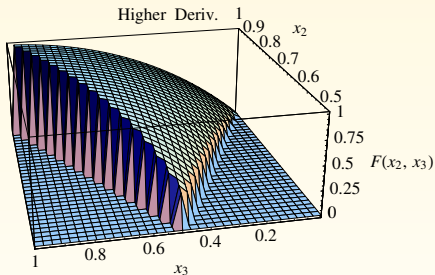
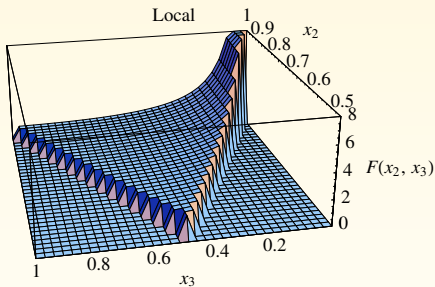


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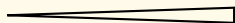
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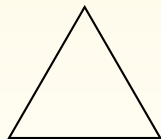
# THE MANY SHAPES OF NON-GAUSSIANITY (CONT.)

$$\langle \Phi(\vec{k}_1)\Phi(\vec{k}_2)\Phi(\vec{k}_3) \rangle = f_{\text{NL}} B(\vec{k}_1, \vec{k}_2, \vec{k}_3) \delta^3(\vec{k}_1 + \vec{k}_2 + \vec{k}_3) + \text{connected}$$

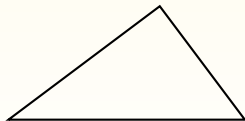
$$B_{\text{loc}}(\vec{k}_i) = \frac{1}{k_1^3 k_2^3} + \frac{1}{k_1^3 k_3^3} + \frac{1}{k_2^3 k_3^3}$$



$$B_{\dot{\sigma}^3}(\vec{k}_i) = \frac{1}{k_1 k_2 k_3 k_t^3}, \quad k_t = k_1 + k_2 + k_3$$



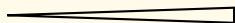
$$B_{\dot{\sigma}(\nabla\sigma)^2}(\vec{k}_i) = \frac{\vec{k}_1 \cdot \vec{k}_2}{k_1^3 k_2^3 k_3} \left( \frac{1}{k_t} + \frac{k_1 + k_2}{k_t^2} + 2 \frac{k_1 k_2}{k_t^3} \right) + \text{permutations}$$



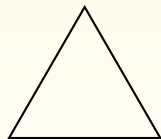
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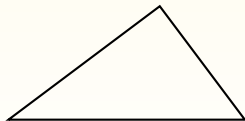
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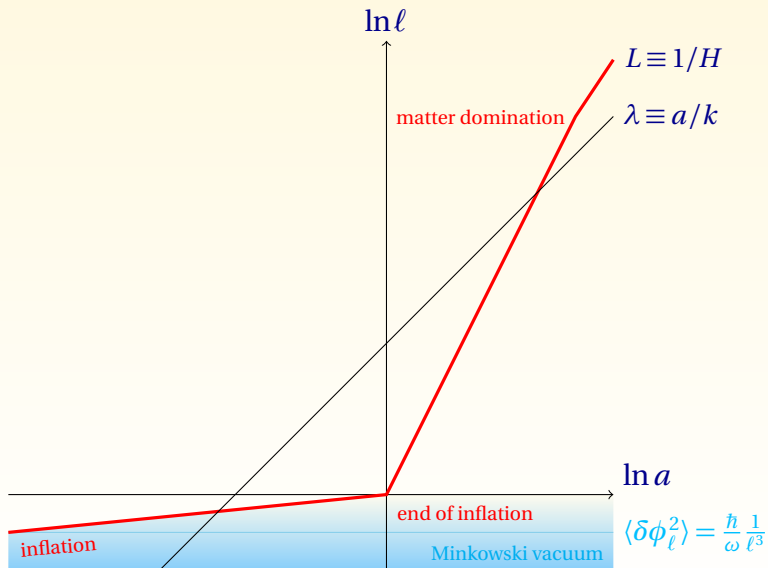
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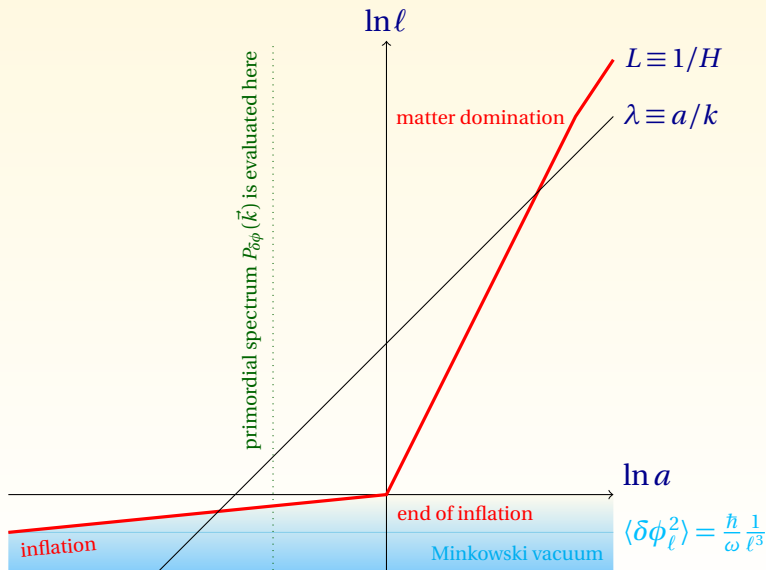
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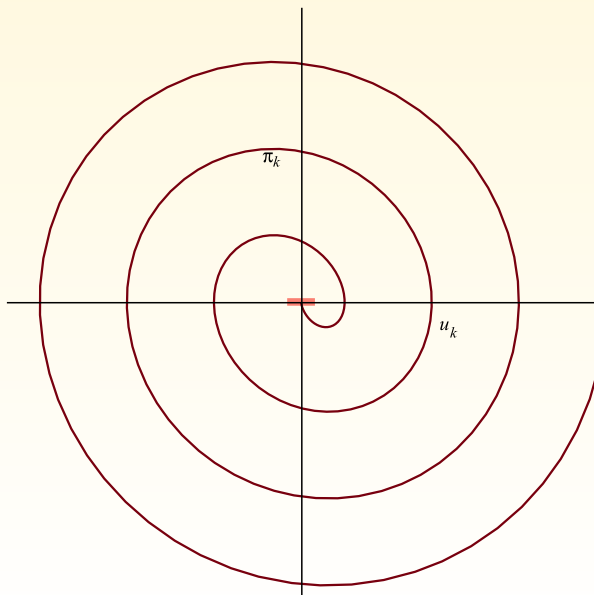
# QUANTUM WHAT & WHERE, OR MAP TO INFLATION

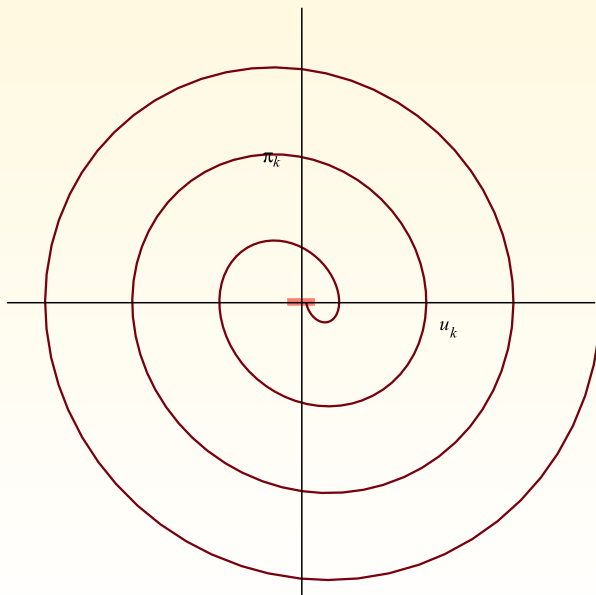


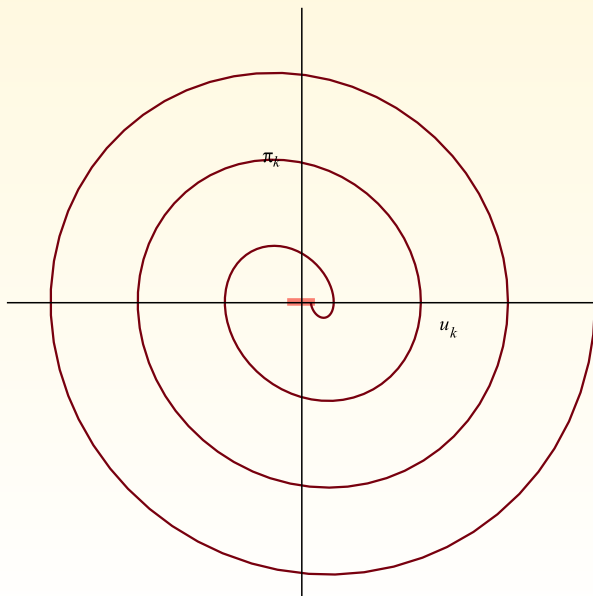
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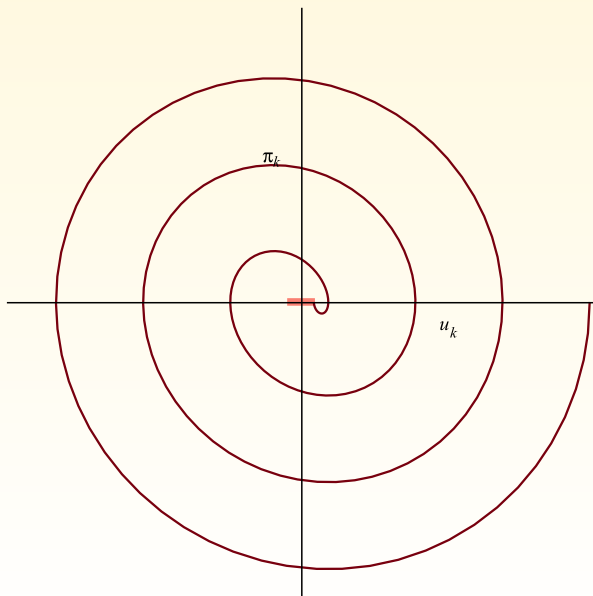


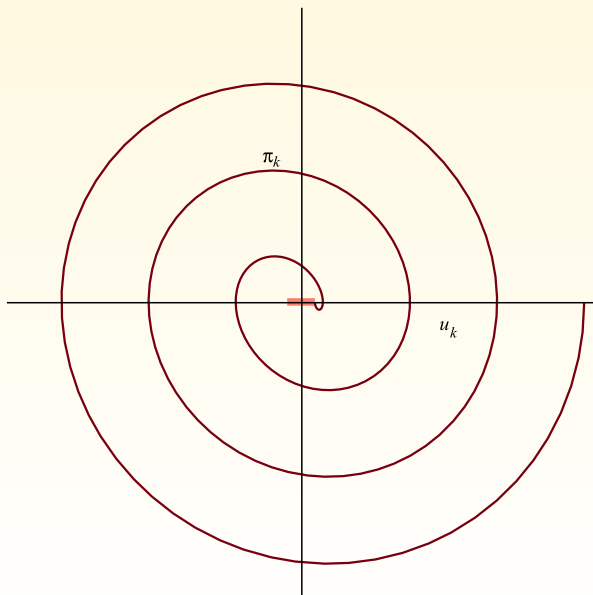


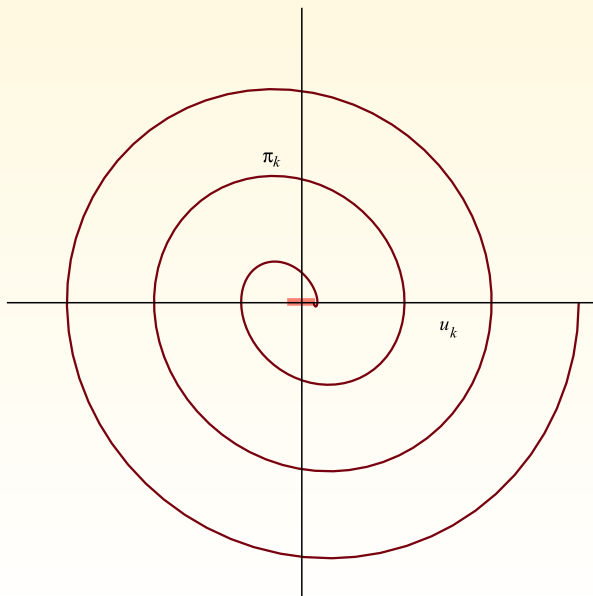


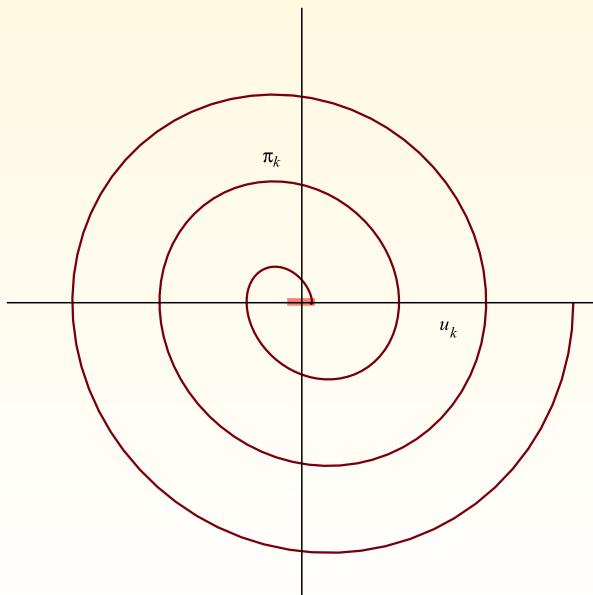


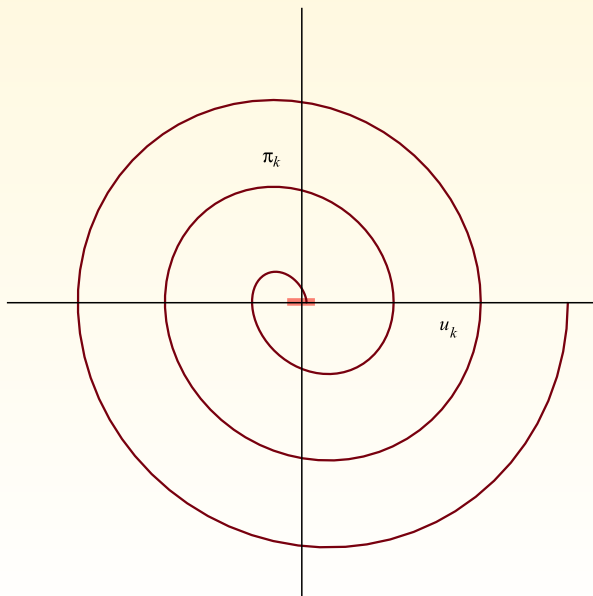




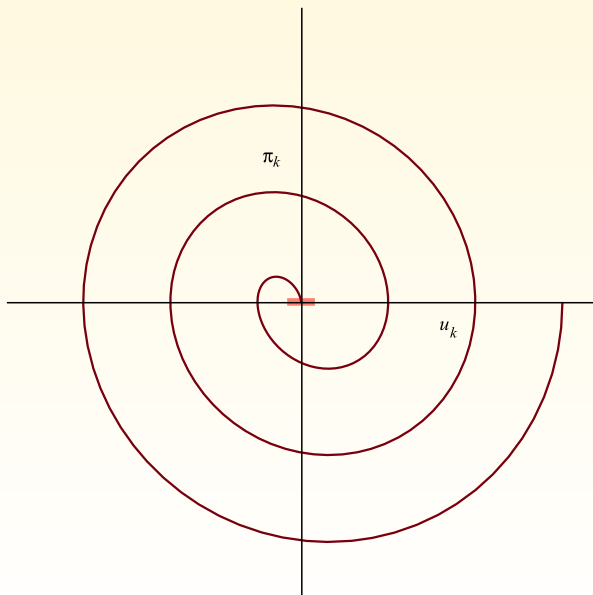












Wave equation for mode  $\varphi_k$  with mass  $m$ , wavevector  $k$  in FRW metric:

$$\ddot{\varphi}_k + 3H\dot{\varphi}_k + \left(m^2 + \frac{k^2}{a^2}\right)\varphi_k = 0$$

Use conformal time  $H\eta = -1/a$ , switch variables to  $x = -k\eta = p/H$ :

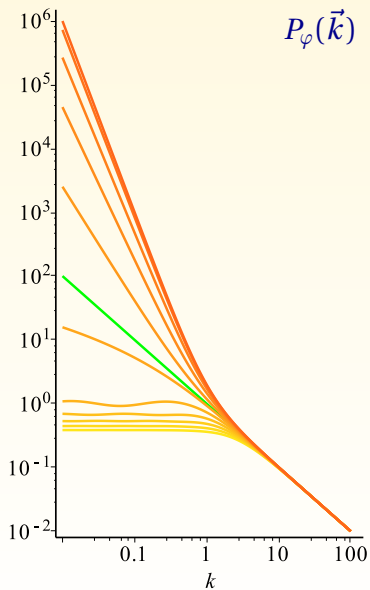
$$x^2\varphi_k'' - 2x\varphi_k' + \left(\frac{m^2}{H^2} + x^2\right)\varphi_k = 0$$

$$\varphi_k(x) = c_1 x^{\frac{3}{2}} H_\mu^{(1)}(x) + c_2 x^{\frac{3}{2}} H_\mu^{(2)}(x), \quad \mu^2 = \frac{9}{4} - \frac{m^2}{H^2}$$

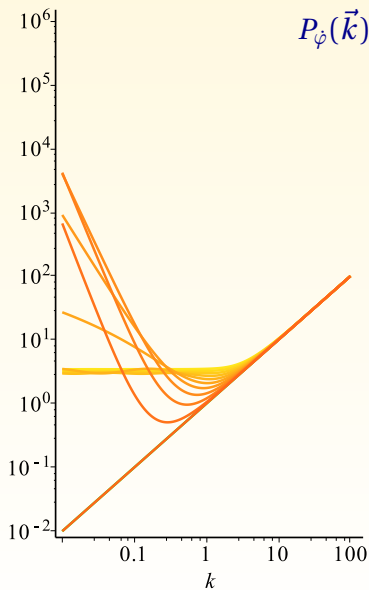
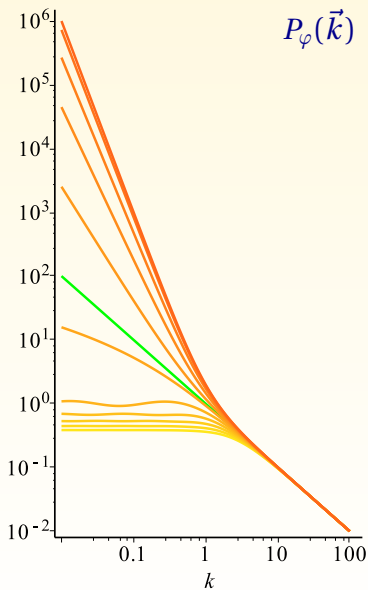
Match to Minkowski vacuum  $\frac{e^{-ipt}}{\sqrt{2p}}$  as  $p \rightarrow \infty$  to fix values of  $c_1$  and  $c_2$ :

$$\varphi_k(\eta) = \frac{\pi^{\frac{1}{2}}}{2} H(-\eta)^{\frac{3}{2}} c_\mu H_\mu^{(1)}(-k\eta), \quad c_\mu = e^{i\frac{\pi}{2}(\mu+\frac{1}{2})} \text{ is pure phase}$$

# DE SITTER VACUUM SPECTRA



# DE SITTER VACUUM SPECTRA, WITH A SURPRISE!



# NON-LINEARITY SOURCES NON-GAUSSIANITY

- Vacuum state deformation
- Mode freeze-out on (sound) horizon exit
- Transfer to curvature fluctuations

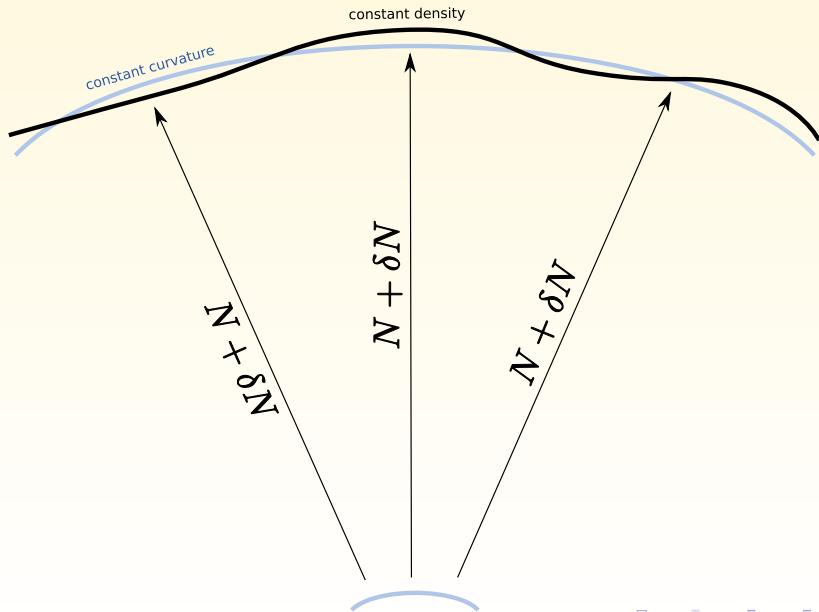
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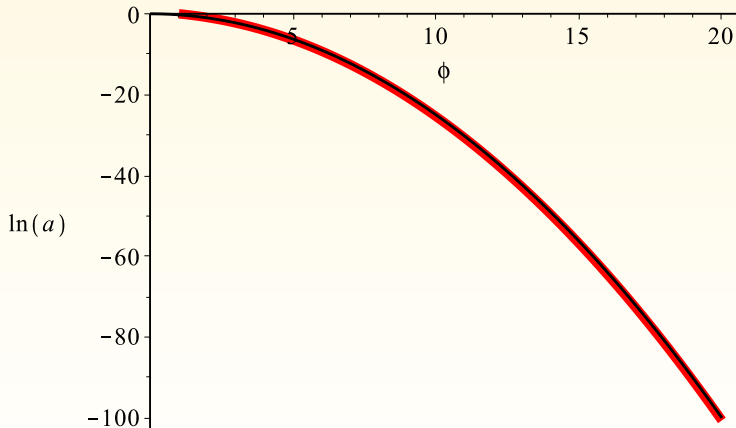
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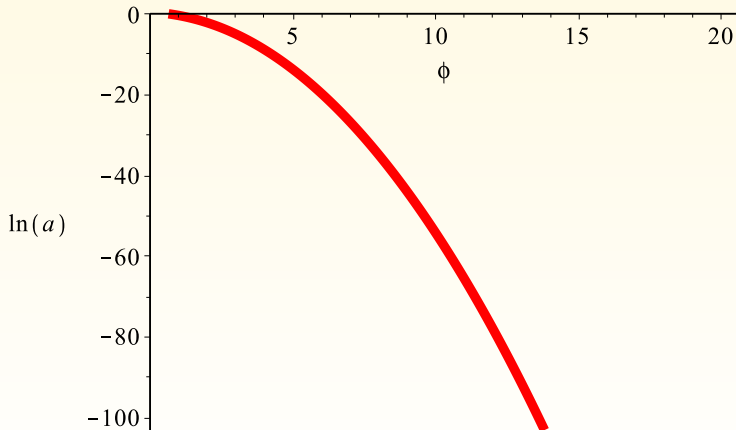
# TRANSFER FIELD TO CURVATURE FLUCTUATIONS



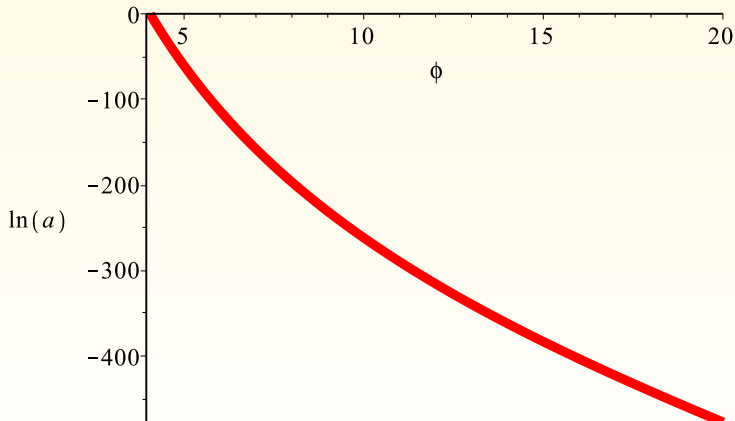
$$\Phi(\vec{x}) = F(\varphi(\vec{x}))$$



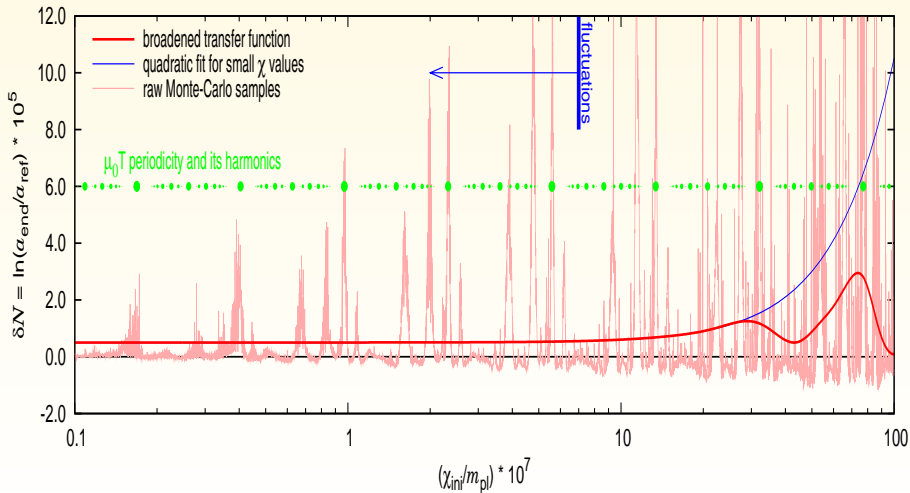
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Non-linear transfer function can be expanded (if smooth enough):

$$F(\varphi) = \sum_{n=1}^{\infty} f_n \varphi^n, \quad f_n \equiv \frac{F^{(n)}}{n!}$$

Odd-point correlators of Gaussian random field vanish:

$$\langle \varphi(\vec{x}) \varphi(\vec{y}) \rangle = \xi(\vec{x}, \vec{y})$$

$$\langle \varphi(\vec{x}) \varphi(\vec{y}) \varphi(\vec{z}) \rangle = 0$$

But non-linearity generates them (from even-point correlators of GRF):

$$\langle \Phi(\vec{x}) \Phi(\vec{y}) \Phi(\vec{z}) \rangle \neq 0$$

# THIS IS HOW IT WORKS...

$$\begin{aligned}\langle \Phi(\vec{x})\Phi(\vec{y})\Phi(\vec{z}) \rangle &= \left( \sum_{a=1}^{\infty} f_a \varphi^a(\vec{x}) \right) \left( \sum_{b=1}^{\infty} f_b \varphi^b(\vec{y}) \right) \left( \sum_{c=1}^{\infty} f_c \varphi^c(\vec{z}) \right) \\ &= \sum_{n=1}^{\infty} \sum_{a+b+c=n} f_a f_b f_c \underbrace{\langle \varphi^a(\vec{x}) \varphi^b(\vec{y}) \varphi^c(\vec{z}) \rangle}_{\text{only even } n \text{ contribute}} \\ &= f_1^2 f_2 \langle \varphi^2(\vec{x}) \varphi(\vec{y}) \varphi(\vec{z}) \rangle + \text{permutations} \\ &\quad + f_2^3 \langle \varphi^2(\vec{x}) \varphi^2(\vec{y}) \varphi^2(\vec{z}) \rangle \\ &\quad + f_1 f_2 f_3 \langle \varphi^3(\vec{x}) \varphi^2(\vec{y}) \varphi(\vec{z}) \rangle + \text{permutations} \\ &\quad + f_1^2 f_4 \langle \varphi^4(\vec{x}) \varphi(\vec{y}) \varphi(\vec{z}) \rangle + \text{permutations} \\ &\quad + \dots\end{aligned}$$

# BACK TO FOURIER DOMAIN

$$\langle \varphi(\vec{k}_1) \varphi(\vec{k}_2) \rangle = P(\vec{k}_1) \delta^3(\vec{k}_1 + \vec{k}_2)$$

$$\langle \Phi(\vec{k}_x) \Phi(\vec{k}_y) \Phi(\vec{k}_z) \rangle = \iiint_{\vec{x}, \vec{y}, \vec{z}} \langle \Phi(\vec{x}) \Phi(\vec{y}) \Phi(\vec{z}) \rangle e^{i(\vec{k}_x \vec{x} + \vec{k}_y \vec{y} + \vec{k}_z \vec{z})}$$

$$\langle \varphi^2(\vec{x}) \varphi(\vec{y}) \varphi(\vec{z}) \rangle \rightarrow \iiint \iiint_{\vec{k}_1, \vec{k}_2, \vec{k}_3, \vec{k}_4} \langle \varphi(\vec{k}_1) \varphi(\vec{k}_2) \rangle \langle \varphi(\vec{k}_3) \varphi(\vec{k}_4) \rangle \times \\ \delta^3(\vec{k}_1 + \vec{k}_2 - \vec{k}_x) \delta^3(\vec{k}_3 - \vec{k}_y) \delta^3(\vec{k}_4 - \vec{k}_z)$$

$$= \iiint \iiint_{\vec{k}_1, \vec{k}_2, \vec{k}_3, \vec{k}_4} P(\vec{k}_1) P(\vec{k}_3) \times \delta^3(\vec{k}_1 + \vec{k}_2) \delta^3(\vec{k}_3 + \vec{k}_4) \\ \delta^3(\vec{k}_1 + \vec{k}_2 - \vec{k}_x) \delta^3(\vec{k}_3 - \vec{k}_y) \delta^3(\vec{k}_4 - \vec{k}_z)$$

$$= P(k_x) P(k_y) \delta^3(\vec{k}_x + \vec{k}_y + \vec{k}_z)$$

+ permutations + connected pieces



# MORE SHAPES OF LOCAL NON-GAUSSIANITY!

$$\langle \varphi^4(\vec{x})\varphi(\vec{y})\varphi(\vec{z}) \rangle \rightarrow \int d^3k P(\vec{k}) (P(k_x)P(k_y) + \text{permutations})$$

$$\langle \varphi^3(\vec{x})\varphi^2(\vec{y})\varphi(\vec{z}) \rangle \rightarrow \int d^3k P(\vec{k}) (P(\vec{k}_x - \vec{k})P(k_y) + \text{permutations})$$

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How do these shapes look like?

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# IS NON-"LOCAL" SHAPE REALLY LOCAL?

$$B_{\sigma^3}(\vec{k}_i) = \frac{1}{k_1 k_2 k_3 k_t^3}, \quad k_t = k_1 + k_2 + k_3$$

