

Evolution of the most massive AGB star as a progenitor for ECSNe

Koh Takahashi¹, Takashi Yoshida², Hideyuki Umeda¹

¹Department of Astronomy, The University of Tokyo

²Yukawa Institute for Theoretical Physics, Kyoto University

Abstract

The most massive asymptotic giant branch (**AGB**) star can form a critical mass ONe core at its center and core collapse of such a critical ONe core ends up as an electron capture supernova (**ECSN**).

In this talk, I will present evolutionary properties of the most massive AGB star as a progenitor system of ECSNe.

Abstract

The model of ECSNe has plenty of important suggestions on both observation and theory. Explosion properties which can be inferred by considering the pre-explosion structure of an AGB star provide a possible explanation for some **low-energetic supernovae** (e.g. SN2005cs) and also **a part of type II_n supernovae** (e.g. SN2008S). Several authors discuss that **the Crab** supernova, SN1054, can be a likely candidate of ECSN as well. As for theory, ECSN model is **the only exception** for which a one dimensional simulation can model a successful explosion a priori. However, there had been only one evolutionary calculation for ECSN, **done by Nomoto in 1987**, and thus a modern progenitor calculation as a basis of theoretical investigation is strongly demanded. Recently, we have accomplished a progenitor calculation for ECSN **for the first time in twenty six years** and updated a pre-explosion structure of the model. Evolutionary sequences for this massive AGB star, which is located at (Abridged.)

Abstract

What I want to say...

Abstract

1. ECSN : interesting object

Abstract

2. only 1 progenitor calculation

Abstract

3. New calculation is needed

Abstract

4. We have done it.

Abstract

published in ApJ 771 28
(arXiv : **1302.6402**)

THE ASTROPHYSICAL JOURNAL, 771:28 (13pp), 2013 July 1

doi:[10.1088/0004-637X/771/1/28](https://doi.org/10.1088/0004-637X/771/1/28)

© 2013. The American Astronomical Society. All rights reserved. Printed in the U.S.A.

EVOLUTION OF PROGENITORS FOR ELECTRON CAPTURE SUPERNOVAE

KOH TAKAHASHI¹, TAKASHI YOSHIDA², AND HIDEYUKI UMEDA¹

¹ Department of Astronomy, The University of Tokyo, Tokyo 113-0033, Japan; ktakahashi@astron.s.u-tokyo.ac.jp, umeda@astron.s.u-tokyo.ac.jp

² Yukawa Institute for Theoretical Physics, Kyoto University, Kyoto 606-8502, Japan; yoshida@yukawa.kyoto-u.ac.jp

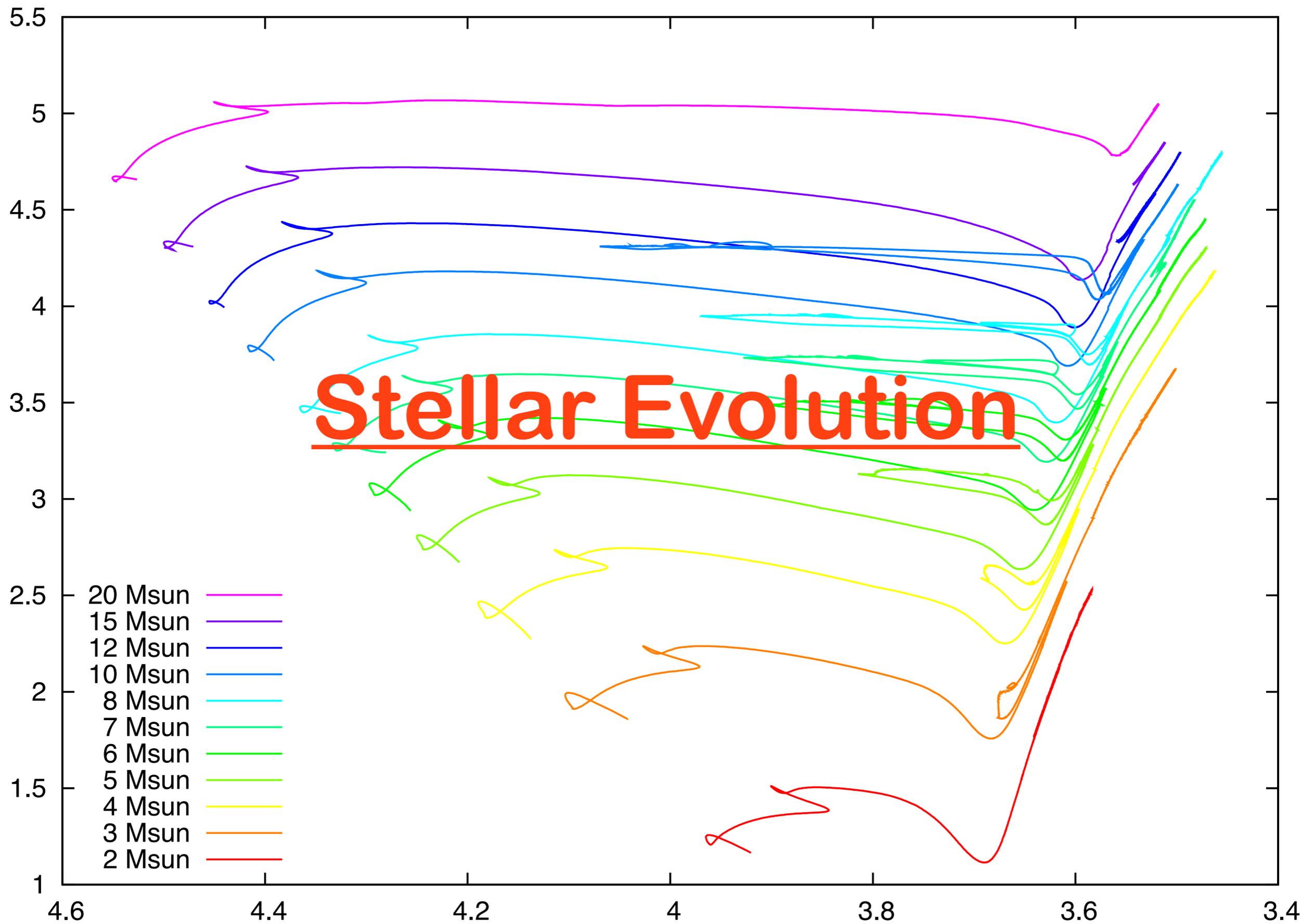
Received 2013 February 26; accepted 2013 May 6; published 2013 June 12

ABSTRACT

We provide progenitor models for electron capture supernovae (ECSNe) with detailed evolutionary calculation. We include minor electron capture nuclei using a large nuclear reaction network with updated reaction rates. For electron capture, the Coulomb correction of rates is treated and the contribution from neutron-rich isotopes is taken into account in each nuclear statistical equilibrium (NSE) composition. We calculate the evolution of the most massive super asymptotic giant branch stars and show that these stars undergo off-center carbon burning and form ONe cores at the center. These cores become heavier up to the critical mass of $1.367 M_{\odot}$ and keep contracting even after the initiation of O+Ne deflagration. Inclusion of minor electron capture nuclei causes convective URCA cooling during the contraction phase, but the effect on the progenitor evolution is small. On the other hand, electron capture by neutron-rich isotopes in the NSE region has a more significant effect. We discuss the uniqueness of the critical core mass for ECSNe and the effect of wind mass loss on the plausibility of our models for ECSN progenitors.

Key words: nuclear reactions, nucleosynthesis, abundances – stars: evolution – stars: interiors – supernovae: general

Online-only material: color figures



Stellar Evolution

○ Basic Equations

- mass conservation

$$\frac{dr}{dM} = \frac{1}{4\pi r^2 \rho}$$

- momentum conservation

$$\frac{dp}{dM} = -\frac{GM}{4\pi r^4} - \frac{M}{4\pi r^2 \rho} \frac{d^2 r}{dt^2}$$

- energy conservation

$$\frac{dL}{dM} = -T \frac{ds}{dt} - \mu \frac{dn}{dt} + \epsilon_{\text{nuc}} - \epsilon_{\text{v}} + \epsilon_{\text{ec}}$$

- energy transport

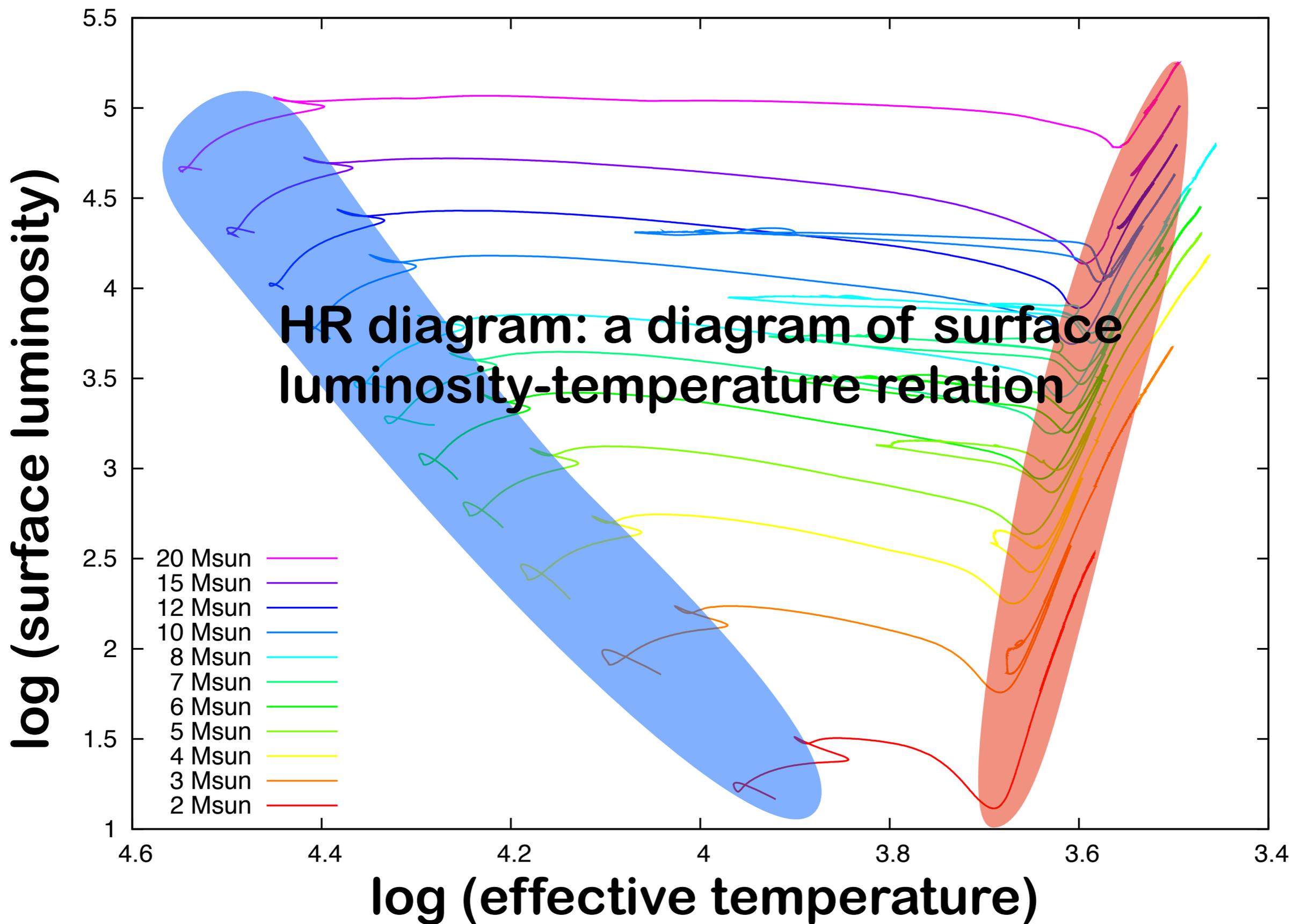
$$\frac{dT}{dM} = -\frac{GMT}{4\pi r^4 \rho} [\nabla_{\text{rad}} \text{ or } \nabla_{\text{conv}}]$$

+ nuclear reaction network

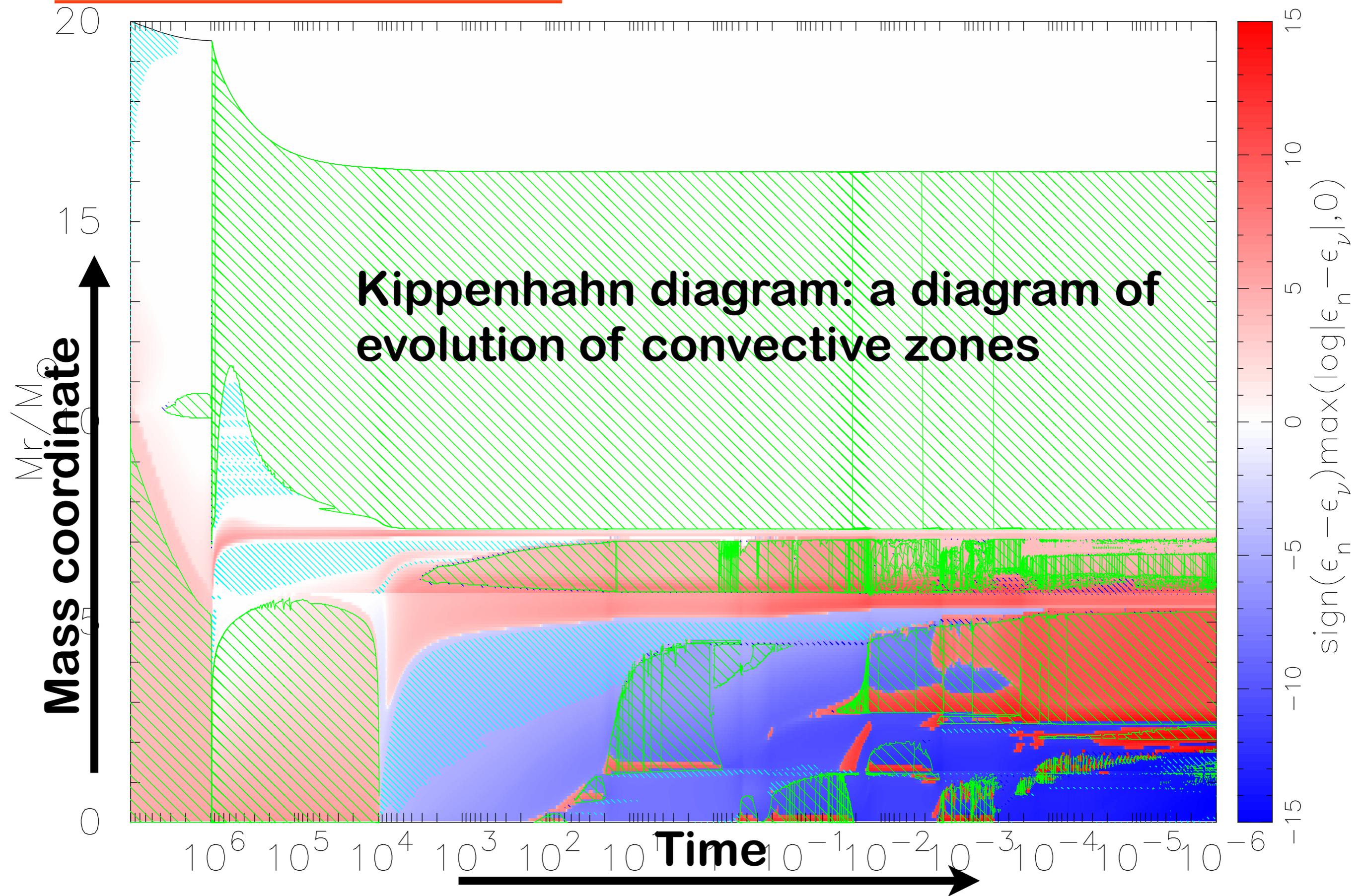
most stellar evolutionary calculations solve these equations iteratively.

then we get evolving structure as a solution.

Stellar Evolution



Stellar Evolution



Stellar Evolution

Massive stars
($M \gtrsim 10 M_{\text{sun}}$)

**AIM: to get a pre-explosion
structure of a core-collapse
supernova**

INITIAL MASS

8

10

20

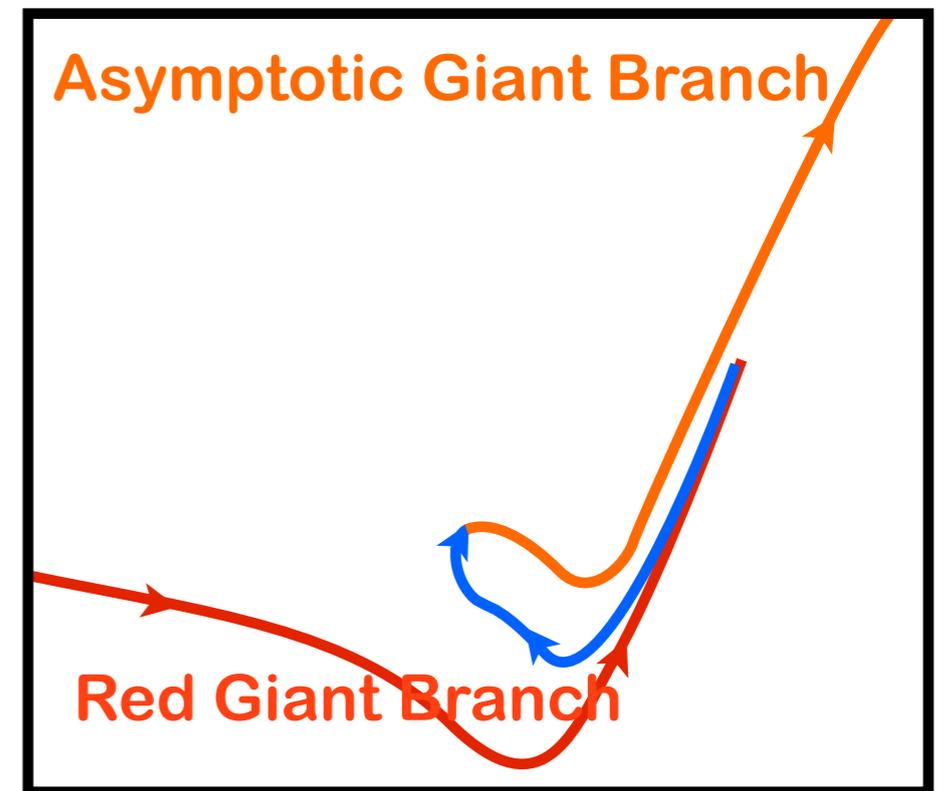
200



Stellar Evolution

Intermediate mass stars ($M \sim 2-10 M_{\text{sun}}$)

evolve as asymptotic giant
branch (**AGB**) stars
AIM: s-process nucleosynthesis,
dust formation, etc.



INITIAL MASS

2

4

8

10

Stellar Evolution

ECSN progenitor ($M \sim 10 M_{\text{sun}}$)

is massive AGB stars,
so, it is from the mass range of...

- **the most massive** intermediate-mass star
- **the least massive** massive star



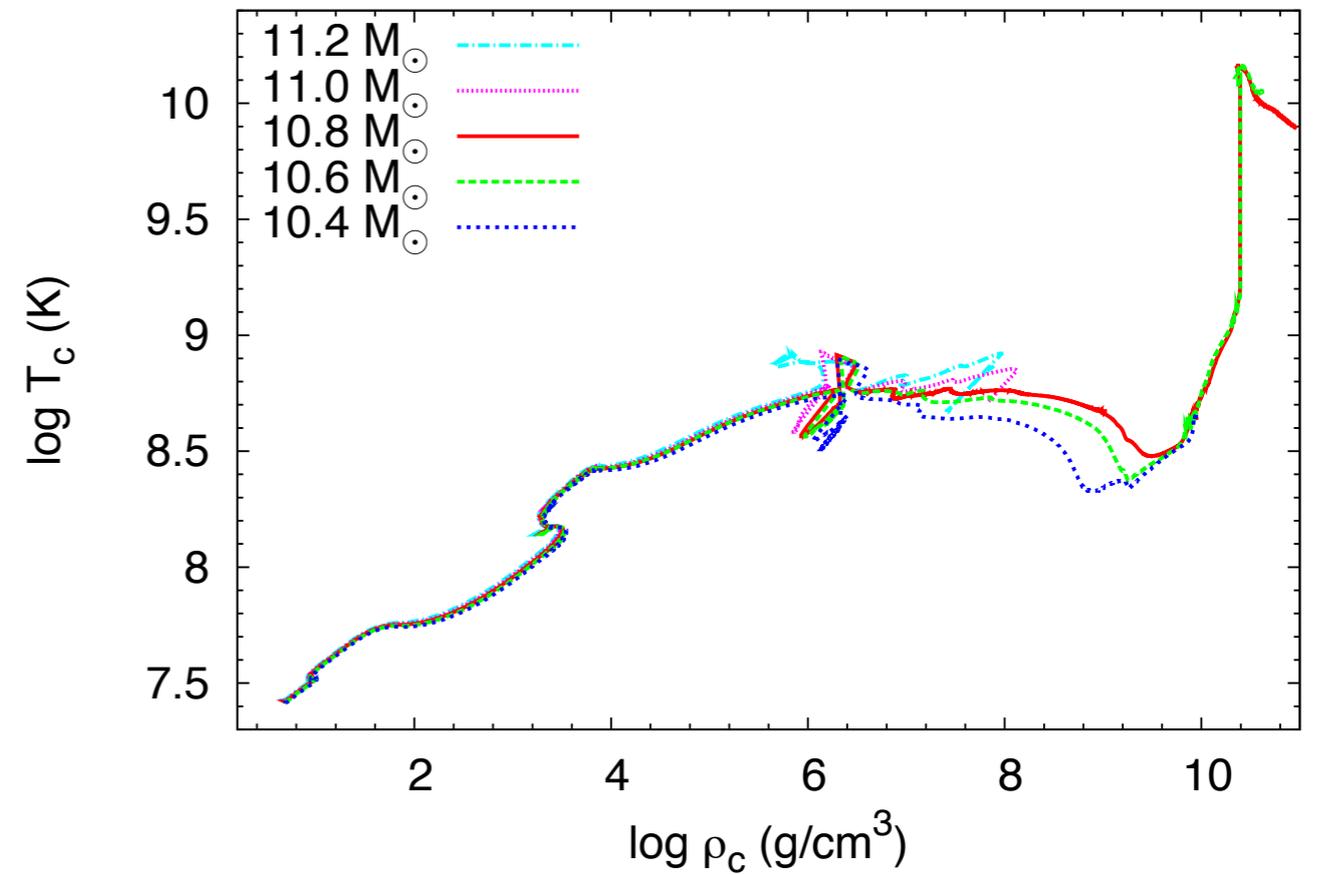
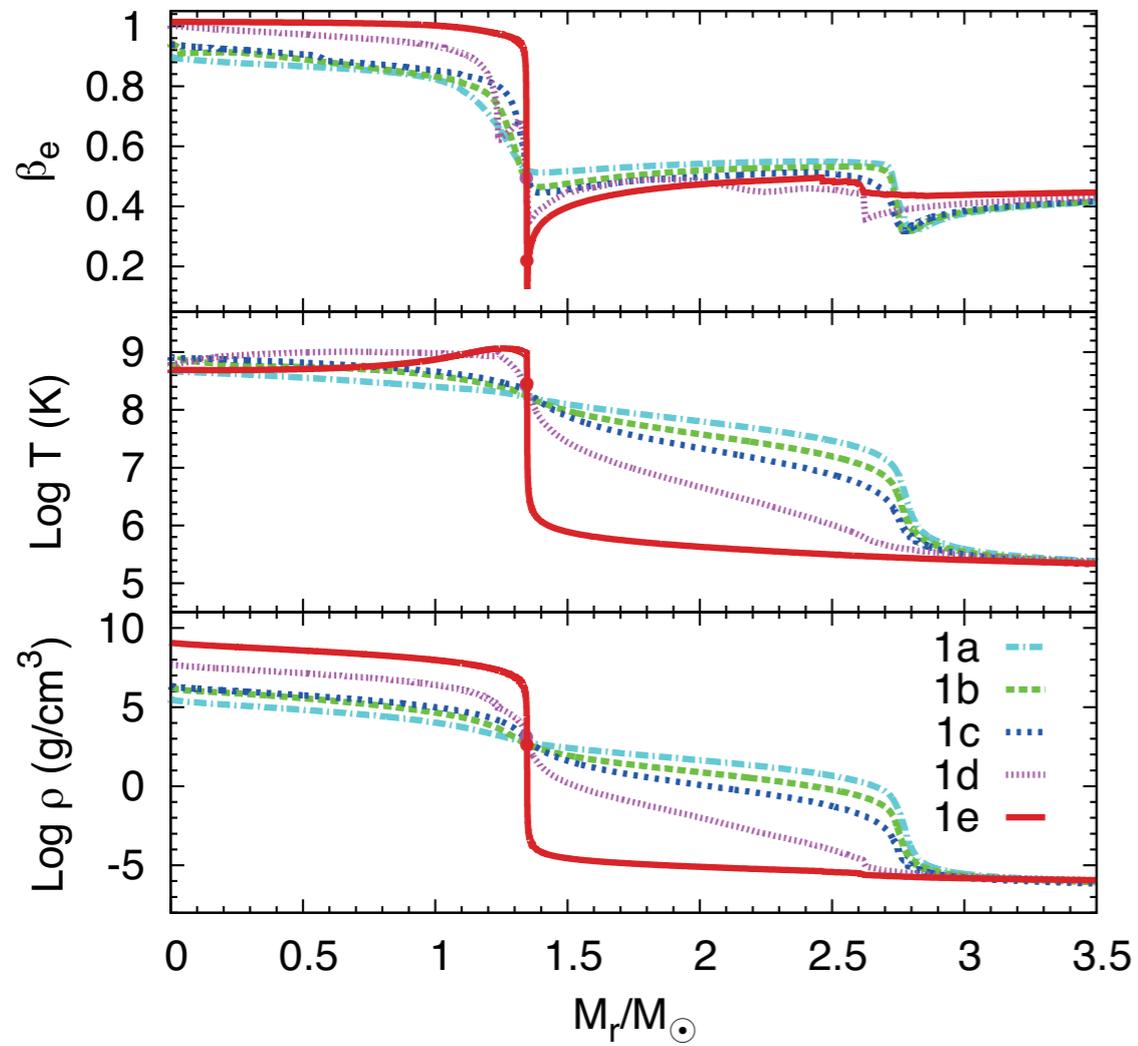
Stellar Evolution

It means that

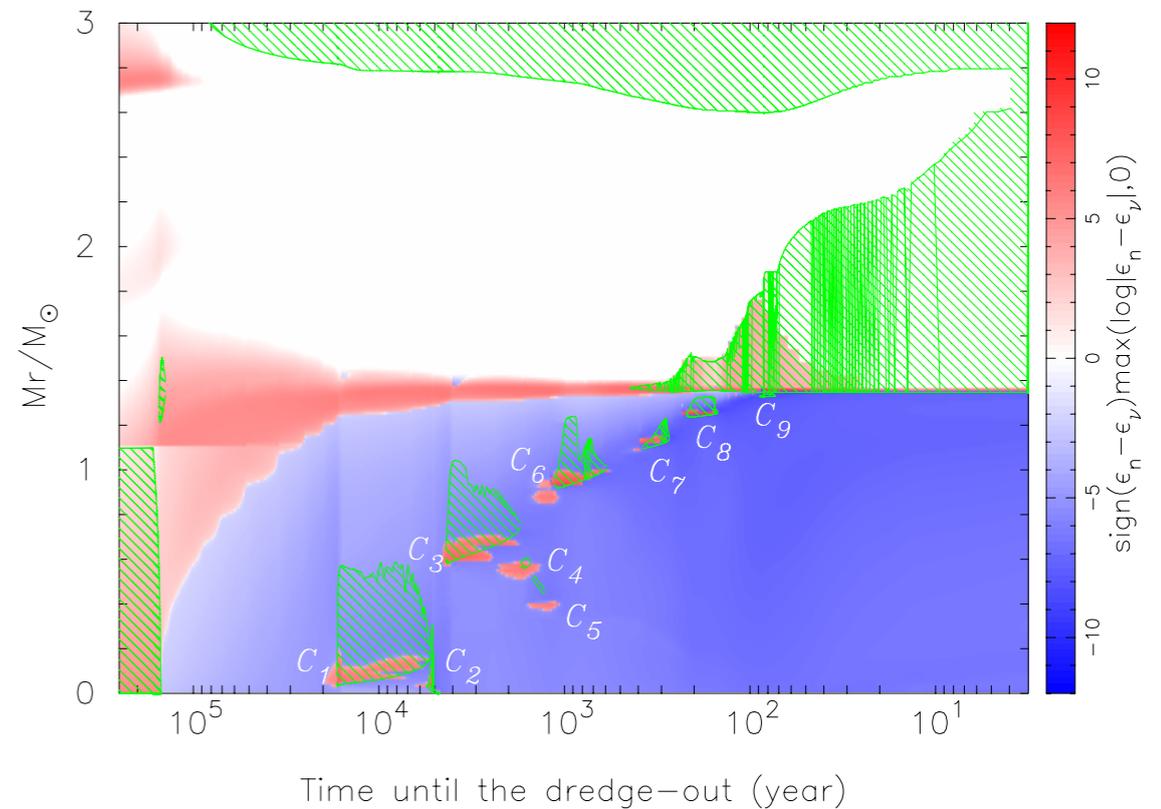
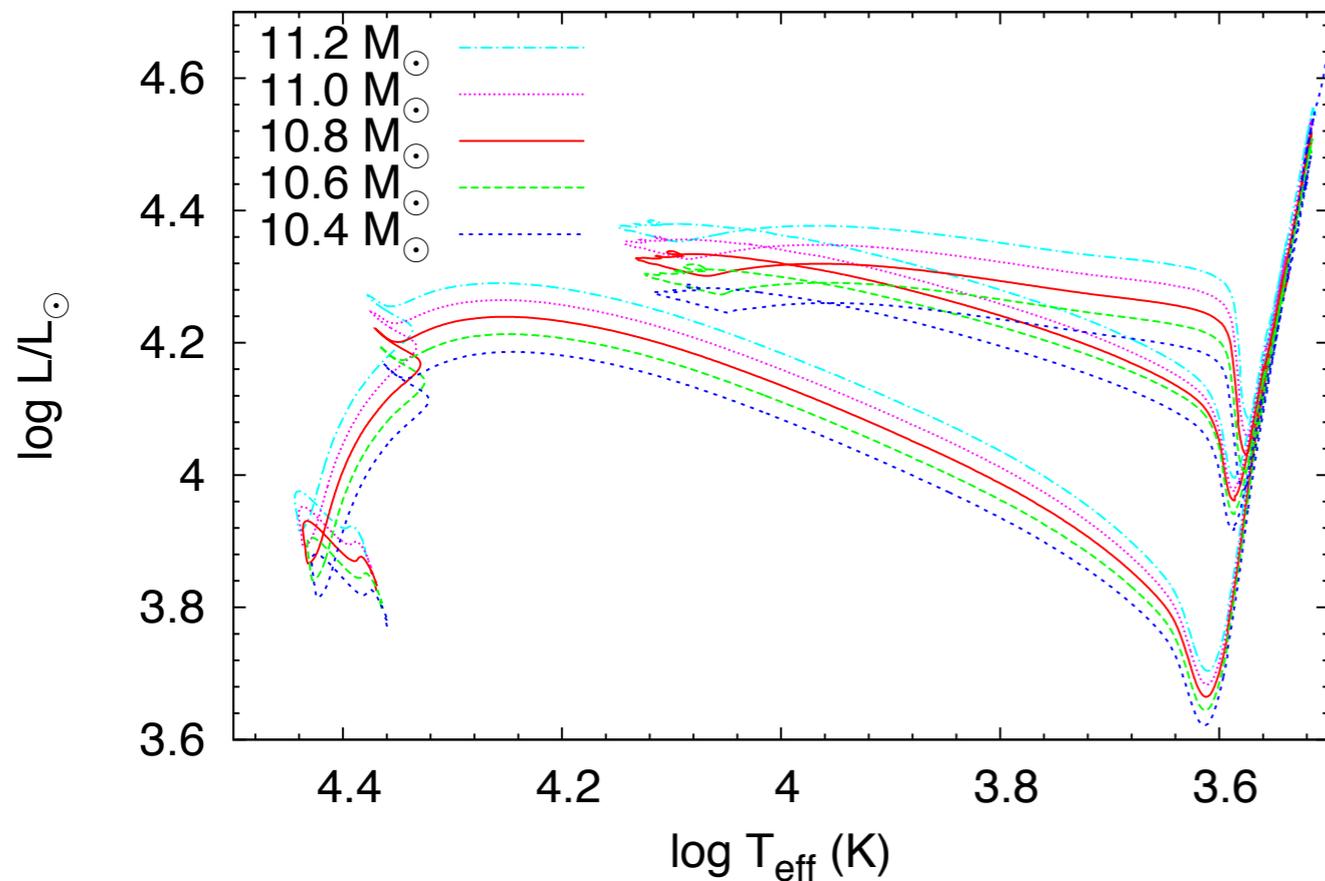
Progenitor evolution of ECSN is quite distinct from other (both massive and intermediate mass) stars.

- we should treat some particular physics





SAGB evolution

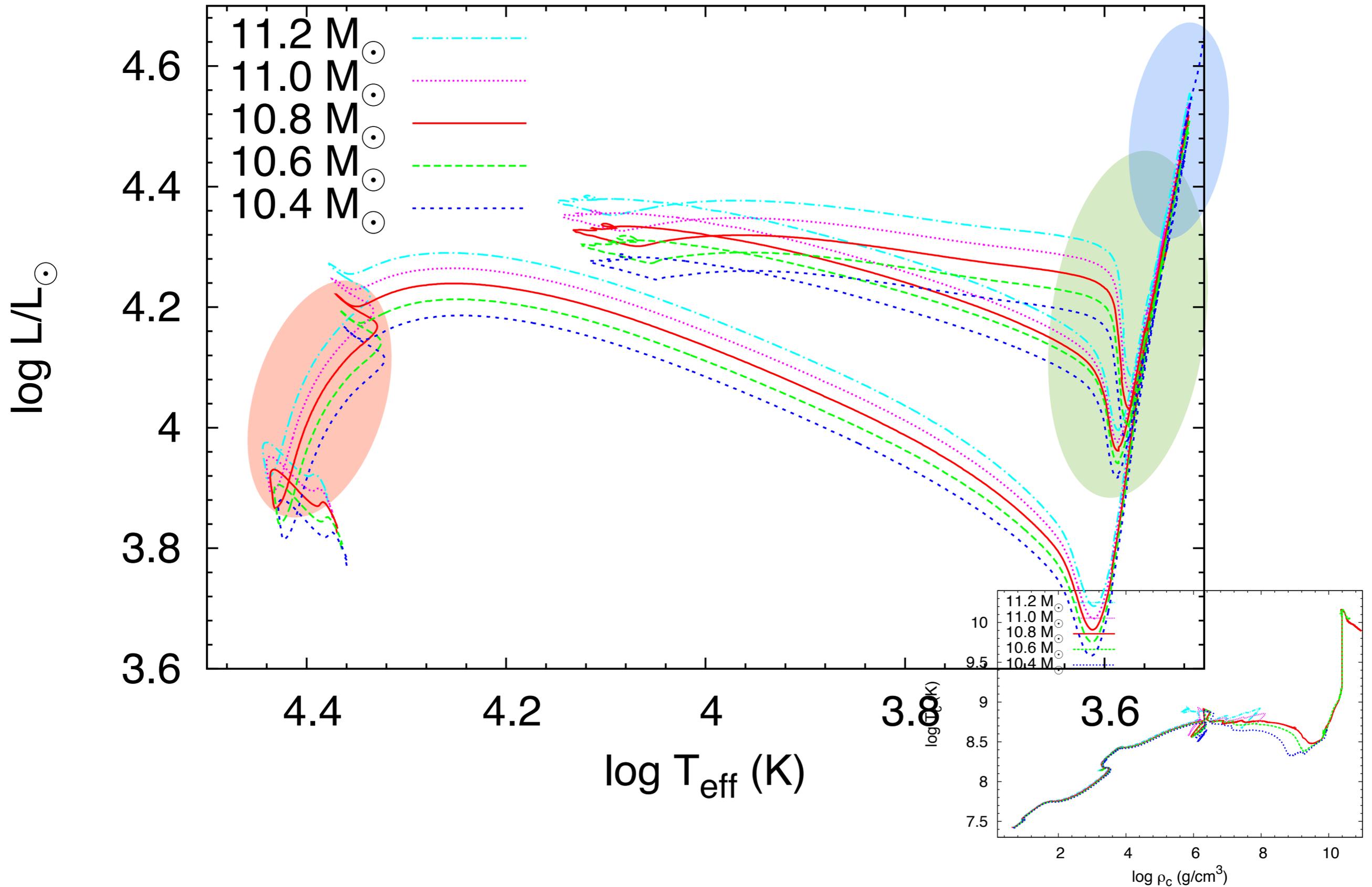


SAGB Evolution

- What is the particularity?
- Qualitative expectation of ECSN properties

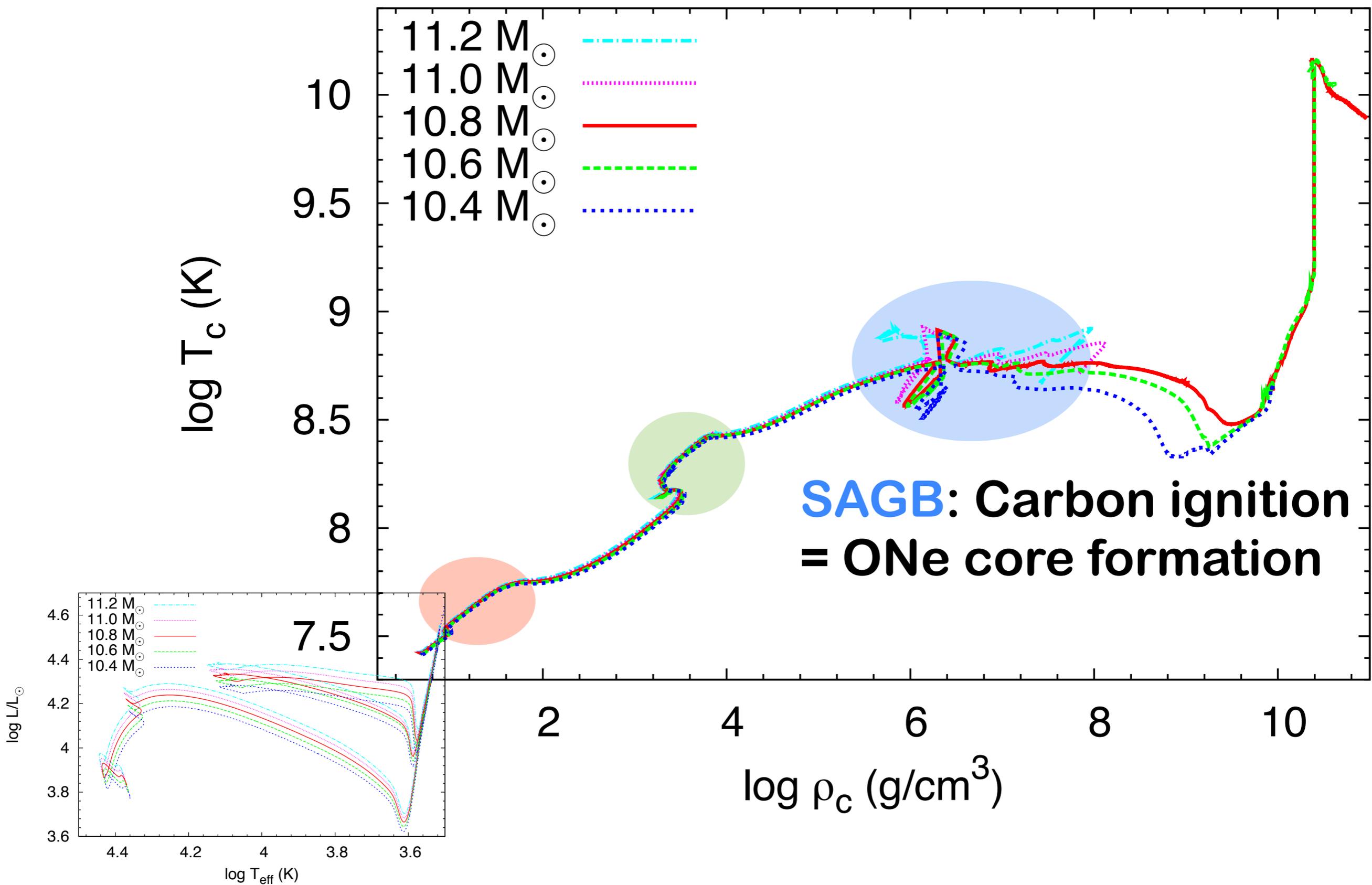
SAGB Evolution

HR diagram: a diagram of surface luminosity-temperature relation



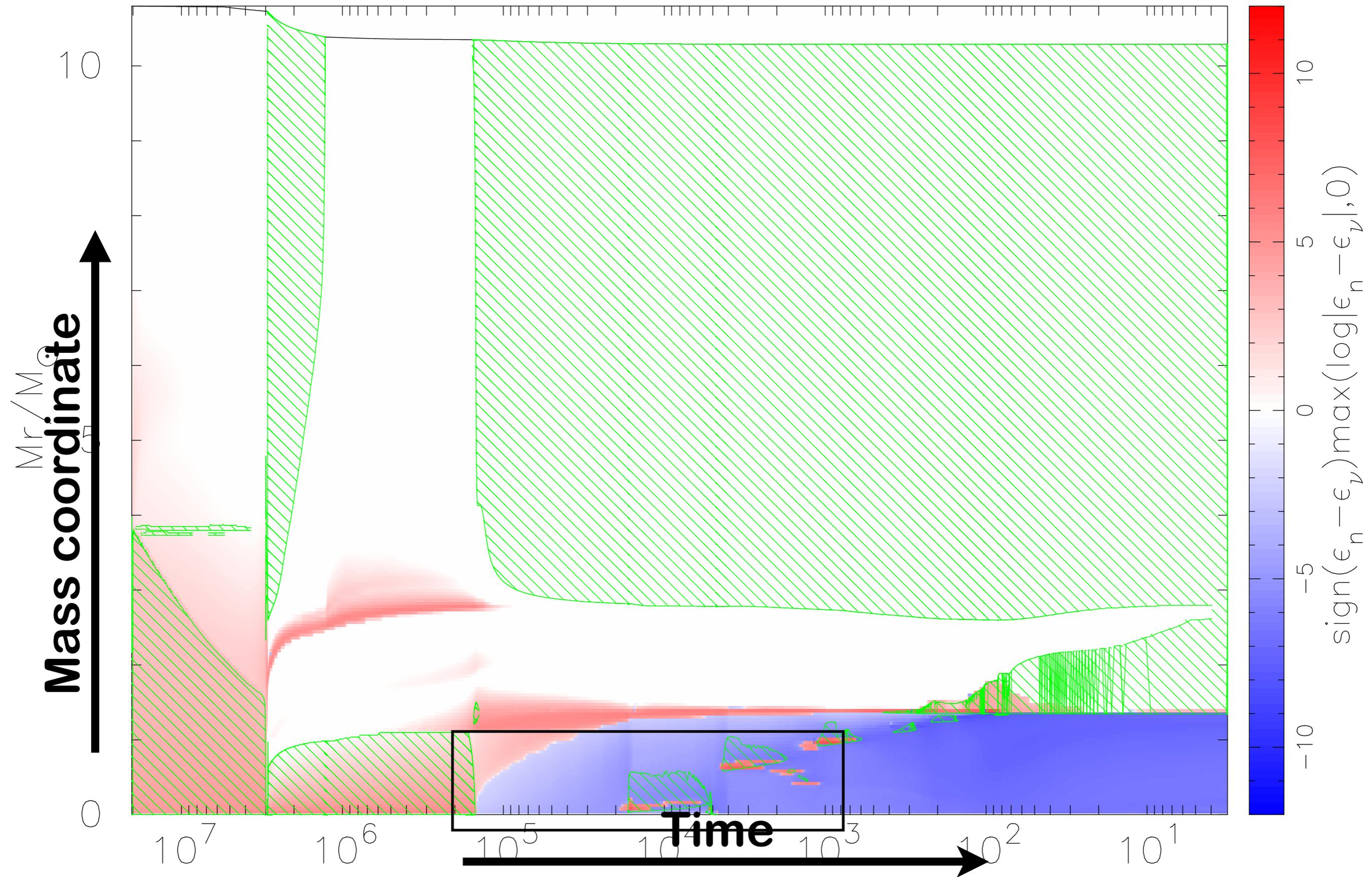
SAGB Evolution

ρ_c - T_c diagram: a diagram of central density-temperature relation

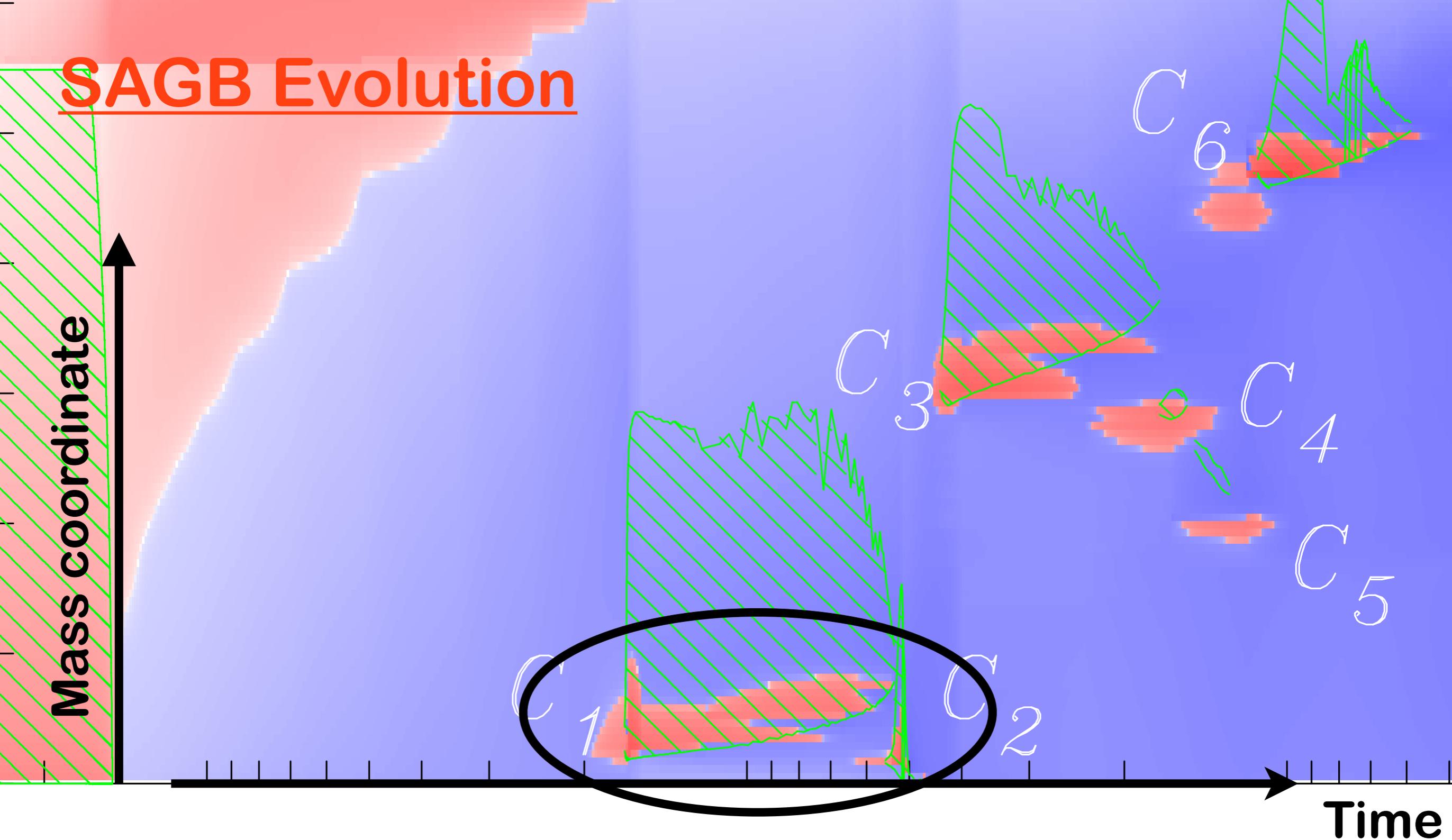


SAGB Evolution

Kippenhahn diagram: a diagram evolution of convective zones



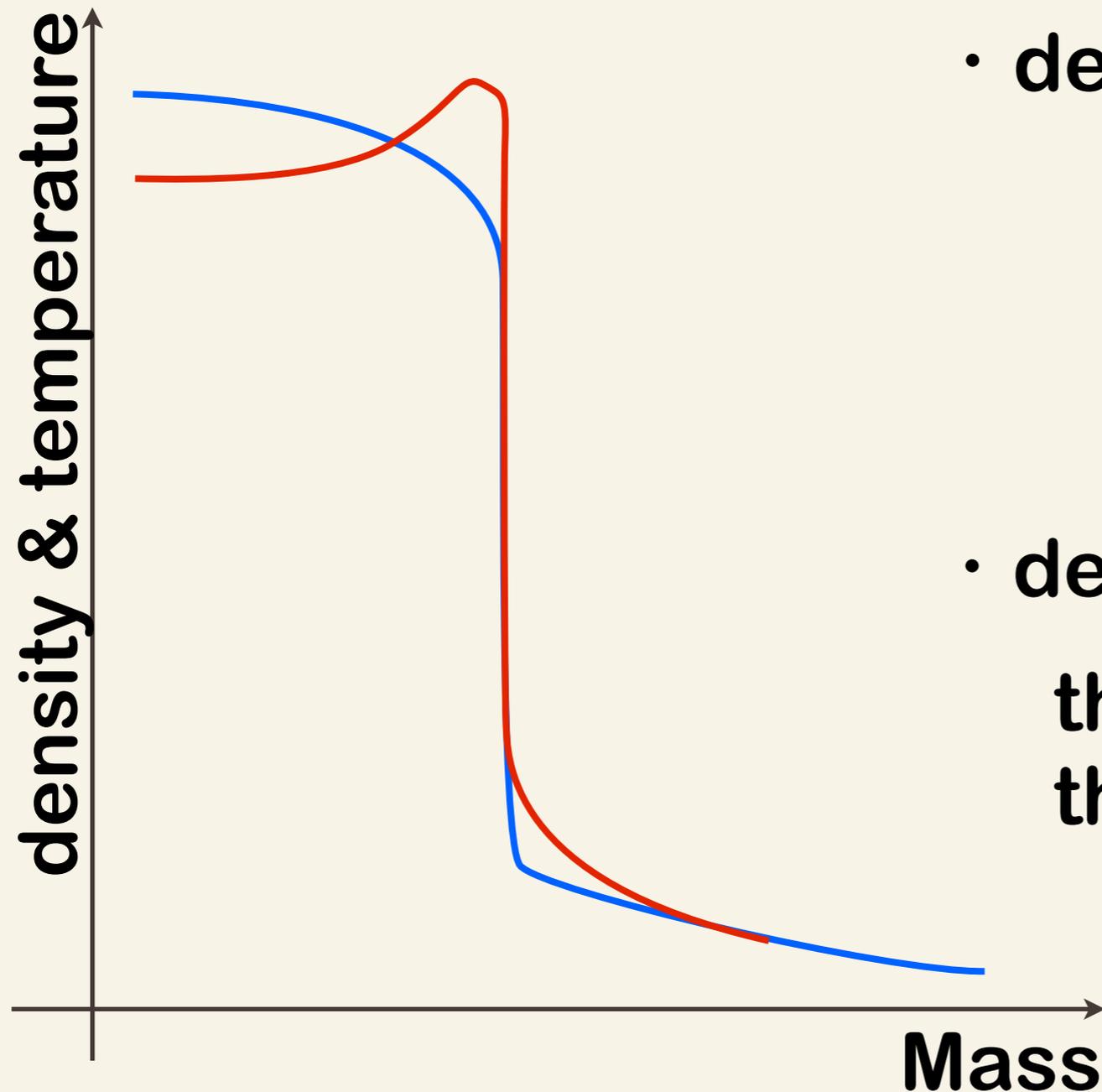
SAGB Evolution



Carbon ignites
off-centrally

- **temperature inversion**
due to degeneracy of
electrons

Degeneracy of electrons



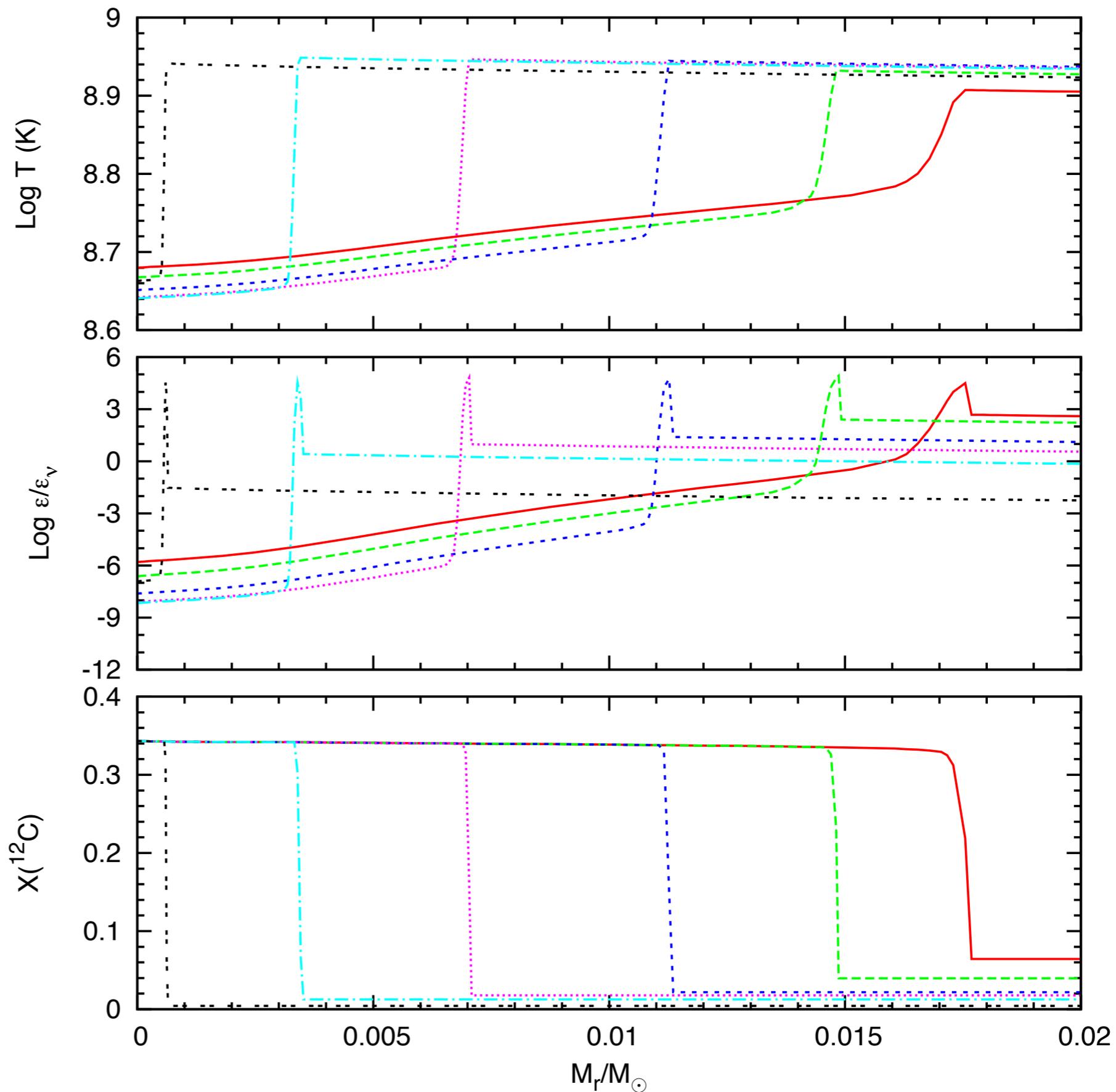
- degenerate electrons

$$\begin{aligned} p_e &= \frac{1}{3} \int \frac{p_\epsilon v_\epsilon}{\exp\left(\frac{\epsilon - \mu}{kT}\right) + 1} d\epsilon \\ &= p_e(\rho Y_e) \end{aligned}$$

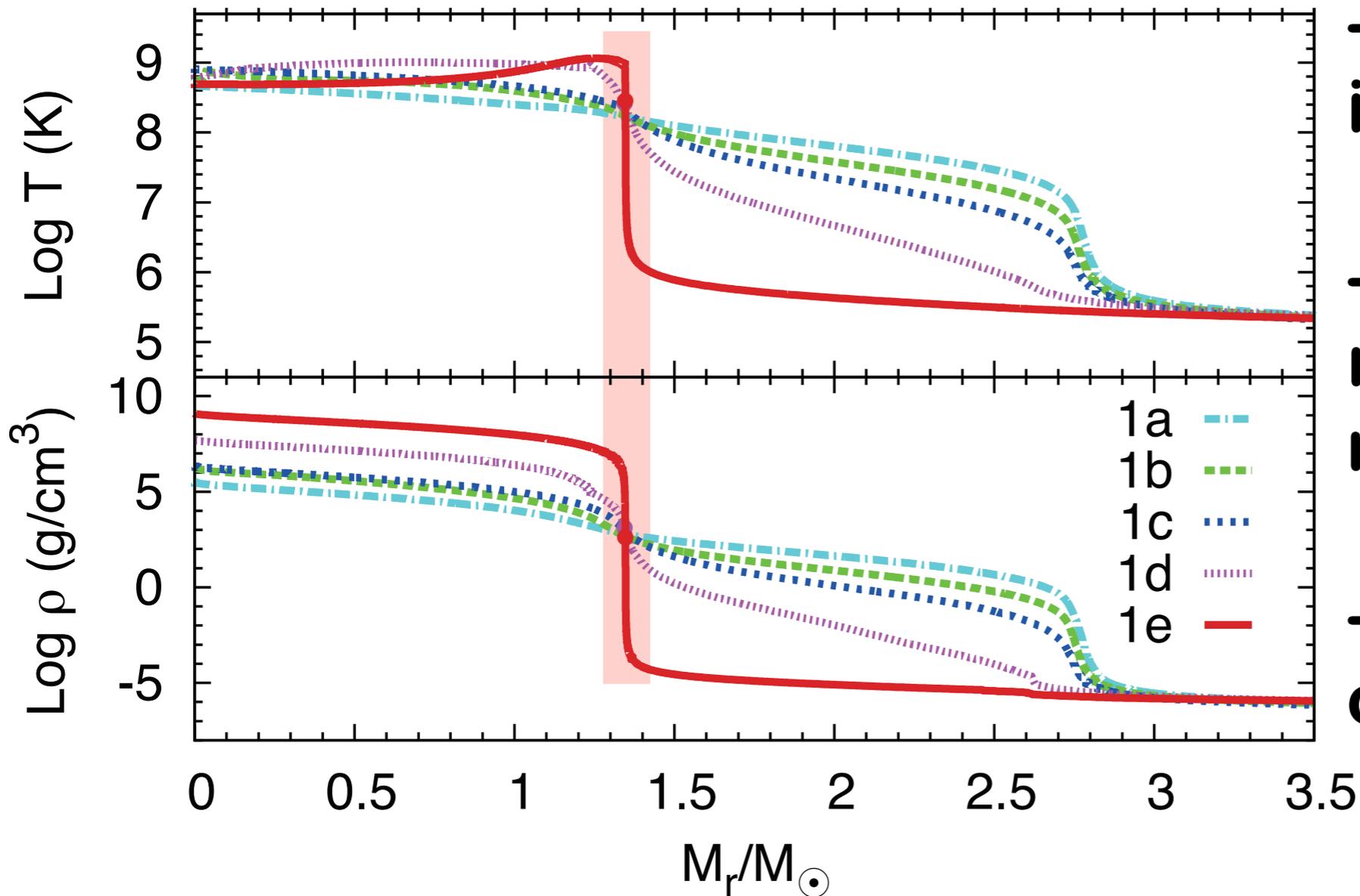
- decoupling of
the **dynamical** structure
the **thermal** structure

SAGB Evolution

- following flame propagation needs high resolution



SAGB Evolution



- degenerate core is compact

- surrounding region becomes much hotter

- energetic reaction occurs at the edge

- Energetic shell burning results in a highly contrasting structure

- A $\sim 1.37 M_{\text{sun}}$ ONe core with diffuse H+He envelope

SAGB Evolution

What comes next?

- The core becomes heavier up to the critical mass.
- OR loses its whole envelope by intense mass loss.

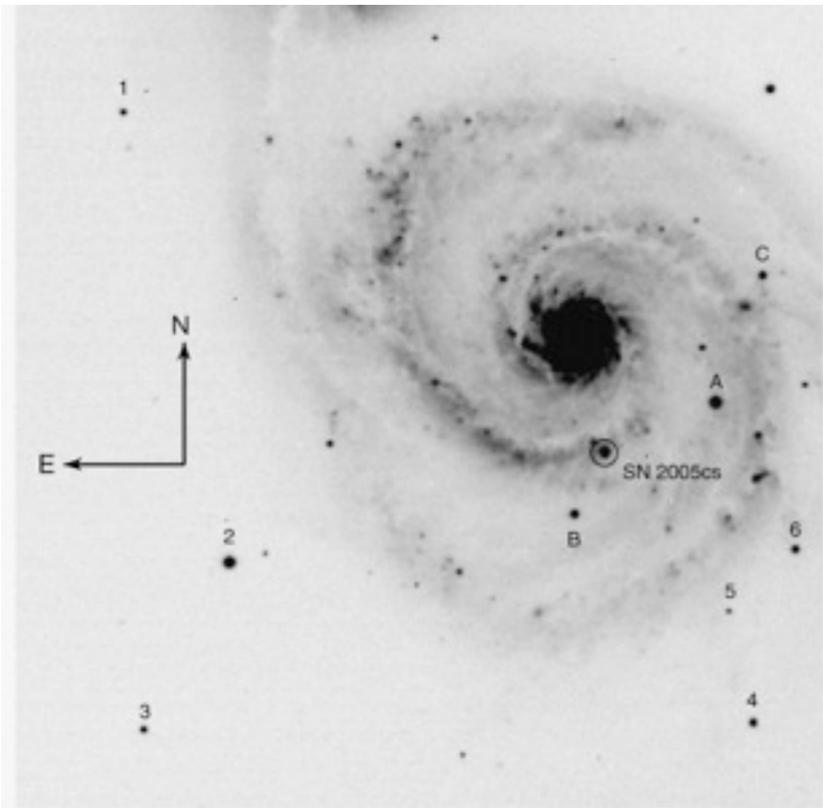
SAGB Evolution

Short summary

(the answer for the first question)

- SAGB star is an AGB star with an ONe core.
- Electron degeneracy permit carbon ignites off-centrally.
- Energetic shell burning occur around the degenerate core.
- A SAGB star has a contrasting structure.

ECSN candidates ?



Pastorello A et al. MNRAS
2006;370:1752-1762

SN 2008S (Botticella et al. 2009)

- SN 2005cs (**SN IIP**) may be an ECSN
 - subluminous and subenergetic SN
 - **the most less massive core**
- SN 2008S (**SN IIn**) may be an ECSN
 - the progenitor star of $\sim 10 M_{\text{sun}}$
 - surrounded by dense CSM
 - **the large mass loss**
- ECSN can explain **the Crab**
 - small explosion energy (~ 0.1 foe)
 - large $X(\text{He})$, small $X(\text{O})$
 - **Not have any substructure**

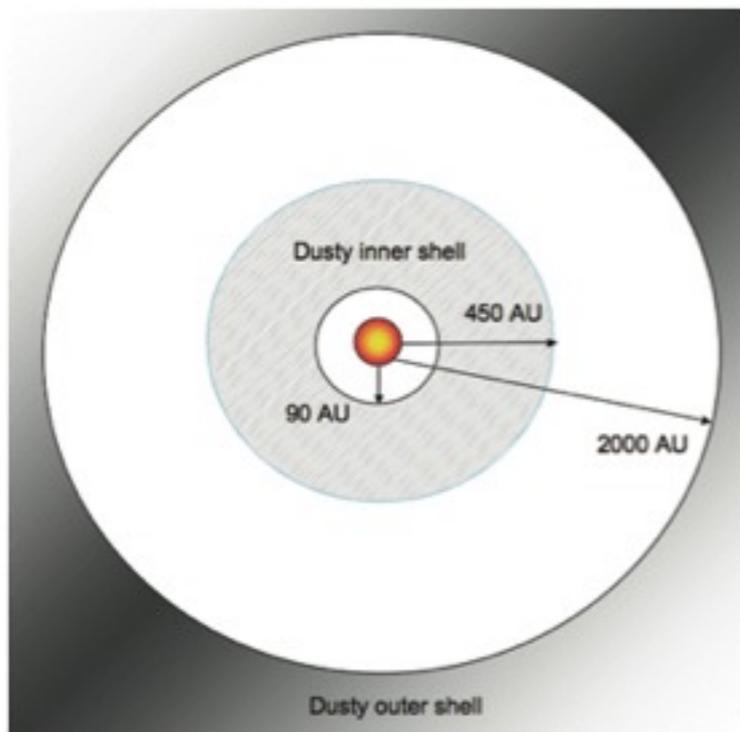
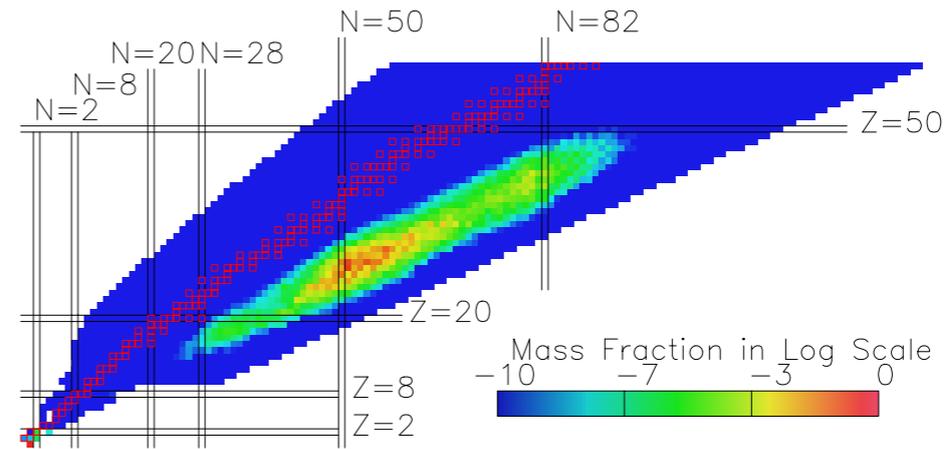
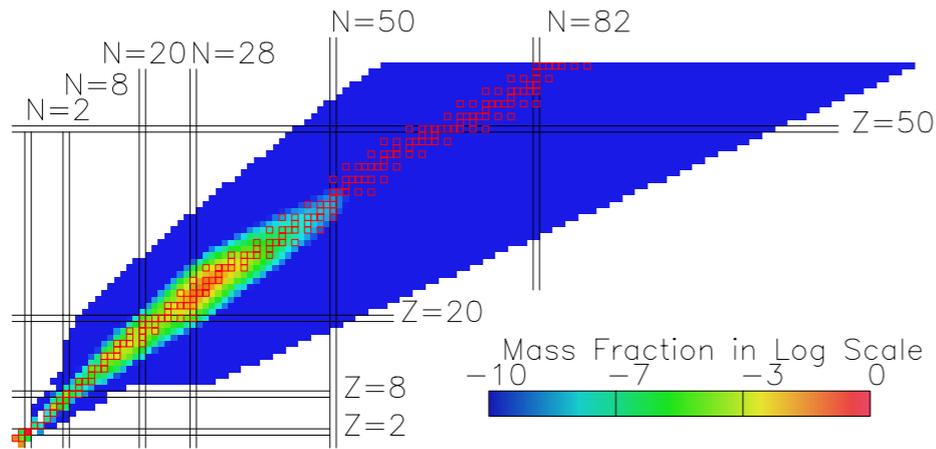


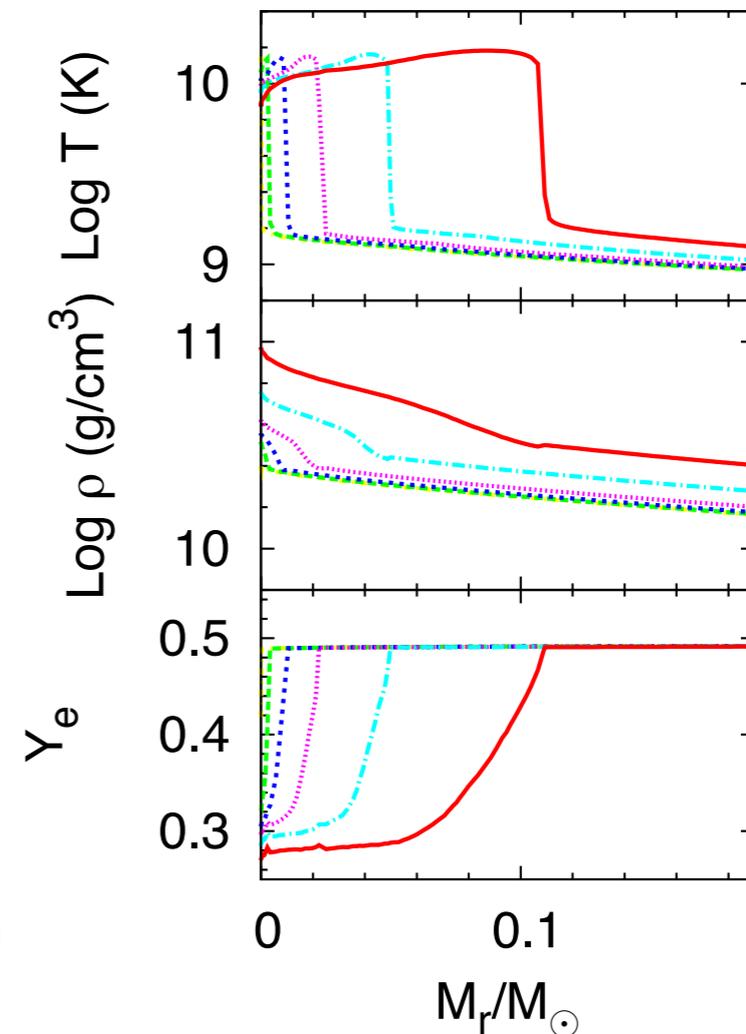
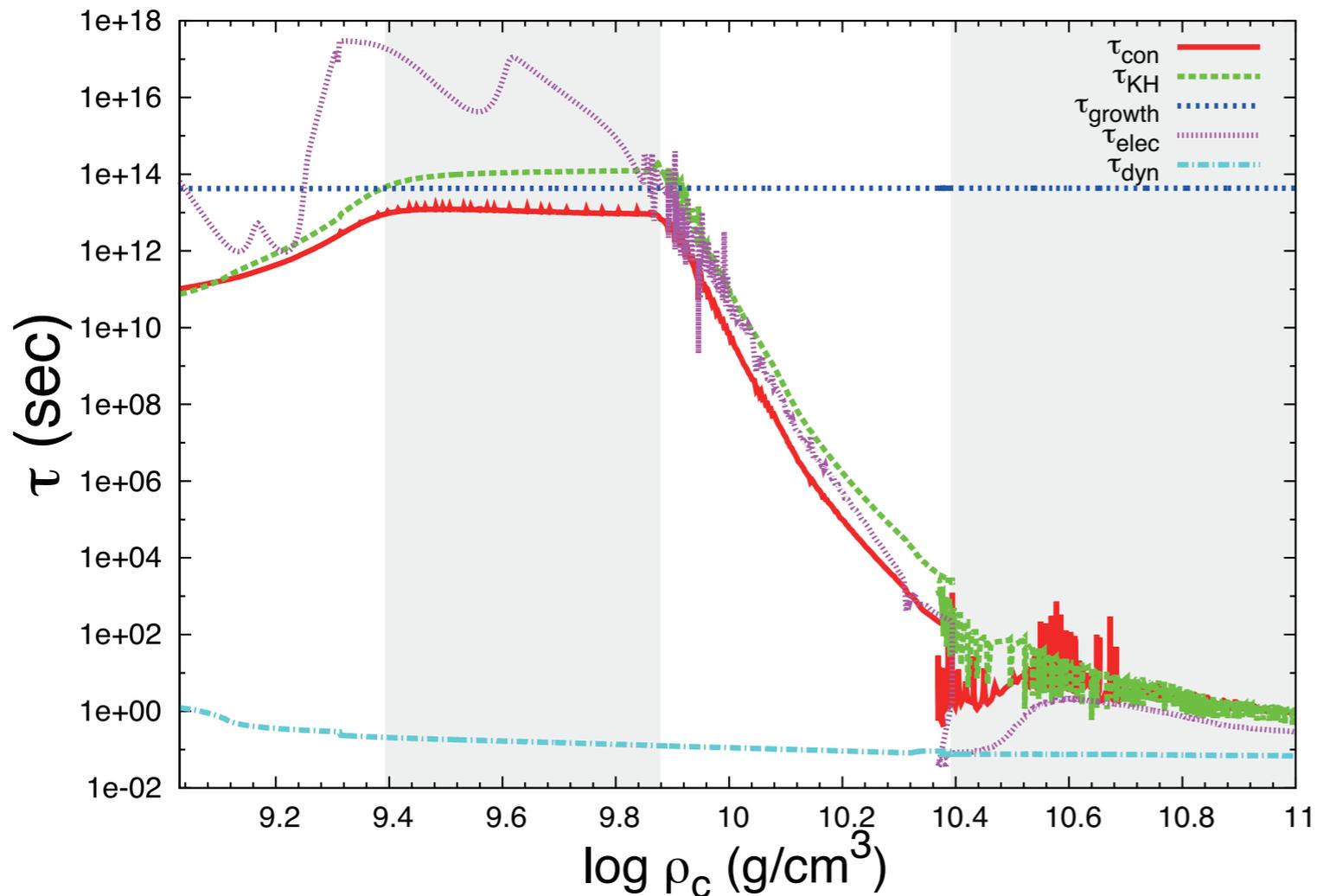
Figure 20. Schematic illustration of the pre-explosion geometry of SN 2008S.

include 3091 species of nuclei
 $T_g = 1.00E+01$; $\rho = 1.00E+09 \text{g/cm}^3$; $Y_e = 4.90E-01$
 $Y_p = 1.90E-02$; $Y_n = 1.28E-02$; $Y_\alpha = 1.65E-01$
 $Y_h = 5.84E-03$; $Y_i = 2.02E-01$

include 3091 species of nuclei
 $T_g = 1.00E+01$; $\rho = 1.00E+11 \text{g/cm}^3$; $Y_e = 2.70E-01$
 $Y_p = 2.13E-09$; $Y_n = 2.34E-01$; $Y_\alpha = 7.11E-07$
 $Y_h = 9.36E-03$; $Y_i = 2.44E-01$



Until core collapse



Until core collapse

Critical mass objects

- CO WD -> Thermonuclear explosion
- Fe core -> Core collapse
- ONe core -> Core collapse (or AIC)

Until core collapse

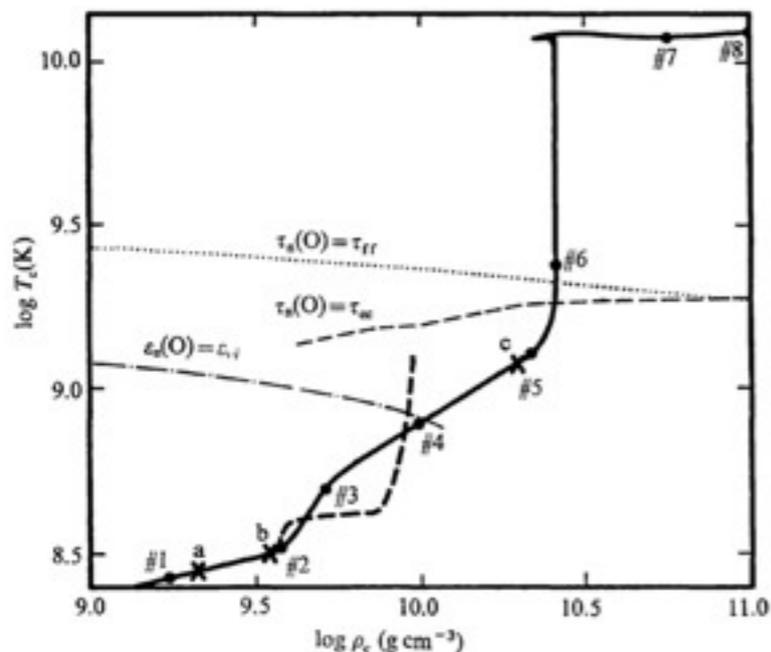
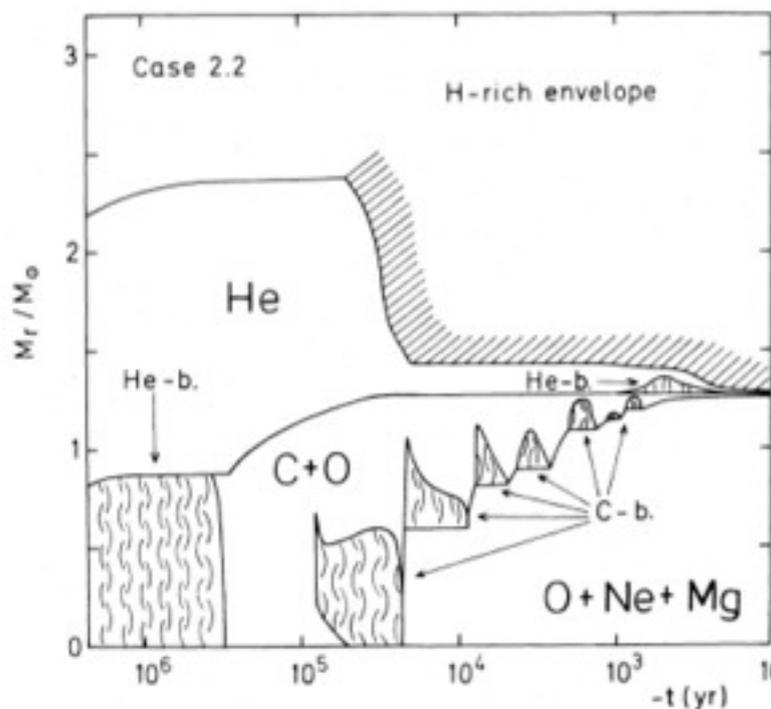


Fig. 5. Evolutionary change in the central temperature and the central density. See the text for stages a, b, and c. If the occurrence of the convective core were neglected, it would run as indicated by the thick dashed curve. Thin lines are concerned with the conditions for oxygen burning (see the text for detail).

(Miyaji et al. 1980)

Evolution of the critical mass
ONe core was firstly solved by
Miyaji et al. in 1980.

In 1987, Nomoto showed that
such ONe core collapse can
happen in a star
which have a He
core of $\sim 2.2 M_{\text{sun}}$.



(Nomoto 1987)

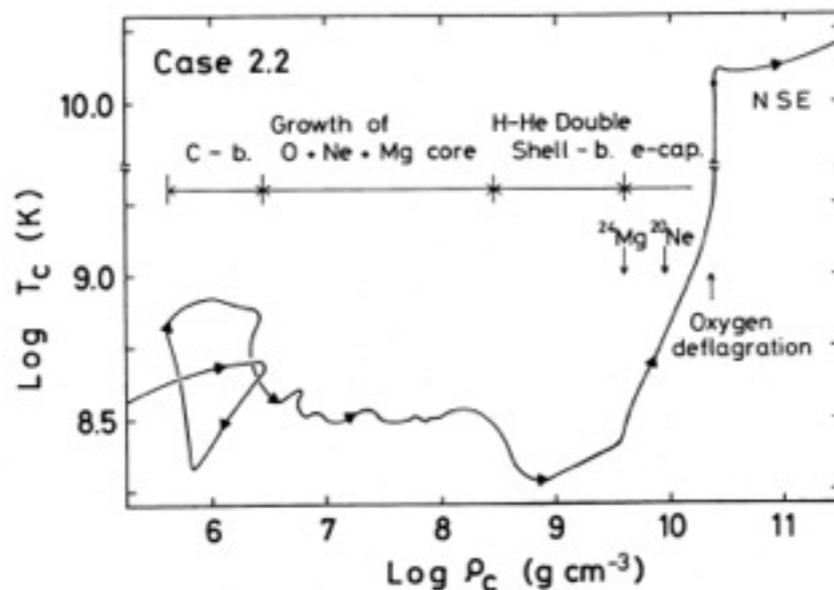


FIG. 4.—Evolutionary track in the central density and temperature diagram

Until core collapse

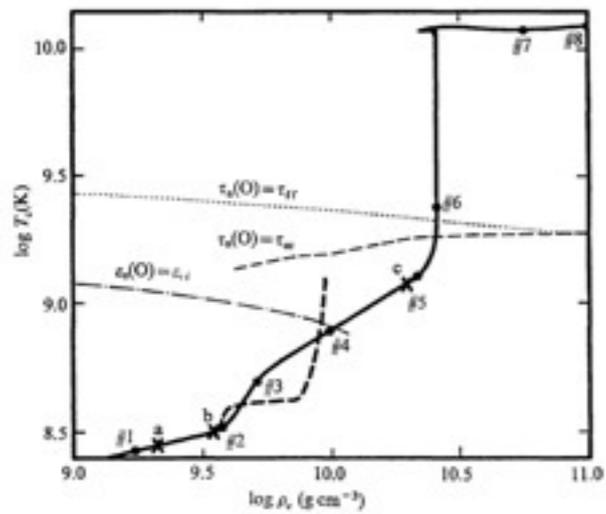


Fig. 5. Evolutionary change in the central temperature and the central density. See the text for stages a, b, and c. If the occurrence of the convective core were neglected, it would run as indicated by the thick dashed curve. Thin lines are concerned with the conditions for oxygen burning (see the text for detail).

(Miyaji et al. 1980)

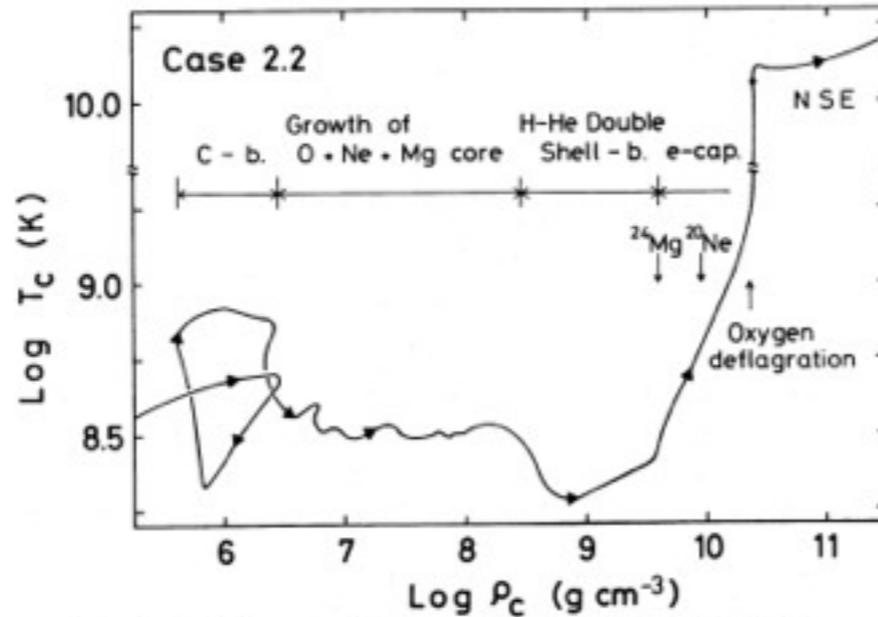
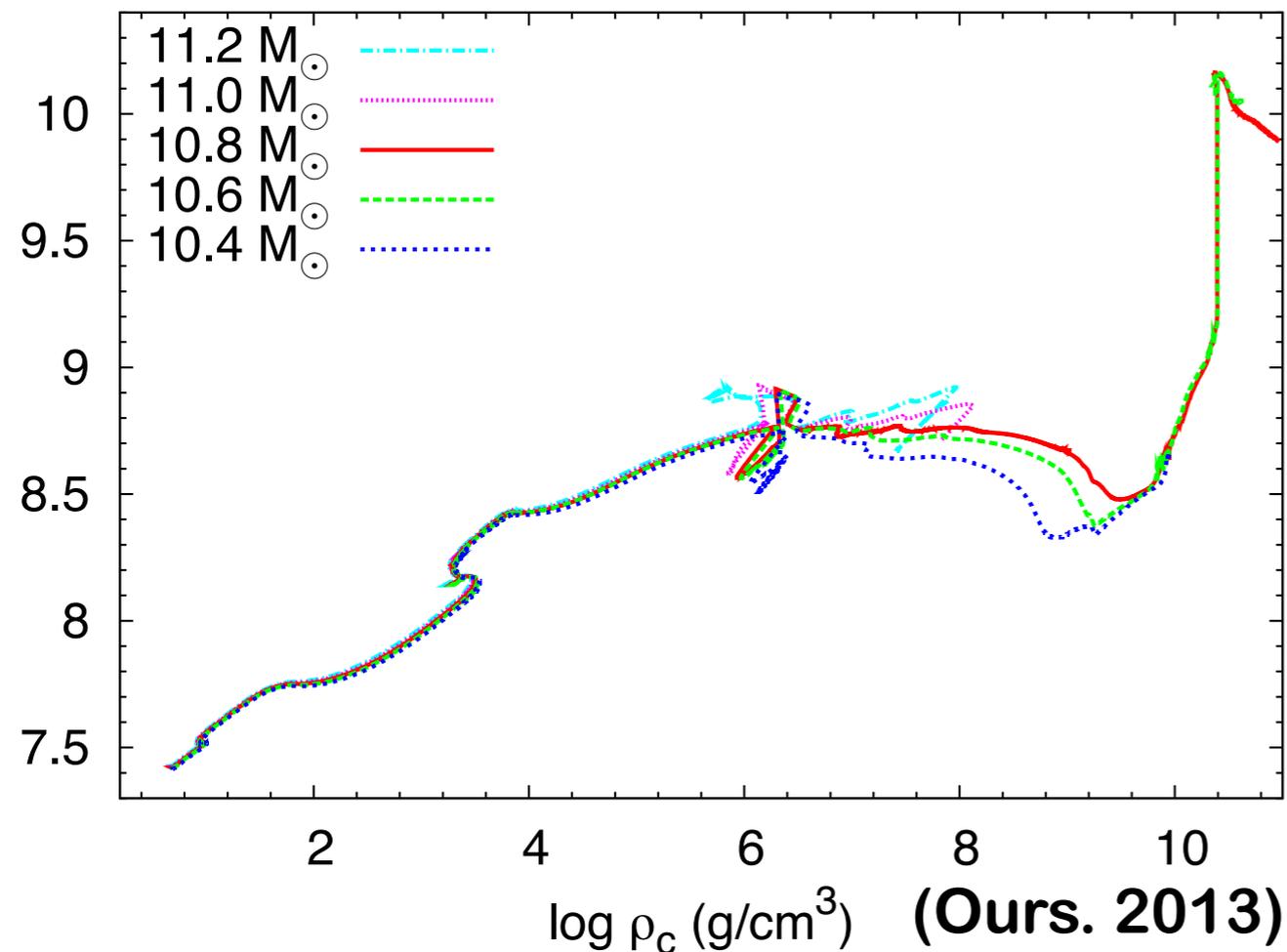


FIG. 4.—Evolutionary track in the central density and temperature diagram

(Nomoto 1987)

We have calculated the collapse of an ONe core after 26 years absence of other works.



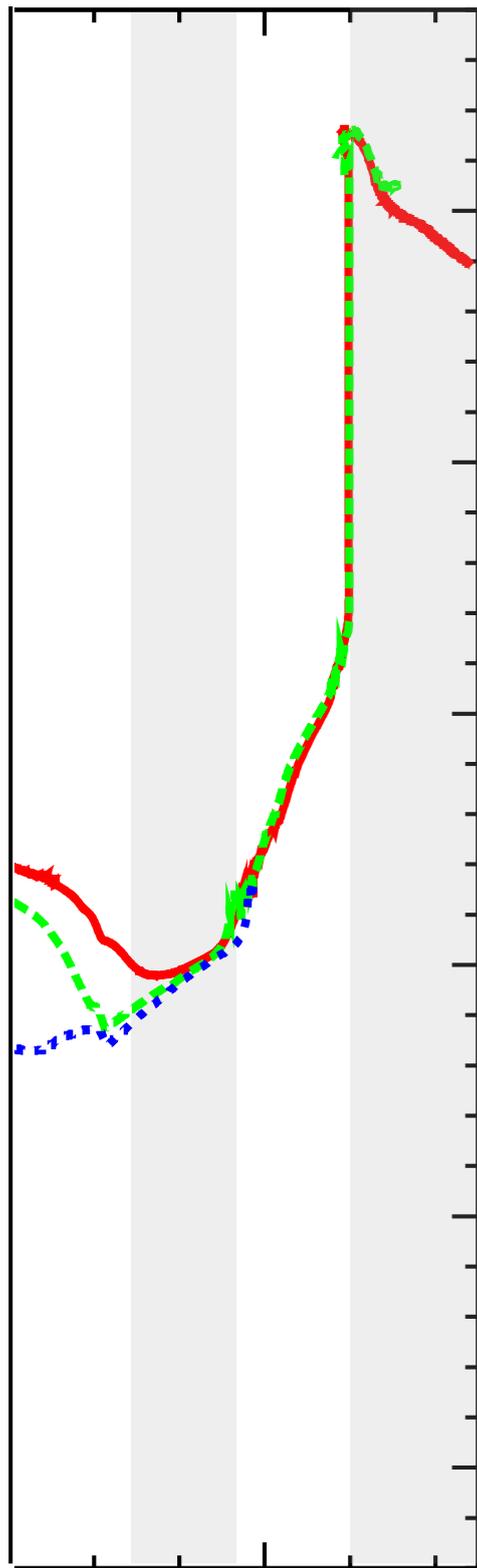
(Ours. 2013)

Our result shows generally good agreements with previous works.

Until core collapse

- How does the core end up as core collapse?
- What physics is important?

Until core collapse

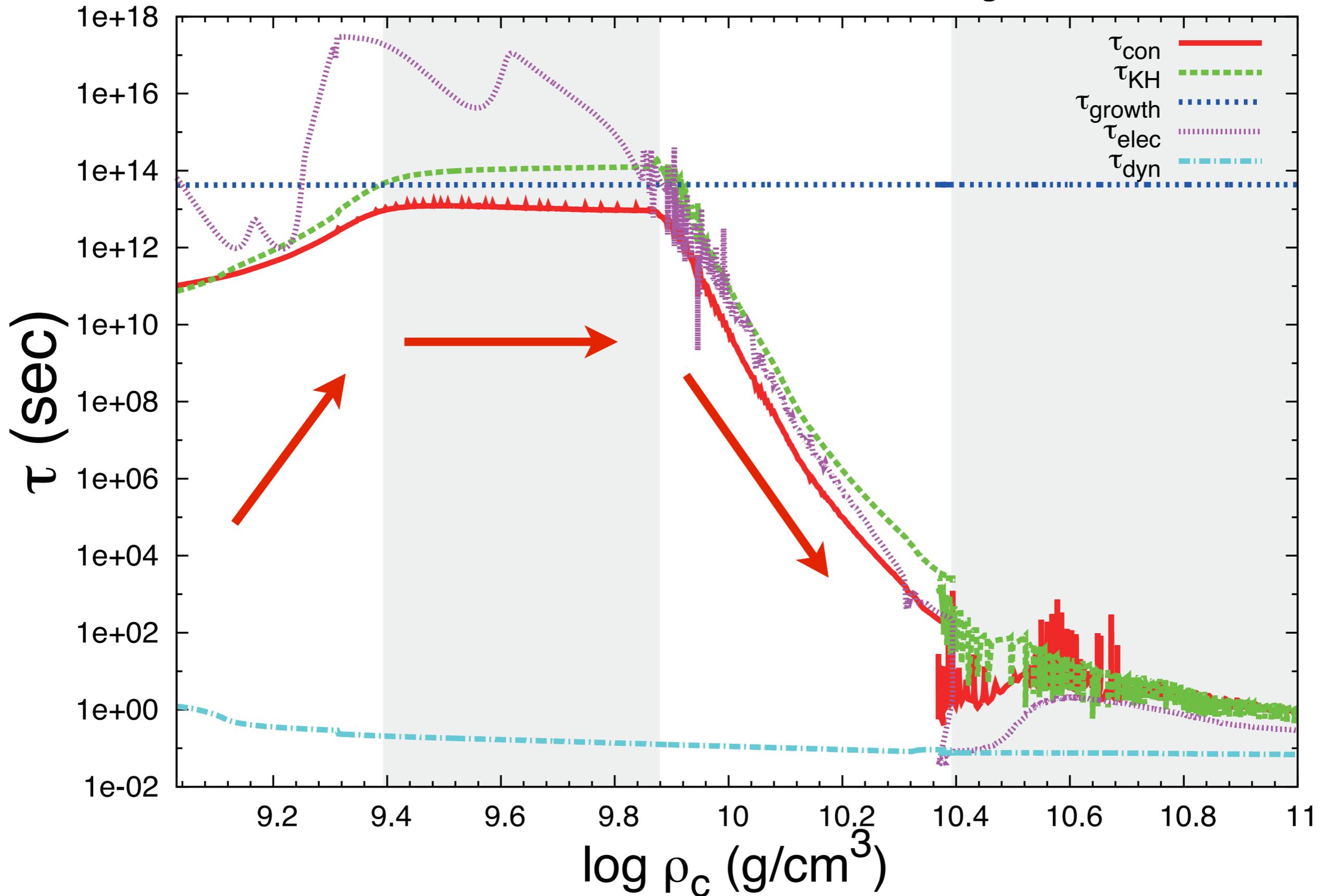


Core contraction can be divided into 4 sub-phases.

1. ν cooling phase
2. core growth phase
3. electron capture phase
4. deflagration phase

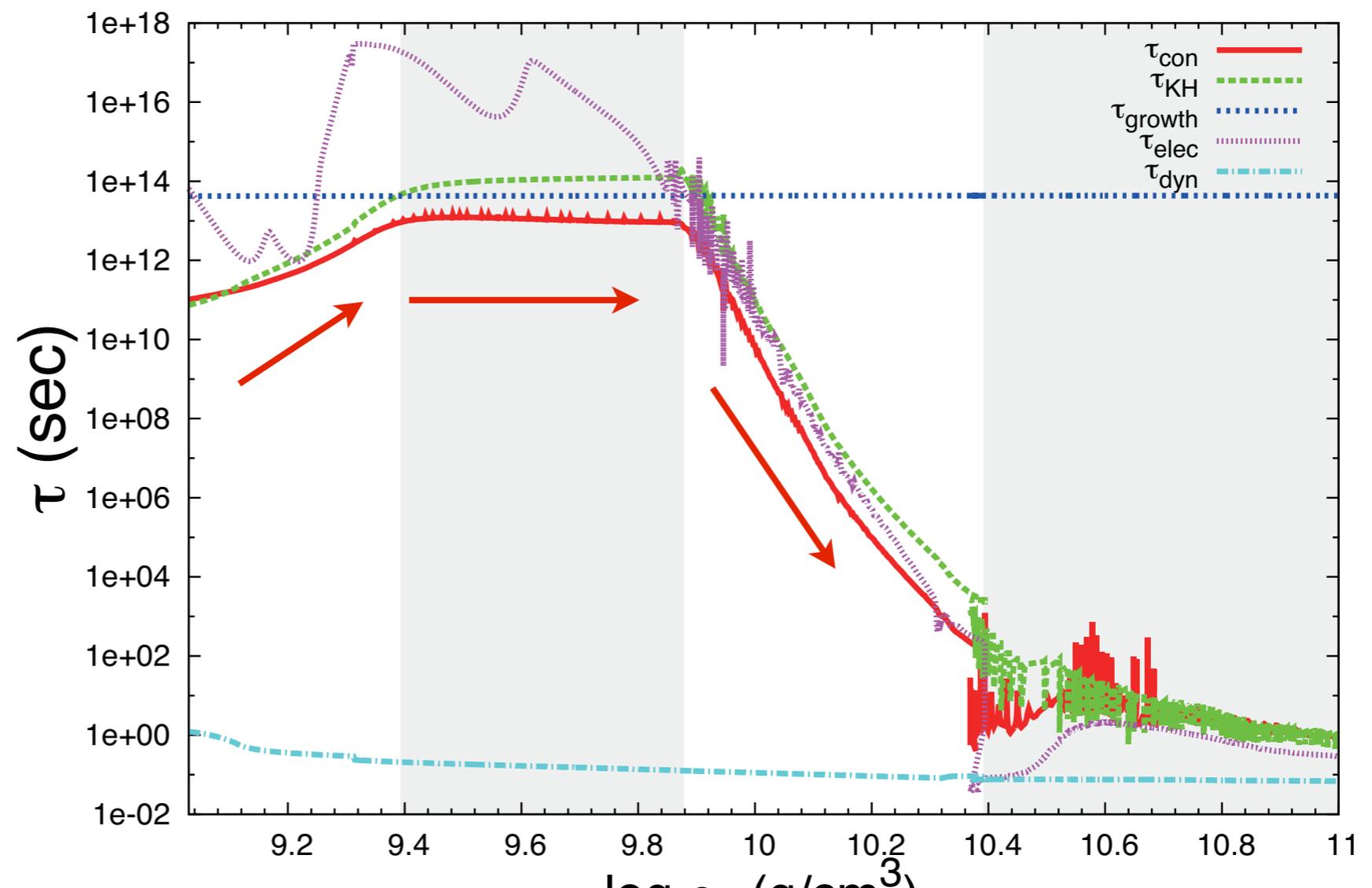
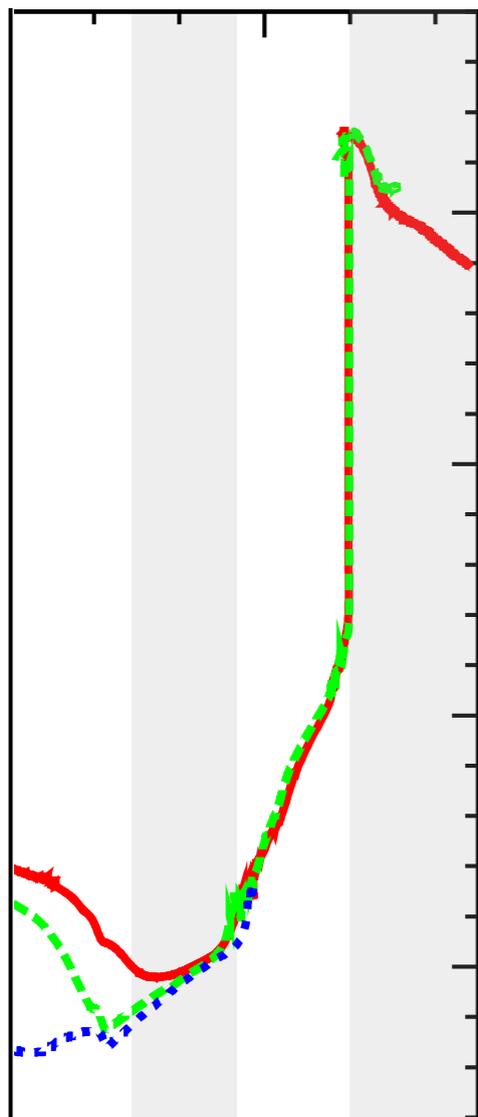
Until core collapse

diagram of
Evolutionary timescales

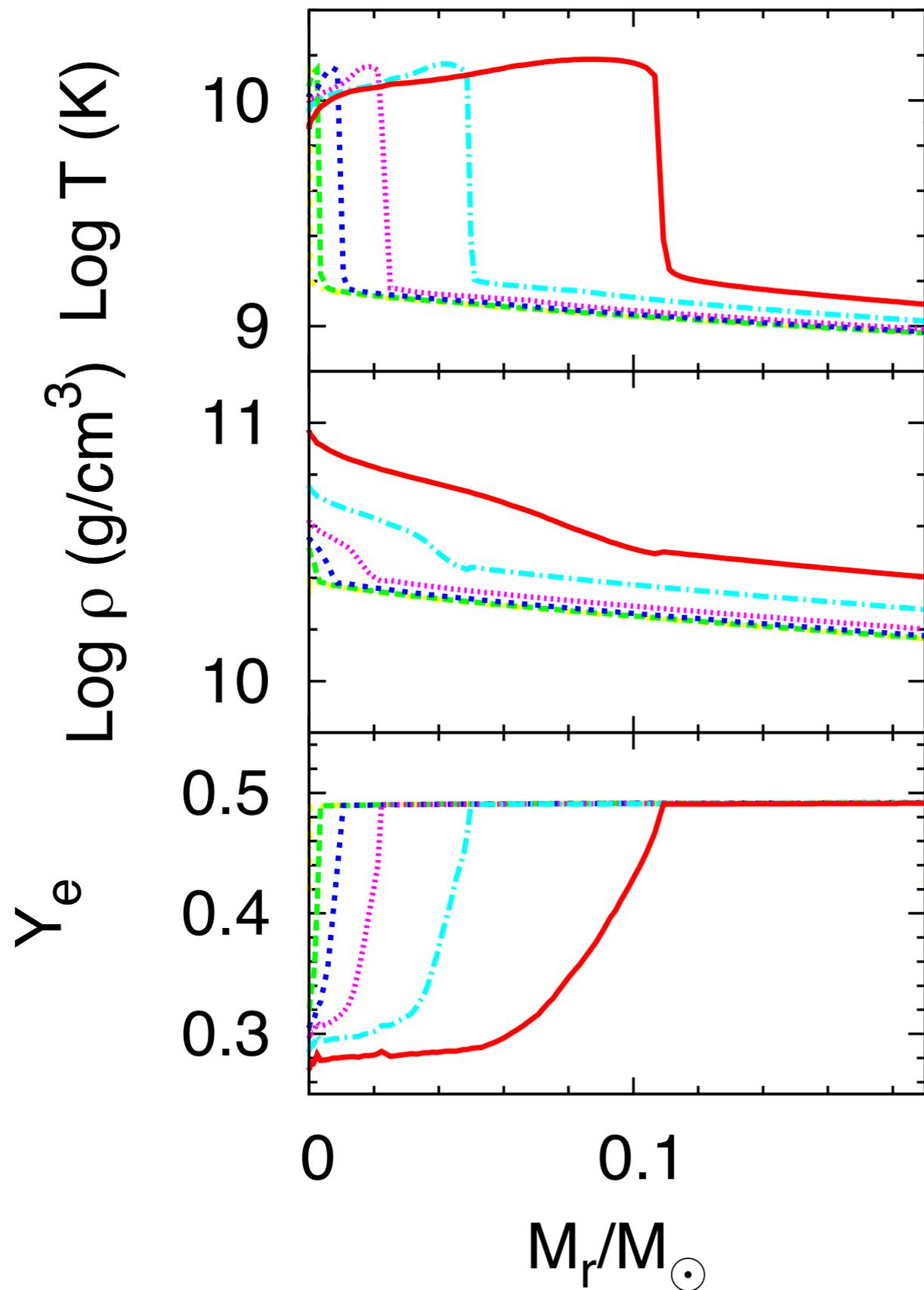


Until core collapse

- ν cooling phase (\nearrow) : self-regulating
- core growth phase (\rightarrow) : constant
- electron capture phase (\searrow) : self-enhancing
- deflagration phase : complicated



Until core collapse



- O+Ne burning causes deflagration in the degenerate core.
- convective energy transport is solved.
- **NSE** is achieved behind the deflagration front.
- in the NSE region, intense electron capture reduces Y_e .
- Y_e reduction cause core contraction.

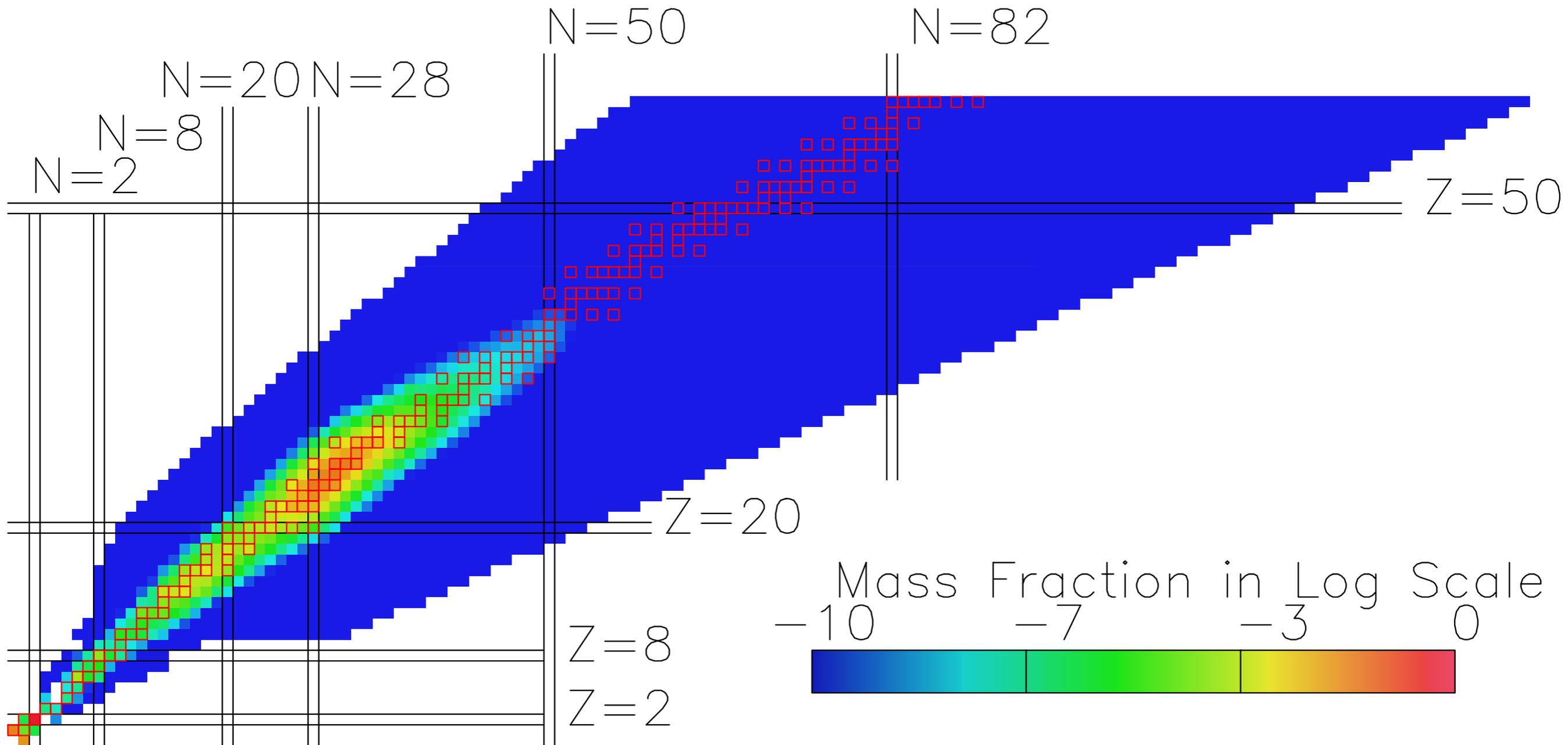
Nuclear Statistical Equilibrium

include 3091 species of nuclei

$T_g = 1.00E+01$; $\rho = 1.00E+09 \text{ g/cm}^3$; $Y_e = 4.90E-01$

$Y_p = 1.90E-02$; $Y_n = 1.28E-02$; $Y_\alpha = 1.65E-01$

$Y_h = 5.84E-03$; $Y_i = 2.02E-01$



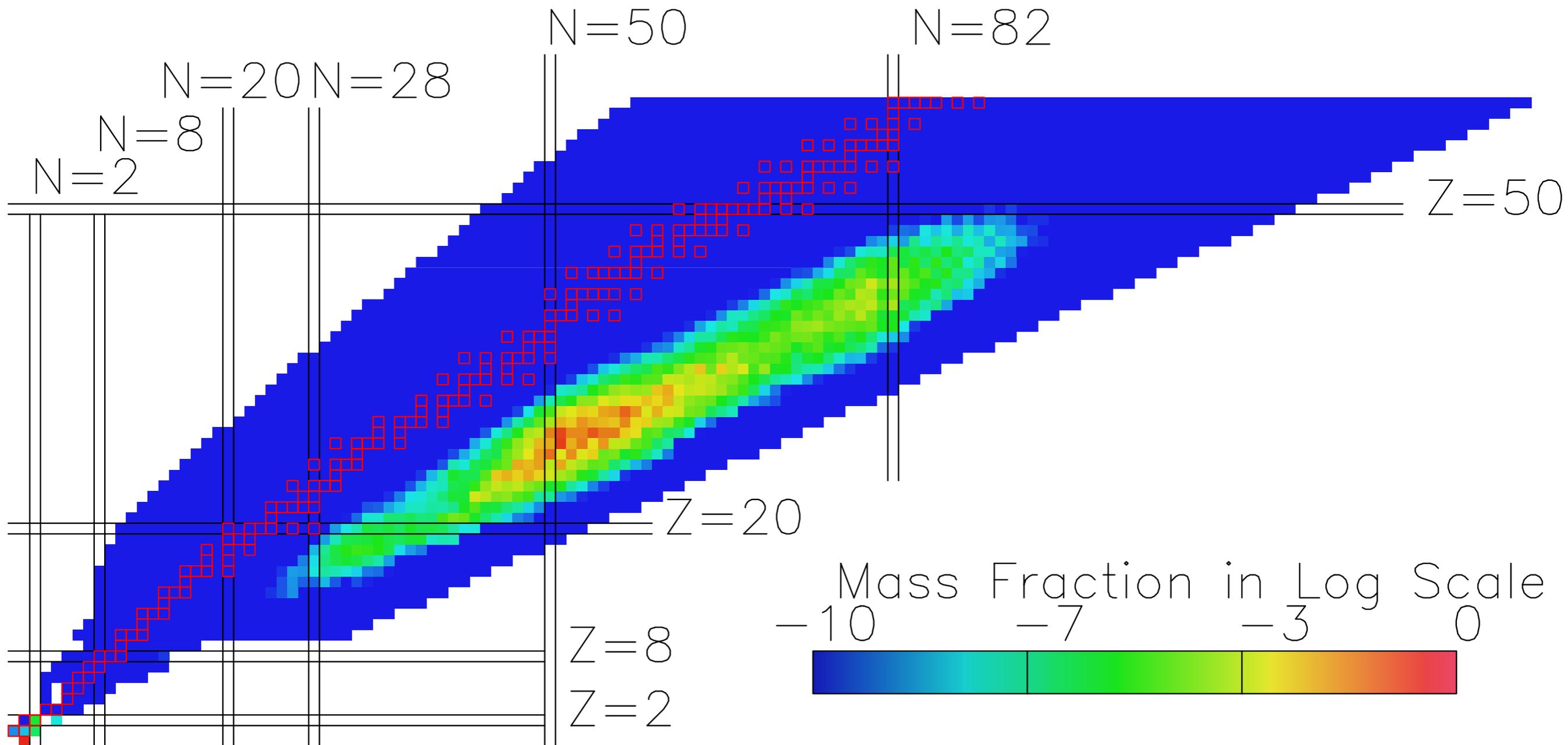
Nuclear Statistical Equilibrium

include 3091 species of nuclei

$T_g = 1.00E+01$; $\rho = 1.00E+11 \text{ g/cm}^3$; $Y_e = 2.70E-07$

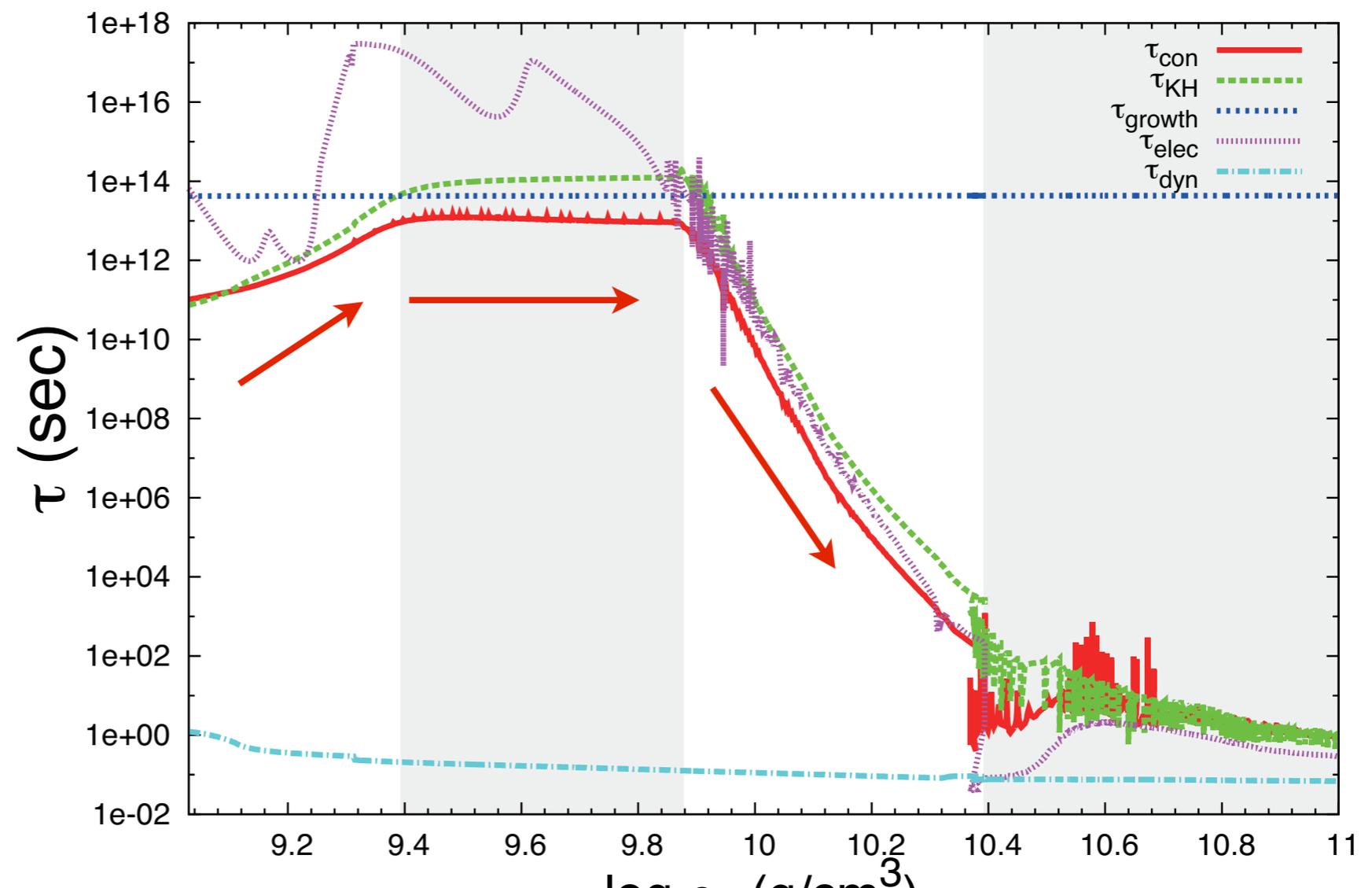
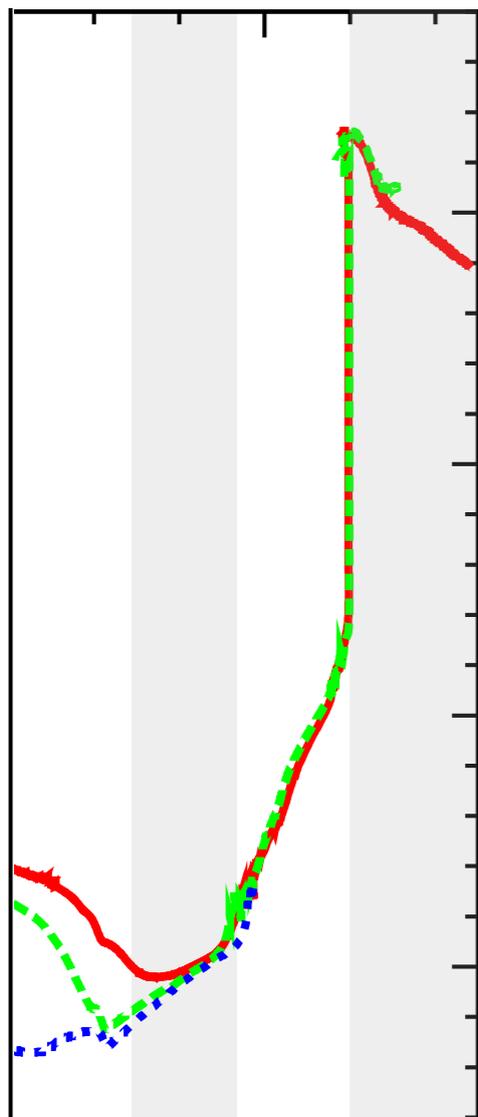
$Y_p = 2.13E-09$; $Y_n = 2.34E-01$; $Y_\alpha = 7.11E-07$

$Y_h = 9.36E-03$; $Y_i = 2.44E-01$



Until core collapse

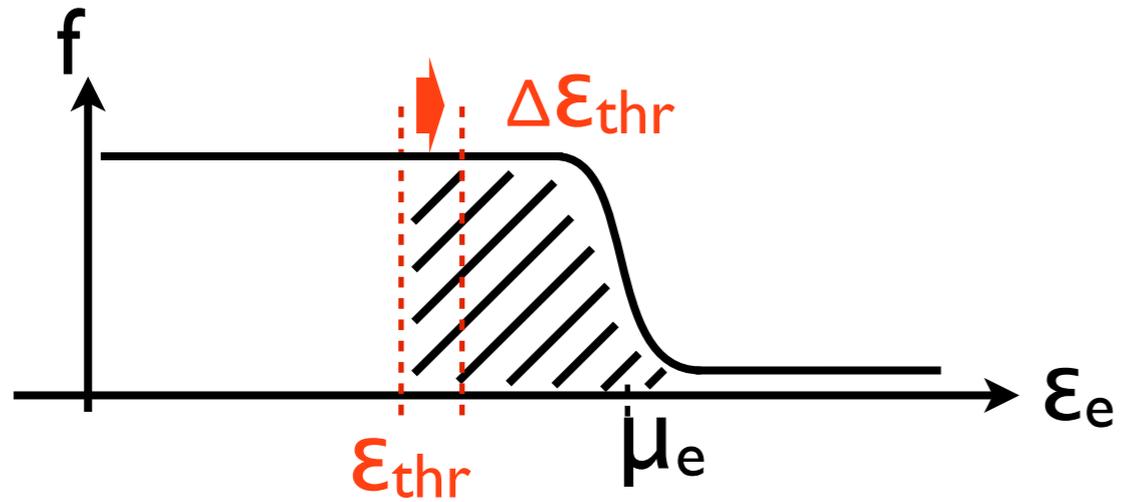
- deflagration phase
 - **propagation** of the burning front
 - electron capture reaction in the NSE region
 - reduction of **Ye**
 - reduction of **energy**



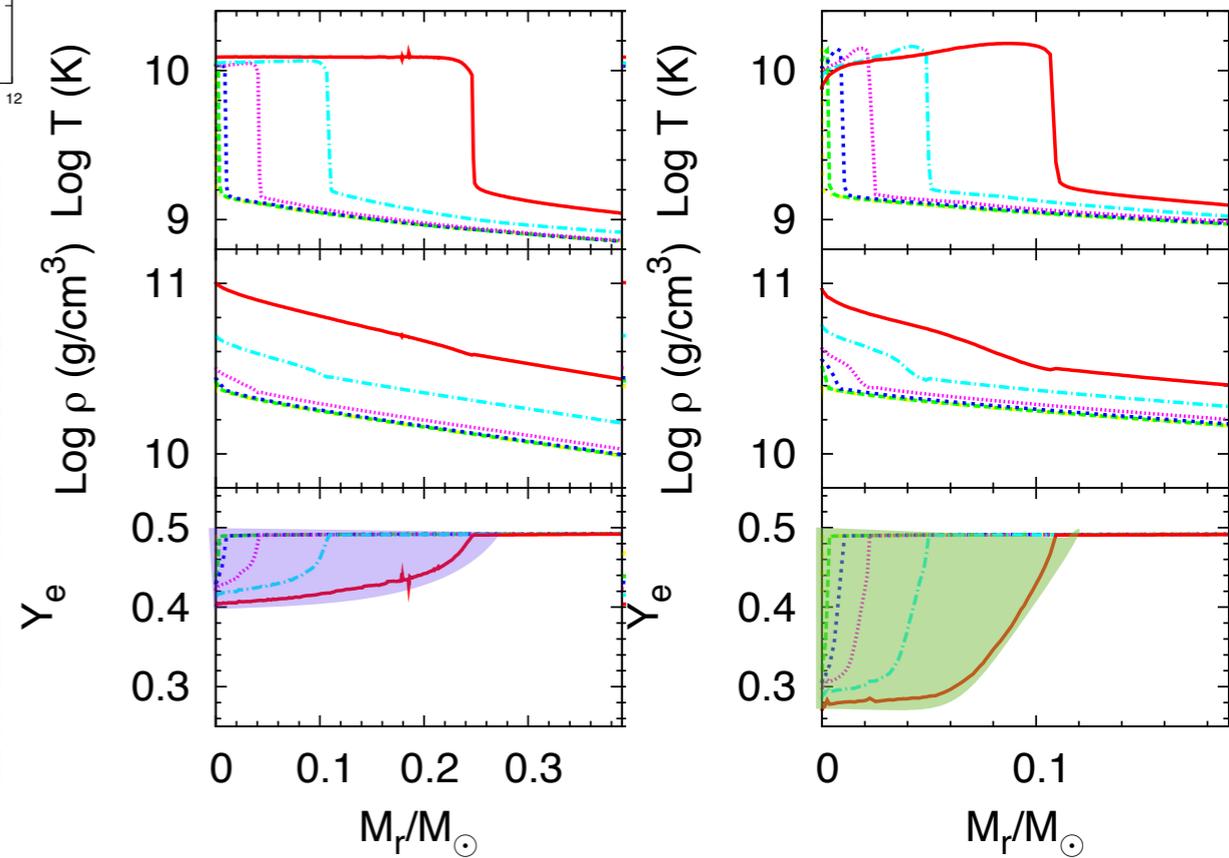
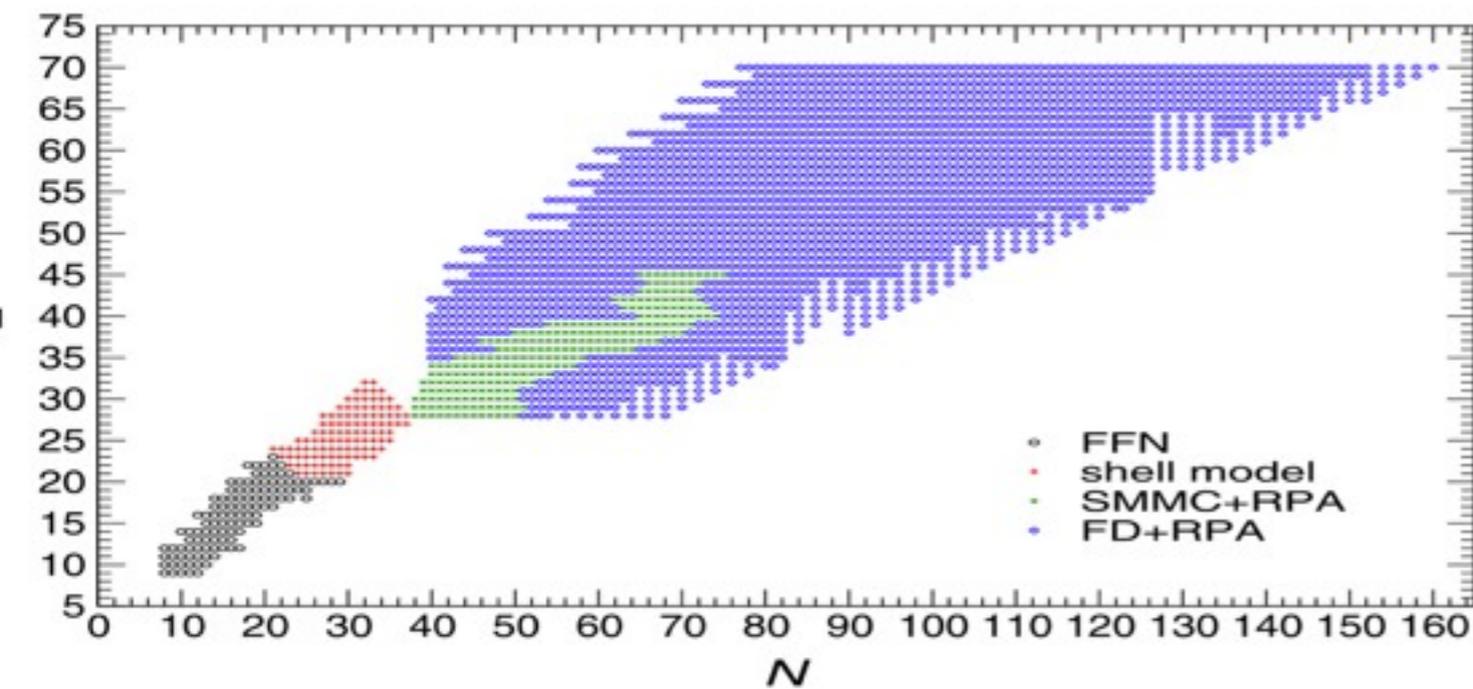
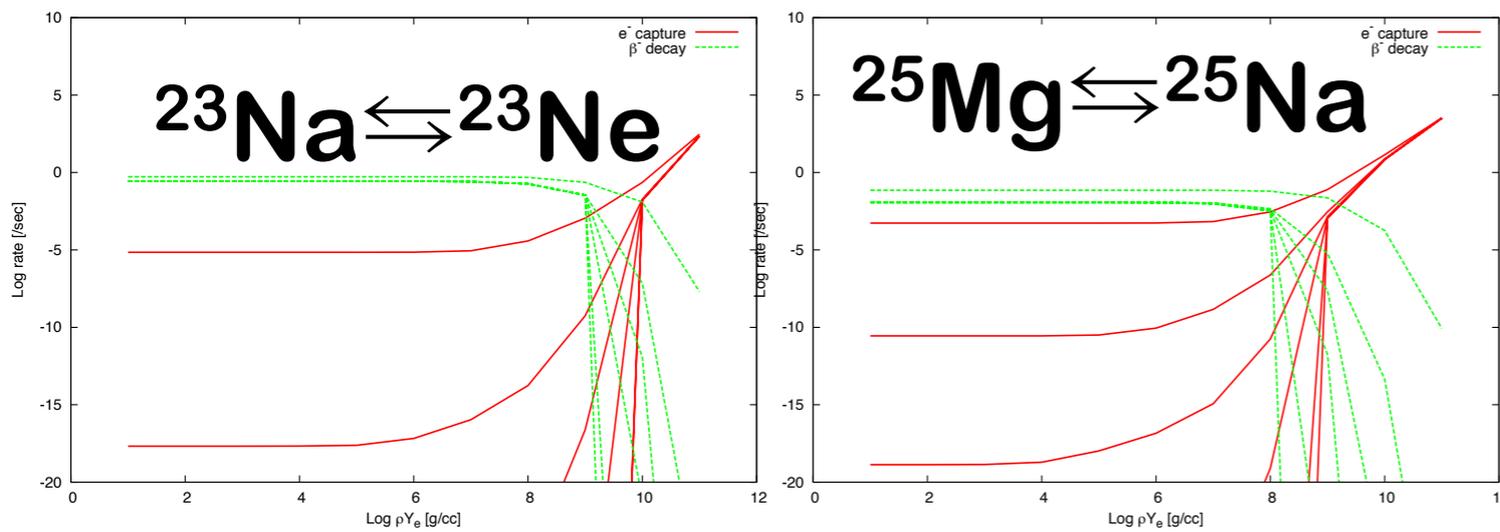
Until core collapse

Short summary 2

- **Contraction of an ONe core has 4 sub-phases.**
 - self-regulating ν -cooling phase
 - core mass growth phase
 - self-enhancing electron capture phase
 - deflagration phase
- **Deflagration phase is much complicated.**
 - propagation of the flame front
 - electron capture in the NSE region
- **Core collapse is finally induced by electron capture in the NSE region.**



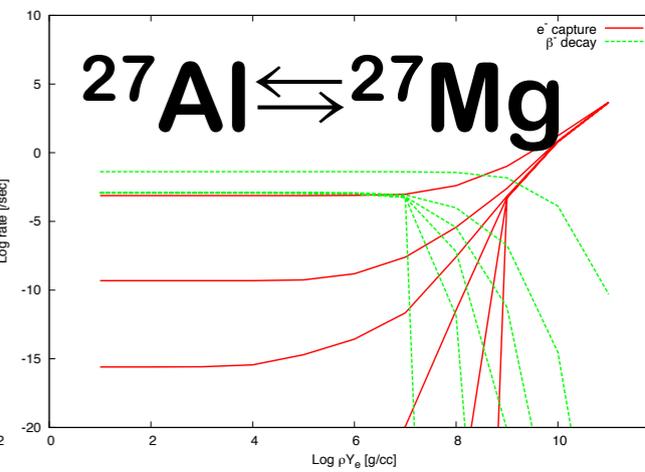
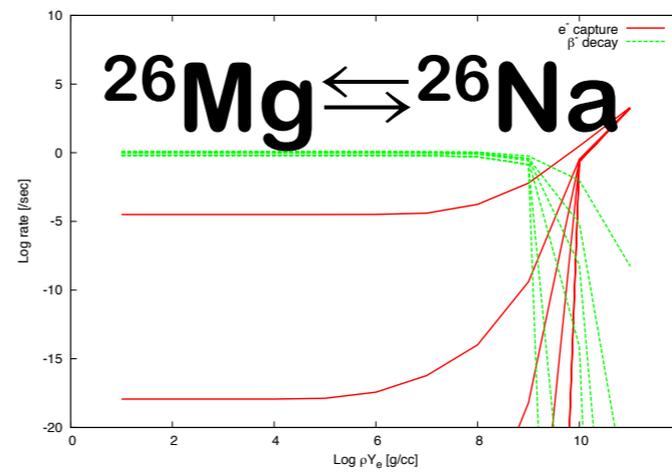
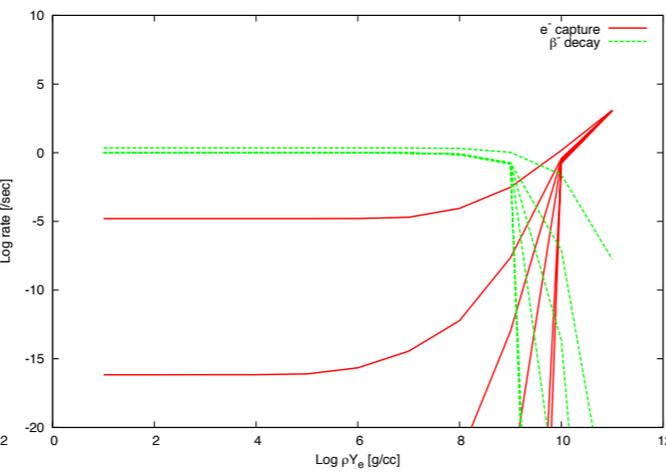
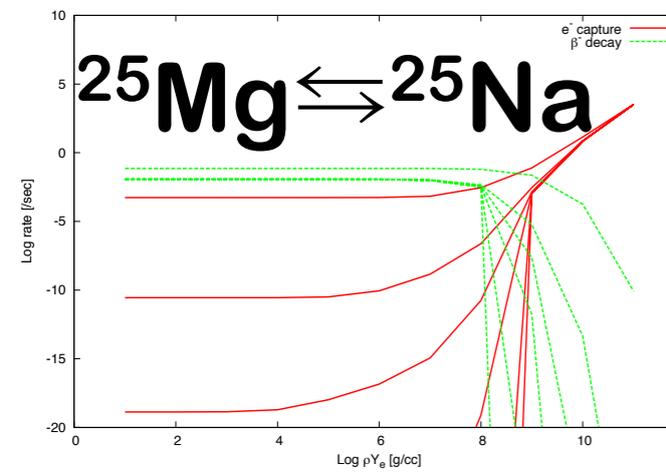
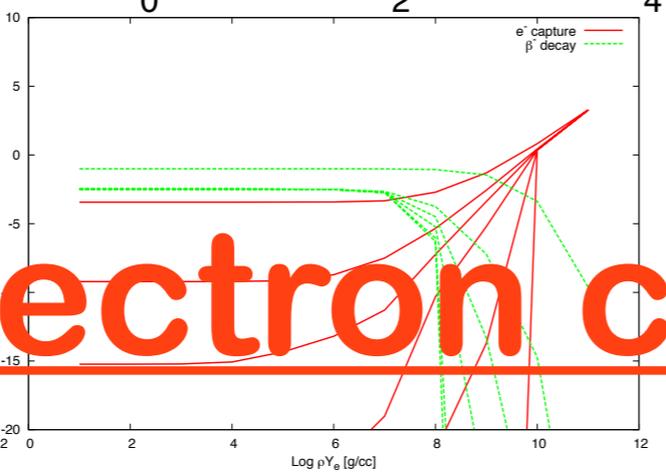
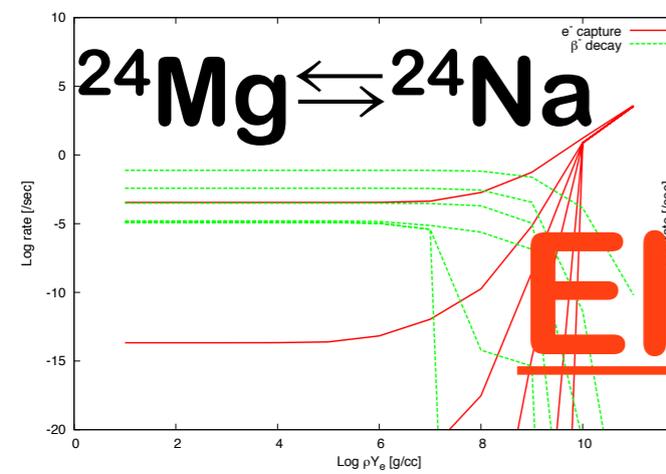
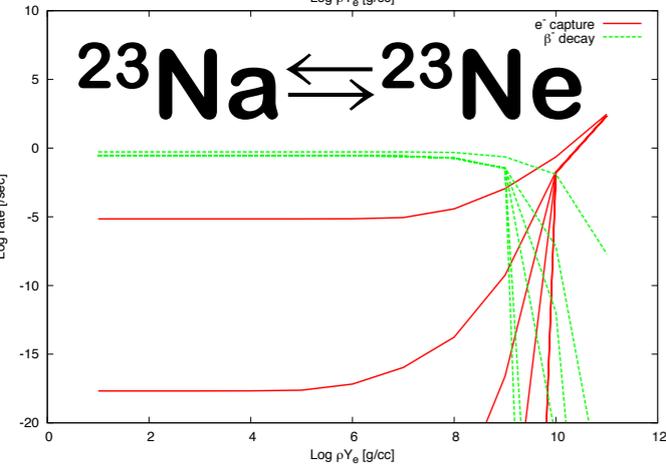
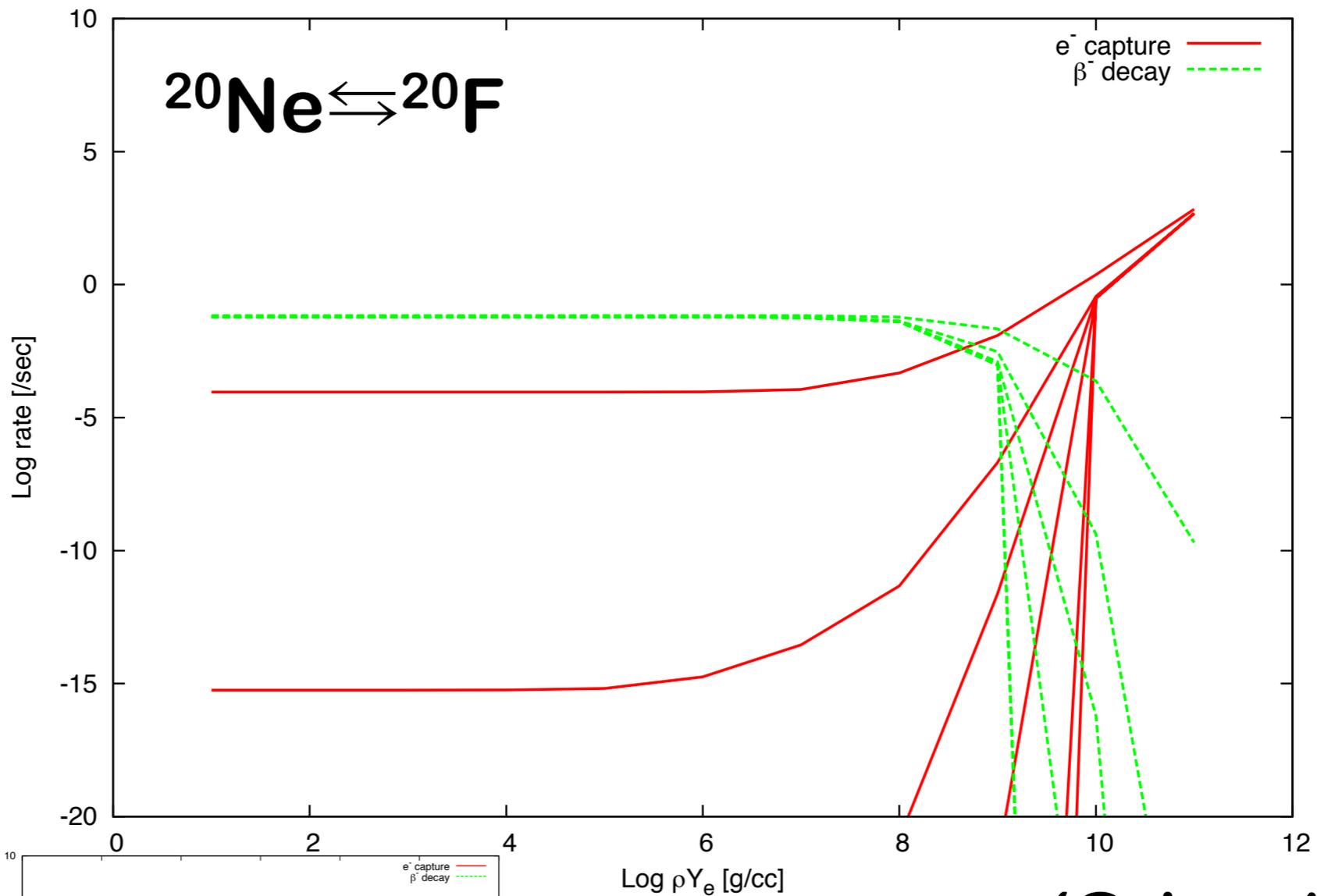
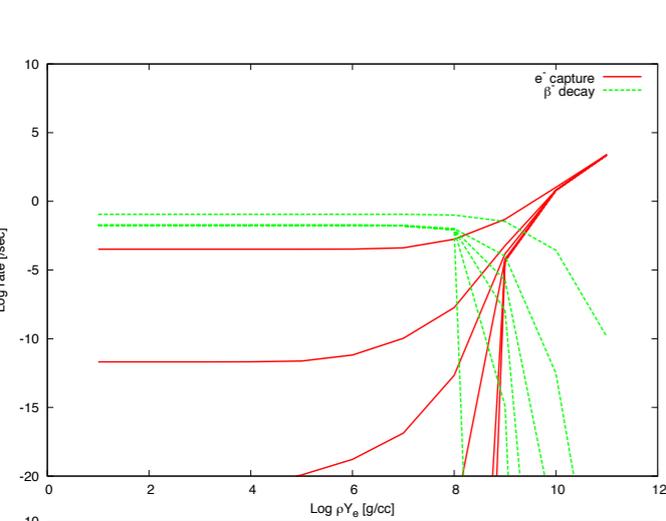
New physics



New physics

Importance of improved input physics is investigated.

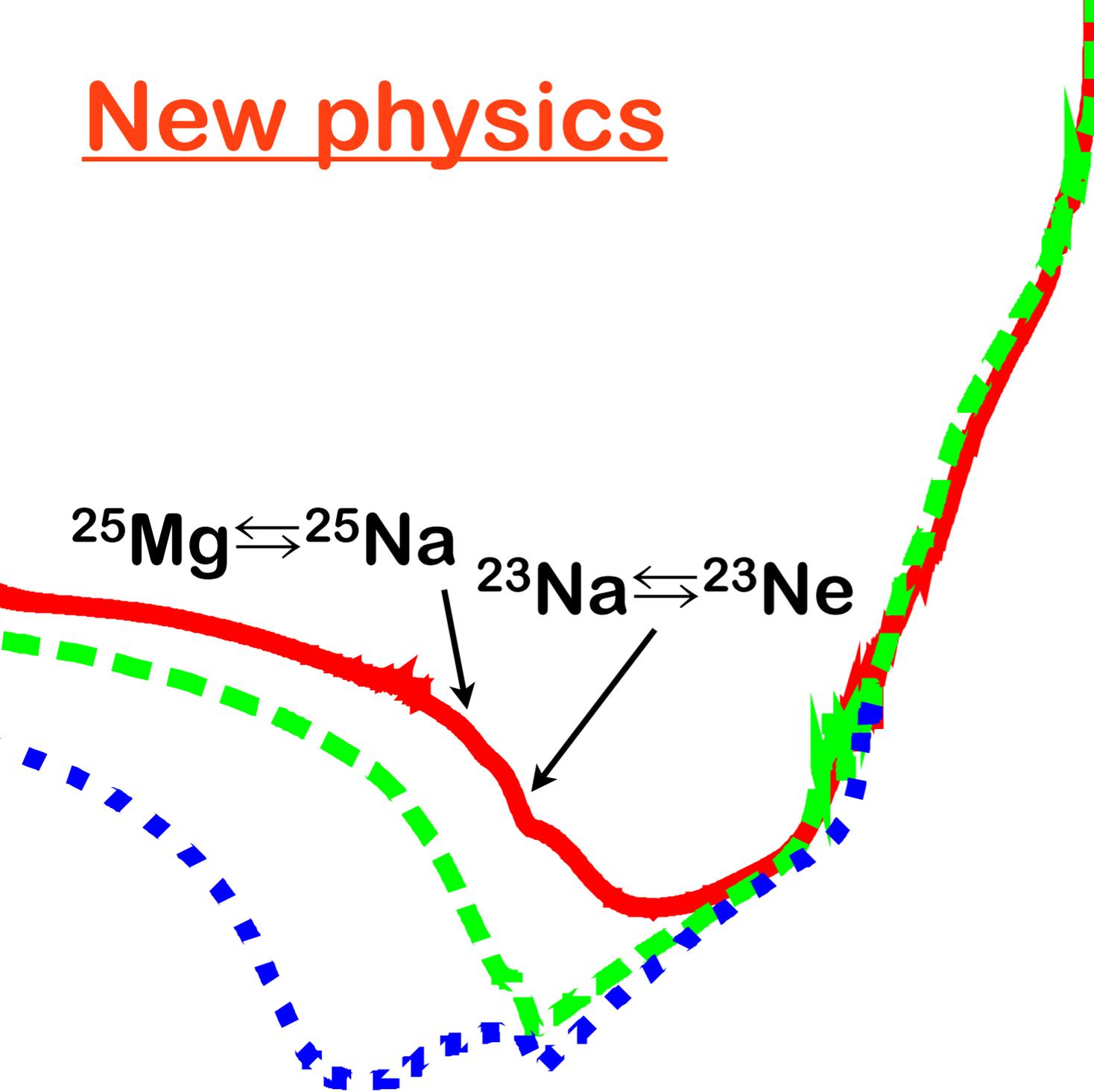
- electron capture by odd proton number isotopes
- the Coulomb correction on the electron capture rates
- electron capture by neutron-rich isotopes in NSE compositions



(Oda+ 1994)

Electron capture reactions

New physics



The effect can be seen.
 But less important.
 - small mass fraction
 - endothermic reaction

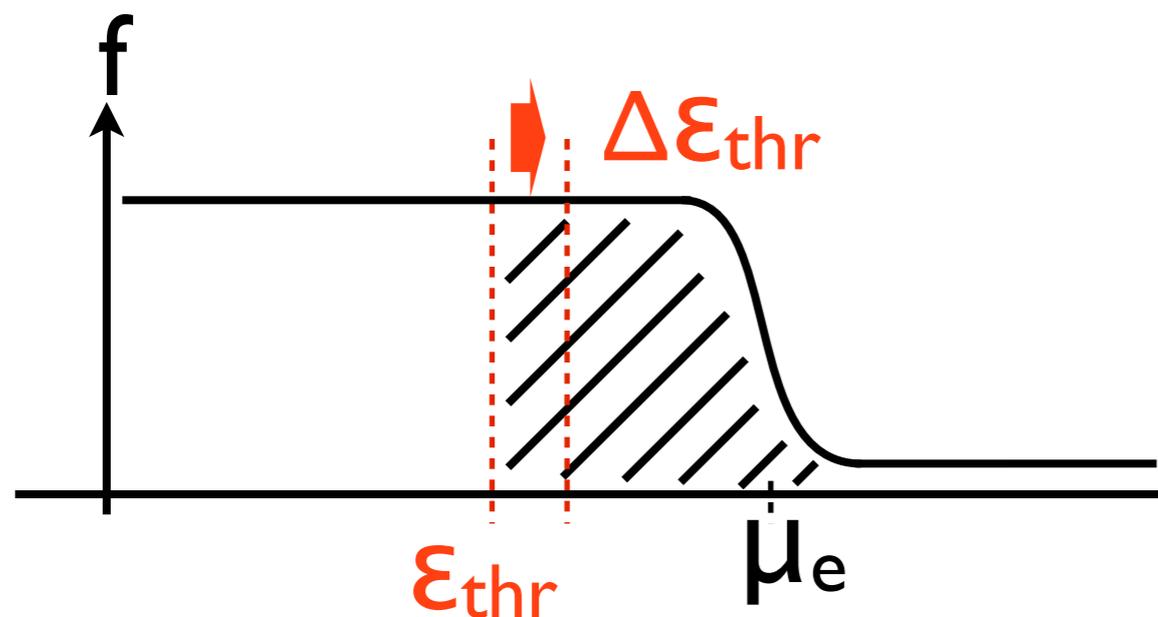
Table 2

Central Mass Fractions of the 14 Most Abundant Isotopes at the End of the Core Carbon Burning Phase

^{16}O	^{20}Ne	^{24}Mg	^{23}Na	^{25}Mg	^{26}Mg	^{27}Al
4.783×10^{-1}	4.074×10^{-1}	4.255×10^{-2}	3.217×10^{-2}	1.451×10^{-2}	7.952×10^{-3}	7.245×10^{-3}
^{22}Ne	^{28}Si	^{12}C	^{21}Ne	^{30}Si	^{29}Si	^{32}S
2.836×10^{-3}	2.330×10^{-3}	1.277×10^{-3}	8.384×10^{-4}	3.621×10^{-4}	3.362×10^{-4}	1.463×10^{-4}

The Coulomb correction for electron capture rates

$$\lambda_{ec} = \frac{1}{\pi^2 h^2} \sum \int_{\epsilon_{thr} + \Delta\epsilon_{thr}}^{\infty} p_e^2 \sigma_{ec} f(\epsilon_e - \mu_e) d\epsilon_e$$



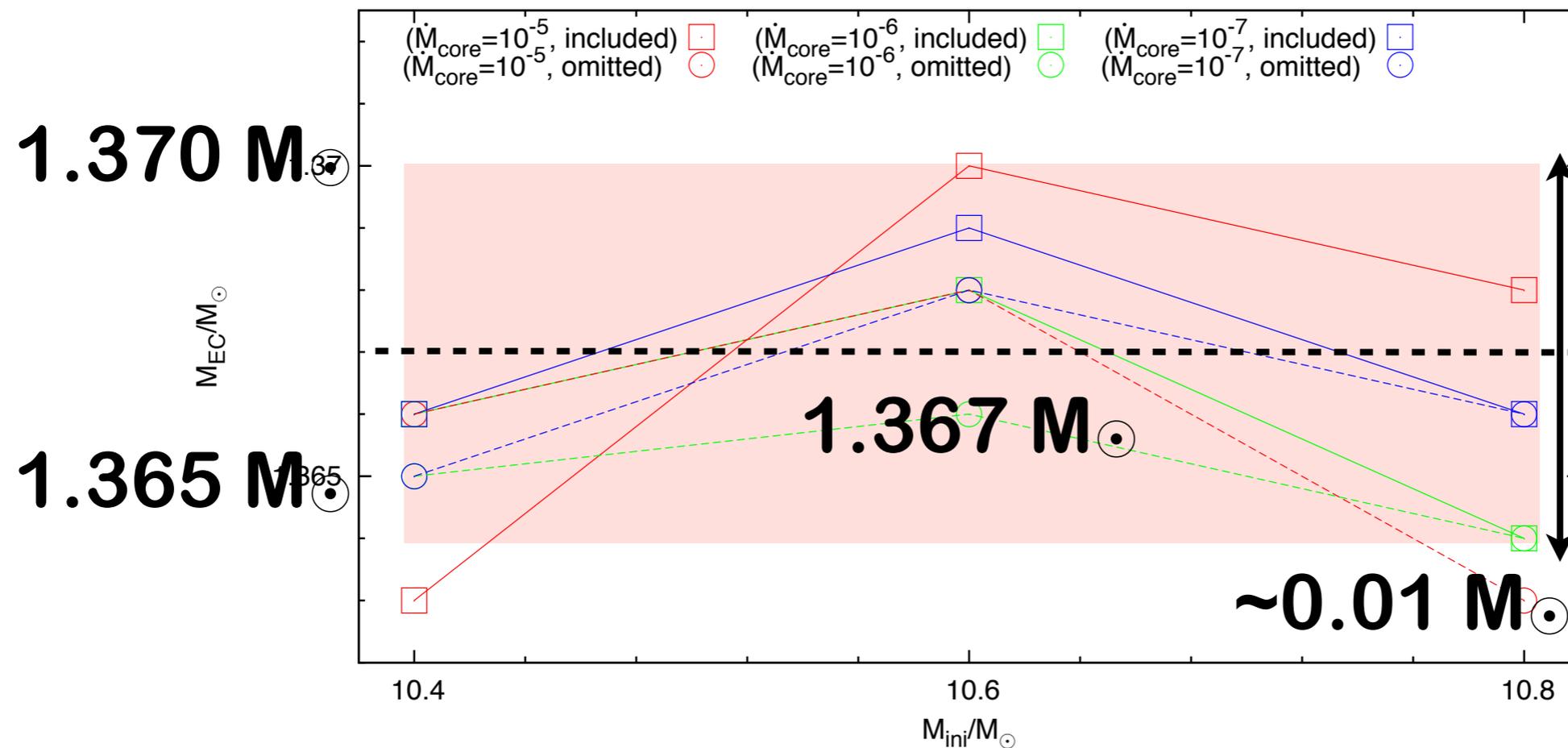
$$\Delta\epsilon_{thr} = \mu_{ion}(Z-1) - \mu_{ion}(Z)$$

$$\mu_{ion}(Z) = -kT \left(\frac{Z}{\bar{Z}} \right)$$

$$\left\{ \Gamma_Z \left[0.9 + c_1 \left(\frac{\bar{Z}}{Z} \right)^{1/3} + c_2 \left(\frac{\bar{Z}}{Z} \right)^{2/3} \right] + [d_0 + d_1 \left(\frac{\bar{Z}}{Z} \right)^{1/3}] \right\}$$

(Couch & Loumos 1974; DeWitt et al. 1973)

New physics

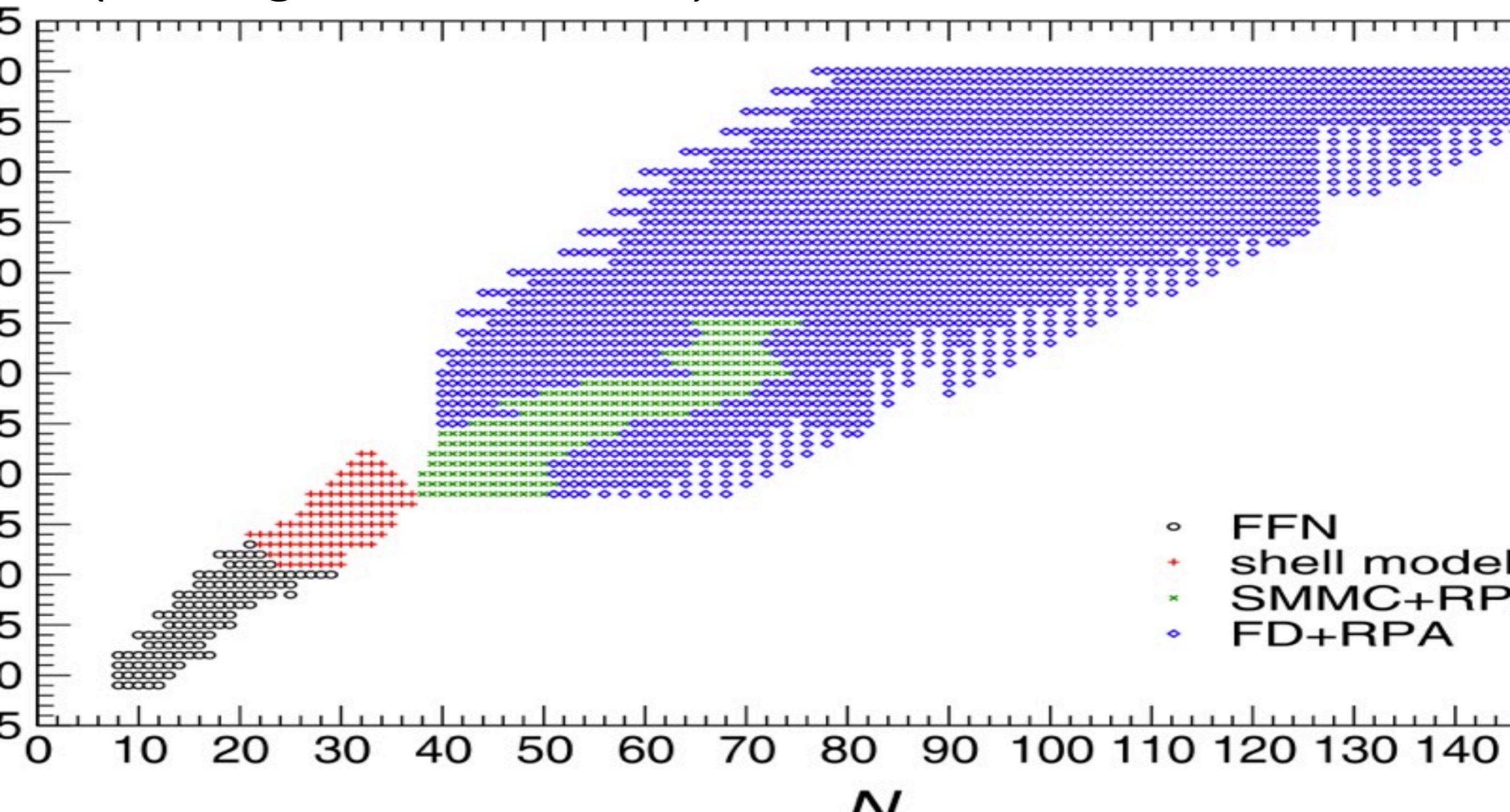


- ONe core evolution is independent from parameters ($3 M_{\text{ZAMS}}$ & $3 \dot{M}_{\text{core}}$)

- **Extremely degenerate ONe core** should have a unique structure

New physics

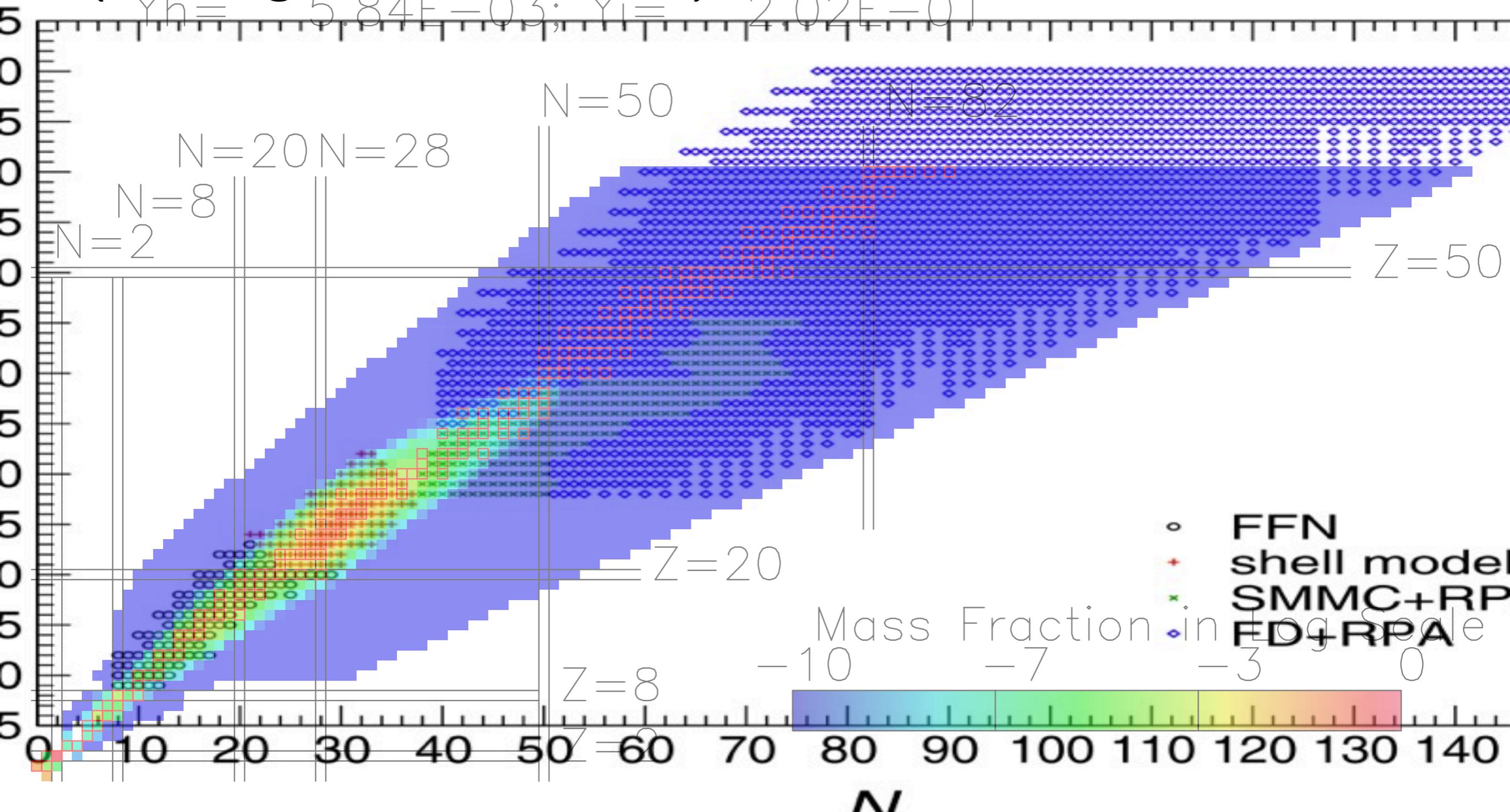
- electron capture by neutron-rich isotopes in NSE compositions
(Juodagalvis et al. 2010)



New physics

- electron capture by neutron-rich isotopes in NSE compositions
(Juodagalvis et al; 2010)

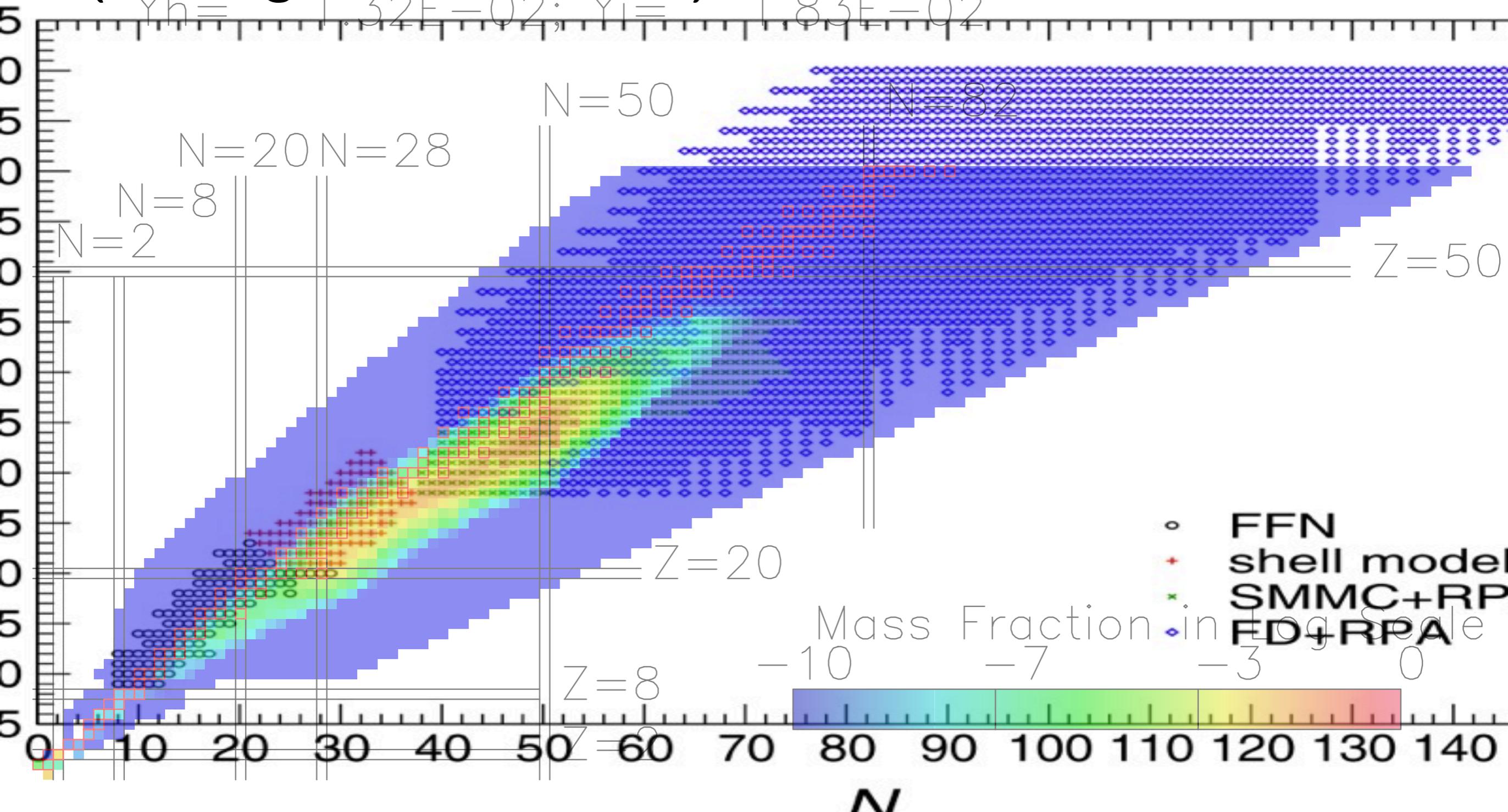
$T_g = 1.00E+01$ g/cm³; $Y_e = 4.90E-01$
 $Y_p = 1.00E-02$; $Y_n = 1.28E-02$; $Y_\alpha = 1.65E-01$
 $Y_h = 5.84E-03$; $Y_i = 2.02E-01$



New physics

- electron capture by neutron-rich isotopes in NSE compositions
(Juodagalvis et al. 2010)

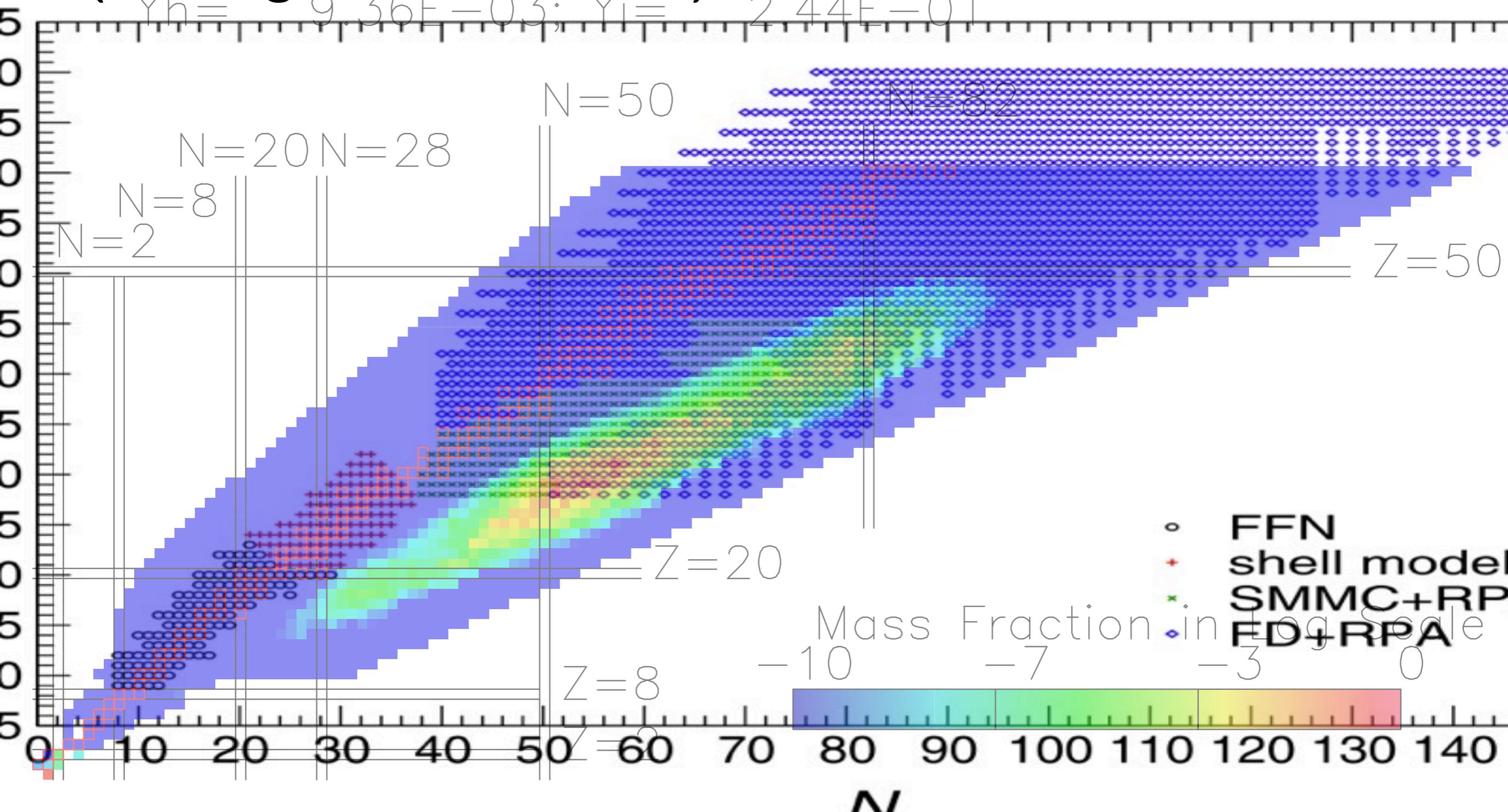
$T_9 = 1.00E+01$; $\rho = 1.00E+10 \text{ g/cm}^3$; $Y_e = 4.10E-01$
 $Y_p = 8.50E-06$; $Y_r = 4.52E-03$; $Y_\alpha = 5.26E-04$
 $Y_h = 1.32E-02$; $Y_i = 1.83E-02$

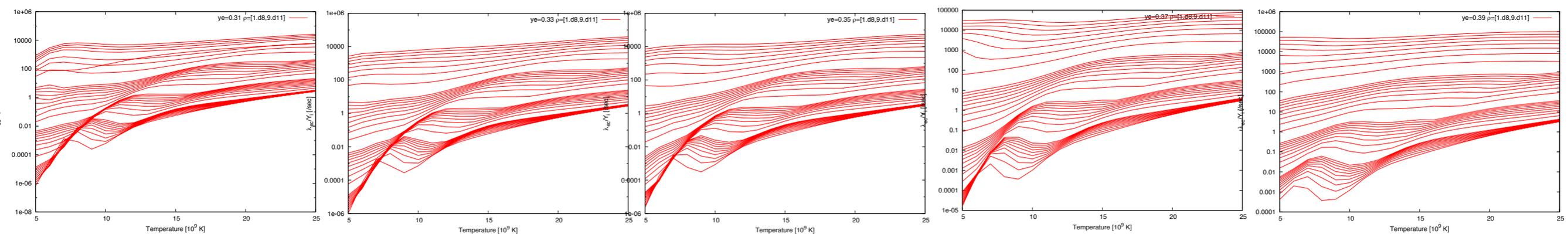


New physics

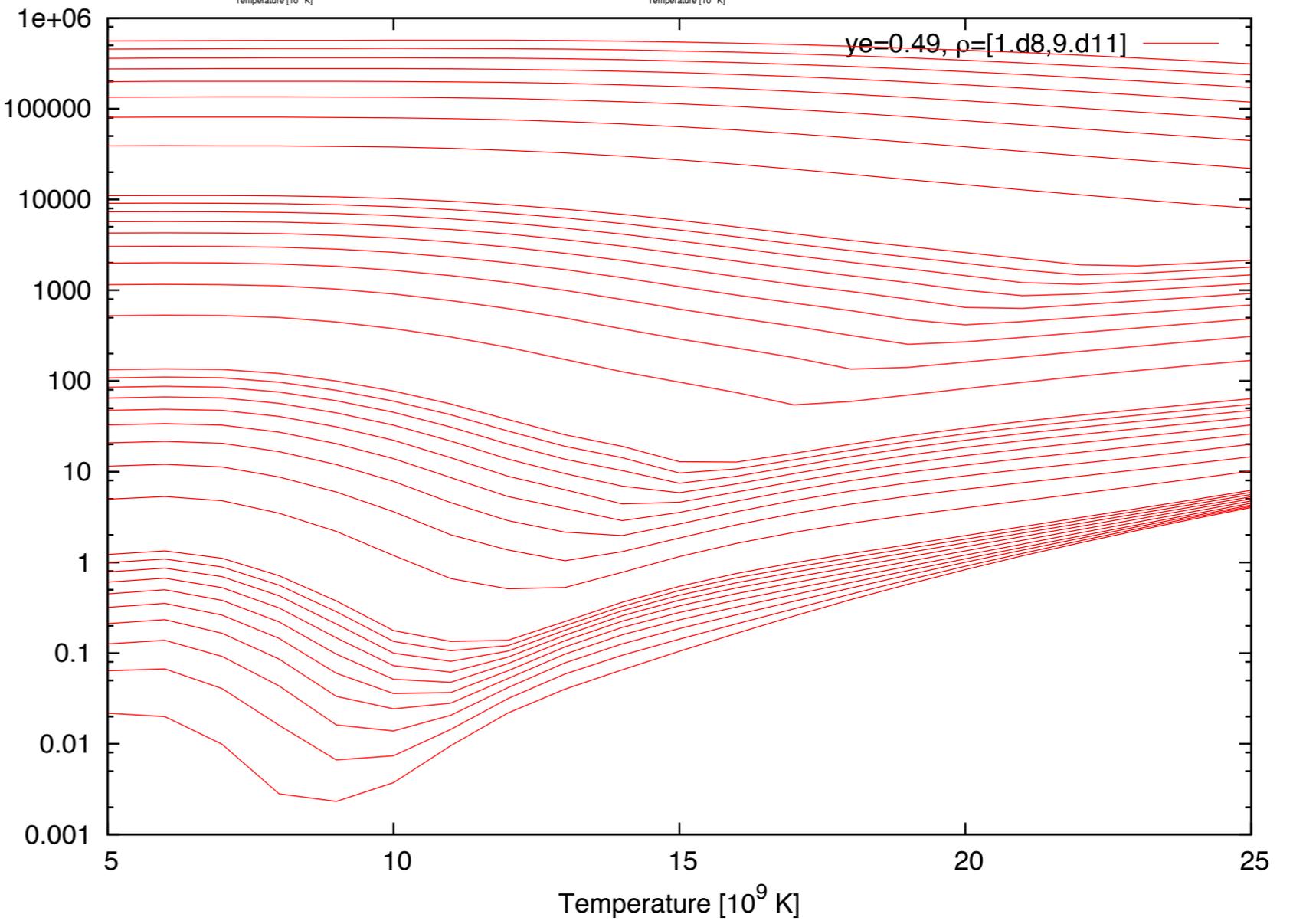
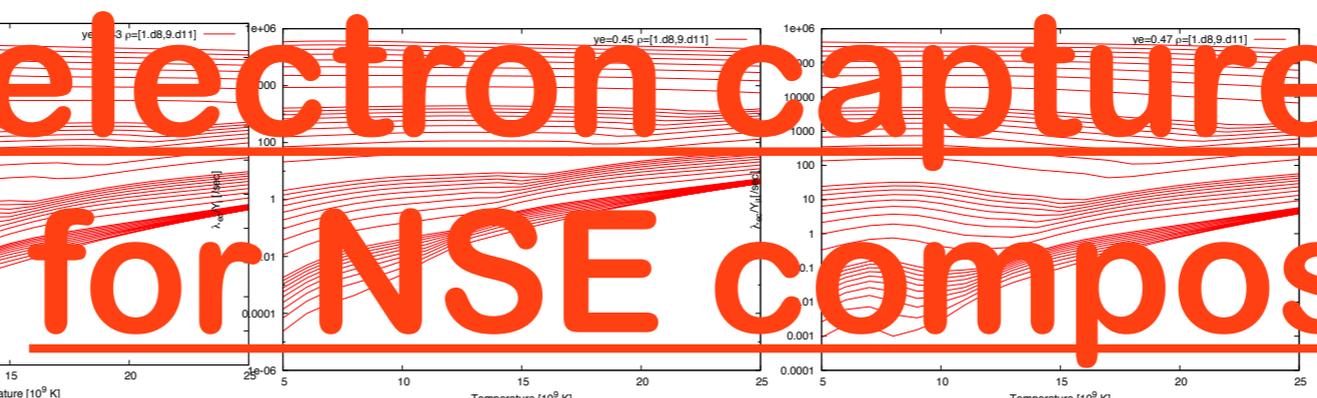
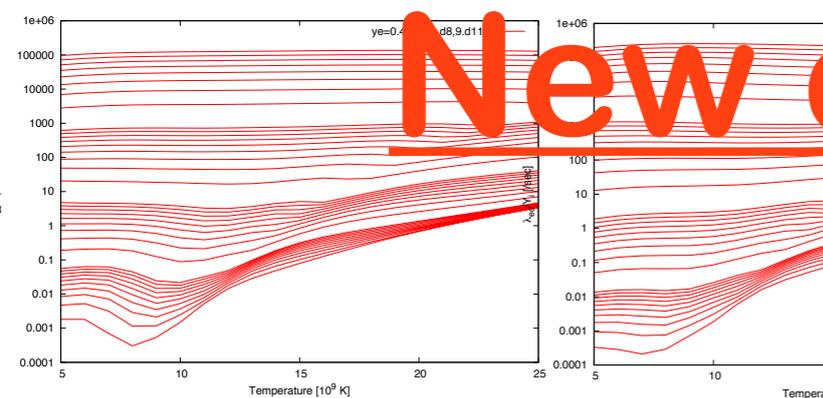
- electron capture by neutron-rich isotopes in NSE compositions
(Juodagalvis et al. 2010)

$\rho = 1.00E+11 \text{ g/cm}^3$; $Y_e = 2.70E-01$
 $Y_p = 2.34E-01$; $Y_\alpha = 7.11E-07$
 $Y_h = 9.36E-03$; $Y_i = 2.44E-01$





New electron capture rates for NSE compositions



- electron capture by
neutron-rich isotopes

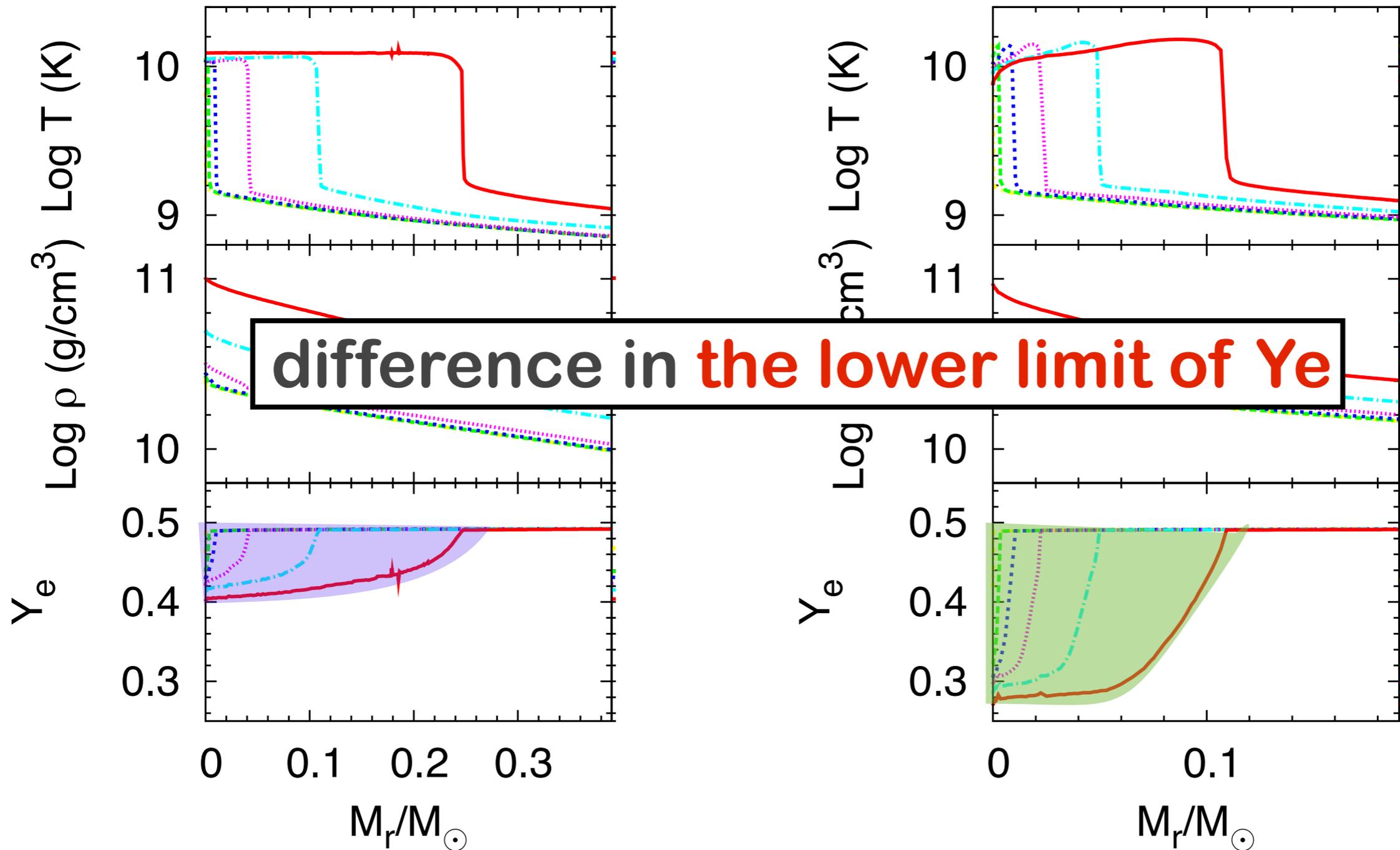
λ_{ec}/Y_i [1/sec]

(Juodagalvis+ 2010)

New physics

only by
free-protons

by protons
+ n-rich isotopes



Until core collapse

Short summary 3

- odd proton number isotopes
 - > **minor** (convective URCA?)
- the Coulomb correction
 - > **minor**
- neutron-rich isotopes
 - > **important**

Conclusion

- ECSNe arise from **the most massive SAGB stars**.
- In the star, a highly degenerate ONe core forms.
- Some robust properties of ECSN can be expected from the SAGB structure.

- Core evolution can be divided into 4 sub-phases.
- **Deflagration phase** is much complicated.
- Core collapse is finally induced by electron capture in the NSE region.

- Electron capture by odd proton number nuclei and the Coulomb correction on the rates have minor effect.
- **Electron capture by neutron-rich isotopes** in NSE compositions should be treated.

Evolution of the most
massive AGB star
The End
as a progenitor for ECSNe

Thank you for your attention!

Koh Takahashi¹, Takashi Yoshida², Hideyuki Umeda¹

¹Department of Astronomy, The University of Tokyo

²Yukawa Institute for Theoretical Physics, Kyoto University