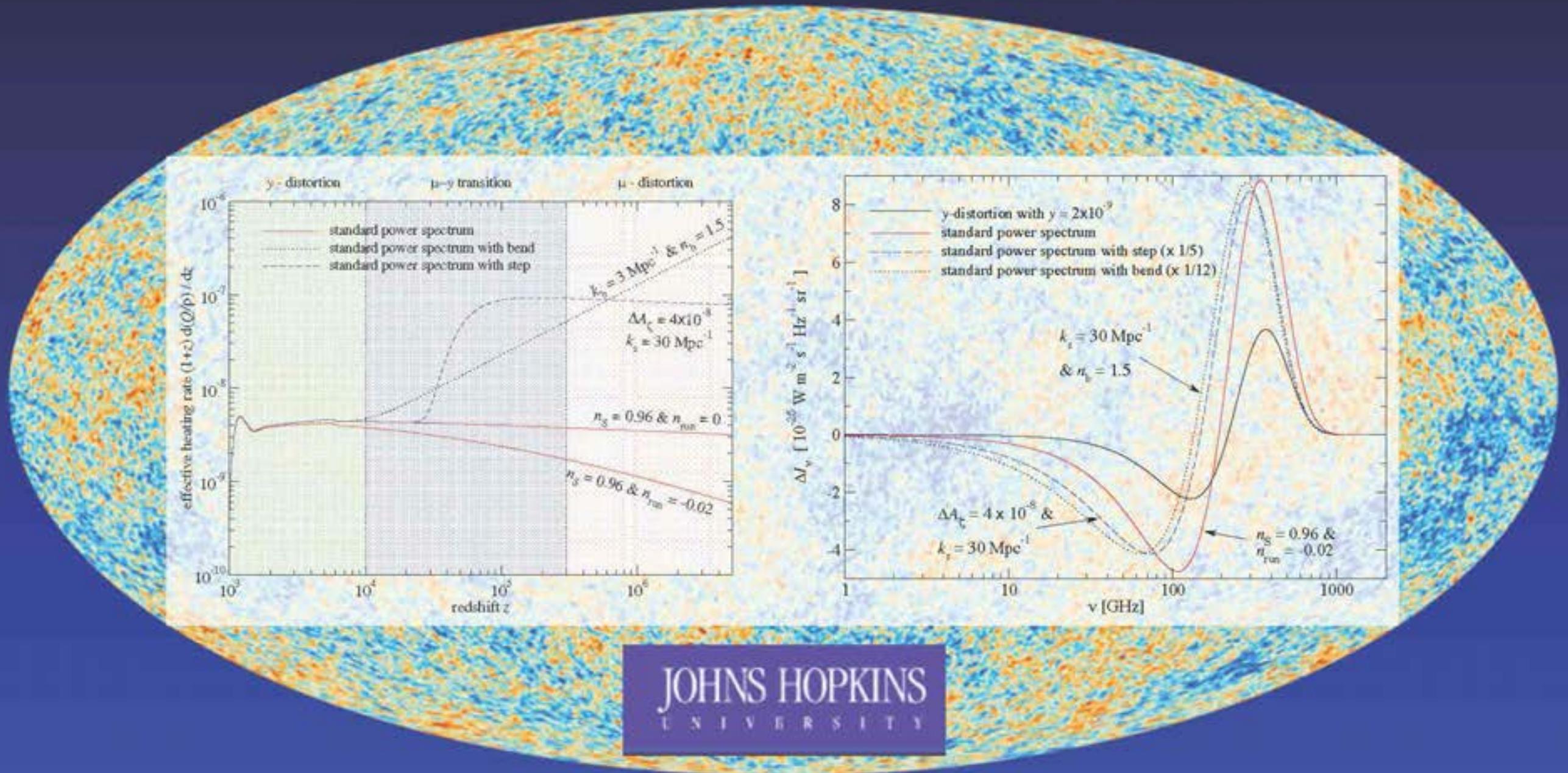
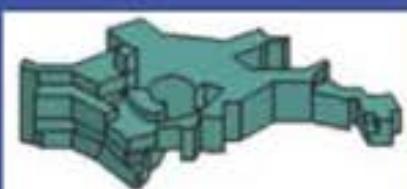


Science with CMB Spectral Distortions: a New Window to Early-Universe Physics



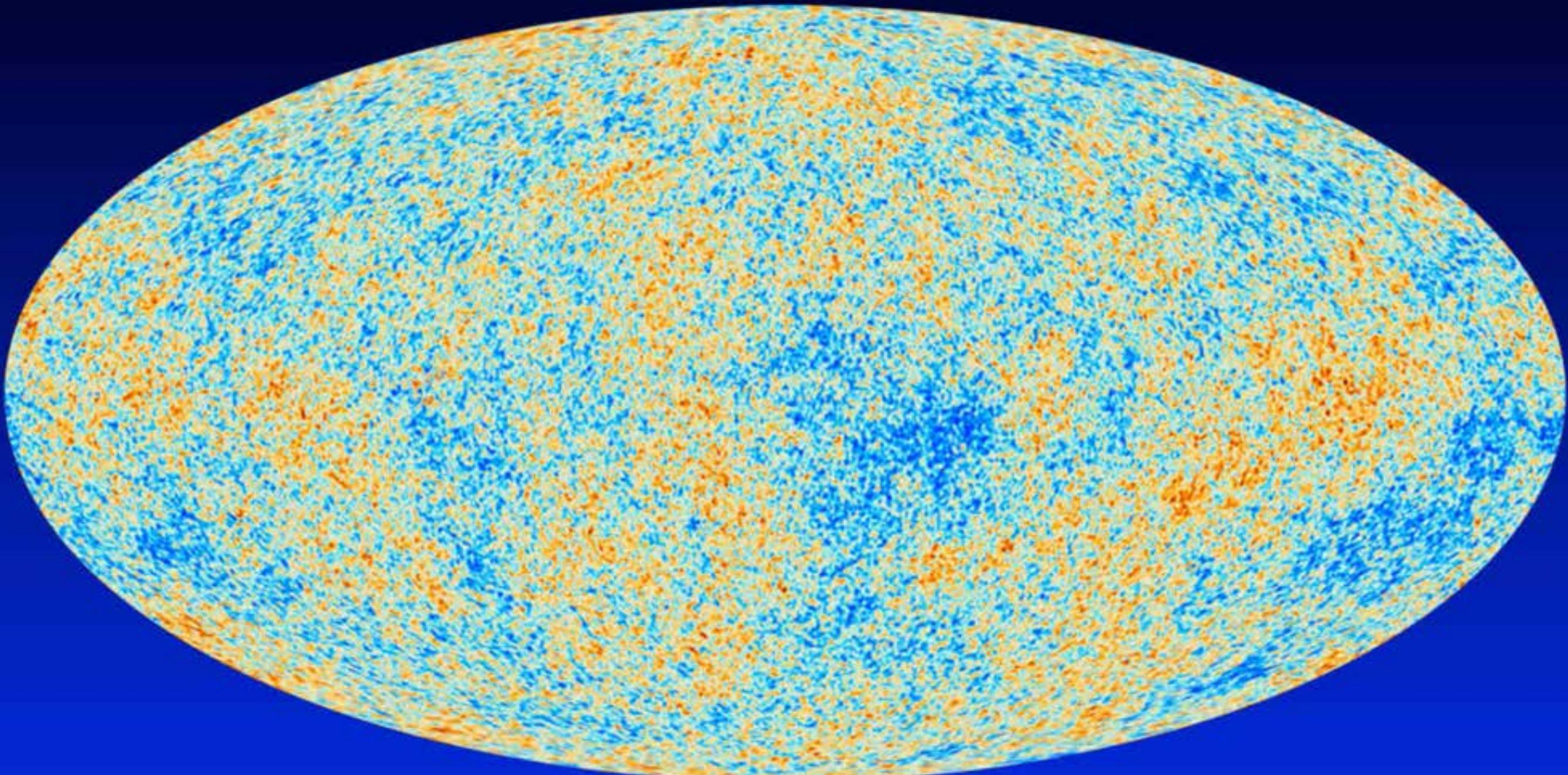
Jens Chluba
APC Seminar

IPMU, Kashiwa, Japan, Oct 4th, 2013



Collaborators: M. Kamionkowski & D. Jeong (JHU), R.A. Sunyaev & R. Khatri (MPA), A.L. Erickcek & I. Ben-Dayan (CITA), D. Grin (IAS)

Cosmic Microwave Background Anisotropies



Planck all sky map

- CMB has a blackbody spectrum in every direction
- tiny variations of the CMB temperature $\Delta T/T \sim 10^{-5}$

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Here we are Interested in the CMB Monopole Signal!!!

COBE/FIRAS

$$T_0 = (2.726 \pm 0.001) \text{ K}$$

Mather et al., 1994, ApJ, 420, 439
Fixsen et al., 1996, ApJ, 473, 576
Fixsen, 2003, ApJ, 594, 67
Fixsen, 2009, ApJ, 707, 916

- CMB monopole is 10000 - 100000 times larger than fluctuations!

COBE / FIRAS (Far InfraRed Absolute Spectrophotometer)



$$T_0 = 2.725 \pm 0.001 \text{ K}$$

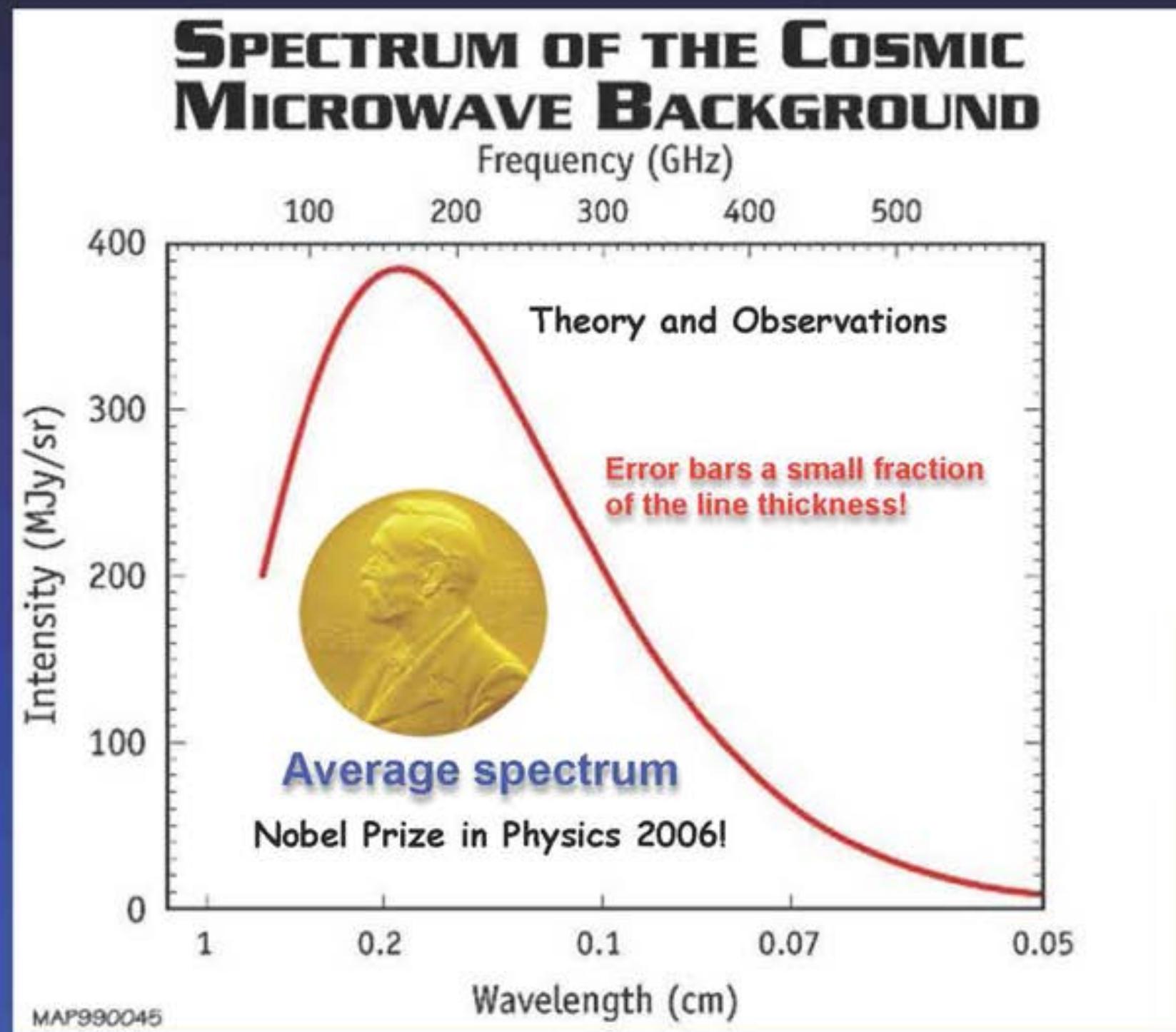
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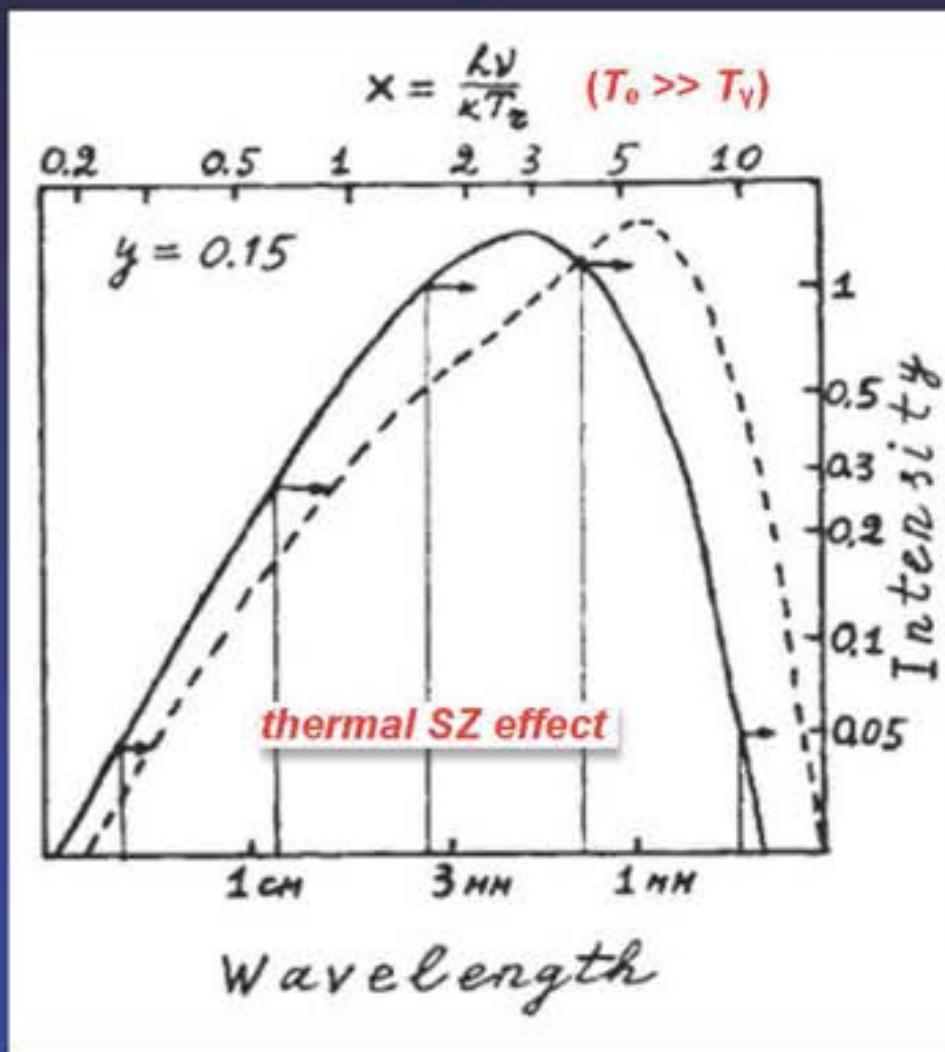
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Small Sneak Preview....

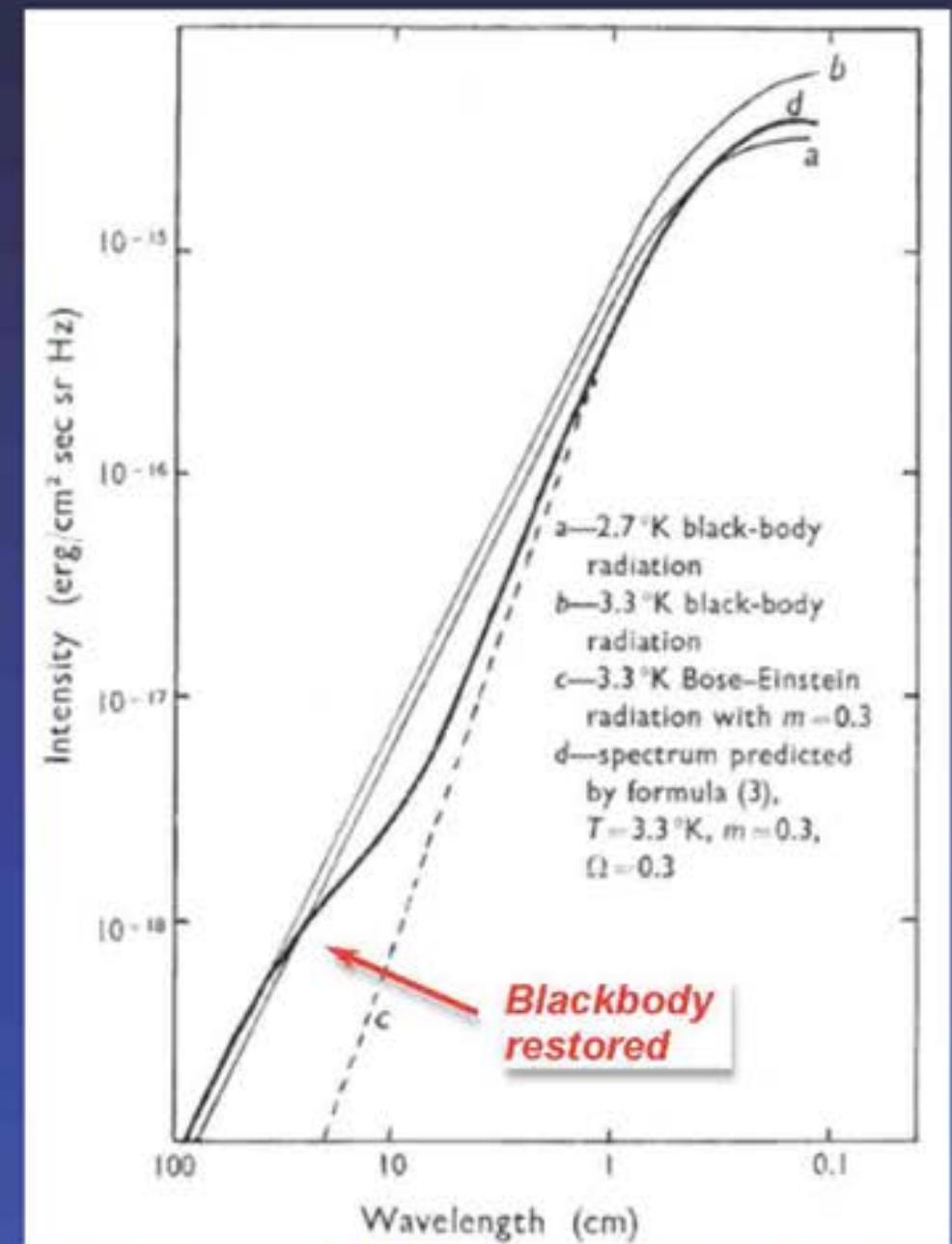
Compton γ -distortion



Sunyaev & Zeldovich, 1980, ARAA, 18, 537

- also known from thSZ effect
- up-scattering of CMB photon
- important at late times ($z < 50000$)
- scattering inefficient

Chemical potential μ -distortion



Sunyaev & Zeldovich, 1970, ApSS, 2, 66

- important at very times ($z > 50000$)
- scattering very efficient

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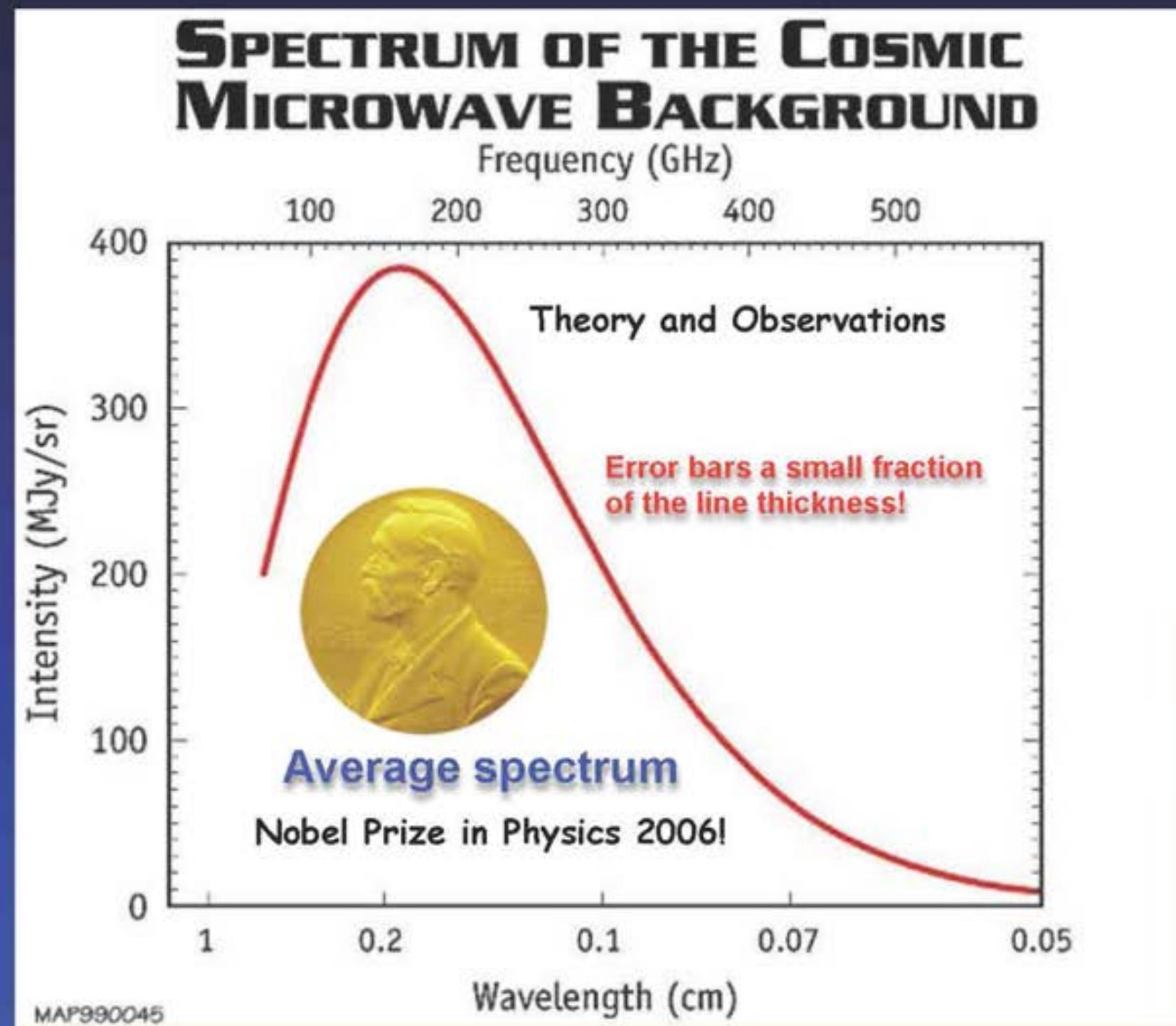


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Only very small distortions of CMB spectrum are still allowed!

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Full thermodynamic equilibrium (certainly valid at very high redshift)

- CMB has a blackbody spectrum at every time (not affected by expansion)
- Photon number density and energy density determined by temperature T_γ

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Perturbing full equilibrium by

- Energy injection (interaction *matter* \leftrightarrow *photons*)
- Production of (energetic) photons and/or particles (i.e. change of entropy)

→ CMB spectrum deviates from a pure blackbody

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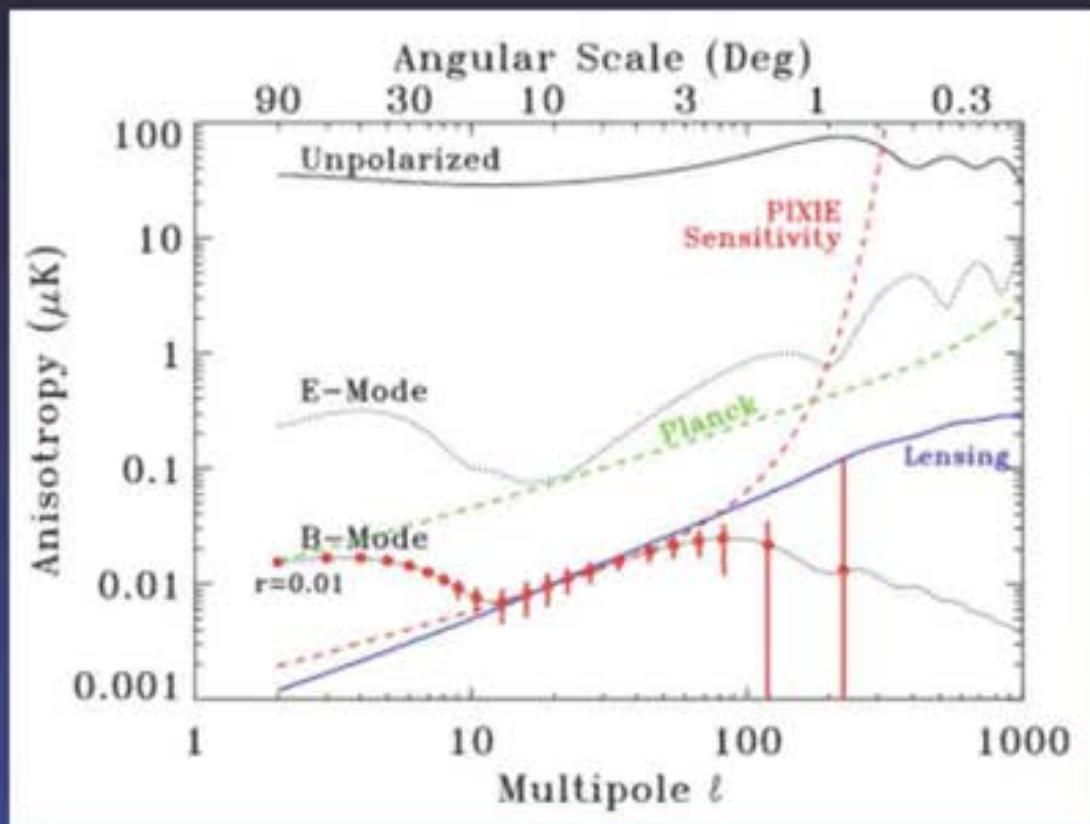
Measurements of CMB spectrum place very tight limits on the thermal history of our Universe!

Why bother? No distortion detected so far!??

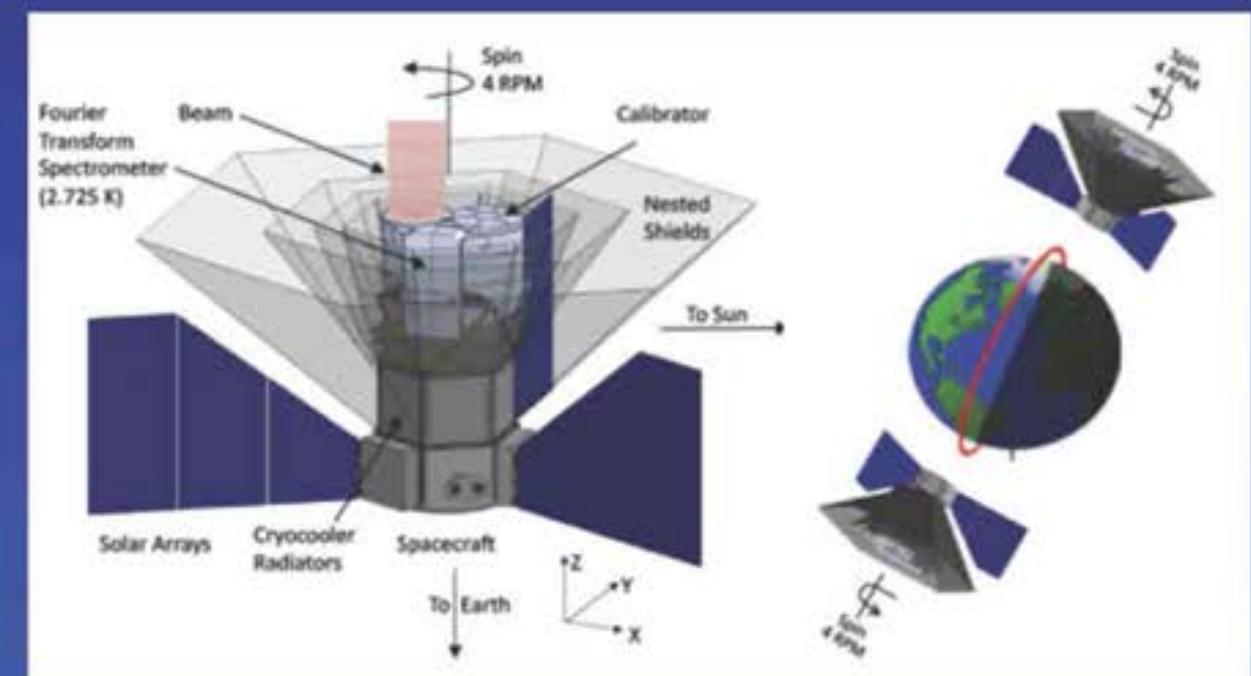
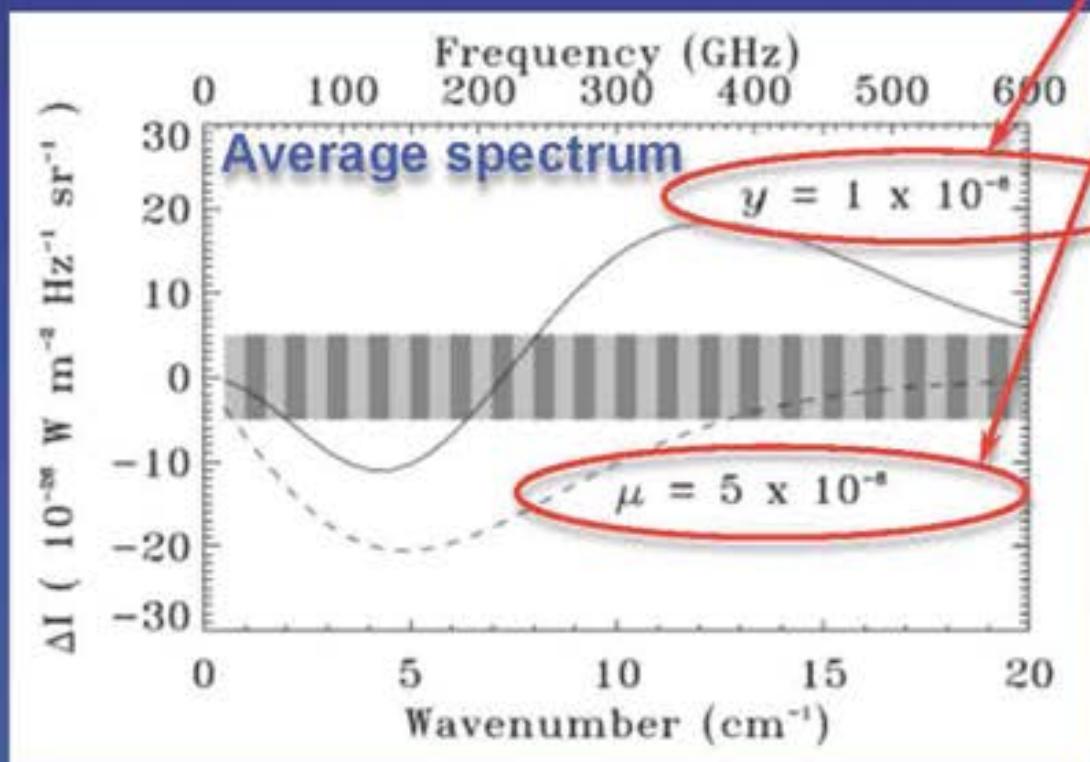
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PIXIE: Primordial Inflation Explorer

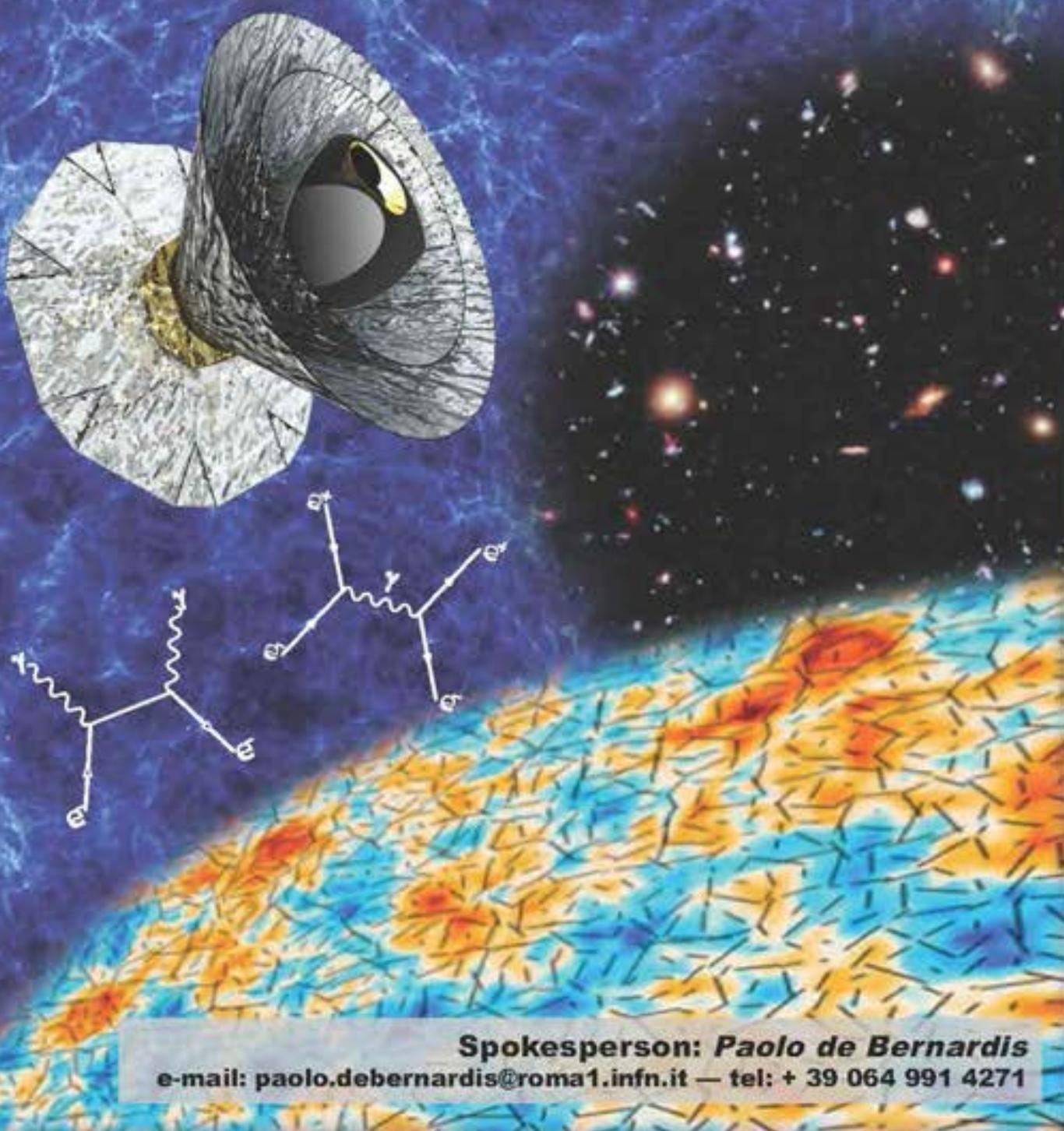


- 400 spectral channel in the frequency range 30 GHz and 6THz ($\Delta\nu \sim 15\text{GHz}$)
- about 1000 (!!!) times more sensitive than COBE/FIRAS
- B-mode polarization from inflation ($r \approx 10^{-3}$)
- improved limits on μ and y
- was proposed 2011 as NASA EX mission (i.e. cost $\sim 200\text{ M\$}$)



PRISM

Probing cosmic structures and radiation
with the ultimate polarimetric spectro-imaging
of the microwave and far-infrared sky



Instruments:

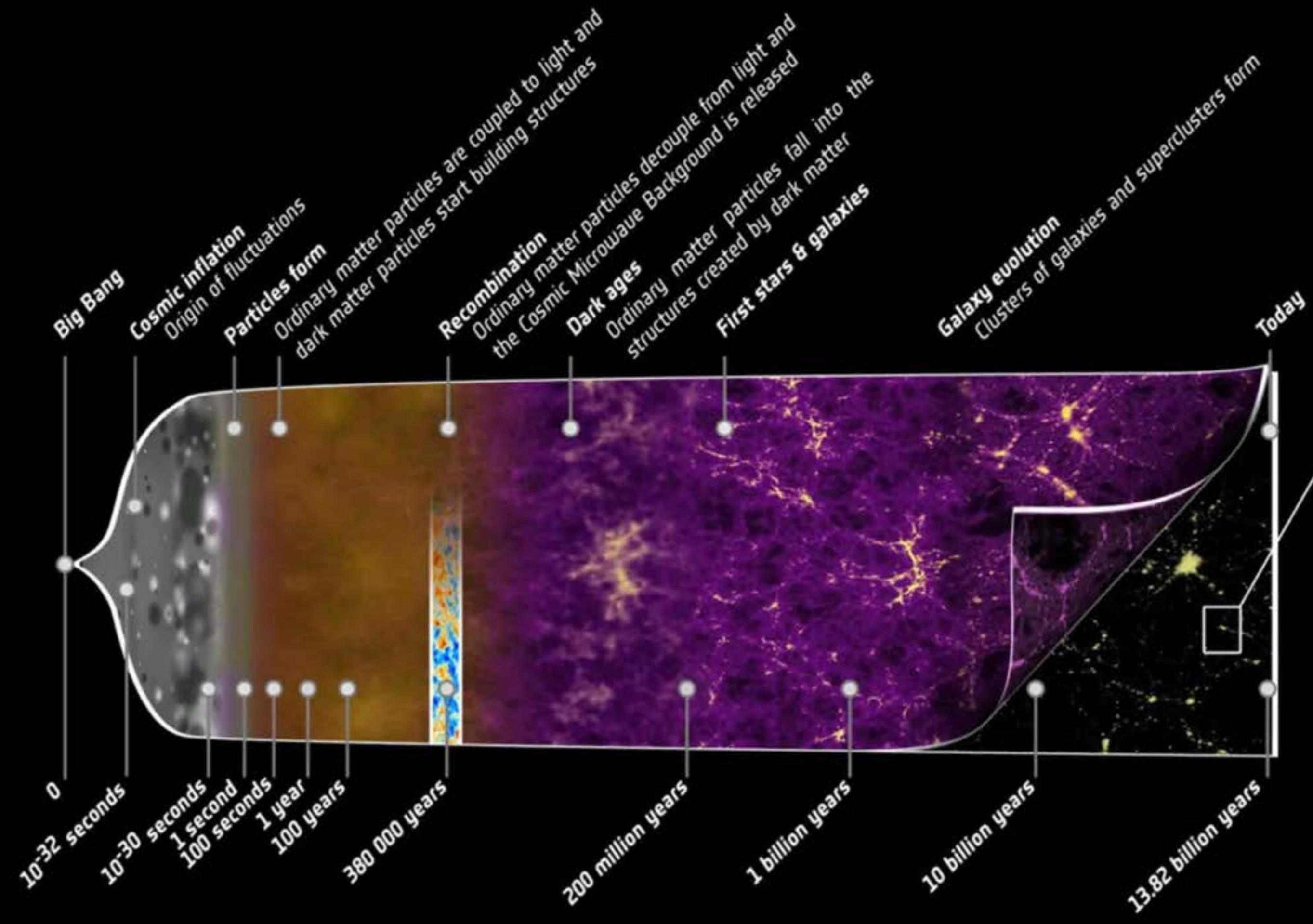
- L-class ESA mission
- White paper, May 24th, 2013
- Imager:
 - polarization sensitive
 - 3.5m telescope [arcmin resolution at highest frequencies]
 - 30GHz-6THz [30 broad ($\Delta\nu/\nu \sim 25\%$) and 300 narrow ($\Delta\nu/\nu \sim 2.5\%$) bands]
- Spectrometer:
 - FTS similar to PIXIE
 - 30GHz-6THz ($\Delta\nu \sim 15$ & 0.5 GHz)

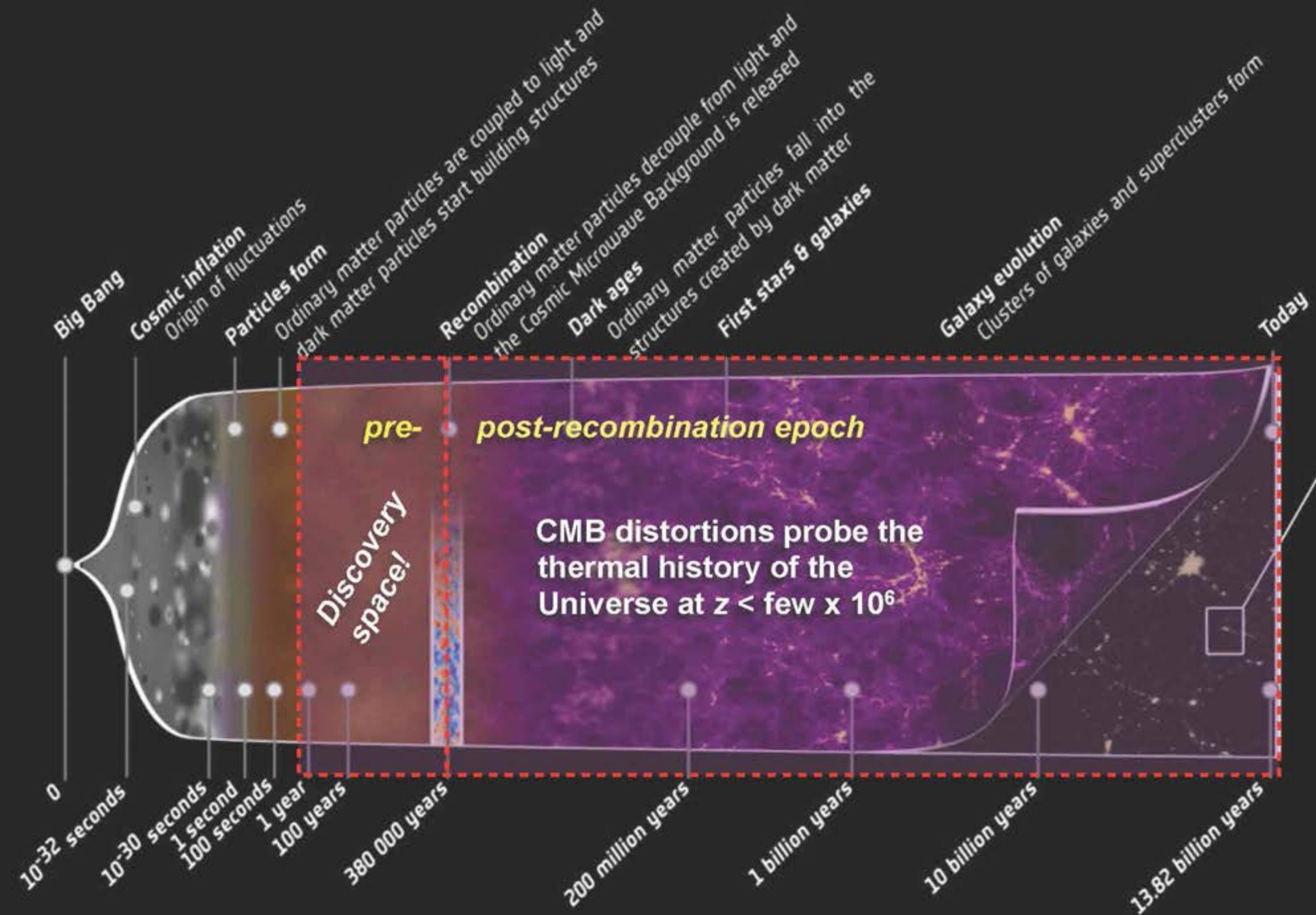
Some of the science goals:

- B-mode polarization from inflation ($r \approx 5 \times 10^{-4}$)
- count all SZ clusters $> 10^{14} M_{\odot}$
- CIB/large scale structure
- Galactic science
- *CMB spectral distortions*

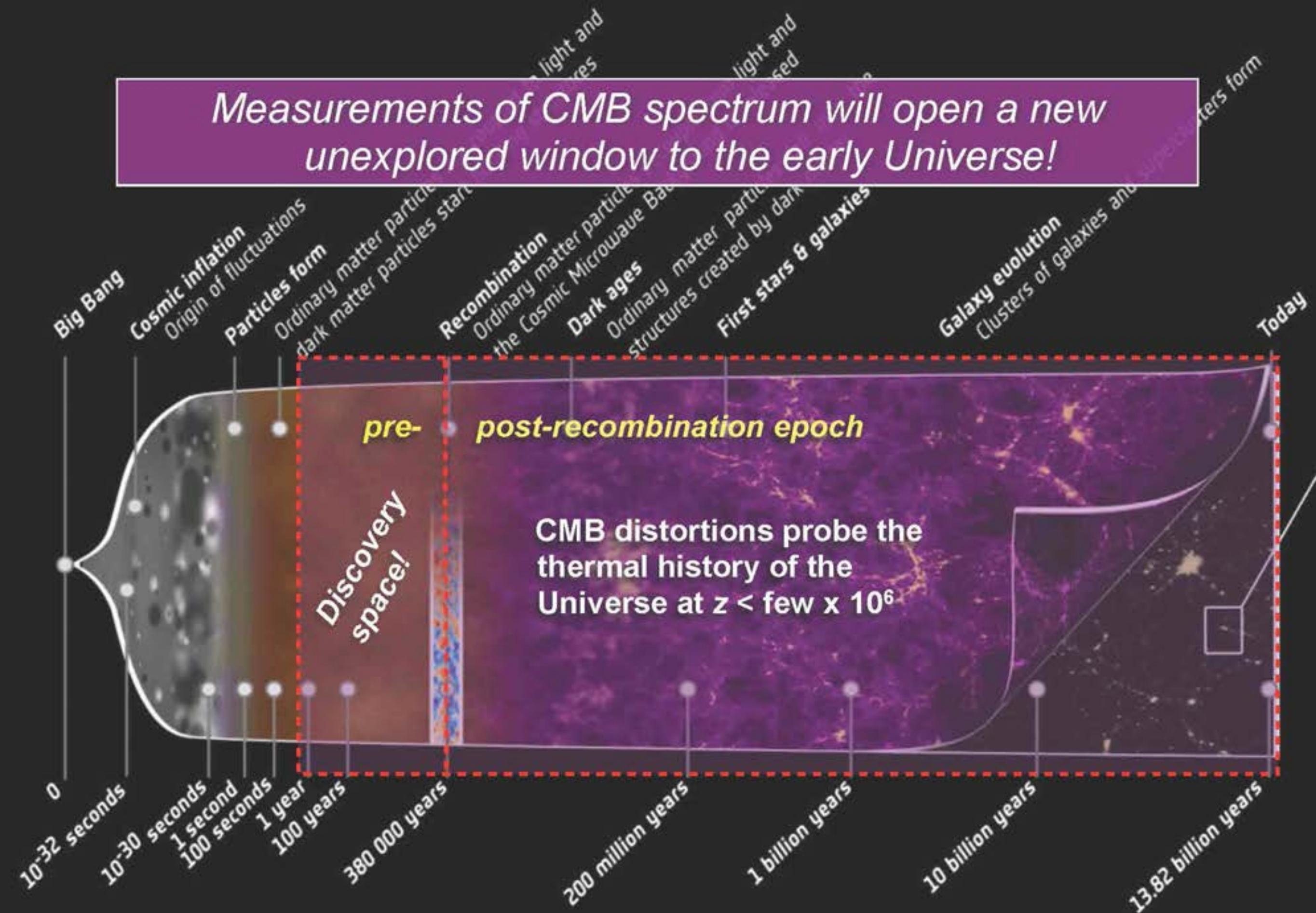
Sign up at:

<http://www.prism-mission.org/>





Measurements of CMB spectrum will open a new unexplored window to the early Universe!



How does the thermalization process work?

Some important conditions

- Plasma fully ionized before recombination ($z \sim 1000$)
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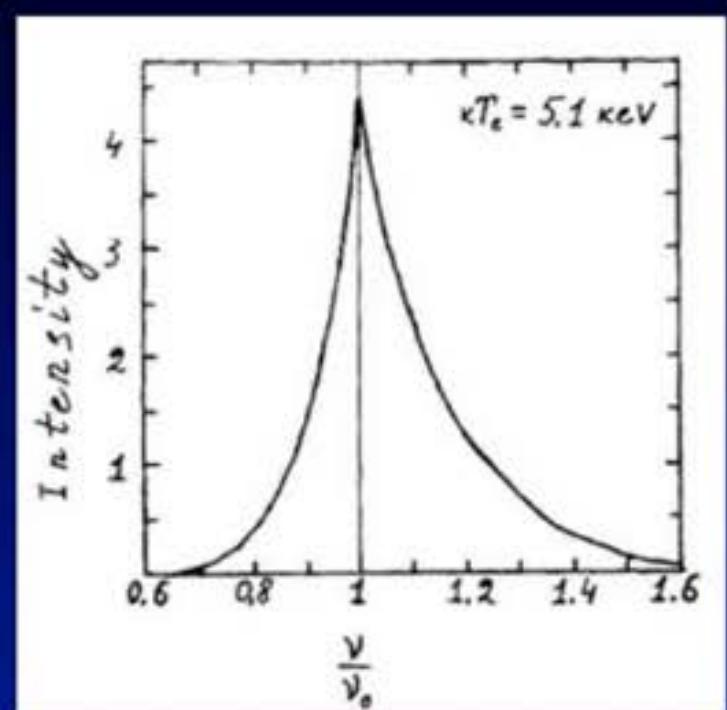
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- Hubble expansion
 - adiabatic cooling of photons [$T_\gamma \sim (1+z)$] and ordinary matter [$T_m \sim (1+z)^2$]
 - redshifting of photons

Redistribution of photons by Compton scattering

- Compton scattering $e + \gamma \leftrightarrow e' + \gamma'$

→ *redistribution of photons in frequency*

- up-scattering due to the **Doppler** effect for $h\nu < 4kT_e$
- down-scattering because of **recoil** (and stimulated recoil) for $h\nu > 4kT_e$
- **Doppler** broadening $\frac{\Delta\nu}{\nu} \simeq \sqrt{\frac{2kT_e}{m_e c^2}}$



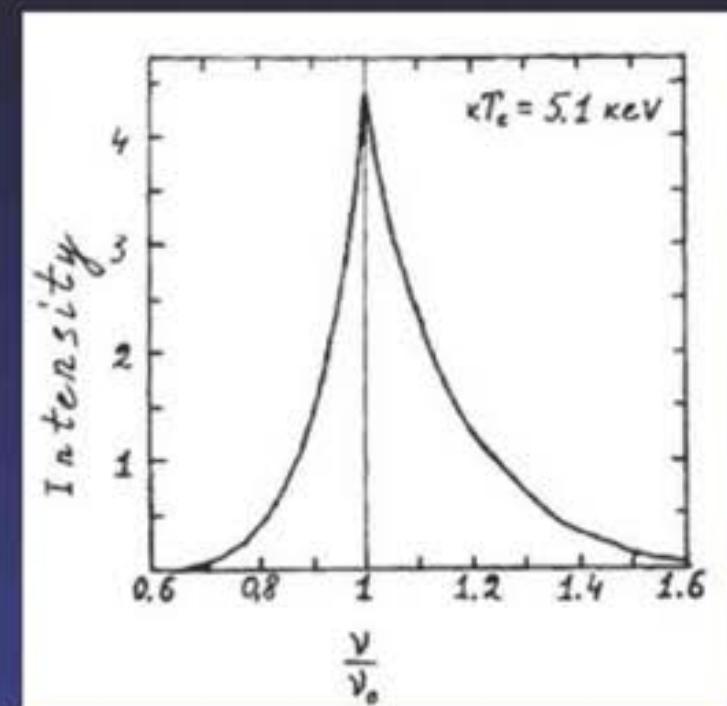
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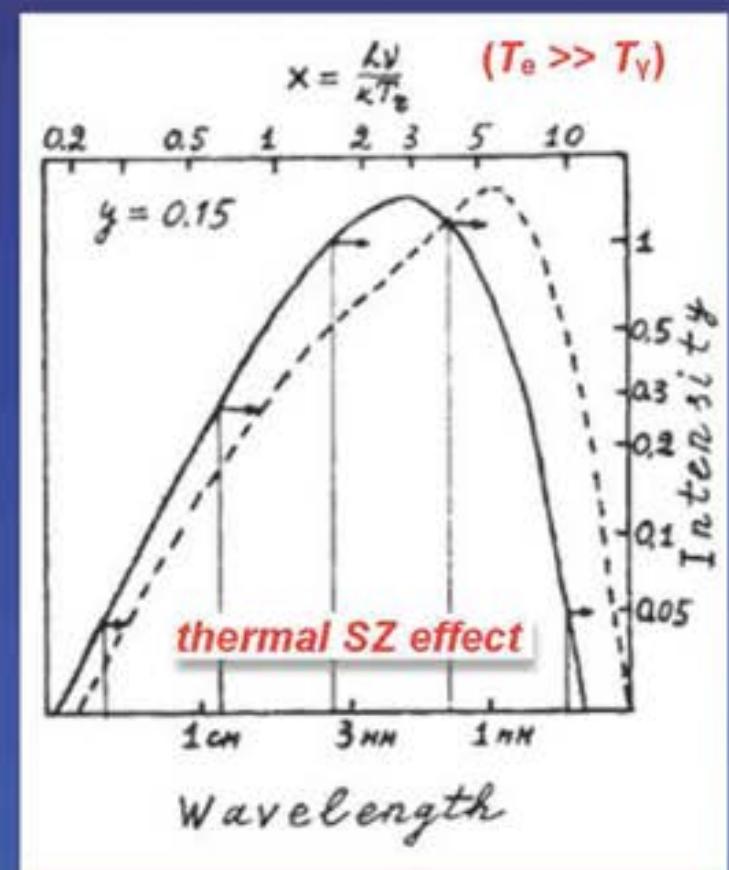
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- Kompaneets Equation → ‘pure’ **y-distortion**

$$\frac{\Delta I_\nu}{I_\nu} \simeq y \frac{x e^x}{e^x - 1} \left[x \frac{e^x + 1}{e^x - 1} - 4 \right],$$

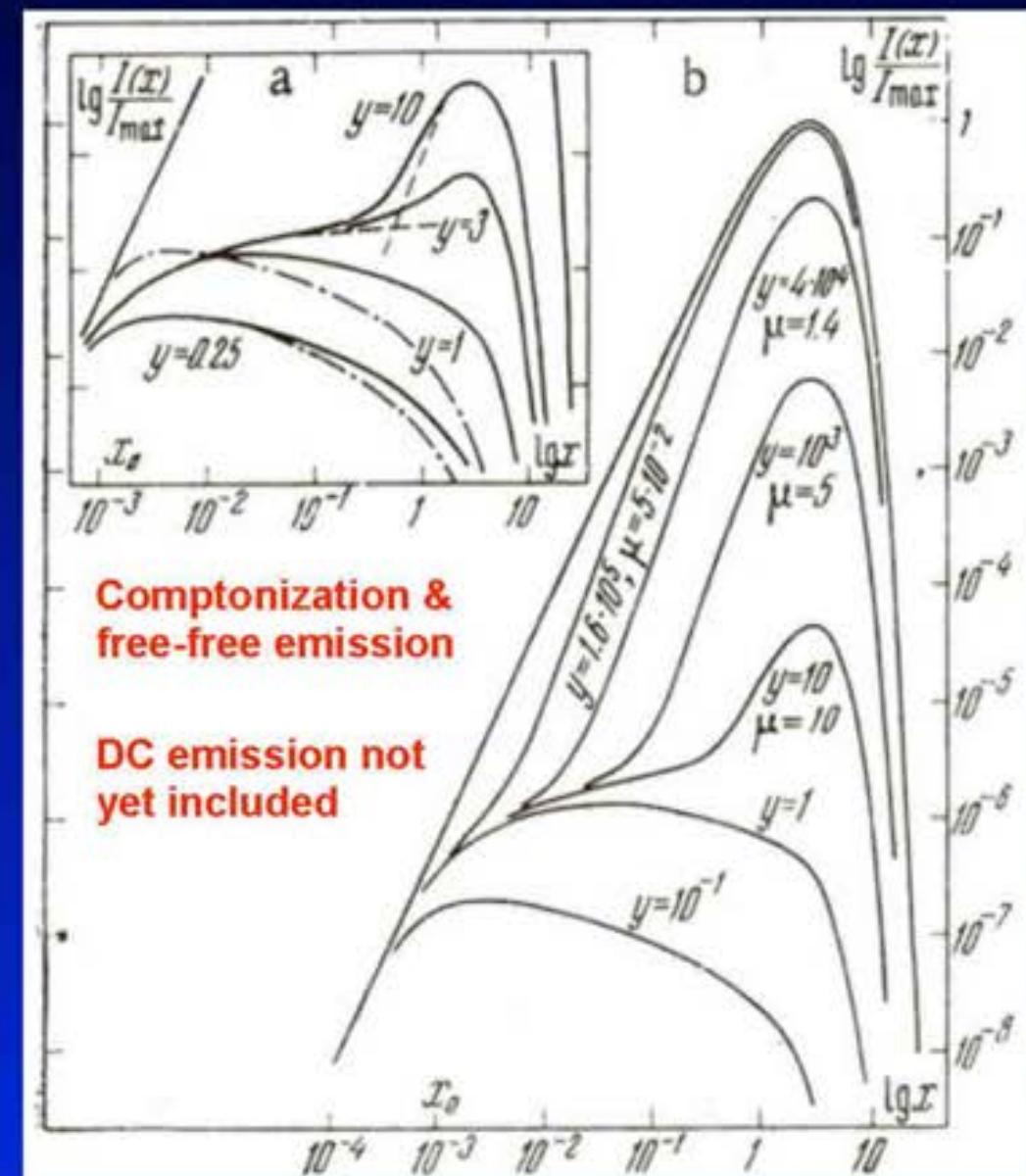
Temperature difference

where $x = \frac{h\nu}{kT_\gamma}$ and $y = \int \frac{k(T_e - T_\gamma)}{m_e c^2} \sigma_T n_e dl \ll 1$



Adjusting the photon number

- Bremsstrahlung $e + p \leftrightarrow e' + p + \gamma$
 - 1. order α correction to *Coulomb* scattering
 - production of low frequency photons
 - important for the evolution of the distortion at low frequencies and late times ($z < 2 \times 10^5$)
- Double Compton scattering
(Lightman 1981; Thorne, 1981)
$$e + \gamma \leftrightarrow e' + \gamma' + \gamma_2$$
 - 1. order α correction to *Compton* scattering
 - was only included later (Danese & De Zotti, 1982)
 - production of low frequency photons
 - very important at high redshifts ($z > 2 \times 10^5$)



Illarionov & Sunyaev, 1975, Sov. Astr, 18, pp.413

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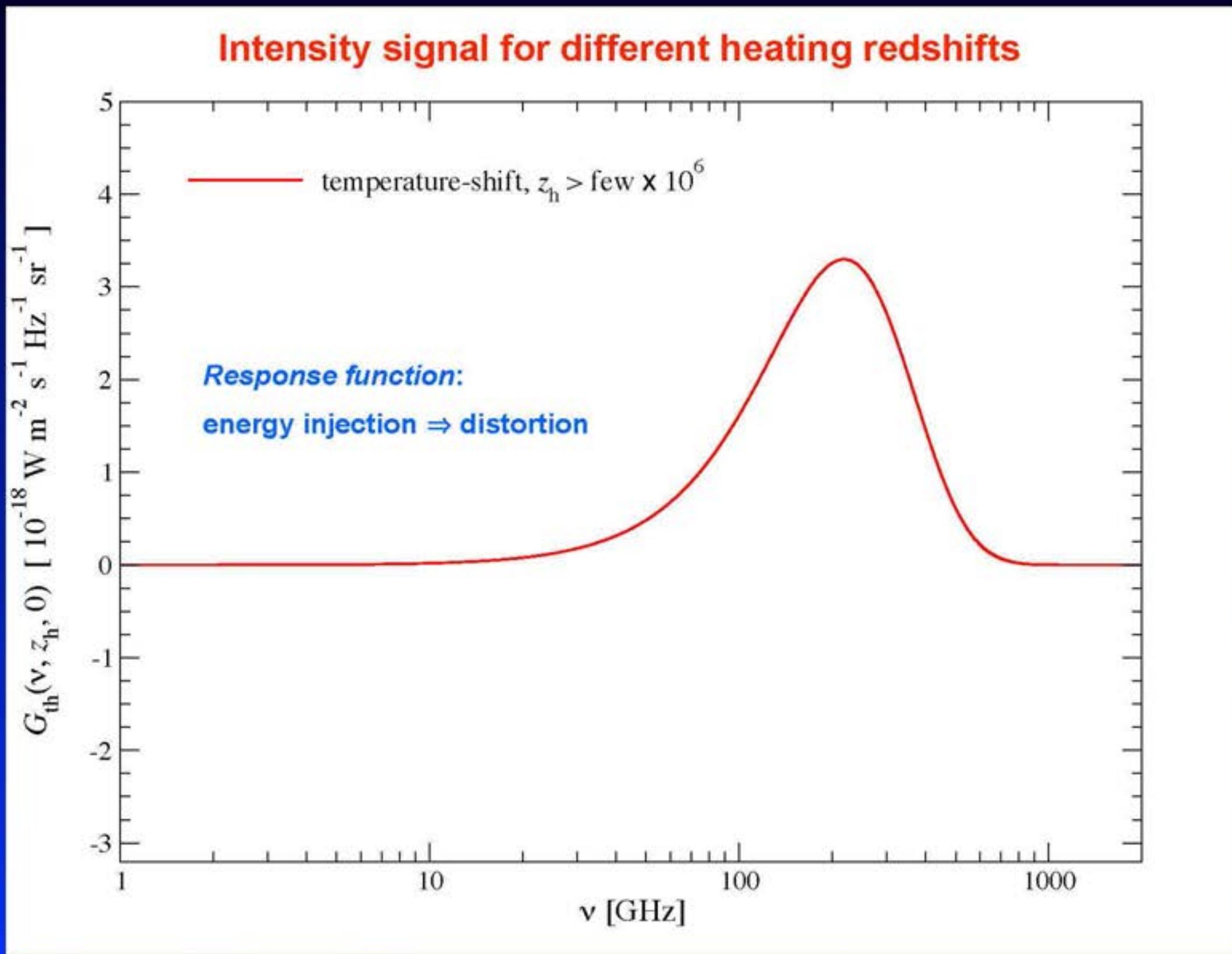
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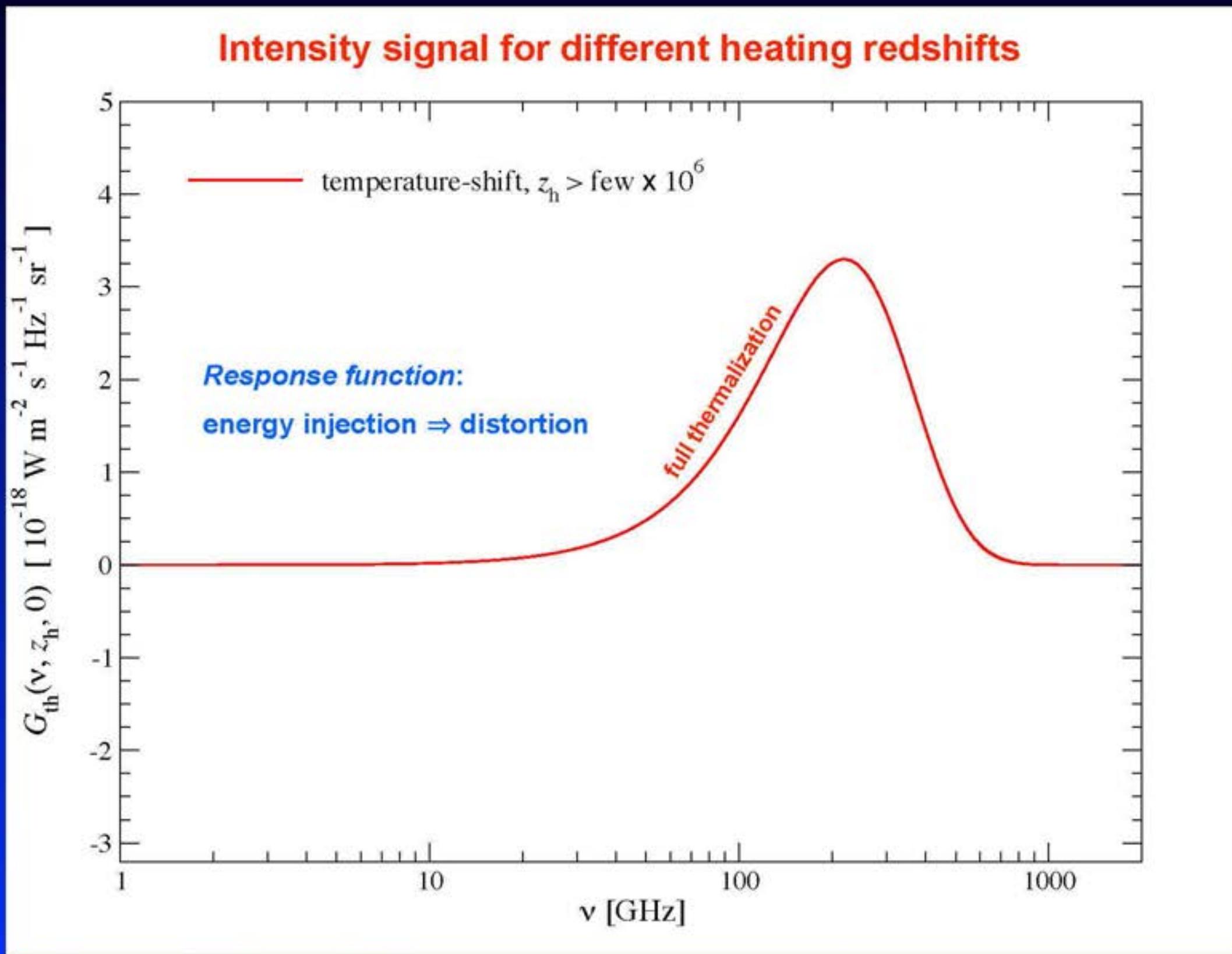
Thermalization Green's function

- Fast and quasi-exact! No additional approximations!

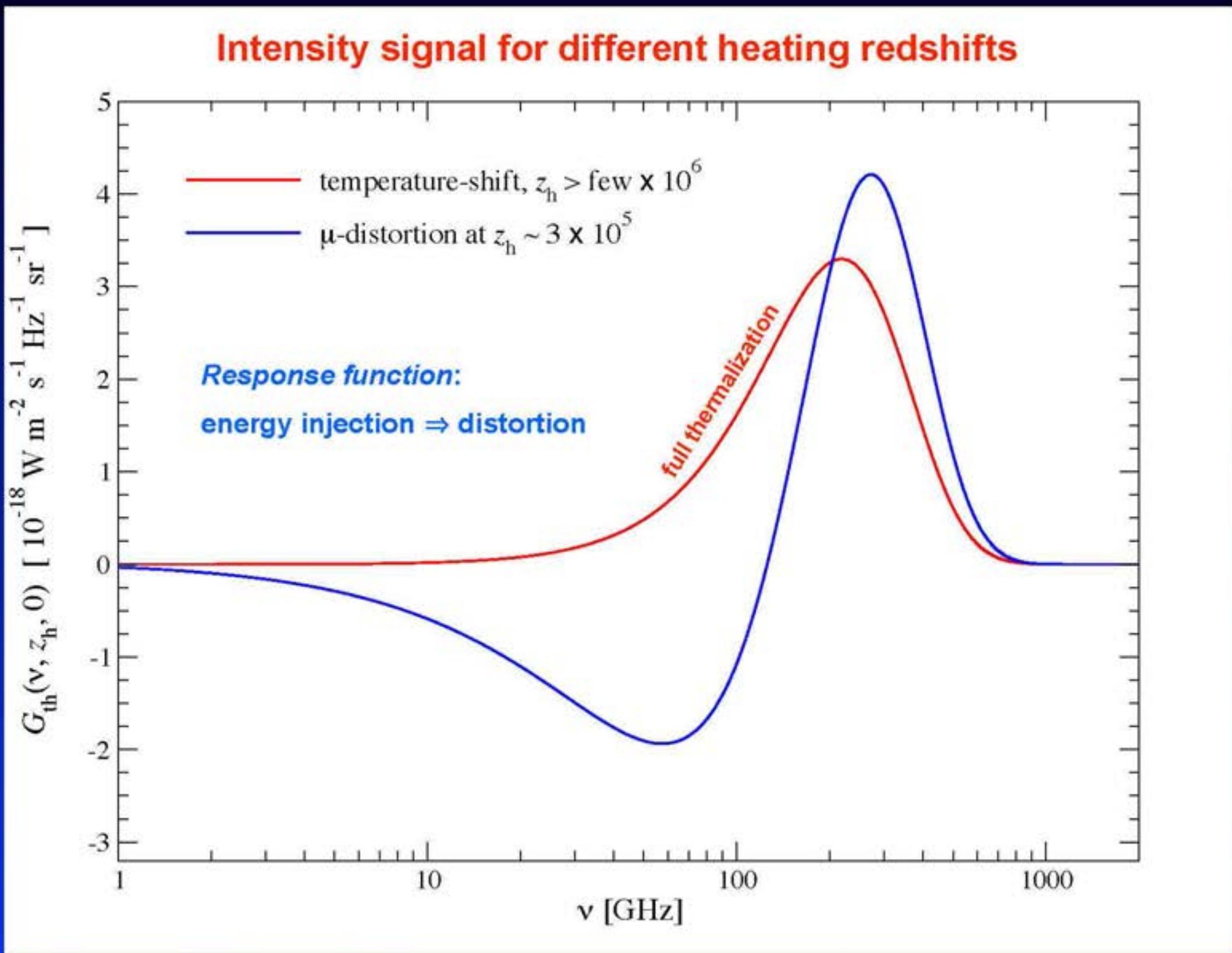
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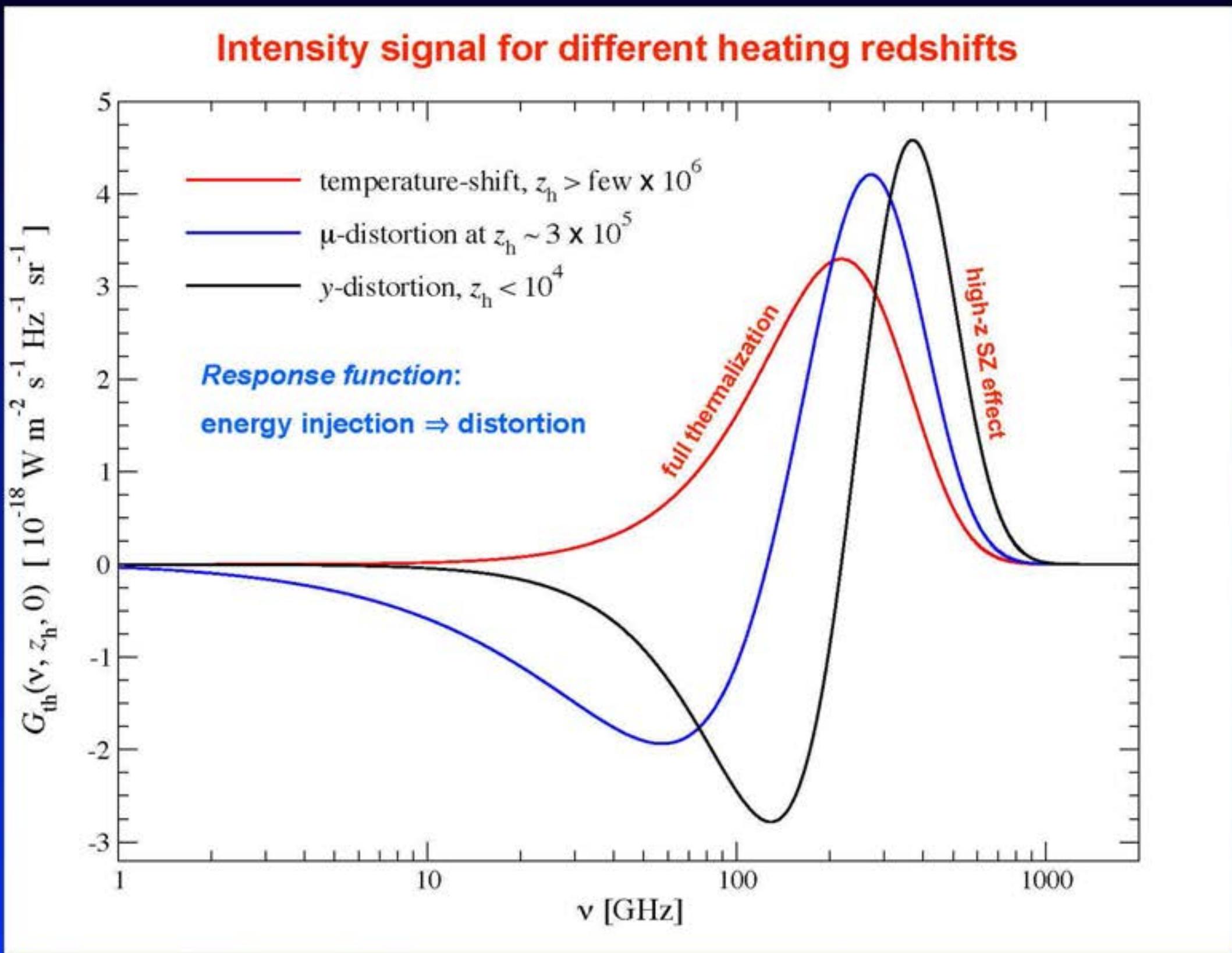
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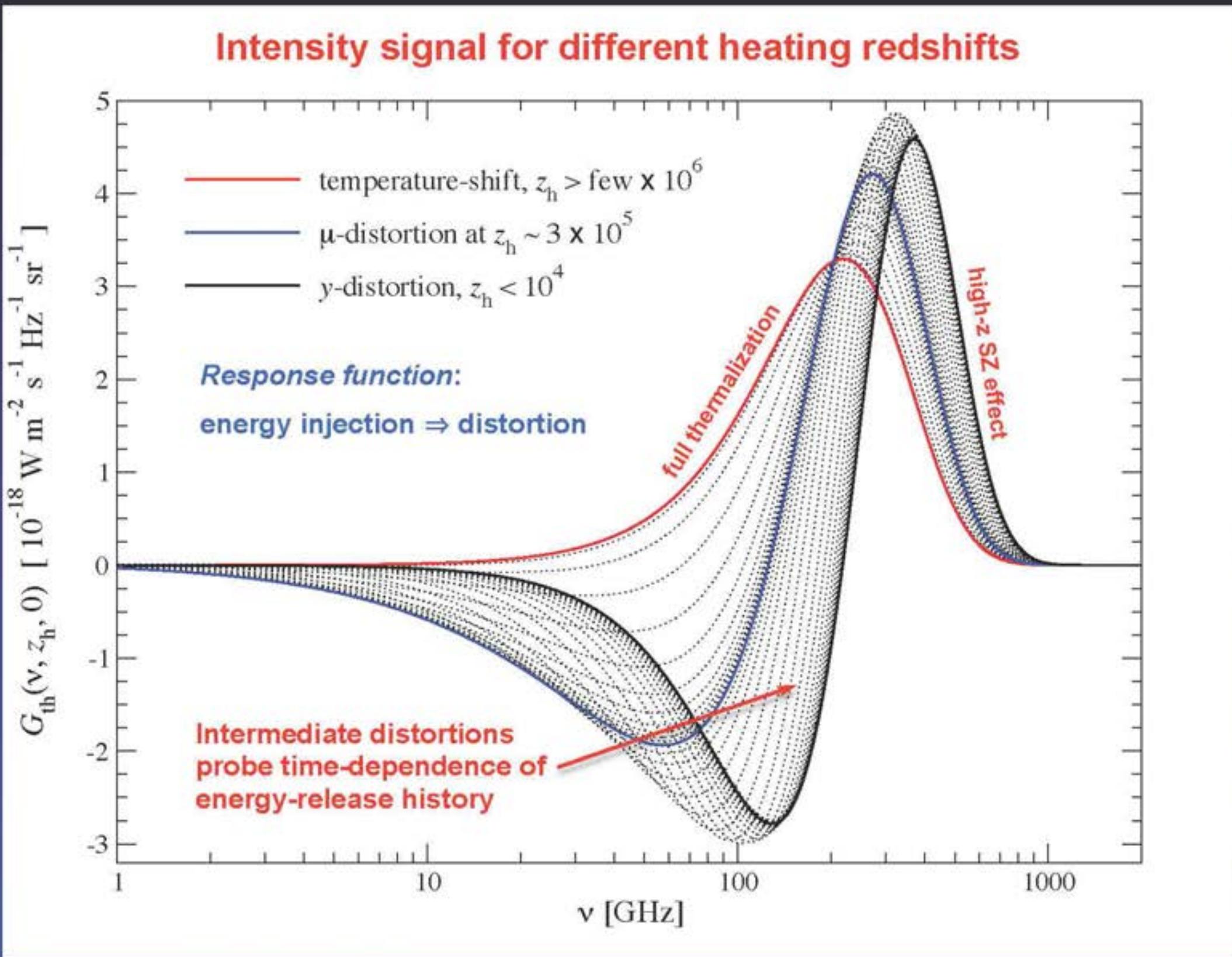
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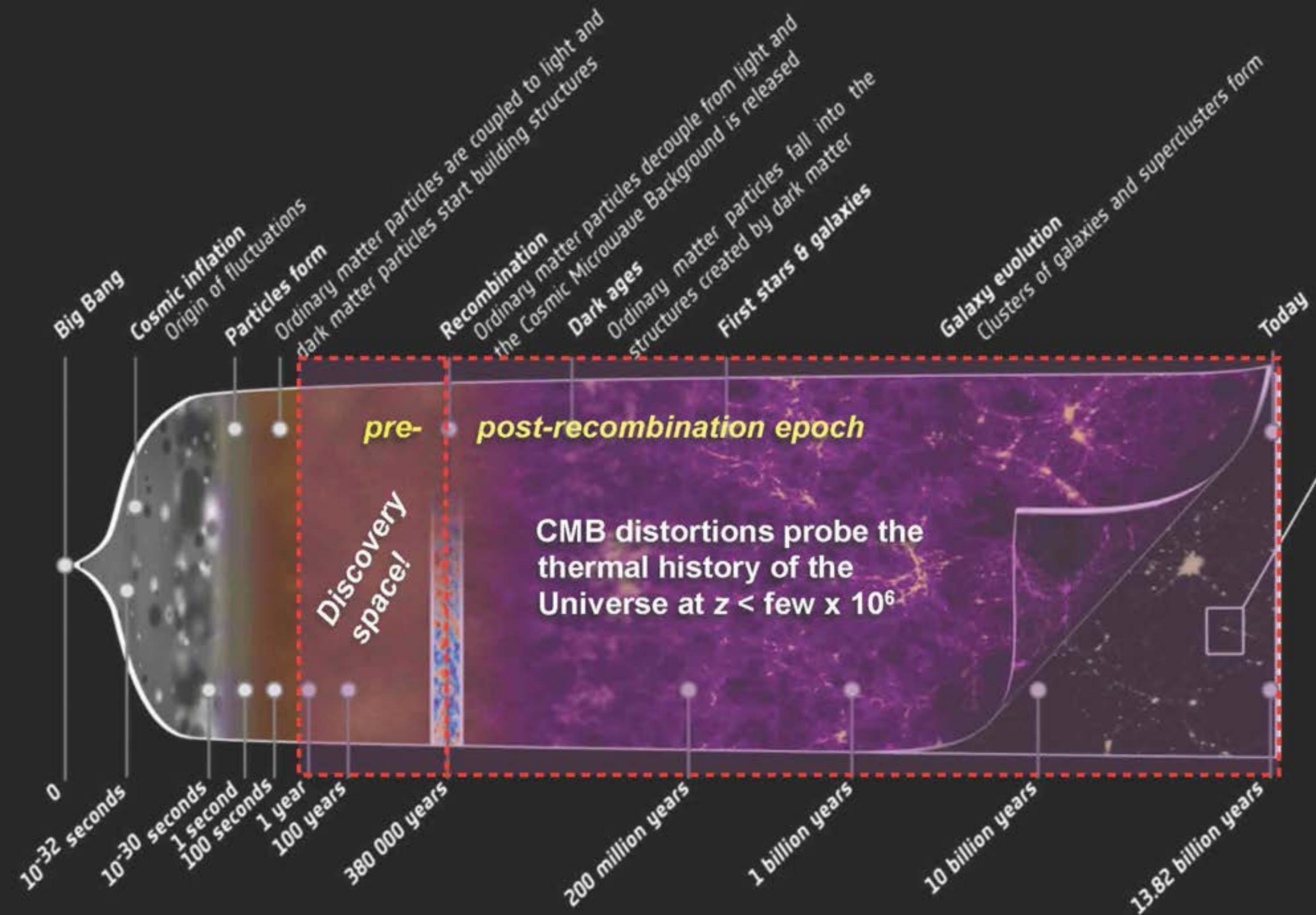


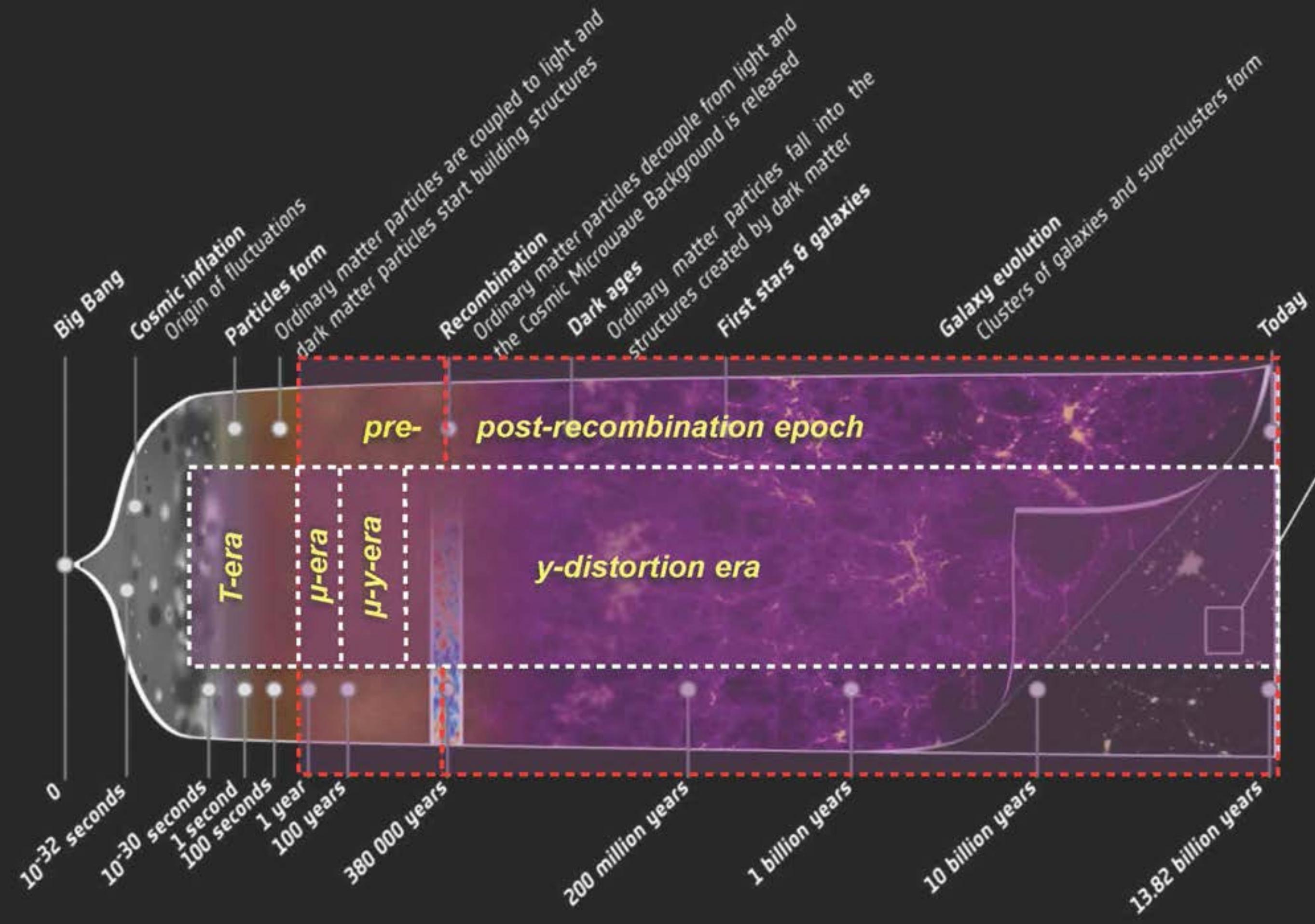
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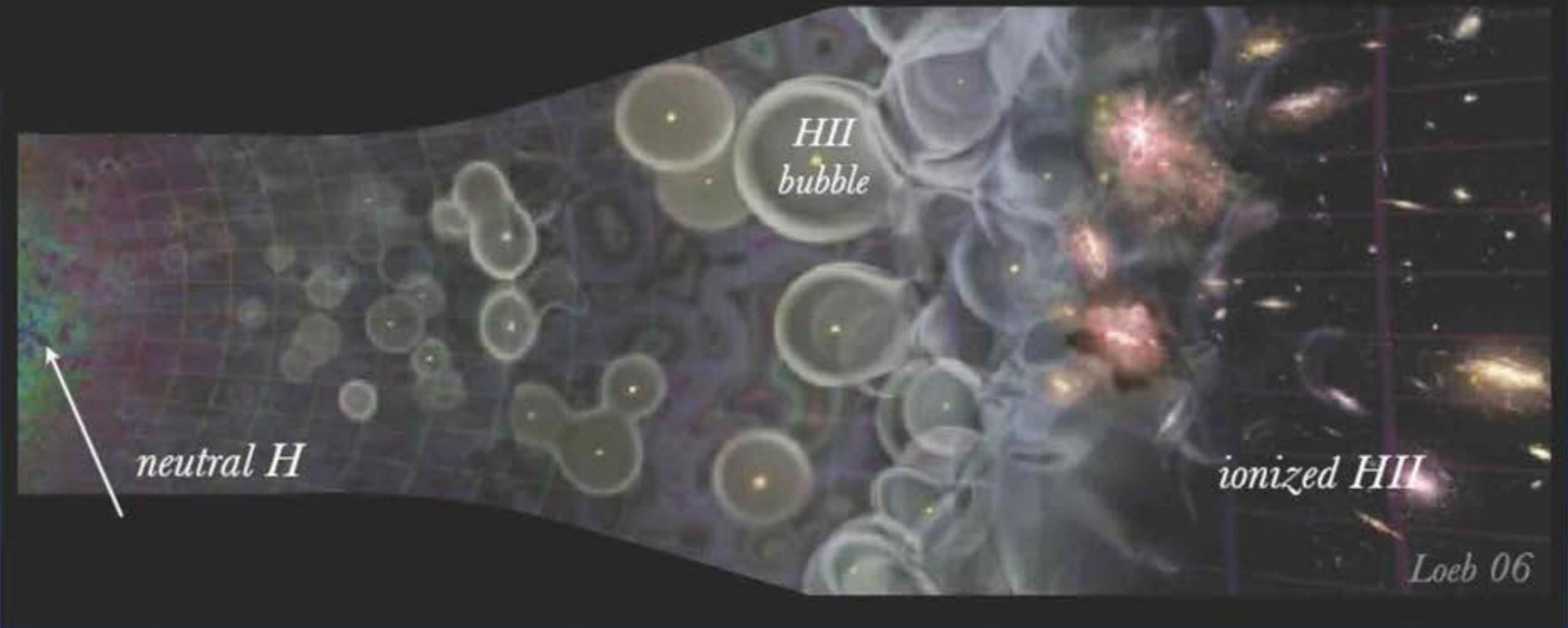


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Reionization and structure formation

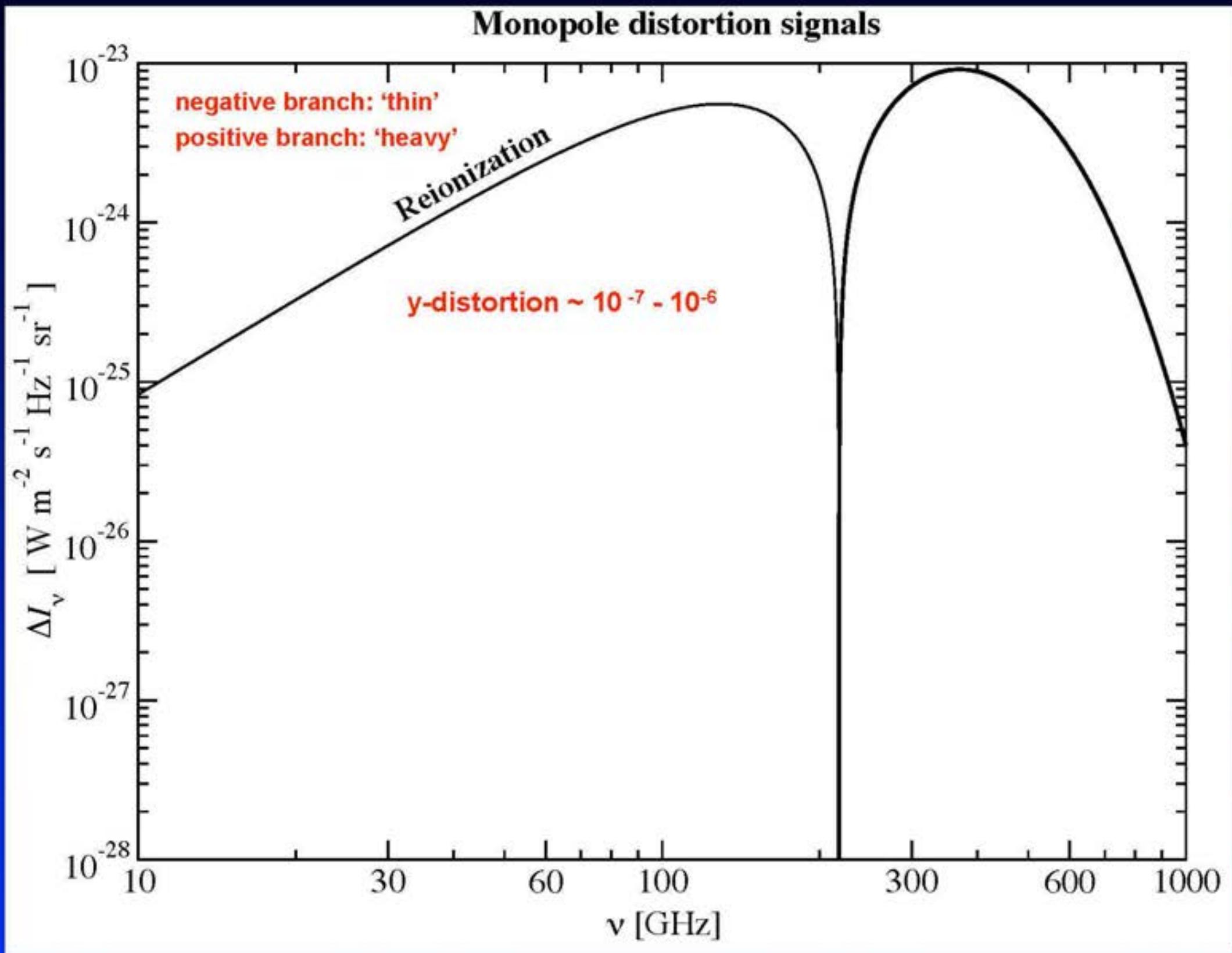
Simple estimates for the distortion



- Gas temperature $T \approx 10^4 \text{ K}$ $\Rightarrow y \approx \frac{kT_e}{m_e c^2} \approx 2 \times 10^{-7}$
- Thomson optical depth $\tau \approx 0.1$
- second order Doppler effect $y \approx \text{few} \times 10^{-8}$
- structure formation / SZ effect (e.g., Refregier et al., 2003) $y \approx \text{few} \times 10^{-7}\text{-}10^{-6}$

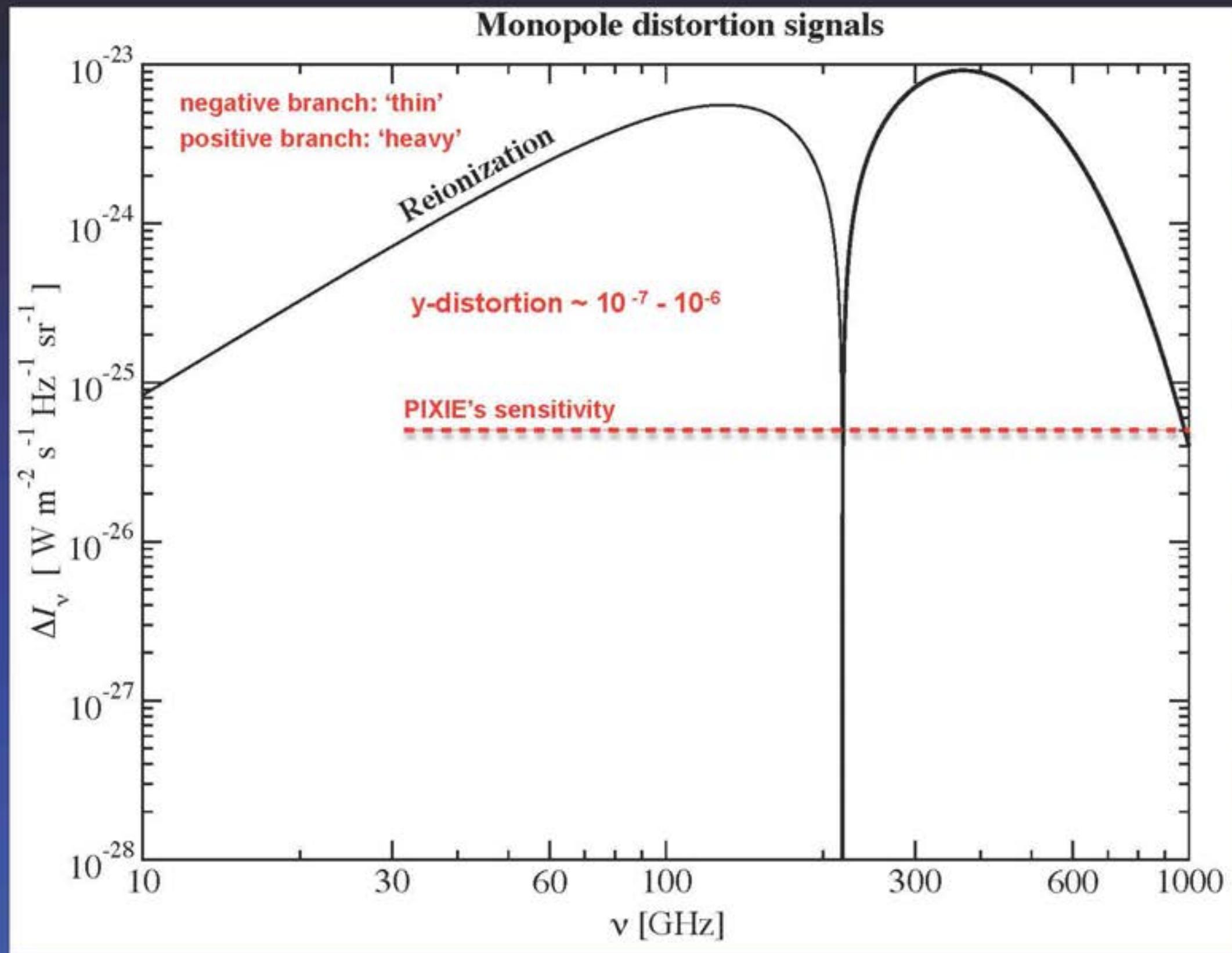
Average CMB spectral distortions

Absolute value of intensity signal



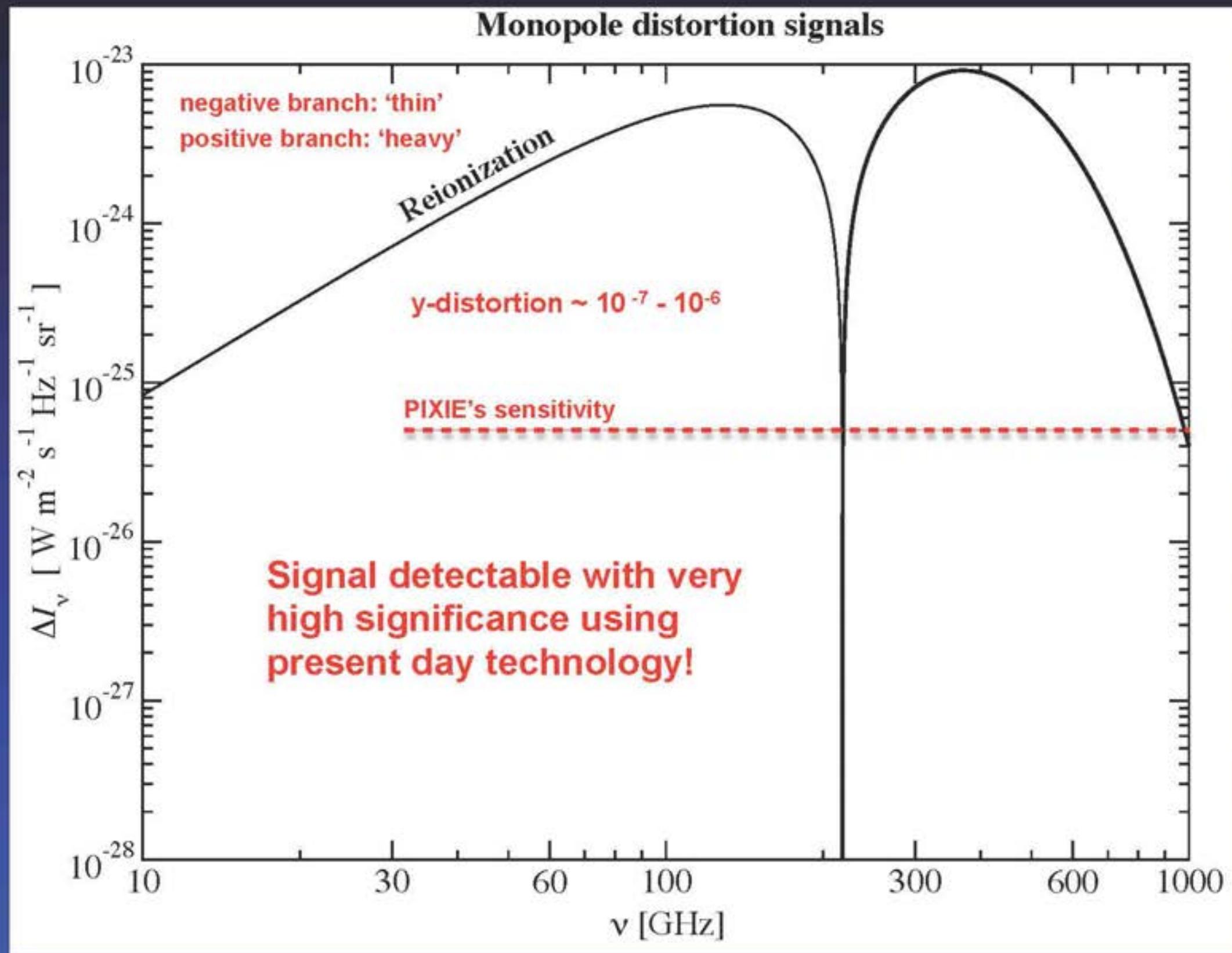
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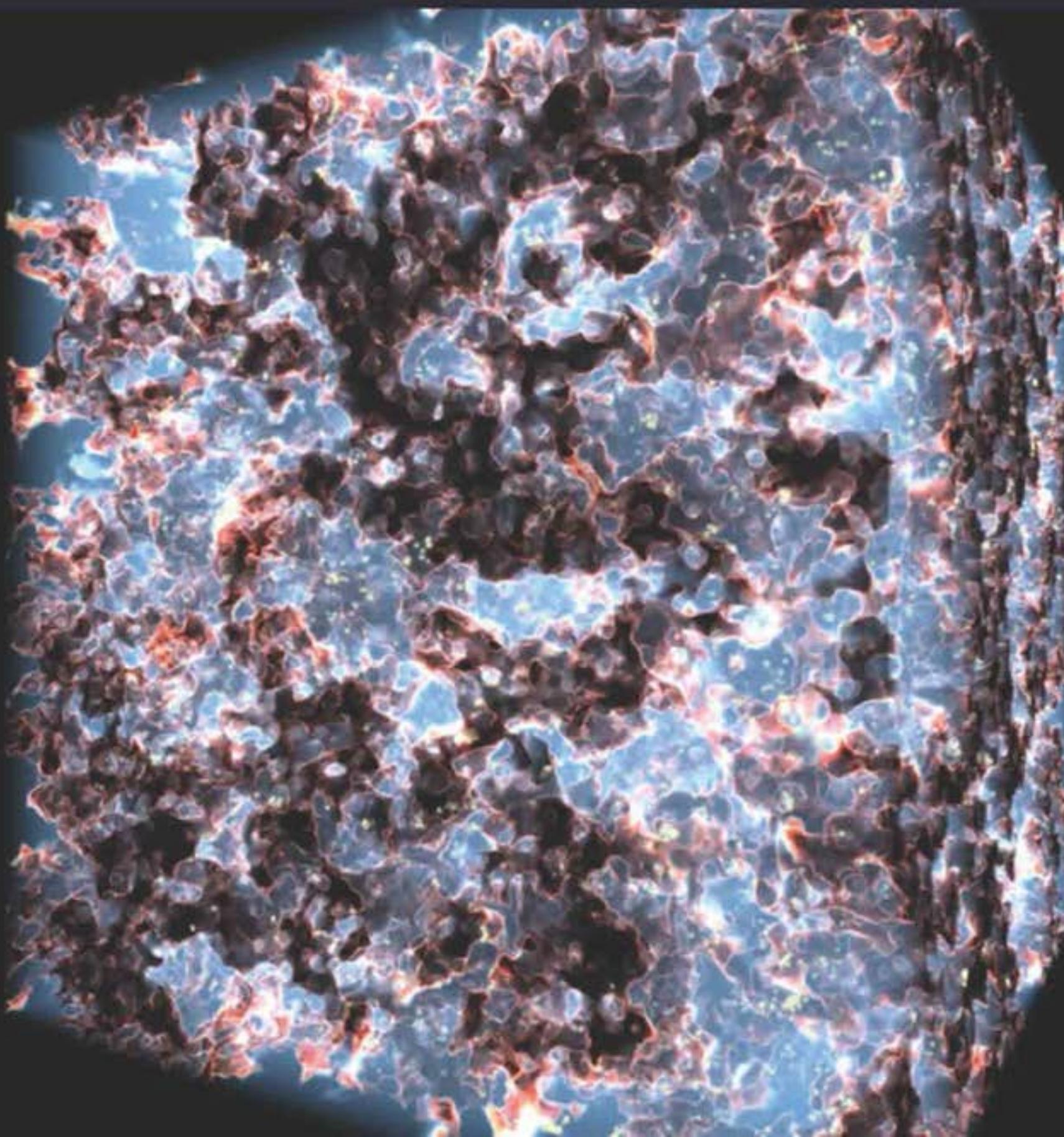


Average CMB spectral distortions

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Fluctuations of the y -parameter at large scales



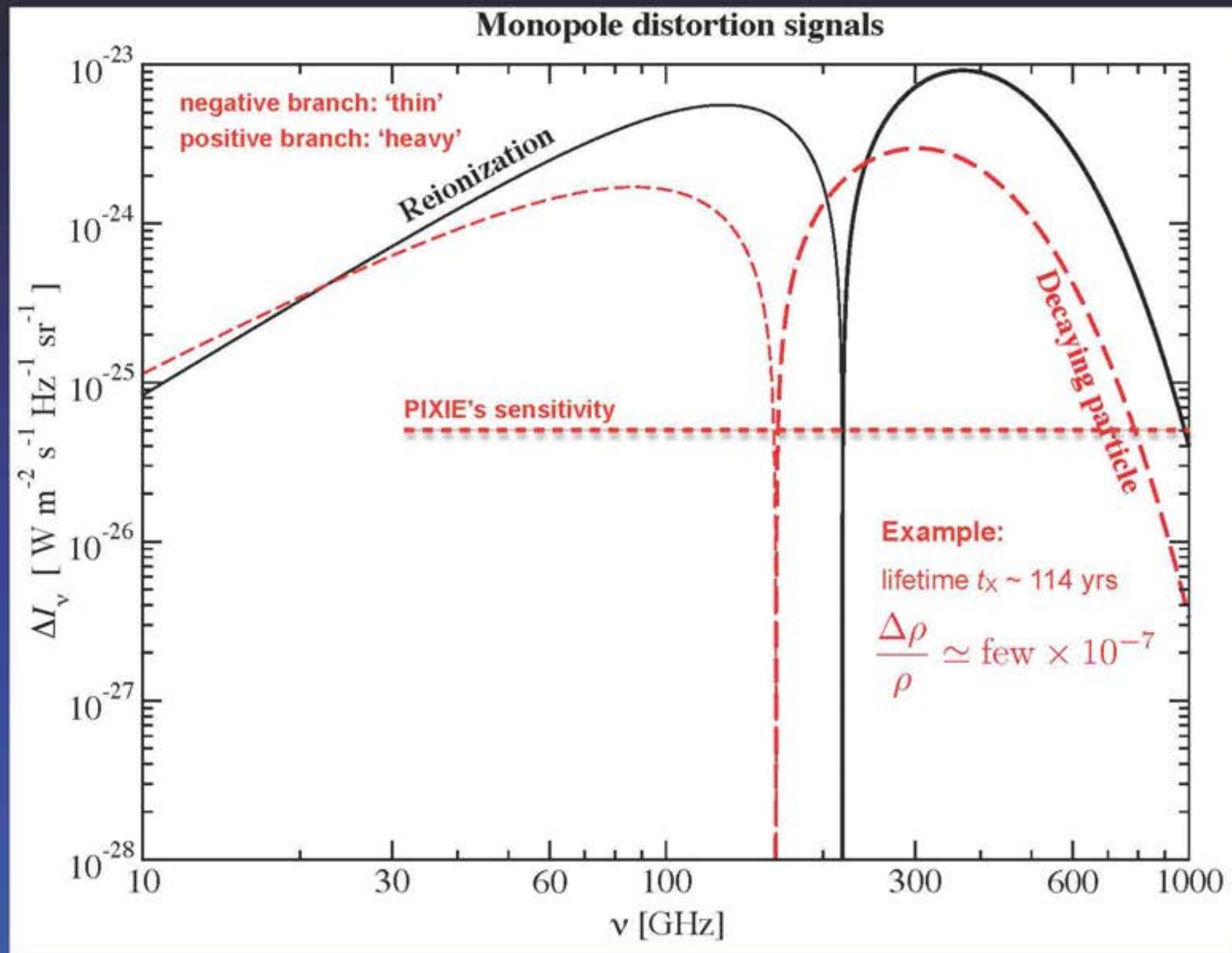
- spatial variations of the optical depth and temperature cause small-spatial variations of the y -parameter at different angular scales
- could tell us about the reionization sources and structure formation process
- additional independent piece of information!
- Cross-correlations with other signals

Example:
Simulation of reionization process
(1Gpc/h) by Alvarez & Abel

Decaying particles

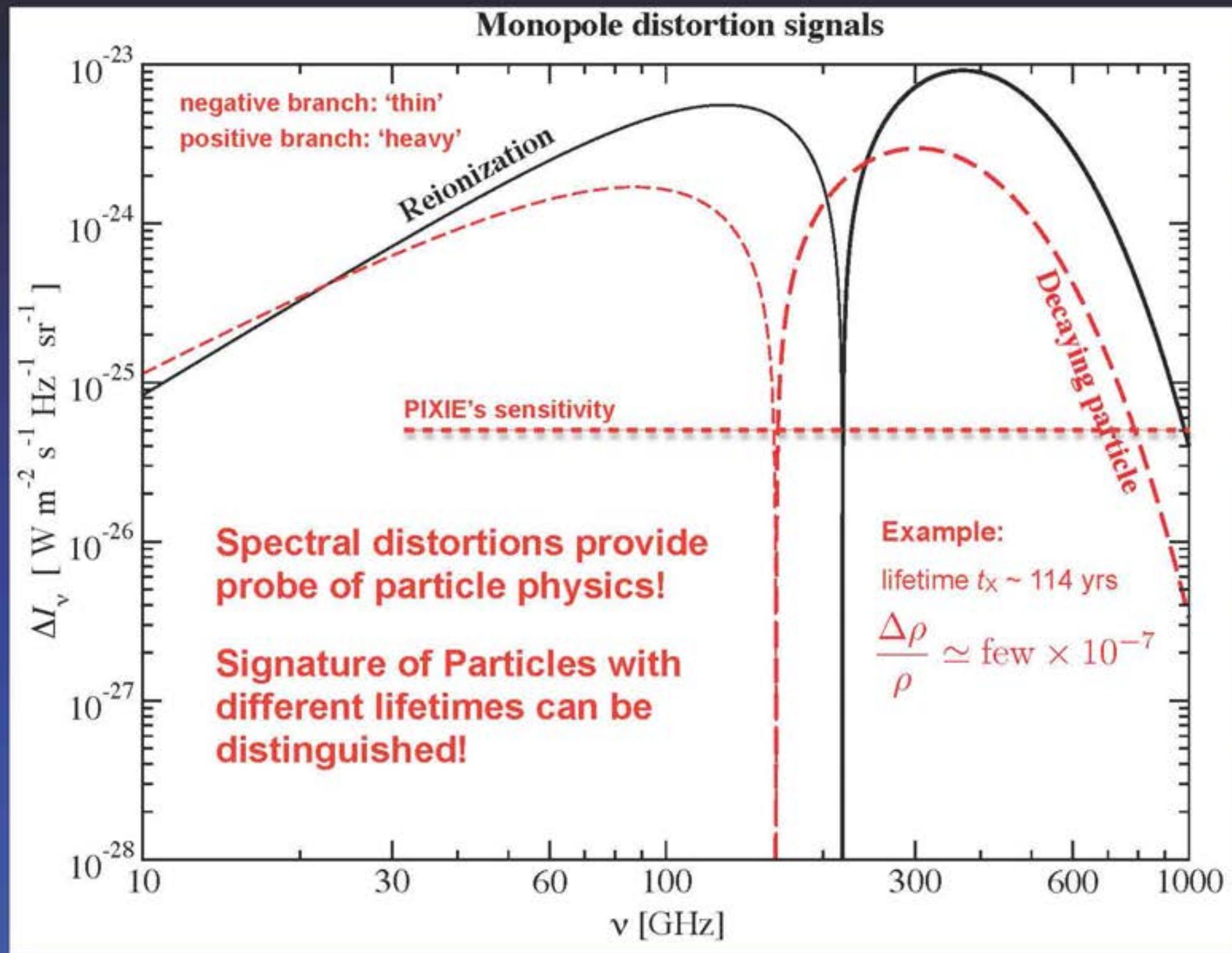
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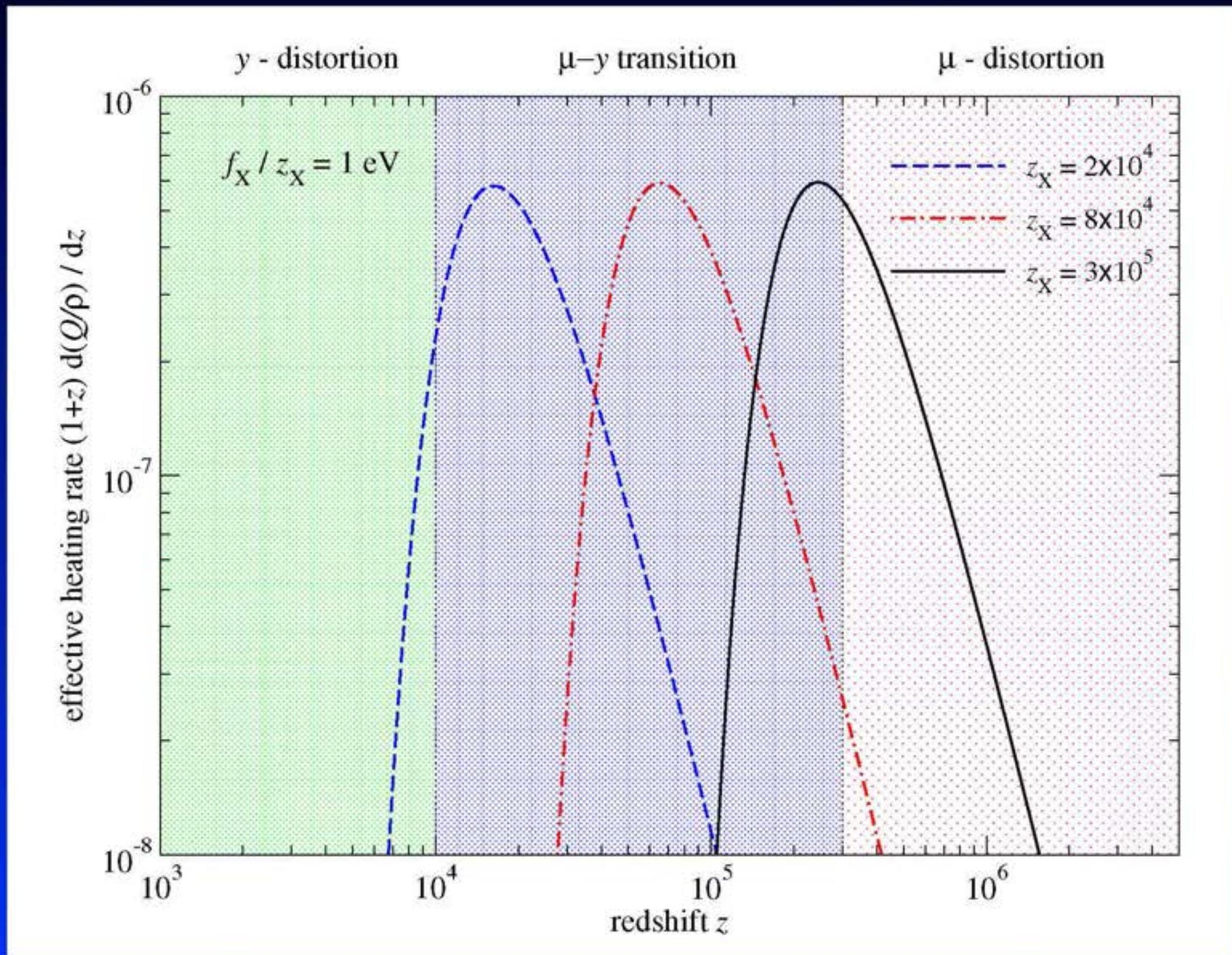


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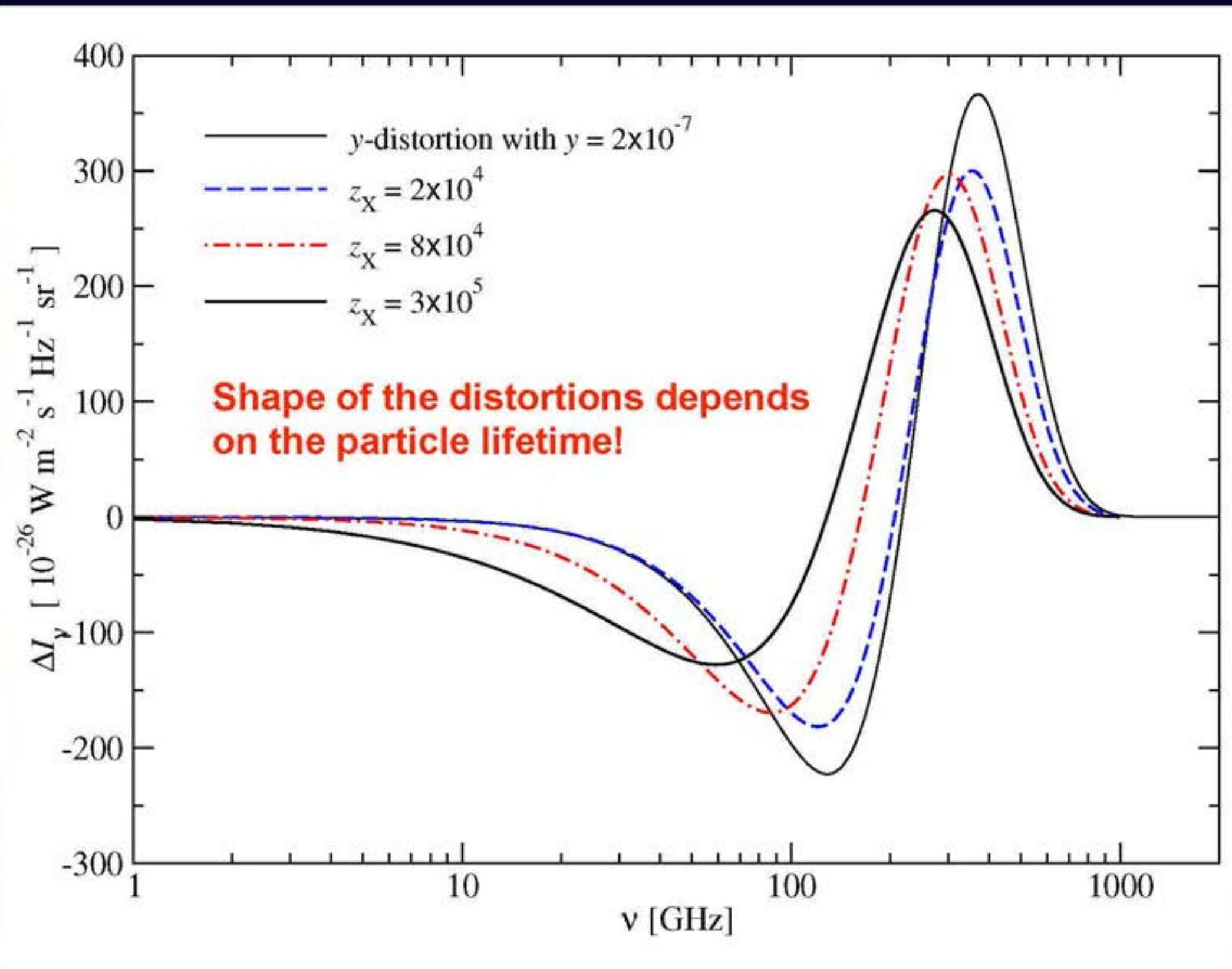
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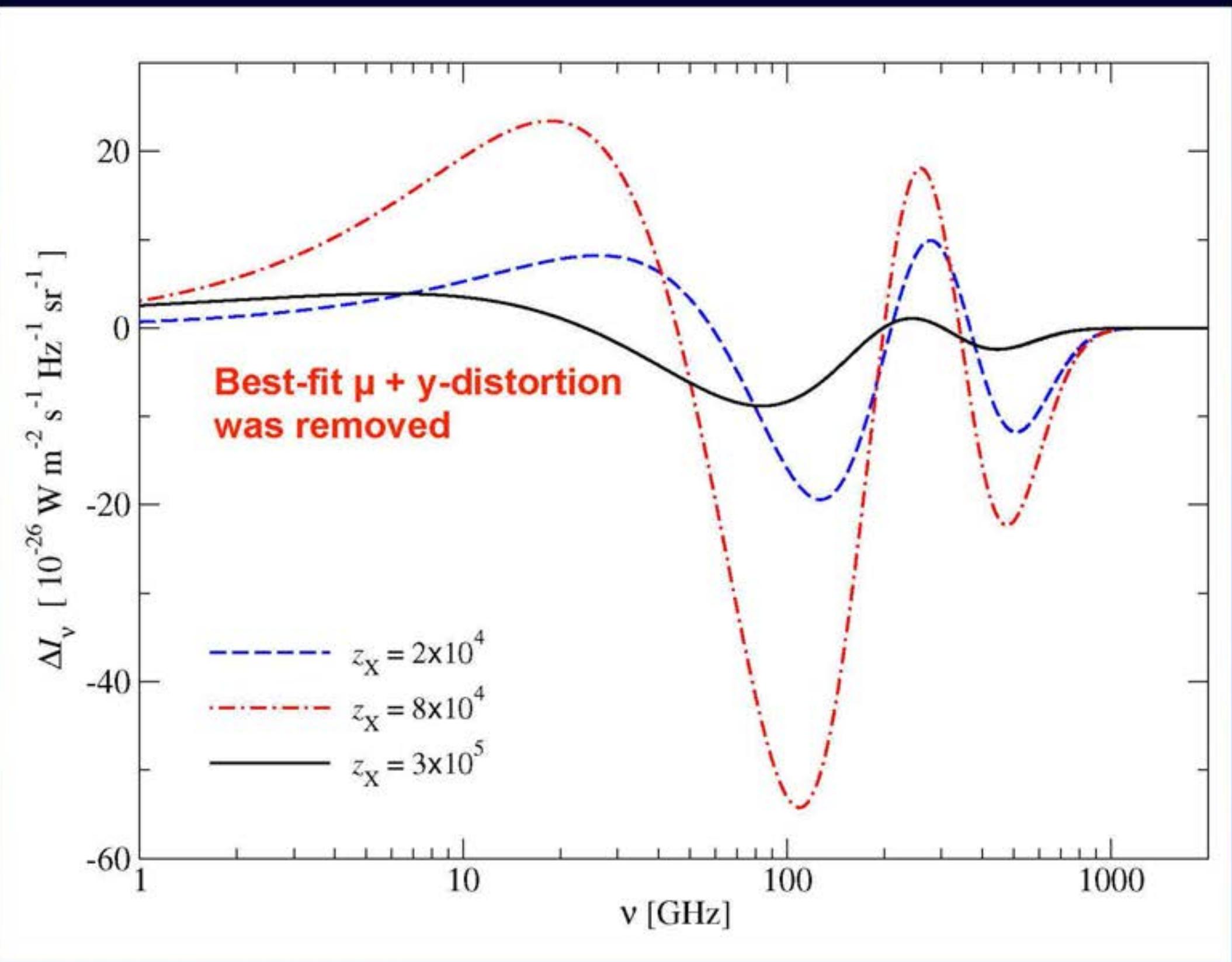
Decaying particle scenarios



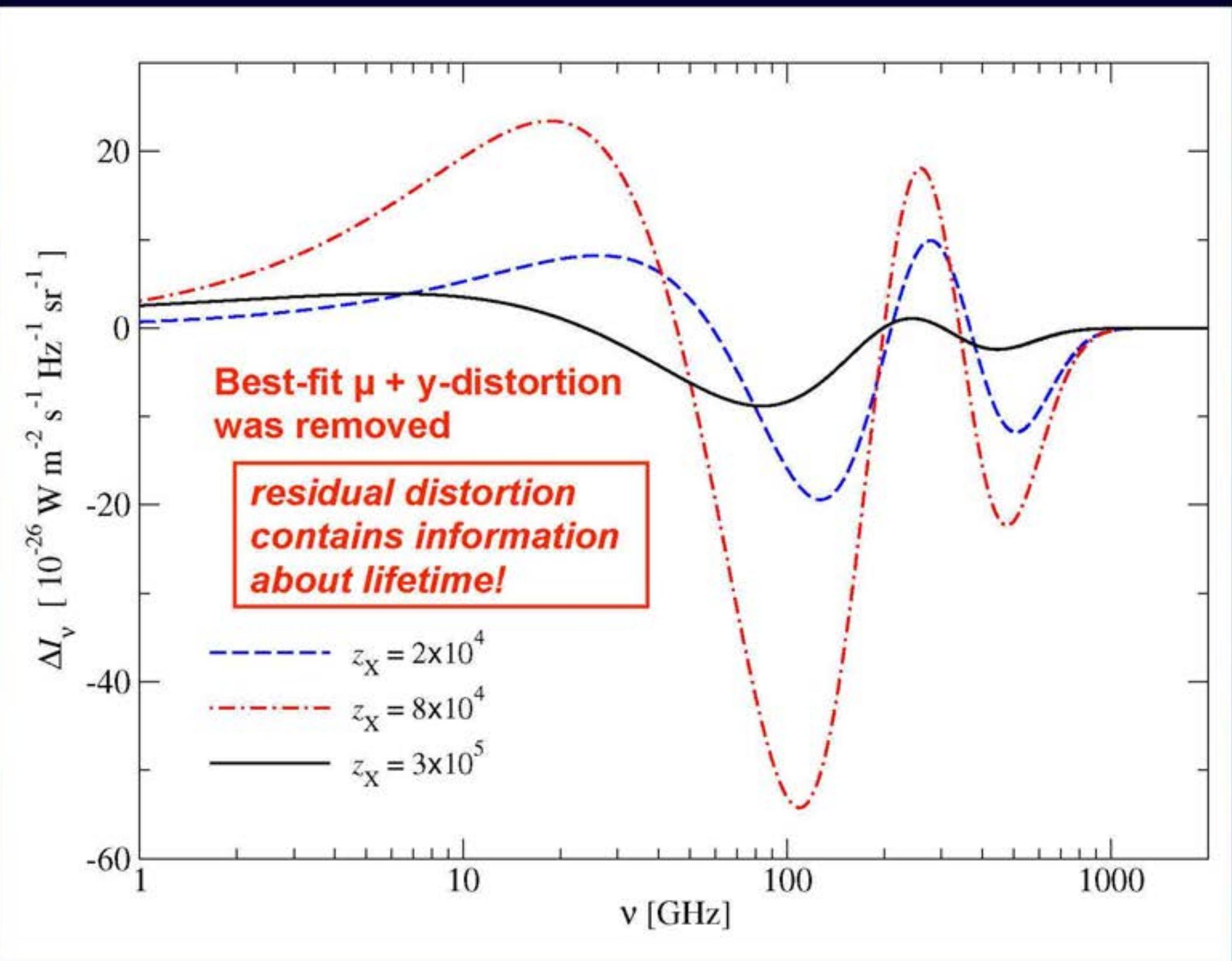
Decaying particle scenarios



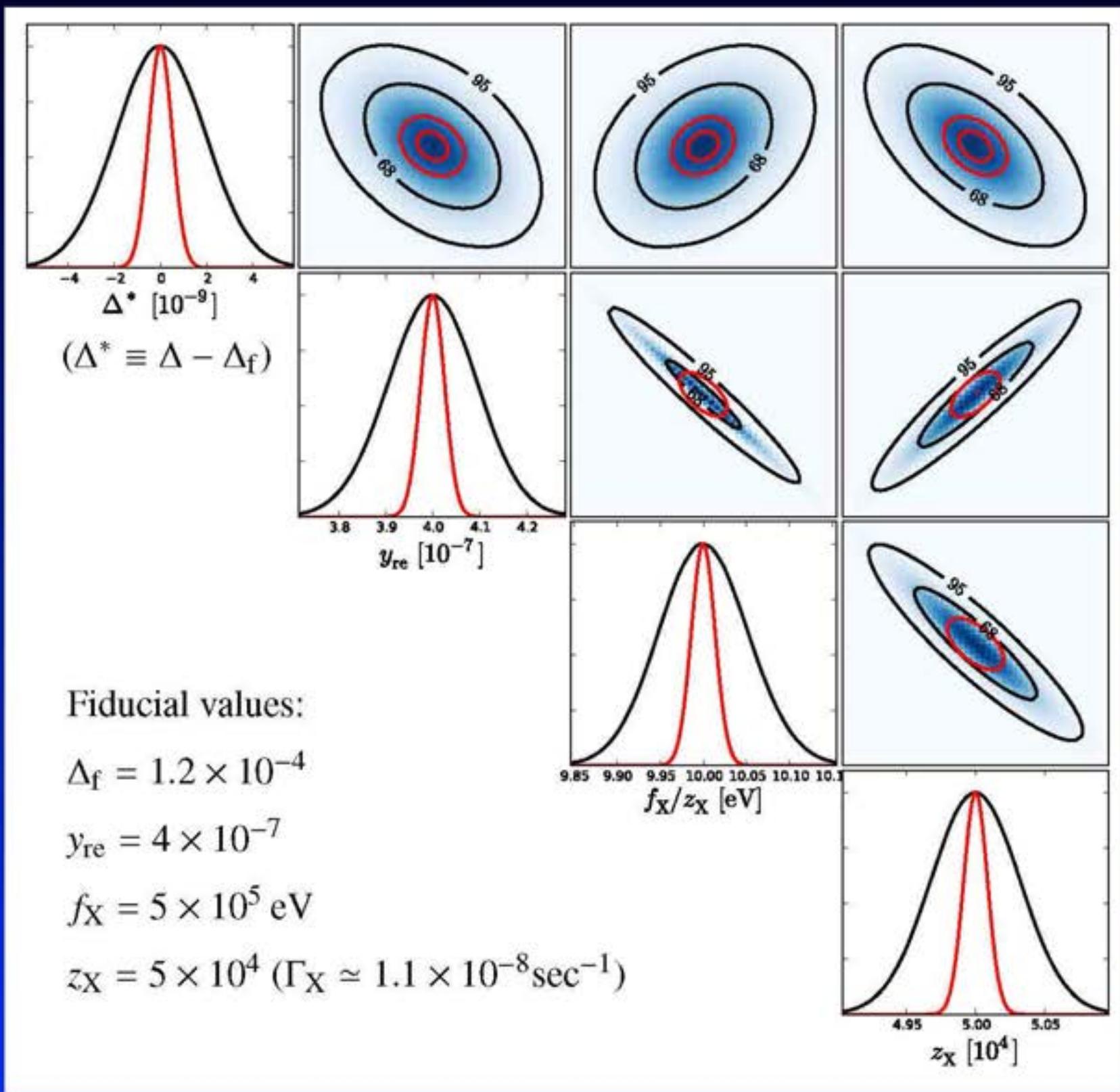
Decaying particle scenarios (information in residual)



Decaying particle scenarios (information in residual)



Decaying particle scenarios



Fiducial values:

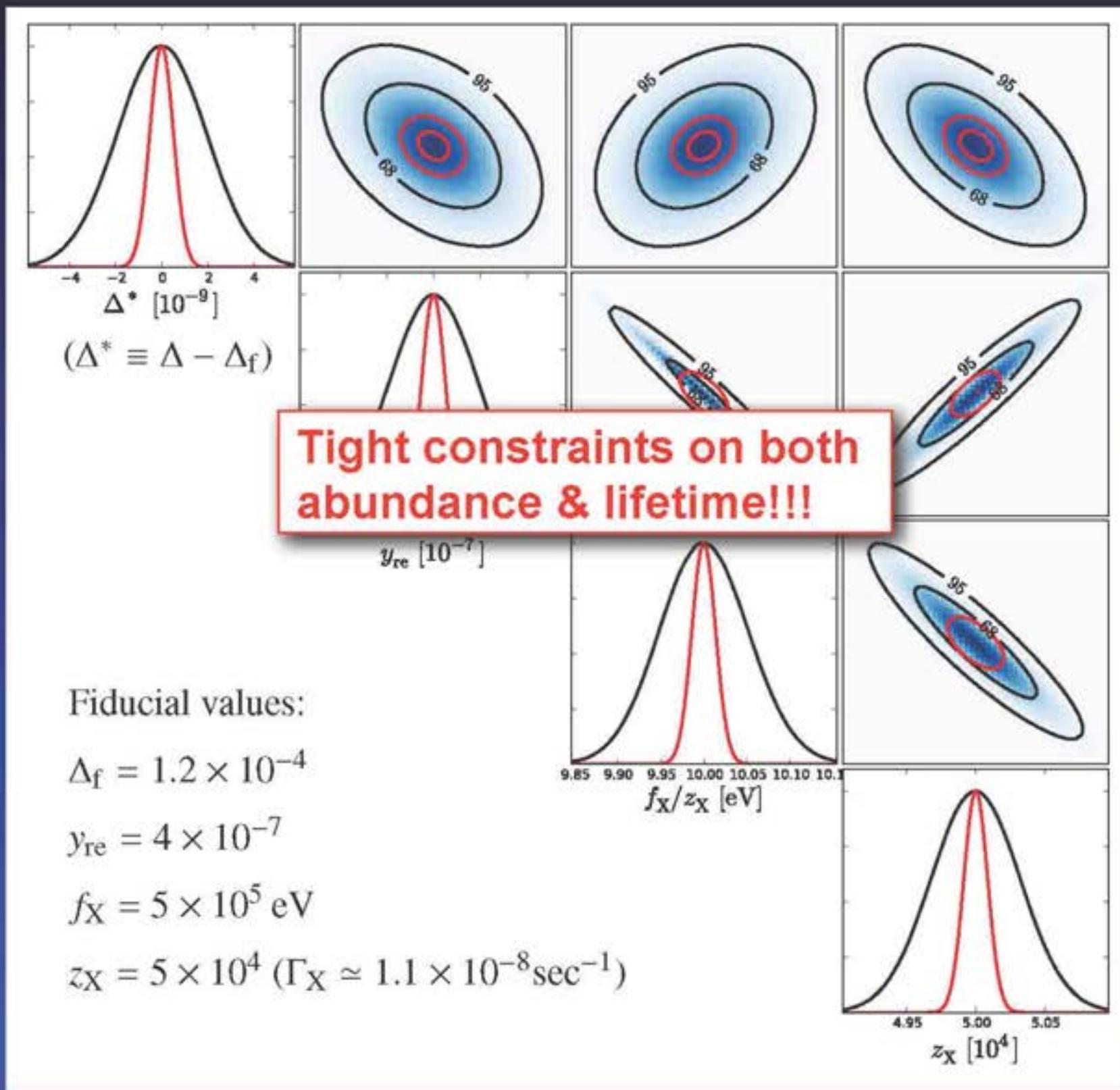
$$\Delta_f = 1.2 \times 10^{-4}$$

$$y_{re} = 4 \times 10^{-7}$$

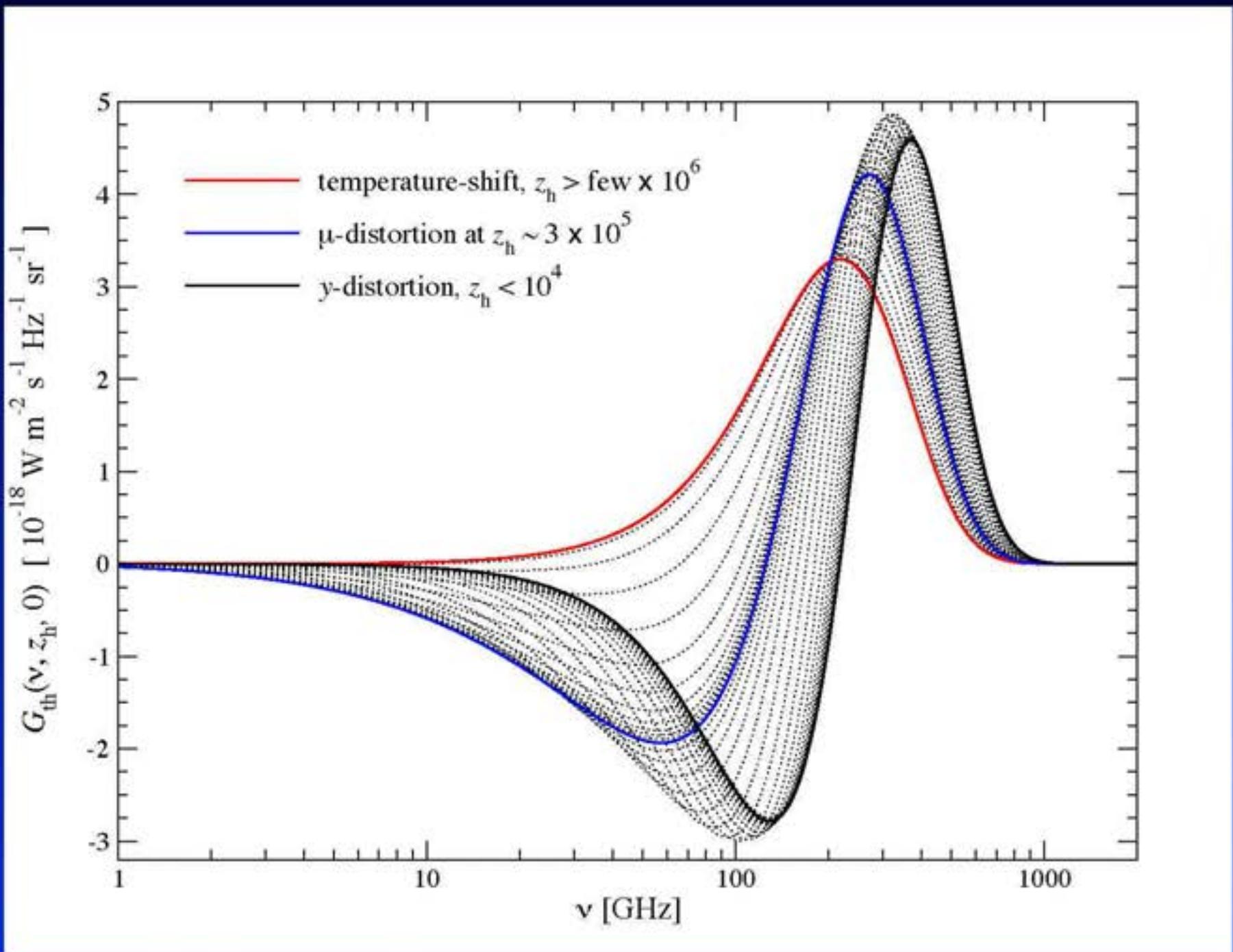
$$f_X = 5 \times 10^5 \text{ eV}$$

$$z_X = 5 \times 10^4 \quad (\Gamma_X \simeq 1.1 \times 10^{-8} \text{ sec}^{-1})$$

Decaying particle scenarios

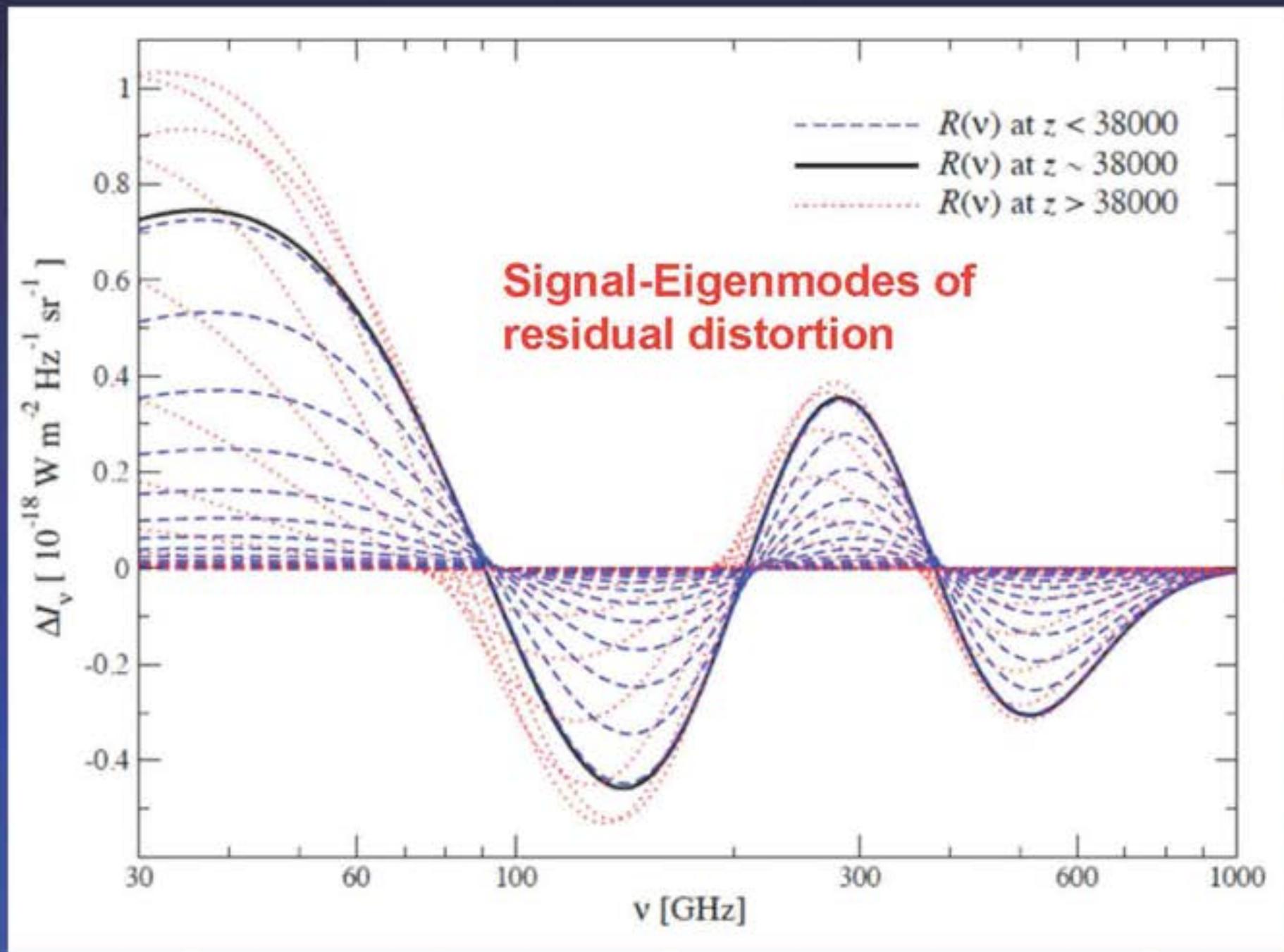


Using signal eigenmodes to compress the distortion data



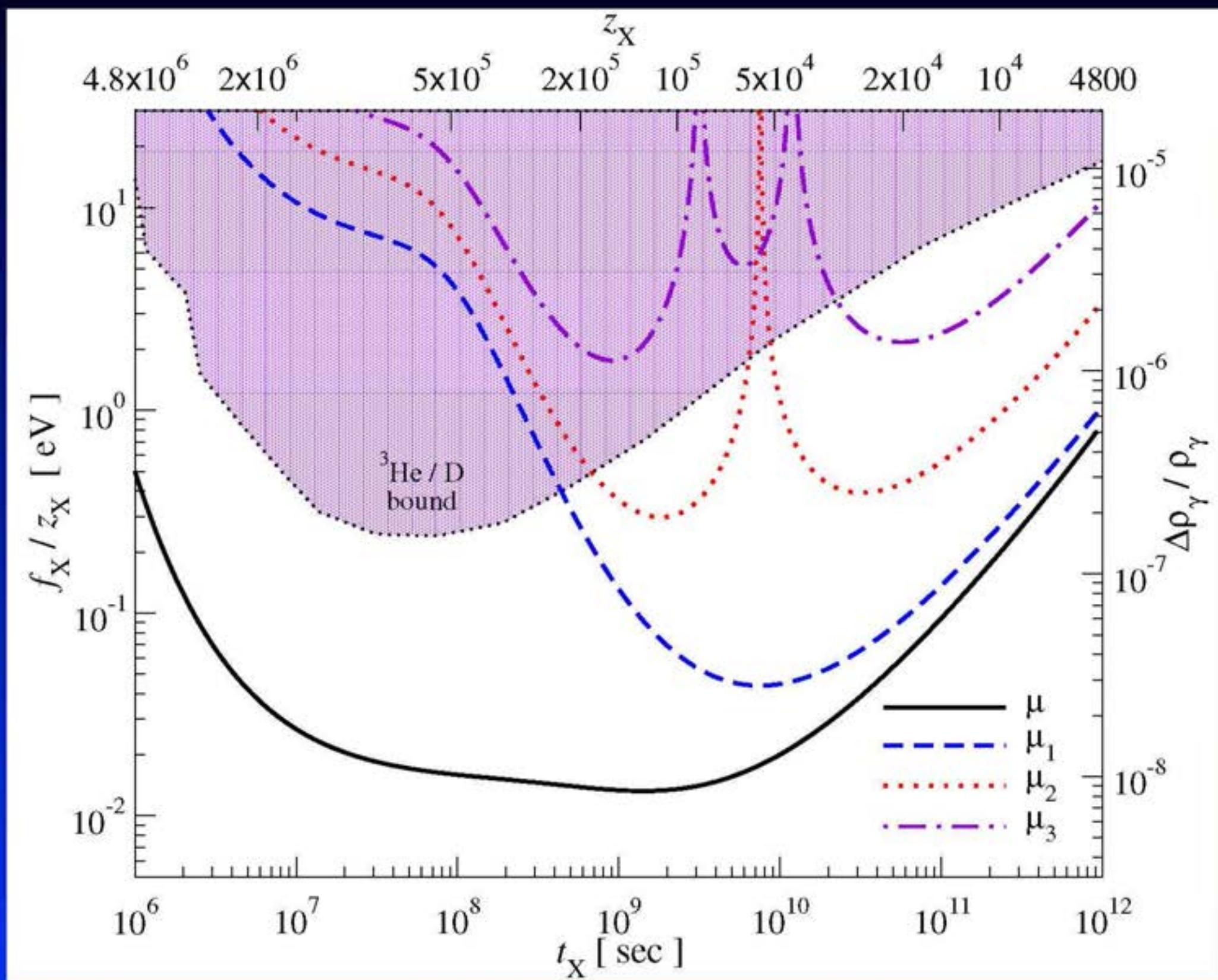
- Principle component decomposition of the distortion signal
- compression of the useful information given instrumental settings

Using signal eigenmodes to compress the distortion data

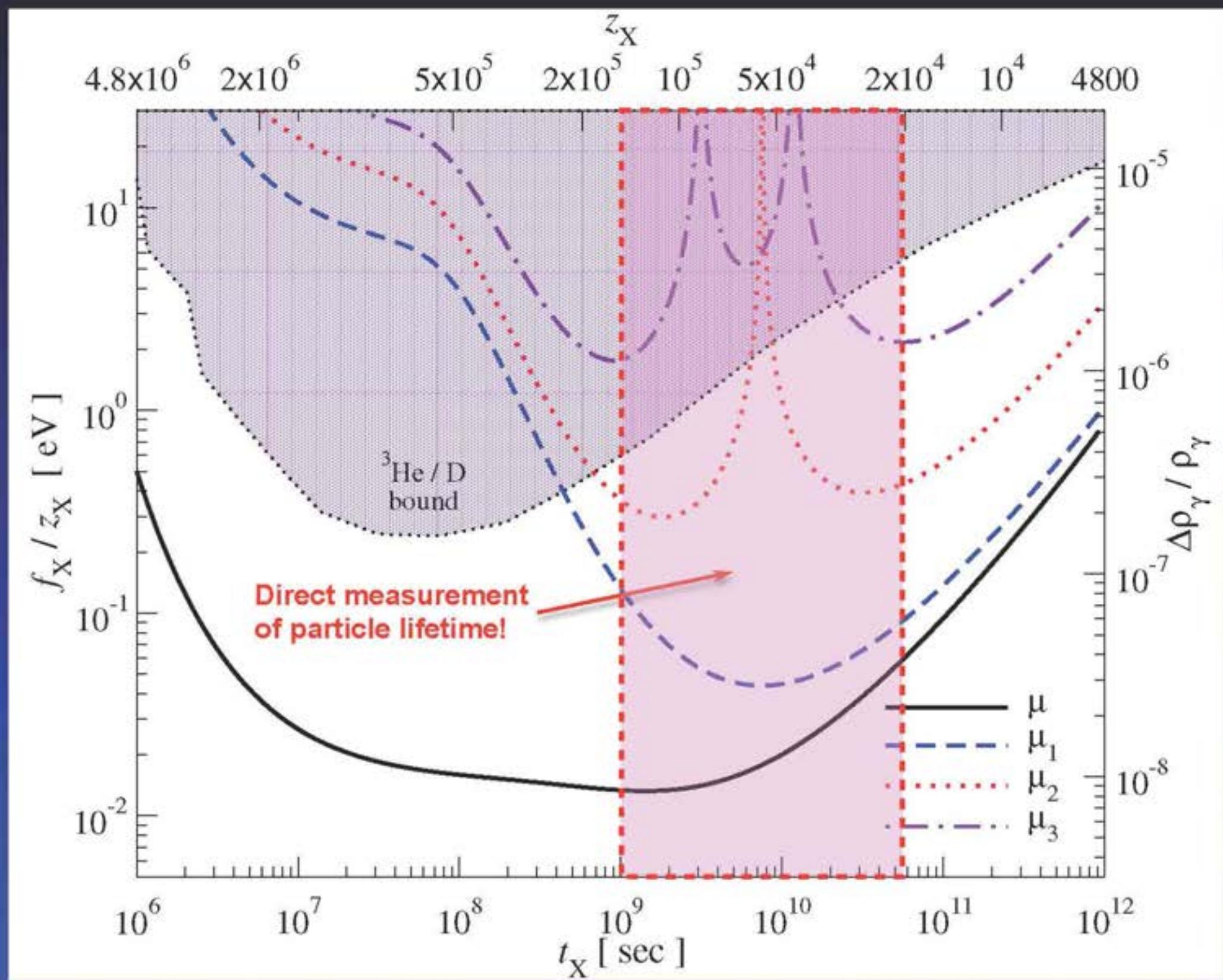


- Principle component decomposition of the distortion signal
- compression of the useful information given instrumental settings
- new set of observables
 $p=\{y, \mu, \mu_1, \mu_2, \dots\}$
- model-comparison + forecasts of errors very simple!

Decaying particle 1σ -detection limits for PIXIE

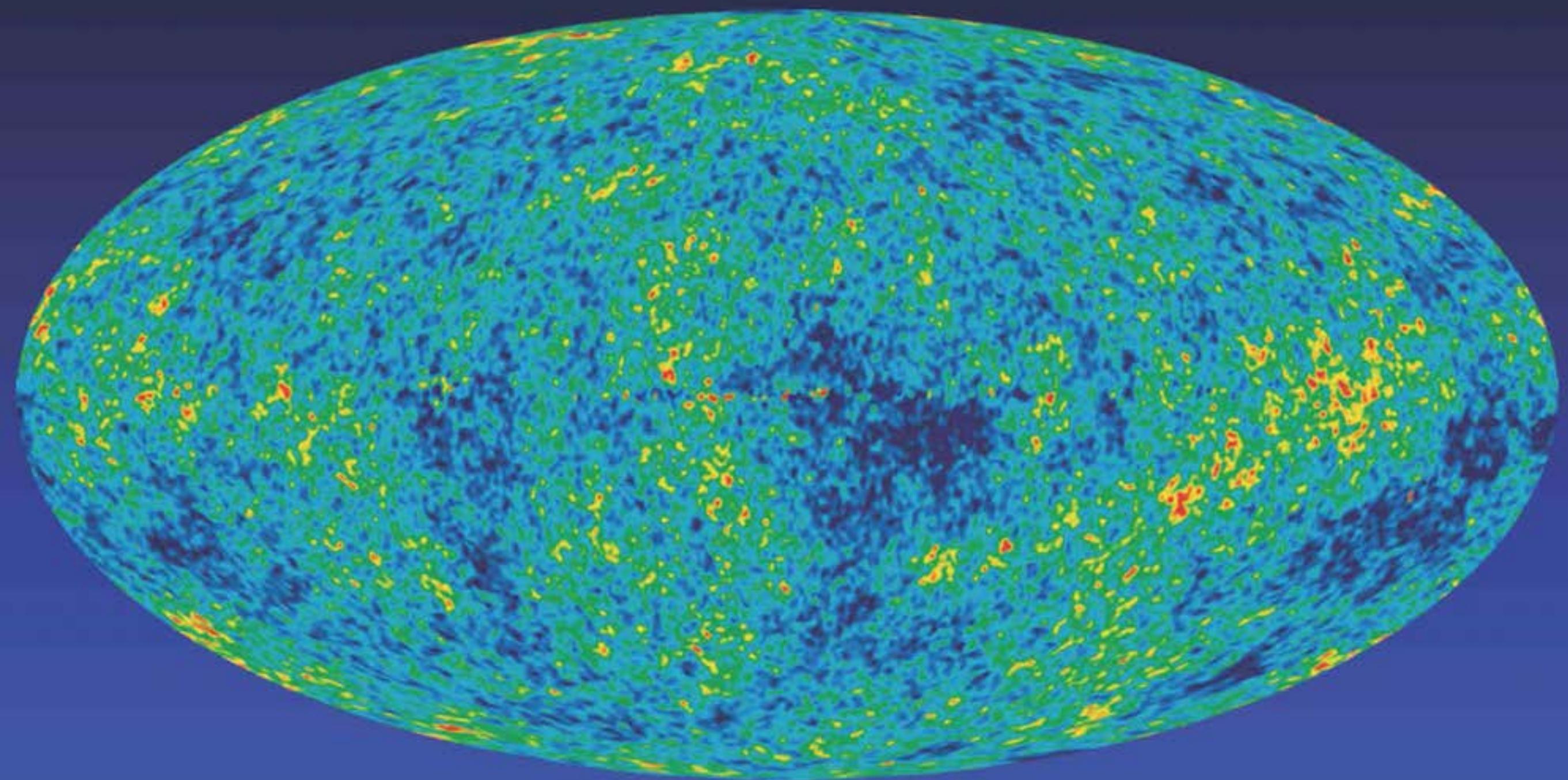


Decaying particle 1σ -detection limits for PIXIE

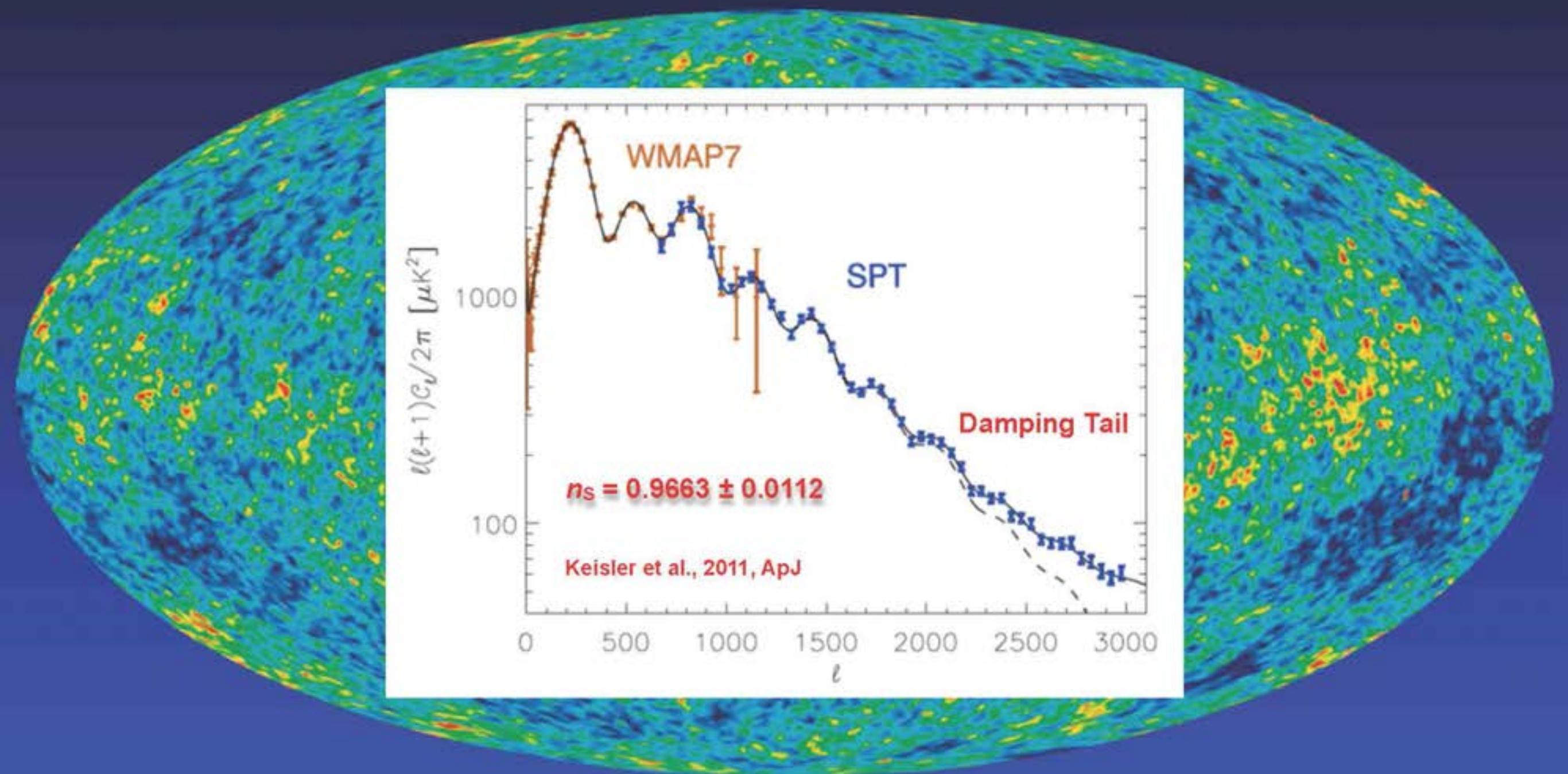


The dissipation of small-scale acoustic modes

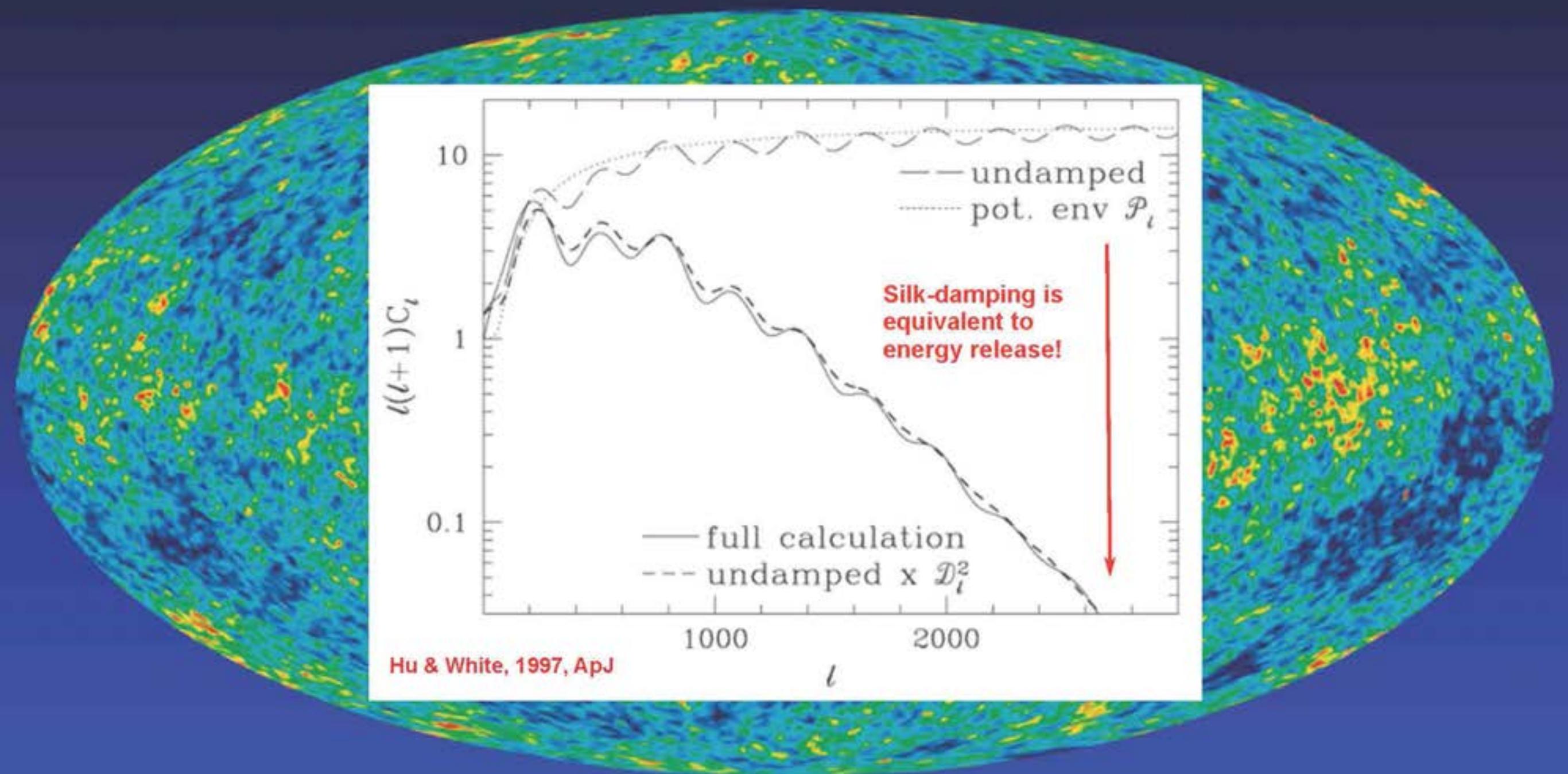
Dissipation of small-scale acoustic modes



Dissipation of small-scale acoustic modes



Dissipation of small-scale acoustic modes



Energy release caused by dissipation process

‘Obvious’ dependencies:

- *Amplitude* of the small-scale power spectrum
- *Shape* of the small-scale power spectrum
- *Dissipation scale* $\rightarrow k_D \sim (H_0 \Omega_{\text{rel}}^{1/2} N_{e,0})^{1/2} (1+z)^{3/2}$ at early times

not so ‘obvious’ dependencies:

- *primordial non-Gaussianity* in the squeezed limit
(Pajer & Zaldarriaga, 2012; Ganc & Komatsu, 2012)
- *Type of the perturbations* (adiabatic \leftrightarrow isocurvature)
(Barrow & Coles, 1991; Hu et al., 1994; Dent et al, 2012, JC & Grin, 2012)
- *Neutrinos* (or any extra relativistic degree of freedom)

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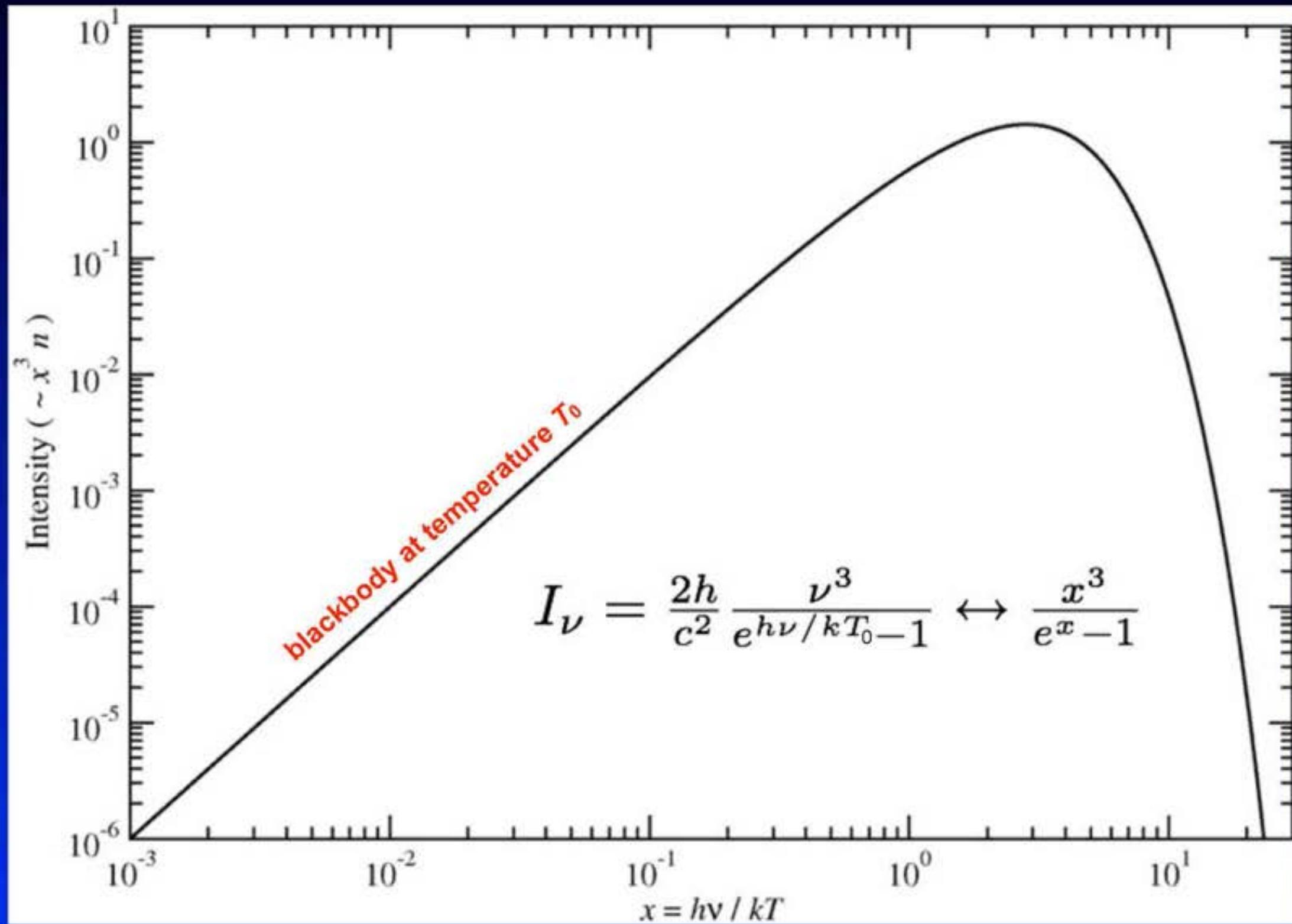
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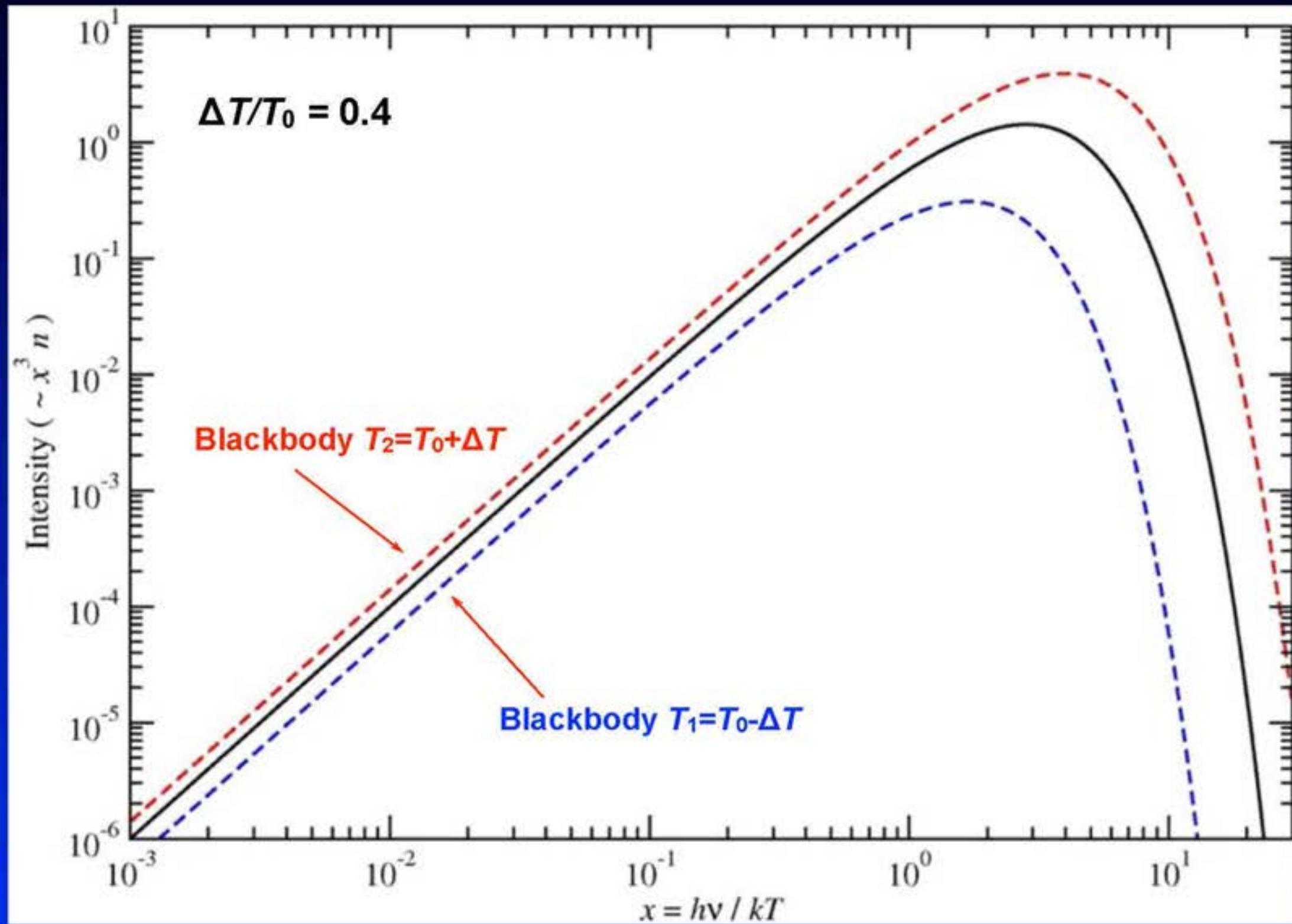
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CMB Spectral distortions provide probe of Inflation physics!!!

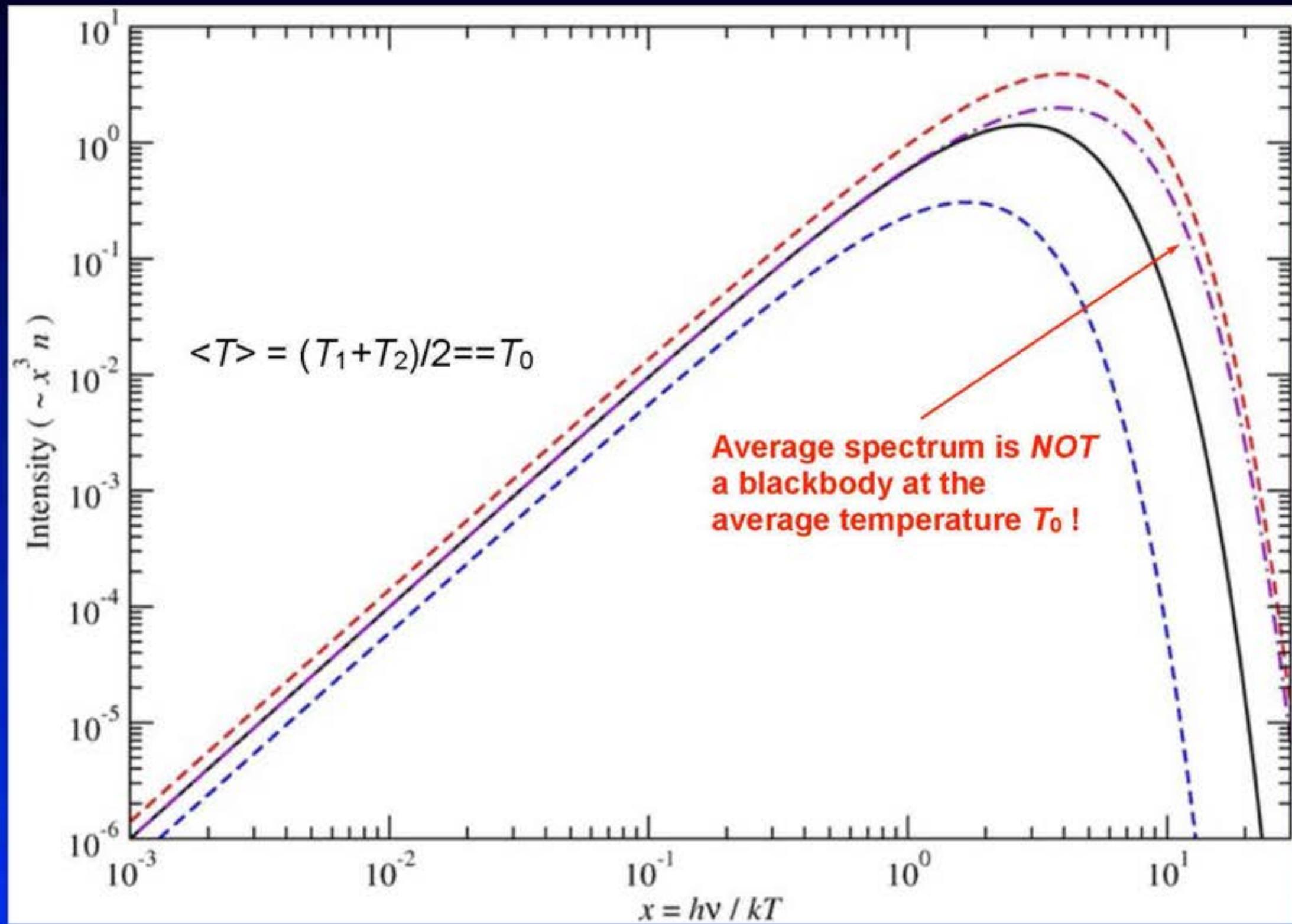
Superpositions of blackbody spectra



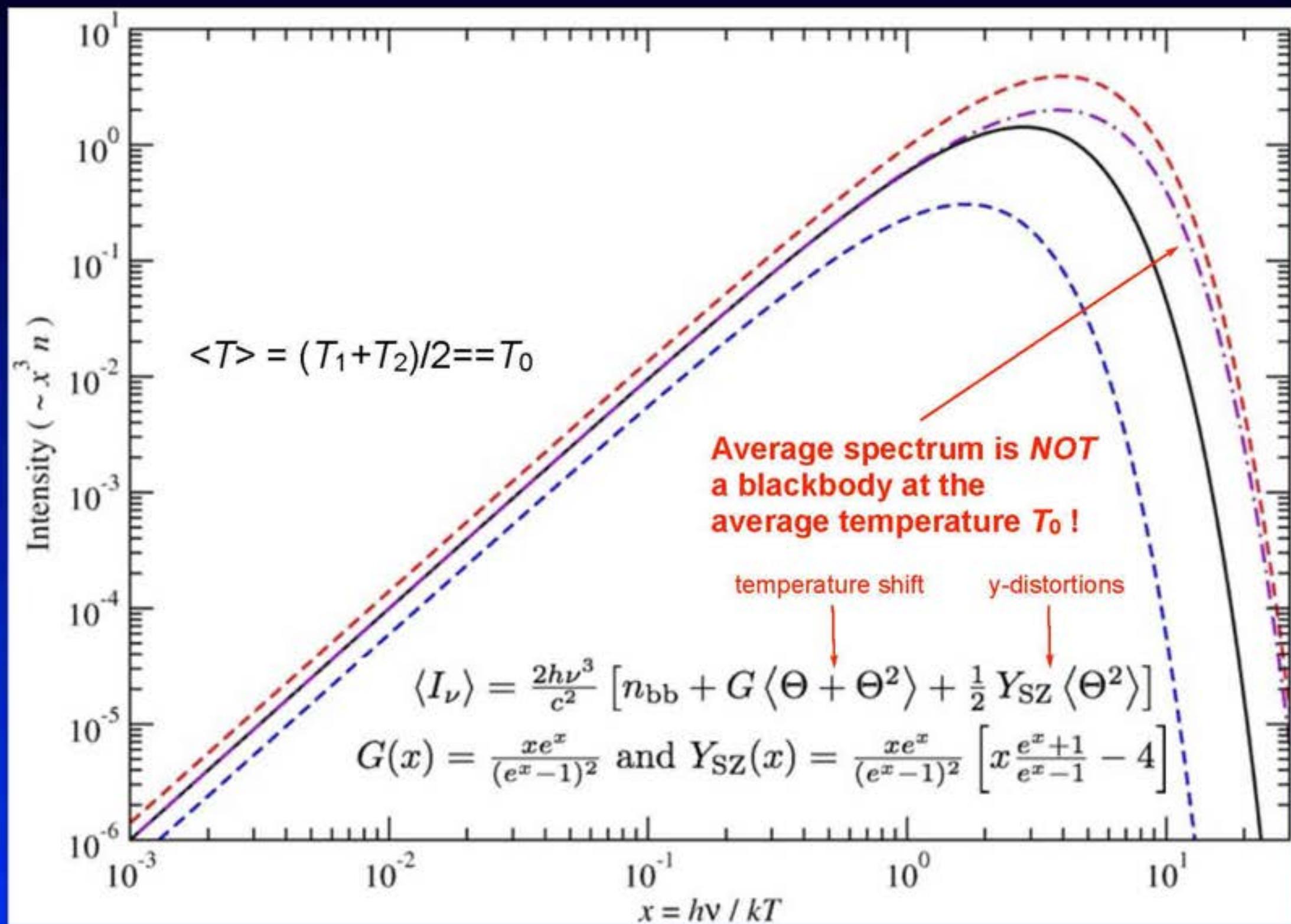
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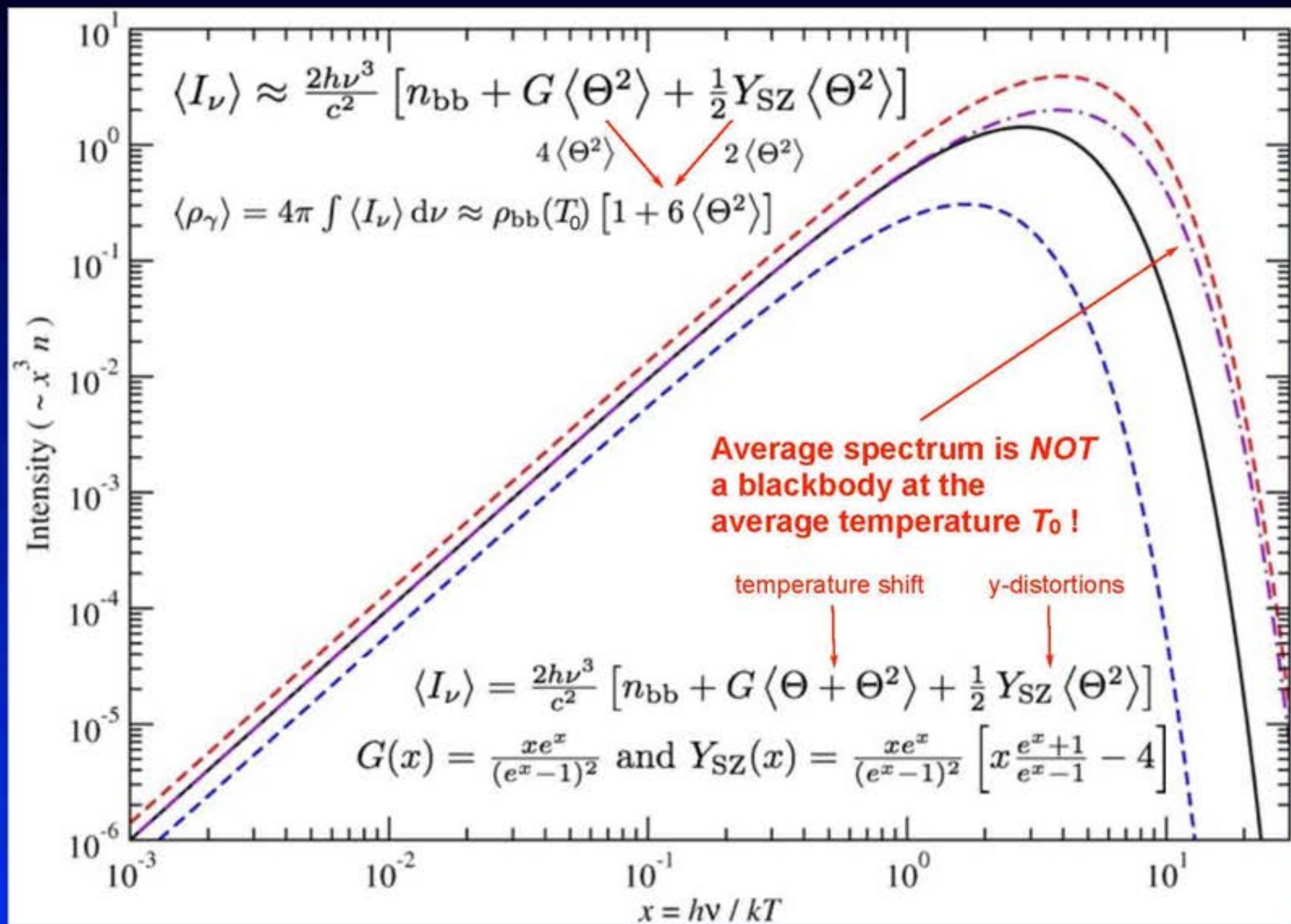
Superpositions of blackbody spectra



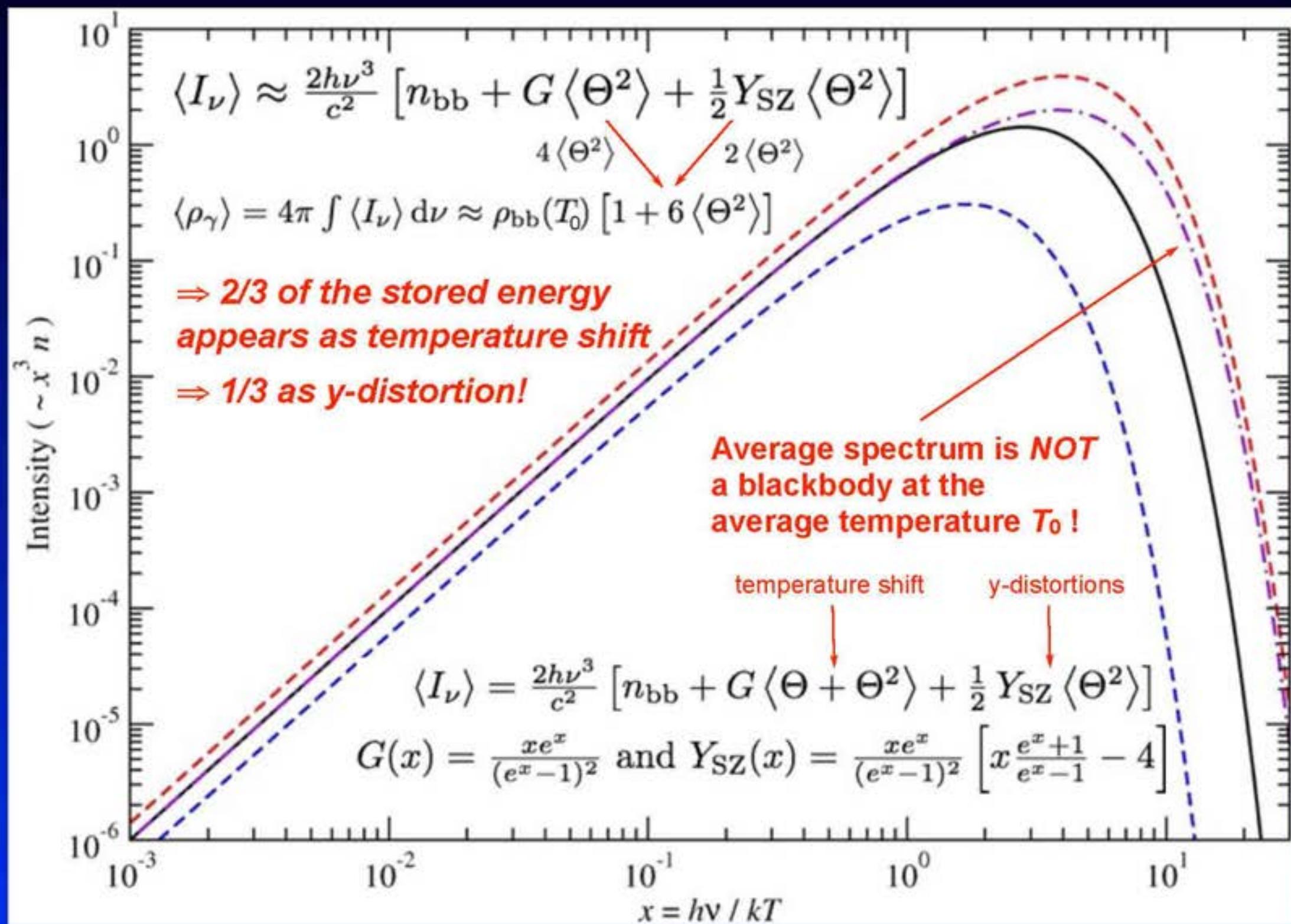
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Superpositions of blackbody spectra



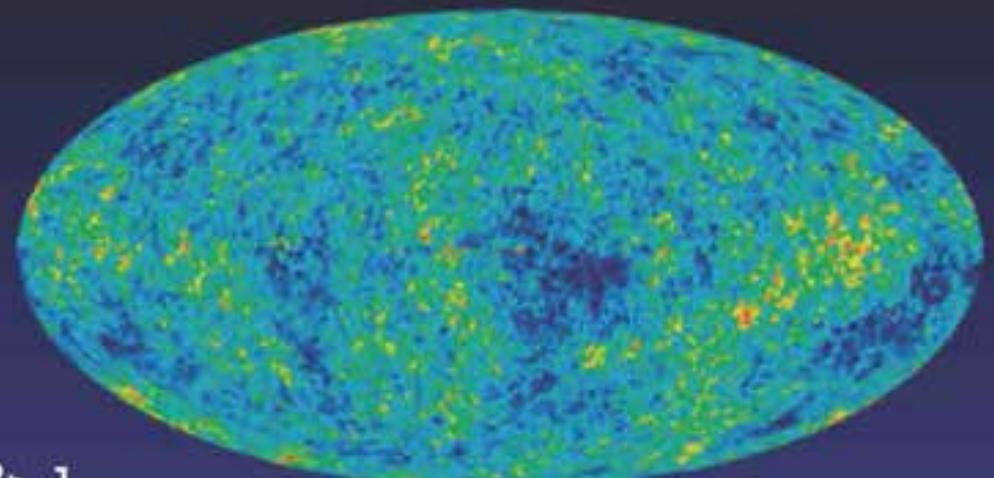
Superpositions of blackbody spectra



Dissipation of acoustic modes: ‘microscopic picture’

- after inflation: photon field has spatially varying temperature T
- average energy stored in photon field at any given moment

$$\langle \rho_Y \rangle = a_R \langle T^4 \rangle \approx a_R \langle T \rangle^4 [1 + 4\langle \Theta \rangle + 6\langle \Theta^2 \rangle] \\ == 0$$

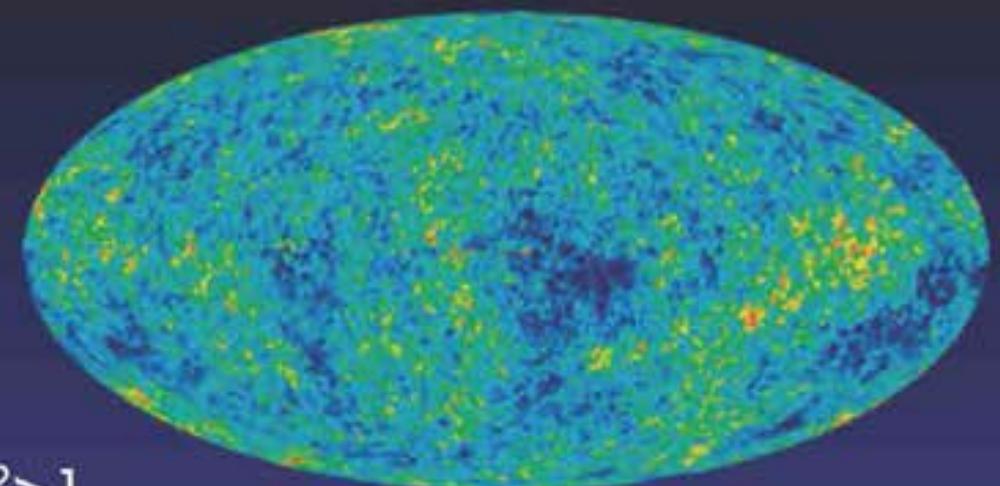


E.g., our snapshot at $z=0$

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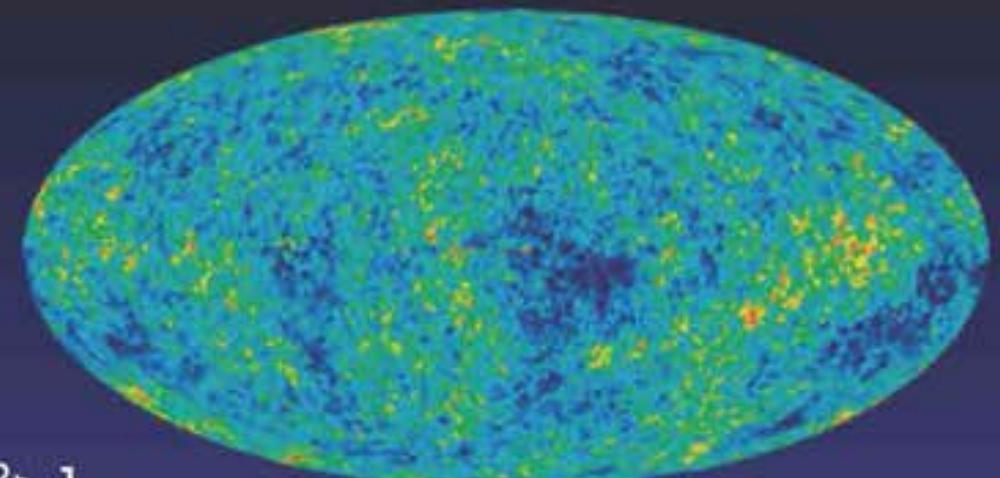
$$\Rightarrow (a^4 \rho_Y)^{-1} da^4 Q_{ac} / dt = -6 d\langle \Theta^2 \rangle / dt$$

- Monopole actually **drops** out of the equation!
- In principle ***all*** higher multipoles contribute to the energy release

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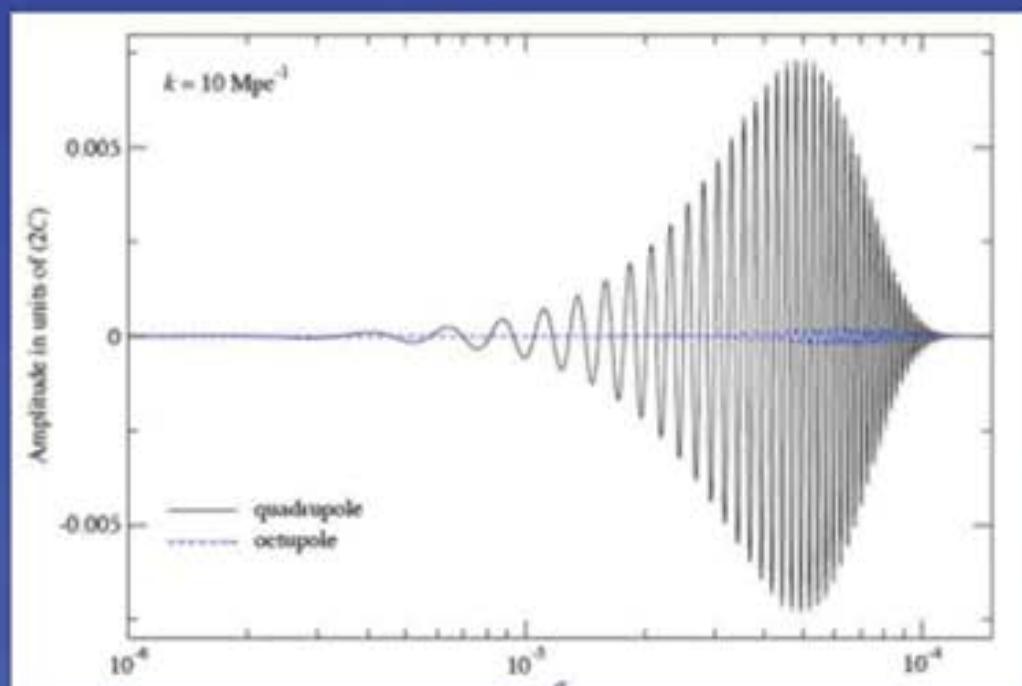
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- Monopole actually **drops** out of the equation!
- In principle ***all*** higher multipoles contribute to the energy release
- At high redshifts ($z \geq 10^4$):
 - ▶ *net (gauge-invariant) dipole and contributions from higher multipoles are negligible*
 - ▶ *dominant term caused by quadrupole anisotropy*

$$\Rightarrow (a^4 \rho_Y)^{-1} da^4 Q_{ac} / dt \approx -12 d\langle \Theta_0^2 \rangle / dt$$

↑
9/4 larger than classical estimate



Effective energy release caused by damping effect

- Effective heating rate from full 2x2 Boltzmann treatment (JC, Khatri & Sunyaev, 2012)

$$\frac{1}{a^4 \rho_\gamma} \frac{da^4 Q_{ac}}{dt} = 4\sigma_T N_e c \left\langle \frac{(3\Theta_1 - \beta)^2}{3} + \frac{9}{2}\Theta_2^2 - \frac{1}{2}\Theta_2(\Theta_0^P + \Theta_2^P) + \sum_{l \geq 3} (2l+1)\Theta_\ell^2 \right\rangle$$

$\Theta_\ell = \frac{1}{2} \int \Theta(\mu) P_\ell(\mu) d\mu$

↑ gauge-independent dipole ↑ effect of polarization ↑ higher multipoles

$$\langle XY \rangle = \int \frac{k^2 dk}{2\pi^2} P(k) X(k) Y(k)$$

↑ Primordial power spectrum

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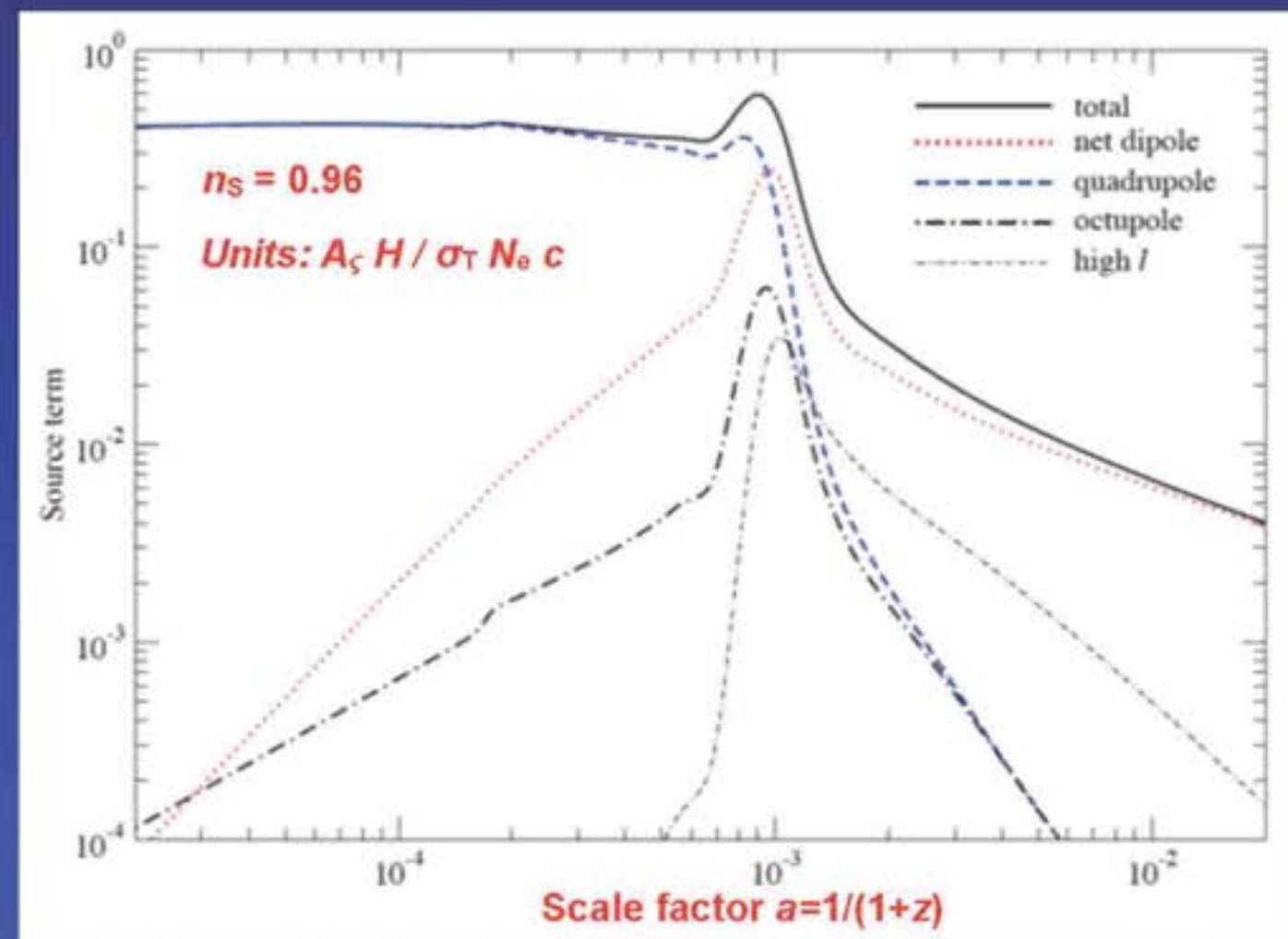
$$\Theta_\ell = \frac{1}{2} \int \Theta(\mu) P_\ell(\mu) d\mu$$

gauge-independent dipole **effect of polarization** **higher multipoles**

$$\langle XY \rangle = \int \frac{k^2 dk}{2\pi^2} P(k) X(k) Y(k)$$

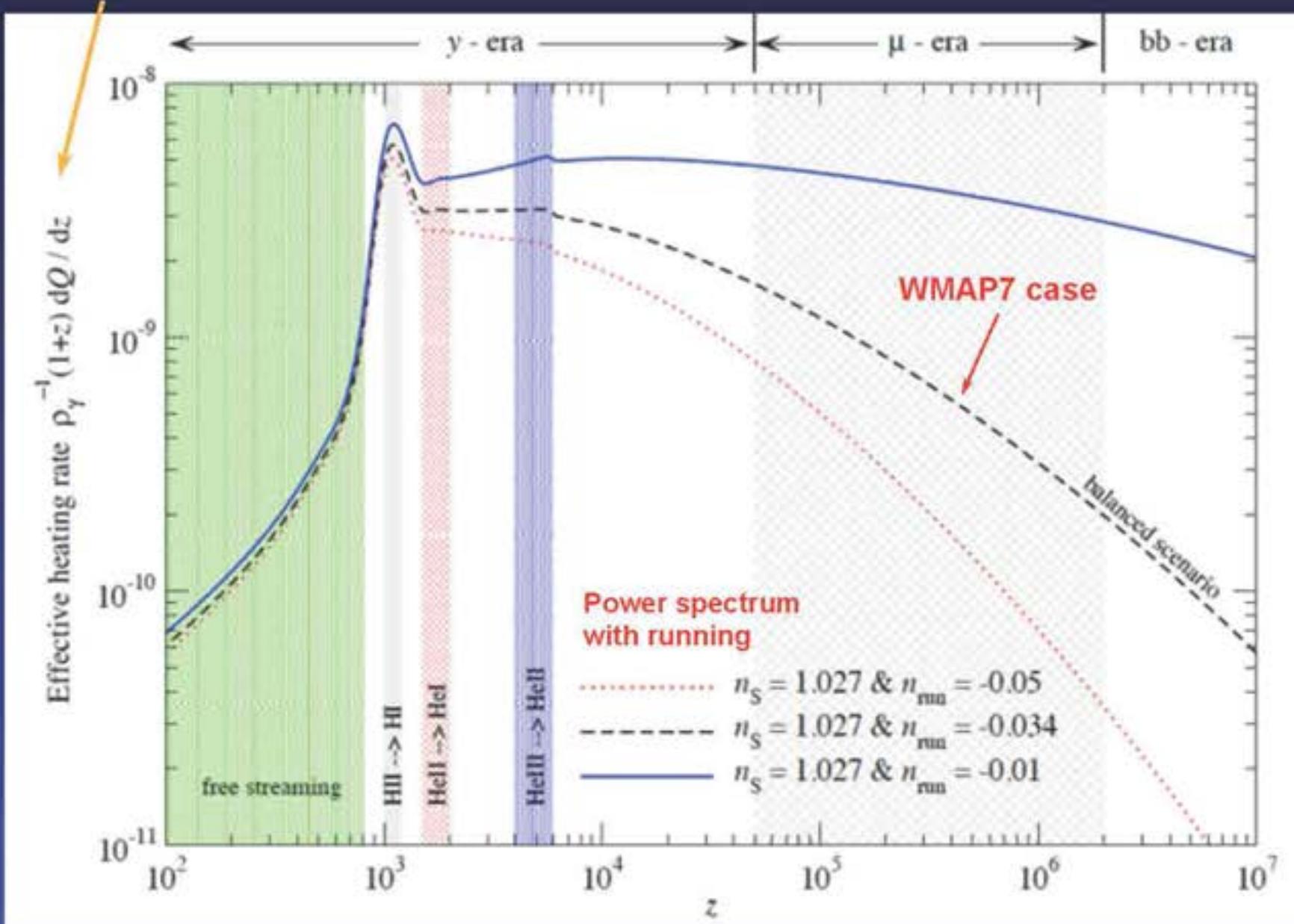
Primordial power spectrum

- quadrupole dominant at high z
- net dipole important only at low redshifts
- polarization ~5% effect
- contribution from higher multipoles rather small



Our computation for the effective energy release

scaled such that constant for $n_s = 1$



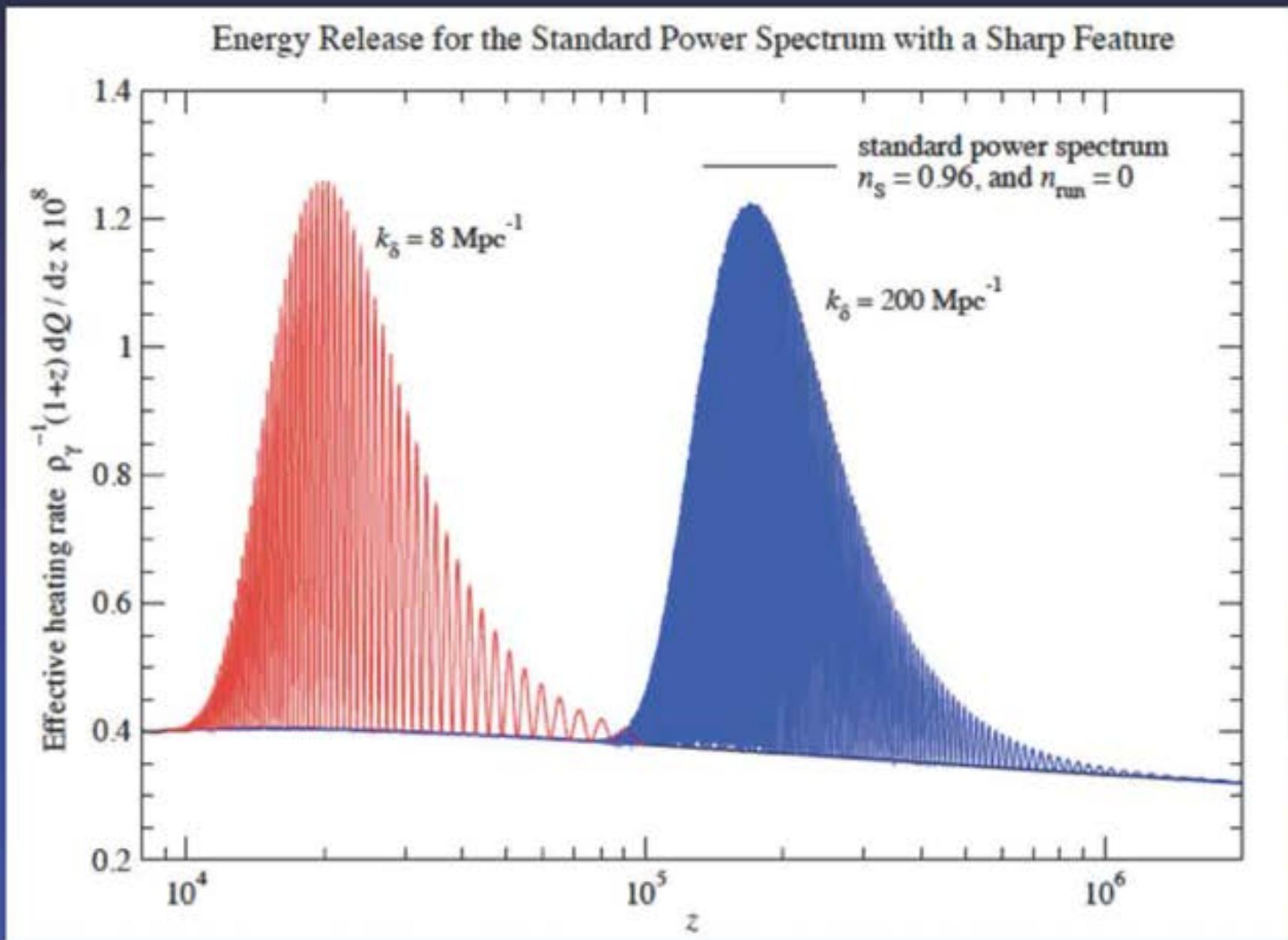
- Our 2. order perturbation calculation showed that the *classical* picture was slightly inconsistent
- Amplitude of the distortion depends on the small-scale power spectrum
- Computation carried out with **CosmoTherm** (JC & Sunyaev 2011)

JC, Khatri & Sunyaev, 2012

$$P_\zeta(k) = 2\pi^2 A_\zeta k^{-3} (k/k_0)^{n_S - 1 + \frac{1}{2} n_{run} \ln(k/k_0)}$$

Primordial power spectrum of curvature perturbations is input for the calculation

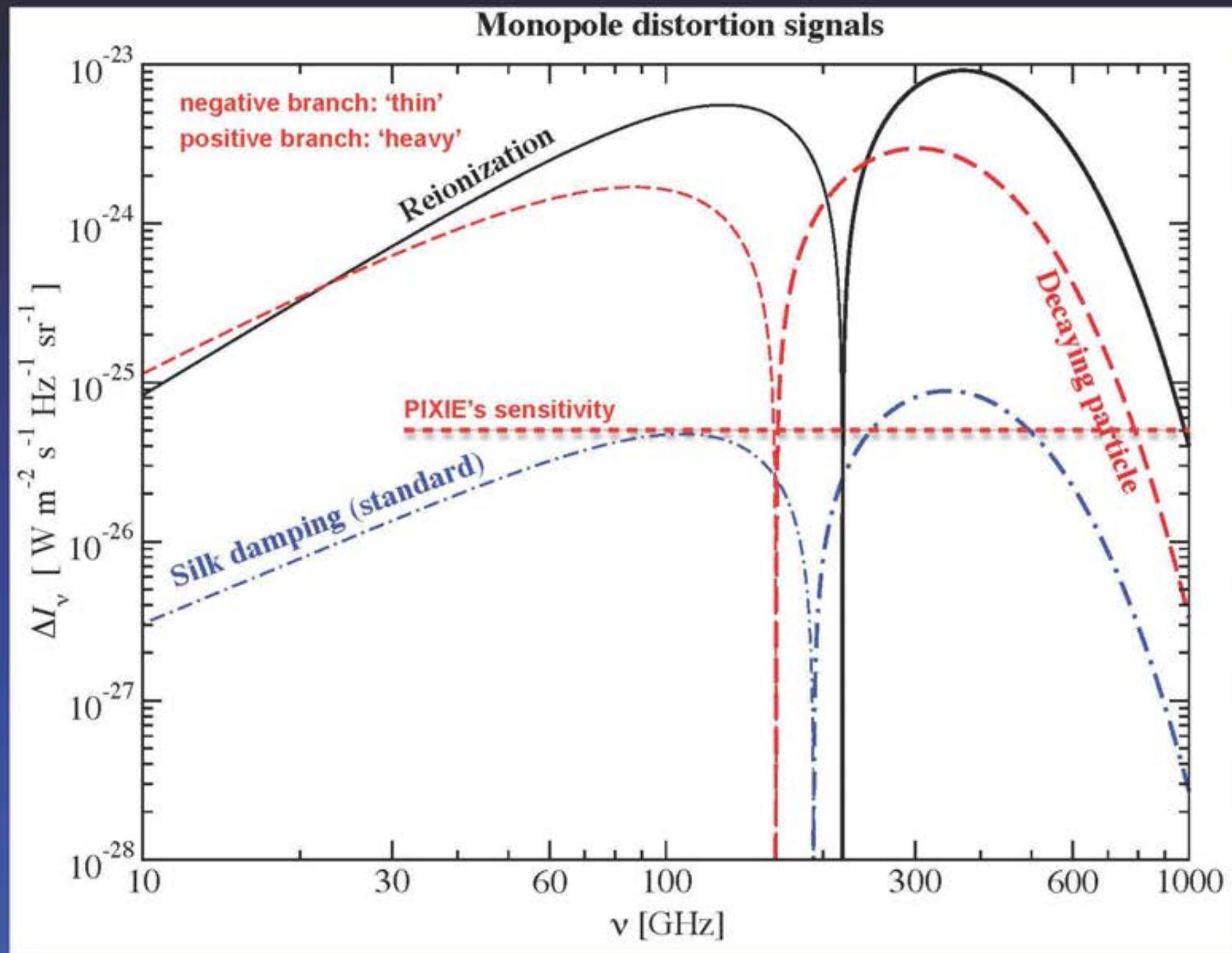
Which modes dissipate in the μ and γ -eras?



- Single mode with wavenumber k dissipates its energy at $z_d \sim 4.5 \times 10^5 (k \text{ Mpc}/10^3)^{2/3}$
- Modes with wavenumber $50 \text{ Mpc}^{-1} < k < 10^4 \text{ Mpc}^{-1}$ dissipate their energy during the μ -era
- Modes with $k < 50 \text{ Mpc}^{-1}$ cause γ -distortion

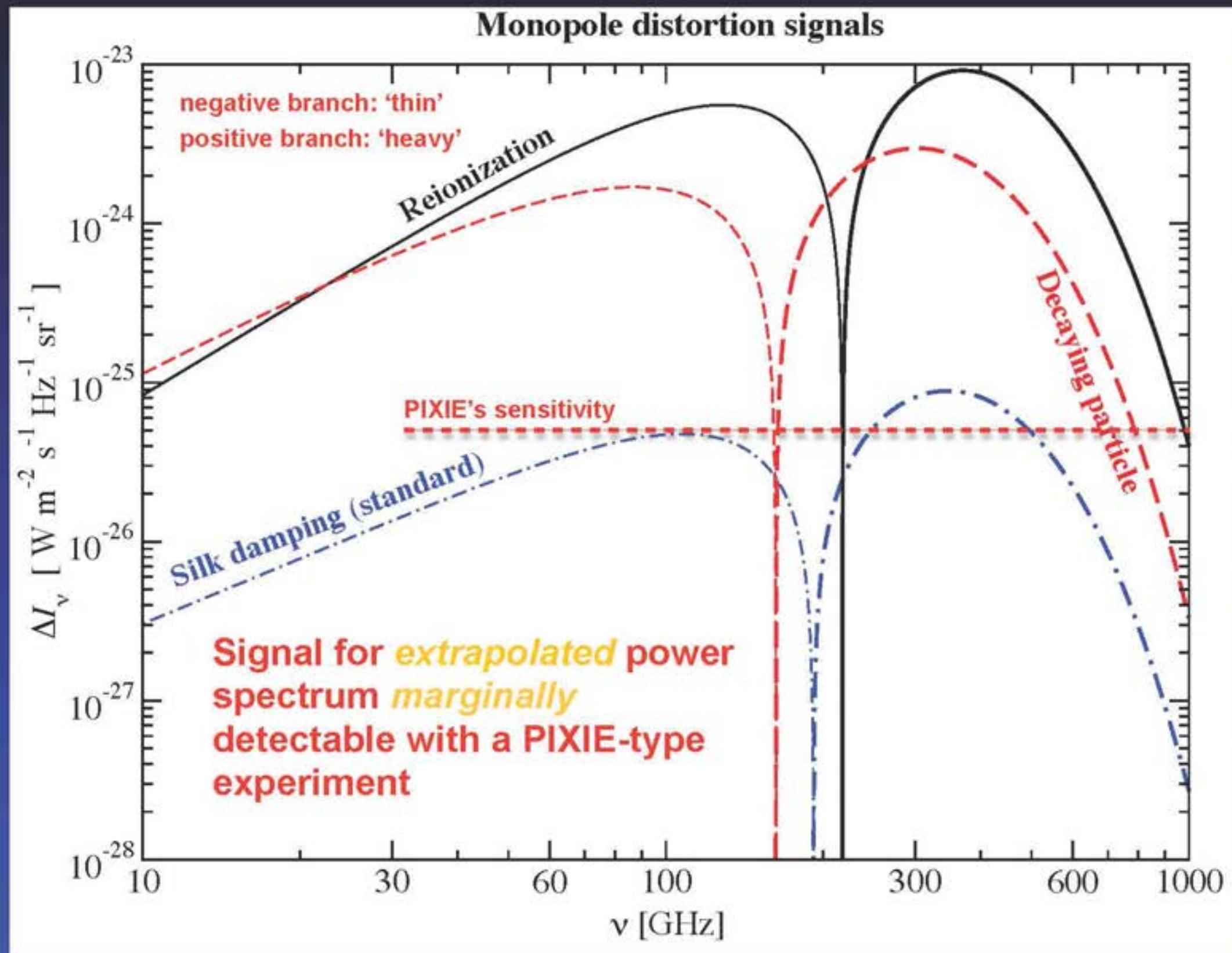
Average CMB spectral distortions

Absolute value of Intensity signal



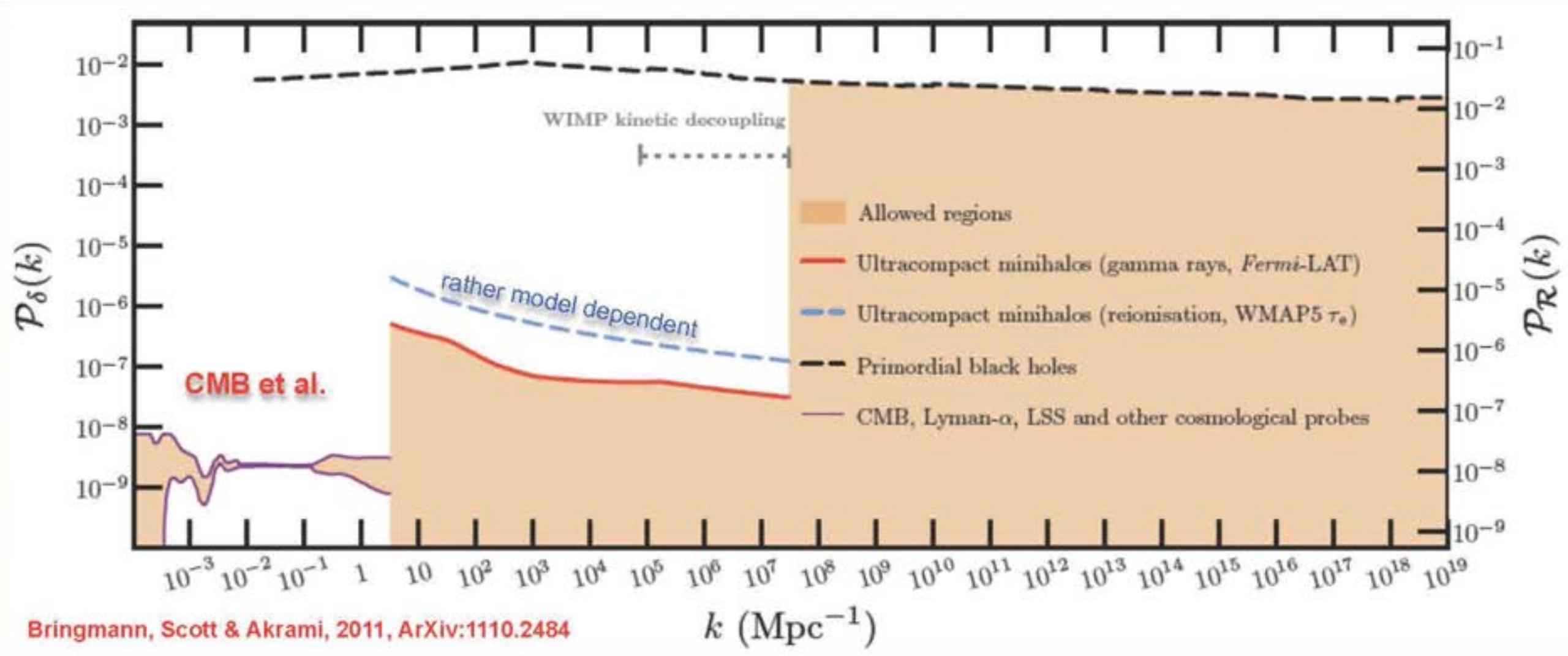
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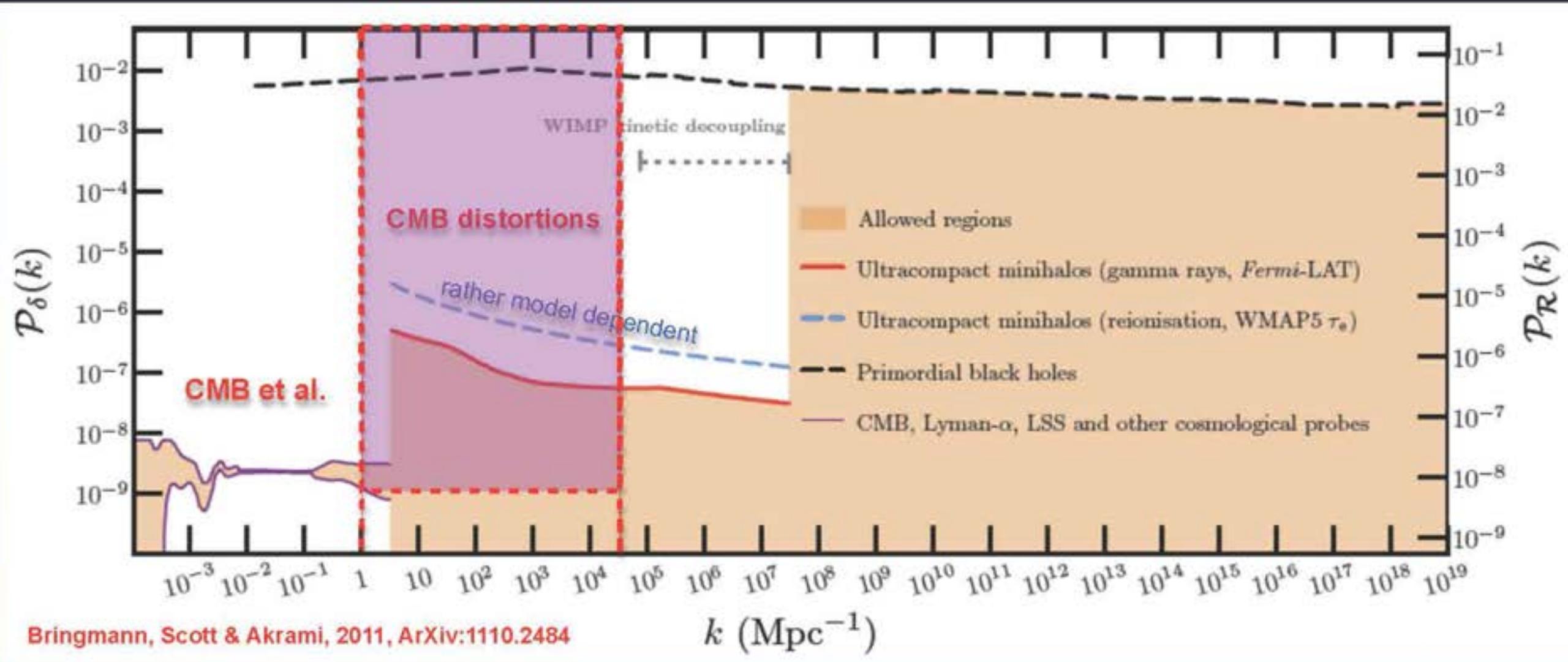
But this is not all that one could look at !!!

Power spectrum constraints



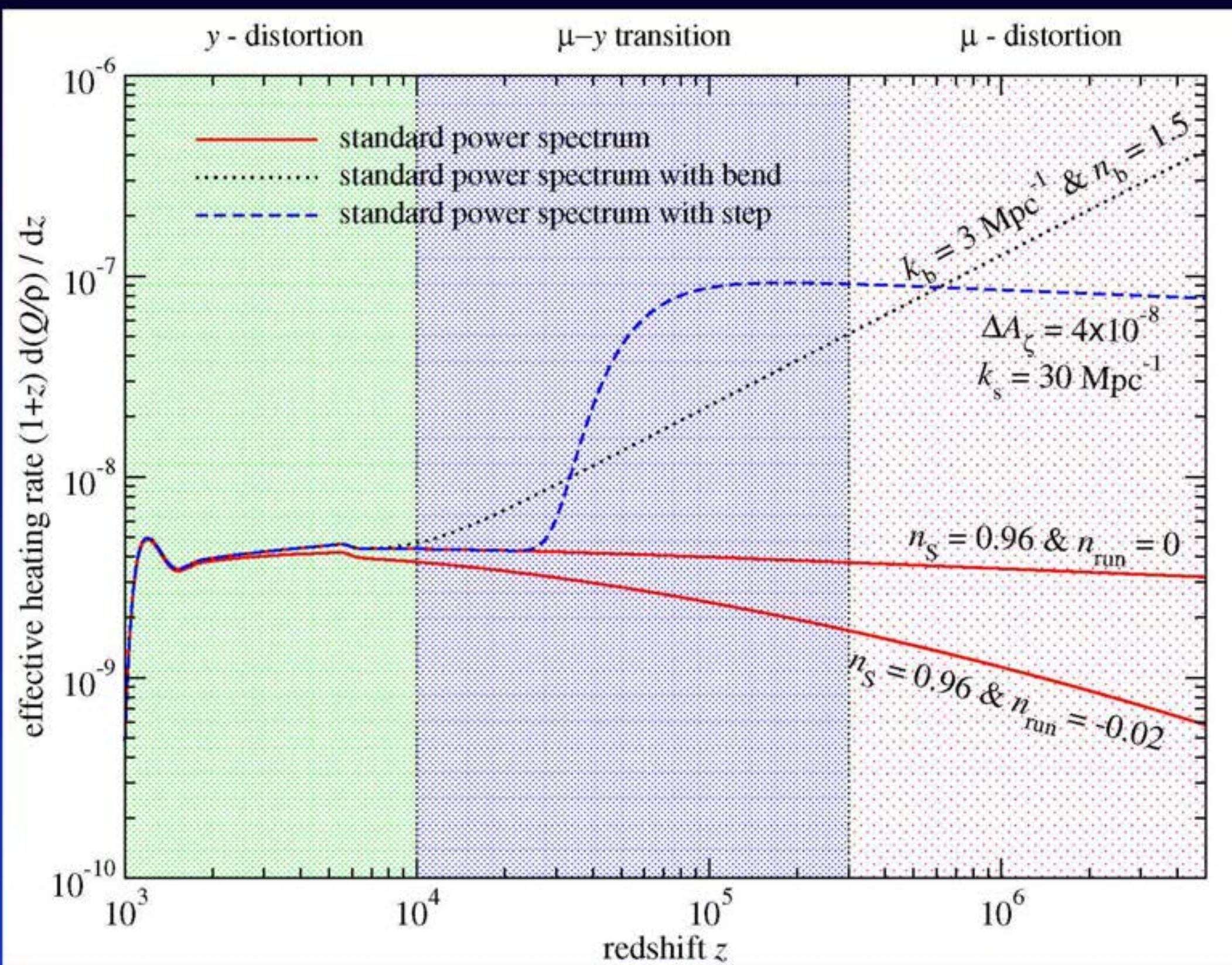
- Amplitude of power spectrum rather uncertain at $k > 3 \text{ Mpc}^{-1}$
- improving limits at smaller scales would constrain inflationary models

Power spectrum constraints

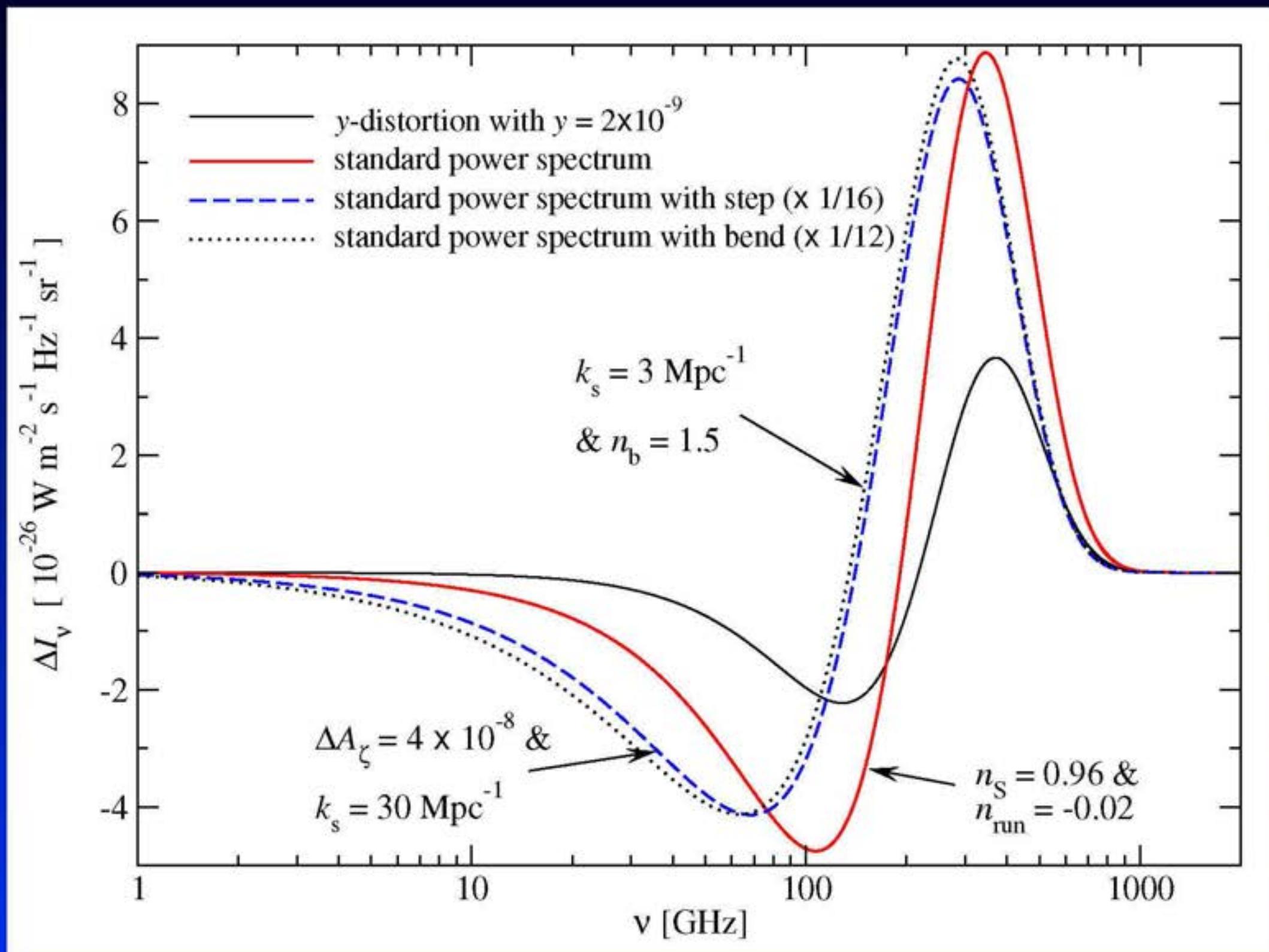


- Amplitude of power spectrum rather uncertain at $k > 3 \text{ Mpc}^{-1}$
- improving limits at smaller scales would constrain inflationary models
- CMB spectral distortions could allow extending our lever arm to $k \sim 10^4 \text{ Mpc}^{-1}$
- See JC, Erickcek & Ben-Dayan, 2012 for constraints on more general $P(k)$

Probing the small-scale power spectrum

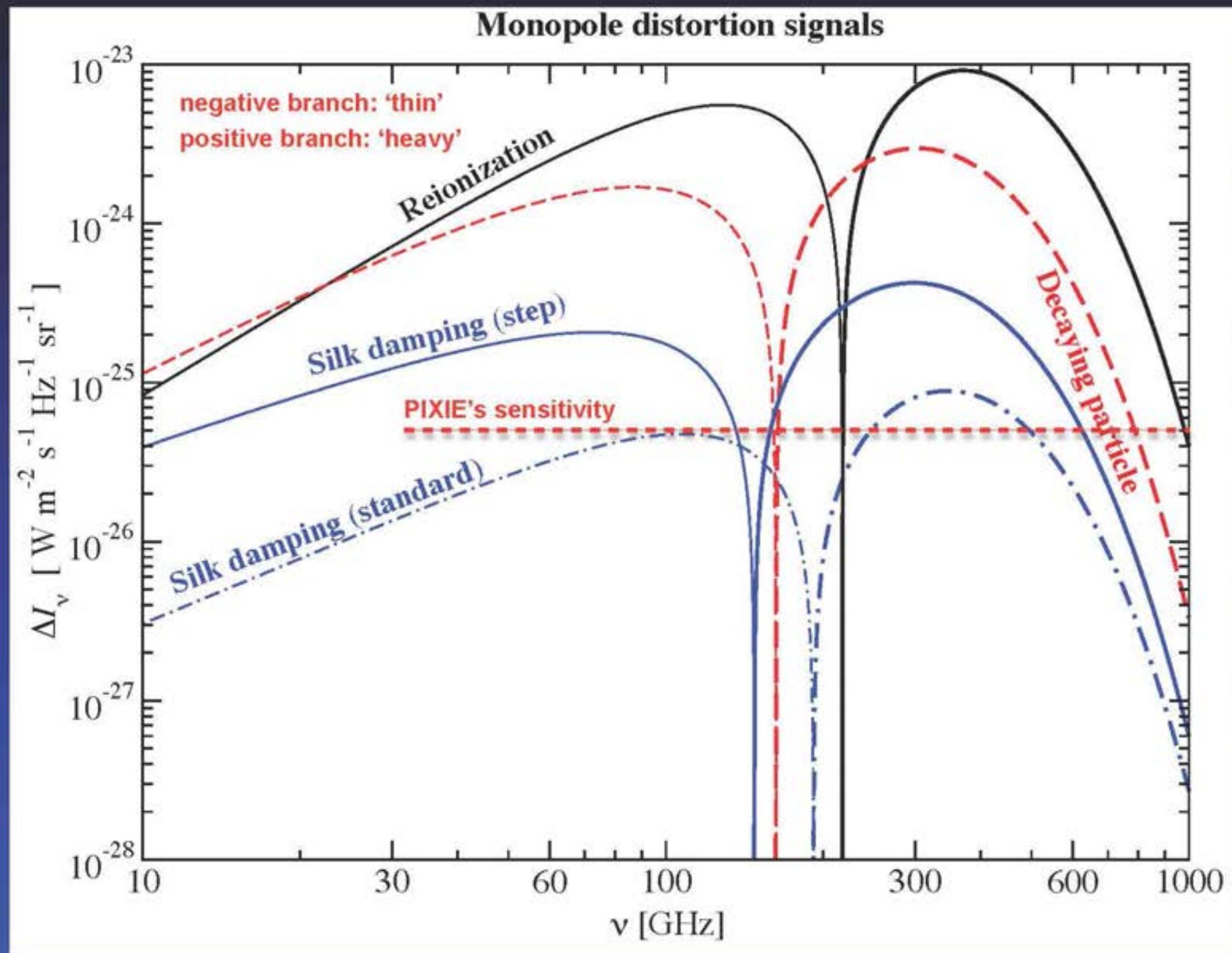


Probing the small-scale power spectrum



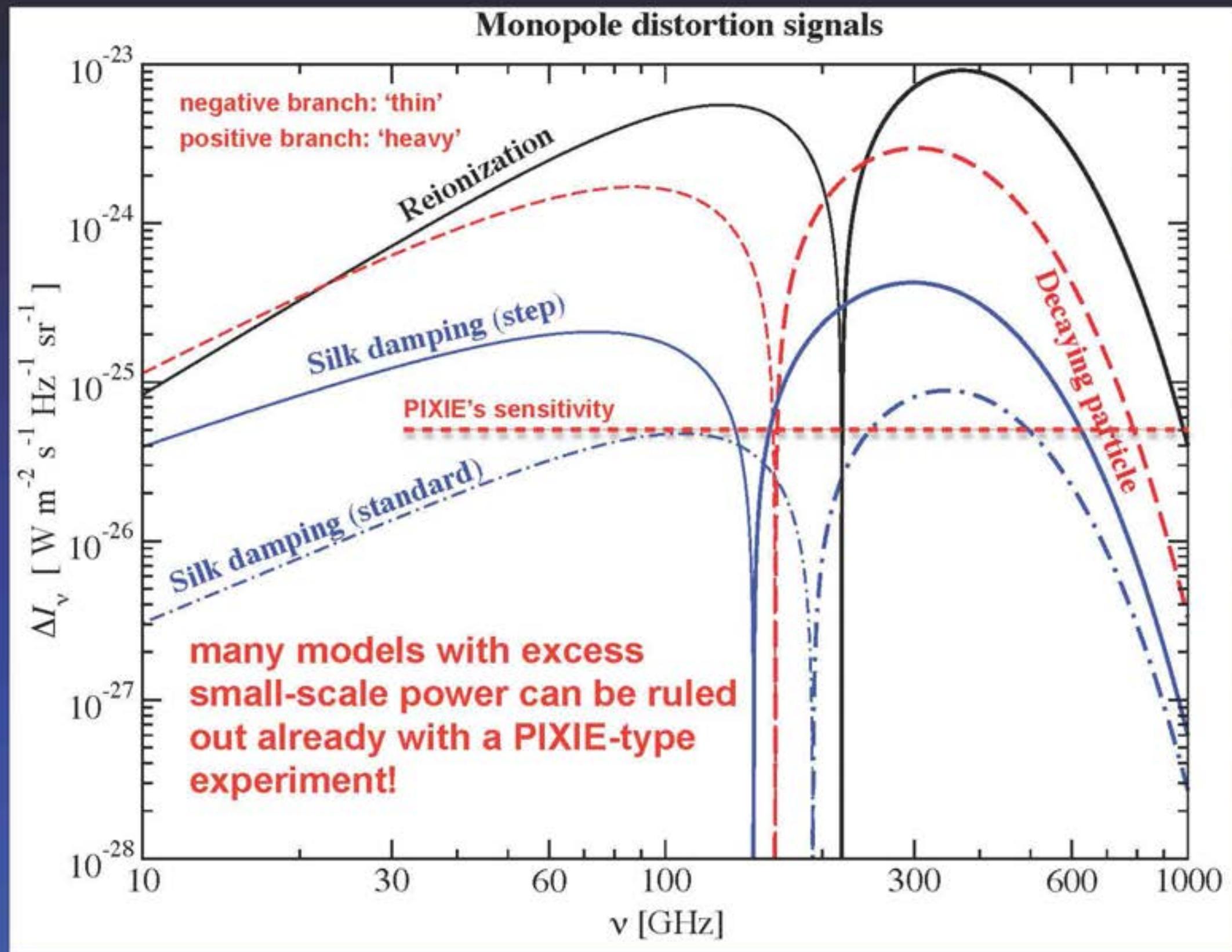
Average CMB spectral distortions

Absolute value of Intensity signal

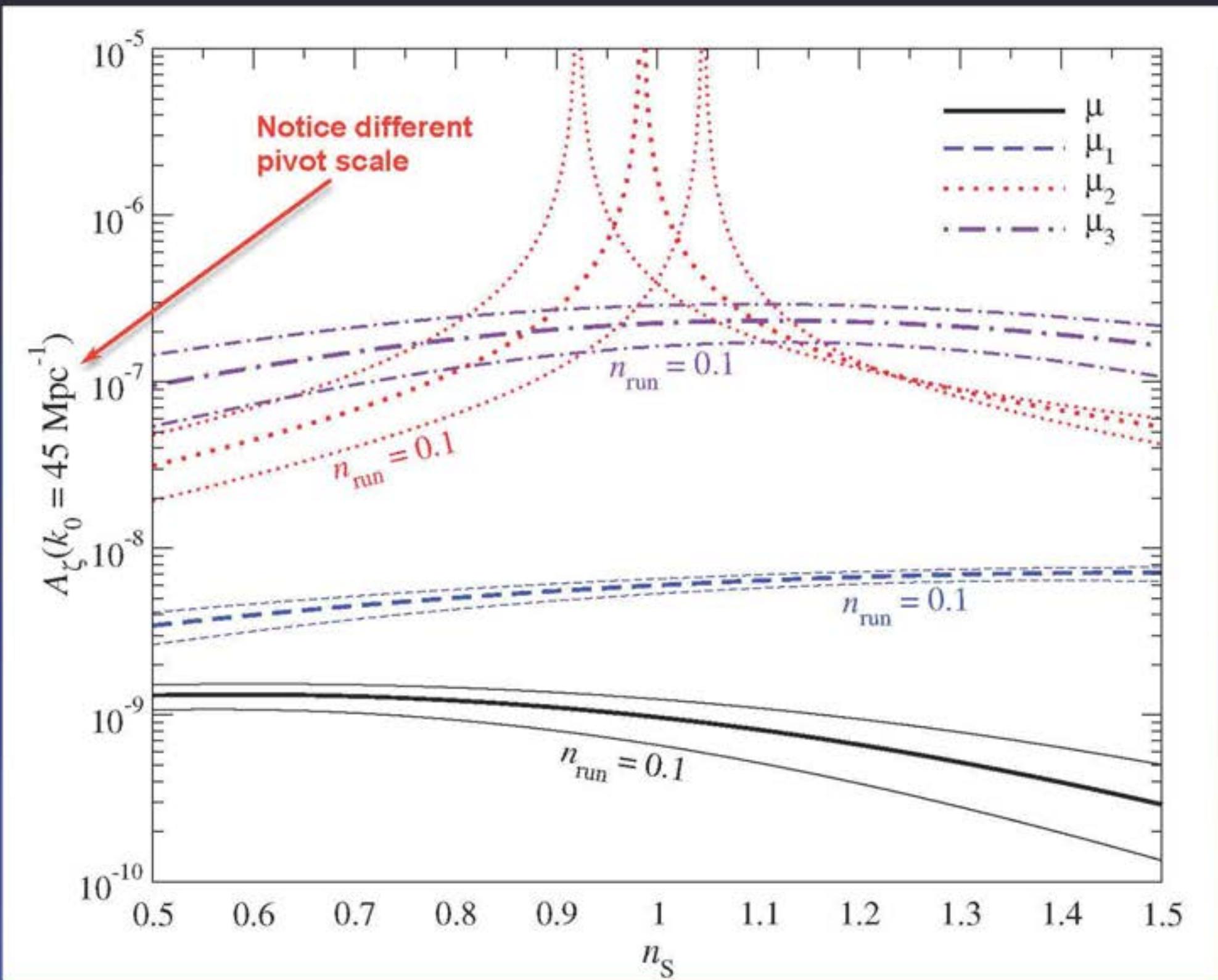


Average CMB spectral distortions

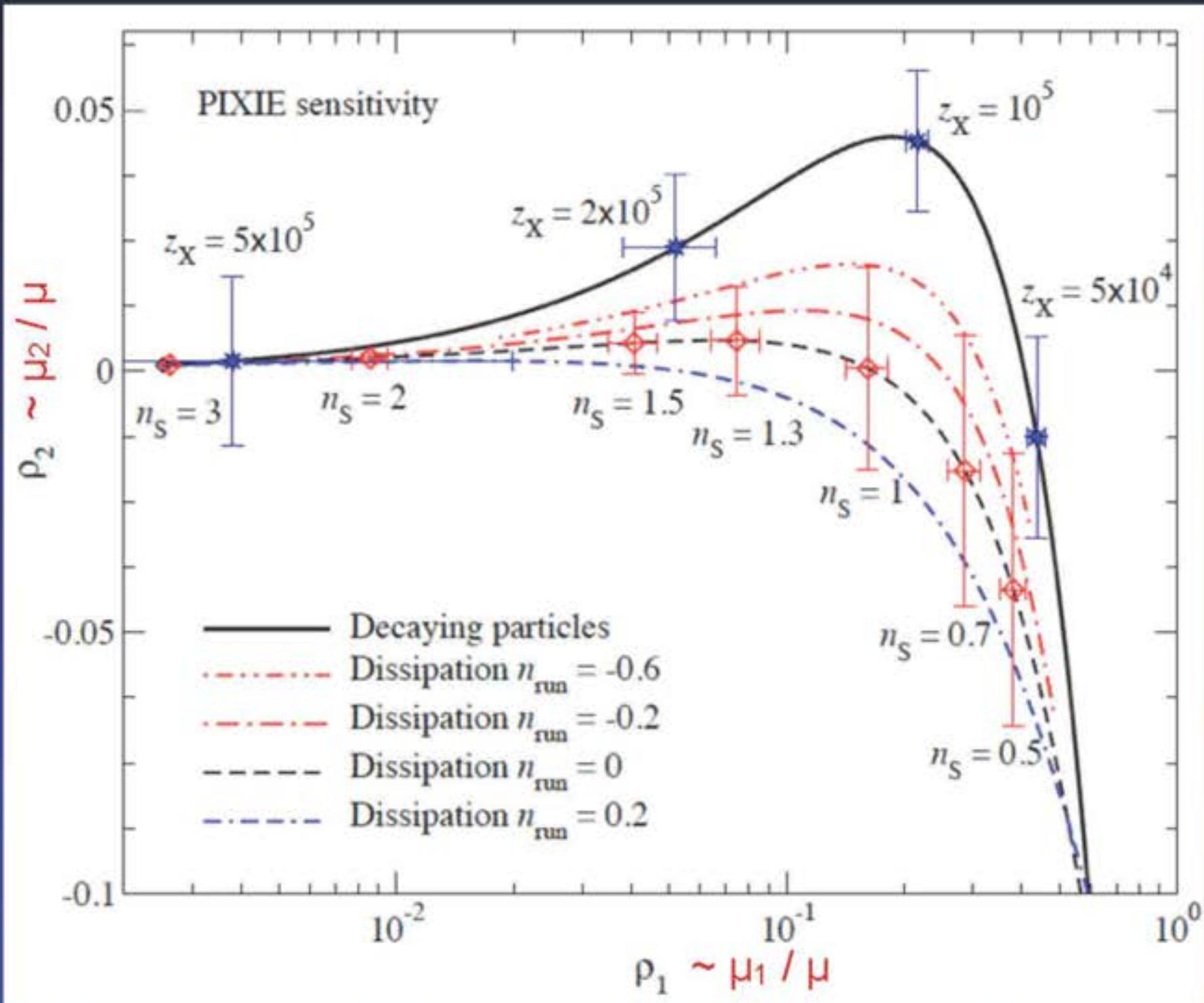
Absolute value of Intensity signal



Dissipation scenario: 1σ -detection limits for PIXIE



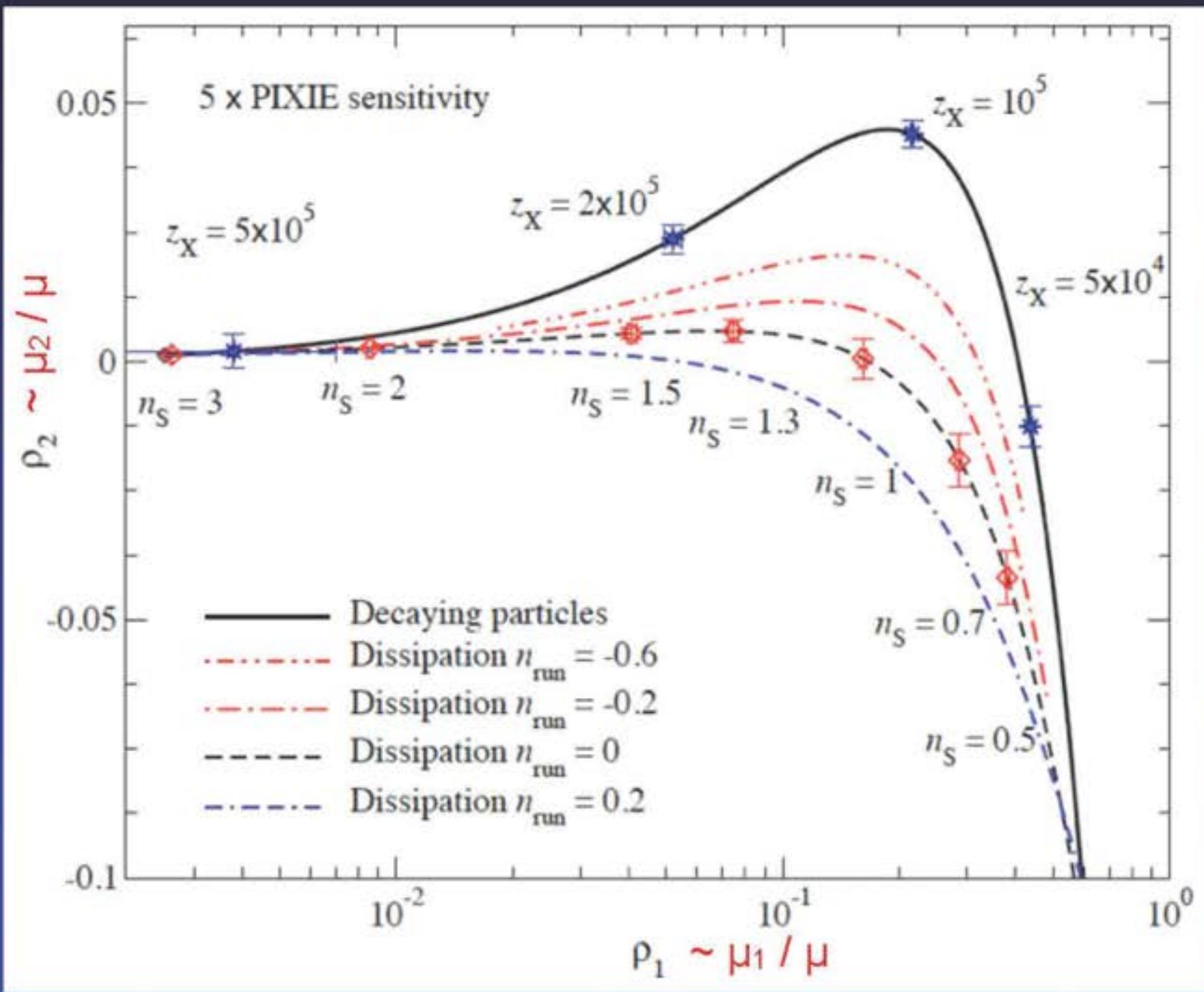
Distinguishing dissipation and decaying particle scenarios



- measurement of μ , μ_1 & μ_2
- trajectories of decaying particle and dissipation scenarios differ!
- scenarios can in principle be distinguished

$$A_\zeta = 5 \times 10^{-8}$$

Distinguishing dissipation and decaying particle scenarios



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The cosmological recombination radiation

Simple estimates for hydrogen recombination

Hydrogen recombination:

- per recombined hydrogen atom an energy of ~ 13.6 eV in form of photons is released
- at $z \sim 1100 \rightarrow \Delta\epsilon/\epsilon \sim 13.6$ eV $N_b / (N_\gamma 2.7kT_r) \sim 10^{-9} - 10^{-8}$

- recombination occurs at redshifts $z < 10^4$
- At that time the *thermalization* process doesn't work anymore!
- There should be some *small* spectral distortion due to additional Ly- α and 2s-1s photons!

(Zeldovich, Kurt & Sunyaev, 1968, ZhETF, 55, 278; Peebles, 1968, ApJ, 153, 1)

- In 1975 **Viktor Dubrovich** emphasized the possibility to observe the recombinational lines from $n > 3$ and $\Delta n \ll n$!

First recombination computations completed in 1968!



Moscow



Vladimir Kurt
(UV astronomer)



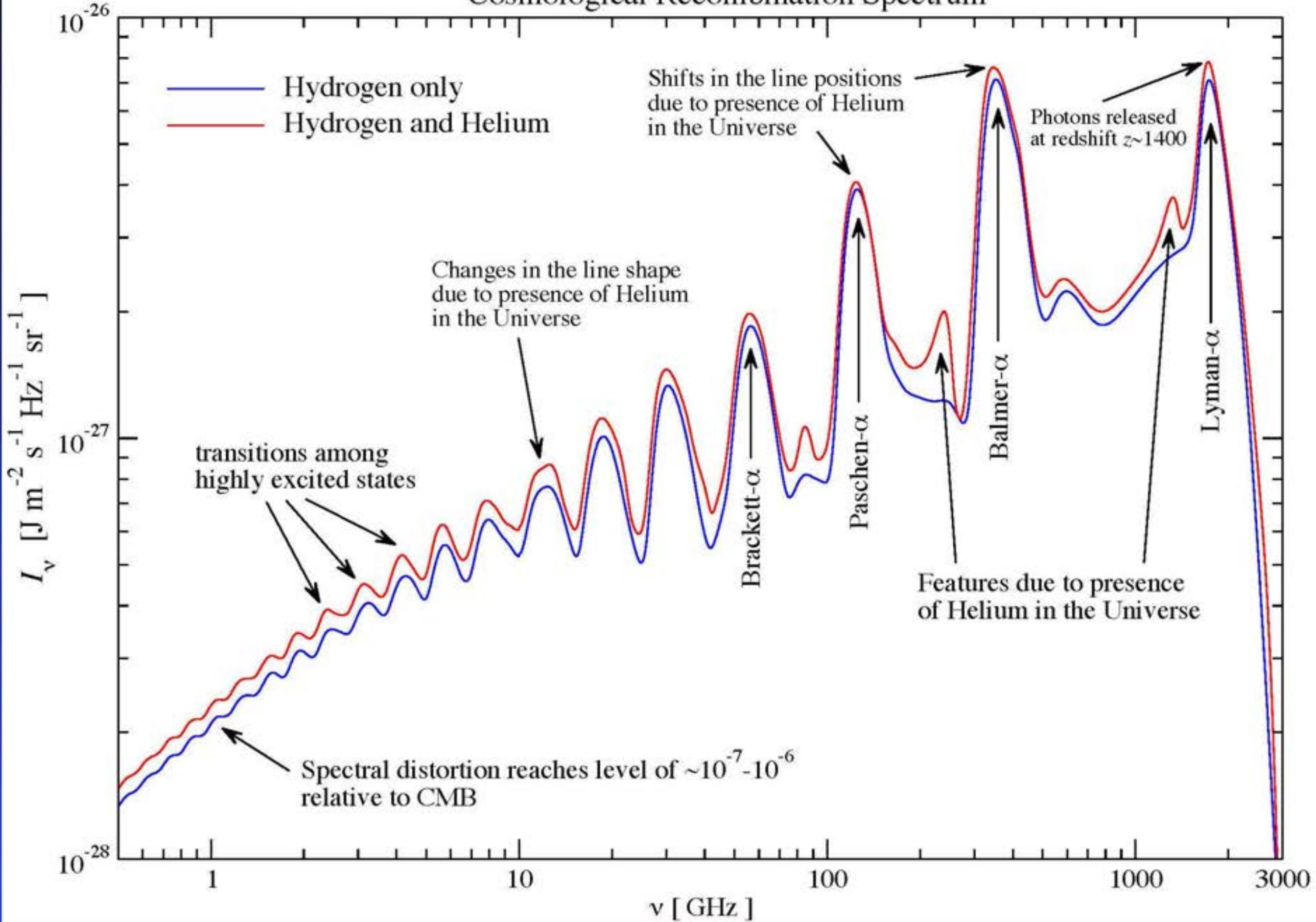
Rashid Sunyaev



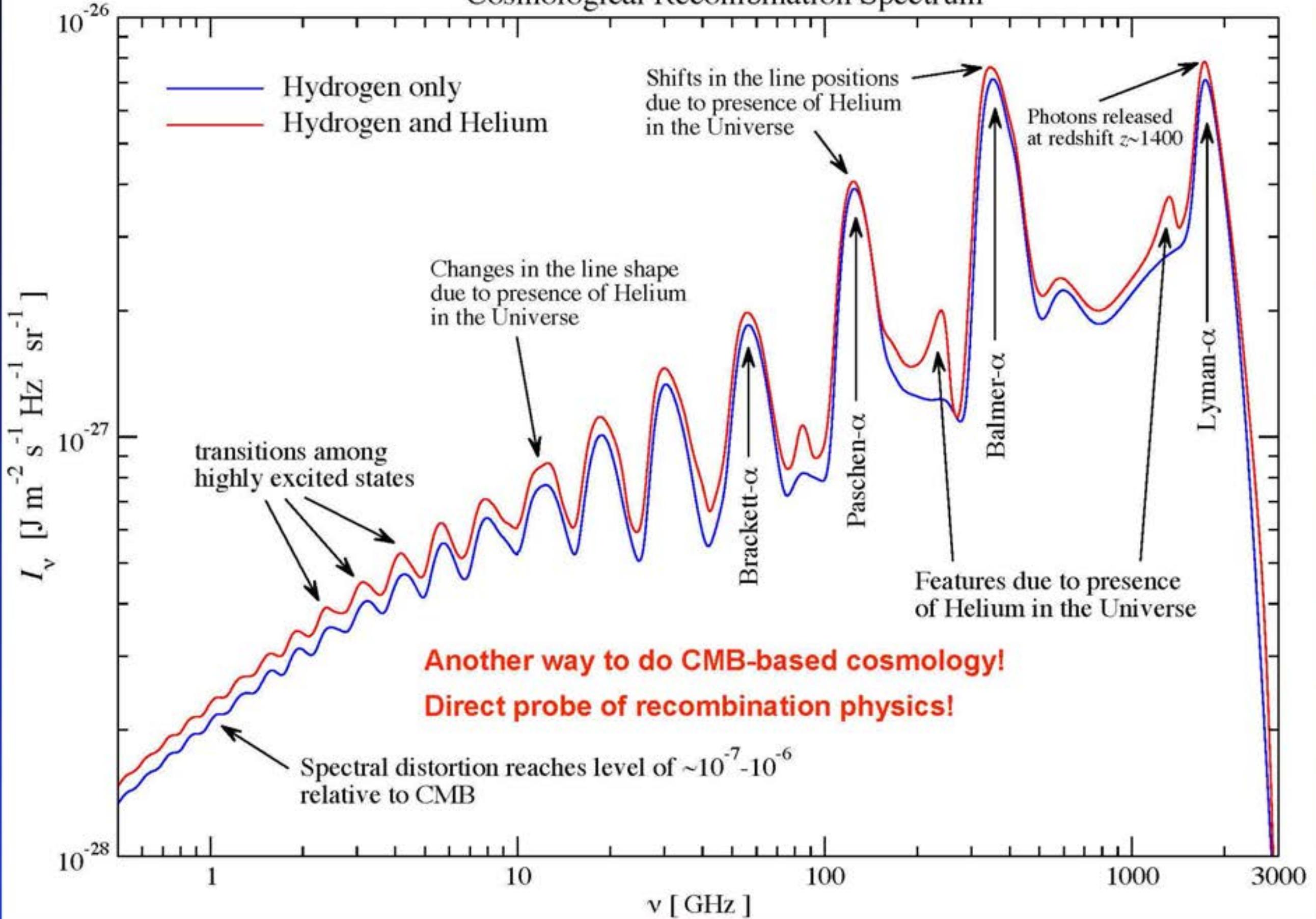
Jim Peebles

Princeton

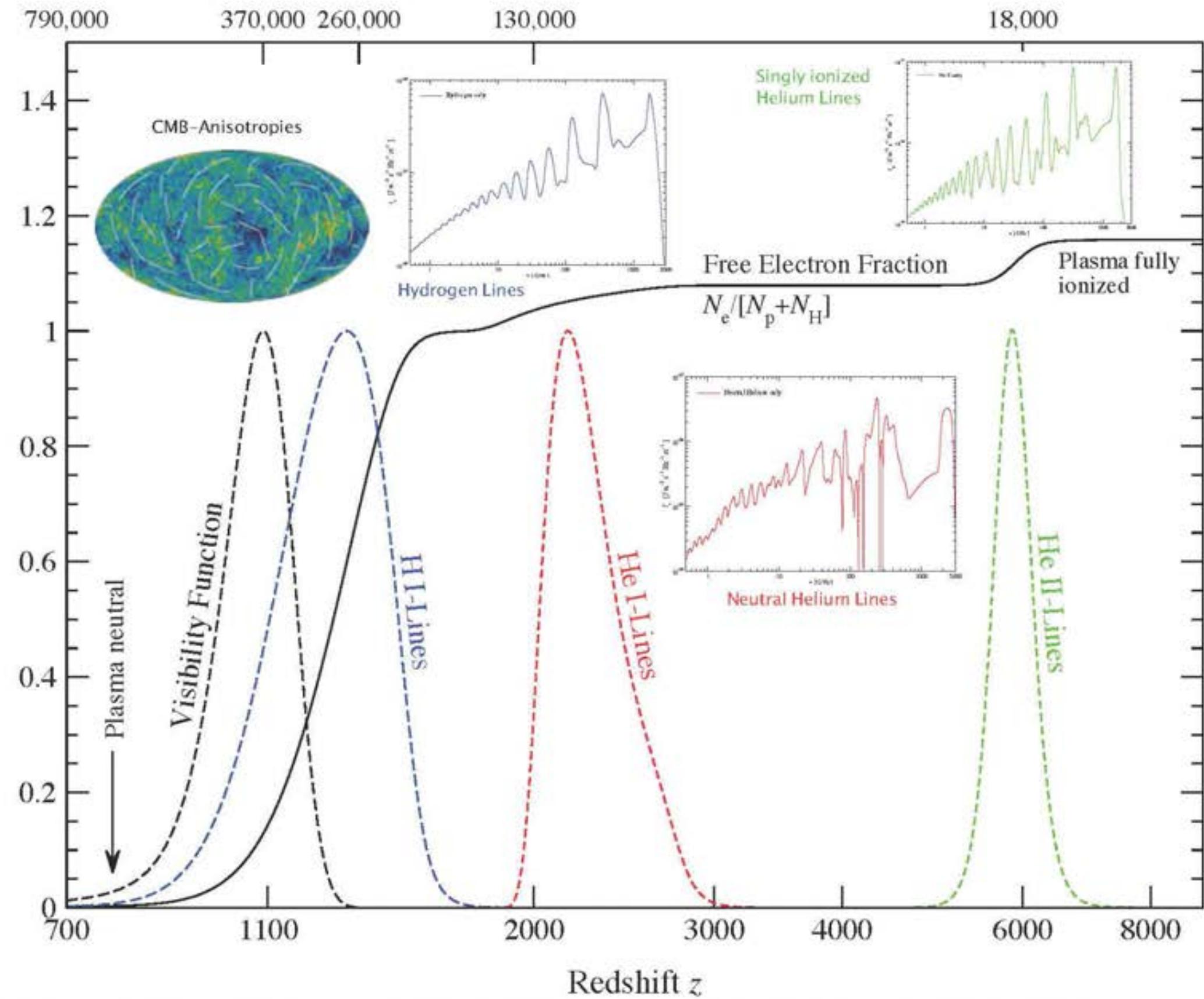
Cosmological Recombination Spectrum



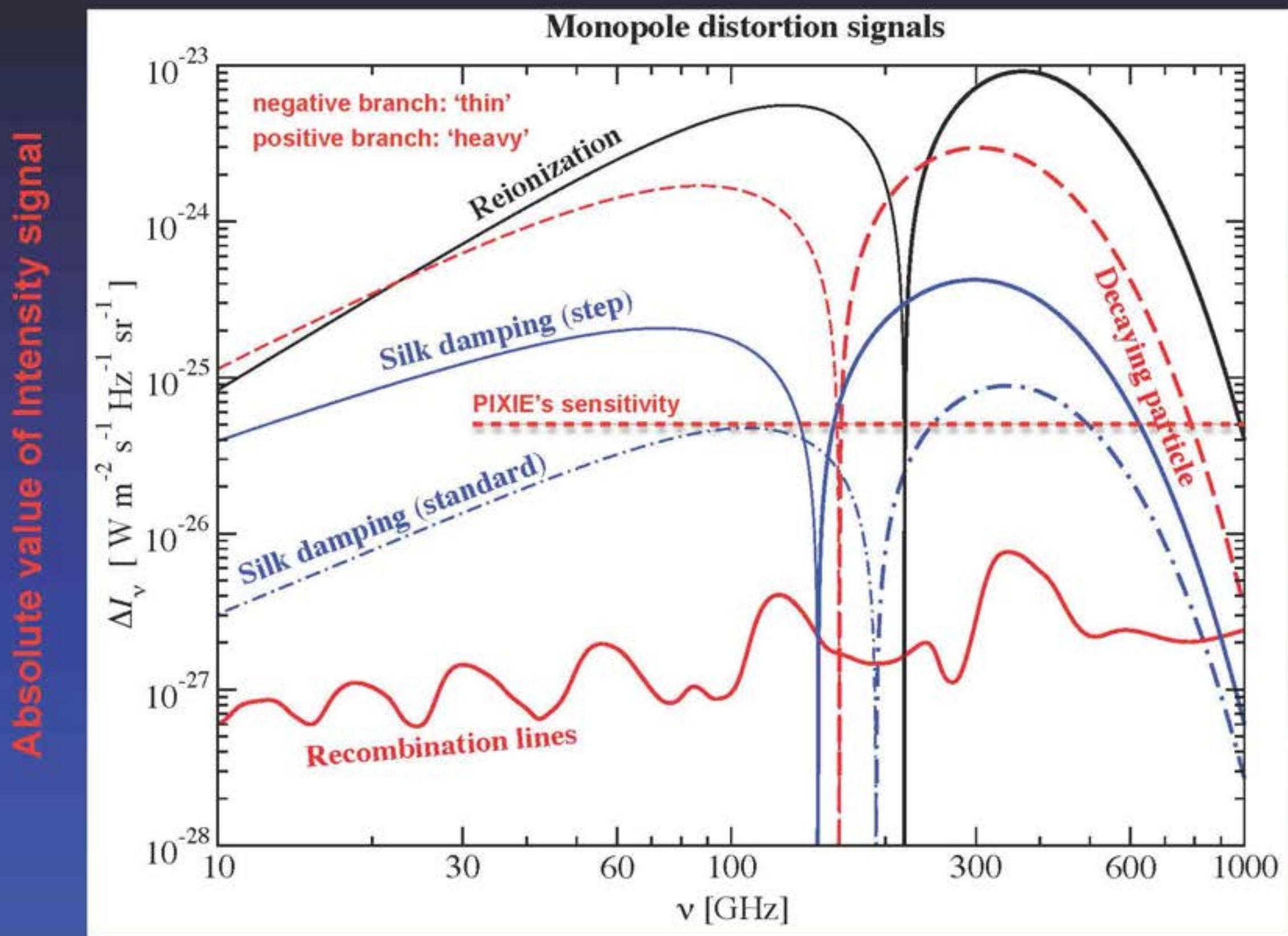
Cosmological Recombination Spectrum



Cosmological Time in Years

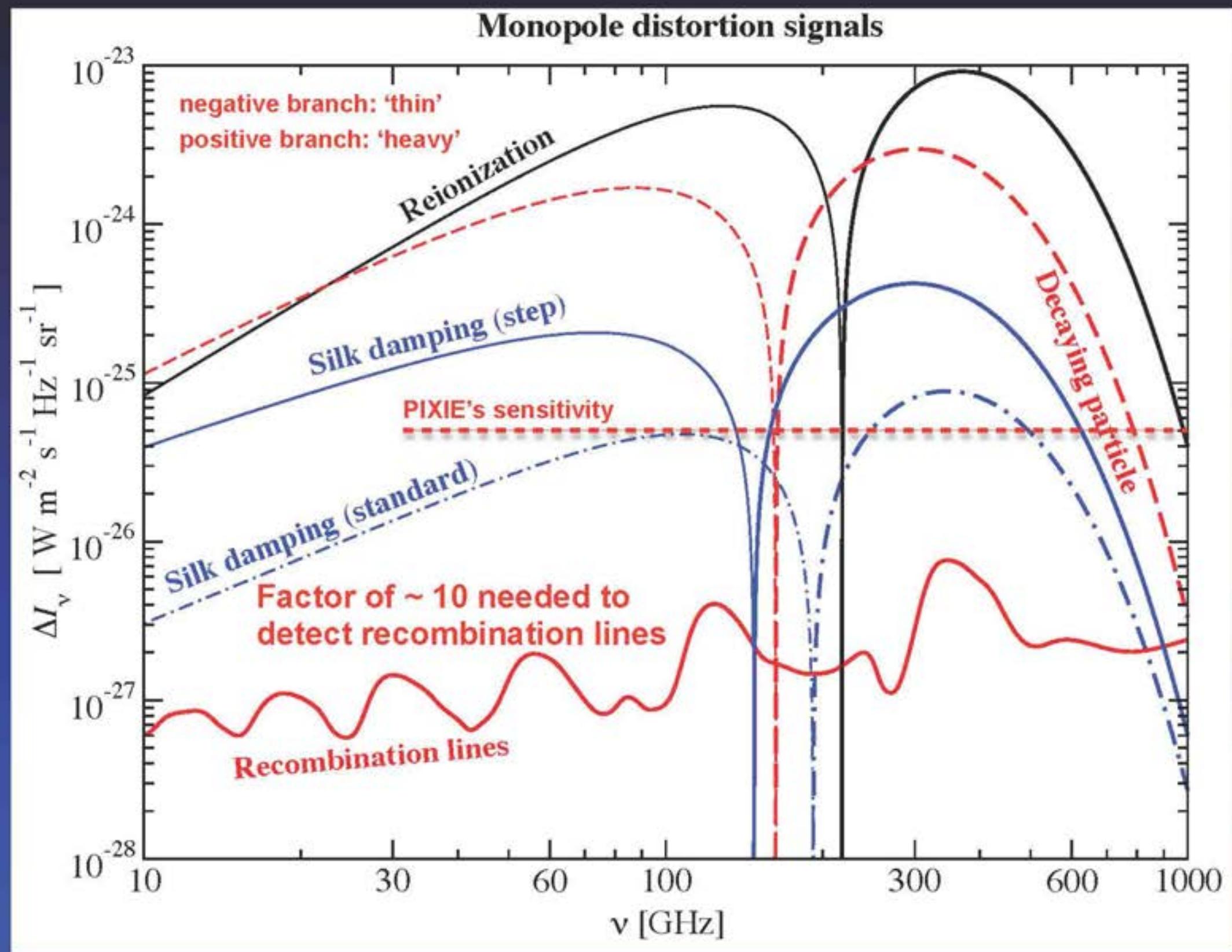


Average CMB spectral distortions

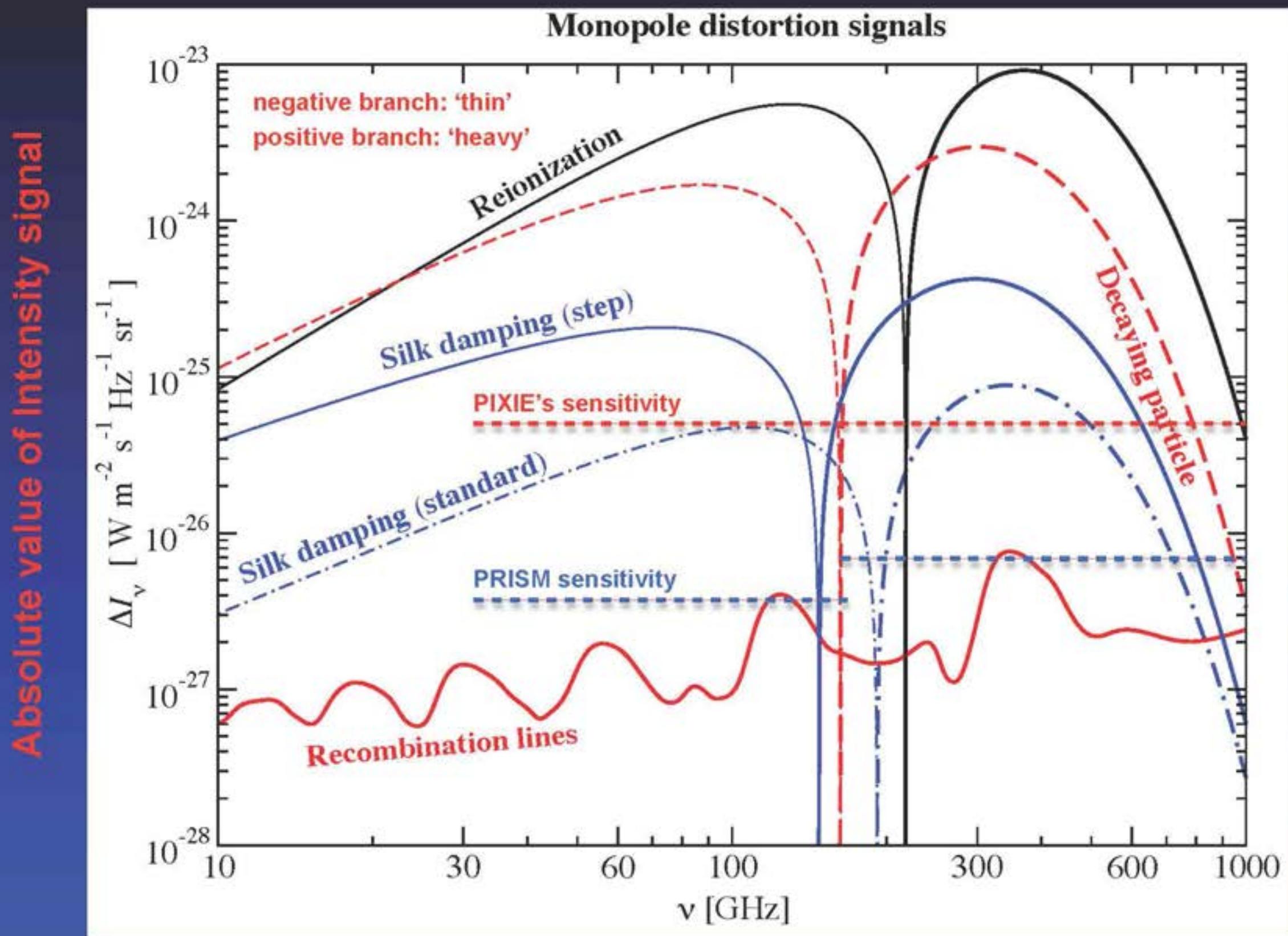


Average CMB spectral distortions

Absolute value of Intensity signal



Average CMB spectral distortions



What would we actually learn by doing such hard job?

Cosmological Recombination Spectrum opens a way to measure:

- the specific **entropy** of our universe (related to $\Omega_b h^2$)
- the CMB **monopole** temperature T_0
- **the pre-stellar abundance of helium Y_p**
- **If recombination occurs as we think it does, then the lines can be predicted with very high accuracy!**
- **In principle allows us to directly check our understanding of the standard recombination physics**

computations prepared by Chad Fendt
in 2009 using detailed recombination code

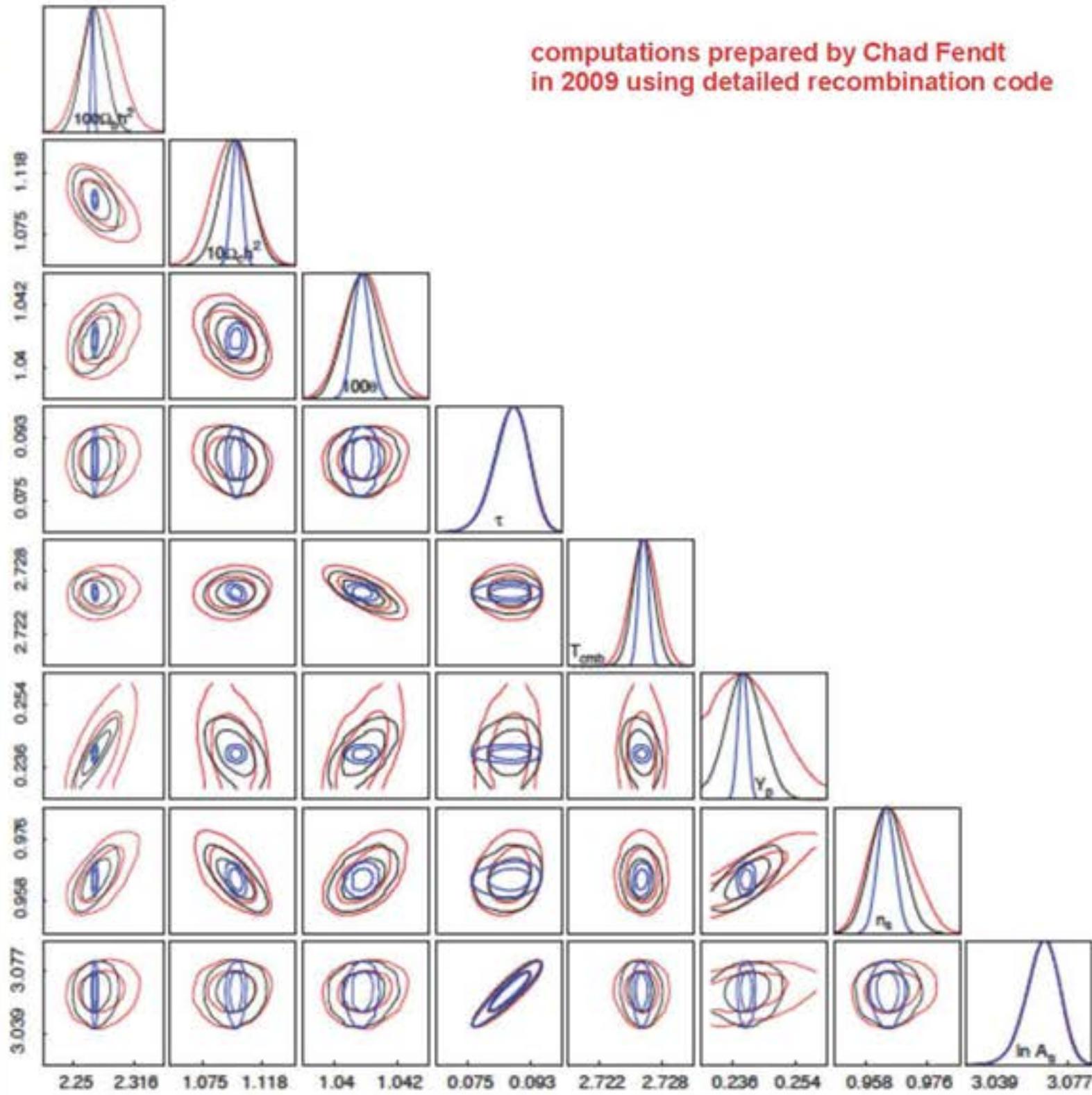
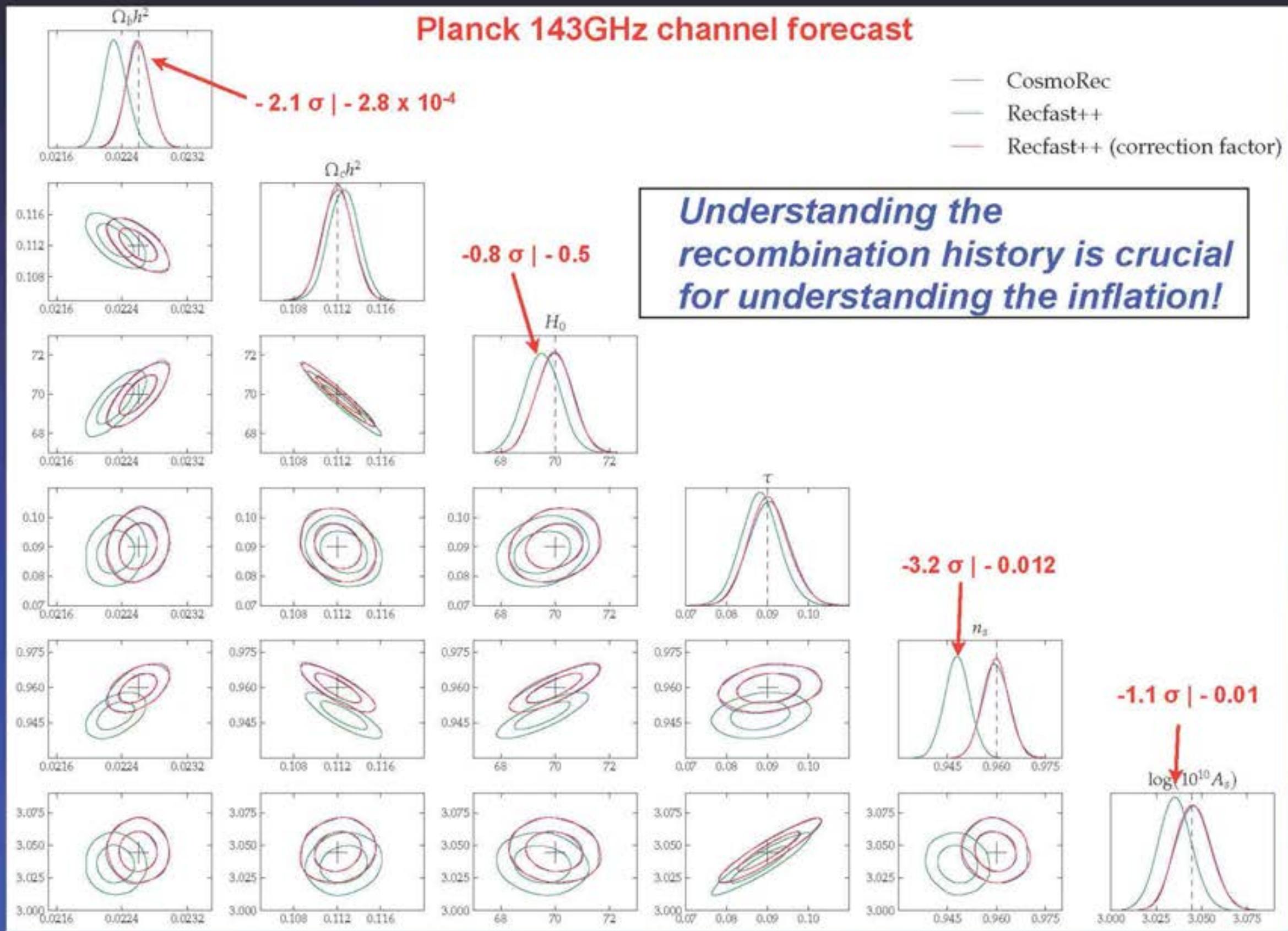


Figure 7.3: The 1 and 2 dimensional marginalized parameter posterior using the CMB spectral distortions. All three cases constrain the CMB power spectrum using a Gaussian likelihood based on Planck noise levels. The black line adds constraints due to a 10% measurement of the spectral distortions, while the blue line assumes a 1% measurement. The red line does not include the data from the spectral distortions.

- CMB based cosmology alone
- Spectrum helps to break some of the parameter degeneracies
- Planning to provide a module that computes the recombination spectrum in a fast way
- detailed forecasts: which lines to measure; how important is the absolute amplitude; how accurately one should measure; best frequency resolution;

Importance of recombination



What would we actually learn by doing such hard job?

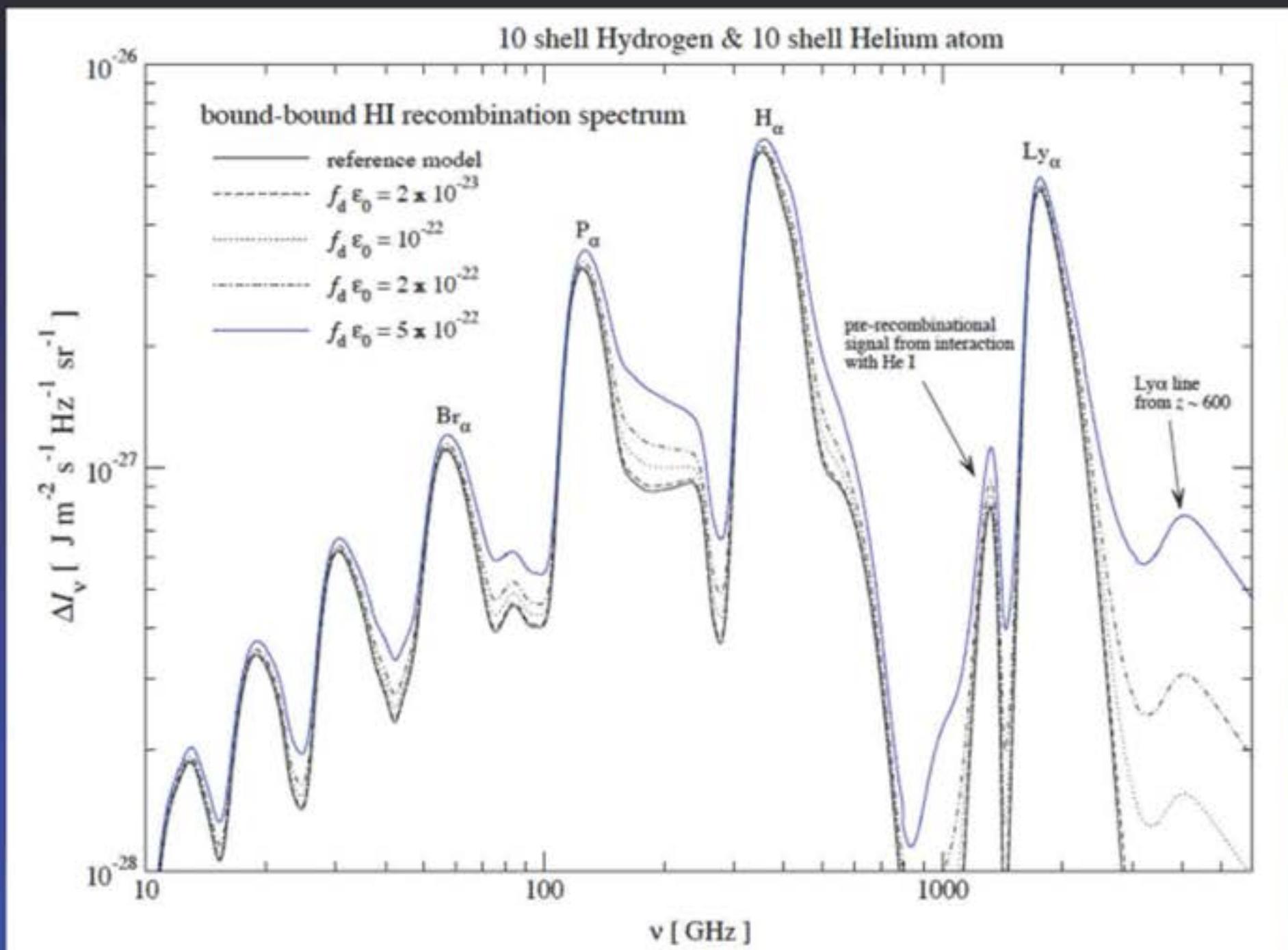
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If something unexpected or non-standard happened:

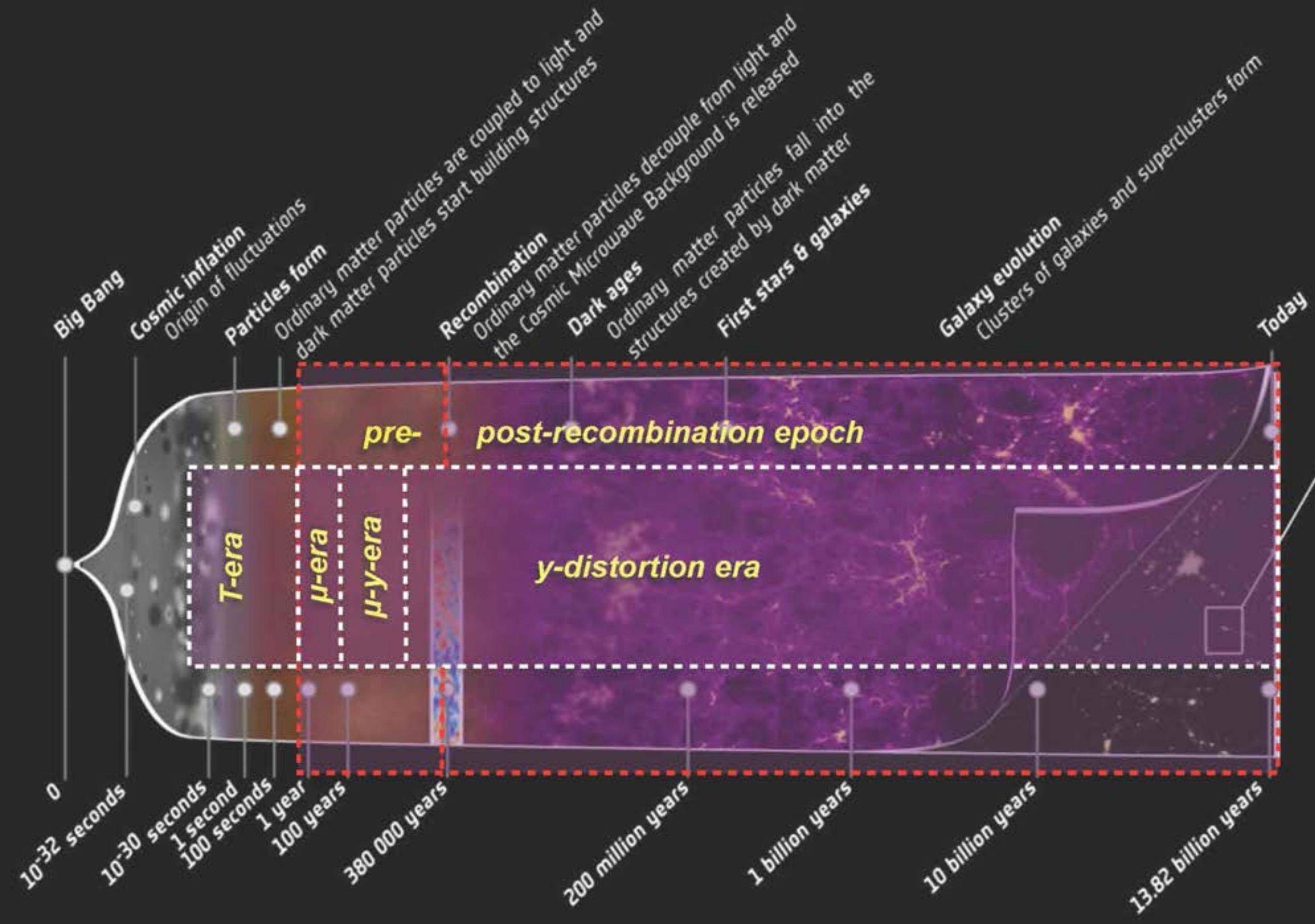
- **non-standard thermal histories should leave some measurable traces**
- **direct way to measure/reconstruct the recombination history!**
- **possibility to distinguish pre- and post-recombination y-type distortions**
- **sensitive to energy release during recombination**
- **variation of fundamental constants**

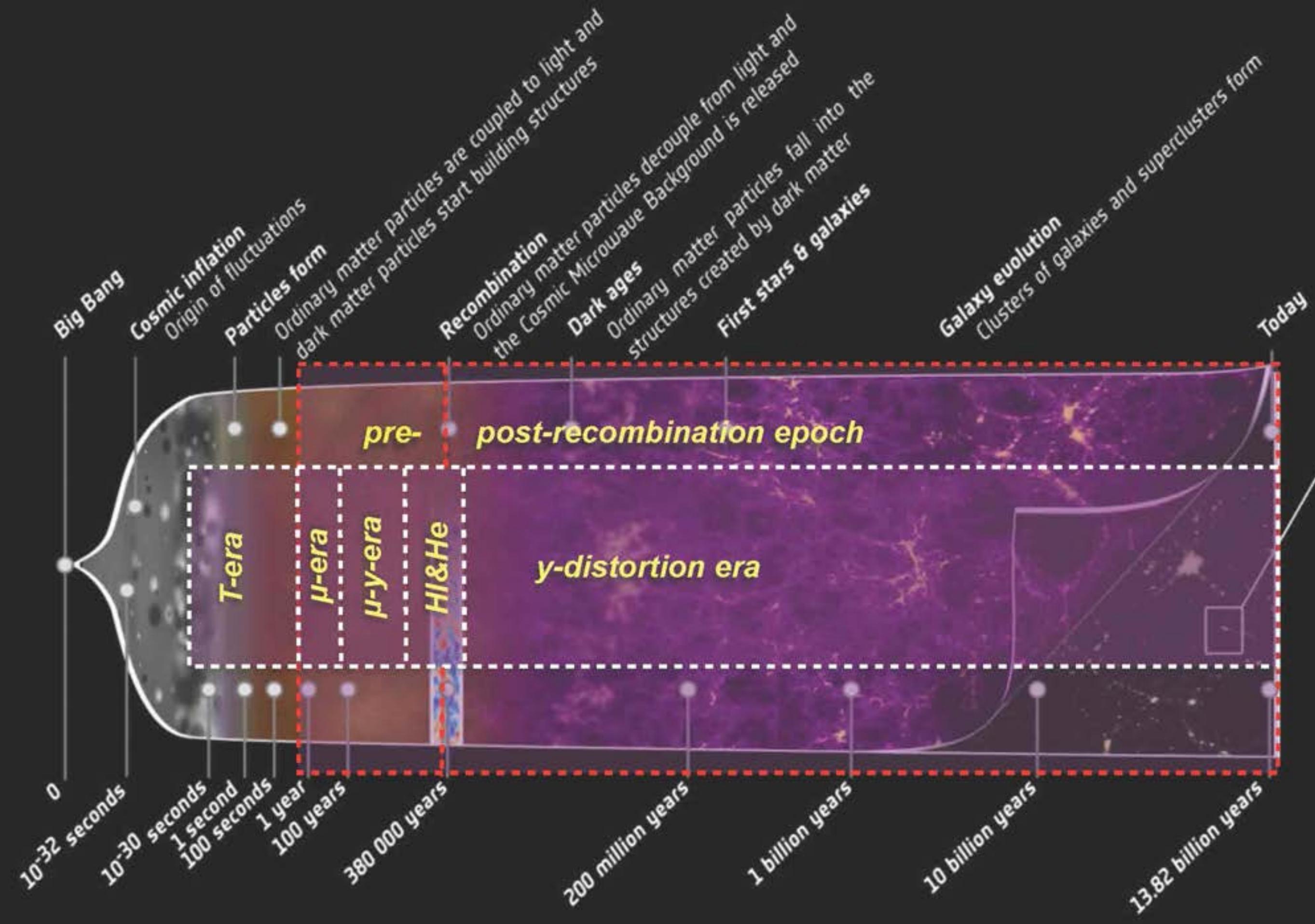
Dark matter annihilations / decays



- Additional photons at all frequencies
- Broadening of spectral features
- Shifts in the positions

JC, 2009, arXiv:0910.3663





Conclusions

- CMB spectral distortions open a *new window* to the early Universe and inflationary epoch
- *complementary* and *independent* source of information about our Universe *not* just confirmation
- simplicity of thermalization physics allows making very *precise predictions* for the distortions caused by different heating mechanisms
- in *standard cosmology* several processes lead to *early energy release* at a level that will be detectable in the future
- extremely interesting *future* for CMB based science!