

Determination of the Mass of an Ultralight Gravitino at LHC

Koichi Hamaguchi (Tokyo U.)

with S. Shirai and T.T. Yanagida

today!

arXiv:0705.0219, 0707.2463, 0712.2462

@ "IPMU focus week: Facing LHC data", Dec. 2007

Summary: (beer is waiting for us....)

Main Message

SUSY models with an
ultralight gravitino is interesting!

$$(m_{\tilde{G}} \lesssim 10 \text{ eV})$$

- No Cosmological Problem! at all!
- LSP (gravitino) \neq DM, but a natural DM candidate.
- It can be tested at LHC!
(gravitino mass can be determined!)

Gravitino

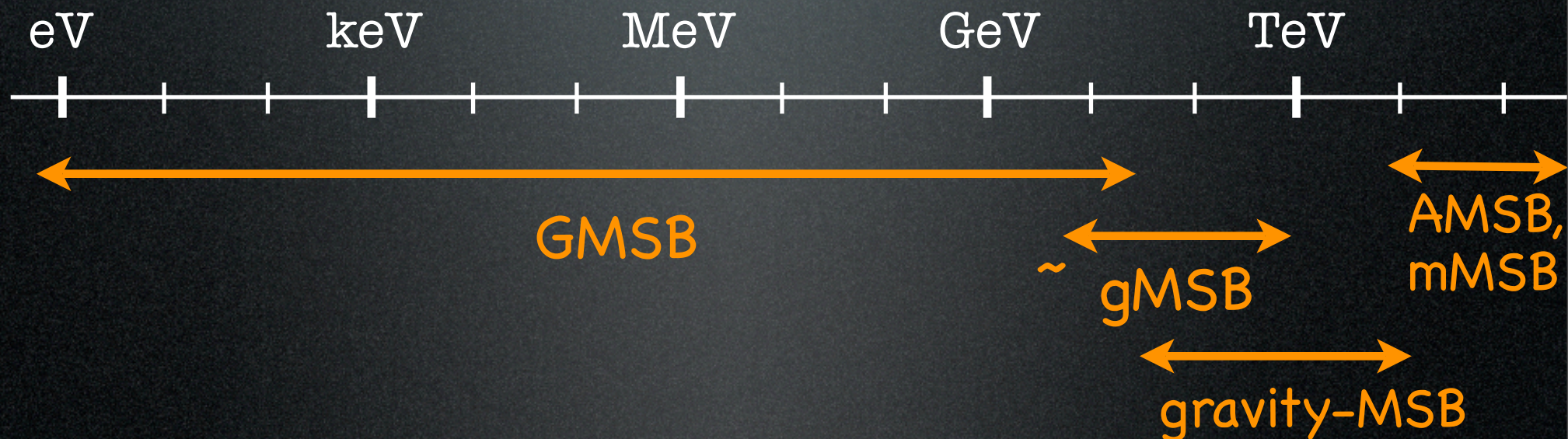
- If **SUSY** is in nature, it is (probably) a spontaneously broken **local symmetry** (i.e., not an accidental global symmetry).
- And the **gravitino** is an inevitable prediction of **local SUSY (= SUGRA)**.

Gravitino

- Gravitino Interaction: extremely weak

suppressed by $\sim \frac{1}{M_{\text{P}}}$ (or $\sim \frac{1}{F} \sim \frac{1}{M_{\text{P}} m_{\tilde{G}}}$)

- Gravitino Mass: model dependent

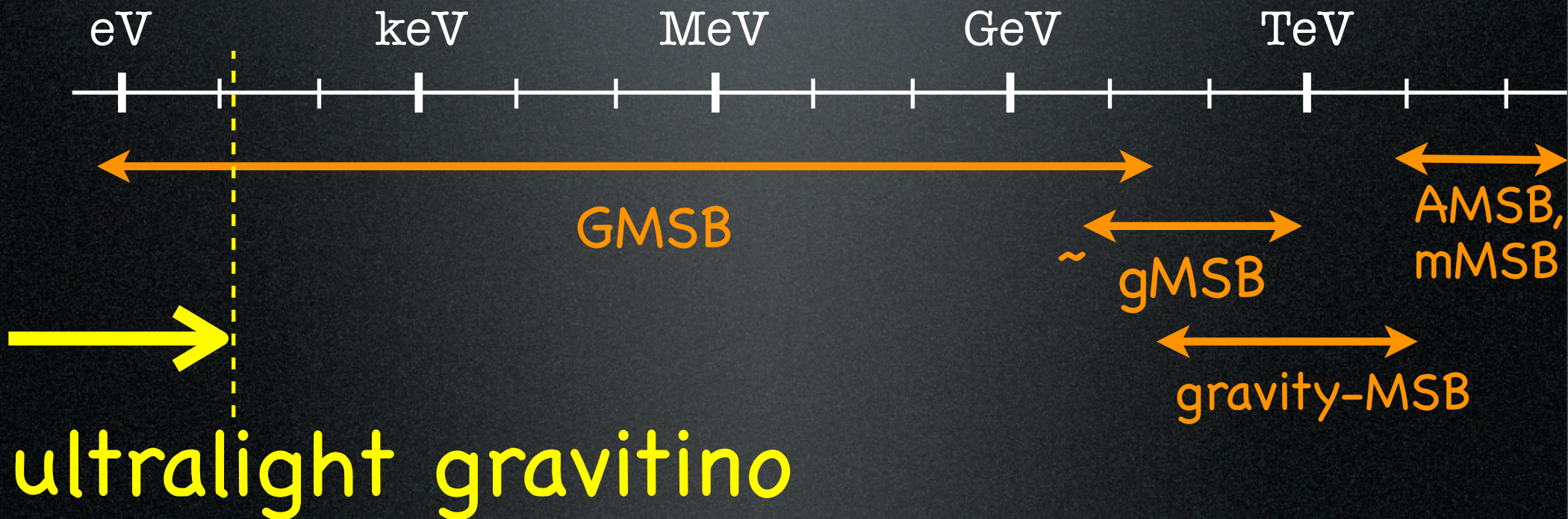


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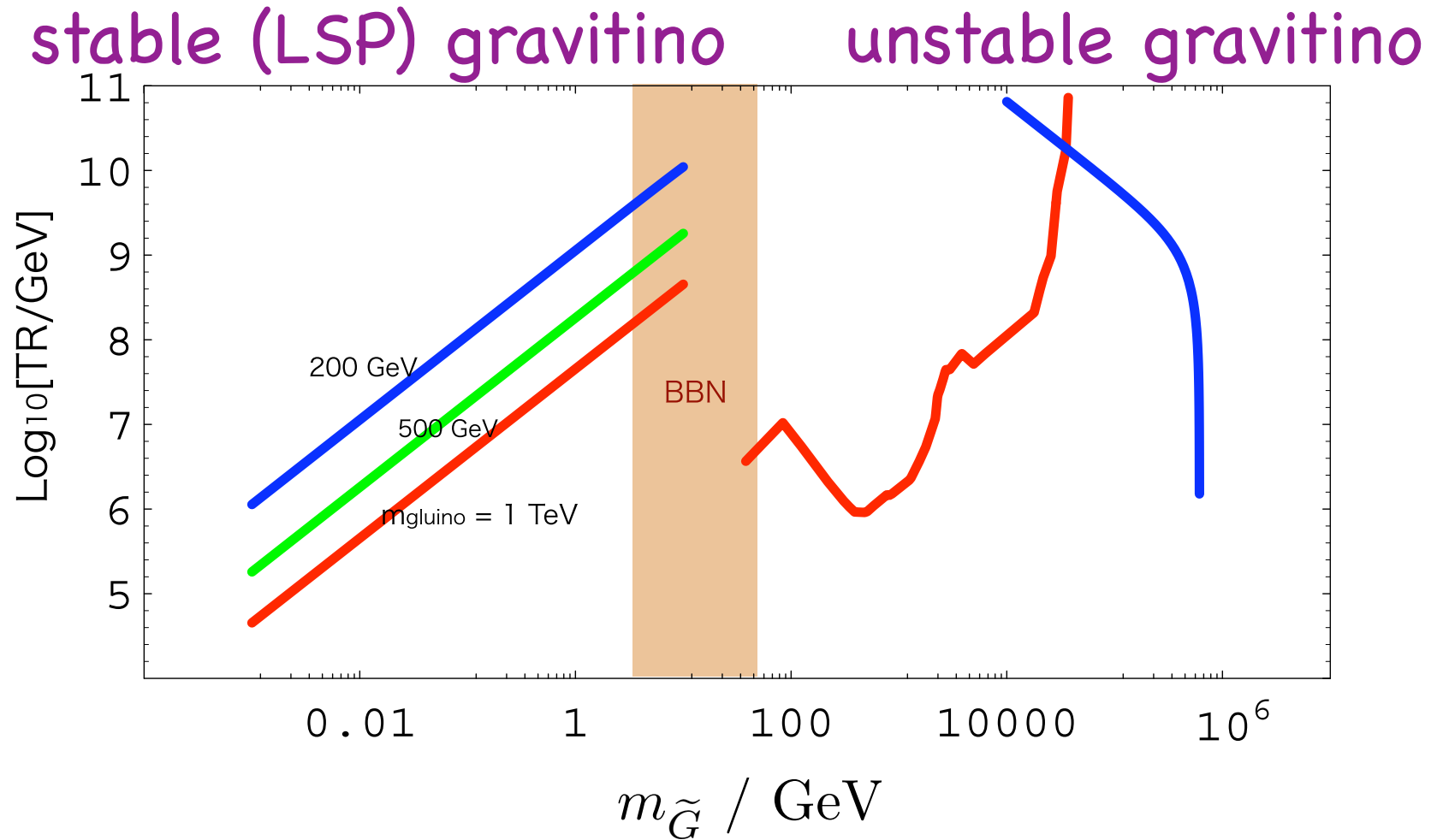
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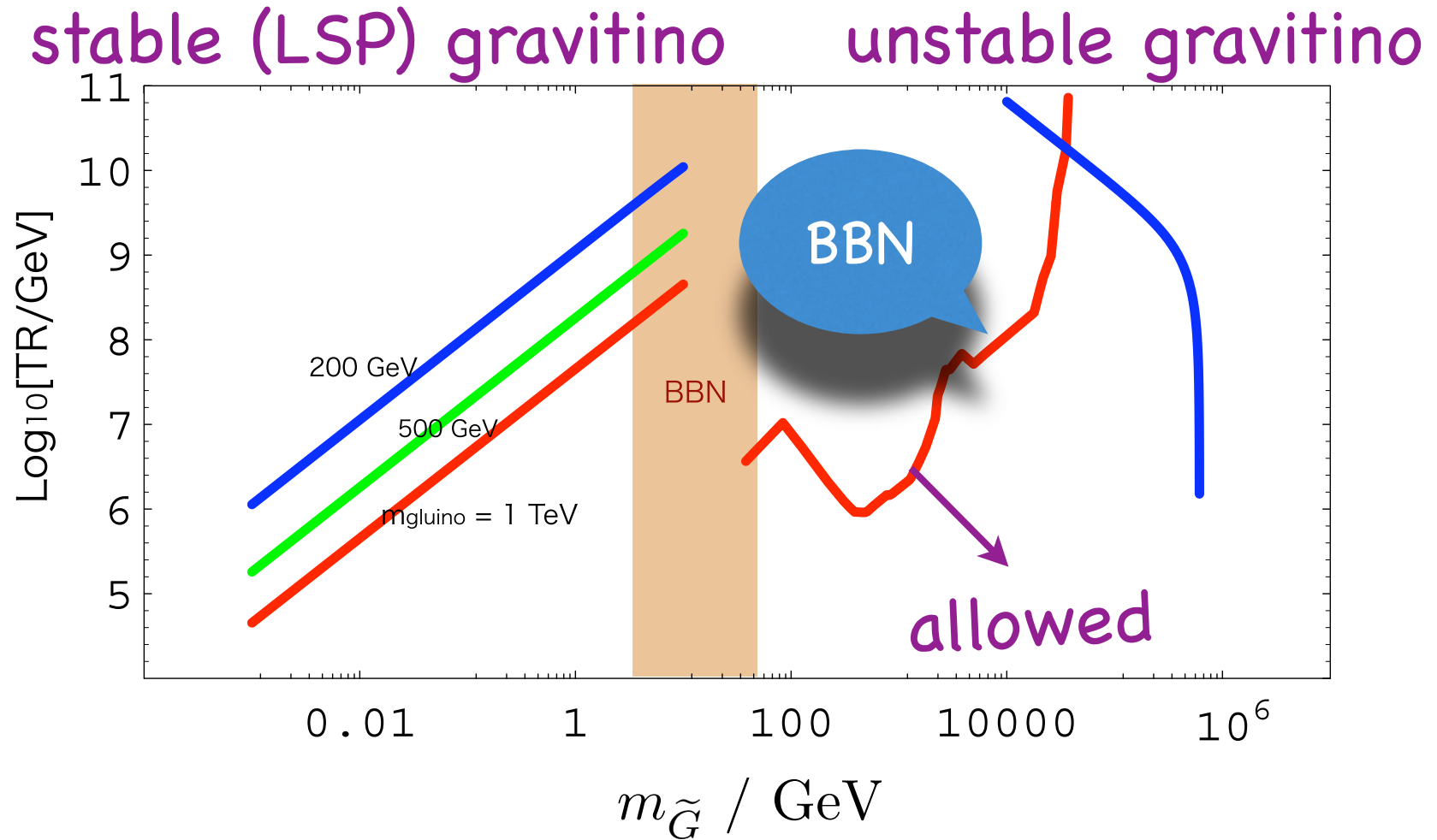
Gravitino Problems



(NOTE: precise line positions in this figure may be out-dated.)

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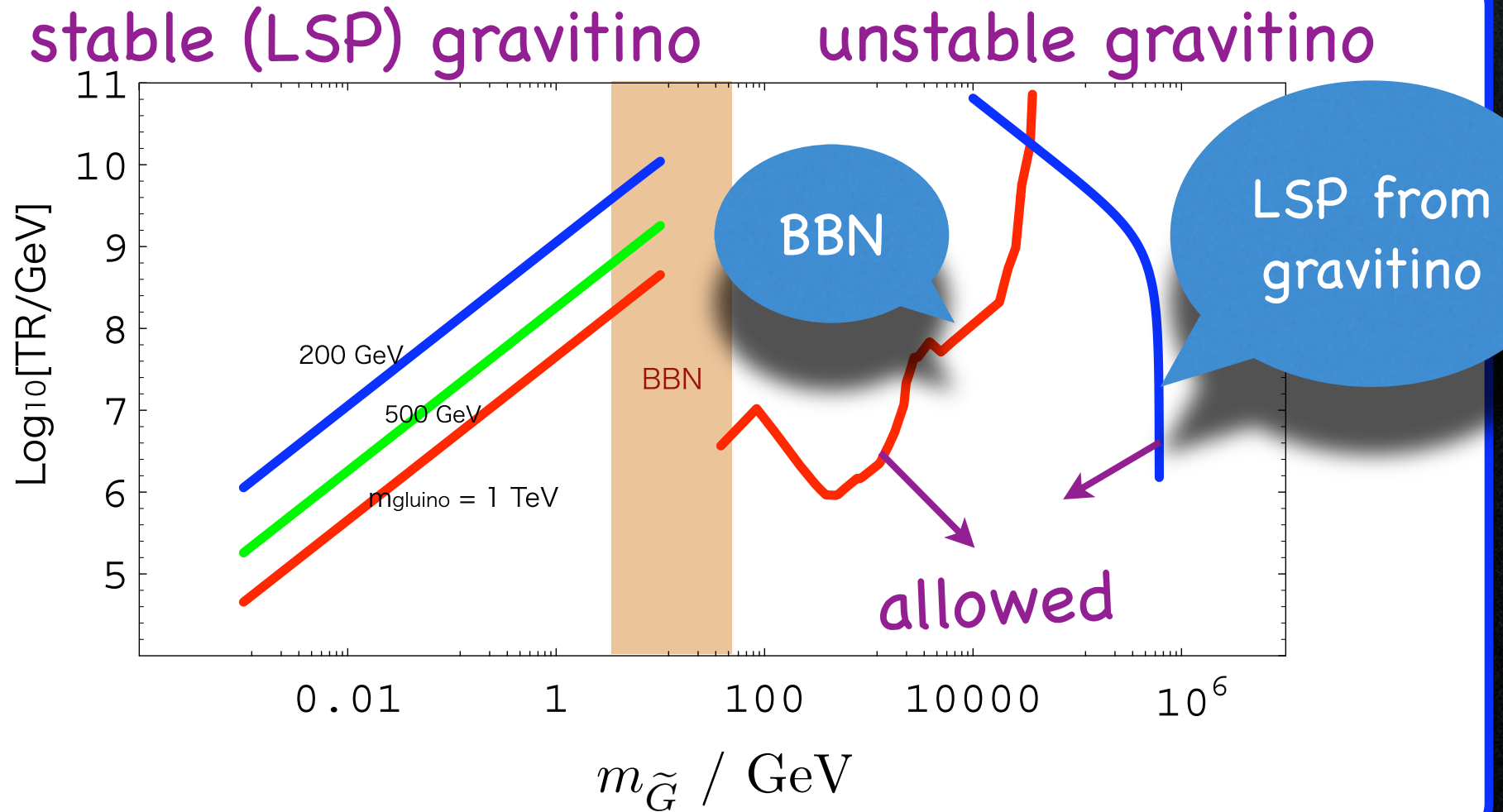
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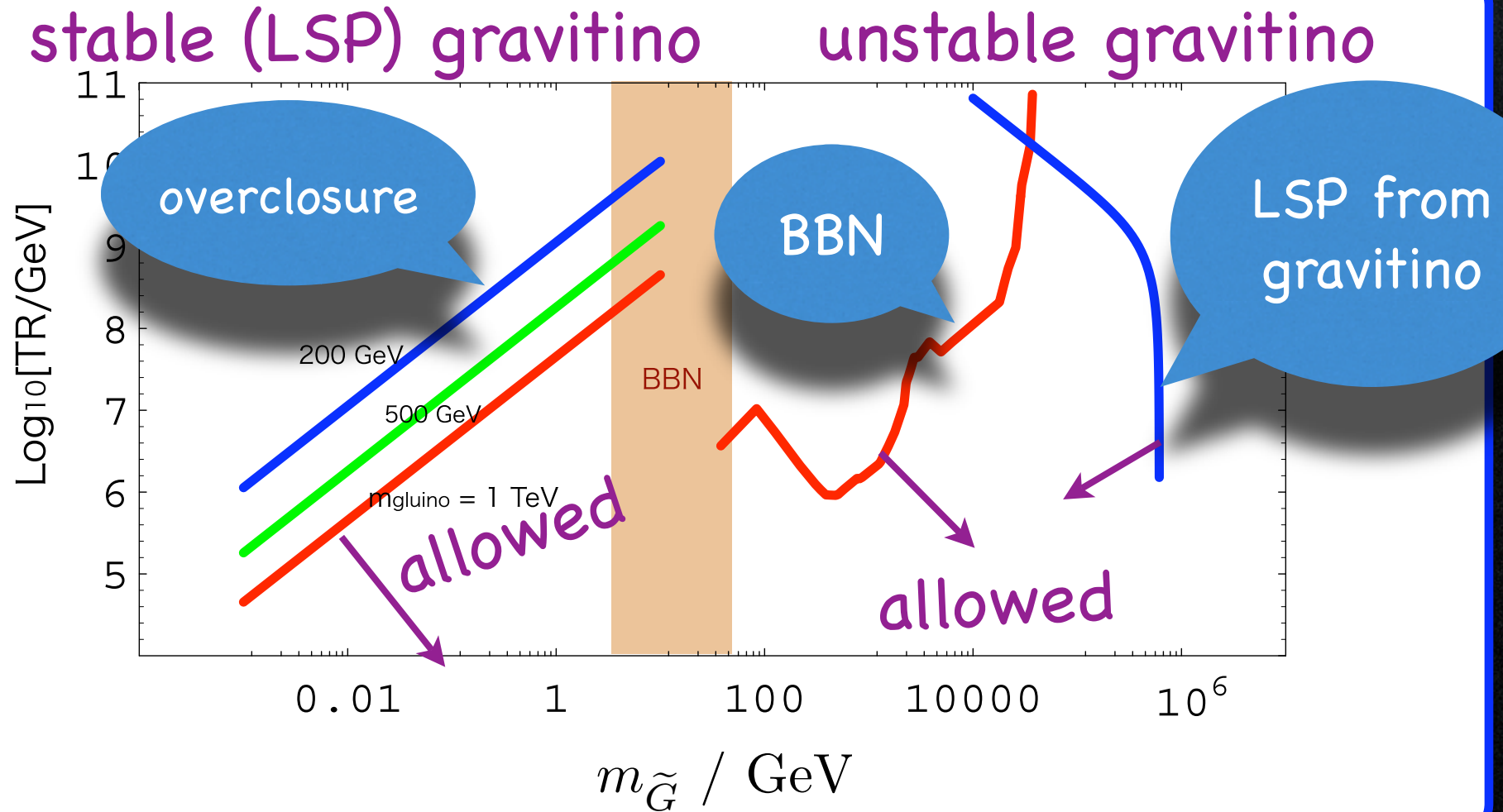
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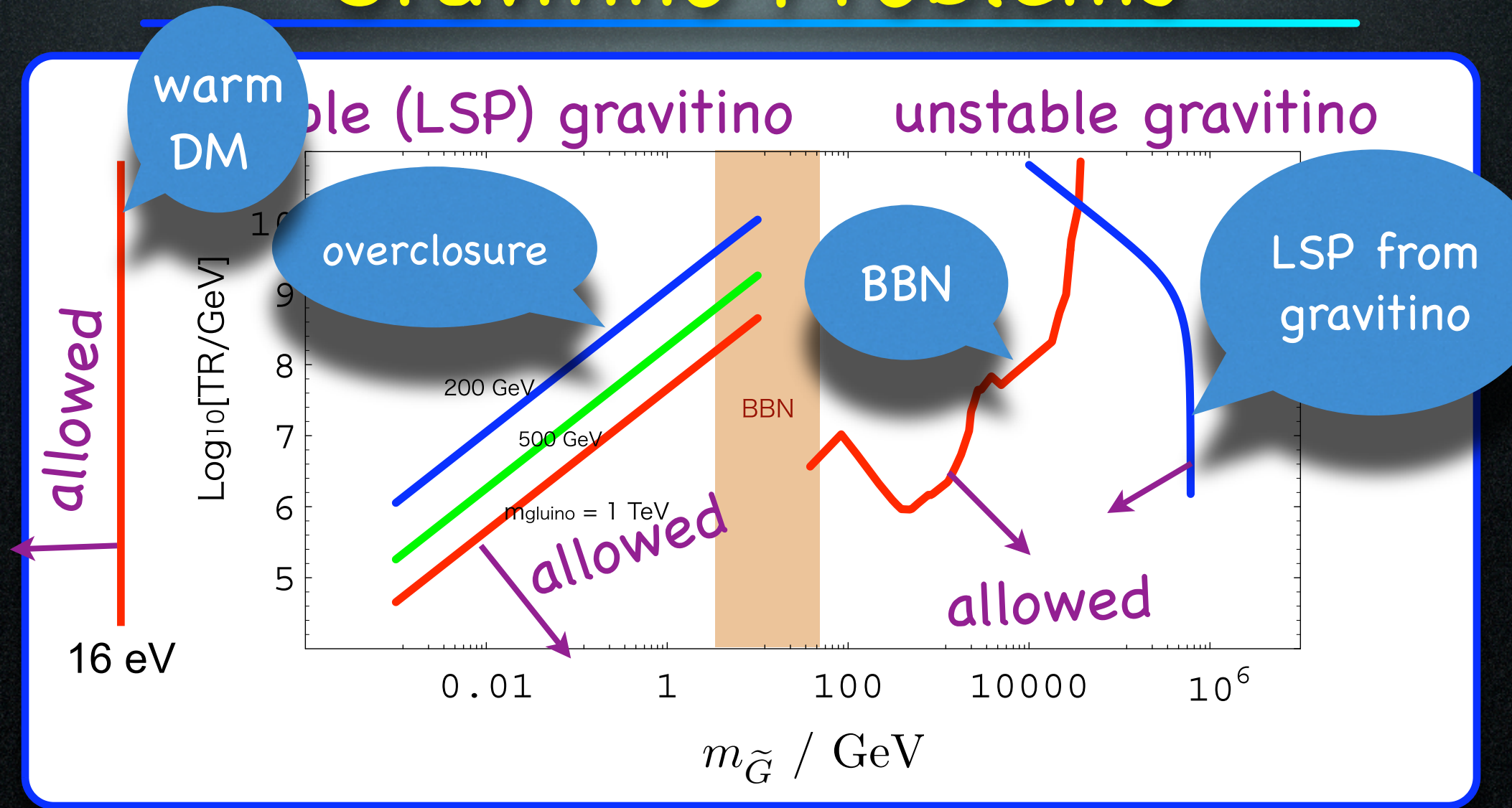
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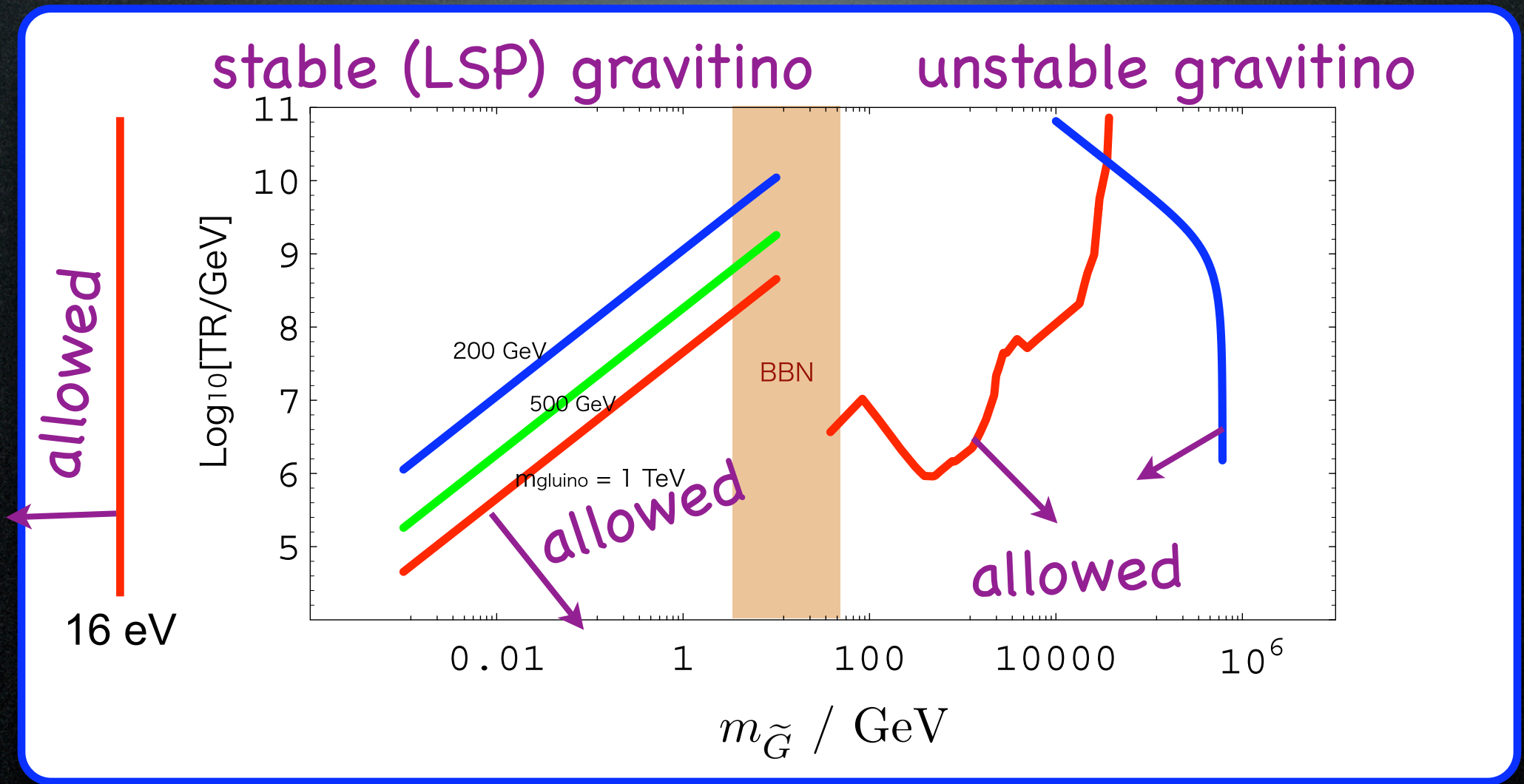
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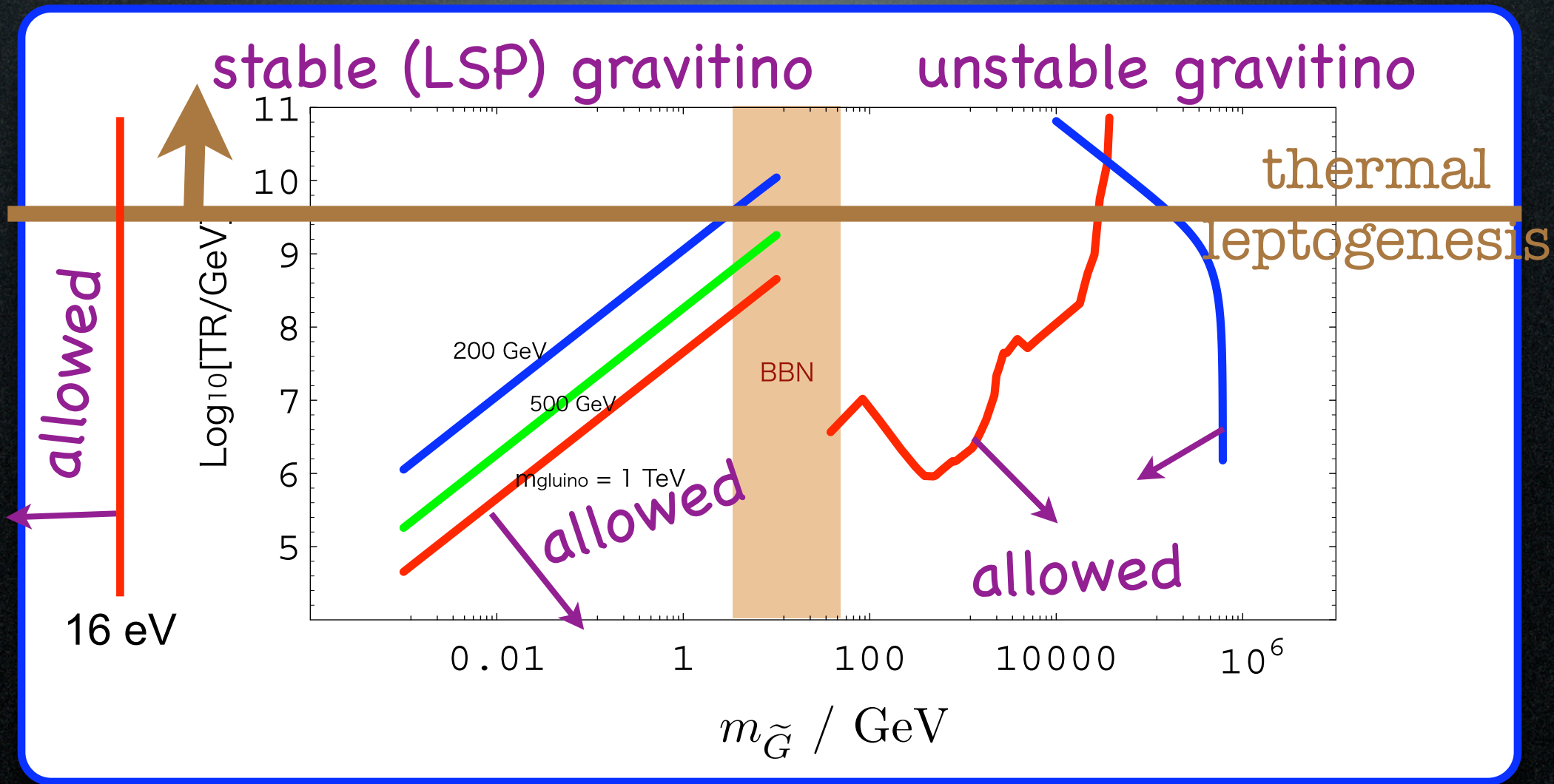
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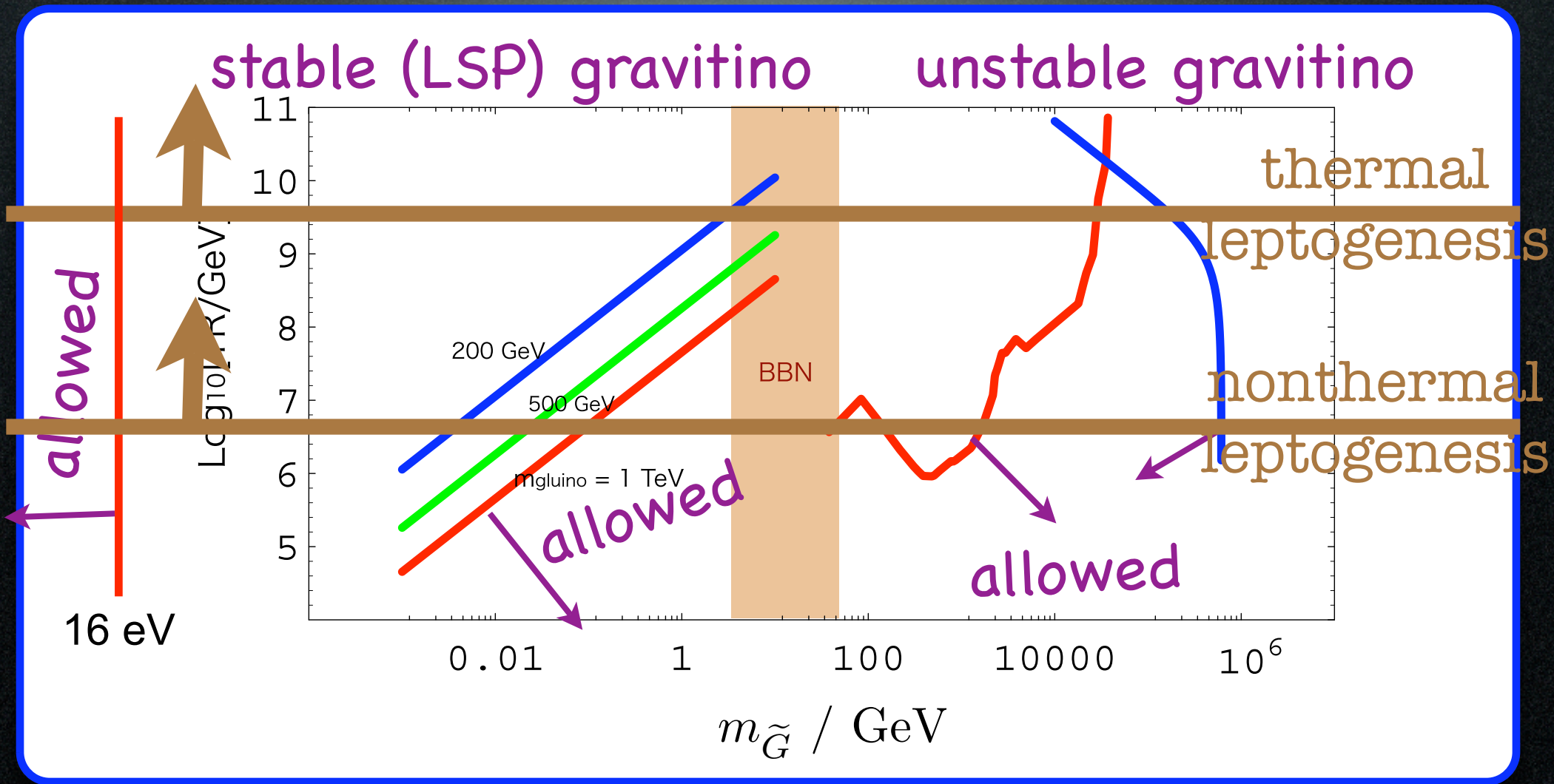
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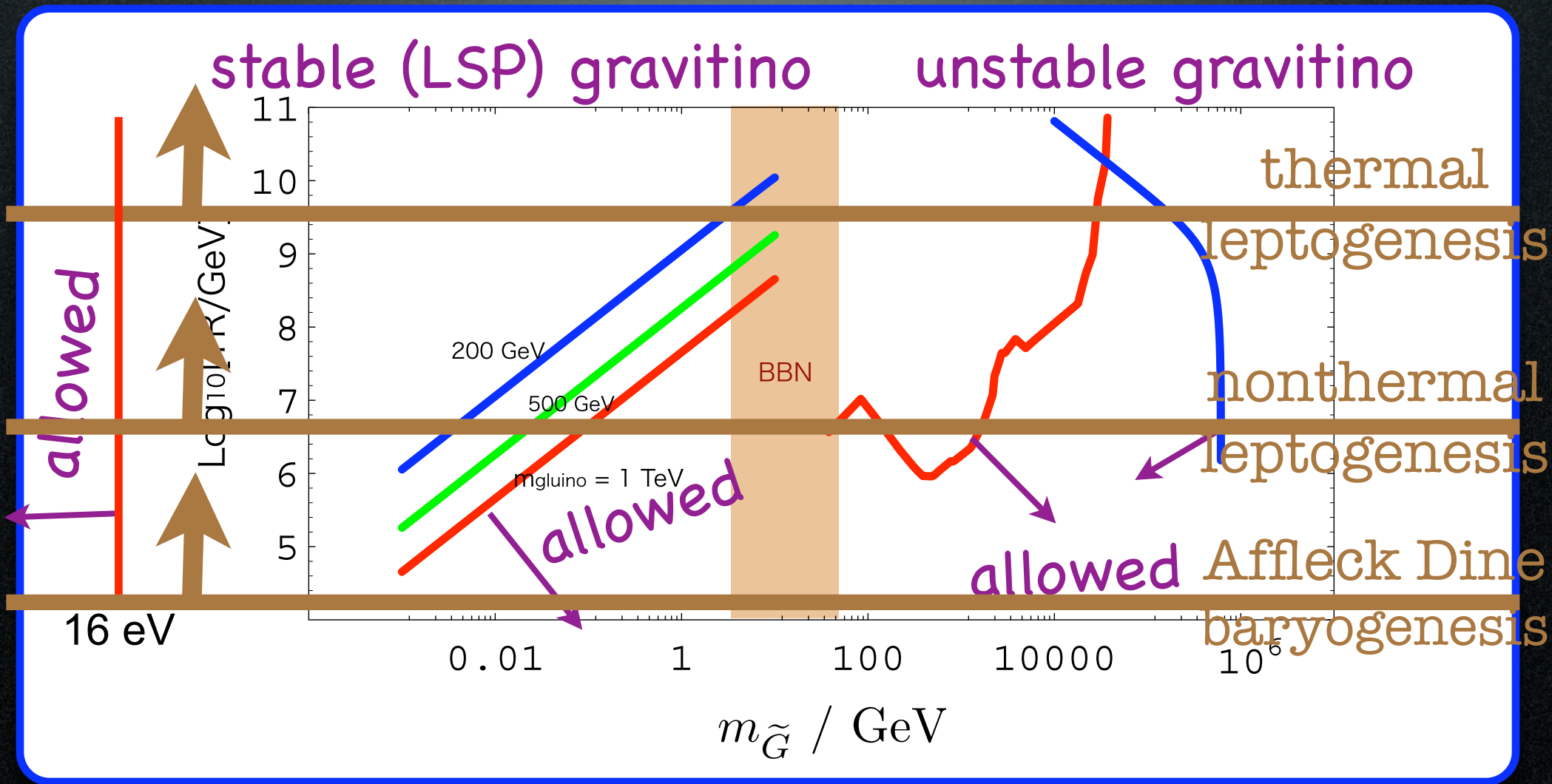
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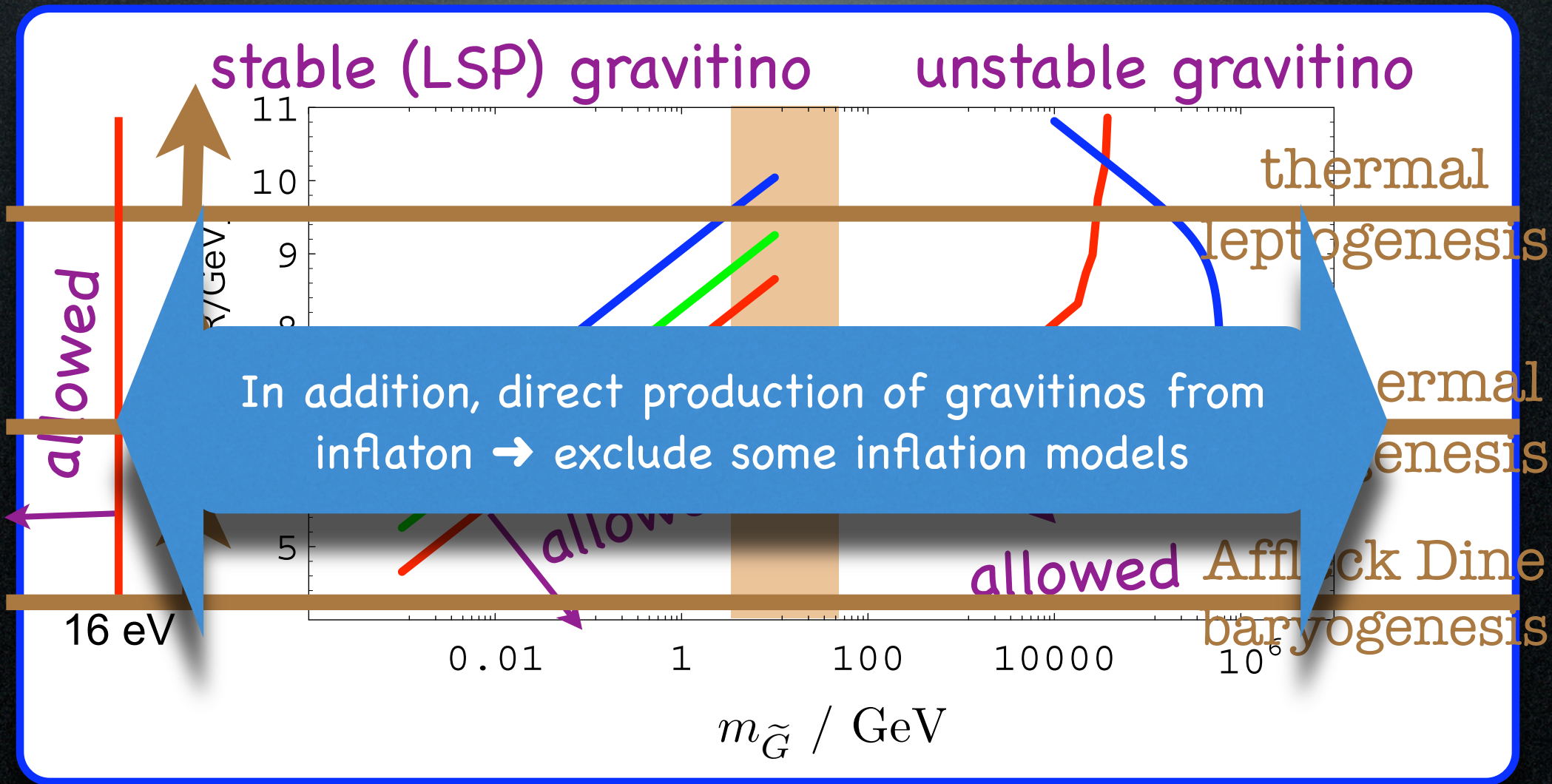
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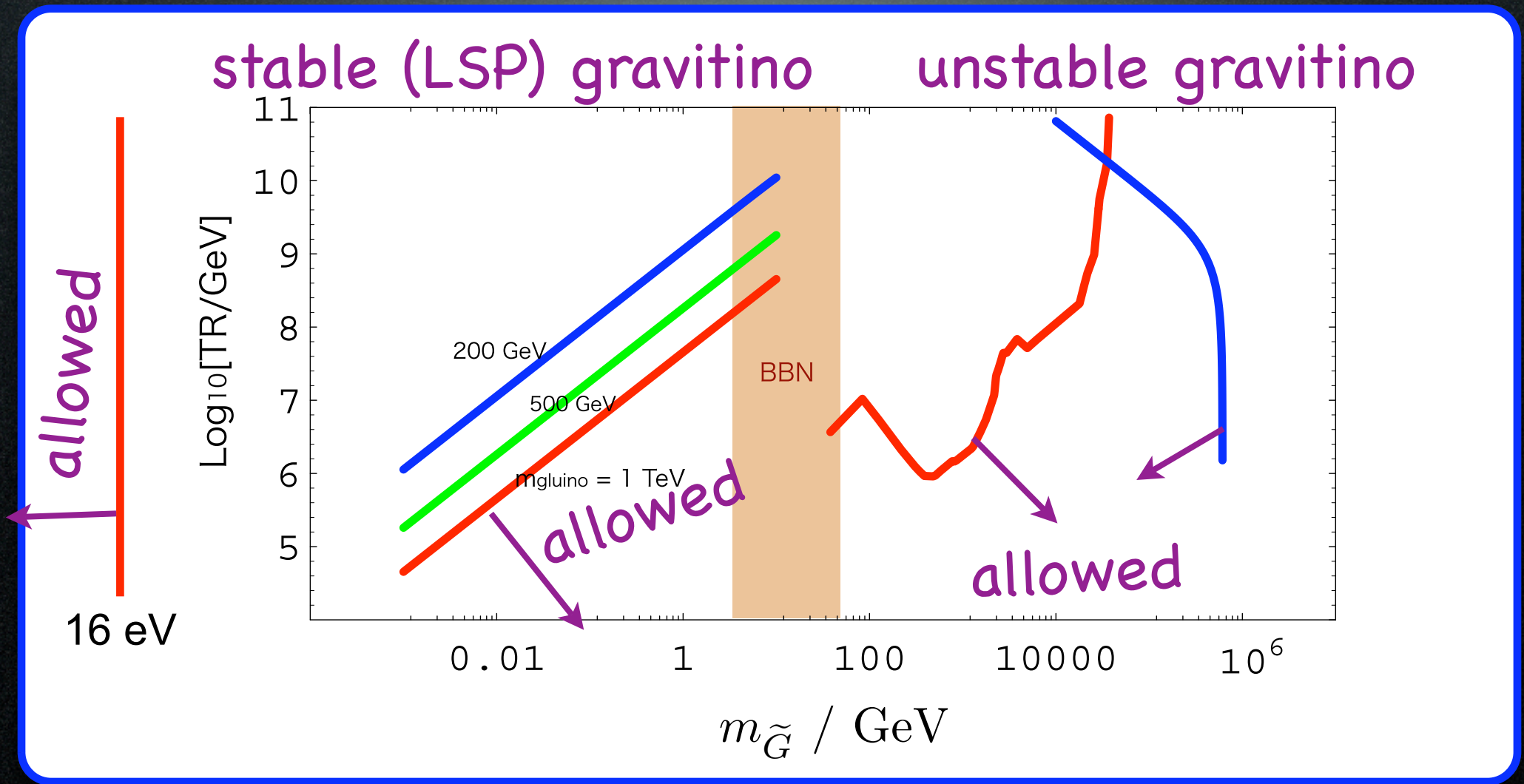
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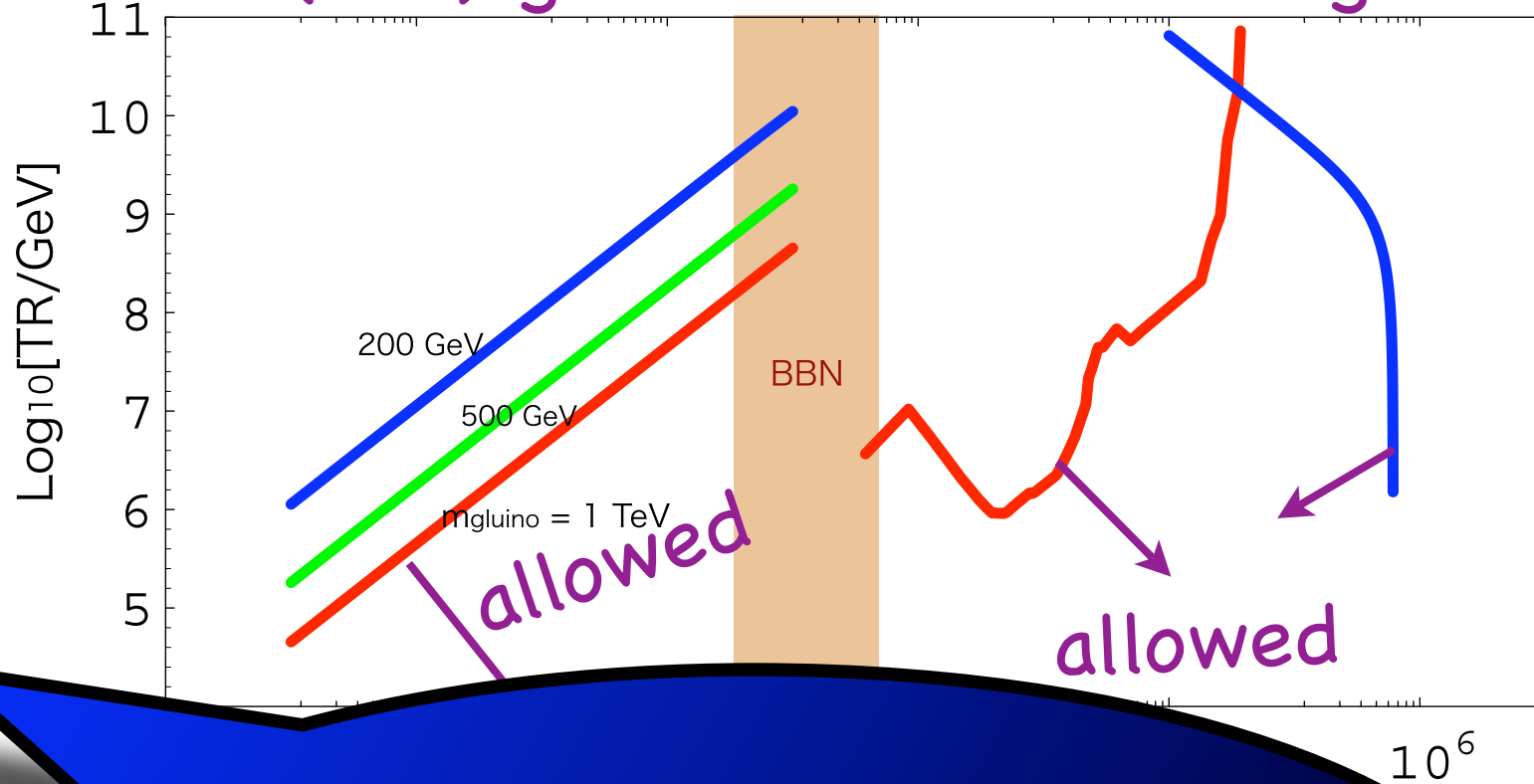
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Gravitino Problems

stable (LSP) gravitino

unstable gravitino



This region is completely free from cosmological gravitino problems!!

(this is out-dated.)

Main Message

SUSY models with an
ultralight gravitino is interesting!

$$(m_{\tilde{G}} \lesssim 10 \text{ eV})$$

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- LSP (gravitino) \neq CDM (too light \rightarrow hot DM), but...
 $m_{\tilde{G}} \sim 10 \text{ eV} \implies F = \Omega^2 \sim (100 \text{ TeV})^2$
100 TeV DM \rightarrow natural thermal relic DM **if strongly interacting**

DM may be **100 TeV composite "baryon"** made from **strongly self-interacting** hidden-sector/messenger particles

Dimopoulos, Giudice, Pomarol '96 / KH, Shirai, Yanagida '07

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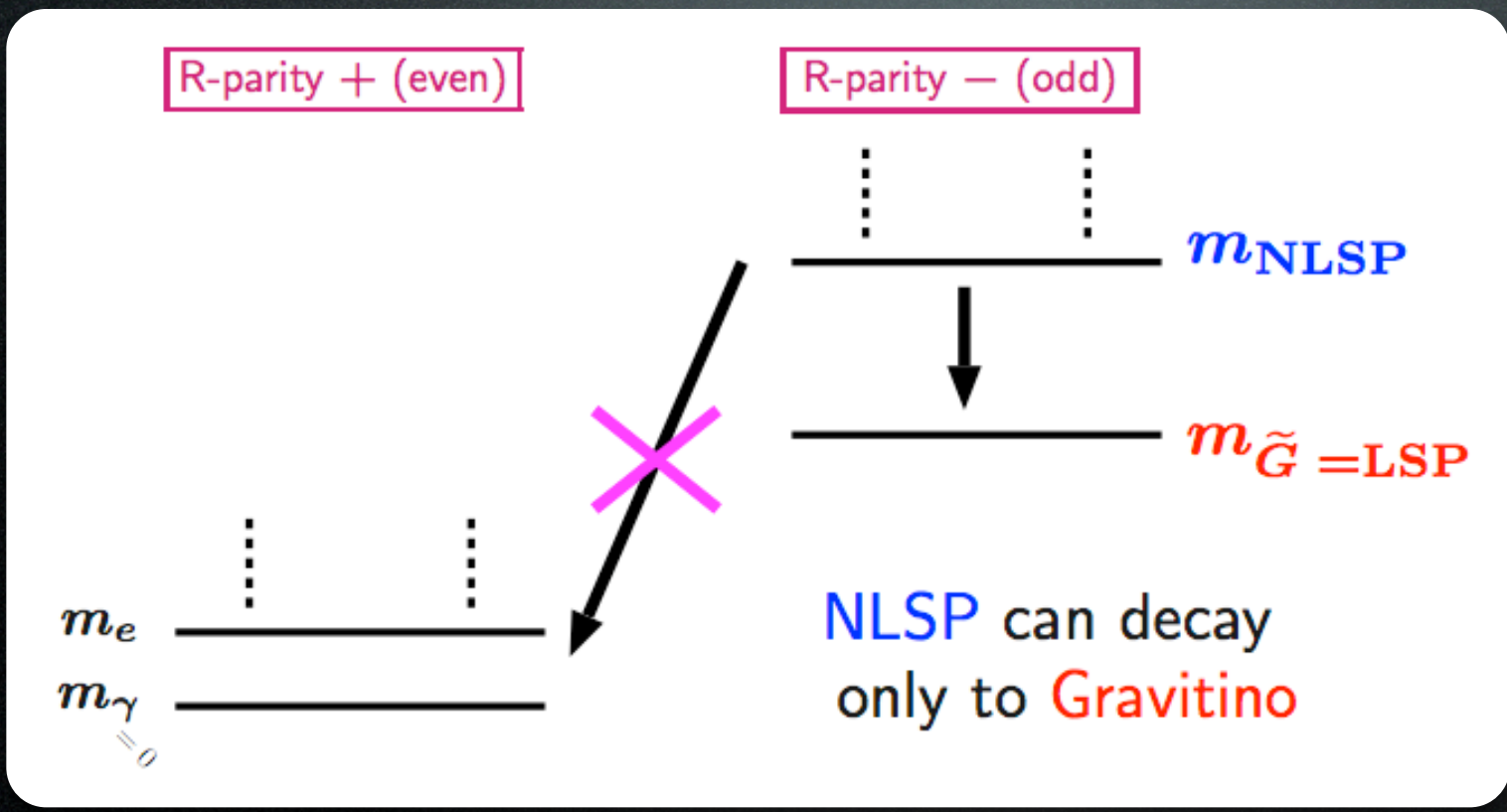
- **No Cosmological Problem!** at all!
- LSP (gravitino) \neq DM, but a natural DM candidate exist.

Can it be **tested at LHC?**

(Can the **gravitino mass = SUSY breaking scale**
be **determined?**)

NLSP (Next-to-Lightest SUSY Particle)

In Gravitino LSP scenario, the NLSP decay always include the gravitino.




Interaction

$$\sim \frac{1}{F} \sim \frac{1}{M_P m_{\tilde{G}}}$$

NLSP (Next-to-Lightest SUSY Particle)

In Gravitino LSP scenario, the NLSP decay always include the gravitino.

For a slepton NLSP,.....

$$\Gamma(\tilde{\tau} \rightarrow \tilde{G}\tau) \simeq \frac{m_{\tilde{\tau}}^5}{48\pi m_{\tilde{G}}^2 M_{\text{pl}}^2} \left(1 - \frac{m_{\tilde{G}}^2}{m_{\tilde{\tau}}^2}\right)^4$$


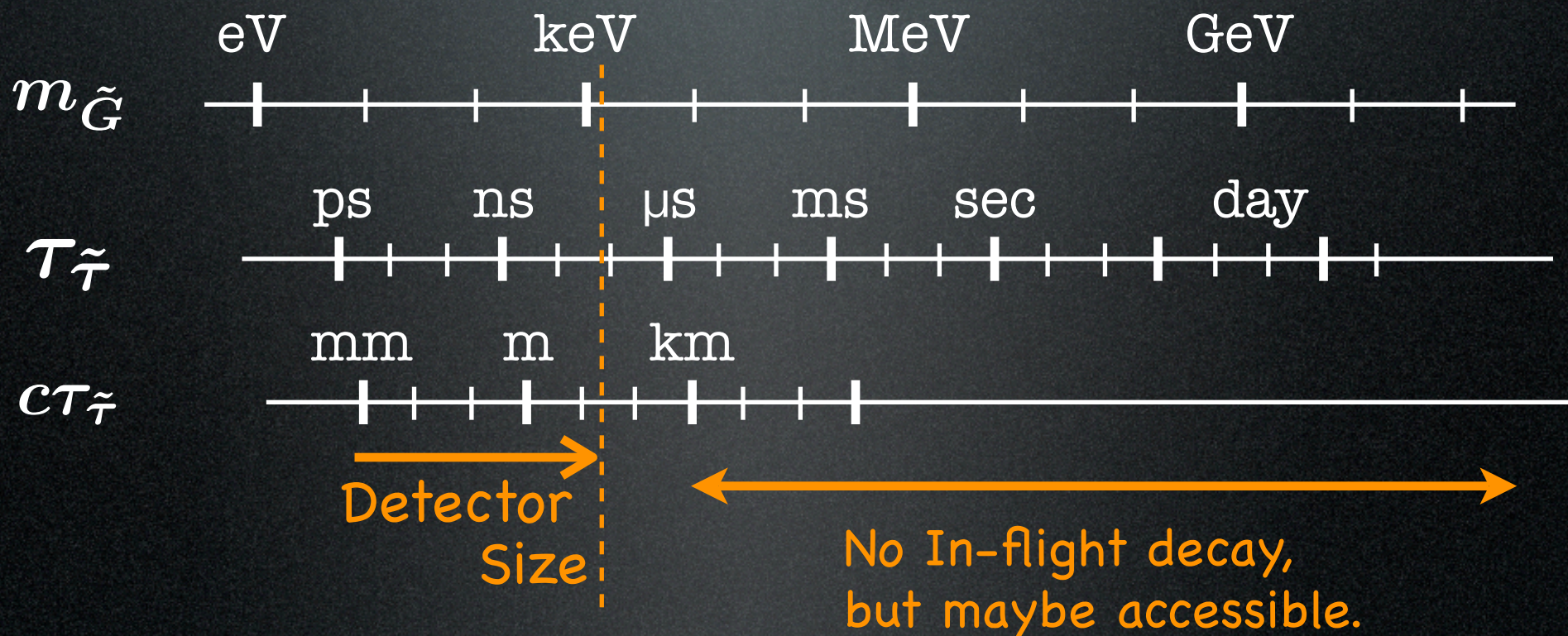
The gravitino mass may be determined by measuring the NLSP decay rate! However,.....

charged sleptons @ LHC

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Lifetime (decay length) of NLSP stau

e.g., for $m_{\tilde{\tau}} = 100 \text{ GeV}$,



charged

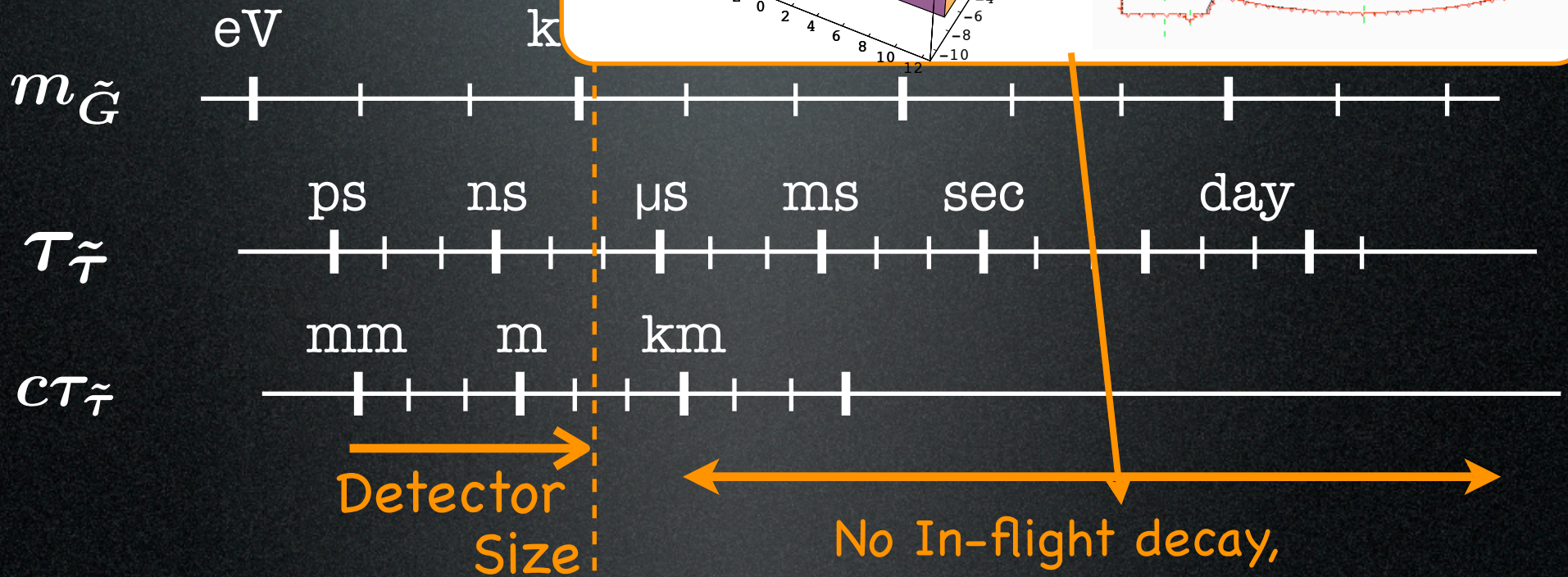
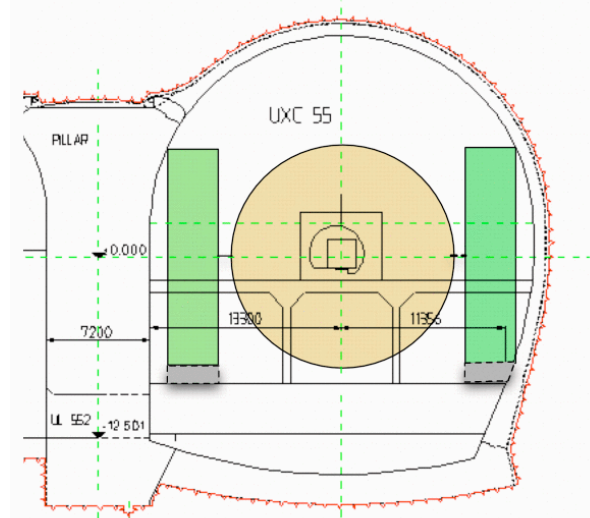
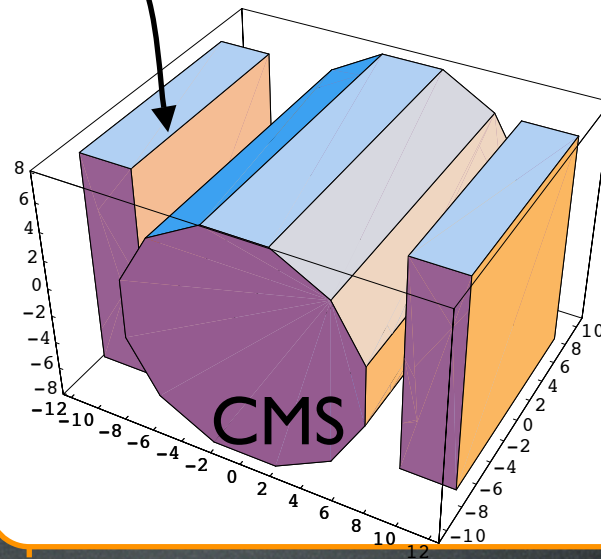
$$\Gamma(\tilde{\tau} \rightarrow \tilde{G}\tau) \simeq$$

Lifetime (decay length)

stopper-detector

KH, Kuno, Nakaya, Nojiri '04

KH, Nojiri, DeRoeck '06



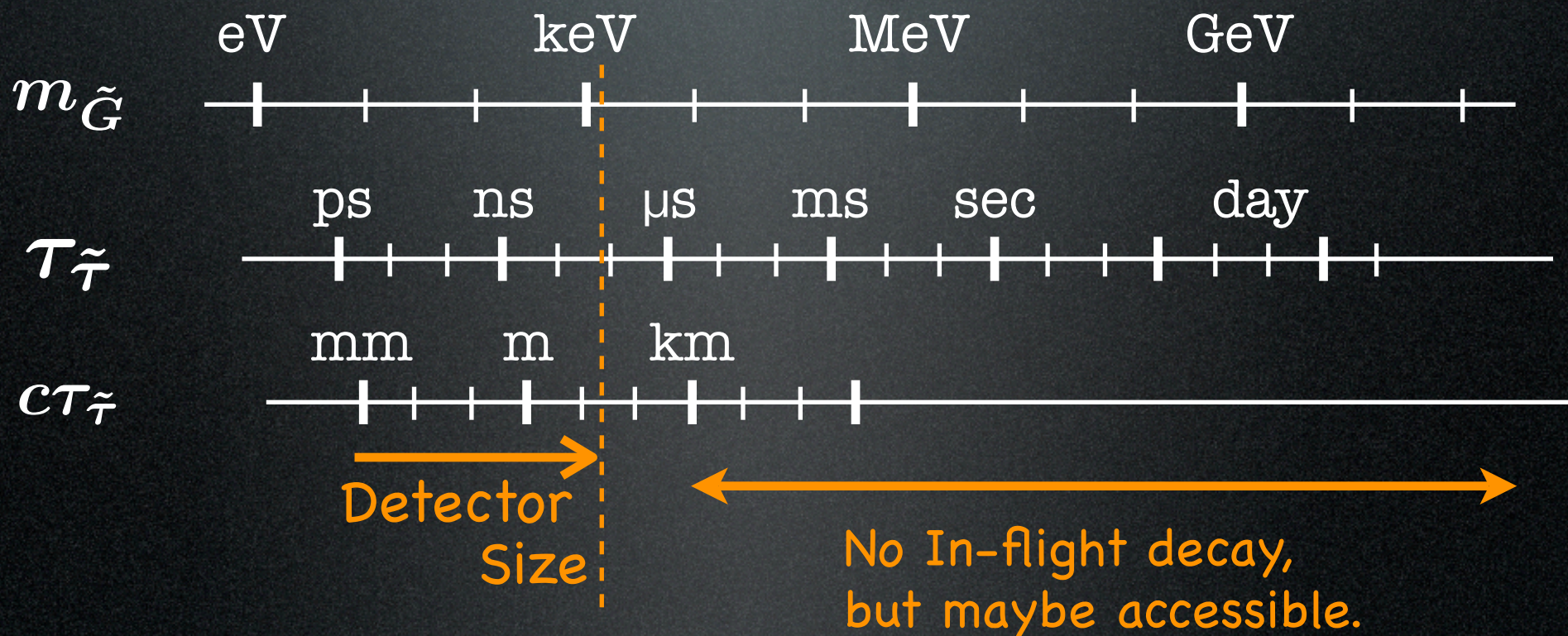
No In-flight decay,
but maybe accessible.

charged sleptons @ LHC

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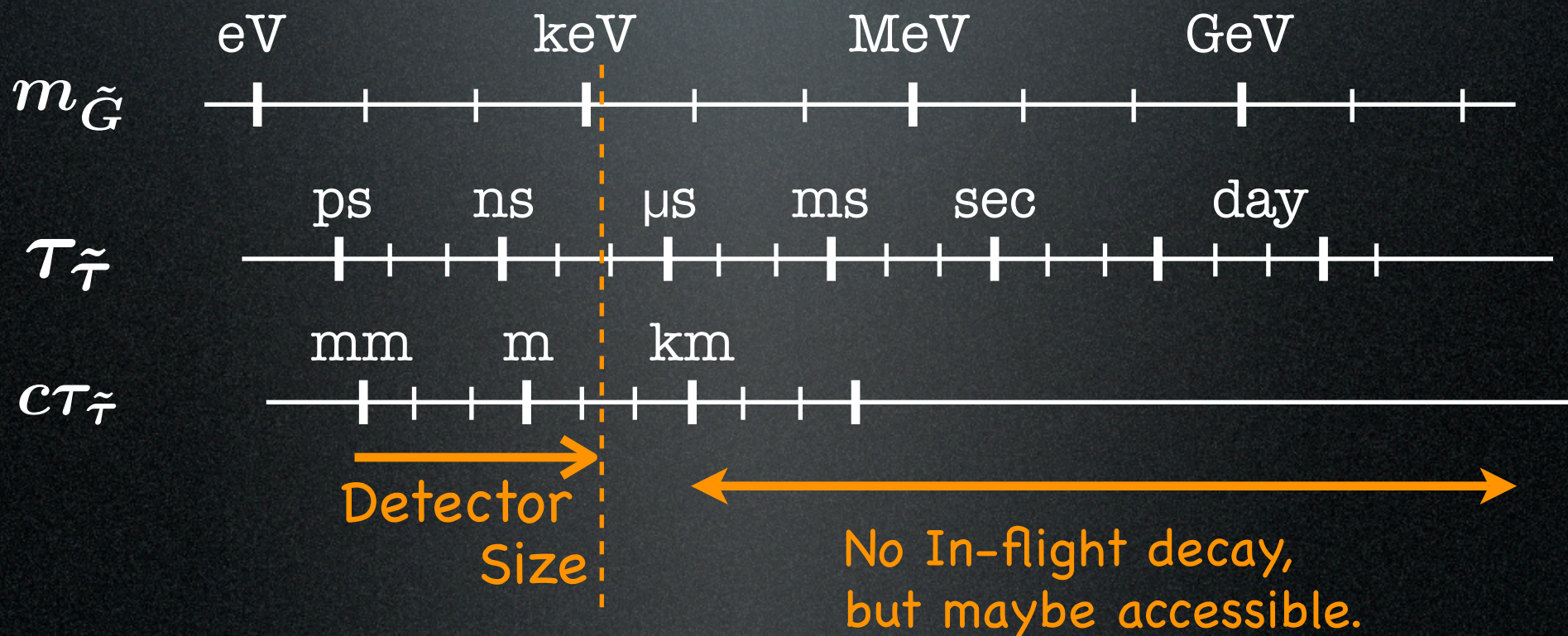


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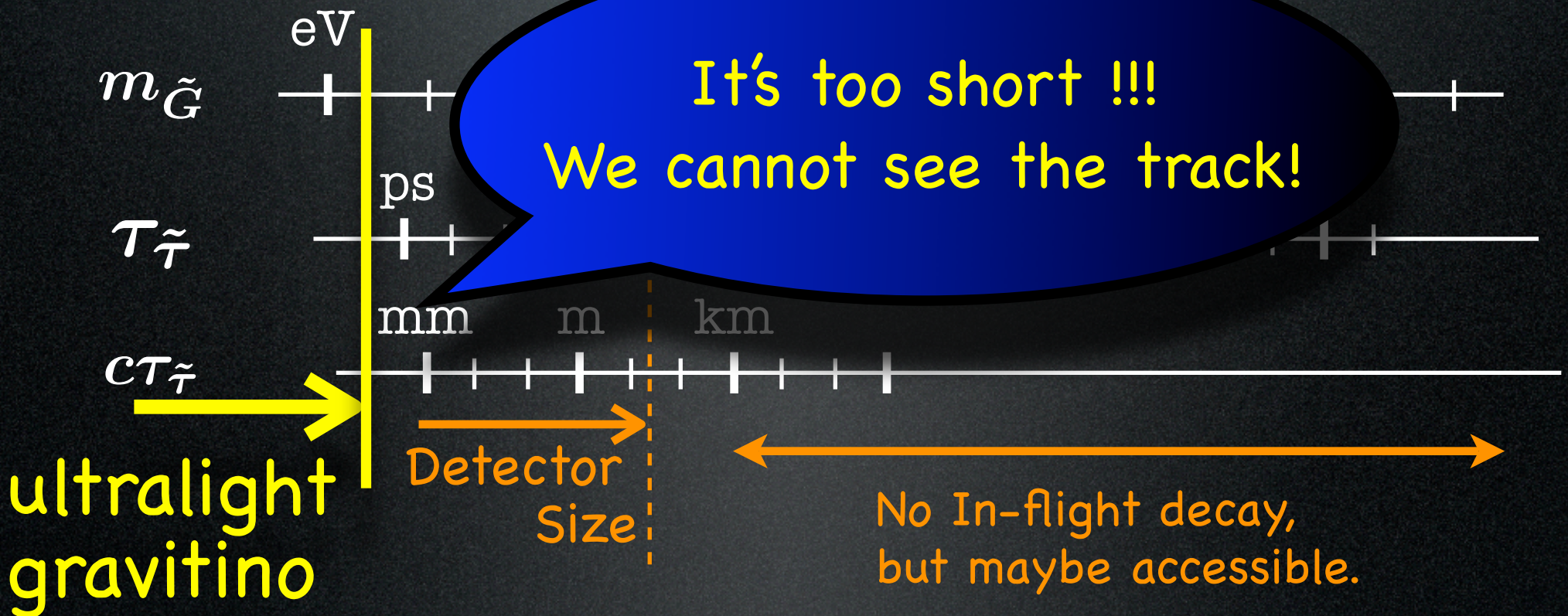


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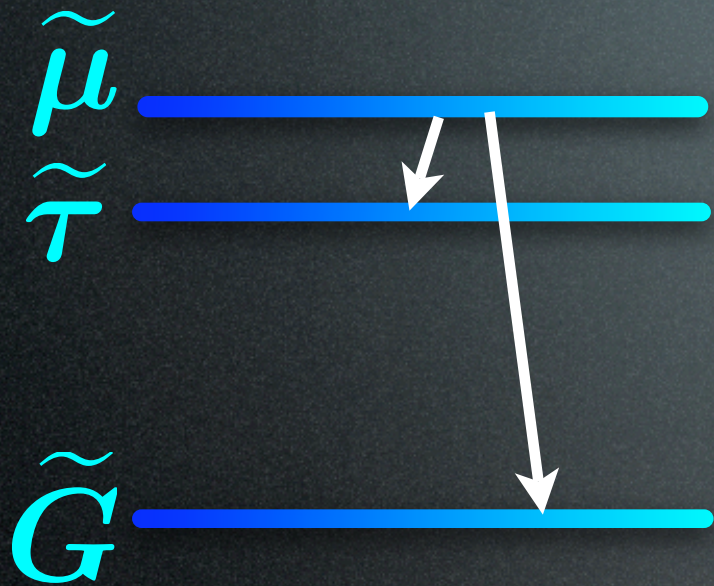
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Idea: Look at NNLSP sleptons



two decay modes

$$\tilde{\mu} \rightarrow \mu \tau \tilde{\tau} : \Gamma_{3\text{body}}$$

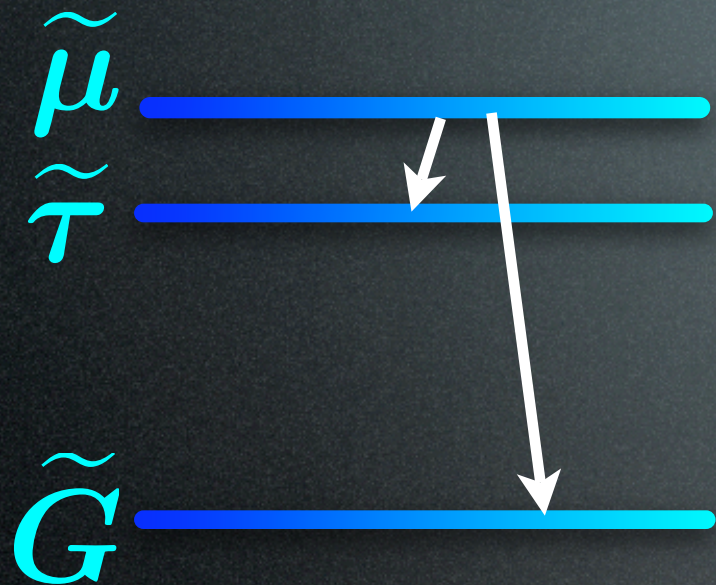
$$\tilde{\mu} \rightarrow \mu \tilde{G} : \Gamma_{2\text{body}}$$

$$m_{\tilde{G}}^2 \propto \frac{1}{\Gamma_{2\text{body}}}$$

difficult to measure

(too short decay length)

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$$m_{\tilde{G}}^2 \propto \frac{1}{\Gamma_{2\text{body}}} = \frac{1}{\Gamma_{3\text{body}}} \left(\frac{\Gamma_{3\text{body}}}{\Gamma_{2\text{body}}} \right)$$

difficult to measure

(too short decay length)

calculable if other

SUSY masses are known

measurable!

Idea: Look at NNLSP sleptons

$\tilde{\mu}$
 $\tilde{\tau}$
 \tilde{G}

$$\Gamma(\tilde{\ell} \rightarrow \ell + \tilde{\tau} + \tau) \approx \frac{2}{15\pi} \left(\frac{\alpha_{\text{EM}}}{\cos^2 \theta_W} \right)^2 m_{\tilde{\ell}} \frac{(1 + r_{\tilde{\chi}_1^0}^2)}{(r_{\tilde{\chi}_1^0}^2 - 1)^2} \left(\frac{\Delta m}{m_{\tilde{\ell}}} \right)^5$$

$$= 4.4 \text{ eV} \left(\frac{m_{\tilde{\ell}}}{100 \text{ GeV}} \right)^{-4} \left(\frac{\Delta m}{10 \text{ GeV}} \right)^5 \frac{1 + r_{\tilde{\chi}_1^0}^2}{(r_{\tilde{\chi}_1^0}^2 - 1)^2}$$

$r_{\tilde{\chi}_1^0} = m_{\tilde{\chi}_1^0} / m_{\tilde{\ell}}$. assumed $\Delta m = m_{\tilde{\ell}} - m_{\tilde{\tau}} \gg m_{\tau}$

idea

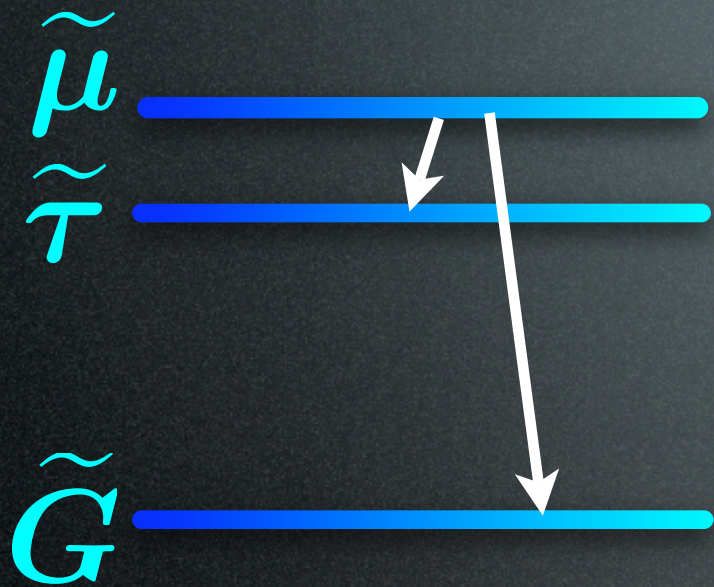
$$m_{\tilde{G}}^2 \propto \frac{1}{\Gamma_{\text{2body}}} = \frac{1}{\Gamma_{\text{3body}}} \left(\frac{\Gamma_{\text{3body}}}{\Gamma_{\text{2body}}} \right)$$

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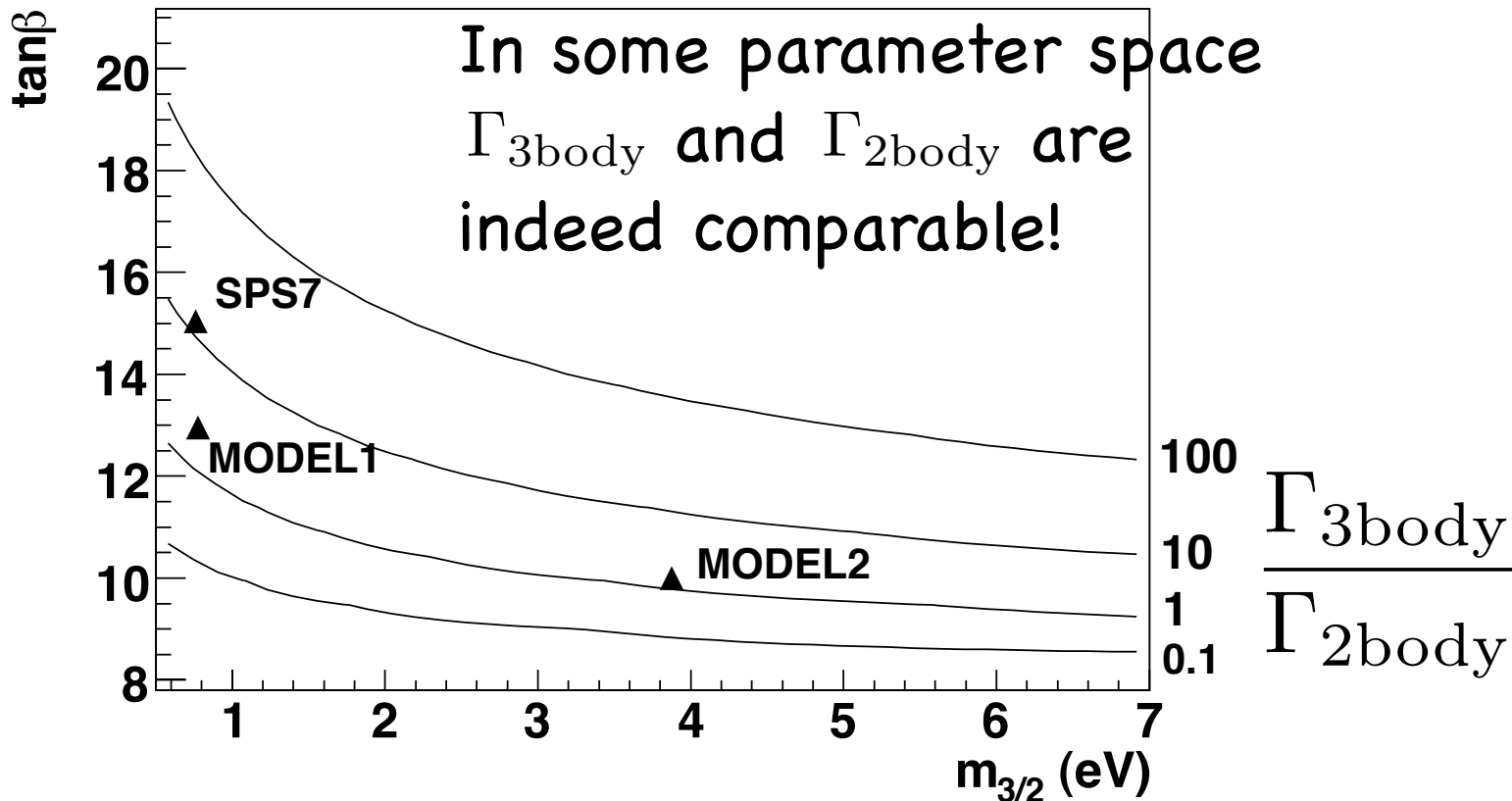
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Idea:

$\tilde{\mu}$

$\tilde{\tau}$

\tilde{G}



idea

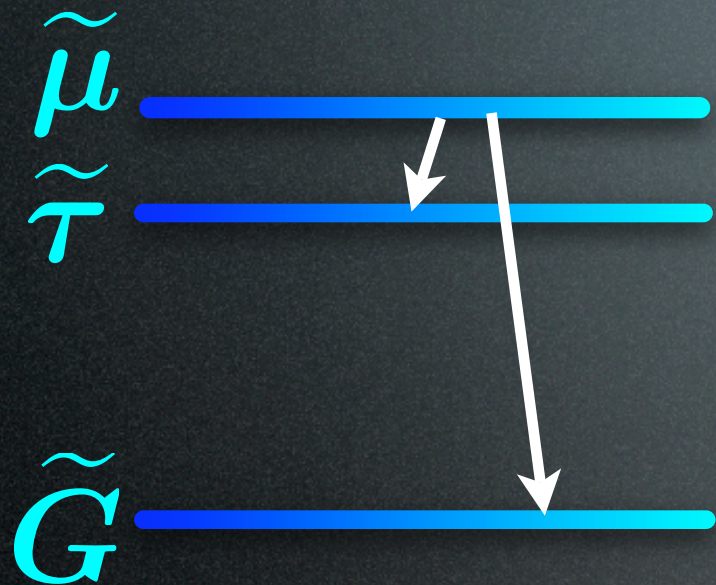
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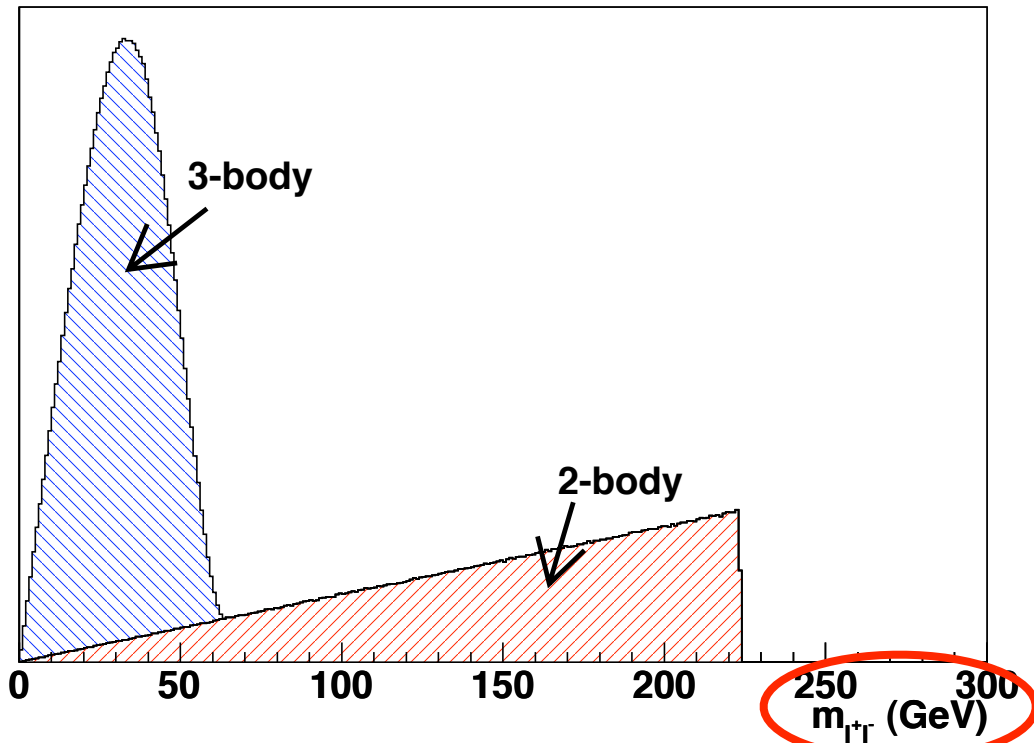
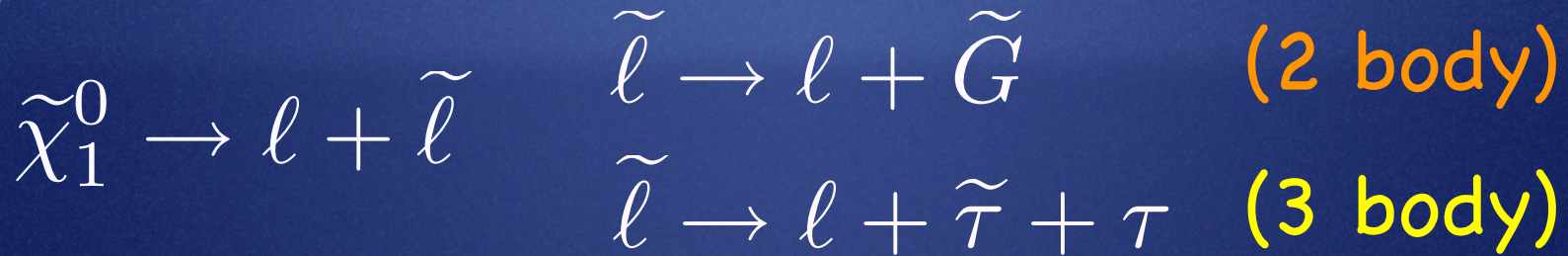
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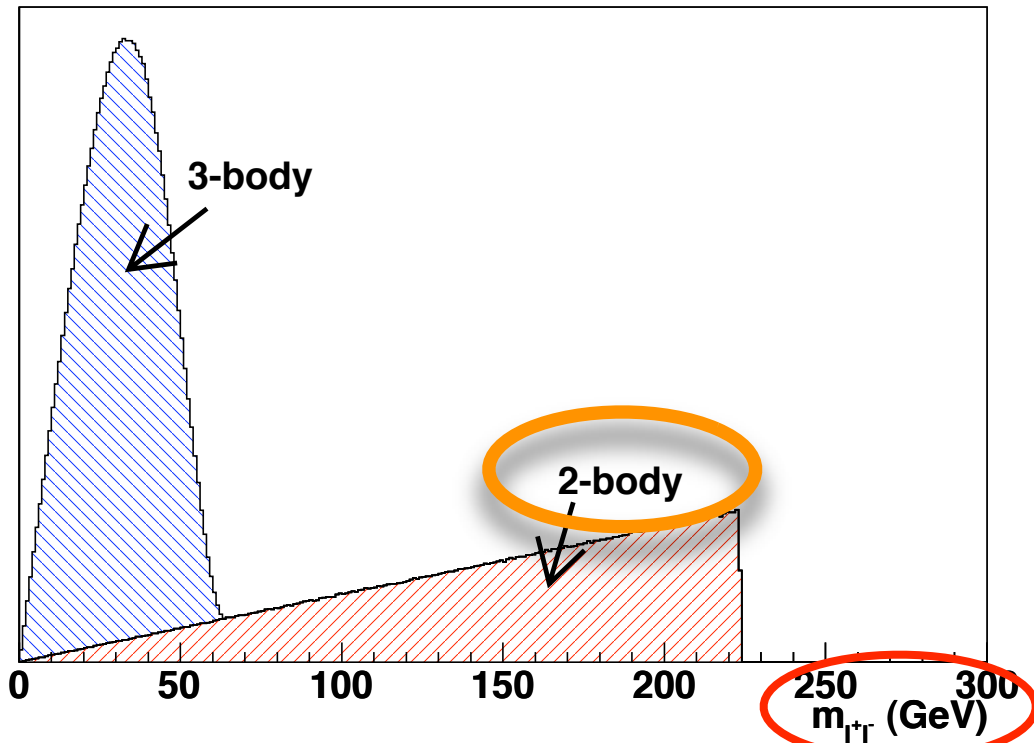
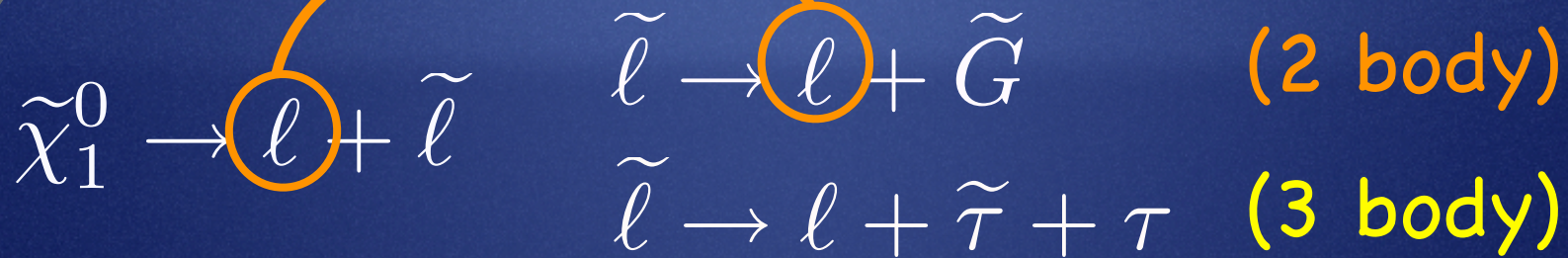
→ look at the dilepton invariant mass $m_{\ell\ell} = (p_{\ell,1} + p_{\ell,2})^2$



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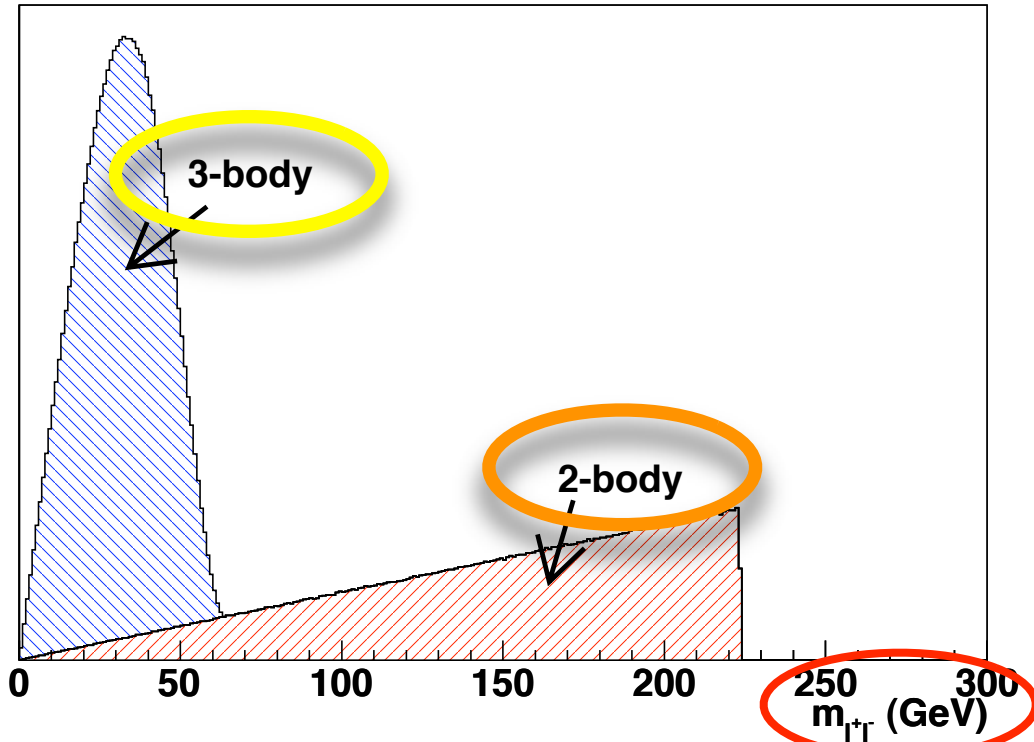
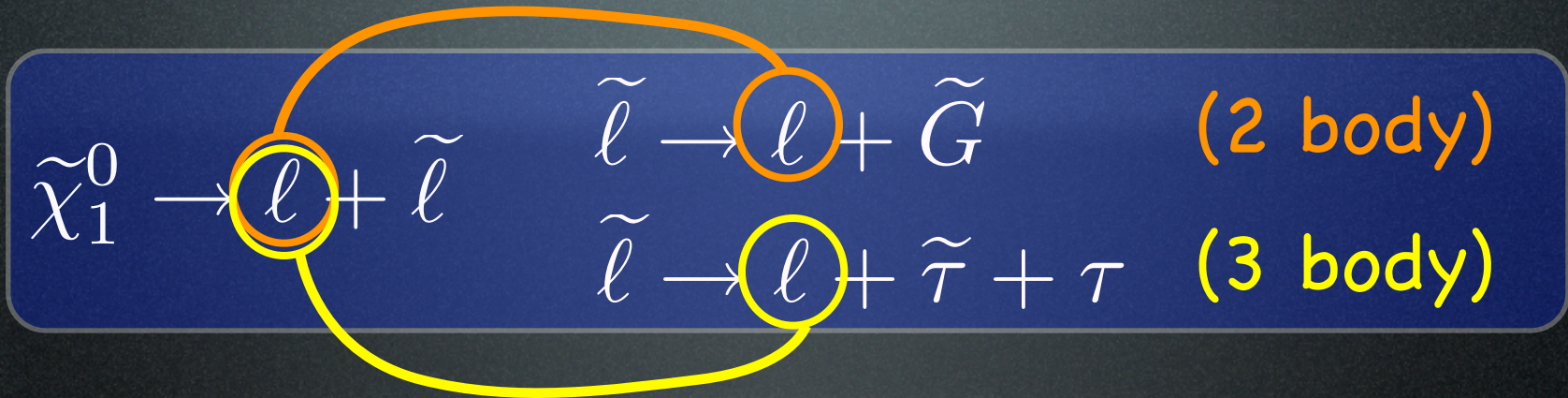
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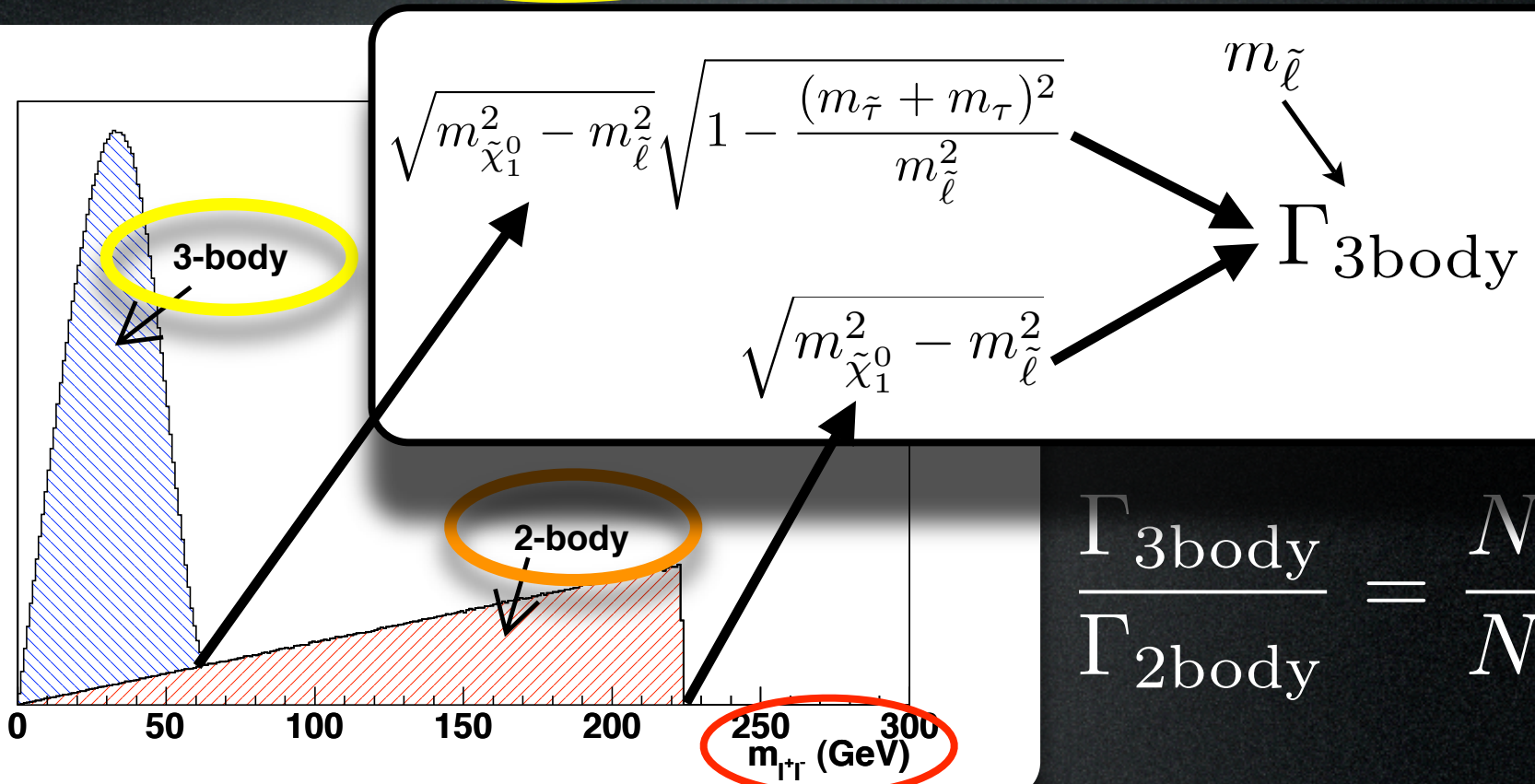
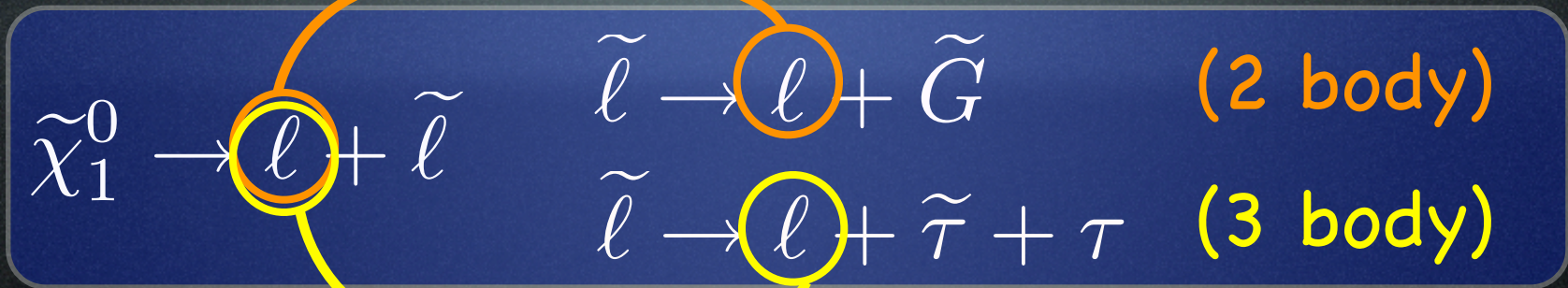
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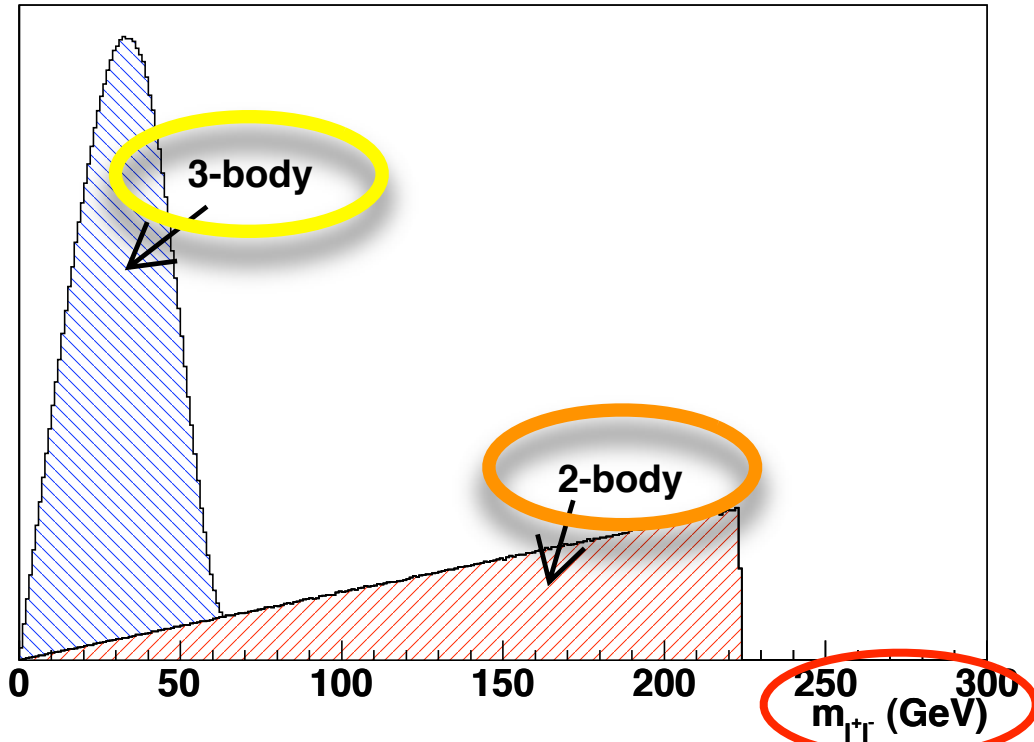
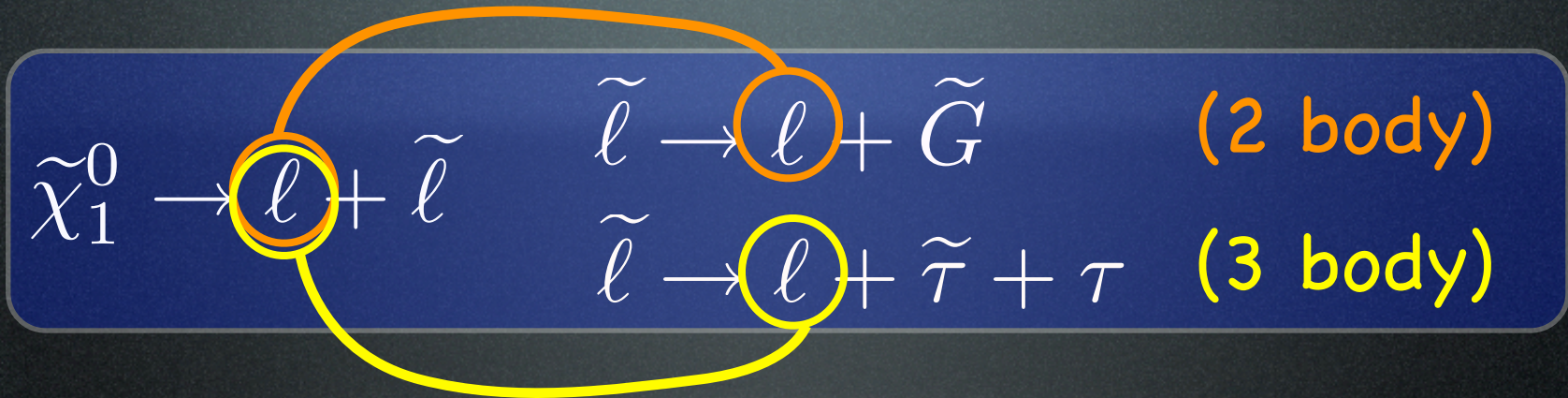
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event selection

- at least four jets with $P_T \geq 25$ GeV, where τ -jets are excluded.
- missing transverse momentum $P_{T,\text{miss}} \geq 100$ GeV.
- $M_{\text{eff}} \geq 500$ GeV, where

$$M_{\text{eff}} = \sum_{\text{jets}(\neq\tau)}^4 P_{T_j} + P_{T,\text{miss}}.$$

- two leptons with $P_T \geq 20$ GeV and $|\eta| < 2.5$.
- the dilepton mass is formed only if one of the two leptons has $P_T \geq 20$ GeV, $|\eta| < 2.5$ and the other has $P_T \geq 6$ GeV, $|\eta| < 2.5$,
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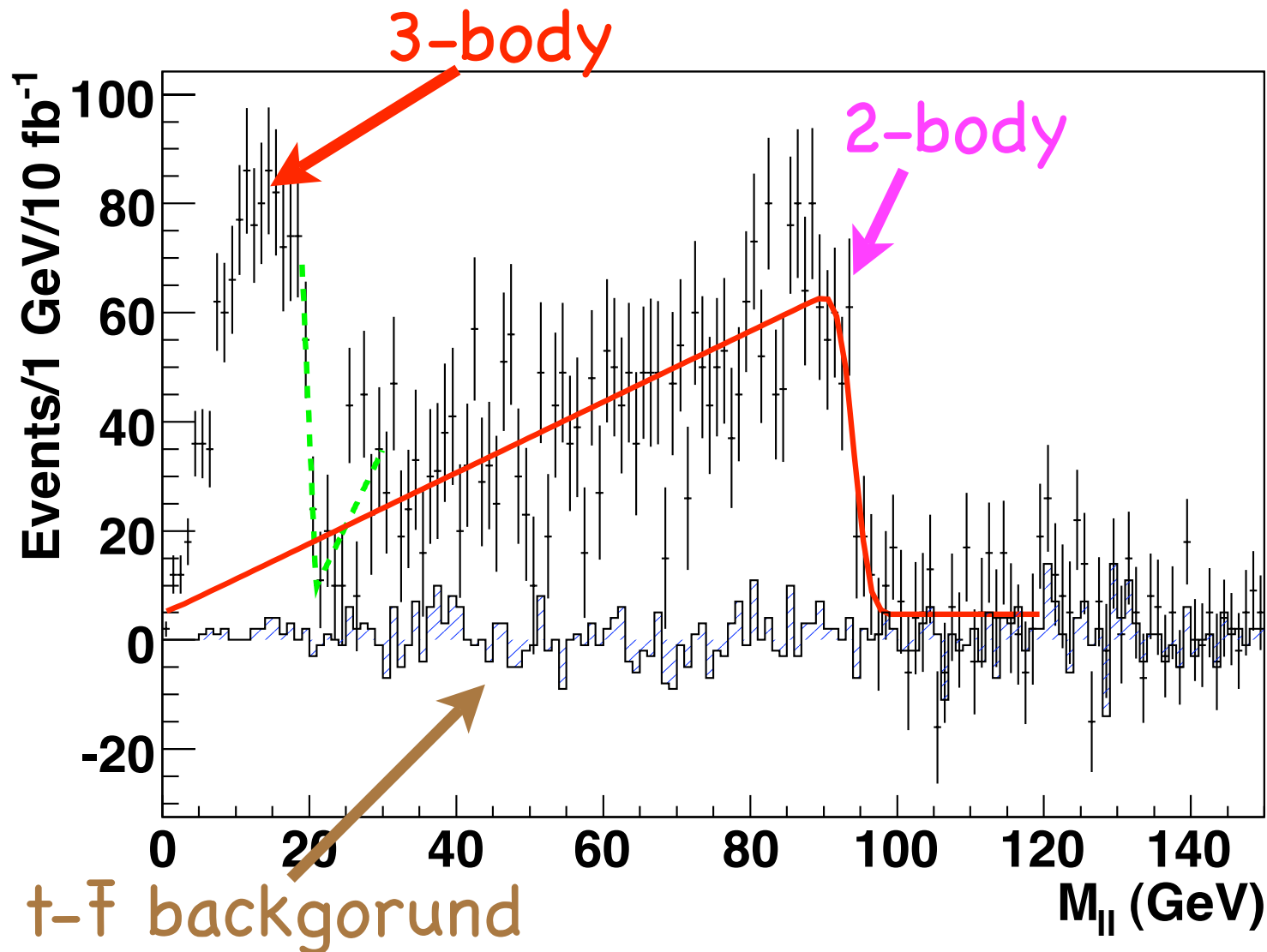
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We have checked that our method works even if a harder cut $P_T > 10$ GeV is taken.

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An example (Model 1): 10 fb^{-1}

(We used ISAJET(mass spectrum) + HERWIG + AcerDET)



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- from the # of events,

$$\frac{\Gamma_{3\text{body}}}{\Gamma_{2\text{body}}} = \frac{N_3/R_3}{N_2/R_2} = (1.50 \pm 0.15) \left(\frac{R_2/R_3}{4.5} \right)$$

$R_2, R_3 \dots$ correction factor due to P_T cut

In this case, $R_2/R_3 \simeq 4.5 \pm 1.1$

- from the endpoints of $M_{\ell\ell}, M_{\ell j}, M_{\ell\ell j}$

$$\Gamma_{3\text{body}} = 0.21^{+0.09}_{-0.07} \text{ eV}$$

- We then obtain the **gravitino mass!!**

$$m_{\tilde{G}} = (0.53^{+0.11}_{-0.10}) \left(\frac{R_2/R_3}{4.5} \right)^{1/2} \text{ eV}$$

(true value : $m_{\tilde{G}} = 0.77 \text{ eV}$)

Events/1 GeV/10 fb⁻¹

t-f

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- It can be **tested at LHC!**
(gravitino mass = SUSY breaking scale
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Backup Slides

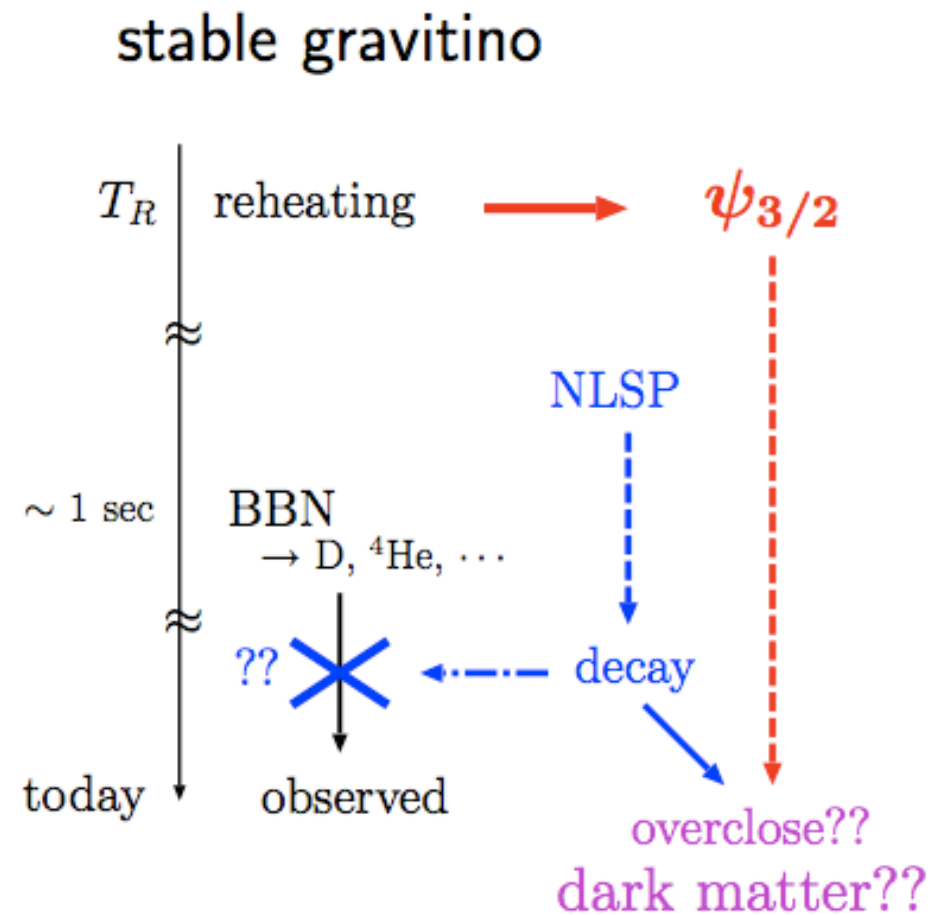
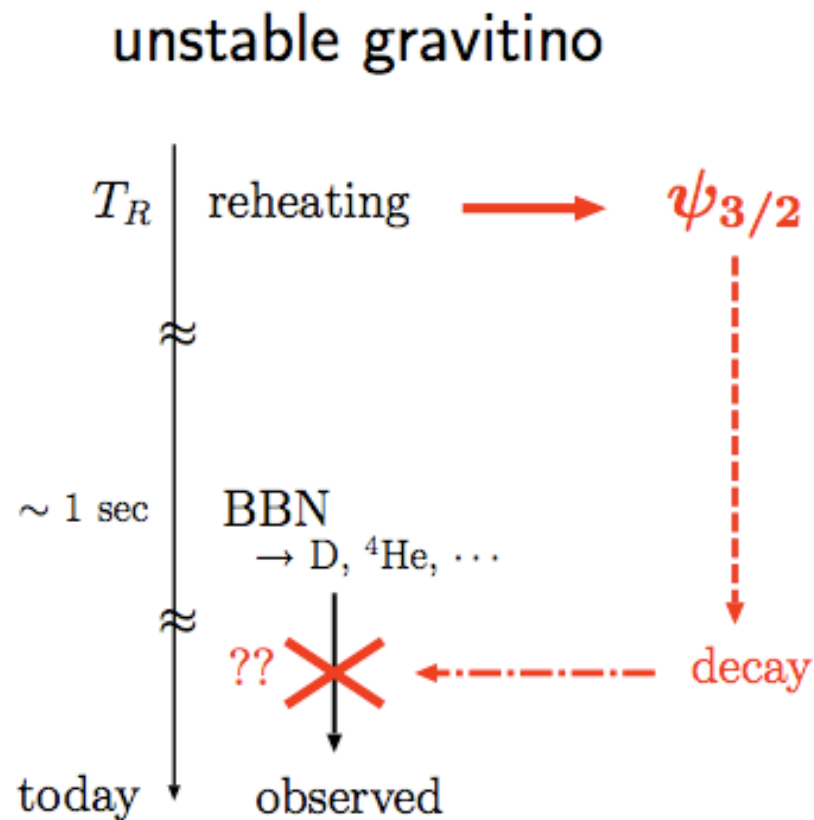
Gravitino Problems

thermal history

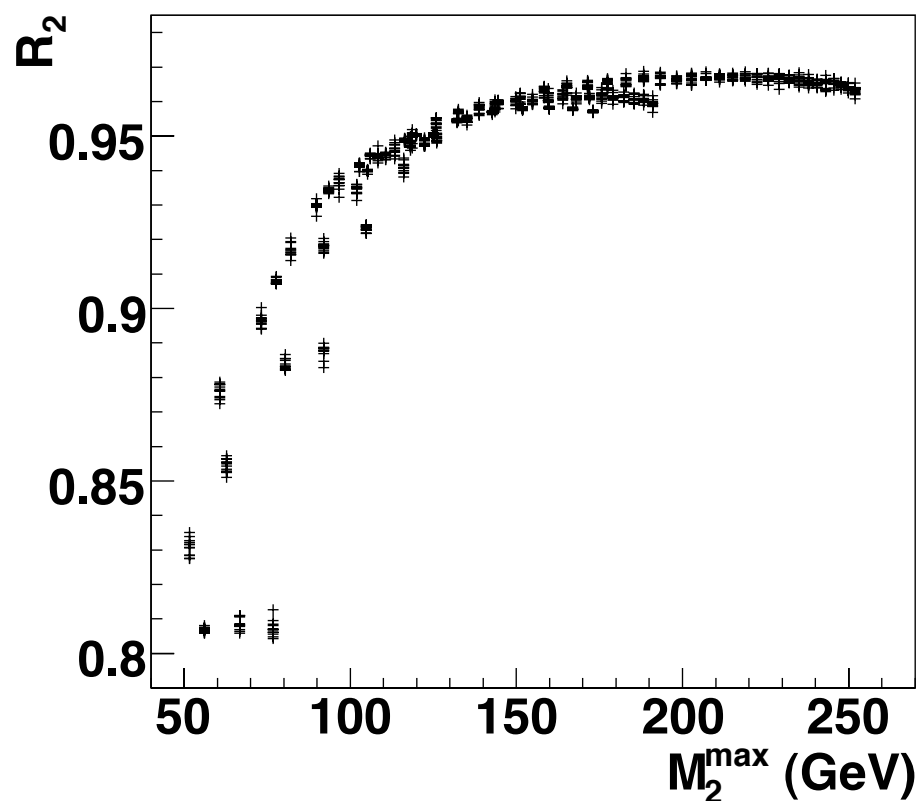
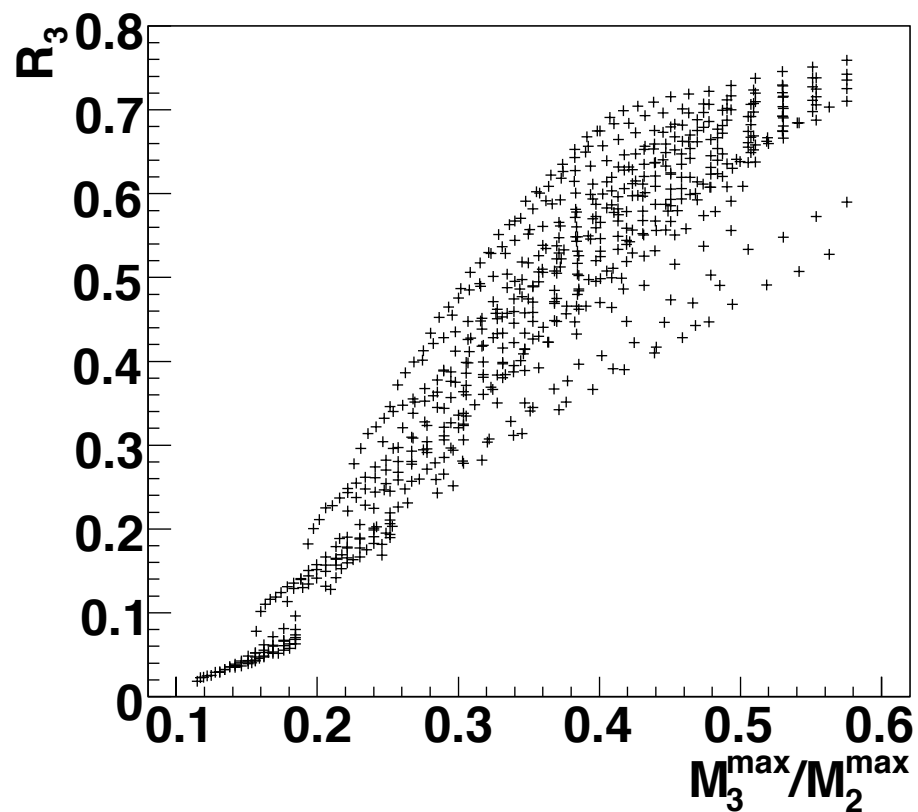
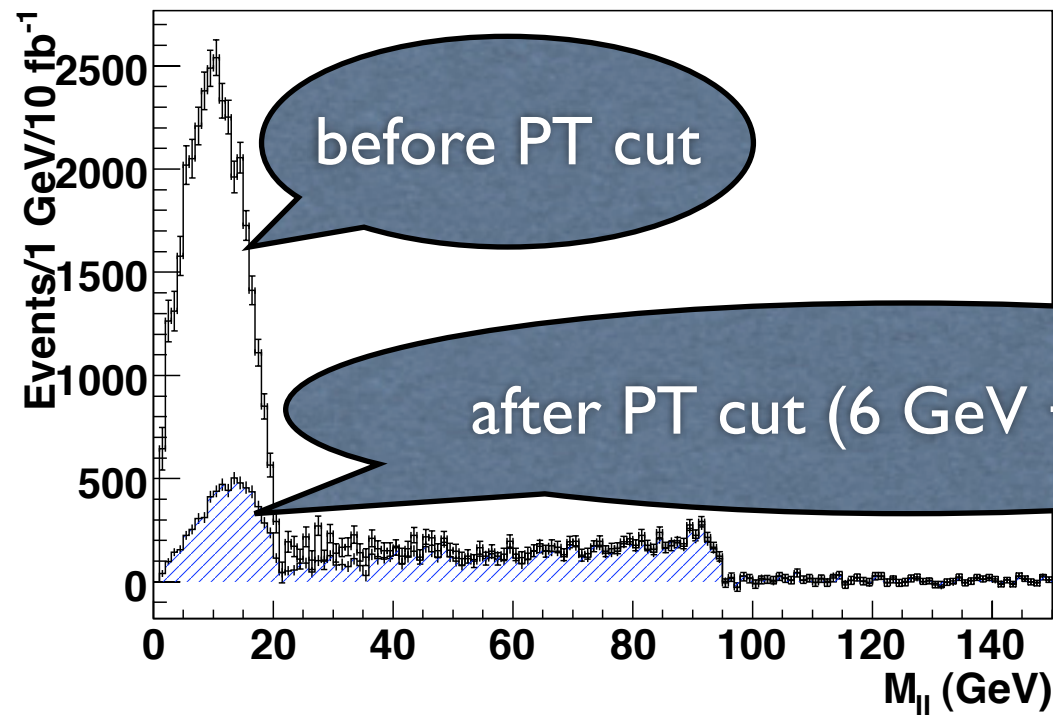
time	temperature	
??	~ 0	inflation
??	T_R	<u>reheating</u>
	\approx	<u>baryogenesis</u> $\rightarrow n_B/s \simeq 10^{-10}$
~ 1 sec	~ 1 MeV	Big Bang Nucleosynthesis $\rightarrow D, {}^4\text{He}, \dots$
	\approx	
14 Gyr	2.7 K	observed

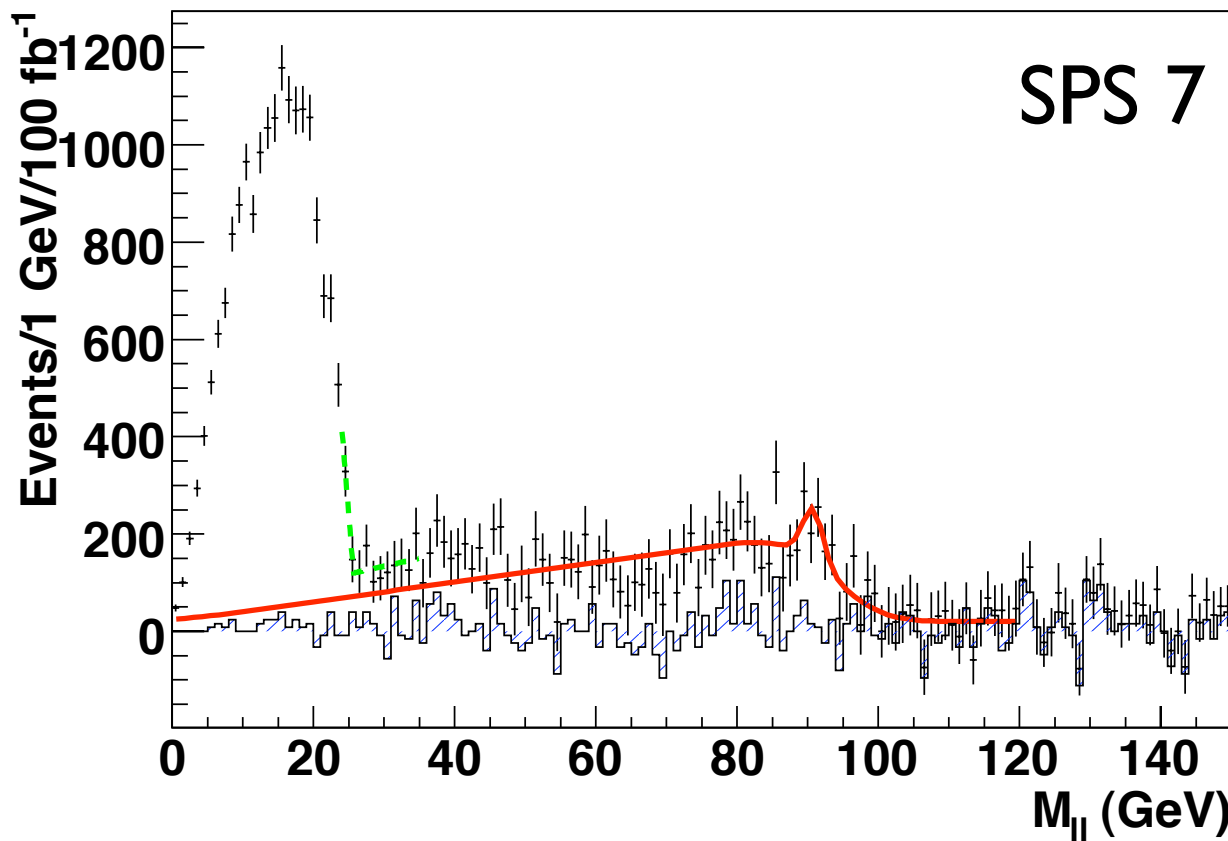
Gravitino Problems

thermal history with gravitino $\psi_{3/2}$



Reduction
factor R





$$\frac{\Gamma_{3\text{-body}}}{\Gamma_{2\text{-body}}} = \frac{N_3 R_2}{N_2 R_3} = (5.01 \pm 0.84) \left(\frac{R_2}{0.90}\right) \left(\frac{R_3}{0.35}\right)^{-1}$$

Assuming that we know $m_{\tilde{\ell}_R} = 129.1 \pm 0.5$ GeV, we estimate

$$\Gamma_{3\text{-body}} = 1.16_{-0.59}^{+1.28} \text{eV} \quad (\text{the true value is } \Gamma_{3\text{-body}} = 1.11 \text{ eV}),$$

which leads to

$$m_{3/2} = (0.41_{-0.12}^{+0.17}) \left(\frac{R_2}{0.90}\right)^{\frac{1}{2}} \left(\frac{R_3}{0.35}\right)^{-\frac{1}{2}} \text{eV}.$$

The expected value is $m_{3/2} = 0.77$ eV.

- from the # of events,

$$\frac{\Gamma_{3\text{body}}}{\Gamma_{2\text{body}}} = \frac{N_3/R_3}{N_2/R_2} = (1.50 \pm 0.15) \left(\frac{R_2/R_3}{4.5} \right)$$

$R_2, R_3 \dots$ correction factor due to P_T cut

In this case, $R_2/R_3 \simeq 4.5 \pm 1.1$

- from the endpoints of $M_{\ell\ell}, M_{\ell j}, M_{\ell\ell j}$

$$\Gamma_{2\text{body}} = 0.21^{+0.09}_{-0.07} \text{ eV}$$

- from the # of events,

$$\frac{\Gamma_{3\text{body}}}{\Gamma_{2\text{body}}} = \frac{N_3/R_3}{N_2/R_2} = (1.50 \pm 0.15) \left(\frac{R_2/R_3}{4.5} \right)$$

$R_2, R_3 \dots$ correction factor due to P_T cut

- from the endpoints of $M_{\ell\ell}, M_{\ell j}, M_{\ell\ell j}$

$$\Gamma_{2\text{body}} = 0.21^{+0.09}_{-0.07} \text{ eV}$$