



15th Sep 2008

@IPMU workshop on quantum BH

Holographic Nuclei

Koji Hashimoto (RIKEN)

[arXiv/0806.3122](https://arxiv.org/abs/0806.3122)(hep-th)

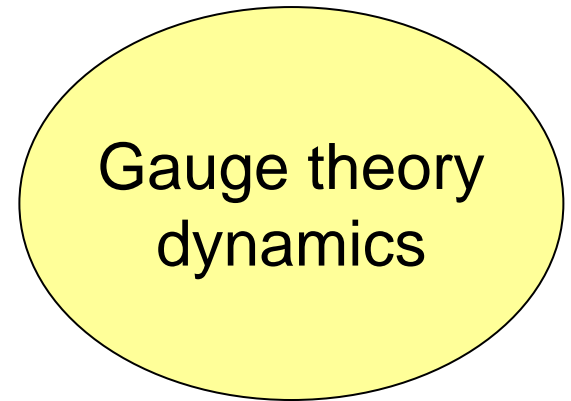
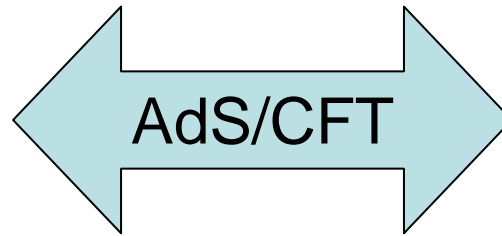
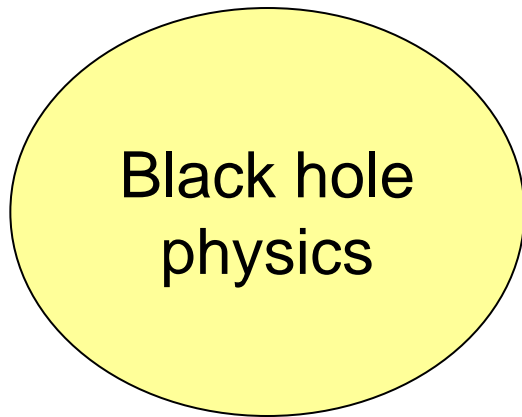
T.Sakai, S.Sugimoto, KH

Work in progress

KH

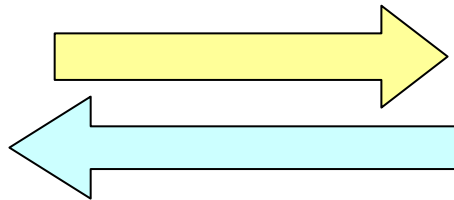
1. A Road to Nuclear Physics





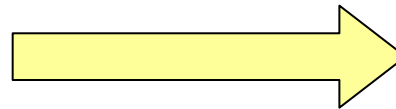
“Quantitative” success
(Large N?)

Holographic QCD



Experimental
data on hadrons

Holographic Nuclei



Nuclear data

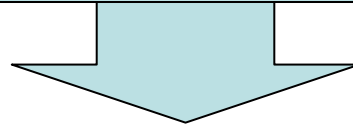
How to explore the Sakai-Sugimoto model

Achievements of the model

Analytic calculations of :

- * Meson spectrum
- * Chiral lagrangian of all mesons
- * Baryon spectrum
- * Finite temperature and phase transition

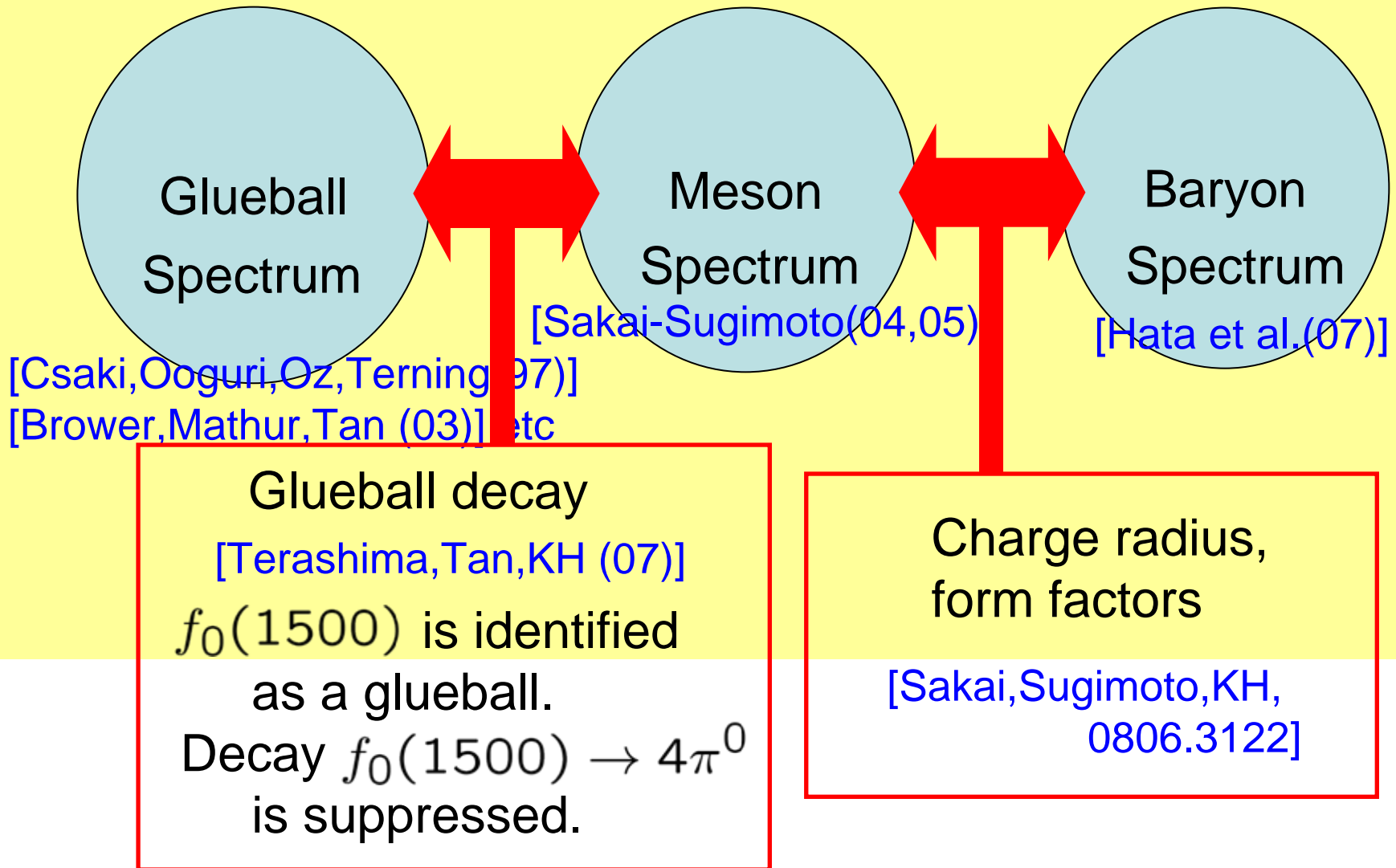
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How to
explore
further?

Compute more quantities in the model
Resolve problems of the model

Compute more in the model



Compute more in the model

Glueball

Meson

Baryon

Exotics

Finite temperature,
Finite density,
Chemical potential

More dynamical quantities :

Jet quenching parameters, viscosity of plasma, ...

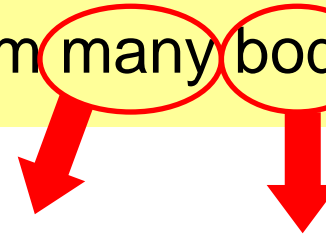
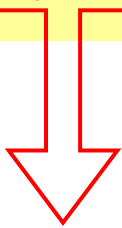
A road to nuclear physics?

Most of developments so far : hadron physics,
not really nuclear physics!

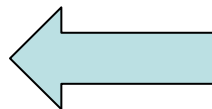
Nuclear physics

Light nuclei : Can be studied by nuclear force

Heavy nuclei : Complicated quantum many body problem



Possible
Gravity dual?



Nucleon number



large A

Nucleon



D-brane

Heavy nuclei have dual geometry description

Plan



1. A Road to Nuclear Physics

4 slides

2. Brief Review of Sakai-Sugimoto

4 slides

3. Holographic Baryons

[T.Sakai, S.Sugimoto and KH, 0806.3122]

7 slides

4. Holographic Nuclei

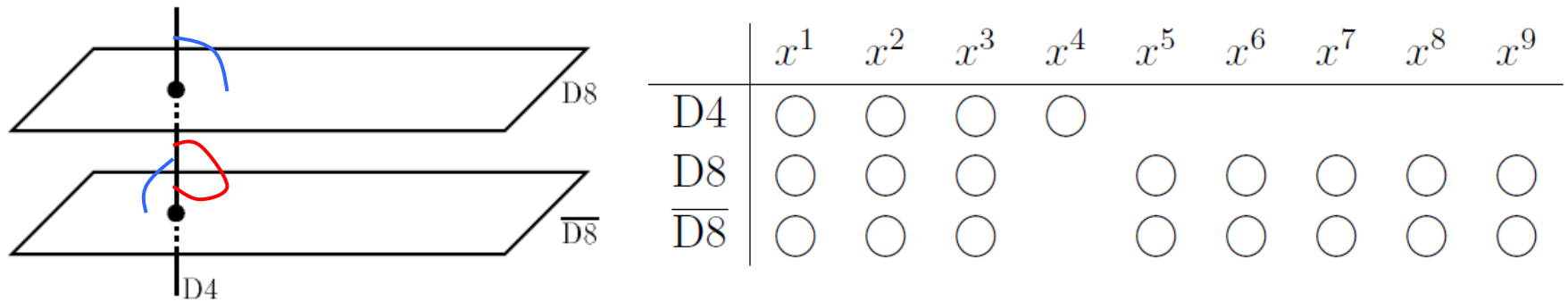
[KH, to appear]

5 slides

2. Sakai-Sugimoto : Review

Best model of holographic QCD : Sakai-Sugimoto model

■ Open string side (D-branes)



- N_c D4 wrap S^1 with radius $2\pi/M_{\text{KK}}$ [Witten(98)]
Gauginos satisfy anti-periodic boundary condition
→ **Massless gluons** at low energy, $SU(N_c)$ gauge
- N_f $\overline{\text{D8}}$ intersecting D4s → **N_f massless L-quarks**
- N_f D8 intersecting D4s → **N_f massless R-quarks**

Massless QCD is brane-engineered at low energy

$\text{D8} \times \overline{\text{D8}}$ gauge group = Chiral symmetry $U(N_f)_L \times U(N_f)_R$

- Closed string side (gravity)

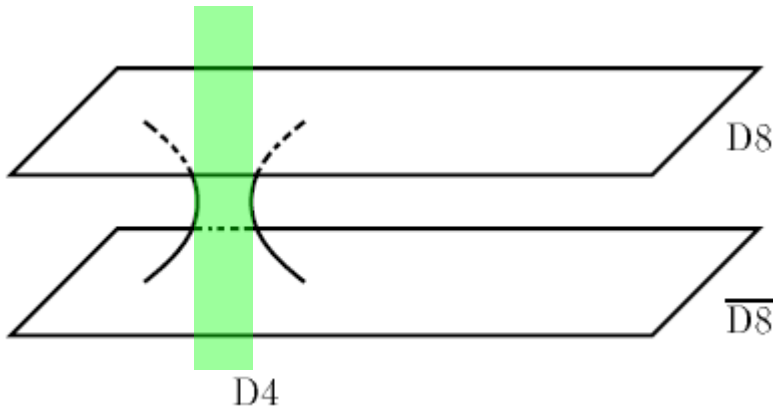
D8s : probe ($N_c \gg N_f$)

Geometry :

[Witten(98)]

$$ds^2 = \left(\frac{U}{R}\right)^{3/2} (\eta_{\mu\nu} dx^\mu dx^\nu + f(U) d\tau^2) + \left(\frac{R}{U}\right)^{3/2} \left(\frac{dU^2}{f(U)} + U^2 d\Omega_4^2 \right)$$

$$e^\phi = g_s \left(\frac{U}{R}\right)^{3/4}, \quad F_4 = \frac{2\pi N_c}{V_4} \epsilon_4, \quad f(U) = 1 - \frac{U_{\text{KK}}^3}{U^3}$$



Geometry truncated at $U = U_{\text{KK}}$

→ D8 and $\overline{\text{D8}}$ are connected

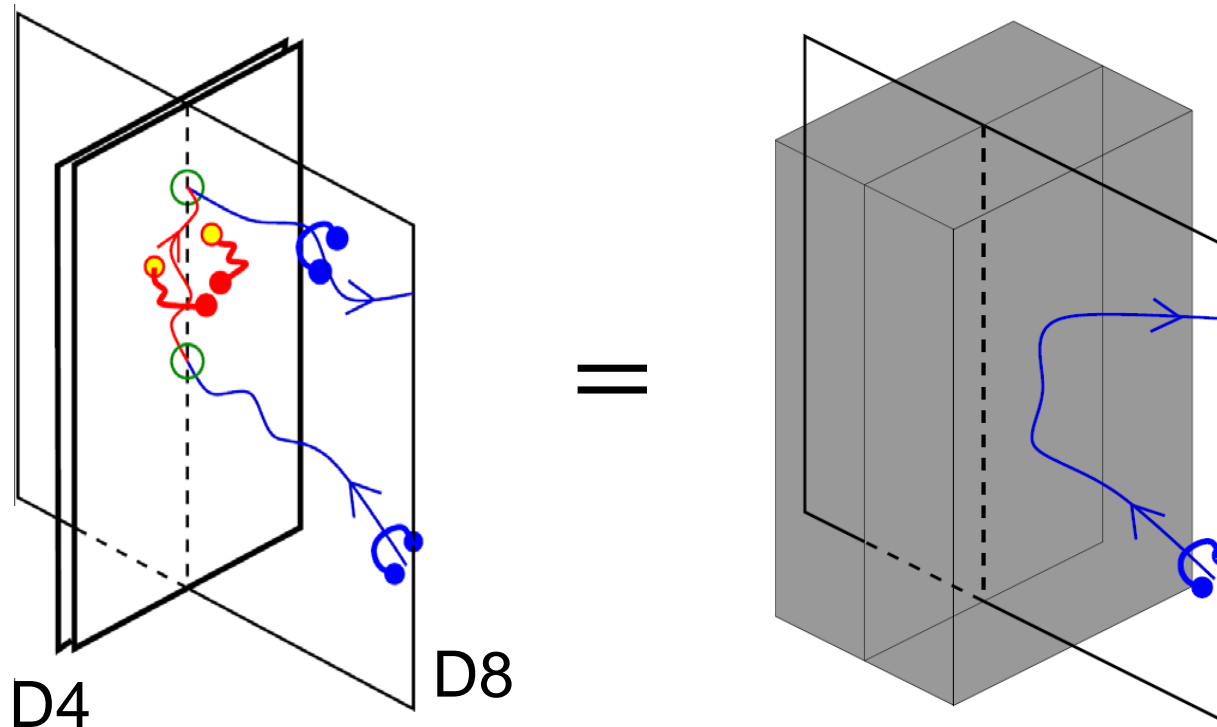
→ **Spontaneous Chiral sym. br.**

$$U(N_f)_L \times U(N_f)_R \rightarrow U(N_f)_V$$

Sugra parameters are

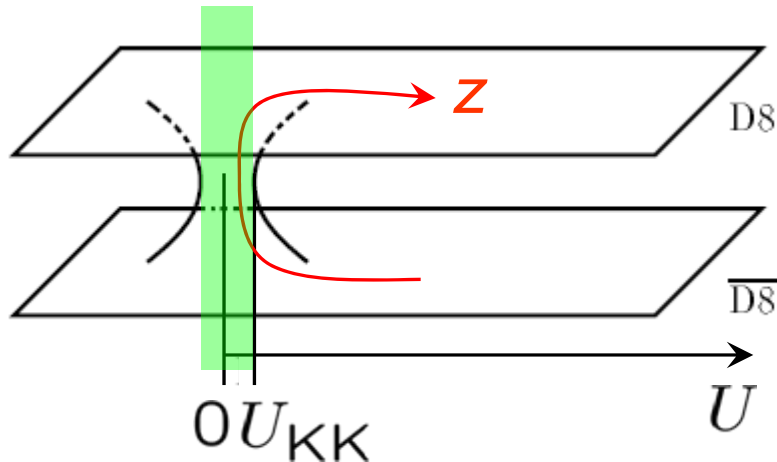
$$R^3 = \frac{g_{\text{YM}}^2 N l_s^2}{2M_{\text{KK}}}, \quad U_{\text{KK}} = \frac{2}{9} g_{\text{YM}}^2 N M_{\text{KK}} l_s^2, \quad g_s = \frac{g_{\text{YM}}^2}{2\pi M_{\text{KK}} l_s}$$

Gauge/Gravity correspondence implies....



- D8 \rightarrow bound states of quarks (Mesons, Baryons)
KK modes of gauge fields on the D8 = Mesons
D8-brane action = Chiral lagrangian
- gravitons = bound states of gluons (Glueballs)

Meson sector in the model



Redefinition of a coordinate:

$$U^3 \equiv U_{\text{KK}}^3 + U_{\text{KK}} z^2$$

D8-brane action on the curved background ($M_{\text{KK}} = 1$)

$$-\frac{\lambda N_c}{216\pi^3} \int d^4x dz \text{tr} \left[\frac{1}{2} (1+z^2)^{-1/3} F_{\mu\nu}^2 + (1+z^2) F_{\mu z}^2 \right] + \frac{N_c}{24\pi^2} \text{tr} \int \left[A F^2 - \frac{i}{2} A^3 F - \frac{1}{10} A^5 \right]$$

KK fluctuation analysis

$A_\mu \rightarrow$ Vector meson tower spectrum
 $A_z \rightarrow$ Massless pion

3. Holographic baryons



arXiv/0806.3122(hep-th)
T.Sakai, S.Sugimoto, KH

Baryons in the model

Baryon = D4 Wrapping S^4 [Witten(98), Gross, Ooguri(98)]
= Instanton in (x^1, x^2, x^3, z) of D8
[Sakai, Sugimoto(04)]

$$-\frac{\lambda N_c}{216\pi^3} \int d^4x dz \text{tr} \left[\frac{1}{2} (1+z^2)^{-1/3} F_{\mu\nu}^2 + (1+z^2) F_{\mu z}^2 \right] \\ + \frac{N_c}{24\pi^2} \text{tr} \int \left[A \boxed{F^2} - \frac{i}{2} A^3 F - \frac{1}{10} A^5 \right]$$

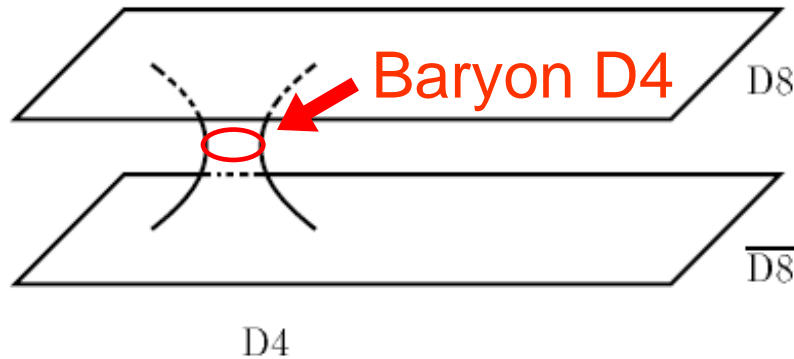
Instanton charge sources $U(1)_v$

“Electric energy” + “Potential in z ”

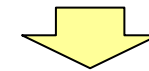
→ Stabilization of the instanton size $\sim \mathcal{O}(\lambda^{-1/2})$

Quantization of instantons → Baryon spectrum

[Hata, Sakai, Sugimoto, Yamato (HSSY) (07)]



Localized at $z = 0$



Geometry can be expanded around flat space

Solution : BPST instanton + electrostatic potential

$$A_M^{\text{cl}} = -if(\xi)g\partial_M g^{-1}, \quad \hat{A}_0^{\text{cl}} = \frac{N_c}{8\pi^2\kappa} \frac{1}{\xi^2} \left[1 - \frac{\rho^4}{(\rho^2 + \xi^2)^2} \right], \quad A_0 = \hat{A}_M = 0$$

$$f(\xi) = \frac{\xi^2}{\xi^2 + \rho^2}, \quad g(x) = \frac{(z - Z) - i(\vec{x} - \vec{X}) \cdot \vec{\tau}}{\xi}, \quad \xi = \sqrt{(z - Z)^2 + |\vec{x} - \vec{X}|^2}$$

Inserting this back to the action leads to a potential

$$U(\rho, Z) = 8\pi^2\kappa \left(1 + \frac{\rho^2}{6} + \frac{N_c^2}{5(8\pi^2\kappa)^2} \frac{1}{\rho^2} + \frac{Z^2}{3} \right) \quad \kappa = \frac{\lambda N_c}{216\pi^3}$$



Size is stabilized, $\rho_{\text{cl}}^2 = \frac{N_c}{8\pi^2\kappa} \sqrt{\frac{6}{5}}, \quad Z_{\text{cl}} = 0$

Moduli space approximation

$$\left\{ \begin{array}{l} \text{Pseudo-moduli : } (\vec{X}, Z), y_I \quad (I = 1, 2, 3, 4) \\ \text{Moduli potential : } L = \frac{M_0}{2} (\dot{\vec{X}}^2 + \dot{Z}^2) + M_0 \dot{y}_I^2 - U(\rho, Z) \end{array} \right.$$

Quantum mechanics for (Z, y^I) $M_0 = 8\pi^2 \kappa$

$SO(4) \simeq (SU(2)_I \times SU(2)_J) / \mathbb{Z}_2$: Isospin and spin

$$I = J = l/2 \quad l = 1, 3, 5, \dots$$

$Z, \rho \rightarrow$ Harmonic-like potential $\rightarrow n_\rho, n_z$

Baryon state : $|\vec{p}, B, s\rangle \quad B \equiv (l, I_3, n_\rho, n_z)$

Proton $B = (1, 1/2, 0, 0) \quad |p \uparrow\rangle \propto R(\rho) \psi_Z(Z) (a_1 + ia_2)$

Neutron $B = (1, -1/2, 0, 0) \quad |n \uparrow\rangle \propto R(\rho) \psi_Z(Z) (a_4 + ia_3)$

$$R(\rho) = \rho^{-1+2\sqrt{1+N_c^2/5}} e^{-\frac{M_0}{\sqrt{6}} \rho^2}, \quad \psi_Z(Z) = e^{-\frac{M_0}{\sqrt{6}} Z^2} \quad a_I \equiv y_I / \rho$$

Chiral currents

$$J_L(r), J_R(r)$$



Static properties

Charge radius
Magnetic moments
Axial radius, coupling

Chiral symmetry = D8 gauge symmetry at $z = \pm\infty$

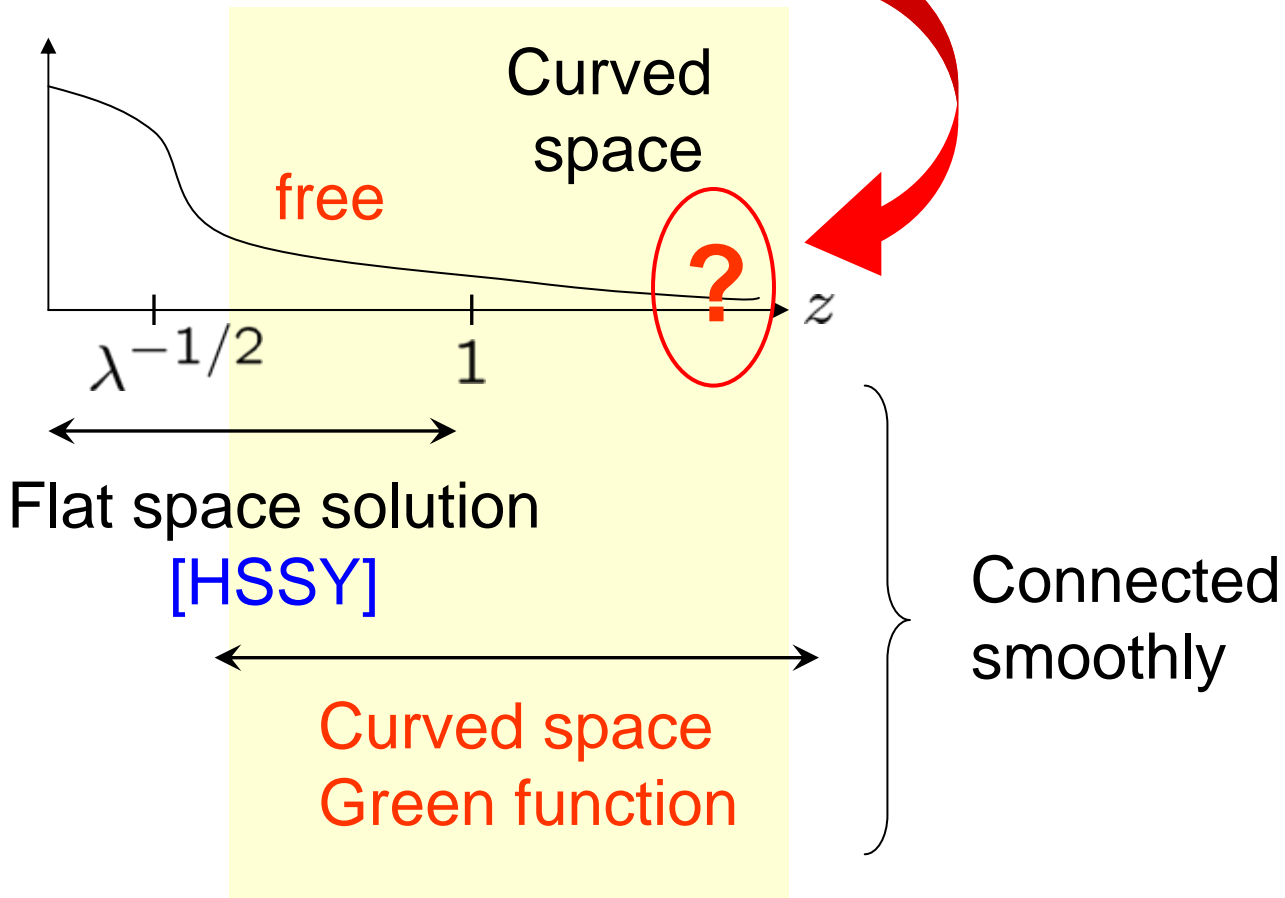
$$S|_{\mathcal{O}(\mathcal{A}_L, \mathcal{A}_R)} = -2 \int d^4x \text{tr} \left(\mathcal{A}_{L\mu} \mathcal{J}_L^\mu + \mathcal{A}_{R\mu} \mathcal{J}_R^\mu \right)$$

$$\left(\begin{array}{l} \mathcal{A}_\alpha(x^\mu, z) = \mathcal{A}_\alpha^{\text{cl}}(x^\mu, z) + \delta \mathcal{A}_\alpha(x^\mu, z) \\ \delta \mathcal{A}_\mu(x^\mu, z \rightarrow +\infty) = \mathcal{A}_{L\mu}(x^\mu), \quad \delta \mathcal{A}_\mu(x^\mu, z \rightarrow -\infty) = \mathcal{A}_{R\mu}(x^\mu) \end{array} \right)$$

$$\Rightarrow J_{L,R}(r) = \mp \frac{\lambda N_c}{216\pi^3} \left[(1 + z^2) F_{\mu z}^{(\text{sol})} \right]_{z=\pm\infty}$$

Evaluation of chiral currents

We need the solution at $z = \pm\infty$



Explicit expression of the chiral currents

Static properties of baryons

1) Baryon number current $J_B(r)$

$$1 = N_B = \int_0^\infty dr \rho_B(r), \quad \rho_B(r) \equiv 4\pi r^2 \langle J_B^0(r) \rangle$$

→ Isoscalar mean square radius $\langle r^2 \rangle_{I=0} = \int_0^\infty dr r^2 \rho_B(r)$

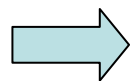
Our result : $\langle r^2 \rangle_{I=0}^{1/2} \simeq 0.742 \text{ fm}$

(Input : $M_{\text{KK}} = 949 \text{ MeV}$, $\kappa = 0.00745$)

2) Vector current $J_V^0(r)$

Isvector charge density turns out to be $\rho_{I=1}(r) = \rho_B(r)$

$$Q_{\text{em}} = I_3 + \frac{N_B}{2}$$



Charge radius $\langle r^2 \rangle_{\text{E,p}} = \langle r^2 \rangle_{I=0}$, $\langle r^2 \rangle_{\text{E,n}} = 0$

Our results

	Holographic QCD	Skyrmion	Experiment
$\langle r^2 \rangle_{I=0}^{1/2}$	0.742 fm	0.59 fm	0.806 fm
$\langle r^2 \rangle_{M, I=0}^{1/2}$	0.742 fm	0.92 fm	0.814 fm
$\langle r^2 \rangle_{E,p}$	$(0.742 \text{ fm})^2$	∞	$(0.875 \text{ fm})^2$
$\langle r^2 \rangle_{E,n}$	0	$-\infty$	-0.116 fm^2
$\langle r^2 \rangle_{M,p}$	$(0.742 \text{ fm})^2$	∞	$(0.855 \text{ fm})^2$
$\langle r^2 \rangle_{M,n}$	$(0.742 \text{ fm})^2$	∞	$(0.873 \text{ fm})^2$
$\langle r^2 \rangle_A^{1/2}$	0.537 fm	—	0.674 fm
μ_p	2.18	1.87	2.79
μ_n	-1.34	-1.31	-1.91
$\left \frac{\mu_p}{\mu_n} \right $	1.63	1.43	1.46
g_A	0.73	0.61	1.27
$g_{\pi NN}$	7.46	8.9	13.2
$g_{\rho NN}$	5.80	—	4.2 ~ 6.5

input : f_π, m_ρ

4. Holographic Nuclei

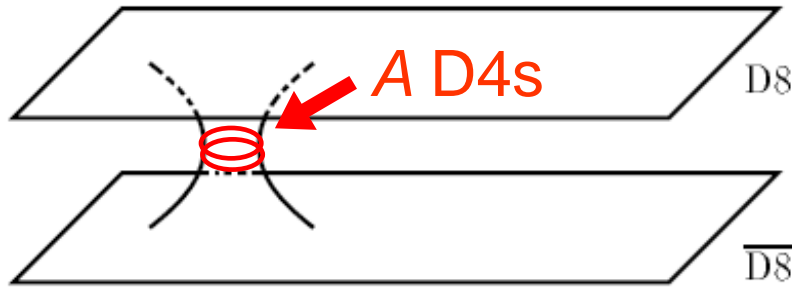


Work in progress

KH

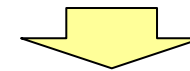
Heavy nuclei have dual geometry?

A : Mass number of nucleus



Nucleon = D4 wrapping S^4

Nucleus = A D4s



Large A

Heavy nucleus = D4 geometry

Uranium, Plutonium, ...

	0	1	2	3	y	z	S^4
N_f D8s	0	0	0	0		0	0 0 0 0
A D4s	0						0 0 0 0

Fields on the A D4s

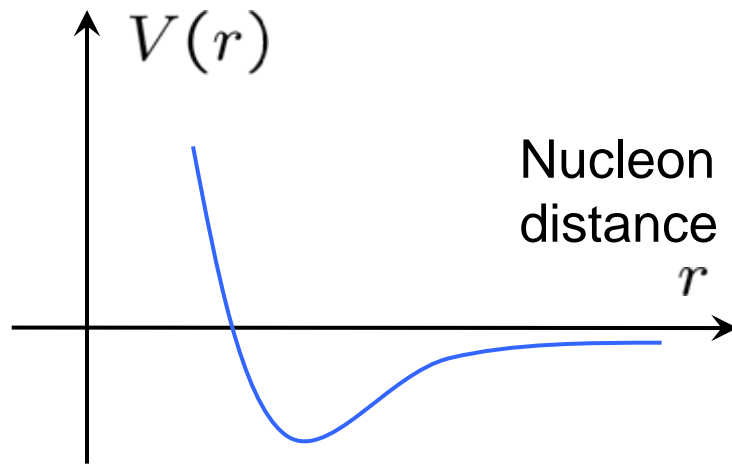
- $\left\{ \begin{array}{l} \Phi_i \ (i = 1, 2, 3, z, y) : \text{Adj. scalars (D4-D4)} \\ \rho_a \ (a = 1, \dots, N_f) : \text{Fund. scalars (D4-D8)} \end{array} \right.$

\rightarrow $U(A)$ "ADHM" gauge theory on S^4

: Complicated. Non-susy, nontrivial potential, ..

Very quick course for Nuclear theory

Nuclear force



Short range : Strong repulsion
[Sakai,Sugimoto,KH] in progress

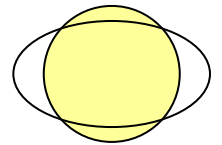
Long range : Pion exchange

→ Nucleons are stabilized at a certain distance

→ Heavy nuclei has radius

$$R_0 \propto A^{1/3}$$

Liquid drop model of heavy nuclei [Gamow]



Nuclei are approximated by incompressible fluid ball

Nucleon binding energy is independent of nucleus size

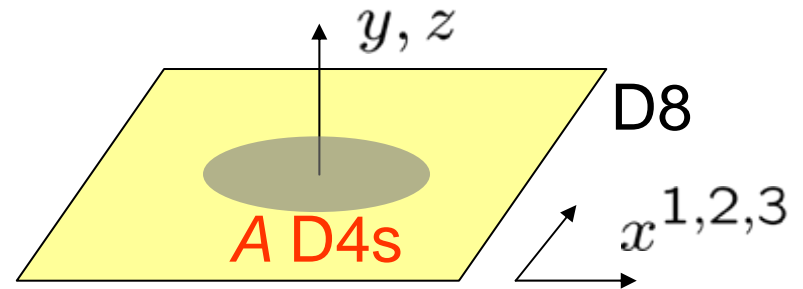
Excitation spectrum of nucleus : $E_n = \omega n$ ($n = 1, 2, \dots$)

Dual geometry?

Backreacted geometry : Gravitating instantons

D4s = Electrically charged A instantons on N_f D8

Let's assume isotropic instanton density on D8 to mimic the nucleus



D4s are localized at the tip of Witten's geometry

$$ds^2 = \left(\frac{U_{KK}}{R}\right)^{3/2} \eta_{\mu\nu} dx^\mu dx^\nu + \frac{4}{9} \left(\frac{R}{U_{KK}}\right)^{3/2} (dy^2 + dz^2) + R^{3/2} U_{KK}^{1/2} d\Omega_4^2$$

→ Dual geometry can be roughly approximated by A D4s distributed on a 3-ball with radius $R_N = \frac{A^{1/3}}{\sqrt{\lambda} M_{KK}}$ in this almost flat asymptotics.

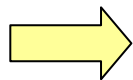
Near horizon geometry, Dictionary

Near horizon?We don't know explicit solution

But we can argue validity of near-horizon requirement.
Validity condition for “would-be AdS radius”: $R_0 > R_N$

Indeed,

$$R_0 = (Ag'_s)^{1/3} l_s / \sqrt{g_{ii}} = \frac{A^{1/3}}{N_c^{1/3} M_{KK}} > R_N = \frac{A^{1/3}}{\sqrt{\lambda} M_{KK}}$$



Near horizon argument is suggested to be fine

Dictionary (partial)

Metric fluctuation $g_{ij} \Leftrightarrow \text{tr}[[\Phi_i, \Phi_k][\Phi_j, \Phi_k]]$

Gauge fluctuation $A_i \Leftrightarrow \rho^\dagger \gamma_i \rho$

Collective motion of nucleons is described by sugra fields

Spectrum and comparison with nuclear data

The “closest approximation” of the sugra solution is

D3s distributed on a 3-ball of radius R_N

$$H(r) = 1 + \int_{B_3} d^3\vec{a} \frac{cR_N^4}{|\vec{r} - \vec{a}|^4}$$

Fluctuation spectrum

$$\frac{1}{\sqrt{G}} \partial_\mu \sqrt{G} G^{\mu\nu} \partial_\nu \phi = 0$$

→ $E \sim \frac{R_N}{R_0^2} n \quad (n = 1, 2, \dots)$

[Freedman, Gubser, Pilch, Warner(99)] [Brandhuber, Sfetsos(99)]

Mass gap due to automatic “truncation” of geometry

$E \propto n$ coincident with the empirical formula for nuclei

$$\frac{R_N}{R_0^2} \sim \frac{N^{2/3}}{A^{1/3} \sqrt{\lambda}} M_{KK} \sim \mathcal{O}(10) [\text{MeV}]$$

Rough agreement

5. Conclusion & Discussions

Conclusion

Holographic baryons :

Static properties of baryons, computed in SS model, nicely match exp. data

Holographic nuclei :

A new AdS/CFT scheme which can be compared with real nuclear data is presented.

Discussions

Using AdS/CFT twice?

“CFT” (QCD) \rightarrow “AdS” + other D-branes
“CFT” \rightarrow “AdS”

Cf) Multi-center BH and Coulomb branch of SYM

Holographic nuclei can be black holes?!

Event horizon may happen to form ?

\rightarrow Deconfinement phase (“BH embedding”)

\rightarrow Continuous spectrum.....

But, it leads to fluid dynamics...

\rightarrow Possible relation to liquid drop model of nuclei ?!

BH formation and hadron collision? [Nastase] & many others...