

Jet angular correlation in vector-boson fusion processes at hadron colliders

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based on the work with **Qiang Li** (Karlsruhe)
and **Kentarou Mawatari** (KIAS → Heidelberg)

KH, Qiang Li, Kentarou Mawatari,
arXiv:0810.xxxx [hep-ph]

Two recent works on massive graviton productions at hadron colliders:

- *Graviton production with 2 jets at the LHC in large extra dimensions*
by KH, P.Konar, Q.Li, K.Mawatari, D.Zeppenfeld,
JHEP 0804:019(2008)[arXiv:0801.1794[hep-ph]]
- *HELAS and MadGraph/MadEvent with spin-2 particles*
by KH, J.Kanzaki, Q.Li, K.Mawatari,
EPJC 56:435-447(2008)[arXiv:0805.2554[hep-ph]]

There were two motivations for the above works:

- **jet physics** as a probe of new physics
- **tools** for simulating new physics models

We found (confirmed the naive expectation) that

$$\sigma(2\text{jets})/\sigma(1\text{jet}) \text{ grows with the missing mass}$$

when the missing p_T is common. This type of studies will be helpful in estimating the **mass** of a singly produced missing particle.

However, we failed to identify interesting angular correlations among two high p_T jets and the missing p_T direction, from which we hoped to obtain information about the **spin** of the produced particle.

This is in contrast to the observation made for $H + 2$ jets production, where the jet angular correlation can reveal its **CP parity**:

- *H + 2 jets via gluon fusion*
by V.Del Duca, W.Kilgore, C.Oleari, C.Schmidt, D.Zeppenfeld,
PRL 87:122001(2001)[arXiv:hep-ph/0105129]
- *Determining the structure of Higgs couplings at the LHC*
by T.Plehn, D.L.Rainwater, D.Zeppenfeld,
PRL 88:051801(2002)[arXiv:hep-ph/0105325]
- *Gluon fusion contributions to H + 2 jet production*
by V.Del Duca, W.Kilgore, C.Oleari, C.Schmidt, D.Zeppenfeld,
NPB 616:367(2001)[arXiv:hep-ph/0108030]

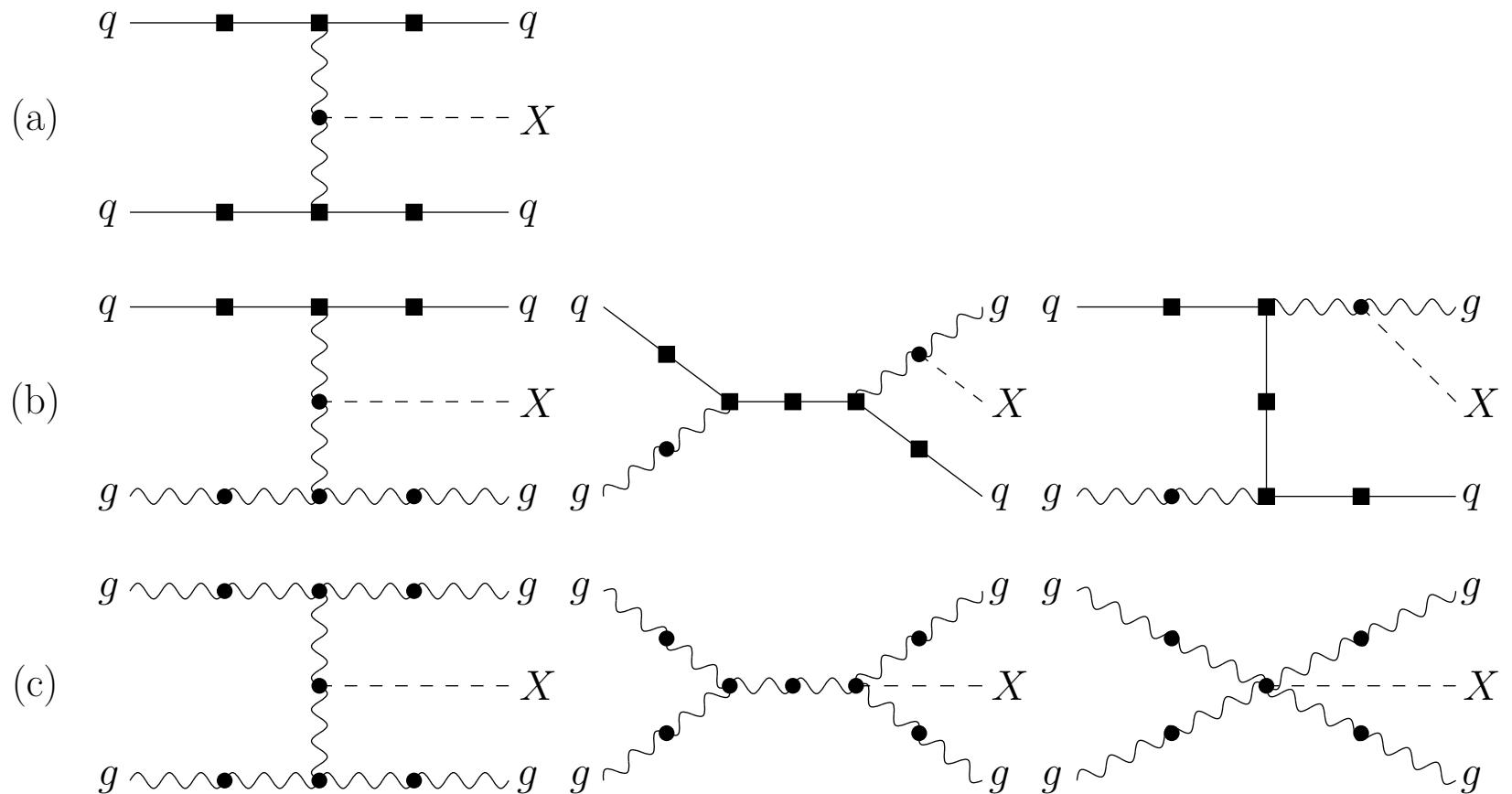
We try to understand the **jet angular correlations** associated with **vector-boson (gluon, $W/Z/\gamma$) fusion** production of a massive particle (X) with spin 0 and 2 ($J_X = 0$ or 2):

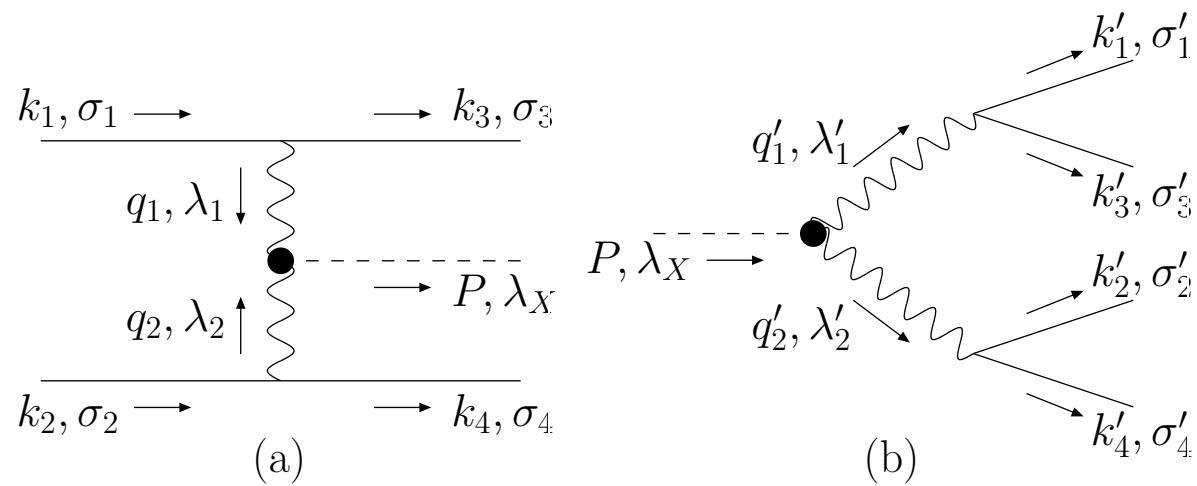
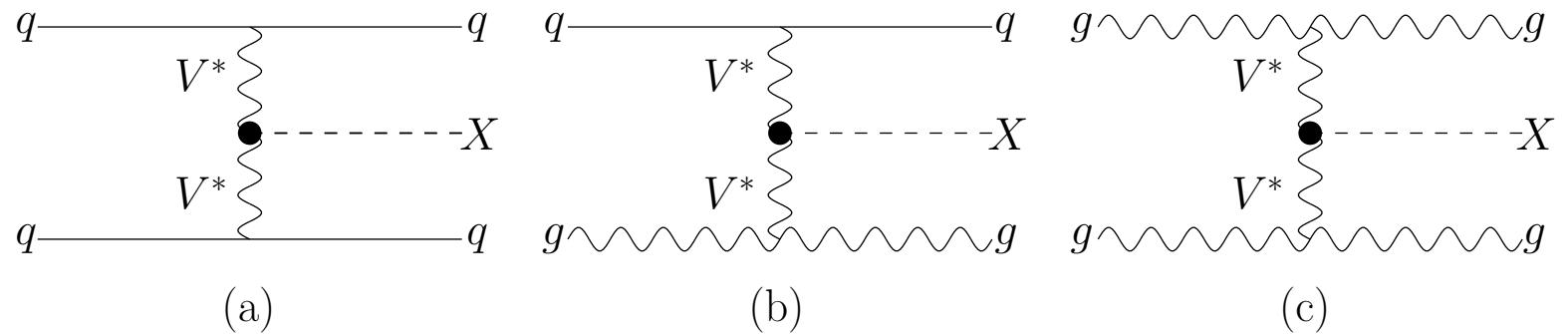
$$\begin{aligned} qq &\rightarrow qqX \quad \text{via} \quad g^*g^*, W^*W^*, Z^*Z^*, Z^*\gamma^*, \gamma^*\gamma^* \rightarrow X \\ qg &\rightarrow qgX \quad \text{via} \quad g^*g^* \rightarrow X \\ gg &\rightarrow ggX \quad \text{via} \quad g^*g^* \rightarrow X \end{aligned}$$

The results are then compared to the correlations in X decays into 4 jets (or 4 leptons):

$$\begin{aligned} X &\rightarrow W^*W^*, Z^*Z^*, Z^*\gamma^*, \gamma^*\gamma^* \rightarrow f\bar{f}f\bar{f} \\ X &\rightarrow g^*g^* \rightarrow q\bar{q}q\bar{q} \\ X &\rightarrow g^*g^* \rightarrow q\bar{q}gg \\ X &\rightarrow g^*g^* \rightarrow gggg \end{aligned}$$

where $f\bar{f}$ can either be a quark-pair or a lepton-pair.





The helicity amplitudes for the VBF processes

$$\begin{aligned} \mathcal{M}_{\sigma_1\sigma_3,\sigma_2\sigma_4}^{\lambda_X} &= \sum_{V_1,V_2} J_{V_1 a_1 a_3}^{\mu'_1}(k_1, k_3; \sigma_1, \sigma_3) J_{V_2 a_2 a_4}^{\mu'_2}(k_2, k_4; \sigma_2, \sigma_4) \\ &\quad \times D_{\mu'_1 \mu_1}^{V_1}(q_1) D_{\mu'_2 \mu_2}^{V_2}(q_2) \Gamma_{X V_1 V_2}^{\mu_1 \mu_2}(q_1, q_2; \lambda_X)^* \end{aligned}$$

can be expressed by using

$$-g_{\mu'\mu} + \frac{q_i \mu' q_i \mu}{q_i^2} = \sum_{\lambda_i=\pm,0} (-1)^{\lambda_i+1} \epsilon_{\mu'}(q_i, \lambda_i)^* \epsilon_\mu(q_i, \lambda_i)$$

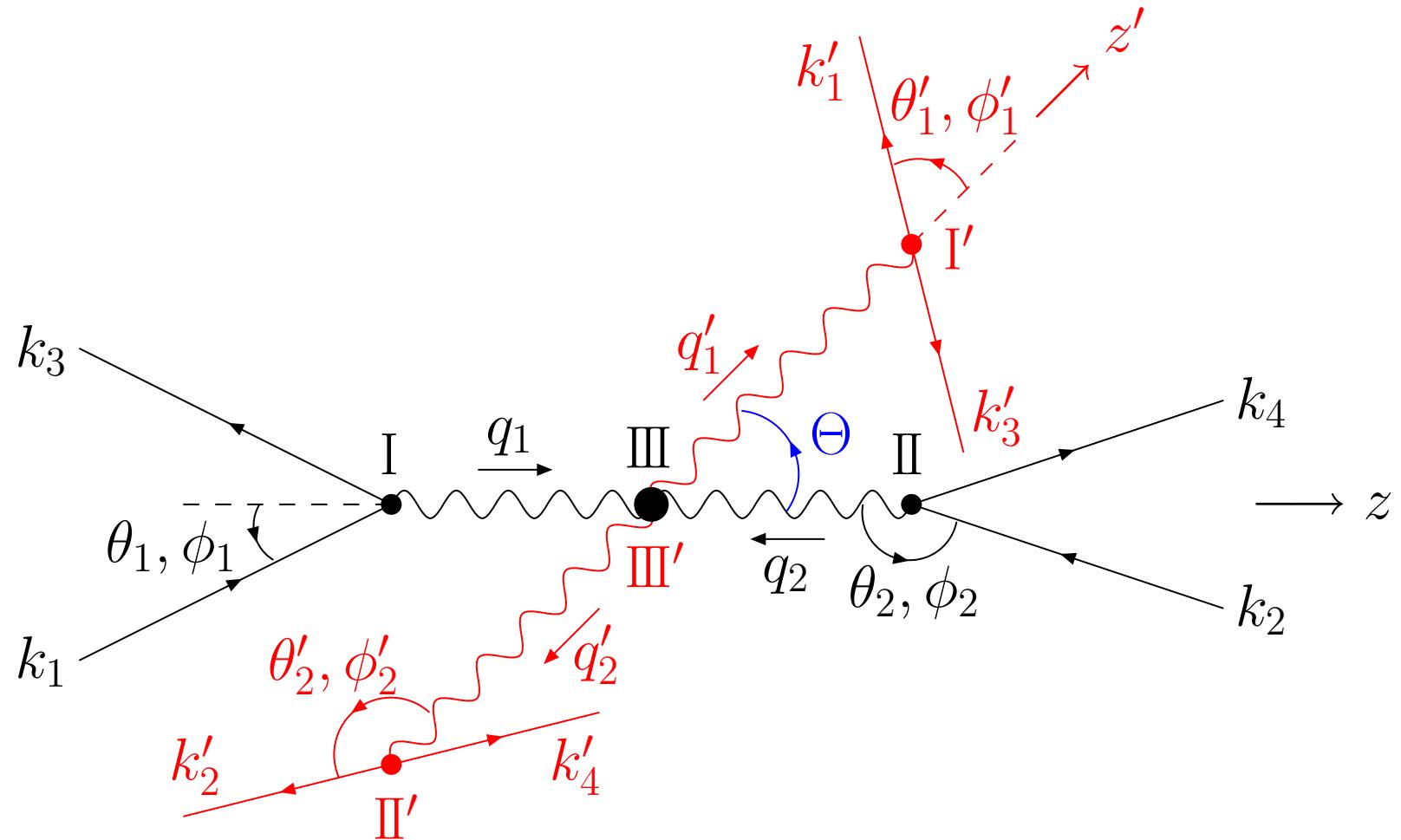
$$q_i \mu J_{V_i a_i a_{i+2}}^\mu(k_i, k_{i+2}; \sigma_i, \sigma_{i+2}) = 0$$

as follows:

$$\begin{aligned} \mathcal{M}_{\sigma_1\sigma_3,\sigma_2\sigma_4}^{\lambda_X} &= \sum_{V_1,V_2} \frac{1}{(q_1^2 - m_{V_1}^2)(q_2^2 - m_{V_2}^2)} \\ &\quad \times \sum_{\lambda_{1,2}=\pm,0} \left(J_{a_1 a_3}^{V_1} \right)_{\sigma_1 \sigma_3}^{\lambda_1} \left(J_{a_2 a_4}^{V_2} \right)_{\sigma_2 \sigma_4}^{\lambda_2} \left(\mathcal{M}_{V_1 V_2}^X \right)_{\lambda_1 \lambda_2}^{\lambda_X} \end{aligned}$$

where

$$\begin{aligned} \left(J_{a_i a_{i+2}}^{V_i} \right)_{\sigma_i \sigma_{i+2}}^{\lambda_i} &= (-1)^{\lambda_i+1} J_{V_i a_i a_{i+2}}^\mu(k_i, k_{i+2}; \sigma_i, \sigma_{i+2}) \epsilon_\mu(q_i, \lambda_i)^* \\ \left(\mathcal{M}_{V_1 V_2}^X \right)_{\lambda_1 \lambda_2}^{\lambda_X} &= \epsilon_{\mu_1}(q_1, \lambda_1) \epsilon_{\mu_2}(q_2, \lambda_2) \Gamma_{X V_1 V_2}^{\mu_1 \mu_2}(q_1, q_2; \lambda_X)^* \end{aligned}$$



I) the q_1 Breit frame ($Q_1 = \sqrt{-q_1^2}$, $0 < \theta_1 < \pi/2$ and $0 < \phi_1 < 2\pi$):

$$q_1^\mu = k_1^\mu - k_3^\mu = (0, 0, 0, Q_1)$$

$$k_1^\mu = \frac{Q_1}{2\cos\theta_1}(1, \sin\theta_1 \cos\phi_1, \sin\theta_1 \sin\phi_1, \cos\theta_1)$$

$$k_3^\mu = \frac{Q_1}{2\cos\theta_1}(1, \sin\theta_1 \cos\phi_1, \sin\theta_1 \sin\phi_1, -\cos\theta_1)$$

II) the q_2 Breit frame ($Q_2 = \sqrt{-q_2^2}$, $\pi/2 < \theta_2 < \pi$ and $0 < \phi_2 < 2\pi$).

$$q_2^\mu = k_2^\mu - k_4^\mu = (0, 0, 0, -Q_2)$$

$$k_2^\mu = -\frac{Q_2}{2\cos\theta_2}(1, \sin\theta_2 \cos\phi_2, \sin\theta_2 \sin\phi_2, \cos\theta_2)$$

$$k_4^\mu = -\frac{Q_2}{2\cos\theta_2}(1, \sin\theta_2 \cos\phi_2, \sin\theta_2 \sin\phi_2, -\cos\theta_2)$$

III) the VBF frame

$$q_1^\mu + q_2^\mu = P^\mu = q'_1{}^\mu + q'_2{}^\mu = (M, 0, 0, 0)$$

$$q'_1{}^\mu = \frac{M}{2}\left(1 - \frac{Q_1^2 - Q_2^2}{M^2}, 0, 0, \beta\right)$$

$$q'_2{}^\mu = \frac{M}{2}\left(1 - \frac{Q_2^2 - Q_1^2}{M^2}, 0, 0, -\beta\right)$$

$$q'_1{}^\mu = \frac{M}{2}\left(1 + \frac{q'_1{}^2 - q'_2{}^2}{M^2}, \beta' \sin\Theta, 0, \beta' \cos\Theta\right)$$

$$q'_2{}^\mu = \frac{M}{2}\left(1 + \frac{q'_2{}^2 - q'_1{}^2}{M^2}, -\beta' \sin\Theta, 0, -\beta' \cos\Theta\right)$$

where $\beta = \bar{\beta}\left(-\frac{Q_1^2}{M^2}, -\frac{Q_2^2}{M^2}\right)$ and $\beta' = \bar{\beta}\left(\frac{q'_1{}^2}{M^2}, \frac{q'_2{}^2}{M^2}\right)$ with
 $\bar{\beta}(a, b) \equiv (1 + a^2 + b^2 - 2a - 2b - 2ab)^{1/2}$.

$\hat{J}_1^{\lambda_1}_{\sigma_1 \sigma_3}(f_{\sigma_1} \rightarrow f_{\sigma_3} V_{\lambda_1}^*)$	$[\cos \theta_1 \rightarrow z_1/(2 - z_1)]$	
$\hat{J}_1^{+}_{++} = -(\hat{J}_1^{-}_{--})^*$	$\frac{1}{2 \cos \theta_1} (1 + \cos \theta_1) e^{-i\phi_1}$	$\frac{1}{z_1} e^{-i\phi_1}$
$\hat{J}_1^0_{++} = \hat{J}_1^0_{--}$	$-\frac{1}{\sqrt{2} \cos \theta_1} \sin \theta_1$	$-\frac{\sqrt{2(1 - z_1)}}{z_1}$
$\hat{J}_1^{-}_{++} = -(\hat{J}_1^{+}_{--})^*$	$-\frac{1}{2 \cos \theta_1} (1 - \cos \theta_1) e^{i\phi_1}$	$-\frac{1 - z_1}{z_1} e^{i\phi_1}$
$\hat{J}_1^{\lambda_1}_{+-} = \hat{J}_1^{\lambda_1}_{-+}$	0	0

$\hat{J}'_1^{\lambda'_1}_{\sigma'_1 \sigma'_3}(V_{\lambda'_1}^* \rightarrow f_{\sigma'_1} \bar{f}_{\sigma'_3})$	$[\cos \theta'_1 \rightarrow 2z'_1 - 1]$	
$\hat{J}'_1^{+}_{+-} = -(\hat{J}'_1^{-}_{-+})^*$	$\frac{1}{2} (1 + \cos \theta'_1) e^{i\phi'_1}$	$z'_1 e^{i\phi'_1}$
$\hat{J}'_1^0_{+-} = \hat{J}'_1^0_{-+}$	$\frac{1}{\sqrt{2}} \sin \theta'_1$	$\sqrt{2z'_1(1 - z'_1)}$
$\hat{J}'_1^{-}_{+-} = -(\hat{J}'_1^{+}_{-+})^*$	$\frac{1}{2} (1 - \cos \theta'_1) e^{-i\phi'_1}$	$(1 - z'_1) e^{-i\phi'_1}$
$\hat{J}'_1^{\lambda'_1}_{++} = \hat{J}'_1^{\lambda'_1}_{--}$	0	0

$\widehat{J}_1^{\lambda_1}_{\sigma_1 \sigma_3}(g_{\sigma_1} \rightarrow g_{\sigma_3} V_{\lambda_1}^*)$	$[\cos \theta_1 \rightarrow z_1/(2-z_1)]$	
$\widehat{J}_1^{+}_{++} = -(\widehat{J}_1^{-}_{--})^*$	$\frac{1}{2 \sin \theta_1 \cos \theta_1} (1 + \cos \theta_1)^2 e^{-i\phi_1}$	$\frac{1}{z_1 \sqrt{1-z_1}} e^{-i\phi_1}$
$\widehat{J}_1^0_{++} = \widehat{J}_1^0_{--}$	$-\frac{1}{\sqrt{2} \cos \theta_1}$	$-\frac{2-z_1}{\sqrt{2} z_1}$
$\widehat{J}_1^{-}_{++} = -(\widehat{J}_1^{+}_{--})^*$	$-\frac{1}{2 \sin \theta_1 \cos \theta_1} (1 - \cos \theta_1)^2 e^{i\phi_1}$	$-\frac{(1-z_1)^2}{z_1 \sqrt{1-z_1}} e^{i\phi_1}$
$\widehat{J}_1^{+}_{+-} = -(\widehat{J}_1^{-}_{-+})^*$	$-\frac{2}{\tan \theta_1} e^{i\phi_1}$	$-\frac{z_1}{\sqrt{1-z_1}} e^{i\phi_1}$
$\widehat{J}_1^{0/-}_{+-} = \widehat{J}_1^{0/+}_{-+}$	0	0

$\widehat{J}'^{\lambda'_1}_{1\sigma'_1\sigma'_3}(V_{\lambda'_1}^* \rightarrow g_{\sigma'_1} g_{\sigma'_3})$	$[\cos \theta'_1 \rightarrow 2z'_1 - 1]$	
$\widehat{J}'^{+}_{1+-} = -(\widehat{J}'^{-}_{-+})^*$	$-\frac{1}{2 \sin \theta'_1} (1 + \cos \theta'_1)^2 e^{i\phi'_1}$	$-\frac{z'_1{}^2}{\sqrt{z'_1(1-z'_1)}} e^{i\phi'_1}$
$\widehat{J}'^0_{1+-} = \widehat{J}'^0_{-+}$	$-\frac{1}{\sqrt{2}} \cos \theta'_1$	$-\frac{2z'_1-1}{\sqrt{2}}$
$\widehat{J}'^{-}_{1+-} = -(\widehat{J}'^{+}_{-+})^*$	$\frac{1}{2 \sin \theta'_1} (1 - \cos \theta'_1)^2 e^{-i\phi'_1}$	$\frac{(1-z'_1)^2}{\sqrt{z'_1(1-z'_1)}} e^{-i\phi'_1}$
$\widehat{J}'^{+}_{1++} = -(\widehat{J}'^{-}_{--})^*$	$\frac{2}{\sin \theta'_1} e^{-i\phi'_1}$	$\frac{1}{\sqrt{z'_1(1-z'_1)}} e^{-i\phi'_1}$
$\widehat{J}'^{0/-}_{1++} = \widehat{J}'^{0/+}_{1--}$	0	0

X	(λ_X)	V_i	$\Gamma_{XV_1V_2}^{\mu_1\mu_2}(q_1, q_2; \lambda_X)/g_{XV_1V_2}(q_1, q_2)$
H	(0)	W, Z	$g^{\mu_1\mu_2}$
H	(0)	$\gamma, Z/\gamma, g$	$q_1 \cdot q_2 g^{\mu_1\mu_2} - q_2^{\mu_1} q_1^{\mu_2}$
A	(0)	$\gamma, Z/\gamma, g$	$\epsilon^{\mu_1\mu_2\alpha\beta} q_{1\alpha} q_{2\beta}$
G	$(\pm 2, \pm 1, 0)$	W, Z, γ, g	$\epsilon_{\alpha\beta}(p_X, \lambda_X) \hat{\Gamma}_{GVV}^{\alpha\beta, \mu_1\mu_2}(q_1, q_2)$

λ_X	$(\lambda_1 \lambda_2)$	CP -even		CP -odd
		$H(\text{WBF})$	$H(\text{loop-induced})$	A
0	($\pm\pm$)	-1	$-\frac{1}{2}(M^2 + Q_1^2 + Q_2^2)$	$\mp\frac{i}{2}\sqrt{(M^2 + Q_1^2 + Q_2^2)^2 - 4Q_1^2Q_2^2}$
0	(00)	$\frac{(M^2 + Q_1^2 + Q_2^2)}{2Q_1Q_2}$	Q_1Q_2	0

λ_X	$(\lambda_1 \lambda_2)$	G
± 2	($\pm\mp$)	$-(M^2 + Q_1^2 + Q_2^2 + 2m_V^2)$
± 1	(± 0)	$\frac{1}{\sqrt{2}MQ_2}[Q_2^2(M^2 - Q_1^2 + Q_2^2) - m_V^2(M^2 + Q_1^2 - Q_2^2)]$
± 1	($0\mp$)	$\frac{1}{\sqrt{2}MQ_1}[Q_1^2(M^2 + Q_1^2 - Q_2^2) - m_V^2(M^2 - Q_1^2 + Q_2^2)]$
0	($\pm\pm$)	$\frac{1}{\sqrt{6}M^2}[(Q_1^2 - Q_2^2)^2 + M^2(Q_1^2 + Q_2^2 - 2m_V^2)]$
0	(00)	$-\frac{1}{\sqrt{6}Q_1Q_2}\left[4Q_1^2Q_2^2 + 2m_V^2(M^2 + Q_1^2 + Q_2^2) - \frac{m_V^2}{M^2}\{(M^2 + Q_1^2 + Q_2^2)^2 - 4Q_1^2Q_2^2\}\right]$

With the minimal cuts:

$$p_{T_j} > 20 \text{ GeV}, \quad |\eta_j| < 5, \quad R_{jj} = \sqrt{\Delta\eta_{jj}^2 + \Delta\phi_{jj}^2} > 0.6$$

plus the VBF cuts:

$$\eta_{j_1} > 0 > \eta_{j_2}, \quad \Delta\eta_{jj} = \eta_{j_1} - \eta_{j_2} > \Delta\eta_{jj\min}$$

we find

$\sigma_{\text{VBF}}/\sigma_{\text{exact}}$	$\Delta\eta_{jj} > 3$	$\Delta\eta_{jj} > 4$	$\Delta\eta_{jj} > 5$
$qq \rightarrow qqH/A/G$	1.00/1.00/1.58	1.00/1.00/1.43	1.00/1.00/1.25
$qg \rightarrow qgH/A/G$	1.07/1.05/1.30	1.04/1.03/1.18	1.02/1.02/1.11
$gg \rightarrow ggH/A/G$	1.07/1.06/1.16	1.04/1.04/1.11	1.02/1.02/1.07

In addition, if we impose the p_{T_j} slicing cut:

$$20 \text{ GeV} < p_{T_j} < 100 \text{ GeV}$$

we find

$\sigma_{\text{VBF}}/\sigma_{\text{exact}}$	$\Delta\eta_{jj} > 3$	$\Delta\eta_{jj} > 4$	$\Delta\eta_{jj} > 5$
$qq \rightarrow qqH/A/G$	1.00/1.00/1.02	1.00/1.00/1.02	1.00/1.00/1.02
$qg \rightarrow qgH/A/G$	1.04/1.04/1.07	1.03/1.03/1.06	1.02/1.02/1.04
$gg \rightarrow ggH/A/G$	1.05/1.05/1.09	1.04/1.04/1.07	1.02/1.02/1.05

$$\begin{aligned}\mathcal{M}_{\sigma_1\sigma_3,\sigma_2\sigma_4}^{\lambda_X=0} &= \sum_{V_1,V_2} \frac{1}{(q_1^2 - m_{V_1}^2)(q_2^2 - m_{V_2}^2)} \sum_{\lambda_{1,2}=\pm,0} \left(J_{a_1a_3}^{V_1}\right)_{\sigma_1\sigma_3}^{\lambda_1} \left(J_{a_2a_4}^{V_2}\right)_{\sigma_2\sigma_4}^{\lambda_2} \left(\mathcal{M}_{V_1V_2}^X\right)_{\lambda_1\lambda_2}^0 \\ &\sim \hat{J}_1^+_{\sigma_1\sigma_3} \hat{J}_2^+_{\sigma_2\sigma_4} \hat{\mathcal{M}}_X^0_{++} + \hat{J}_1^0_{\sigma_1\sigma_3} \hat{J}_2^0_{\sigma_2\sigma_4} \hat{\mathcal{M}}_X^0_{00} + \hat{J}_1^-_{\sigma_1\sigma_3} \hat{J}_2^-_{\sigma_2\sigma_4} \hat{\mathcal{M}}_X^0_{--}\end{aligned}$$

$$\begin{aligned}\mathcal{M}_{\sigma_1,\sigma_2}^{\lambda_X=0} &\sim -\hat{J}_1^+_{\sigma_1}(\theta_1) \hat{J}_2^+_{\sigma_2}(\theta_2) \hat{\mathcal{M}}_X^0_{++} e^{-i\Delta\phi_{12}} \\ &\quad - \hat{J}_1^0_{\sigma_1}(\theta_1) \hat{J}_2^0_{\sigma_2}(\theta_2) \hat{\mathcal{M}}_X^0_{00} \\ &\quad - \hat{J}_1^-_{\sigma_1}(\theta_1) \hat{J}_2^-_{\sigma_2}(\theta_2) \hat{\mathcal{M}}_X^0_{--} e^{i\Delta\phi_{12}}\end{aligned}$$

