

Argyres-Seiberg duality and the Higgs branch

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in collaboration with
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[arXiv:0810.4541]

IPMU, December, 2008

Introduction

- Montonen-Olive S-duality in $\mathcal{N} = 4$ theory
- Seiberg-duality in $\mathcal{N} = 1$ theory

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Argyres-Seiberg, Nov 2007

a totally **new class** of S-duality in $\mathcal{N} = 2$ theory

- Richer structure !
- [Argyres-Wittig] [Aharony-YT] [Shapere-YT]

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- Seiberg-duality in $\mathcal{N} = 1$ theory

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- Richer structure !
- [Argyres-Wittig] [Aharony-YT] [Shapere-YT]
- Today: Higgs branch side of the story

1. Argyres-Seiberg duality

2. Higgs branch

3. Summary

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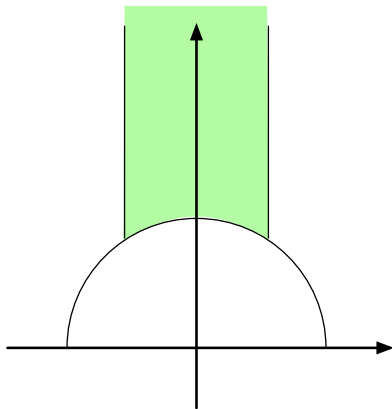
Montonen-Olive S-duality

$$\mathcal{N} = 4 \text{ SU}(N)$$

$$\tau = \frac{\theta}{2\pi} + \frac{4\pi i}{g^2}$$

$$\tau \rightarrow \tau + 1, \quad \tau \rightarrow -\frac{1}{\tau}$$

- Exchanges monopoles
W-bosons
- Comes from S-duality of
Type IIB



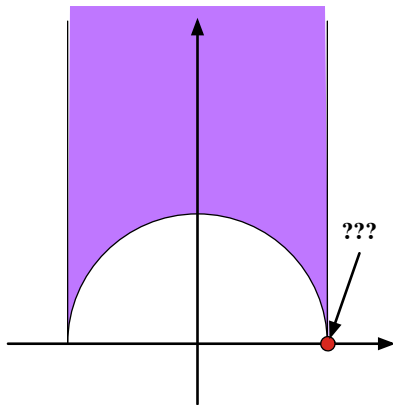
S-duality in $\mathcal{N} = 2$

SU(3) with $N_f = 6$

$$\tau = \frac{\theta}{\pi} + \frac{8\pi i}{g^2}$$

$$\tau \rightarrow \tau + 2, \quad \tau \rightarrow -\frac{1}{\tau}$$

- Exchanges monopoles and **quarks**
- Infinitely Strongly coupled at $\tau = 1$



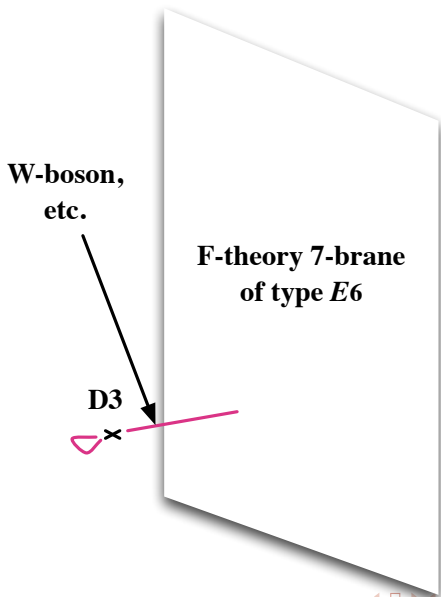
SU(3) + 6 flavors

at coupling τ

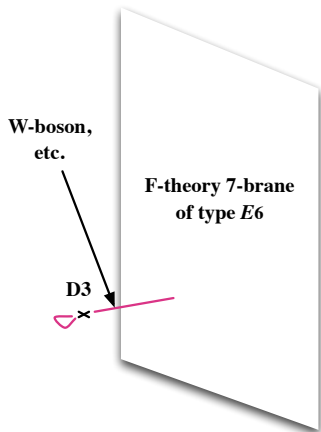


SU(2) + 1 flavor + SCFT[E_6]

at coupling $\tau' = 1/(1 - \tau)$, $SU(2) \subset E_6$ is gauged



- a 3-brane probing F-theory singularity of type E_6 .
- Gauge symmetry on 7-brane
 → Flavor symmetry on 3-brane
- Motion transverse to 7-brane
 → Vector multiplet moduli u ,
- Motion parallel to 7-brane
 → free hypermultiplet, discard
- Conformal when $u = 0$, $\dim(u) = 3$
- Family of Seiberg-Witten curve is the elliptic fibration of F-theory.



Argyres-Seiberg: Dimensions

SU(3) + 6 flavors

$$\dim(\text{tr } \phi^2) = 2,$$

$$\dim(\text{tr } \phi^3) = 3$$



SU(2) + 1 flavor + SCFT[E_6]

$$u \text{ of SU(2) : } \dim = 2,$$

$$u \text{ of SCFT}[E_6] : \dim = 3$$

Argyres-Seiberg: Flavor symmetry

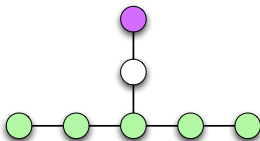
$SU(3) + 6$ flavors

- Flavor symmetry: $U(6) = U(1) \times SU(6)$



$SU(2) + 1$ flavor + SCFT[E_6]

- $SO(2)$ acts on 1 flavor = 2 half-hyper of $SU(2)$ doublet
- $SU(2) \subset E_6$ is gauged
- $SU(2) \times SU(6) \subset E_6$ is a maximal regular subalgebra



Current Algebra Central Charge

Normalize s.t. a free hyper in the fund. of $\mathbf{SU}(N)$ contributes **2** to k_G

$$J_\mu^a(x) J_\nu^b(0) = \frac{3}{4\pi^2} k_G \delta^{ab} \frac{x^2 g_{\mu\nu} - 2x_\mu x_\nu}{x^8} + \dots$$

A bifundamental hyper under $\mathbf{SU}(N) \times \mathbf{SU}(M)$

$$\rightarrow k_{\mathbf{SU}(N)} = 2M, \quad k_{\mathbf{SU}(M)} = 2N$$

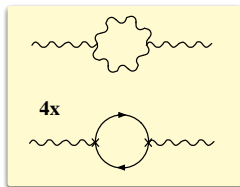
SU(3) + 6 flavors

$$k_{\mathbf{SU}(6)} = \mathbf{6}$$

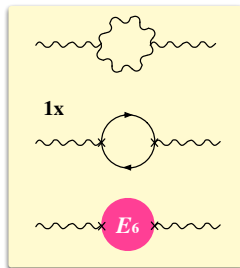
k for SCFT[E_6]

$SU(2) \subset E_6$ central charge:

$SU(2) + 4$ flavors



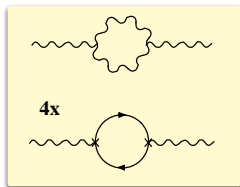
$SU(2) + 1$ flavor + SCFT[E_6]



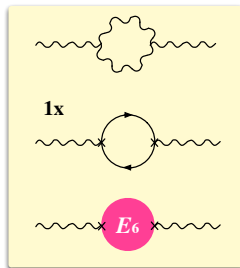
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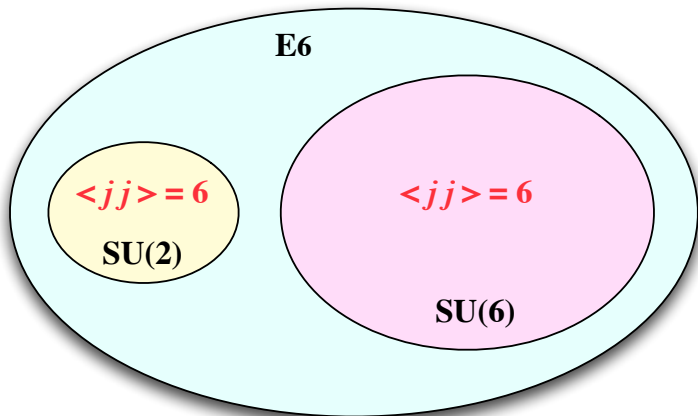


$SU(2) + 1$ flavor + SCFT[E_6]



$$\begin{aligned}
 & \text{fermion loop} = \langle j_\mu j_\nu \rangle_{\text{free hyper}} \\
 & \text{E}_6 \text{ line} = \langle j_\mu j_\nu \rangle_{\text{SCFT}[E_6]}
 \end{aligned}
 \quad \Rightarrow \quad
 \langle j_\mu j_\nu \rangle_{\text{SCFT}[E_6]} = 3 \langle j_\mu j_\nu \rangle_{\text{free hyper}} = 6$$

k for SCFT[E_6]



SU(3) + 6 flavors

- $k_{U(1)} = 3 \times 6 = 18$

**SU(2) + 1 flavor + SCFT[E_6]**

- $k_{U(1)} = 2 \times 1 \times q^2 \rightarrow q = 3$

Matching Seiberg-Witten curves

Argyres & Seiberg studied the SW curve on the both sides, and found

Curve of **SU(3)** theory with generic masses $\sum_i m_i Q^i \tilde{Q}_i$
at $\tau \rightarrow \mathbf{1}$

⊃ Curve of **SCFT**[E_6] with masses to **SU(6)** $\subset E_6$

I won't (can't) explain it today ...

Another example: E_7

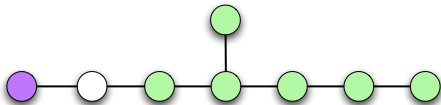
$\mathrm{USp}(4) + 12$ half-hypers in 4

- $\dim(\mathrm{tr} \phi^2) = 2$, $\dim(\mathrm{tr} \phi^4) = 4$
- $k_{\mathrm{SO}(12)} = 8$



$\mathrm{SU}(2)$ w/ $\mathrm{SCFT}[E_7]$

- $\dim(\mathrm{tr} \phi^2) = 2$ from $\mathrm{SU}(2)$, $\dim(u) = 4$ from $\mathrm{SCFT}[E_7]$
- $\mathrm{SU}(2) \times \mathrm{SO}(12) \subset E_7$
- $k_{\mathrm{SU}(2) \subset E_7} = 8$



	\mathfrak{g}	w/ \mathfrak{r}	$= \tilde{\mathfrak{g}}$	$w/ \tilde{\mathfrak{r}}$	\oplus	SCFT $[d: \mathfrak{h}]$
1.	$\mathfrak{sp}(3)$	$14 \oplus 11 \cdot 6$	$= \mathfrak{sp}(2)$			$[6: E_8]$
2.	$\mathfrak{su}(6)$	$20 \oplus 15 \oplus \overline{15} \oplus 5 \cdot 6 \oplus 5 \cdot \overline{6}$	$= \mathfrak{su}(5)$	$5 \oplus \overline{5} \oplus 10 \oplus \overline{10}$		$[6: E_8]$
3.	$\mathfrak{so}(12)$	$3 \cdot 32 \oplus 32' \oplus 4 \cdot 12$	$= \mathfrak{so}(11)$	$3 \cdot 32$		$[6: E_8]$
4.	G_2	$8 \cdot 7$	$= \mathfrak{su}(2)$	2		$[6: \mathfrak{sp}(5)]$
5.	$\mathfrak{so}(7)$	$4 \cdot 8 \oplus 6 \cdot 7$	$= \mathfrak{sp}(2)$	$5 \cdot 4$		$[6: \mathfrak{sp}(5)]$
6.	$\mathfrak{su}(6)$	$21 \oplus \overline{21} \oplus 20 \oplus 6 \oplus \overline{6}$	$= \mathfrak{su}(5)$	$10 \oplus \overline{10}$		$[6: \mathfrak{sp}(5)]$
7.	$\mathfrak{sp}(2)$	$12 \cdot 4$	$= \mathfrak{su}(2)$			$[4: E_7]$
8.	$\mathfrak{su}(4)$	$2 \cdot 6 \oplus 6 \cdot 4 \oplus 6 \cdot \overline{4}$	$= \mathfrak{su}(3)$	$2 \cdot 3 \oplus 2 \cdot \overline{3}$		$[4: E_7]$
9.	$\mathfrak{so}(7)$	$6 \cdot 8 \oplus 4 \cdot 7$	$= G_2$	$4 \cdot 7$		$[4: E_7]$
10.	$\mathfrak{so}(8)$	$6 \cdot 8 \oplus 4 \cdot 8' \oplus 2 \cdot 8''$	$= \mathfrak{so}(7)$	$6 \cdot 8$		$[4: E_7]$
11.	$\mathfrak{so}(8)$	$6 \cdot 8 \oplus 6 \cdot 8'$	$= G_2$			$[4: E_7] \oplus [4: E_7]$
12.	$\mathfrak{sp}(2)$	$6 \cdot 5$	$= \mathfrak{su}(2)$			$[4: \mathfrak{sp}(3) \oplus \mathfrak{su}(2)]$
13.	$\mathfrak{sp}(2)$	$4 \cdot 4 \oplus 4 \cdot 5$	$= \mathfrak{su}(2)$	$3 \cdot 2$		$[4: \mathfrak{sp}(3) \oplus \mathfrak{su}(2)]$
14.	$\mathfrak{su}(4)$	$10 \oplus \overline{10} \oplus 2 \cdot 4 \oplus 2 \cdot \overline{4}$	$= \mathfrak{su}(3)$	$3 \oplus \overline{3}$		$[4: \mathfrak{sp}(3) \oplus \mathfrak{su}(2)]$
15.	$\mathfrak{su}(3)$	$6 \cdot 3 \oplus 6 \cdot \overline{3}$	$= \mathfrak{su}(2)$	$2 \cdot 2$		$[3: E_6]$
16.	$\mathfrak{su}(4)$	$4 \cdot 6 \oplus 4 \cdot 4 \oplus 4 \cdot \overline{4}$	$= \mathfrak{sp}(2)$	$6 \cdot 4$		$[3: E_6]$
17.	$\mathfrak{su}(3)$	$3 \oplus \overline{3} \oplus 6 \oplus \overline{6}$	$= \mathfrak{su}(2)$	$n \cdot 2$		$[3: \mathfrak{h}]$

Table 2: Predicted dualities with one marginal operator.

Central charges:

$$\langle T_{\mu}^{\mu} \rangle = a \cdot \text{Euler} + c \cdot \text{Weyl}^2$$

calculable for **SCFT**[$E_{6,7}$] using

- **SU(3)** w/ 6 flavors \leftrightarrow **SU(2)** + 1 flavor + **SCFT**[E_6]
- **USp(4)** w/ 12 flavors \leftrightarrow **SU(2)** + **SCFT**[E_7]

We performed holographic calculation for **SCFT**[$E_{6,7,8}$]

G	E_6	E_7	E_8
k_G	6	8	12
$24a$	41	59	95
$6c$	13	19	31

It was done **before** publication of [Argyres-Wittig]
 Perfectly agreed ! Power of string theory.

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Objective

- Argyres & Seiberg studied the story on the Coulomb branch side.
- I wanted to know the Higgs branch side of the story.

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- Argyres & Seiberg studied the story on the Coulomb branch side.
- I wanted to know the Higgs branch side of the story.
- I have two friends who are experts on hyperkähler things !



$\mathcal{N} = 1$ Seiberg duality

SU(N) w/ N_f flavors q^i, \tilde{q}_j

- $m_j^i = q_a^i \tilde{q}_j^a$
- $b^{i_1 i_2 \dots i_{N_f}} = \epsilon^{a_1 a_2 \dots a_N} q_{a_1}^{i_1} q_{a_2}^{i_2} \dots q_{a_N}^{i_{N_f}}$
- $\tilde{b}_{i_1 i_2 \dots i_{N_f}} = \epsilon_{a_1 a_2 \dots a_N} \tilde{q}_{i_1}^{a_1} \tilde{q}_{i_2}^{a_2} \dots \tilde{q}_{i_N}^{a_N}$



SU(N') w/ N_f flavors Q_i, \tilde{Q}^i

- w/ singlets $M_j^i, W = Q_i \tilde{Q}^j M_j^i$
- $N' = N - N_f$
- $B_{j_1 j_2 \dots j_{N'}} = \epsilon_{a_1 a_2 \dots a_{N'}} Q_{j_1}^{a_1} Q_{j_2}^{a_2} \dots Q_{j_{N'}}^{a_{N'}}$
- $\tilde{B}^{j_1 j_2 \dots j_{N'}} = \epsilon^{a_1 a_2 \dots a_{N'}} \tilde{Q}_{a_1}^{j_1} \tilde{Q}_{a_2}^{j_2} \dots \tilde{Q}_{a_{N'}}^{j_{N'}}$

$\mathcal{N} = 1$ Seiberg duality

Mapping of operators:

$$\begin{aligned}m_j^i &= q^i \tilde{q}_j \leftrightarrow M_j^i \\b^{i_1 i_2 \dots i_N} &\leftrightarrow \epsilon^{i_1 i_2 \dots i_N j_1 \dots j_{N'}} B_{j_1 j_2 \dots j_{N_f - N_c}} \\ \tilde{b}_{i_1 i_2 \dots i_{N_f}} &\leftrightarrow \epsilon_{i_1 i_2 \dots i_{N_f} j_1 \dots j_{N'}} \tilde{B}^{j_1 j_2 \dots j_{N'}}\end{aligned}$$

Mapping of constraints:

$$\begin{aligned}m_j^{[i} b^{i_1 \dots i_N]} &= 0 & (\tilde{q}_j^a q_{[a}^{i_1} \dots q_{a_N]}^{i_N]} &= 0) \\ \rightarrow M_j^{i_1} B_{i_1 \dots i_{N'}} &= 0 & (M_j^i Q_i^a &= 0)\end{aligned}$$

etc.

Computation of moduli space, $\mathcal{N} = 1$

$$\frac{\{F = 0, D = 0\}}{G} = \frac{\{F = 0\}}{G_{\mathbb{C}}}$$

Basically:

- List gauge-invariant chiral operators
- Impose F-term = 0
- Study the constraints

Computation of moduli space, $\mathcal{N} = 2$

$$\frac{\{F^A = 0, D^A = 0\}}{G} = \frac{\{F^A = 0\}}{G_{\mathbb{C}}}$$

- $W = Q\Phi\tilde{Q} \rightarrow F^A = t_{a\bar{a}}^A Q^a \tilde{Q}^{\bar{a}}$
- $F^A = D^A = 0$ imposes **3 dim G** conditions
- $/G$ removes another **dim G**
- \rightarrow loose **4 dim G** dimensions

Dimensions

SU(3) + 6 flavors Q^i, \tilde{Q}_i

$$4 \times 3 \times 6 - 4 \dim \text{SU}(3) = 40$$



SU(2) + 1 flavor q, \tilde{q} + SCFT[E_6]

$$4 \times 2 + \text{????} - 4 \dim \text{SU}(2)$$

Higgs branch of SCFT[E_6]

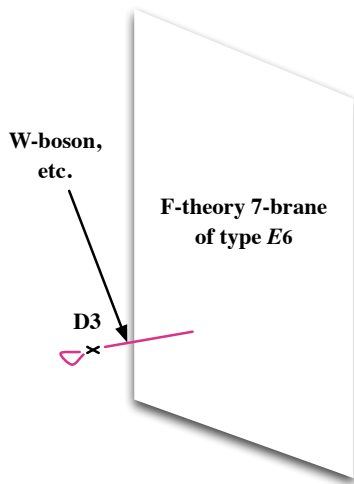
D3 brane absorbed into the 7-brane

→ becomes an instanton of type E_6 !

Center of mass along the 7-brane
decoupled .

$$\begin{aligned} \dim(\text{centered } k\text{-instanton moduli}) \\ = 4h_{E_6}k - 4 \end{aligned}$$

$$k = 1, h_{E_6} = 12 \rightarrow \dim = 44$$



SU(3) + 6 flavors Q^i, \tilde{Q}_i

$$4 \times 3 \times 6 - 4 \dim \text{SU}(3) = 40$$



SU(2) + 1 flavor q, \tilde{q} + SCFT[E_6]

$$4 \times 2 + 44 - 4 \dim \text{SU}(2) = 40$$

SU(3) + 6 flavors Q^i, \tilde{Q}_i

- $M_j^i = Q_a^i \tilde{Q}_j^a$
- $B^{[ijk]} = \epsilon^{abc} Q_a^i Q_b^j Q_c^k$
- $\tilde{B}_{[ijk]} = \epsilon_{abc} \tilde{Q}_i^a \tilde{Q}_j^b \tilde{Q}_k^c$
- Lots of constraints.



SU(2) + 1 flavor q, \tilde{q} + SCFT[E_6]

- ??? → Need to know more about the E_6 instanton moduli space.
- But how? We don't have ADHM for exceptional groups...

One-instanton moduli

- Any one-instanton of G is an embedding of the $\mathbf{SU}(2)$ BPST instanton via $\mathbf{SU}(2) \subset G$
- Space equivalent to a subspace of $\mathfrak{g}_{\mathbb{C}}$, **minimal nilpotent orbit**

$$G_{\mathbb{C}} \cdot T^{\rho}, \quad \rho : \text{highest root}$$

- Equations explicitly known. Just quadratic. [Joseph, Kostant]
- Let $V(\alpha)$: irrep with highest weight α , and $\mathfrak{g}_{\mathbb{C}} = V(\rho)$.
- Decompose

$$\mathbf{Sym}^2 V(\rho) = V(2\rho) \oplus \mathcal{I}$$

then

$$\{G_{\mathbb{C}} \cdot T^{\rho}\} = \{\mathbb{X} \in \mathfrak{g}_{\mathbb{C}} \mid (\mathbb{X} \otimes \mathbb{X})|_{\mathcal{I}} = \mathbf{0}\}$$

One-instanton moduli of $SU(2)$

- Equations explicitly known. Just quadratic. [Joseph, Kostant]
- Let $V(\alpha)$: irrep with highest weight α , and $\mathfrak{g}_{\mathbb{C}} = V(\rho)$.
- Decompose

$$\mathrm{Sym}^2 V(\rho) = V(2\rho) \oplus \mathcal{I}$$

then

$$\{G_{\mathbb{C}} \cdot T^{\rho}\} = \{\mathbb{X} \in \mathfrak{g}_{\mathbb{C}} \mid (\mathbb{X} \otimes \mathbb{X})|_{\mathcal{I}} = \mathbf{0}\}$$

- Parametrize $\mathfrak{su}(2)$ by $x_{1,2,3}$
- $\mathrm{Sym}^2 \mathfrak{3} = \mathfrak{5} \oplus \mathfrak{1}$
- $x_1^2 + x_2^2 + x_3^2 = 0$: just $\mathbb{C}^2/\mathbb{Z}_2$.

One-instanton moduli of E_6

$$\mathbb{X} \in \text{adj}(E_6), \quad (\mathbb{X} \otimes \mathbb{X})|_{\mathcal{I}} = 0$$

- We couple $\mathbf{SU}(2)$ gauge field Φ to \mathbb{X} .
- Convenient to decompose under $\mathbf{SU}(2) \times \mathbf{SU}(6) \subset \mathbb{X}$

$$X_j^i, \quad Y_\alpha^{[ijk]}, \quad Z_{(\alpha\beta)}$$

- Superpotential is $W = \Phi^{\alpha\beta} Z_{\alpha\beta}$.

SU(3) + 6 flavors Q^i, \tilde{Q}_i

- $M_j^i = Q_a^i \tilde{Q}_j^a \rightarrow \text{tr } M, \hat{M}_j^i$
- $B^{[ijk]} = \epsilon^{abc} Q_a^i Q_b^j Q_c^k, \quad \tilde{B}_{[ijk]} = \epsilon_{abc} \tilde{Q}_i^a \tilde{Q}_j^b \tilde{Q}_k^c$



SU(2) + 1 flavor v, \tilde{v} + SCFT[E_6]

- $X_j^i, Y_\alpha^{[ijk]}, Z_{(\alpha\beta)}$
- F-term=0 $\rightarrow Z_{(\alpha\beta)} + v_{(\alpha} \tilde{v}_{\beta)} = 0$
- $\hat{M}_j^i = X_j^i$
- $M = \epsilon^{\alpha\beta} v_\alpha \tilde{v}_\beta$
- $B^{[ijk]} = \epsilon^{\alpha\beta} v_\alpha Y_\beta^{ijk}, \quad \tilde{B}^{[ijk]} = \epsilon^{\alpha\beta} \tilde{v}_\alpha Y_\beta^{ijk}$

U(1) charge

$$B^{[ijk]} = \epsilon^{abc} Q_a^i Q_b^j Q_c^k$$

$$B^{[ijk]} = \epsilon^{\alpha\beta} v_\alpha Y_\beta^{ijk}$$

$$\tilde{B}_{[ijk]} = \epsilon_{abc} \tilde{Q}_i^a \tilde{Q}_j^b \tilde{Q}_k^c$$

$$\tilde{B}_{[ijk]} = \epsilon^{\alpha\beta} \tilde{v}_\alpha Y_\beta^{ijk}$$

U(1) charge of Q : 1 \rightarrow U(1) charge of v : 3

SU(3) + 6 flavors

- $k_{U(1)} = 3 \times 6 = 18$

SU(2) + 1 flavor + SCFT[E_6]

- $k_{U(1)} = 2 \times 1 \times (\text{charge of } v)^2 \rightarrow \text{charge of } v = 3$

Constraints of min. nilpotent orbit $E_{6,\mathbb{C}} \cdot T^\rho$

$$(\mathbb{X} \otimes \mathbb{X})|_{\mathcal{I}} = \mathbf{0} \text{ where } \mathbf{Sym}^2 V(\rho) = V(2\rho) \oplus \mathcal{I}$$

- We decompose $\mathbb{X} = (X_j^i, Y_\alpha^{[ijk]}, Z_{(\alpha\beta)})$
- Decompose \mathcal{I} under $\mathbf{SU}(2) \times \mathbf{SU}(6)$
- Coefficients laboriously fixed

$$0 = X^i_j Z_{\alpha\beta} + \frac{1}{4} Y_{(\alpha}^{ikl} Y_{jkl)\beta},$$

$$0 = X^l_{\{i} Y_{[jk]\}l\alpha},$$

$$0 = X^{\{i} Y_{\alpha}^{[jk]\}l},$$

$$0 = Y_{\alpha}^{ijk} Z_{\beta\gamma} \epsilon^{\alpha\beta} + X^{[i} Y_{\gamma}^{jk]l},$$

$$0 = (Y_{\alpha}^{ijm} Y_{klm\beta} \epsilon^{\alpha\beta} - 4X^{[i} X^{j]l}) \Big|_{0,1,0,1,0},$$

$$0 = X^i_k X^k_j - \frac{1}{6} \delta^i_j X^k_l X^l_k,$$

$$0 = Y_{\alpha}^{ijk} Y_{ijk\beta} \epsilon^{\alpha\beta} + 24Z_{\alpha\beta} Z_{\gamma\delta} \epsilon^{\alpha\gamma} \epsilon^{\beta\delta},$$

$$0 = X^i_j X^j_i + 3Z_{\alpha\beta} Z_{\gamma\delta} \epsilon^{\alpha\gamma} \epsilon^{\beta\delta}.$$

F-term eq.

$$Z_{(\alpha\beta)} + v_{(\alpha}\tilde{v}_{\beta)} = 0.$$

Identifications

$$\begin{aligned}\hat{M}_j^i &= X_j^i, & \text{tr } M &= \epsilon^{\alpha\beta} v_\alpha \tilde{v}_\beta \\ B^{[ijk]} &= \epsilon^{\alpha\beta} v_\alpha Y_\beta^{ijk}\end{aligned}$$

$$0 = X_k^i X_j^k - \frac{1}{6} \delta_j^i X_l^k X_k^l,$$

$$0 = X_j^i X_i^j + 3Z_{\alpha\beta} Z_{\gamma\delta} \epsilon^{\alpha\gamma} \epsilon^{\beta\delta},$$

$$0 = Y_\alpha^{ijk} Z_{\beta\gamma} \epsilon^{\alpha\beta} + X_l^{[i} Y_\gamma^{jkl]}$$

$$\hat{M}_j^i \hat{M}_k^j = \frac{1}{6} \delta_j^i M_n^m M_m^n,$$

$$\hat{M}_j^i \hat{M}_i^j = \frac{1}{6} (\text{tr } M)^2,$$

$$\hat{M}_j^{[i} B^{kl]j} = \frac{1}{6} (\text{tr } M) B^{ikl},$$

Constraints

$$W = Q^i \Phi \tilde{Q}_i \rightarrow Q_a^i \tilde{Q}_i^b - \frac{1}{3} \delta_a^b Q_c^i \tilde{Q}_i^c = 0 \rightarrow$$

$$M_j^i M_k^j = \frac{1}{3} (\text{tr } M) M_k^i, \quad M_j^{[i} B^{kl]j} = \frac{1}{3} (\text{tr } M) B^{ijk}, \quad \text{etc.}$$

$$\hat{M}_j^i \hat{M}_k^j = \frac{1}{6} \delta_j^i \hat{M}_n^m \hat{M}_m^n,$$

$$\hat{M}_j^i \hat{M}_i^j = \frac{1}{6} (\text{tr } M)^2,$$

$$\hat{M}_j^{[i} B^{kl]j} = \frac{1}{6} (\text{tr } M) B^{ikl}, \dots$$

1. Argyres-Seiberg duality

2. Higgs branch

3. Summary

Done

- Argyres-Seiberg duality reviewed.
- Higgs branches compared for E_6 → Perfect agreement !

Done

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- Higgs branches compared for E_6 → Perfect agreement !

To do

- Other cases E_7 ?
- String theoretic realization of the duality